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GRADUATE SCHOOL OF ARTS AND SCIENCES

Dissertation

**MEASUREMENT OF THE ANOMALOUS MAGNETIC
MOMENT OF THE POSITIVE MUON TO .SOMETHING
PARTS PER BILLION**

by

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requirements for the degree of
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Dedication

I dedicate this thesis to

Acknowledgments

Here go all your acknowledgments. You know, your advisor, funding agency, lab mates, etc., and of course your family.

As for me, I would like to thank Jonathan Polimeni for cleaning up old LaTeX style files and templates so that Engineering students would not have to suffer typesetting dissertations in MS Word. Also, I would like to thank IDS/ISS group (ECE) and CV/CNS lab graduates for their contributions and tweaks to this scheme over the years (after many frustrations when preparing their final document for BU library). In particular, I would like to thank Limor Martin who has helped with the transition to PDF-only dissertation format (no more printing hardcopies – hooray !!!)

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Janusz Konrad
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(Order No.)

NICHOLAS BRENNAN KINNAIRD

Boston University, Graduate School of Arts and Sciences, 2019

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ABSTRACT

Have you ever wondered why this is called an *abstract*? Weird thing is that its legal to cite the abstract of a dissertation alone, apart from the rest of the manuscript.

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- C.1 Table of doublets above threshold. Here E_1 and E_2 are the energies of the two singlets, E_{th} is the energy threshold, and N_i are the number of doublets above threshold for the different combinations of E_1 and E_2 . (N_1 is assumed above threshold here.) The numbers in the parentheses indicate the number of counts gained or lost in the pileup histogram. 53

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List of Abbreviations

ppm	parts per million
ppb	parts per billion
BNL	Brookhaven National Laboratory
FNAL	Fermi National Accelerator Laboratory
SM	Standard Model
SiPM	Silicon Photo-Multiplier
Geane	Geometry and Error Propagation
Geant4	Geometry and Tracking 4
WFD	waveform digitizer
PMT	photo-multiplier tube
IBMS	inflector beam monitoring system
CBO	coherent betatron oscillation

Chapter 1

Introduction

The prevailing theory for particle physics, the Standard Model (SM), has had tremendous success in describing our universe. It has been used to predict and explain a wide variety of phenomena and particles, their properties and interactions, to great precision. However, in spite of its success in explaining nearly all experimental results, there exist unanswered questions about our universe. Some of these include the matter-antimatter asymmetry, the source of mass for the neutrinos, the existence of dark matter, and an inability to fully incorporate our best theory of gravitation. Many particle physics experiments are being devised and conducted around the world in order to shed light on these questions and improve our understanding of reality. One such particular experiment is the Fermilab Muon $g - 2$ Experiment (E989) underway at the Fermi National Accelerator Laboratory (FNAL) located in Batavia, Illinois.

I have been a part of the E989 experiment since I began my graduate degree six years ago. Three years ago I moved from Boston to Batavia to be where the action is. This dissertation will describe in detail the work which I have done for the experiment. Chapter 1 will provide experimental and theoretical background to the experiment, as well its motivation. Chapter 2 will describe the experimental principle. Chapter 3 will describe the magnetic field portion of the experiment, and magnetic field simulations I conducted. Chapter 4 will describe the straw tracking detectors and their measurements, including the track fitting I wrote. Chapter 5 will describe

the frequency measurement portion of the experiment, and detail the analysis results from data taken in the first half of 2018. Finally, Chapter 6 will conclude the thesis and the important results contained within.

1.1 Magnetic Dipole Moments

In order to understand the purpose of the Fermilab Muon $g - 2$ Experiment, first we need to understand what the g in $g - 2$ is, which is what the experiment is measuring. All particles have intrinsic properties which describe those particles. One of those properties is the so called magnetic dipole moment. This property of a particle is related to its spin through the equation

$$\vec{\mu} = g \frac{Qe}{2m} \vec{s}, \quad (1.1)$$

where $\vec{\mu}$ is the magnetic dipole moment of a particle, \vec{s} is its spin vector, m is its mass, e is the elementary charge, $Q = \pm 1$, and g is the so called "g-factor". (Here and following c and \hbar have been set to 1.) g is some measureable and predictable constant, which as shown in Equation 1.1 simply relates the magnetic moment of a particle to its spin angular momentum.

In a Dirac theory, g is equal to 2 for spin-1/2 particles with no internal structure [1]. See Appendix A for a derivation of this result. It turns out however, that g is not quite equal to 2 even for these types of simple particles. Motivated by early experimental discrepancies such as the measurements of the hyperfine structure in hydrogen [2], in 1948 Schwinger calculated the first "radiative correction" to the electron magnetic moment [3]. In a quantum field theory, interactions of the particle with virtual particles in loops will contribute to the value of g . In this context it is

nicer to recast the magnetic moment formula as

$$\vec{\mu} = 2(1 + a) \frac{Qe}{2m} \vec{s},$$

$$a = \frac{g - 2}{2}, \quad (1.2)$$

where a is called the "anomalous" part of the magnetic moment, and contains all higher order corrections. The first correction calculated to a by Schwinger was $a = \alpha/2\pi \approx 0.00116$, where α is the fine structure constant. By measuring a , the SM theory (and extensions to it) can be tested. The measurement of the anomalous piece of the muon is indeed where the Fermilab Muon $g - 2$ Experiment gets its name.

1.1.1 Why the muon and not the electron?

The magnetic moment of the electron has been measured extraordinarily precisely, to .28 parts per trillion (ppt) [4]. It has been used to test SM theory extensively, and specifically the QED portion of it. Because the electron is so light, the contributions to a_e come almost exclusively from QED. This is because the various contributions from both the weak and hadronic sectors contain couplings which depend on the mass of the particle squared. It is for this reason that the magnetic moment of the muon is such an interesting quantity to measure. Since the muon is approximately 200 times heavier than the electron, the sensitivity of a_μ to radiative corrections is approximately 4000 times greater than that for the electron.

1.2 Standard Model Contributions to a_μ

Before experimental results have any real meaning, they need a theory with which to compare. The latest theoretical predictions for the muon magnetic moment will be presented here. The contributions to a_μ can be summed from separate pieces relating to different parts of the SM. These include the quantum-electrodynamics (QED)

corrections purely from other leptons and photons, the electroweak (EW) corrections from interactions with the weak force bosons W^\pm and Z^0 , and the hadronic corrections from interactions with hadrons:

$$a_\mu^{\text{SM}} = a_\mu^{\text{QED}} + a_\mu^{\text{EW}} + a_\mu^{\text{Had}} \quad (1.3)$$

The QED corrections are very well understood and have been calculated to very high order. The electroweak corrections are harder to calculate, but due to the heavy masses of the force carrying bosons and by extension weaker couplings, are only required to be known up to two loop level. And indeed they are. Finally, the hadronic contributions are where most of the theoretical uncertainties lie, and where most active work is going on. Indeed there are also some data-driven approaches to the hadronic contributions which are interesting and will be discussed in a following section.

-just put words down on the page, keep it short and simple, I can go back and add detail where necessary, I'll have to do a lot of editing anyways -dont worry about the ordering of sections exactly either -check Saskia's thesis for the latest theoretical results

- see papers cited in my HEP2 class paper - and then look for new ones - Fred Jegerlehner's book

1.2.1 QED

The QED contributions to a_μ stem solely from loops with virtual leptons and photons. They have been calculated up to 5 loop levels from over 12,000 Feynman diagrams [5, 6]. The first couple of diagrams including the Dirac $g = 2$ and Schwinger diagrams are shown in Figure 1.1. The value is

$$a_\mu^{\text{QED}} = (11658471.8971 \pm 0.007) \times 10^{10}. \quad (1.4)$$



Figure 1.1: The first of many QED diagrams contributing to a . B is a magnetic field. Feynman diagrams made with [7, 8].

Over 99% of the value of a_μ comes from the QED parts, but the error is much smaller than the other contributions and than the experimental uncertainty.

1.2.2 Electroweak

The electroweak contributions to a_μ are known to two loop level, with some 3 loop parts estimated. The contributions stem from couplings with the heavy weak gauge bosons. Because the couplings to these heavy bosons are small, the electroweak contributions to a_μ are small. The different one loop diagrams and an example two loop diagram are shown in Figure 1.2. The value of the electroweak contributions is

$$a_\mu^{\text{EW}} = (15.12 \pm 0.1) \times 10^{10}. \quad (1.5)$$

with improvements having been made recently [9, 10]. Again the error on these contributions are small compared to the hadronic contributions discussed next, as well as the experimental uncertainties.

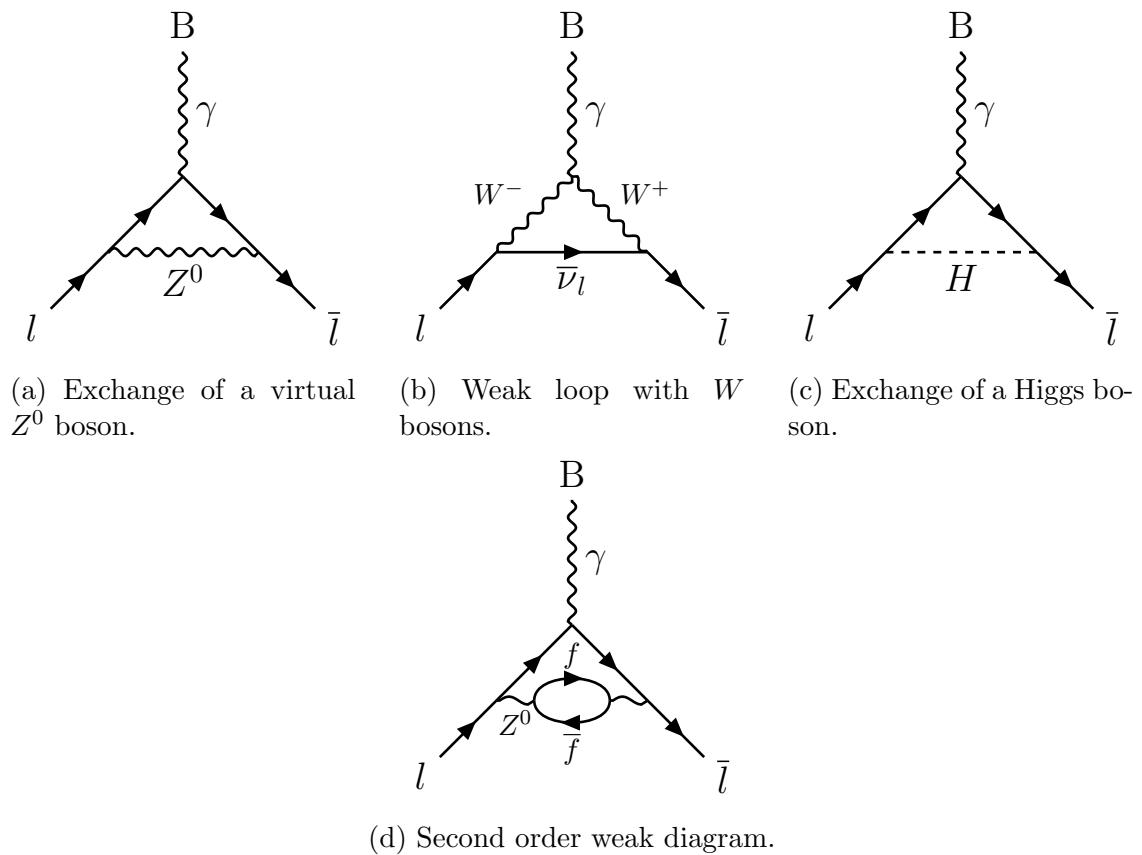


Figure 1.2: First order (and one second) weak diagrams contributing to a . B is a magnetic field. Feynman diagrams made with [7, 8].

1.2.3 Hadronic

The hadronic contributions to a_μ stem from interactions with hadrons. Because they cannot be calculated perturbatively at low energies due to the QCD nature of these particles, these calculations comprise the dominant uncertainty in the SM calculation, and make their error estimation extra important when comparing to experiment. These contributions can be separated into two parts.

