

Quantum Hamiltonian Descent for Rigid Image Registration

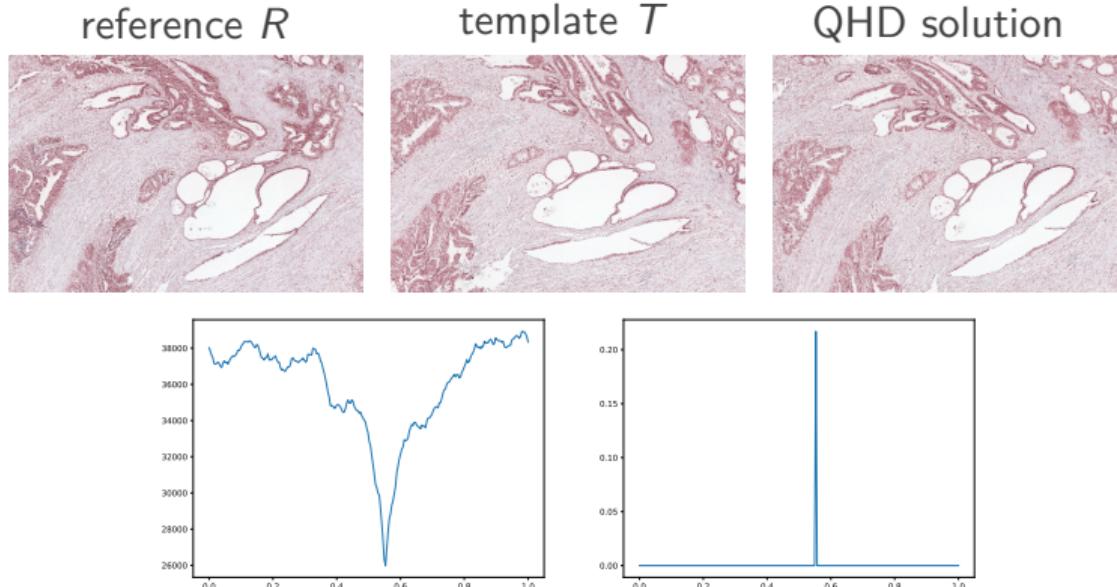
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Rigid Image Registration

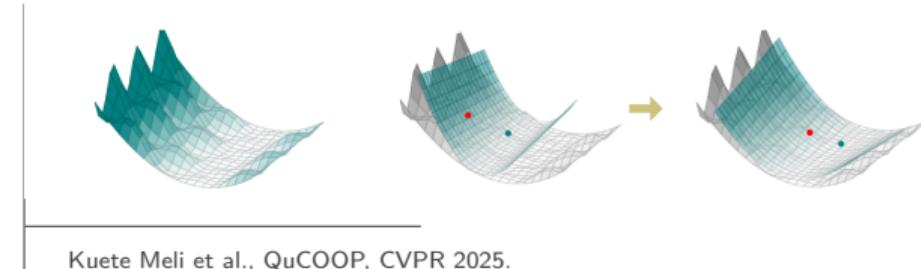


$$\min_{\phi: \mathbb{R}^d \rightarrow \mathbb{R}^d} \mathcal{D}(R, T \circ \phi) + \mathcal{R}(\phi), \quad \phi(x; t, \alpha)^{-1} = R_\alpha x + t, \quad \mathcal{R}(\phi) = 0.$$



Motivation

Optimizing non convex functions is challenging: Traditional methods only iteratively solve a **simple local** (quadratic) surrogate.



Kuete Meli et al., QuCOOP, CVPR 2025.

Quantum computers use superposition: They can “see” the **global** landscape of a function.

$$|\psi\rangle = \alpha_0 |0\rangle + \dots + \alpha_{2^n-1} |2^n - 1\rangle$$
$$|\psi(f)\rangle = e^{i\theta f(0)} |0\rangle + \dots + e^{i\theta f(2^n-1)} |2^n - 1\rangle$$

Can this global view ability of quantum computers be leveraged to escape local minima in optimization?

Background

QCs operate on quantum bits.

