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ν Mass & $0\nu\beta\beta$ in EFT Framework

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Hao-Lin Li, Zhe Ren, Ming-Lei Xiao, **JHY**, Yu-Hui Zheng, 2007.07899

Hao-Lin Li, Zhe Ren, Ming-Lei Xiao, **JHY**, Yu-Hui Zheng, 2012.09188

Hao-Lin Li, Zhe Ren, Ming-Lei Xiao, **JHY**, Yu-Hui Zheng, 2105.09329

Yong Du, Hao-Lin Li, Yu-Han Ni, Ming-Lei Xiao, **JHY**, in preparation

Gang Li, Hao Sun, **JHY**, in collaboration

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Outline

- Why $0\nu b\bar{b}$ in EFT approach?
- SMEFT: $0\nu b\bar{b}$ and $n\nu$ masses, UV physics
- LEFT: quark currents and weak sources
- ChiPT: short-range, pion-range, long-range
- Summary and outlook

Introduction

Why 0vbb in EFT approach?

Search For New Physics

ATLAS SUSY Searches* - 95% CL Lower Limits

March 2021

ATLAS Preliminary

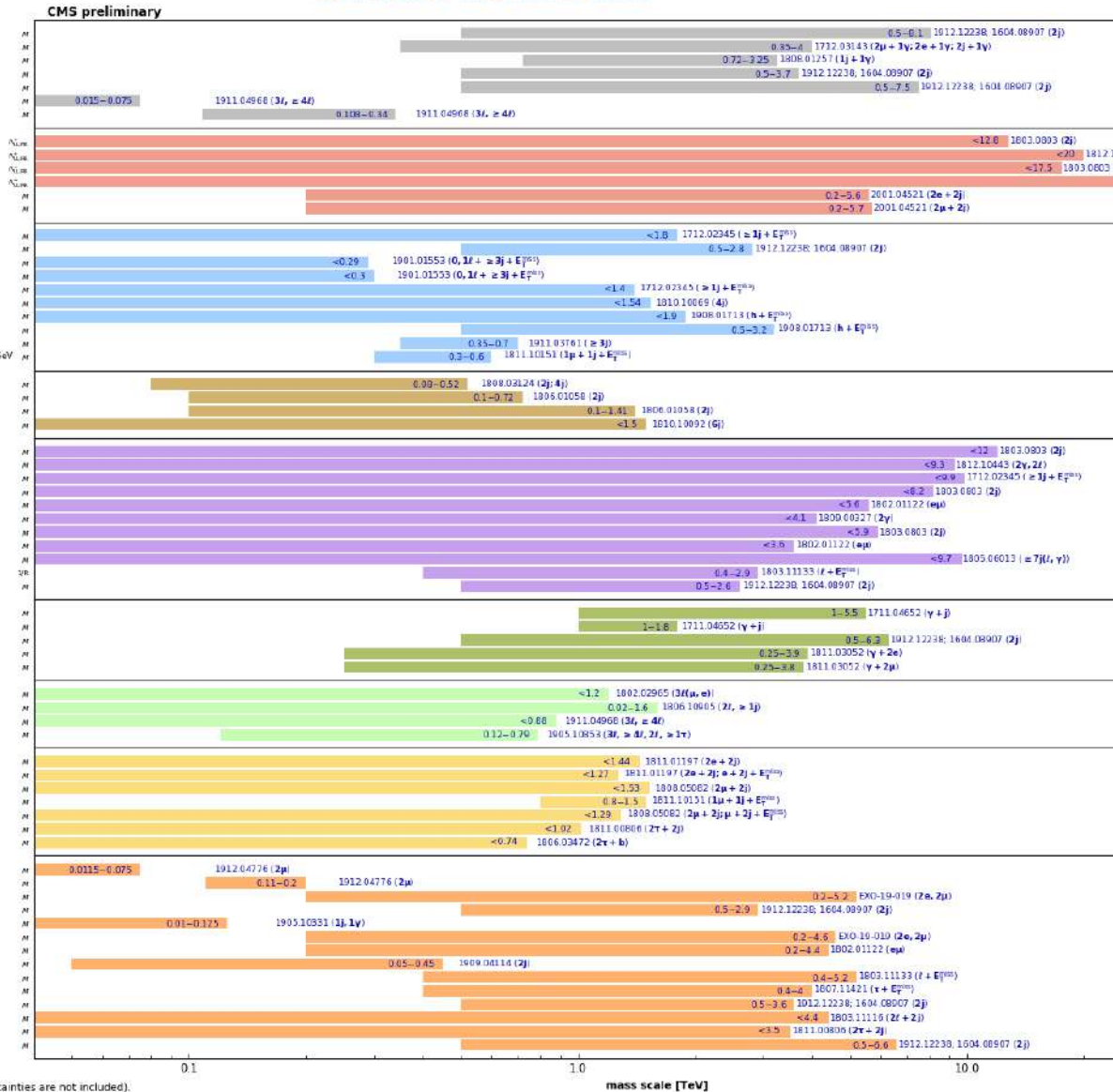
$\sqrt{s} = 13$ TeV

Model	Signature	$\int \mathcal{L} dt$ [fb ⁻¹]	Mass limit	Reference							
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0 e, μ mono-jet	2-6 jets 1-3 jets	E_T^{miss} E_T^{miss}	139 36.1	\tilde{q} [1x, 8x Degen.] \tilde{q} [8x Degen.]	1.0 0.9	1.85	$m(\tilde{\chi}_1^0) < 400$ GeV $m(\tilde{q}) - m(\tilde{\chi}_1^0) = 5$ GeV	2010.14293 2102.10874	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0 e, μ	2-6 jets	E_T^{miss}	139	\tilde{g} \tilde{g}	Forbidden	2.3 1.15-1.95	$m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{\chi}_1^0) = 1000$ GeV	2010.14293 2010.14293	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}W\tilde{\chi}_1^0$	1 e, μ	2-6 jets	E_T^{miss}	139	\tilde{g}		2.2	$m(\tilde{\chi}_1^0) < 600$ GeV	2101.01629	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	$ee, \mu\mu$	2 jets	E_T^{miss}	36.1	\tilde{g}		1.2	$m(\tilde{g}) - m(\tilde{\chi}_1^0) = 50$ GeV	1805.11381	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}WZ\tilde{\chi}_1^0$	0 e, μ	7-11 jets	E_T^{miss}	139	\tilde{g}		1.97	$m(\tilde{\chi}_1^0) < 600$ GeV	2008.06032	
		SS e, μ	6 jets	E_T^{miss}	139	\tilde{g}		1.15	$m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200$ GeV	1909.08457	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0-1 e, μ SS e, μ	3 b 6 jets	E_T^{miss} E_T^{miss}	79.8 139	\tilde{g} \tilde{g}		2.25 1.25	$m(\tilde{\chi}_1^0) < 200$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300$ GeV	ATLAS-CONF-2018-041 1909.08457	