Hadronic Vacuum Polarization

The first of these hadronic contribution parts is the hadronic vacuum polarization part (HVP), the first order diagram of which is shown in Figure 1.3a. There are two main prescriptions for calculating these contributions. The first is to relate the virtual hadron production within these loops to real hadron production through dispersion theory [11]. The leading order contribution can be written as

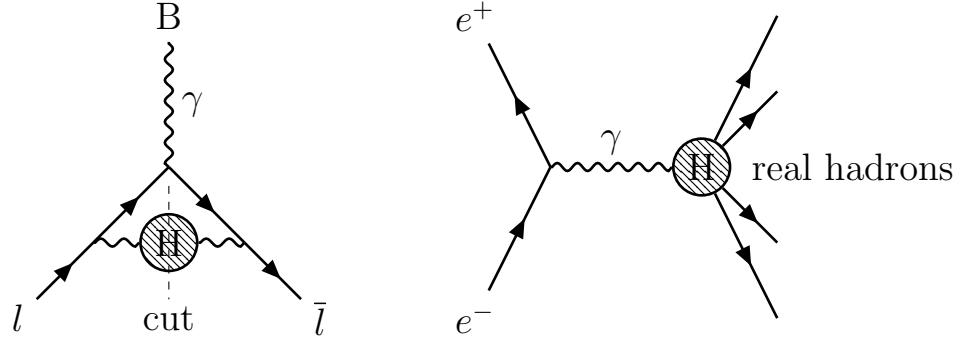
$$a_\mu^{\text{HVP;LO}} = \left(\frac{\alpha m_\mu}{3\pi} \right)^2 \int_{m_\pi^2}^{\infty} \frac{ds}{s^2} K(s) R(s) \quad (1.6)$$

where $K(s)$ is some kinematic factor, and $R(s)$ is a ratio of cross-sections

$$R(s) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}. \quad (1.7)$$

The cross-section data for this relation has been measured in parts by various experiments, including KLOE, BaBar, and BESIII [12, 13, 14]. (Name other experiments? (Belle, VEPP-2000) would then need to cite them somehow.) The analysis by Reference [15] gives results as

$$\begin{aligned} a_\mu^{\text{HVP;LO}} &= (693.26 \pm 2.46) \times 10^{10}, \\ a_\mu^{\text{HVP;NLO}} &= (-9.82 \pm 0.04) \times 10^{10}, \end{aligned} \quad (1.8)$$



(a) The first order HVP diagram. The blob H in the middle indicates any combination of hadrons which satisfy the Feynman rules. By making a 'cut' at the virtual hadrons part, this diagram can be related to the one on the right.

(b) The Feynman diagram for electron positron annihilation to hadrons. This can be related to the diagram on the left.

Figure 1.3: Feynman diagrams made with [7, 8].

where $a_\mu^{\text{HVP;NLO}}$ is the next to leading order calculation. This calculation is consistent with Reference [16].

-should I mention the optical theorem?

The second prescription is to lattice ...

Hadronic Light-by-Light

The second of these hadronic contributions parts is a higher order 4 photon blahblah

1.2.4 BSM

-why the muon mdm is important to measure

-talk about chiral flipping stuff

1.3 Background / experiment history / new experiment

It has the goal of measuring the magnetic moment of the muon, proportional to the g in $g - 2$, to high precision in order to compare to SM theoretical predictions. Because

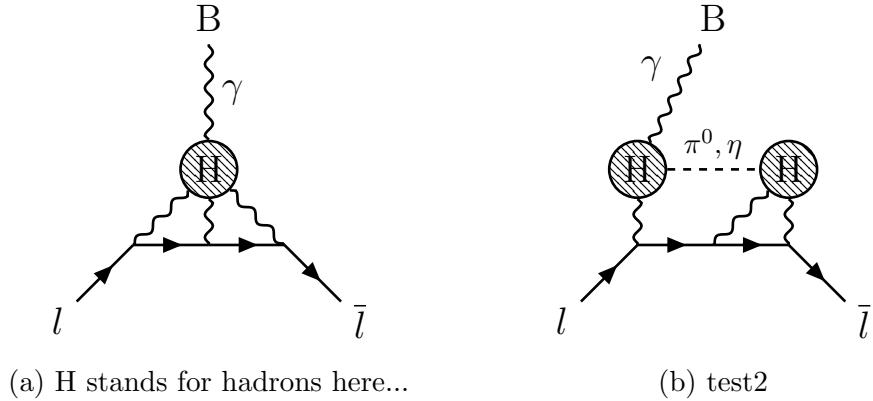


Figure 1.4: Feynman diagrams made with [7, 8].

the magnetic moment of particles couple to all existing particles, known or unknown, (source this? reference to a later section?) this provides an avenue through which theories might be constrained, and new physics narrowed down. Indeed this experiment is the latest in a line of such experiments which have measured the magnetic moment of the muon over the past several decades, the last of which measured the magnetic moment of the muon to .54 parts per million (ppm) at Brookhaven National Laboratory (BNL) in 2001 [17].

The previous $g - 2$ experiment at BNL measured a discrepancy in the magnetic moment of the muon between theory and experiment with a 2.2 - 2.7 standard deviation. (Cite the final report again?) That disagreement has since grown above 3σ [15], depending on the theoretical analysis approaches used.

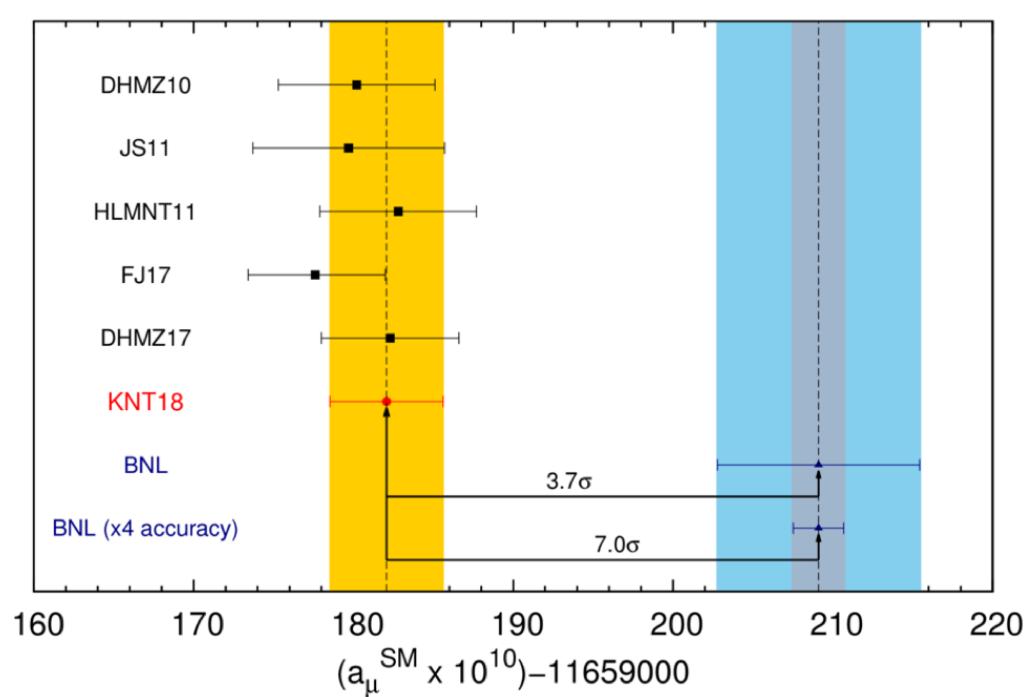


Figure 1.5: any figures that are directly lifted from someone else's work needs to be cited everytime they're used I believe, even if I cite that work in the text somewhere - [15]

Chapter 2

Muon g-2 at Fermilab, E989

2.1 Principle Technique

In a dipole magnetic field, particles will orbit at the cyclotron frequency

$$\omega_c = -\frac{Qe}{\gamma m}B, \quad (2.1)$$

and their spins will turn at the precession frequency

$$\omega_s = -g\frac{Qe}{2m}B - (1 - \gamma)\frac{Qe}{\gamma m}B, \quad (2.2)$$

where $Q = \pm 1$ and $e > 0$. The difference between these two frequencies gives

$$\omega_a = \omega_s - \omega_c = -\frac{g - 2}{2}\frac{Qe}{m}B = -a\frac{Qe}{m}B, \quad (2.3)$$

a measurable frequency that is directly proportional to the property of significance, the anomaly a . By measuring the spin difference frequency for muons and the magnetic field B , a_μ can be determined. In the presence of an electric field, which is useful in storing the muon beam within a dipole magnetic field, this expands to

$$\vec{\omega}_a = -\frac{Qe}{m}[a_\mu \vec{B} - (a_\mu - \frac{1}{\gamma^2 - 1})(\vec{\beta} \times \vec{E})], \quad (2.4)$$

where now the measurable quantities are vector quantities. Finally, for realistic cases of muon momentum which is non-orthogonal to the magnetic field, the spin difference

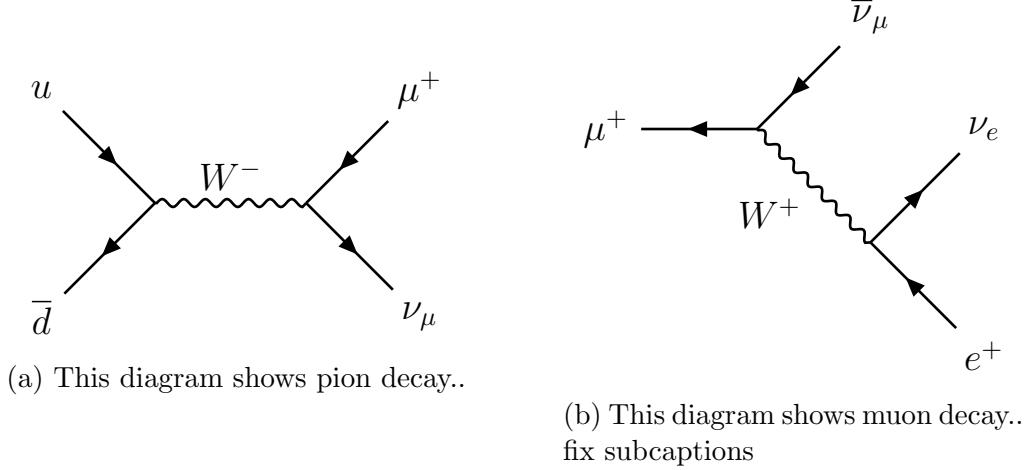


Figure 2.1: These diagrams show pion and muon decay...

frequency becomes

$$\vec{\omega}_a = -\frac{Qe}{m} [a_\mu \vec{B} - a_\mu \left(\frac{\gamma}{\gamma + 1}\right) (\vec{\beta} \cdot \vec{B}) \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1}\right) (\vec{\beta} \times \vec{E})]. \quad (2.5)$$

If the motion of the muons is largely perpendicular to the magnetic field, then the second term is small and can be corrected for. If the particles have a momentum of approximately 3.09 GeV/c, the so called “magic momentum,” then the third term is small and can be corrected for. These will be talked about later.

In order to measure the spin difference frequency of the muon, a clever technique is used. Decay muons in the pion rest frame are 100% polarized due to conservation of angular momentum and the fact that the decay neutrino must have a specific helicity. Within a pion beam then the highest and lowest energy decay muons are polarized. Muons will decay to positrons with a lifetime of about 2.2 μ s, and the positrons with the highest energies will be correlated with the muon spin, a so called “self-analyzing” decay. The single available decay state for a maximum energy positron illustrates this in Figure ???. Thus, by aquiring a large sample of polarized muons and injecting them into a storage ring

Muon decay in the rest frame

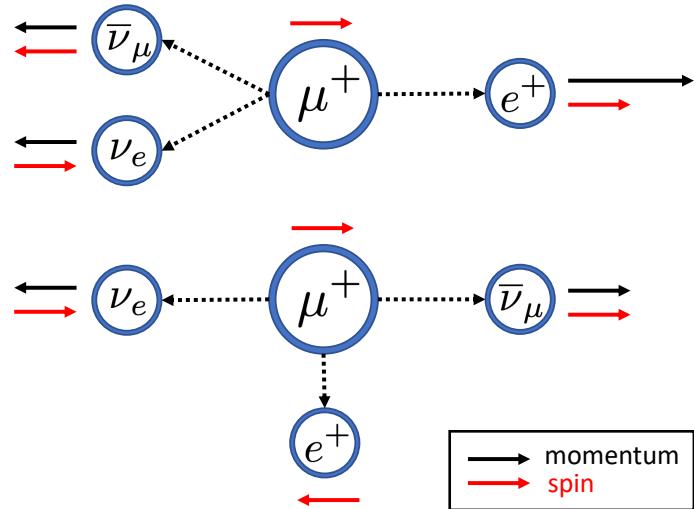


Figure 2.2: improve this caption! The single available decay state for maximum energy decay positrons. Due to the conservation of angular momentum and the single possible helicity states of the decay neutrino and anti-neutrino, the spin of the decay positron is exactly equal to the spin of the muon at the time of the decay.