Qubit

A qubit is a one-particle system whose state vector is expressed as

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle,$$

with

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \text{and} \quad |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix},$$

$\alpha, \beta \in \mathbb{C}$ subject to $|\alpha|^2 + |\beta|^2 = 1$.

On multiple qubits, it holds:

■ Kronecker product:

$$|\psi\rangle = \alpha_{000}|000\rangle + \alpha_{001}|001\rangle + \cdots + \alpha_{111}|111\rangle.$$

$$\begin{array}{c} \uparrow \qquad \qquad \qquad \uparrow \qquad \qquad \qquad \uparrow \\ |0\rangle \otimes |0\rangle \otimes |0\rangle \qquad |0\rangle \otimes |0\rangle \otimes |1\rangle \qquad |1\rangle \otimes |1\rangle \otimes |1\rangle \\ \text{or} \\ \left(\begin{matrix} 1 \\ 0 \end{matrix}\right) \otimes \left(\begin{matrix} 1 \\ 0 \end{matrix}\right) \otimes \left(\begin{matrix} 1 \\ 0 \end{matrix}\right) \qquad \left(\begin{matrix} 1 \\ 0 \end{matrix}\right) \otimes \left(\begin{matrix} 1 \\ 0 \end{matrix}\right) \otimes \left(\begin{matrix} 0 \\ 1 \end{matrix}\right) \\ \text{or 0 in binary basis} \qquad \text{or 1 in binary basis} \qquad \text{or 7 in binary basis} \\ \left(\begin{matrix} 0 \\ 1 \end{matrix}\right) \otimes \left(\begin{matrix} 0 \\ 1 \end{matrix}\right) \otimes \left(\begin{matrix} 0 \\ 1 \end{matrix}\right) \end{array}$$

■ Entanglement:

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 1 \end{pmatrix};$$

no $|\psi_1\rangle, |\psi_2\rangle$ exist such that $|\psi\rangle = |\psi_1\rangle \otimes |\psi_2\rangle$.

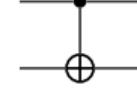
Quantum Operators and Evolution

The state vector $|\psi(t)\rangle$ of a quantum system at time t relates to the state $|\psi(0)\rangle$ via

$$|\psi(t)\rangle = \mathbf{U}(t) |\psi(0)\rangle,$$

where \mathbf{U} is a unitary operator that uniquely depends on t .

Example of unitary operators:

Gate name	Matrix form	Circuit	Example application
Hadamard	$\mathbf{H} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$		$\mathbf{H} 0\rangle = \frac{1}{\sqrt{2}}(0\rangle + 1\rangle)$
Phase gate	$\mathbf{P}(2\theta) = \begin{pmatrix} 1 & 0 \\ 0 & e^{-i\theta} \end{pmatrix}$		$\mathbf{P} 1\rangle = e^{-i\theta} 1\rangle$
Controlled-X on $ q_0q_1\rangle$	$\mathbf{I} \otimes \mathbf{X}^{q_0} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$		$\mathbf{I} \otimes \mathbf{X}^{q_0} 00\rangle = 00\rangle$ $\mathbf{I} \otimes \mathbf{X}^{q_0} 10\rangle = 11\rangle$

Gate-based QC: Nielsen and Chuang, *Quantum Computation and Quantum Information*. Cambridge University Press, 2010.

Sutor. *Dancing with Qubits: How quantum computing works and how it can change the world*. Packt Publishing Ltd, 2019.

Continuous Quantum Computing

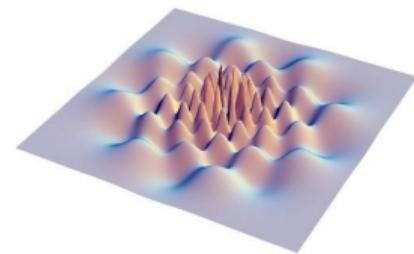
In its continuous form, a quantum state is a wave function $\psi(x, t)$ — a complex-valued function over \mathbb{R}^d , where $|\psi(x, t)|^2$ gives the **probability density** of finding a particle at position x and time t .

This wave function evolves according to the **Schrödinger equation**:

$$i\hbar \frac{\partial}{\partial t} \psi(x, t) = H(t) \psi(x, t),$$

where $H(t)$ is the **Hamiltonian** governing the system dynamics.