	direct production	$\tilde{b}_1\tilde{b}_1$	0 e, μ	2 b	E_T^{miss}	139	\tilde{b}_1 \tilde{b}_1		1.255 0.68	$m(\tilde{\chi}_1^0) < 400$ GeV 10 GeV $< \Delta m(\tilde{b}_1, \tilde{\chi}_1^0) < 20$ GeV	2101.12527 2101.12527
		$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_2^0 \rightarrow bh\tilde{\chi}_1^0$	0 e, μ 2 τ	6 b 2 b	E_T^{miss} E_T^{miss}	139 139	\tilde{b}_1	Forbidden	0.23-1.35	$\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 100$ GeV $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 0$ GeV	1908.03122 ATLAS-CONF-2020-031
		$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	0-1 e, μ	≥ 1 jet	E_T^{miss}	139	\tilde{t}_1		1.25	$m(\tilde{\chi}_1^0) = 1$ GeV	2004.14060, 2012.03799
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$		1 e, μ	3 jets/1 b	E_T^{miss}	139	\tilde{t}_1	Forbidden	0.65	$m(\tilde{\chi}_1^0) = 500$ GeV	2012.03799	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}b\nu, \tilde{\tau}_1 \rightarrow \tau\tilde{G}$		1-2 τ	2 jets/1 b	E_T^{miss}	139	\tilde{t}_1	Forbidden	1.4	$m(\tilde{\tau}_1) = 800$ GeV	ATLAS-CONF-2021-008	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c} \rightarrow c\tilde{\chi}_1^0$		0 e, μ 0 e, μ	2 c mono-jet	E_T^{miss} E_T^{miss}	36.1 139	\tilde{t}_1 \tilde{t}_1		0.85 0.55	$m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 5$ GeV	1805.01649 2102.10874	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z/h\tilde{\chi}_1^0$		1-2 e, μ	1-4 b	E_T^{miss}	139	\tilde{t}_1		0.067-1.18			
$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$		3 e, μ	1 b	E_T^{miss}	139	\tilde{t}_2	Forbidden	0.86			
direct		$\tilde{\chi}_1^0\tilde{\chi}_2^0$ via WZ	3 e, μ $ee, \mu\mu$	≥ 1 jet	E_T^{miss} E_T^{miss}	139 139	$\tilde{\chi}_1^0/\tilde{\chi}_2^0$ $\tilde{\chi}_1^0/\tilde{\chi}_2^0$		0.64 0.205		
		$\tilde{\chi}_1^0\tilde{\chi}_2^0$ via WW	2 e, μ		E_T^{miss}	139	$\tilde{\chi}_1^0$		0.42		
	$\tilde{\chi}_1^0\tilde{\chi}_2^0$ via Wh	0-1 e, μ	2 $b/2 \gamma$	E_T^{miss}	139	$\tilde{\chi}_1^0/\tilde{\chi}_2^0$	Forbidden	0.74			
	$\tilde{\chi}_1^0\tilde{\chi}_1^0$ via $\tilde{L}_R/\tilde{\nu}$	2 e, μ		E_T^{miss}	139	$\tilde{\chi}_1^0$		1.0			
	$\tilde{\tau}\tilde{\tau}, \tilde{\tau} \rightarrow \tau\tilde{\chi}_1^0$	2 τ		E_T^{miss}	139	$\tilde{\tau}$ [$\tilde{\tau}_L, \tilde{\tau}_{R,L}$]		0.16-0.3	0.12-0.39		