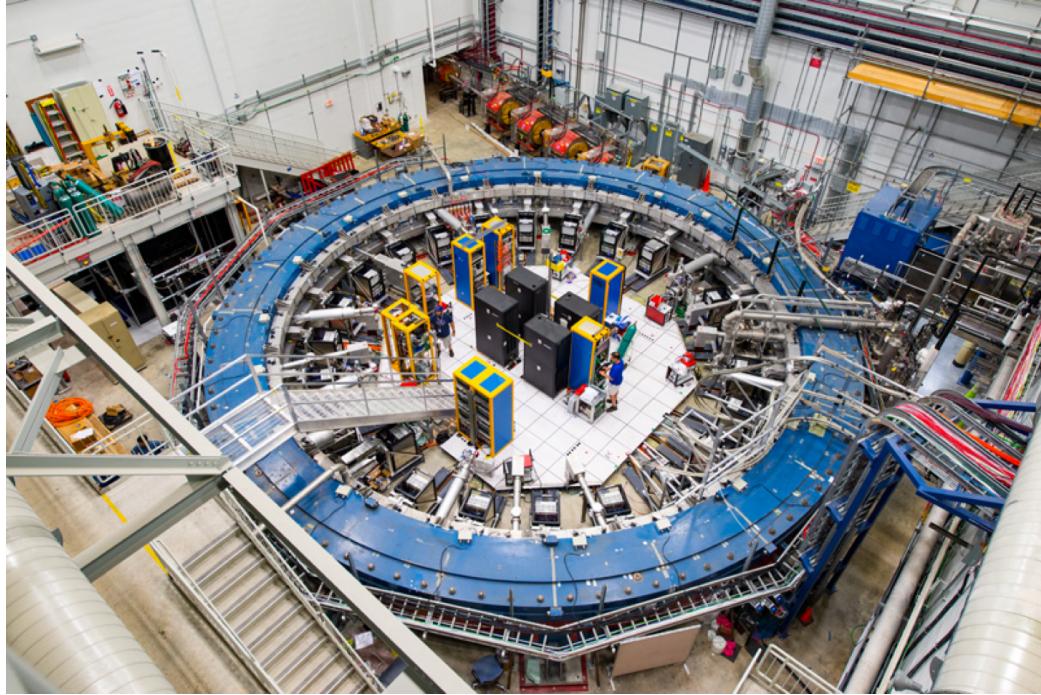


Figure 2·3: clean up and possibly replace

-explain the physics -explain how we get at the physics with our ring and detectors
 -parity violation -actually write out the decay states before explaining some things -
 well shouldn't these have been talked about before?? maybe not -decay probabilities
 and all that -don't measure all decay positrons -By injecting a large ensemble of
 muons and -by measuring a subset of ensemble of muons.... -Careful with spin vs
 polarization

2.2 Accelerator

[18]

2.3 Injection

the inflector

2.4 Storage

kickers and quads

2.5 Detector Systems

-how should I order these sections? perhaps by what the beam sees in order?

2.5.1 T0

In order to align the decay positron spectra in time from fill to fill, an entrance “T0” detector is used. It is made up of a scintillating paddle connected to two (2-inch Hammatsu?) PMTs (photo-multiplier tubes), and is placed just on the outside of the ring before the inflector. See Figure 2·4. One of the PMTs serves as the actual T0 counter which provides an average time of the incoming beam with which to align the fills. The second paddle acts as a magnitude counter to serve as a proxy for the fill intensity. Together they provide a measure of the injected beam profile in time. The varying intensity and bunch shapes can be seen as shown in Figure 2·5.

-might use the same WFD as Calo 1... -cite all this, with the TDR or hopefully something else [19] - can possibly just cite DocDBs as well - need to figure that out -double check all this - just wrote from memory, also add detail - might want to just explain things better too, “one PMT with low pe stats and one with high pe stats” and how that works in the way I’ve already described above - ND (neutral density filters) one at 1% one at 10% -if I use p.e. include it in the abbreviations -get a picture of a T0 output from the midas page, also possibly comparing various bunches -docDB 10911 for details -docDB 10162 for pictures by Hannah

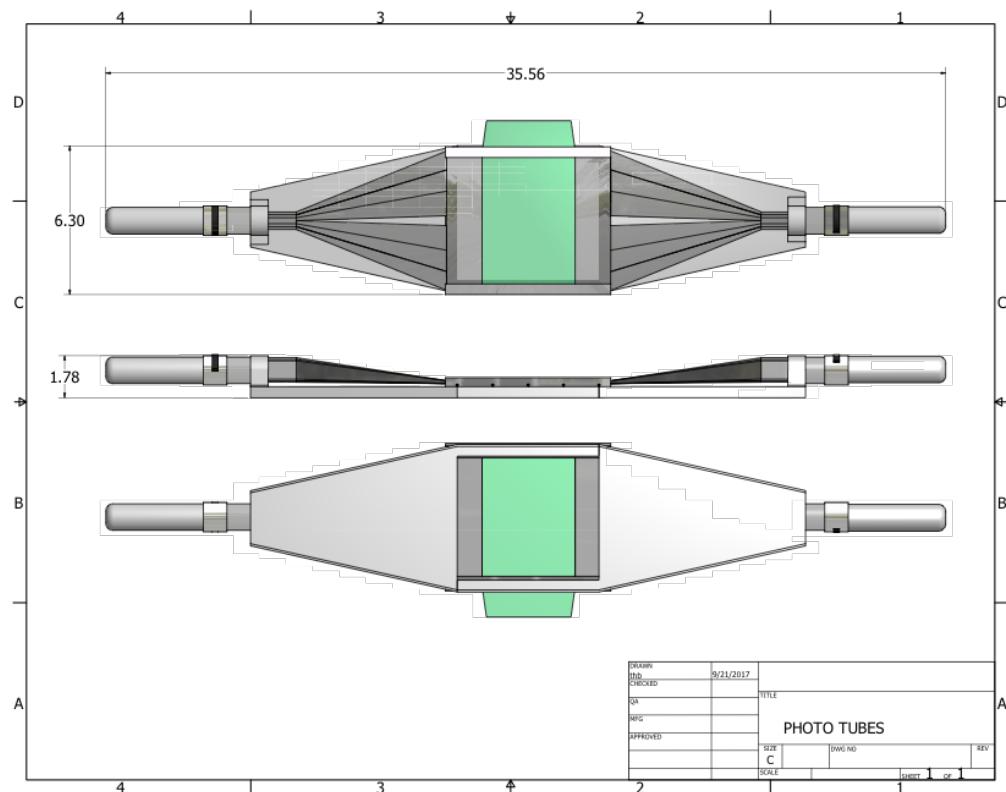


Figure 2.4: Shown is the T0 entrance counter. In the center in green is the scintillator, which connects with light guides to PMTs on the left and right. Each PMT has a separate ND (what's this?) filter to modify the light output into each PMT to configure one more for fill timing, and the other for fill intensity.

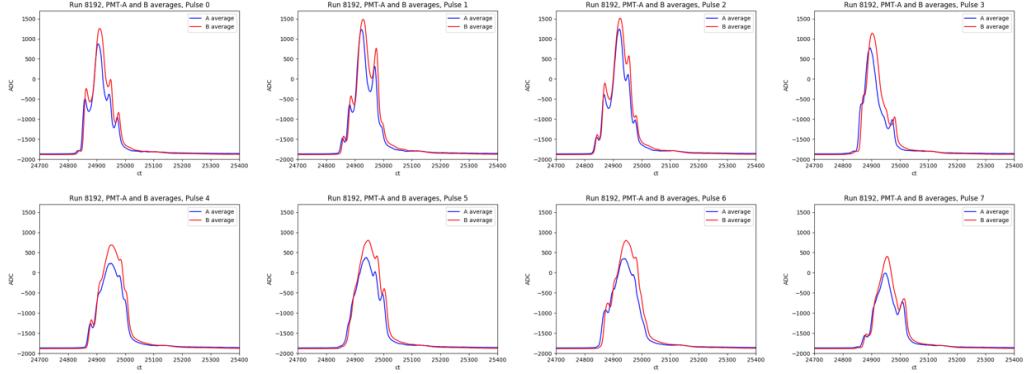


Figure 2.5: Shown are time profiles for the two PMTs (A and B) in the T0 counter for the 8 pulses we receive in an accelerator cycle/bunch/what. (careful here - not quite right) Each profile is an average of 100 such profiles. The x axis is in clock ticks (ct), each of which is 1.25 ns. (I will probably explain cts before this somewhere right?) PMT A is the low stats p.e. phototube and PMT B is the high stats p.e. phototube. Picture is too small, might want to rotate or replace it with something.

2.5.2 IBMS

There are two scintillating fiber detectors which monitor the beam as it passes through the inflector, the so-called inflector beam monitoring system (IBMS). See Figure 2.6. These devices are placed at the outside of the magnet yoke before injection into the back hole of the magnet, and at the entrance to the inflector. A third device is planned to be at or near the downstream end of the inflector (fix this later on). See Figure 2.7. These devices are used to verify the beam optics tune in the muon injection through the inflector and continuously diagnose beam properties as a handle on systematic problems. Probably cite something in this section - not the TDR since the IBMS was not included in the TDR.

-justification document - DocDB 2722 -DocDB 10944

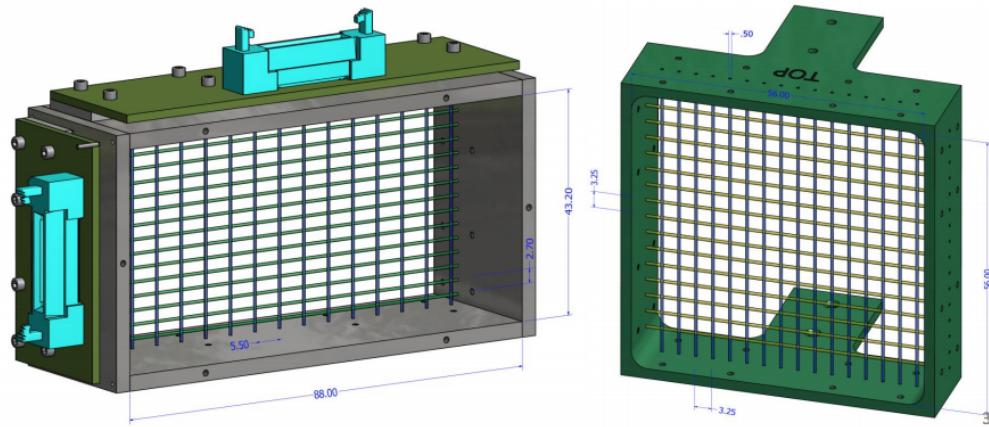


Figure 2·6: Simulation models of the IBMS 1 and 2 detectors. Scintillating fibers form an array which detect particles as they pass through them.



Figure 2·7: The positions of IBMS 1, 2, and the planned 3rd system are shown with respect to the vacuum chamber and inflector.

2.5.3 Fiber Harps

There are four scintillating fiber detectors located within the ring, two at the 180 degree position which measure the vertical and horizontal directions respectively, and similarly at the 270 degree position. One of these devices is shown in Figure 2·8. These ‘fiber harps’ serve to monitor beam properties at all times during a fill, including just after injection and during scraping. This is especially useful as a diagnostic tool as well as a measurement of the central beam position and CBO properties (have I talked about this yet?). Because the fiber harps are a destructive measurement of the beam, they are retractable and are only inserted in special runs in order to make said beam measurements. When they are inserted however they provide a wealth of useful information regarding the beam, as shown in Figure 2·9. Probably cite something in this section.