The goal is to design $H(t)$ to concentrate the wave function around the global minimum of a target function f — enabling optimization via quantum dynamics.



Continuous QC: Feynman et al. *Quantum mechanics and path integrals*, OUP Oxford 1966.

Image source: <https://quantumphysicslady.org/wp-content/uploads/2019/03/Wavefunction-equationsource-wavefunction.png>

Quantum Hamiltonian Descent (QHD)

Quantum Hamiltonian Descent (QHD, Leng et al. 2023, 2025)

Consider the Schrödinger equation

$$i\hbar \frac{\partial}{\partial t} \psi(x, t) = H(t) \psi(x, t) \quad \text{under} \quad H(t) = e^{\chi_t} f(x) + e^{\phi_t} \left(-\frac{1}{2} \Delta \right)$$

with suitable real parameters ϕ_t, χ_t , Δ being the Laplacian operator. Under the condition that $e^{\phi_t}/e^{\chi_t} \xrightarrow{t \rightarrow \infty} 0$ and certain regularity assumptions on f and the Hamiltonian, it holds

$$\lim_{t \rightarrow \infty} \int_x |\psi(x, t)|^2 f(x) dx = \min f.$$

Measuring the state after a sufficiently long time t will yield a near-optimal solution x with high probability.

Continuous-Time Dynamics of Gradient Methods

Most gradient methods are discretizations of continuous-time flows:

Method	Update rule	Continuous limit ($\alpha \rightarrow 0$)
Gradient Descent	$x_{k+1} = x_k - \alpha \nabla f(x_k)$	$\dot{X}_t = -\nabla f(X_t)$
Nesterov	$x_{k+1} = x_k + \mu(x_k - x_{k-1}) - \alpha \nabla f(\dots)$	$\ddot{X}_t + \frac{3}{t} \dot{X}_t = -\nabla f(X_t)$
Newton	$x_{k+1} = x_k - \alpha (\nabla^2 f(x_k))^{-1} \nabla f(x_k)$	$\dot{X}_t = -(\nabla^2 f(X_t))^{-1} \nabla f(X_t)$

These flows X minimize an **action functional** (Wibisono et al., 2016):

$$S(X) = \int_t \mathcal{L}(t, X_t, \dot{X}_t) dt, \quad \mathcal{L} = e^{\alpha_t + \gamma_t} \left(\beta(t) f(X_t) + \text{Dist}(X_t + \alpha(t) \dot{X}_t, X_t) \right),$$

by optimality condition $\nabla_X S(X) = 0$ yielding Euler–Lagrange Eq. $\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{X}} \right) - \frac{\partial \mathcal{L}}{\partial X} = 0$.

Compare with QHD:

$$i\hbar \frac{\partial}{\partial t} \psi(x, t) = H(t) \psi(x, t), \quad H(t) = e^{\chi_t} f(x) + e^{\phi_t} \left(-\frac{1}{2} \Delta \right)$$

The QHD Algorithm

Considering

$$i\hbar \frac{\partial}{\partial t} \psi(x, t) = H(t)\psi(x, t), \quad H(t) = e^{\chi_t} f(x) + e^{\phi_t} \left(-\frac{1}{2} \Delta \right),$$

a spatial, then temporal discretization yield

$$i\hbar \frac{\partial}{\partial t} \psi_t = \left(e^{\chi_t} \hat{F} + e^{\phi_t} \hat{L} \right) \phi_t, \quad \text{then} \quad \psi_{j+1} = \exp(-ia_j dt \hat{L} - ib_j dt \hat{F}) \psi_j,$$

where $a_j = e^{\phi_j \cdot dt}$ and $b_j = e^{\chi_j \cdot dt}$. Further simplifications apply Trotter-Suzuki as

$$\exp(-ia_j dt \hat{L} - ib_j dt \hat{F}) \approx \exp(-ia_j dt \hat{L}) \exp(-ib_j dt \hat{F}),$$

and diagonalization $\hat{L} = \mathcal{F} D \mathcal{F}^{-1}$ for efficient exponentiation, \mathcal{F} being the QFT.