	$\tilde{t}_L\tilde{t}_R, \tilde{t}_L \rightarrow t\tilde{\chi}_1^0$	2 e, μ	0 jets	E_T^{miss}	139	\tilde{t}		0.7			
	$\tilde{t}_L\tilde{t}_R, \tilde{t}_L \rightarrow t\tilde{\chi}_1^0$	$ee, \mu\mu$	≥ 1 jet	E_T^{miss}	139	\tilde{t}		0.256			
	$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	0 e, μ 4 e, μ	$\geq 3 b$ 0 jets	E_T^{miss} E_T^{miss}	36.1 139	\tilde{H} \tilde{H}		0.13-0.23 0.55	0.29-0.88		
	particles	Direct $\tilde{\chi}_1^0\tilde{\chi}_1^0$ prod., long-lived $\tilde{\chi}_1^0$	Disapp. trk	1 jet	E_T^{miss}	139	$\tilde{\chi}_1^0$ $\tilde{\chi}_1^0$		0.66 0.21		
		Stable \tilde{g} R-hadron		Multiple		36.1	\tilde{g}				
Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$			Multiple		36.1	\tilde{g} [$\tau(\tilde{g}) = 10$ ns, 0.2 ns]					
$\tilde{t}\tilde{t}, \tilde{t} \rightarrow t\tilde{G}$		Displ. lep		E_T^{miss}	139	$\tilde{t}, \tilde{\mu}, \tilde{\tau}$		0.7 0.34			
RPV		$\tilde{\chi}_1^0\tilde{\chi}_1^0/\tilde{\chi}_1^0\tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z\ell\ell\ell$	3 e, μ	0 jets	E_T^{miss}	139	$\tilde{\chi}_1^0/\tilde{\chi}_2^0$ [BR(Z τ)=1, BR(Z e)=1]		0.625	1.05	
		$\tilde{\chi}_1^0\tilde{\chi}_1^0/\tilde{\chi}_1^0\tilde{\chi}_2^0 \rightarrow WW/Z\ell\ell\ell\nu\nu$	4 e, μ		E_T^{miss}	139	$\tilde{\chi}_1^0/\tilde{\chi}_2^0$ [$A_{33} \neq 0, A_{133} \neq 0$]		0.95	1.55	
		$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$		4-5 large- R jets		36.1	\tilde{g} [$m(\tilde{\chi}_1^0) = 200$ GeV, 1100 GeV]			1.3	1.1
		$\tilde{t}\tilde{t}, \tilde{t} \rightarrow t\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow tbs$		Multiple		36.1	\tilde{t} [$A'_{133} = 20-4, 10-2$]		0.55	1.05	
		$\tilde{t}\tilde{t}, \tilde{t} \rightarrow b\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow bbs$		$\geq 4b$		139	\tilde{t}	Forbidden	0.95		
		$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$		2 jets + 2 b		36.7	\tilde{t}_1 [$q\tilde{q}, bs$]		0.42	0.61	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$	2 e, μ 1 μ	2 b DV		36.1 136	\tilde{t}_1 \tilde{t}_1			0.4-1.45 1.0	1.6	
	$\tilde{\chi}_1^0/\tilde{\chi}_2^0/\tilde{\chi}_1^0\tilde{\chi}_2^0 \rightarrow tbs, \tilde{\chi}_1^0 \rightarrow bbs$	1-2 e, μ	≥ 6 jets		139	$\tilde{\chi}_1^0$		0.2-0.32			
	Extra Dimensions										