-DocDB 8366 -maybe find another docb/source

2.5.4 Calorimeters

Electromagnetic calorimeters measure the times and energies of decay positrons as they curl inward from the storage region. There are 24 calorimeters located symmetrically around the inside of the ring in close proximity to the vacuum chamber, as shown in Figure 2·10. They lie close to the storage region in order to measure a large fraction of the total number of observable decay positrons, including the high energy decay positrons which curl inward only slightly more than the muons themselves do. (max acceptance)

Each calorimeter consists of 54 channels of PbF_2 crystals arrayed in a 6 high by 9 wide array, which measure Cerenkov light emitted by the incident positrons as they pass through the crystals [20]. (Picture of single calo and its crystals here.) Cerenkov light is naturally fast which improves timing resolution of the incoming hits. Each

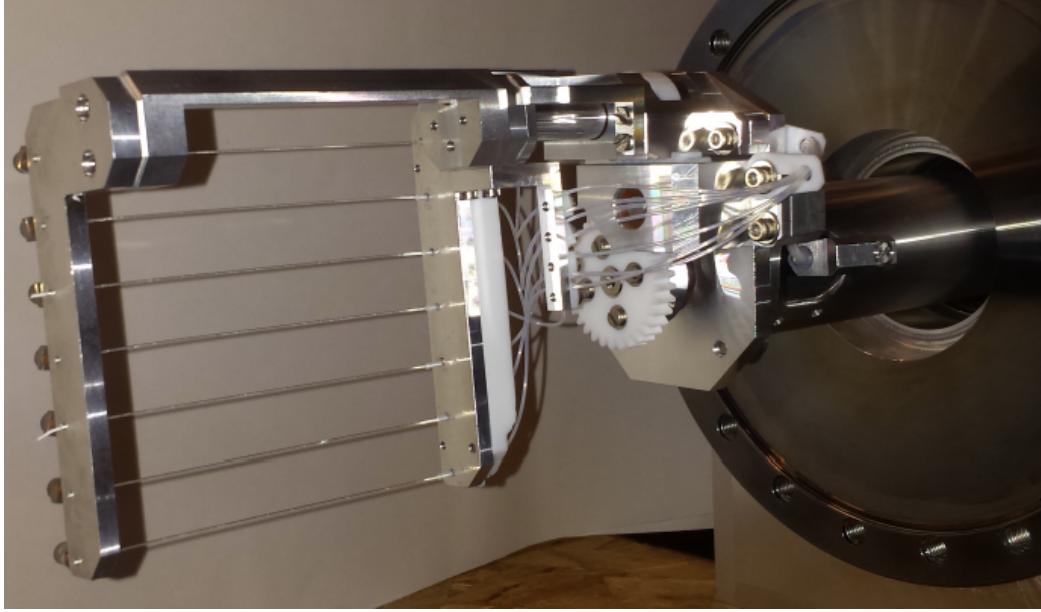


Figure 2.8: Picture of one of the fiber harps. A row of 7 scintillating fibers measures the beam intensity as a function of time at various vertical or horizontal components depending on which fiber harp is inserted.

crystal is $2.5 \times 2.5 \times 14 \text{ cm}^3$ and is wrapped in black Tedlar® foil to reduce light transmission between crystals and improve position reconstruction, as well as reduce pulse width [21]. The light is read out by large area silicon photo-multiplier (SiPM) sensors. Each calorimeter sits on a board extending out from a cart containing the electronic which power the calorimeters and read out the data, as shown in Figure ?? This is to relocate magnetic material away from the field region to avoid perturbing the field and to remove sensitive electronics from the decay path region. Similarly, due to the close proximity to the storage region and by extension the magnetic field, the calorimeters, the encapsulating material, and the SiPMs are made from non-magnetic material.

In order to determine a_μ to the precision goal, there are modest requirements on the performance of the calorimeters. They must have a relative energy resolution of better than 5% at 2 GeV, in order for proper event selection [19]. They must have

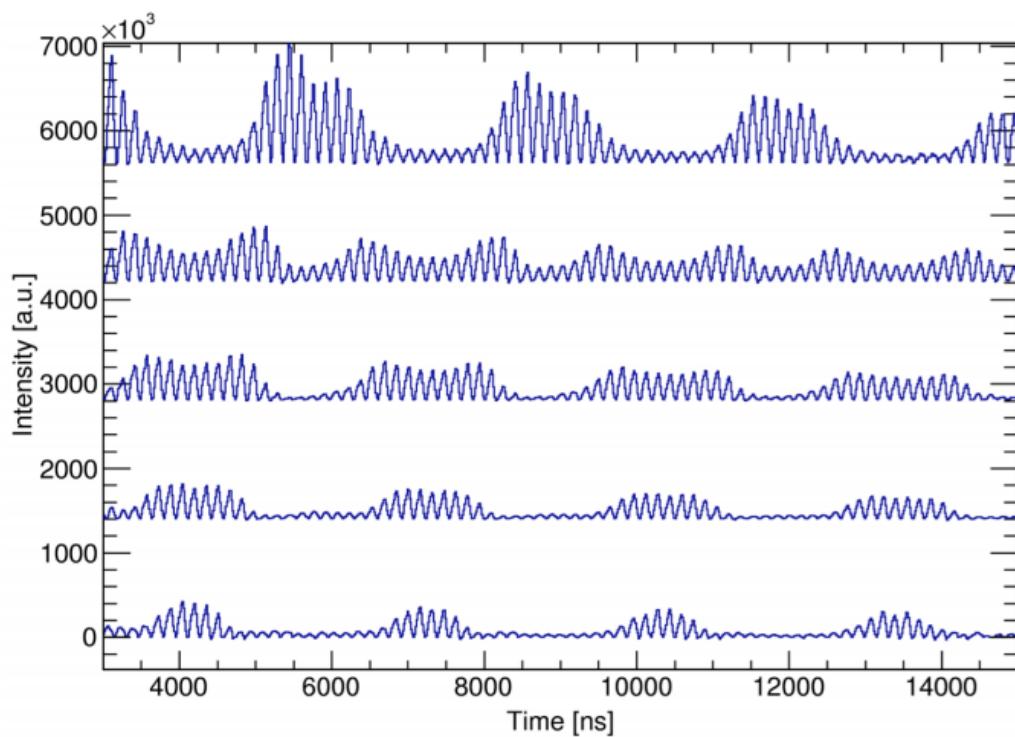


Figure 2.9: The fiber harps measure beam intensity as a function of time and fiber number. The middle spectrum is the output from the middle fiber, the upper spectra being fibers which are at successively greater radial positions, and the lower spectra being fibers which are at successively lower radial positions. The fast oscillations of the cyclotron frequency can be seen along with the slower oscillations of the CBO. This particular set of data was at a time or run when more beam was on the outside part of the storage region. (This last part is conjecture somewhat - maybe I should leave it out. I might honestly also want to leave out this whole picture, not sure..)

a timing resolution of better than 100 ps for positrons with greater than 100 MeV energy. The calorimeters must be able to resolve multiple incoming hits through temporal or spatial separation at 100% efficiency for time separations of greater than 5 ns in order to reduce the pileup systematic error due to the high rate. Finally, the gain of the measured hits must be stable to $< 0.1\%$ over a 200 μs time period within a fill, and unaffected by a pulse arriving in the same channel a few nanoseconds earlier. The long term gain stability over a time period of order seconds must be $< 1\%$. The SiPMs chosen satisfy these requirements (as well as the physical ones), due to their fast rise times and consistent pulse shapes. (Should this part be here? Or elsewhere? Before perhaps? Might need to update the gain numbers...)

(I've condensed quite a bit this section from the TDR - is that okay?)

A 12 bit waveform digitizer (WFD) samples each photodetector channel at a rate of 800 MSPS with a 1 Gb memory buffer and the data are transferred to a bank of GPU processors for on-line data processing [22]. The timing resolution of these WFDs is < 50 ps for most pulse amplitudes.

Should I go into the calorimeter DAQ here? Or make that it's own section - how uniform are the DAQ systems for each detector system? Are the same sorts of modules and crates and WFDs used or what? If so it's own section might make sense, otherwise probably include small pieces at the end of each individual detector about the daq system

2.5.5 Laser Calibration System

In order to satisfy the gain requirements, a laser calibration system was put into place in order to monitor the SiPM responses over short times (fill level), and long term (days to years). By comparing a known signal to the SiPM output. At the front face of each calorimeter is a board containing prisms connected to fibers from the laser system which can be pulsed in-fill or out of fill for each crystal individually.

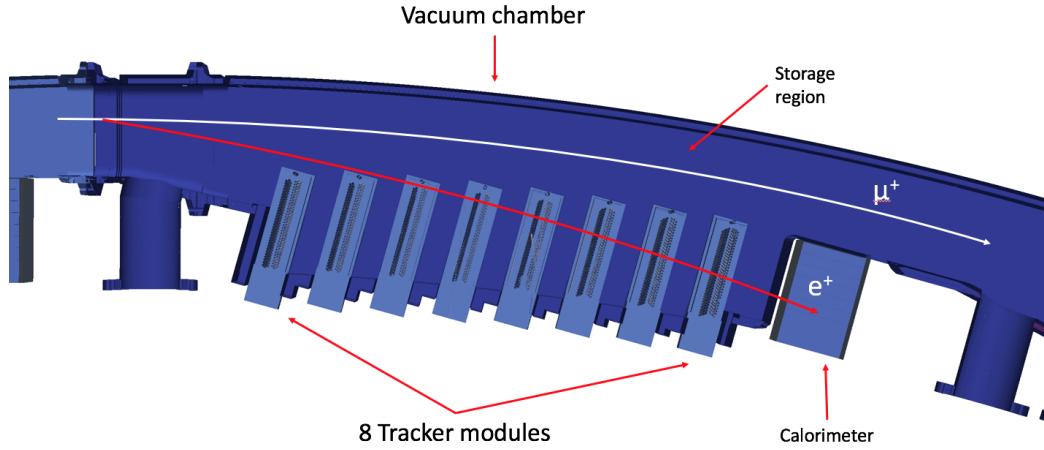


Figure 2.10: Birds eye view of a model of a vacuum chamber containing a tracker station, and the associated calorimeter. Muons pass around the ring in the center of the storage region which is contained within the bulk of the vacuum chamber. The muons decay to positrons, some of which then pass through the trackers or hit calorimeters. Each tracking station consists of 8 tracker modules.

-can pulse laser in fill, out of fill, and double pulse -mention the laser hut at all?
probably not

[23] [24] [25]

2.5.6 Straw Trackers

The Muon $g - 2$ Experiment at Fermilab uses straw tracking detectors to measure decay positron trajectories for the purpose of determining the muon beam distribution and its characteristics (and other things....). By fitting these tracks and extrapolating back to the average decay point, the beam can be characterized in a non-destructive fashion. See section blah. This is important because of the need for matching the average observed magnetic field of the decaying muons and their resulting decay positron directions which result in the ω_a frequency.

The trackers are also useful for determining general beam diagnostics as well as the pitch correction and to a lesser extent the electric field correction (careful here).

Cross-checking separately for pileup removal, hit verification, etc. is a powerful tool. Combining them in order to provide the muon distribution that the calorimeters directly see for the ω_a calculation is perhaps the most important role of the tracker. With three trackers, approximately 5% of decaying muons will result in measureable positron tracks assuming no pileup in the tracker, many of which do not hit the nearest calorimeter.

Each tracker module consists of 4 layers of 32 straws with a stereo angle of 7.5 degrees, the first two “U” layers oriented with the tops of the straws at a greater radial position, and the second two “V” layers oriented with the bottoms of the straws at a greater radial position. A tracking module is shown in Figure 2.11. There are 2 tracker stations located in front of calorimeters 13 and 19, or at approximately 180 and 270 degrees counting clockwise from the top most point of the ring where the inflector resides. Figure ?? shows this. (A third station sits empty after the inflector.) Each station consists of 8 tracking modules arranged in a staircase pattern that follows the curvature of the ring as seen in Figure 2.12.

In order to reduce the amount of multiple scattering within the straw tracking chambers as particles pass through them, the material of the straw trackers is minimized. Each straw is made of mylar foil, within which a $25 \mu\text{m}$ radius tungsten wire resides, and is filled with Argon-Ethane gas [**something**]. Fast moving particles ionize the gas as they pass through it, and the resulting ions are drawn to the wire which is held at high voltage. When they reach the wire (and the mylar) a signal can be read out which tells us that a particle was seen to pass through the straw. By combining many such signals in a brief time span, we are able to construct tracks of incident particles. (See section blah.) The resolution of hits within the straws is approximately $150 \mu\text{m}$ [**something**].

The signals of the straws are read out through...