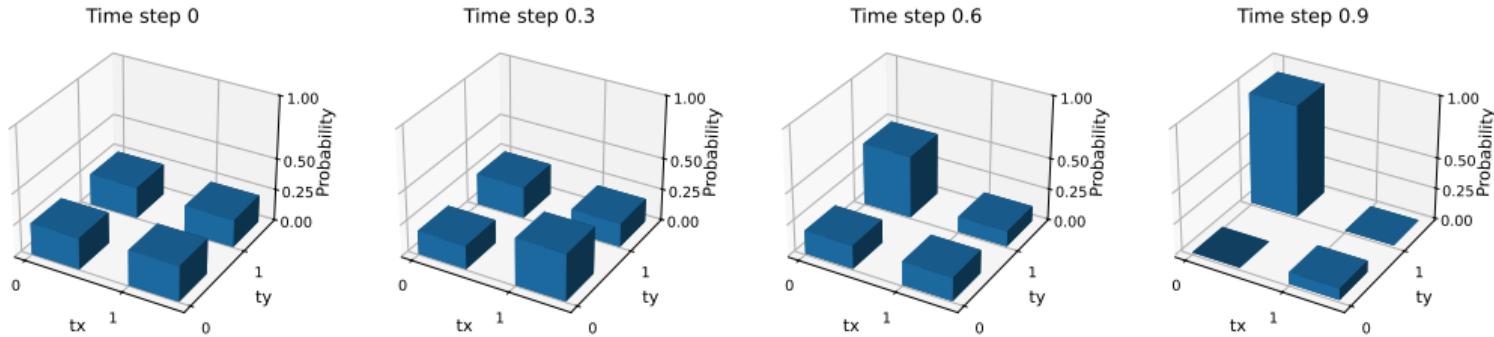
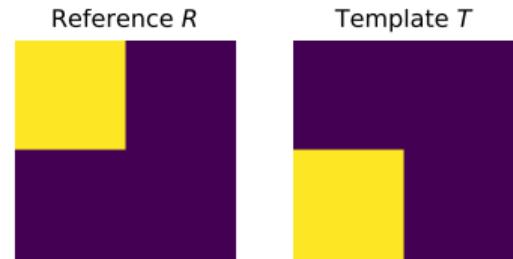
QHD Iteration

$$|\psi_{j+1}\rangle := \mathcal{F} \exp(-ia_j dt \hat{D}) \mathcal{F}^{-1} \exp(-ib_j dt \hat{F}) |\psi_j\rangle, \quad |\psi_0\rangle = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} |k\rangle.$$

Example Implementation

Five registers are initialized as $|\psi_0\rangle = \sum_{tx,ty} |0\rangle_R |0\rangle_T |0\rangle_{SSD} |t_x\rangle_{tx} |t_y\rangle_{ty}$.

- $|R\rangle_R = |1000\rangle, |T\rangle_T = |0010\rangle$.
- $|T \circ \phi(t_x, t_y)\rangle$: controlled shifts on $|T\rangle$.
- SSD phases $e^{-ib_0 \cdot dt \cdot \text{SSD}(t_x, t_y)}$ using:
 - ▶ $|a\rangle |b\rangle |\text{SSD}\rangle_{\text{SSD}} \mapsto |a\rangle |b\rangle |\text{SSD} + (a^2 + b^2 - 2ab)\rangle_{\text{SSD}}$,
 - ▶ $\bigotimes_{i=0}^K P(-b_0 \cdot dt \cdot 2^i) |q_0 q_1 \cdots q_K\rangle_{\text{SSD}}$.