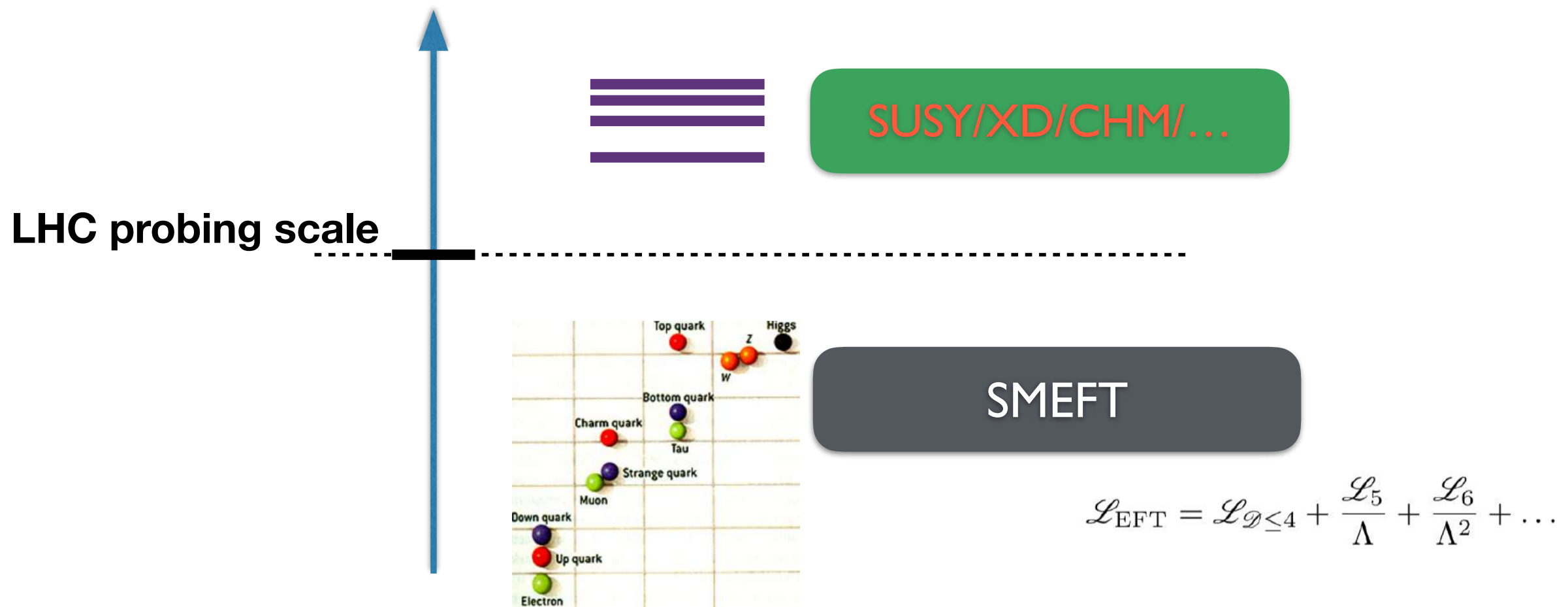
*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