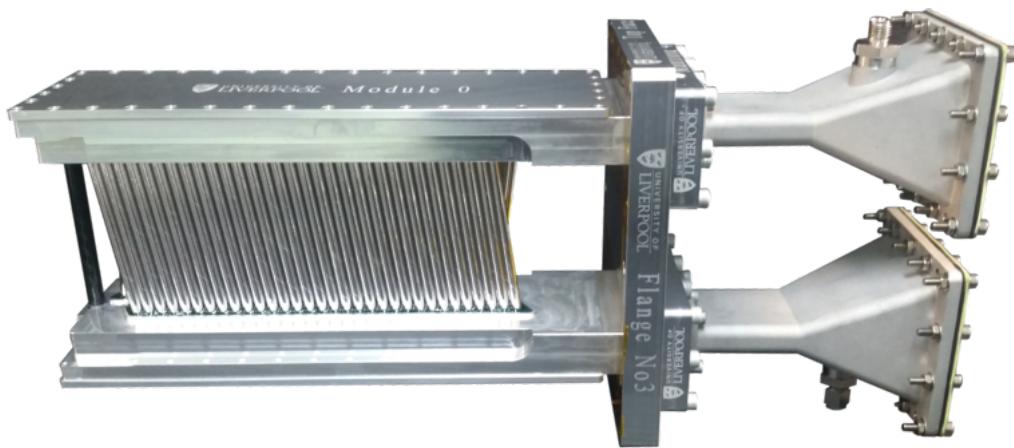


Figure 2.11: Shown is a picture of one of the many tracking modules used in the Muon $g - 2$ experiment. The first layer of straws with a stereo angle of 7.5 degrees can be seen, with the other 3 straw layers hiding behind it. The beam direction is roughly into the page in this picture, to the left of the end of the module, and this view is what the decay positrons will see.

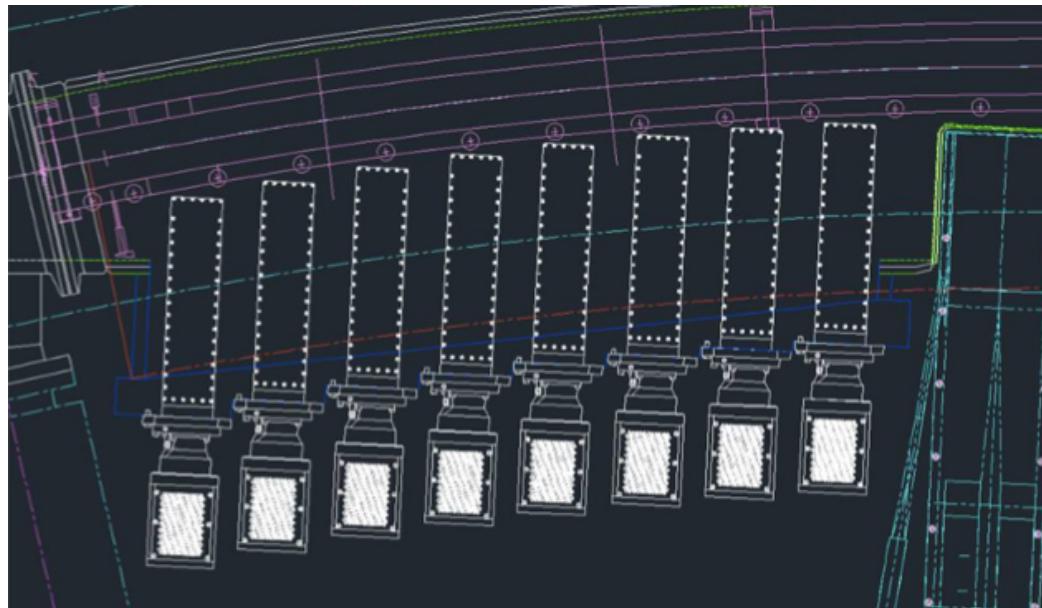


Figure 2.12: Tracker modules are arranged in the shown staircase pattern. In green and dark blue is the edge of the vacuum chamber (where the dark blue identifies the modification that was made to the old vacuum chambers), and it can be seen that vacuum chamber walls lie at the ends of the outside tracking modules. The position of a calorimeter can be seen in cyan at the right. The dark red spots are the locations of the outside magnet pole tips. From the shown geometry one can see that many positrons will hit either the tracker or the calorimeter but not both due to the acceptance differences.

Chapter 3

Magnetic Field Measurement

3.1 Trolley

3.2 Opera Simulations

Where does this section really go?

Chapter 4

Straw Tracking Analysis

4.1 Straw Tracking Intro

As was talked about briefly in section 2.5.6, the straw trackers are used to provide information about the muon beam, as well as info for the calorimeters. The straw track reconstruction is performed in several stages. The “Track Finding” stage takes incoming hits and decides which hits should be grouped together to form a single track for a single positron. The “Track Fitting” stage fits the measured positions of these individual hits and forms a single track describing the trajectory of the incident positron. Finally the “Track Extrapolation” stage takes the fitted track information and extrapolates the position and momentum components to the regions of interest, namely the storage region and the calorimeter.

-see the previous section for the hardware information... -Geane (Geometry and Error Propagation)

4.2 Track Finding

4.3 Track Fitting

The Geane fitting routines originated in Fortran with the EMC collaboration, and was used in the precursor E821 experiment as well as the PANDA experiment with some success [26], [27]. (I’m not actually aware of a useful reference for it’s use in E821, and there are some other instances of its use as well in other experiments. In E821

Run 1863, SubRun 25, Event 281, Island 66, Time 683689

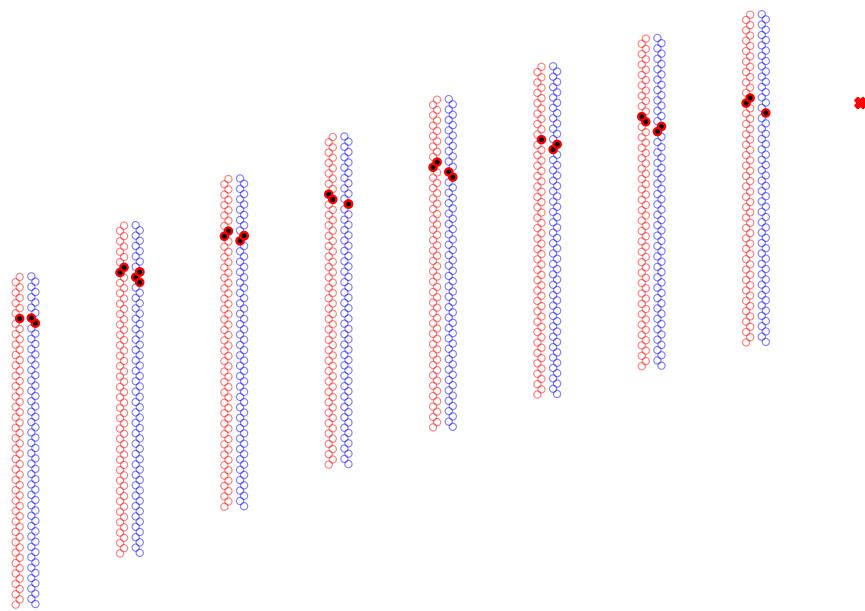


Figure 4·1: clean up and possibly replace

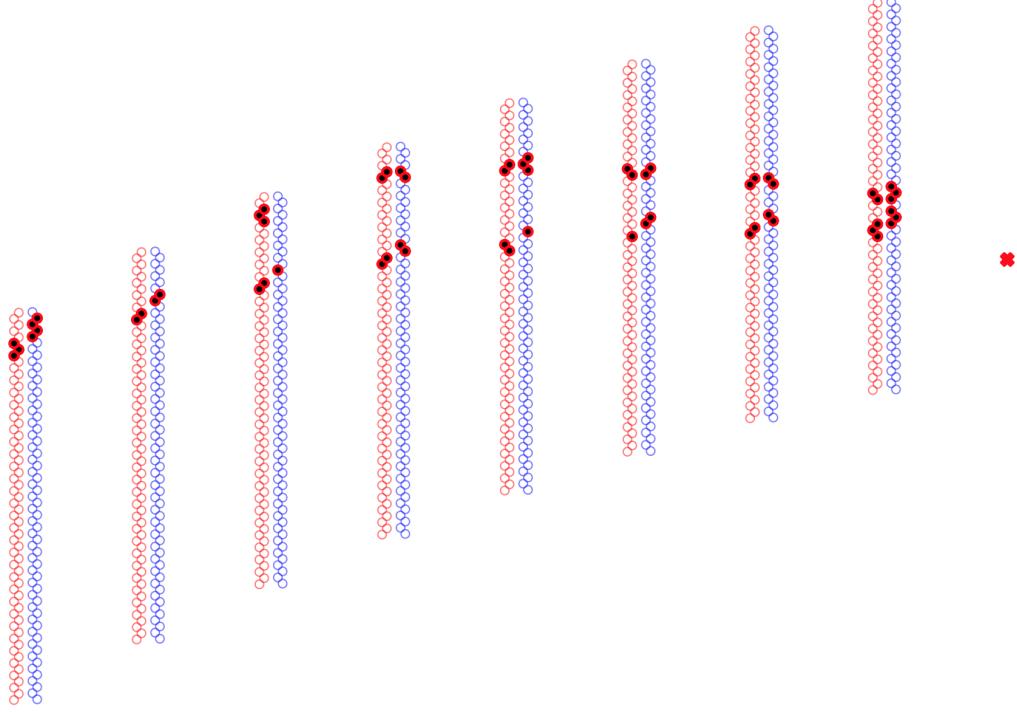


Figure 4·2: clean up and possibly replace

there was a single tracking chamber which was never put to full use.) The core error propagation routines were at some point added to Geant4 under the error_propagation directory which is included in all default installs. The tracking code strengths lie with its direct implementation and access to the Geant4 geometry and field, and its ability to handle the field inhomogeneties. The Geane fitting algorithm code which makes use of the Geant4 error propagation routines follows the structure of [26] and is detailed in the Formalism section in this paper. It is a relatively straight forward least squares global χ^2 minimization algorithm.

Because of the proximity of the trackers to the muon beam, they will lie within a region of varying magnetic field. The radial field of the trackers rises from 0 Tesla at the outer ends to roughly .3 Tesla at the inner top and bottom ends, and the vertical field drops approximately 50% from the storage dipole field of 1.451 Tesla. Shown in Figures 4·3 and 4·4 is the location of the tracker with respect to the horizontal

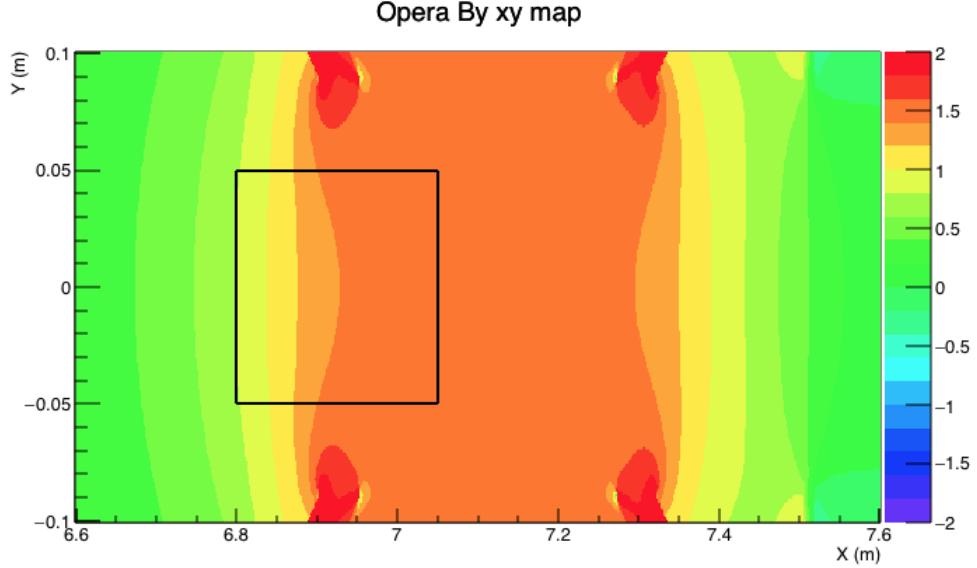


Figure 4.3: Shown is the vertical field of the $g - 2$ magnet in and around the storage region as calculated in Opera 2D. The center of the storage region lies at 7.112 m along the x axis. The black box shows the rough location of the tracker with respect to the field (size exaggerated slightly). It can be seen that there is a large inhomogeneity within the tracker space, going from left to right.

and vertical fields respectively. These large field gradients over the tracking detector region and the long extrapolation distance back to the muon decay point are special to Muon $g - 2$. This is one of the main motivations for using the Geane fitting algorithm and routines, which has direct access to the field.

4.3.1 Track Fitting Formalism

I recommend reading [26], Chapter 4 of [27], and [28] in order to best understand the fitting algorithm. However, due to the at times confusing notation, omitted equations or concepts, and differences between papers, I have attempted to summarize here the different sources and present the material in a more understandable and readable format. The implementation of the fitting algorithm into the code follows this section.