Title

reference R



template T



QHD solution

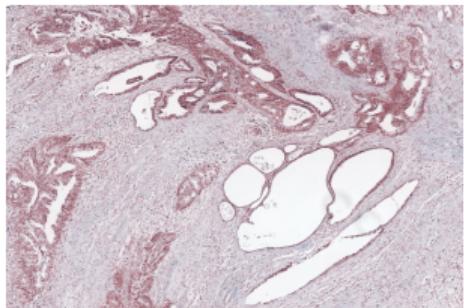


global optimum

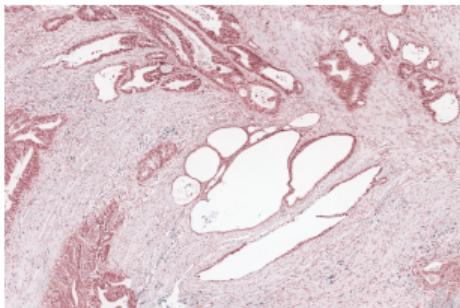


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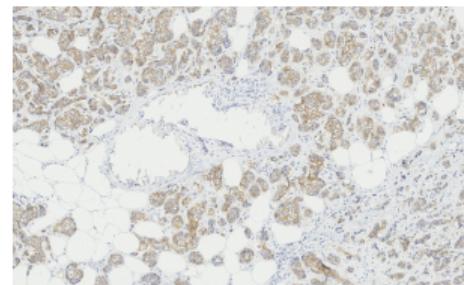
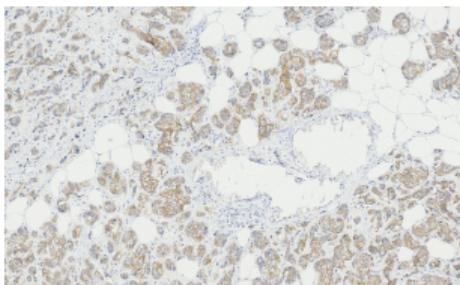
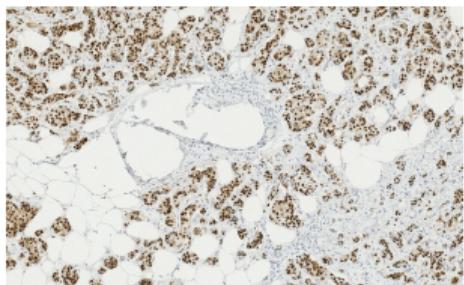
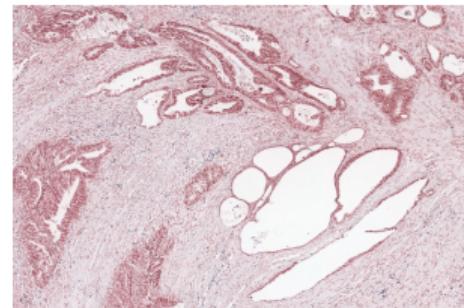
reference R



template T



QHD solution



Title

Dataset	$p(x^*)$	$p(\bar{x})$	$\ x^* - \bar{x}\ _{\ell^1}$
balloon3	52.55%	56.68%	3
playground	4.18%	7.60%	4
dynamicFace	53.31%	57.45%	3
skating	25.38%	41.47%	5
coard17b	79.62%	79.62%	0
breast5a	58.93%	58.93%	0
breast5b	41.61%	41.79%	1
lung-lesion2	78.57%	78.57%	0

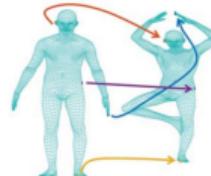
Further Interests

Adiabatic quantum computers can solve QUBO problems:

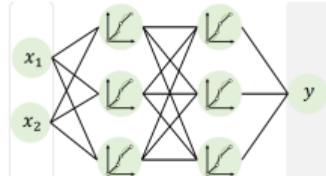
$$\min_{x \in \{0,1\}^n} x^\top Q x$$



Correspondence problems
Kuete Meli et al., CVPR 2025



Matching problems
Benkner et al., 3DV 2020



Quantized neural nets
Li et al., ArXiv 2025

Quantum computers “see” the global landscape of a function:
Useful in optimization?

$$|\psi\rangle = \alpha_0 |0\rangle + \cdots + \alpha_{2^n-1} |2^n - 1\rangle$$
$$|\psi(f)\rangle = f(0) |0\rangle + \cdots + f(2^n - 1) |2^n - 1\rangle$$

Can we enhance computer vision solutions by leveraging quantum methods?
Enhancement can be in term of complexity, time-to-solution, memory demand, ...

Thanks! Questions?