- String resonance
- Z γ resonance
- Higgs γ resonance
- Color Octet Scalar, $k_2^2 = 1/2$
- Scalar Diquark
- $\tilde{t} + \phi$, pseudoscalar (scalar), $g_{\tilde{t}\phi} \times BR(\tilde{t} \rightarrow Z\ell) > -0.03(0.004)$
- $\tilde{t} + \phi$, pseudoscalar (scalar), $g_{\tilde{t}\phi} \times BR(\tilde{t} \rightarrow Z\ell) > -0.03(0.004)$
- quark compositeness ($q\bar{q}$), $\Omega_{\text{LQKK}} = 1$
- quark compositeness (W), $\Omega_{\text{LQKK}} = 1$
- quark compositeness ($q\bar{q}$), $g_{\text{LQKK}} = -1$
- quark compositeness (W), $\Omega_{\text{LQKK}} = -1$
- Excited Lepton Contact Interaction
- Excited Lepton Contact Interaction
- (axial) vector mediator (Z), $g_A = 0.25, g_{\text{SM}} = 1, m_Z = 1$ GeV
- (axial) vector mediator (ϕ), $g_A = 0.22, g_{\text{SM}} = 1, m_\phi = 1$ GeV
- scalar mediator ($+t\bar{t}$), $g_s = 1, g_{\text{SM}} = 1, m_\phi = 1$ GeV
- pseudoscalar mediator ($+t\bar{t}$), $g_A = 1, g_{\text{SM}} = 1, m_\phi = 1$ GeV
- scalar mediator (fermion pair), $A_0 = 1, m_\phi = 1$ GeV
- complex s.c. med. (dark QCD), $m_{\phi_0} = 5$ GeV, $\epsilon\tau_{\phi_0} = 25$ ns
- baryonic Z , $g_A = 0.25, g_{\text{SM}} = 1, m_Z = 1$ GeV
- $Z' - 2HDM$, $g_Z = 0.8, g_{\text{SM}} = 1, \tan\beta = 1, m_{Z'} = 100$ GeV
- Vector resonance, $g_A = 0.25, g_{\text{SM}} = 1, m_{Z'} = 1$ GeV
- Leptoquark mediator, $\beta = 1, B = 0.1, A_{\text{LQ}} = 0.1, 800 < M_{LQ} < 1500$ GeV
- RPV stop to 4 quarks
- RPV squark to 4 quarks
- RPV gluino to 4 quarks
- RPV gluino to 3 quarks
- ADD (1J), $n_{\text{JL}} = 3$
- ADD (1J), $n_{\text{JL}} = 3$
- ADD G_{KK} emission, $n = 2$
- ADD QEH (1J), $n_{\text{JL}} = 6$
- ADD QEH (1J), $n_{\text{JL}} = 6$
- RS G_{KK} (1J), $n_{\text{JL}} = 6$
- RS QEH (1J), $n_{\text{JL}} = 6$
- RS QEH (1J), $n_{\text{JL}} = 6$
- non-relativistic BH, $M_0 = 4$ TeV, $n_{\text{JL}} = 6$
- spin-1/2, $\mu = 4$ TeV
- RS G_{KK} (1J), $n_{\text{JL}} = 6$
- excited light quark ($q\bar{q}$), $f_2 = f = f = 1, \Lambda = m_q^*$
- excited b quark, $f_2 = f = f = 1, \Lambda = m_b^*$
- excited light quark ($q\bar{q}$), $\Lambda = m_q^*$
- excited electron, $f_2 = f = f = 1, \Lambda = m_e^*$
- excited muon, $f_2 = f = f = 1, \Lambda = m_\mu^*$
- MSSM, $|V_{cb}|^2 = 1.8, |V_{ub}|^2 = 1.6$
- MSSM, $|V_{cb}|^2 = 1.8, |V_{ub}|^2 = 1.6$
- Type-II seesaw heavy fermions, Flavor-democratic
- Vector like tau, Doublet
- scalar LQ (pair prod.), coupling to 1st gen. fermions, $\beta = 1$
- scalar LQ (pair prod.), coupling to 1st gen. fermions, $\beta = 0.5$
- scalar LQ (pair prod.), coupling to 2nd gen. fermions, $\beta = 1$
- scalar LQ (pair prod.), coupling to 2nd gen. fermions, $\beta = 0.5$
- scalar LQ (pair prod.), coupling to 3rd gen. fermions, $\beta = 1$
- scalar LQ (single prod.), coupling to 3rd gen. ferm., $\beta = 1, \Lambda = 1$
- Z_ϕ , narrow resonance
- Z_ϕ , narrow resonance
- SSN Z'
- SSN Z' (1J)
- $Z' \rightarrow \phi\phi$
- Superstring Z_ϕ
- UV Z' , BR($e\mu$) = 10%
- Leptoquark Z'
- SSN $W(\mu)$
- SSN $W(\tau)$
- SSN $W(\phi)$
- LRSM $W(\mu)$, $M_W = 0.5M_{W_0}$
- LRSM $W(\tau)$, $M_W = 0.5M_{W_0}$
- Axigluon, Coloron, $\cos\theta = 1$