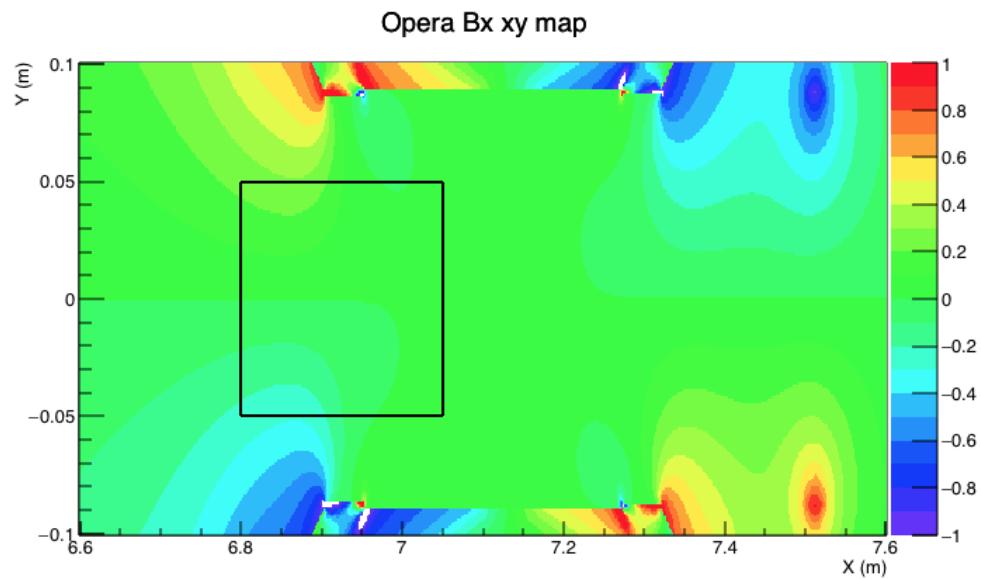


Figure 4.4: Shown is the radial field of the $g - 2$ magnet in and around the storage region as calculated in Opera 2D. The center of the storage region lies at 7.112 m along the x axis. The black box shows the rough location of the tracker with respect to the field (size exaggerated slightly). It can be seen that there is a large homogeneity at the inner upper and lower ends compared to the right center. The shape of the pole pieces and tips can readily be seen.

One can define a χ^2 for a track in the usual way by dividing the residuals of measured and predicted track parameters by their errors:

$$\chi^2 = (\vec{p} - \vec{x})^T (\sigma^{-1}) (\vec{p} - \vec{x}), \quad (4.1)$$

where \vec{p} are predicted track parameters from a fit to the measured track parameters \vec{x} , and σ is a covariance matrix of errors on the fitted parameters. The Geant4 error propagation routines can be used to determine these predicted parameters and error matrices by propagating track parameters from some initial guesses. By minimizing this χ^2 with respect to the track parameters one can then fit and improve the track. The Geant4 error propagation routines propagate particles along their average trajectories neglecting the effects of discrete processes, using a helix equation along small enough steps where the change in the magnetic field is small. The predicted parameters are then a function of path length:

$$p_l = F_{l,l_0}(p_0), \quad (4.2)$$

where the path length can be defined how one wishes. In our system we have tracker planes defined at X positions, and limit path lengths to reach those planes. (From here on the dependence on path length or X position will be neglected, in favor of using plane indices.) In tandem, error matrices describing the expected distribution in true parameters about those predicted parameters due to said discrete process are also calculated:

$$\sigma^{ij} = \langle p^i p^j \rangle - \langle p^i \rangle \cdot \langle p^j \rangle, \quad (4.3)$$

where i and j are track parameter indices. These parameter vectors are 5x1 objects defined in some track representation, as described in the Coordinate Systems section. The propagation of these parameters and error matrices are done using transport

matrices, which express the infinitesimal changes in parameters at some plane (or path length) with respect to the parameters at some previous plane (or previous path length):

$$\delta p_N = T_{N,N-1} \delta p_{N-1}, \quad (4.4)$$

$$\sigma_N = T_{N,N-1} \sigma_{N-1} T_{N,N-1}^T. \quad (4.5)$$

Said transport and error matrices are 5x5 objects since the parameter vectors are 5x1 objects as described above. The calculation of these transport matrices, as well as details on the functional form of 4.2 are shown in [29].

With parameters defined on such planes, one can define the χ^2 as:

$$\chi^2 = \sum_{i=1}^N [(p_i(p) - x_i)^T (\sigma_i^{-1}) (p_i(p) - x_i)], \quad (4.6)$$

where p_i are the average predicted parameters from some general starting parameters p . At first order one can solely include the measurement errors on parameters, which fill in the diagonals of σ_i , if random processes can be neglected. Unmeasured parameters should have measurement errors of infinity (or some large value) along the diagonals in the code, which account for the fact that residuals for unmeasured parameters do not exist. When the error matrix is inverted all rows and columns of the matrix with these large numbers will fall to 0 in the χ^2 .

In order to get the best fit track, the χ^2 should be minimized with respect to the initial track parameters p , and evaluated at some chosen or fitted parameters:

$$\frac{\partial \chi^2}{\partial p} |_{p=p'_0} = 0, \quad (4.7)$$

resulting in

$$\begin{aligned} 0 = \sum_{i=1}^N & [(\frac{\partial p_i(p)}{\partial p}|_{p=p'_0})^T (\sigma_i^{-1})(p_i(p'_0) - x_i) \\ & + (p_i(p'_0) - x_i)^T \frac{\partial(\sigma_i^{-1})}{\partial p}|_{p=p'_0} (p_i(p'_0) - x_i) \\ & + (p_i(p'_0) - x_i)^T (\sigma_i^{-1}) (\frac{\partial p_i(p)}{\partial p}|_{p=p'_0})] \end{aligned} \quad (4.8)$$

where the 1st and 3rd terms are identical, and the 2nd term is small if one assumes that the error matrix doesn't change much with respect to the starting parameters. (Fair since most of the error comes from measurement, and as long as the initial guess is decent enough such that the path length through material doesn't change appreciably from one iteration to the next.) This simplifies to:

$$\sum_{i=1}^N T_{i0}^T (\sigma_i^{-1}) (p_i(p'_0) - x_i) = 0, \quad (4.9)$$

which is just the top term with

$$T_{i0} = \frac{\partial p_i(p)}{\partial p}. \quad (4.10)$$

To solve this make the substitution

$$p_i(p'_0) = p_i(p_0) + \frac{\partial p_i(p_0)}{\partial p} \Delta p_0 = p_i(p_0) + T_{i0} \Delta p_0, \quad (4.11)$$

where p'_0 are the improved starting parameters for the next iteration calculated from the previous starting parameters p_0 , and Δp_0 are the changes in the starting parameters to improve the track. This equation can be plugged into the above if one makes the assumption that T_{i0} does not change much from one iteration to the next, which follows from the inherent nature of making small adjustments to the track in order to improve it.

After simplifying one arrives at

$$\Delta p_0 = \sigma_{p_0} \sum_{i=1}^N T_{i0}^T (\sigma_i^{-1}) (x_i - p_i(p_0)), \quad (4.12)$$

where

$$\sigma_{p_0} = \left[\sum_{i=1}^N T_{i0}^T (\sigma_i^{-1}) T_{i0} \right]^{-1}, \quad (4.13)$$

is the 5x5 covariance matrix of fitted parameters on the starting plane, whose diagonals describe the errors in the 5 track parameters on that plane and in the region close to it. (The fit does not directly return fit errors for track parameters on other planes.) Δp_0 along with χ^2 is exactly what we want to determine since that is what allows us to fit and improve the track from iteration to iteration.

However, since random processes should not be neglected for optimal tracking results, it makes more sense to return to the original χ^2 in equation 4.1, only now the included matrix and vector objects are combined into one large linear algebra equation. Instead of a sum over N 5x1 objects multiplying 5x5 error matrices, the vectors are combined into a single 5Nx1 vector multiplying a single 5Nx5N matrix. The 5x5 diagonal blocks of this large error matrix should now include the effects due to material processes as calculated in Geant from equation 4.3 as well as the measurement errors.

Because now parameters at one plane are no longer independent of the parameters at other planes, due to correlations from these random processes, it's necessary to add off-diagonal elements into the large error matrix. These 5x5 blocks come from

$$\sigma_{MN} = T_{MN} \sigma_N, \quad (4.14)$$

for the top diagonals, and the transpose for the bottom diagonals, where M and N are two separate planes within the detector. (σ_N is the error matrix on plane N

calculated from the starting plane.) This follows from equation 4.3 evaluated at plane M with respect to a path length from plane N, and not plane 0, which is equivalent to 4.14.

You can then minimize the χ^2 in the same way, only again with the matrix objects being aggregates of the per plane objects:

$$\Delta \vec{p}_0 = \sigma_{p_0} \tau^T \sigma^{-1} (\vec{x} - \vec{p}), \quad (4.15)$$

$$\sigma_{p_0} = [\tau^T \sigma^{-1} \tau]^{-1}, \quad (4.16)$$

where τ is the combined transport matrices from the individual 5x5 matrices, a 5Nx5 object.

The unmeasured parameter errors of infinity still come into play in the final calculation in the same was as before. Because however these matrix objects are very large, and the tracking must have a certain amount of speed in order to keep up with data, it is useful to reduce the size of these matrices. (It also makes things easier programming wise. Note that there are other some other ways to speed things up, specifically the banded inversion method as described in reference [28]. This method was not used in favor of getting the code working in the simpler form in the first place, but it is a possibility in the future to use this technique to speed things up even more.) It suffices to simply remove all rows and columns where said infinity values exist in the error matrix. This is mathematically equivalent to inverting the error matrix with the infinities included, which make all rows and columns where they exist go to zero. The associated unmeasured parameter rows in the residual vector and transport matrices must similarly be removed. This results in an Nx1 residual vector, NxN error matrix, 5xN combined transport matrix transpose, which multiply against the 5x5 covariance matrix out front to still result in a 5x1 fix to the starting

parameters, and a scalar χ^2 value. (Note that these element removals should be done just before the final calculation, and not higher up in the algebra, otherwise plane correlations are not properly calculated.)

By calculating the last two equations one can fit the track, acquire a χ^2 describing the degree of the fit, determine how the track parameters can be improved at the starting point, and calculate errors on those starting parameters. This algorithm can be iterated a number of times to get a best fit track until successive iterations produce no improvement, where usually 3 or 4 iterations is enough. Note that there is remarkable robustness with respect to the initial starting parameters in fitting the track. Of course if the initial starting parameters are too poor, then the fit will not converge. All of these calculations are completed within the GeaneFitter.cc file within the framework.