Overview of CMS EXO results



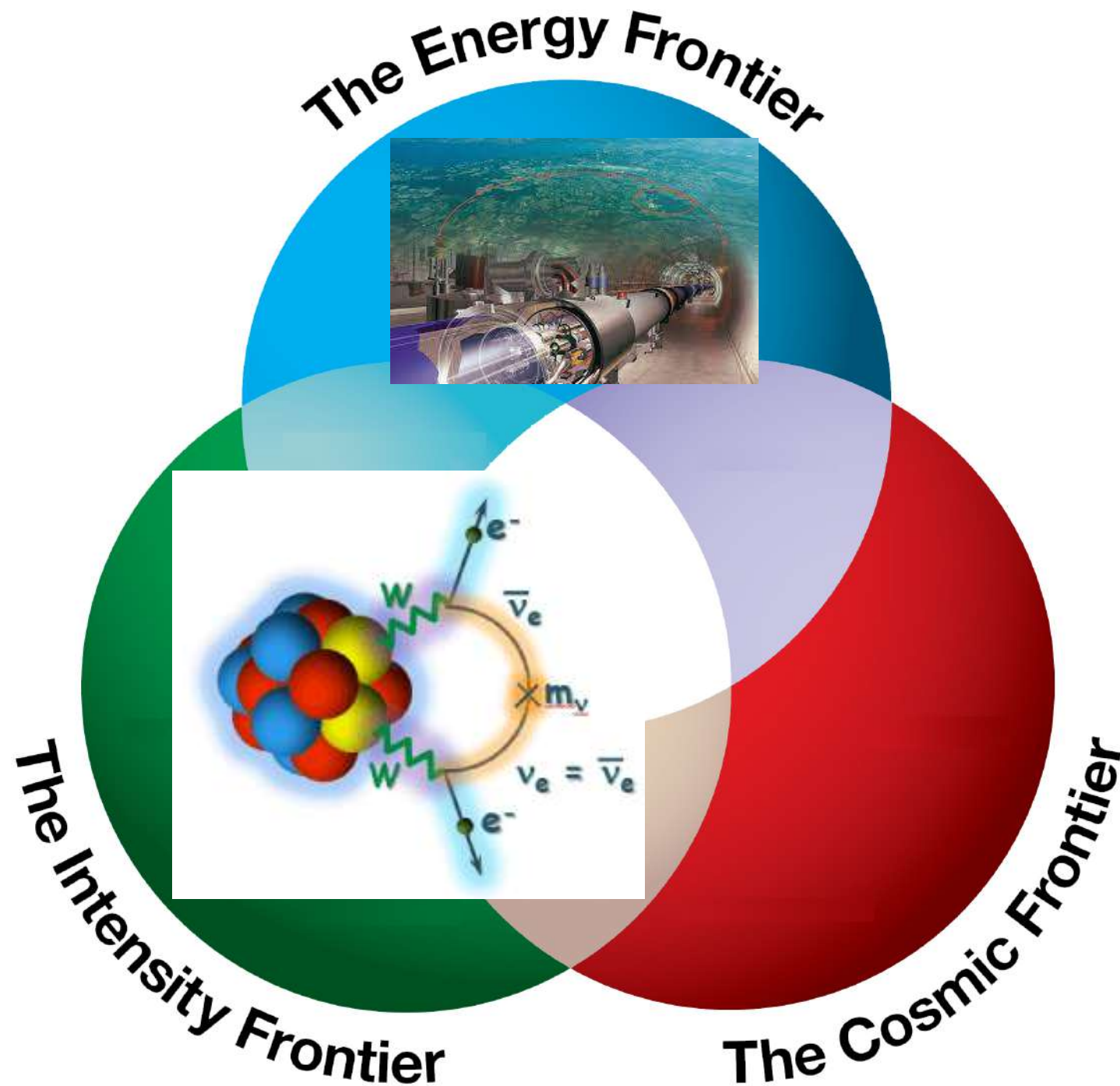
New Physics w/o New Particle



Top-down: Integrate out and matching/running

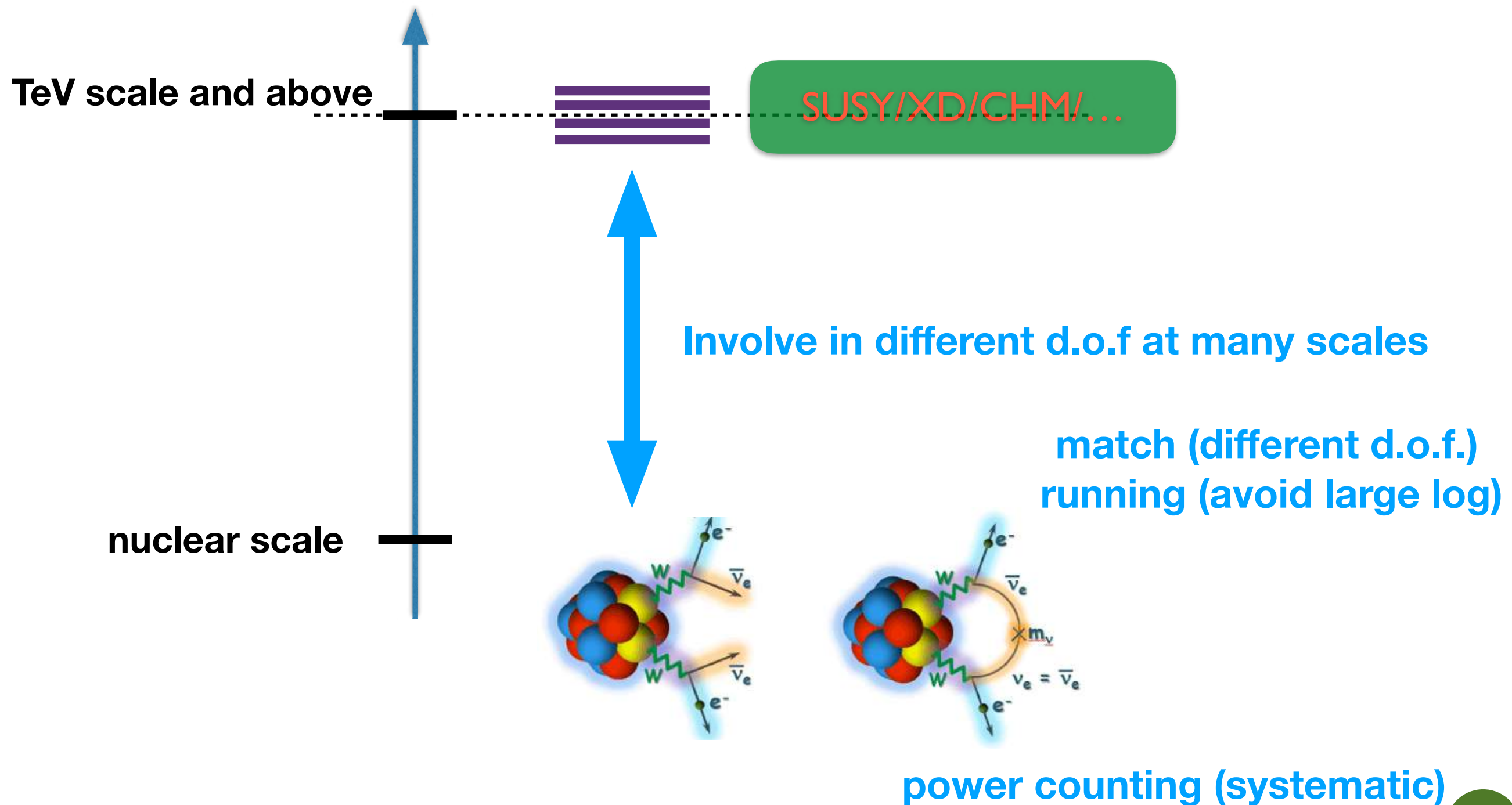
Bottom-up: field d.o.f and symmetry at IR scale