4.4 Track Extrapolation

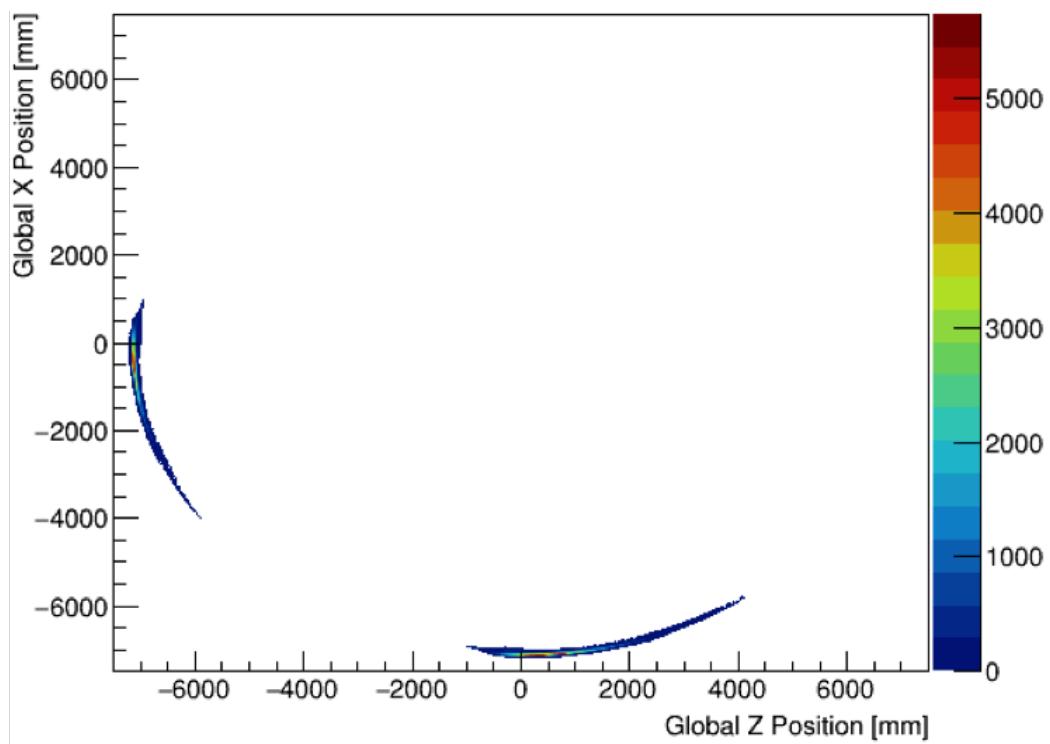


Figure 4·5: preliminary replace at some point

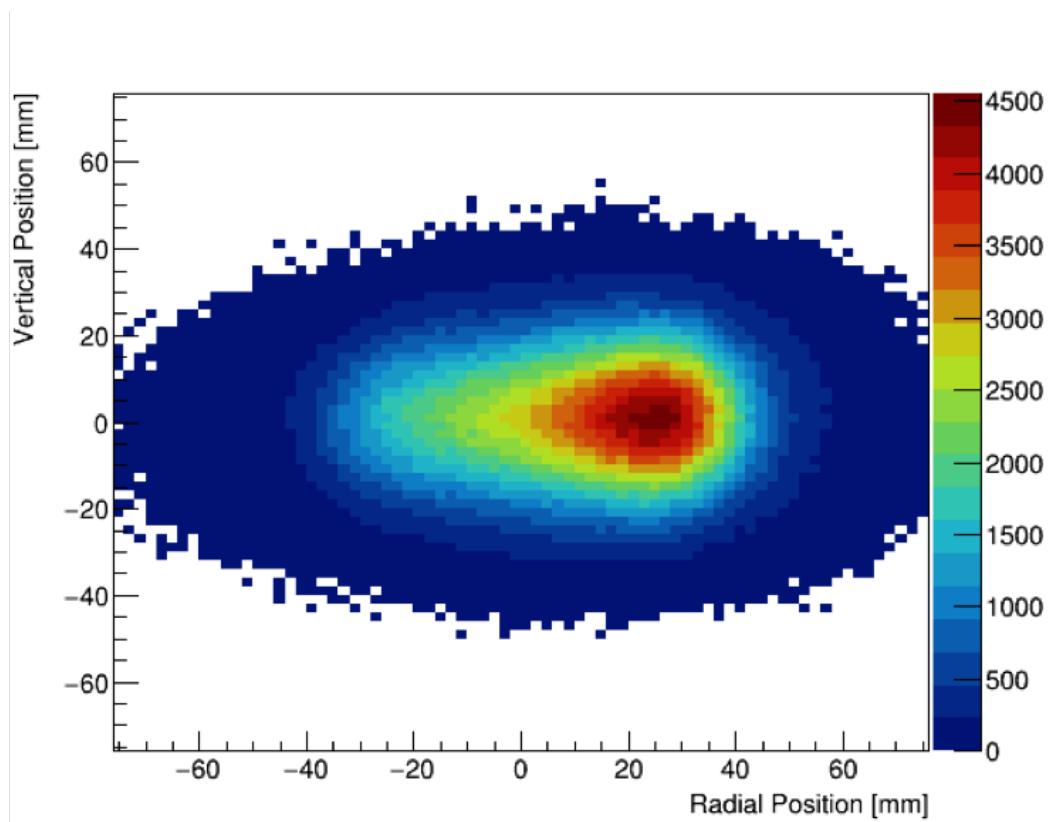


Figure 4·6: preliminary replace at some point

Chapter 5

ω_a Measurement

5.1 Data

5.2 Spectra Making

5.2.1 Clustering

5.2.2 Histogramming

5.3 Fitting

5.4 Systematic Errors

Chapter 6

Conclusion

6.1 Final Value

test4

Appendix A

g for Spin-1/2 Particles and Beyond

This was taken from my old HEP2 class report paper - go back and clean this up/improve it.

The derivation contained here is taken and simplified from Reference [30]. Starting with the Dirac equation

$$(i\gamma^\mu \partial_\mu - m - e\gamma^\mu A_\mu)\psi = 0 \quad (\text{A.1})$$

and multiplying by

$$(i\gamma^\mu \partial_\mu + m - e\gamma^\mu A_\mu)\psi = 0, \quad (\text{A.2})$$

where the sign on m is reversed, you arrive at the equation

$$[(i\partial_\mu - eA_\mu)(i\partial_\nu - eA_\nu)\gamma^\mu\gamma^\nu - m^2]\psi = 0. \quad (\text{A.3})$$

This can be split this into its symmetric and antisymmetric parts:

$$\begin{aligned} & \left(\frac{1}{4}\{i\partial_\mu - eA_\mu, i\partial_\nu - eA_\nu\}\{\gamma^\mu, \gamma^\nu\}\right. \\ & \left. + \frac{1}{4}[i\partial_\mu - eA_\mu, i\partial_\nu - eA_\nu][\gamma^\mu, \gamma^\nu] - m^2\right)\psi = 0 \end{aligned} \quad (\text{A.4})$$

Using the identities

$$\frac{i}{2}[\gamma^\mu, \gamma^\nu] = \sigma^{\mu\nu} \quad (\text{A.5})$$

and

$$\begin{aligned} & [i\partial_\mu - eA_\mu, i\partial_\nu - eA_\nu] \\ &= -ie[\partial_\mu A_\nu - \partial_\nu A_\mu] = -ieF_{\mu\nu} \end{aligned} \quad (\text{A.6})$$

where $\sigma^{\mu\nu}$ is related to the spin of the particle and $F_{\mu\nu}$ is the electromagnetic field tensor, one arrives at the form

$$((i\partial_\mu - eA_\mu)^2 - \frac{e}{2}F_{\mu\nu}\sigma^{\mu\nu} - m^2)\psi = 0. \quad (\text{A.7})$$

Expanding out the tensor objects

$$\frac{e}{2}F_{\mu\nu}\sigma^{\mu\nu} = -e \begin{pmatrix} (\vec{B} + i\vec{E}) \cdot \vec{\sigma} \\ (\vec{B} - i\vec{E}) \cdot \vec{\sigma} \end{pmatrix} \quad (\text{A.8})$$

and forming a new covariant derivative

$$\not{D}^2 = D_\mu^2 + \frac{e}{2}F_{\mu\nu}\sigma^{\mu\nu} \quad (\text{A.9})$$

where D_μ^2 is your ordinary covariant derivative, by moving to momentum space you can arrive at the equation

$$\frac{(H + eA)^2}{2m}\psi = \left(\frac{m}{2} + \frac{(\vec{p} + e\vec{A})^2}{2m}\right) - 2\frac{e}{2m}\vec{B} \cdot \vec{s} \pm i\frac{e}{m}\vec{E} \cdot \vec{s}\psi. \quad (\text{A.10})$$

Lo and behold, you have arrived at the Dirac $g = 2$ result, contained in front of the magnetic piece in the form of Equation ??.

How then does such a term change at loop level? Most generally the vertex of a particle interacting with a magnetic field through the mediation of a photon can be represented by

$$iM^\mu = \bar{u}(q_2)(f_1\gamma^\mu + f_2p^\mu + f_3q_1^\mu + f_4q_2^\mu)u(q_1) \quad (\text{A.11})$$

where q_1 and q_2 are the ingoing and outgoing four-momenta respectively, which can be constrained on-shell, and p is the four-momenta of the photon, which is off-shell. The f_i are in general contractions of four-momenta and gamma matrices. By using the Gordon identity

$$\begin{aligned} & \bar{u}(q_2)(q_1^\mu + q_2^\mu)u(q_1) \\ &= (2m)\bar{u}(q_2)\gamma^\mu u(q_1) + i\bar{u}(q_2)\sigma^{\mu\nu}(q_1^\nu - q_2^\nu)u(q_1) \end{aligned} \quad (\text{A.12})$$

any Feynman diagram can be reorganized into the form

$$iM^\mu = (-ie)\bar{u}\left[F_1\left(\frac{p^2}{m^2}\right)\gamma^\mu + \frac{i\sigma^{\mu\nu}}{2m}p_\nu F_2\left(\frac{p^2}{m^2}\right)\right]u, \quad (\text{A.13})$$

where F_1 and F_2 are form factors. One notices that the F_2 piece is reminiscent of our magnetic dipole moment form that we derived from the Dirac equation. So the problem now becomes for any Feynman diagram calculation in any theory, at any order, to solve for this F_2 to determine the contribution to the magnetic dipole moment.

Appendix B

Ratio Method Derivation

B.1 Ratio Form and Function

Consider the 5 parameter function:

$$N_5(t) = N_0 e^{-t/\tau} (1 + A \cos(\omega_a t + \phi)), \quad (\text{B.1})$$

which describes some ideal dataset in histogram format. Here ϕ will be set to zero for simplicity. Now define the variables $u_+(t)$, $u_-(t)$, $v_1(t)$, and $v_2(t)$ as

$$\begin{aligned} u_+(t) &= \frac{1}{4} N_5(t + T/2) \\ u_-(t) &= \frac{1}{4} N_5(t - T/2) \\ v_1(t) &= \frac{1}{4} N_5(t) \\ v_2(t) &= \frac{1}{4} N_5(t), \end{aligned} \quad (\text{B.2})$$

where the $1/4$ out front reflects randomly splitting the whole dataset into 4 equally weighted sub-datasets, and T is the g-2 period known to high precision, $\mathcal{O}(10^{-6})$. This corresponds to a weighting of 1:1:1:1 between the datasets. To be explicit here regarding the signs, the counts that are filled into the histogram described by u_+ have their times shifted as $t \rightarrow t - T/2$, which is what the function $N_5(t + T/2)$ describes,

and vice versa for u_- . To form the ratio define the variables:

$$\begin{aligned} U(t) &= u_+(t) + u_-(t) \\ V(t) &= v_1(t) + v_2(t) \\ R(t) &= \frac{V(t) - U(t)}{V(t) + U(t)}. \end{aligned} \tag{B.3}$$

Plugging in and dividing the common terms ($N_0 e^{-t/\tau}/4$),

$$R(t) = \frac{2(1 + A \cos(\omega_a t)) - e^{-T/2\tau}(1 + A \cos(\omega_a t + \omega_a T/2)) - e^{T/2\tau}(1 + A \cos(\omega_a t - \omega_a T/2))}{2(1 + A \cos(\omega_a t)) + e^{-T/2\tau}(1 + A \cos(\omega_a t + \omega_a T/2)) + e^{T/2\tau}(1 + A \cos(\omega_a t + \omega_a T/2))}. \tag{B.4}$$

Now set $\omega_a T/2 = \delta$, and note that T is really

$$\begin{aligned} T &= T_{guess} = \frac{2\pi}{\omega_a} + \Delta T, \\ \Delta T &= T_{guess} - T_{true}. \end{aligned} \tag{B.5}$$

Being explicit,

$$\delta = \frac{\omega_a}{2} T_{guess} = \frac{\omega_a}{2} \left(\frac{2\pi}{\omega_a} + \Delta T \right) = \pi + \pi \frac{\Delta T}{T_{true}} = \pi + \pi(\delta T), \tag{B.6}$$

and δ can be redefined as

$$\delta = \pi(\delta T), \tag{B.7}$$

by flipping the sign of any cosine terms that contain δ .

Then, using the trig identity

$$\cos(a \pm b) = \cos(a) \cos(b) \mp \sin(a) \sin(b) \tag{B.8}$$

so that

$$\begin{aligned}
\cos(\omega_a t \pm \delta) &= \cos(\omega_a t) \cos \delta \mp \sin(\omega_a t) \sin \delta \\
&\approx \cos(\omega_a t)(1 - \delta^2) \mp \sin(\omega_a t)\delta \\
&\approx \cos(\omega_a t),
\end{aligned} \tag{B.9}$$

since $\delta \sim O(10^{-5})$, the ratio becomes

$$R(t) \approx \frac{2(1 + A \cos(\omega_a t)) - (1 - A \cos(\omega_a t))(e^{-T/2\tau} + e^{T/2\tau})}{2(1 + A \cos(\omega_a t)) + (1 - A \cos(\omega_a t))(e^{-T/2\tau} + e^{T/2\tau})}. \tag{B.10}$$

Expanding

$$e^{\pm T/2\tau} = 1 \pm \frac{T}{2\tau} + \frac{1}{2} \left(\frac{T}{2\tau} \right)^2 \pm \dots, \tag{B.11}$$

replacing and simplifying,

$$R(t) \approx \frac{A \cos(\omega_a t) - C(1 - A \cos(\omega_a t))}{1 + C(1 - A \cos(\omega_a t))}, \tag{B.12}$$

where

$$C = \frac{1}{16} \left(\frac{T}{\tau} \right)^2 \approx 2.87 * 10^{-4}. \tag{B.13}$$

Using the expansion

$$f(x) = \frac{1}{1+x} = 1 - x + x^2 - \dots, \quad |x| < 1, \tag{B.14}$$

and since C is small, the denominator can be manipulated such that

$$\begin{aligned}
R(t) &\approx (A \cos(\omega_a t)) - C(1 - A \cos(\omega_a t))(1 - C(1 - A \cos(\omega_a t))) \\
&\approx A \cos(\omega_a t) - C + CA^2 \cos^2(\omega_a t),
\end{aligned} \tag{B.15}$$

after dropping terms of $\mathcal{O}(C^2)$ and higher. In practice the last term is omitted since

it has a minimal effect on the fitted value of ω_a [cite], and one arrives at

$$R(t) \approx A \cos(\omega_a t) - C, \quad (\text{B.16})$$

the conventional 3 parameter ratio function.

In order to avoid approximations one can instead weight the counts in the histograms as

$$u_+(t) : u_-(t) : v_1(t) : v_2(t) = e^{T/2\tau} : e^{-T/2\tau} : 1 : 1, \quad (\text{B.17})$$

so that

$$\begin{aligned} u_+(t) &= \frac{e^{T/2\tau}}{2 + e^{T/2\tau} + e^{-T/2\tau}} N_5(t + T/2) \\ u_-(t) &= \frac{e^{-T/2\tau}}{2 + e^{T/2\tau} + e^{-T/2\tau}} N_5(t - T/2) \\ v_1(t) &= \frac{1}{2 + e^{T/2\tau} + e^{-T/2\tau}} N_5(t) \\ v_2(t) &= \frac{1}{2 + e^{T/2\tau} + e^{-T/2\tau}} N_5(t). \end{aligned} \quad (\text{B.18})$$

(These factors out front aren't so far off from 1/4 since $e^{\pm T/2\tau} \approx e^{\pm 4.35/2*64.4} \approx 1.034, .967$.) Then instead $R(t)$ becomes

$$R(t) = \frac{2(1 + A \cos(\omega_a t)) - (1 - A \cos(\omega_a t + \delta)) - (1 - A \cos(\omega_a t - \delta))}{2(1 + A \cos(\omega_a t)) + (1 - A \cos(\omega_a t + \delta)) + (1 - A \cos(\omega_a t - \delta))}, \quad (\text{B.19})$$

where the $e^{\pm T/2\tau}$ terms out front now cancel. Using Equation B.9 again and this time avoiding approximations in δ ,

$$R(t) = \frac{2A \cos(\omega_a t)(1 + \cos \delta)}{4 + 2A \cos(\omega_a t)(1 - \cos \delta)}, \quad (\text{B.20})$$

after simplifying. In the limit that

$$\delta = \pi(\delta T) \rightarrow 0 \quad (\text{B.21})$$

since δT is small,

$$R(t) \approx A \cos(\omega_a t), \quad (\text{B.22})$$

with the only approximation being made at $\mathcal{O}(\delta^2) \sim \mathcal{O}(10^{-10})$.

Finally, while the 3 parameter ratio function suffices for fits to data containing slow modulations, it does not suffice for faster oscillation features. In that case it is more useful to fit with the non-approximated or simplified version of the ratio,

$$\begin{aligned} R(t) &= \frac{v_1(t) + v_2(t) - u_+(t) - u_-(t)}{v_1(t) + v_2(t) + u_+(t) + u_-(t)}, \\ &= \frac{2f(t) - f_+(t) - f_-(t)}{2f(t) + f_+(t) + f_-(t)}, \end{aligned} \quad (\text{B.23})$$

where

$$\begin{aligned} f(t) &= C(t)(1 + A \cos(\omega_a t + \phi)) \\ f_{\pm}(t) &= f(t \pm T_a/2), \end{aligned} \quad (\text{B.24})$$

and $C(t)$ can encode any other effects in the data that need to be fitted for, such as the CBO,

$$C(t) = 1 + A_{cbo} \cdot e^{-t/\tau_{cbo}} \cdot \cos(\omega_{cbo}t + \phi_{cbo}). \quad (\text{B.25})$$

Additionally, any other fit parameters such as A or ϕ can be made a function of t . Using the non-approximated form for the final fit function gives greater confidence in the fit results for the high precision ω_a extraction necessary for the experimental measurement.

B.2 Ratio Point Errors

In order to determine the errors on the points in the formed ratio, Equation B.3, we use standard error propagation:

$$\sigma_R(t)^2 = \left(\frac{\partial R(t)}{\partial V(t)} \right)^2 \delta V(t)^2 + \left(\frac{\partial R(t)}{\partial U(t)} \right)^2 \delta U(t)^2 \quad (\text{B.26})$$

This works because $V(t)$ and $U(t)$ are statistically independent datasets. Using standard error propagation again,

$$\begin{aligned} \delta V(t)^2 &= \delta v_1(t)^2 + \delta v_2(t)^2 = v_1(t) + v_2(t) = V(t), \\ \delta U(t)^2 &= \delta u_+(t)^2 + \delta u_-(t)^2 = u_+(t) + u_-(t) = U(t). \end{aligned} \quad (\text{B.27})$$

Calculating out and simplifying the partial derivatives, (and this time dropping the t 's),

$$\begin{aligned} \frac{\partial R}{\partial V} &= \frac{2U}{(V+U)^2}, \\ \frac{\partial R}{\partial U} &= \frac{-2V}{(V+U)^2}. \end{aligned} \quad (\text{B.28})$$

Combining and simplifying, we arrive at the error formula:

$$\sigma_R^2 = \frac{4UV}{(V+U)^3} = \frac{1-R^2}{(V+U)} \quad (\text{B.29})$$

Appendix C

Pileup Modified Errors

In the pileup subtraction method detailed in Section ??, pileup events are statistically constructed and then subtracted from the data. Because of this, the errors on the bins need to be adjusted appropriately. Reference [31] describes the modified errors, but is not quite correct. Here is provided an improved calculation that I believe is easier to understand. While we are mainly interested in the errors on the histogram bins after pileup subtraction, it first helps to examine the errors of the pileup histogram itself. Here we only consider doublets.

In the asymmetric shadow window pileup method, shadow doublets are constructed from two singlets. The pileup histogram is then filled as the sum of the doublets minus the singlets,

$$P = D - S, \tag{C.1}$$

where D or S are only added or subtracted when they are above some energy threshold. If the threshold is set to 0, then for every doublet one entry will be added and two will be subtracted. Since these entries are exactly correlated, the error in each time bin will be

$$\sigma_P = \sqrt{N_D}, \tag{C.2}$$

where N_D is the number of doublets in that time bin. If the energy threshold is above 0, then we can determine whether the counts in the pileup histogram increase

	$E_1 < E_{th}$	$E_1 > E_{th}$
$E_2 < E_{th}$	$N_1(+1)$	$N_2(0)$
$E_2 > E_{th}$	$N_3(0)$	$N_4(-1)$

Table C.1: Table of doublets above threshold. Here E_1 and E_2 are the energies of the two singlets, E_{th} is the energy threshold, and N_i are the number of doublets above threshold for the different combinations of E_1 and E_2 . (N_1 is assumed above threshold here.) The numbers in the parentheses indicate the number of counts gained or lost in the pileup histogram.

or decrease based on whether the singlets and doublets are above threshold or not.

Table C.1 shows the different combinations of counts put into the pileup histogram.

The counts that go into P will be

$$\begin{aligned}
 P &= \sum_i N_i - \text{singlets above threshold} \\
 &= (N_1 + N_2 + N_3 + N_4) - (N_2 + N_4) - (N_3 + N_4) \\
 &= N_1 - N_4
 \end{aligned} \tag{C.3}$$

and the errors are

$$\sigma_P = \sqrt{N_1 + N_4}. \tag{C.4}$$

This makes sense considering the cases individually. In the cases for N_1 , you will gain a count from the doublet above threshold, and lose no counts since both singlets are below threshold. In the cases for N_2 and N_3 , you will gain a count from the doublet, and lose a count from one of the singlets which is above threshold. In the cases for N_4 , you will gain a count from the doublet and lose two counts from the singlets which are both above threshold. Since the doublet and singlets are exactly correlated, the N_1 and N_4 cases naturally result in a single weight being added into the error, while the N_2 and N_3 cases result in no additions to the error.

Now what about the pileup subtracted time spectrum? Our corrected spectrum can be written as

$$N_{\text{corrected}} = N_{\text{measured}} - P. \quad (\text{C.5})$$

What is in N_{measured} doesn't matter exactly. What we care about is what is in N_{measured} that is also within P , for that is where the correlations come from. Since N_{measured} is the sum of all singlets above threshold, we can write it as

$$N_{\text{measured}} = N_{\text{other}} + N_2 + N_3 + 2N_4 \quad (\text{C.6})$$

since we know that those cases N_i listed come from singlets above threshold, and N_{other} is anything in the measured hits that was not included in the pileup shadow construction. We can then replace P and simplify to get

$$N_{\text{corrected}} = N_{\text{other}} - N_1 + N_2 + N_3 + 3N_4. \quad (\text{C.7})$$

The error on the corrected histogram is then

$$\sigma_{N_{\text{corrected}}} = \sqrt{N_{\text{other}} + N_1 + N_2 + N_3 + 9N_4}. \quad (\text{C.8})$$

Replacing N_{other} as

$$N_{\text{other}} = N_{\text{corrected}} + N_1 - N_2 - N_3 - 3N_4, \quad (\text{C.9})$$

we can remove the dependence of the corrected histogram errors on the unknown quantity and arrive at

$$\begin{aligned} \sigma_{N_{\text{corrected}}} &= \sqrt{N_{\text{corrected}} + 2N_1 + 6N_4}, \\ &= \sqrt{N_{\text{corrected}}} \cdot \sqrt{1 + (2N_1 + 6N_4)/N_{\text{corrected}}}. \end{aligned} \quad (\text{C.10})$$

(This argument might seem circular at the end, but it works because of the squaring that occurs when calculating the error.) In the end we have a form for the bin errors of the pileup corrected histogram which only depend on N_1 and N_4 in addition to the number of counts in the corrected histogram. As shown it can be refactored into a form equal to the naive errors (just the bin content) times some correction factor. Since N_1 and N_4 are much smaller than $N_{\text{corrected}}$ at all times, and because they decay away at about twice the rate as the pileup diminishes, the change to the errors is small, of the order 1 or 2% at 30 μs .

C.1 For the ratio function

Equation C.10 applies to the corrected errors for a pileup subtracted histogram, but what about the modifications to the ratio errors? If we parameterize that equation as

$$\sigma_{N_{\text{corrected}}} = \sqrt{N_{\text{corrected}}} \cdot \sqrt{\gamma(t)}, \quad (\text{C.11})$$

where the correction factor $\gamma(t) \approx \gamma e^{-t/\tau_\mu}$ is small and decays at approximately the muon lifetime, we can recast the errors on the individual ratio sub-datasets as

$$\begin{aligned} \delta V(t)^2 &= \delta v_1(t)^2 \cdot \gamma(t) + \delta v_2(t)^2 \cdot \gamma(t) = (v_1(t) + v_2(t)) \cdot \gamma(t) = V(t) \cdot \gamma(t), \\ \delta U(t)^2 &= \delta u_+(t)^2 \cdot \gamma(t + T/2) + \delta u_-(t)^2 \cdot \gamma(t - T/2) \\ &\approx u_+(t) \cdot \gamma(t) e^{-T/2\tau} + u_-(t) \cdot \gamma(t) e^{+T/2\tau} \\ &\approx (u_+(t) + u_-(t)) \cdot \gamma(t) \cdot \left(1 + \frac{1}{2} \left(\frac{T}{2\tau}\right)^2\right) \\ &\approx U(t) \cdot \gamma(t), \end{aligned} \quad (\text{C.12})$$

where in the last step the $\frac{1}{2} \left(\frac{T}{2\tau}\right)^2$ term has been neglected because it's small. With these approximations having been made, the modified errors on the ratio points simply

become

$$\sigma_R^2 \rightarrow \sigma_R^2 \cdot \gamma(t), \quad (\text{C.13})$$

with the correction being the same as that on the pileup subtracted histogram. Credit to Reference [32] for this derivation.

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