New Physics w/o New Particle



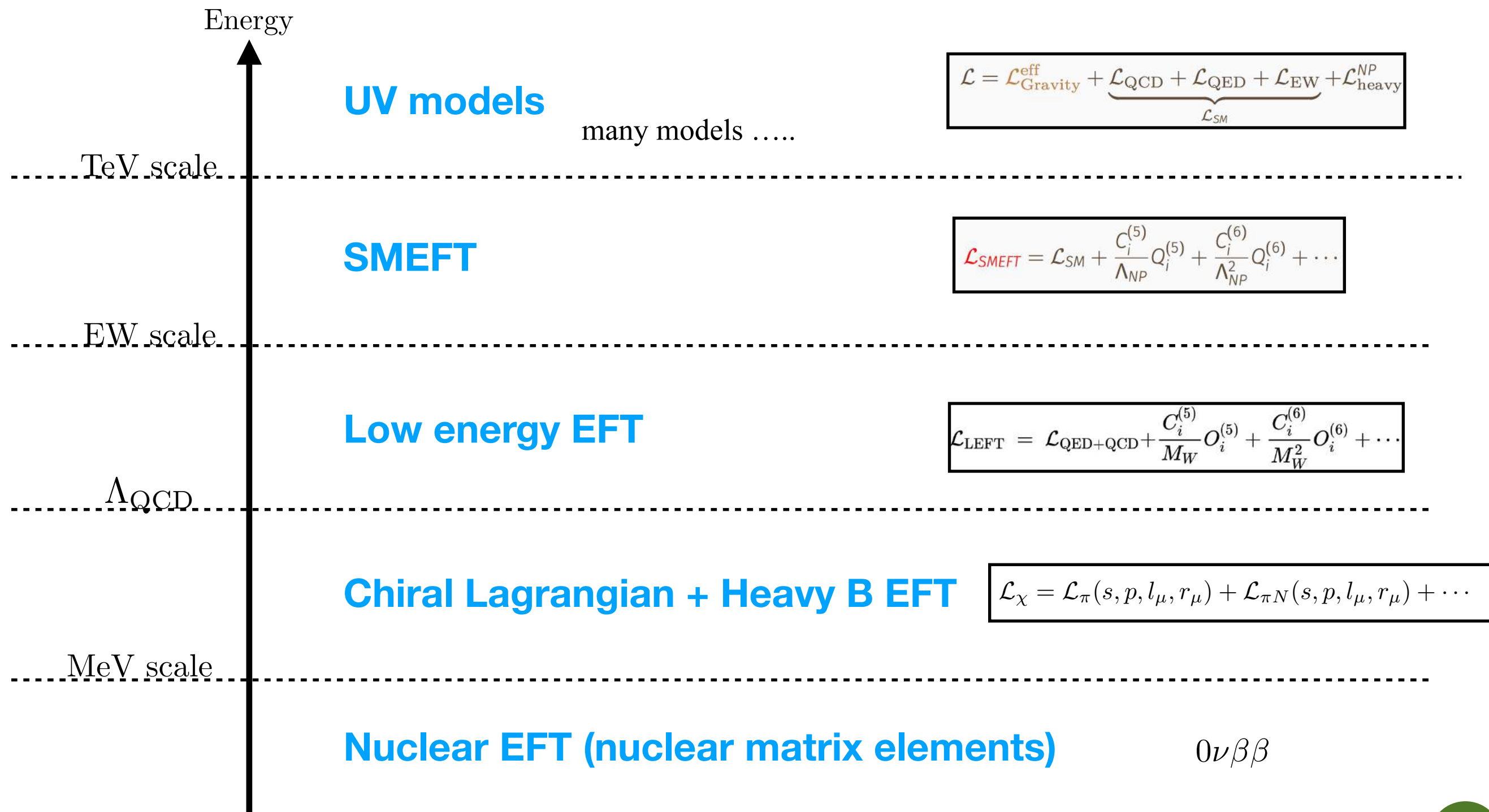
Low Energy Probe of HEP

Low energy probes of high energy physics



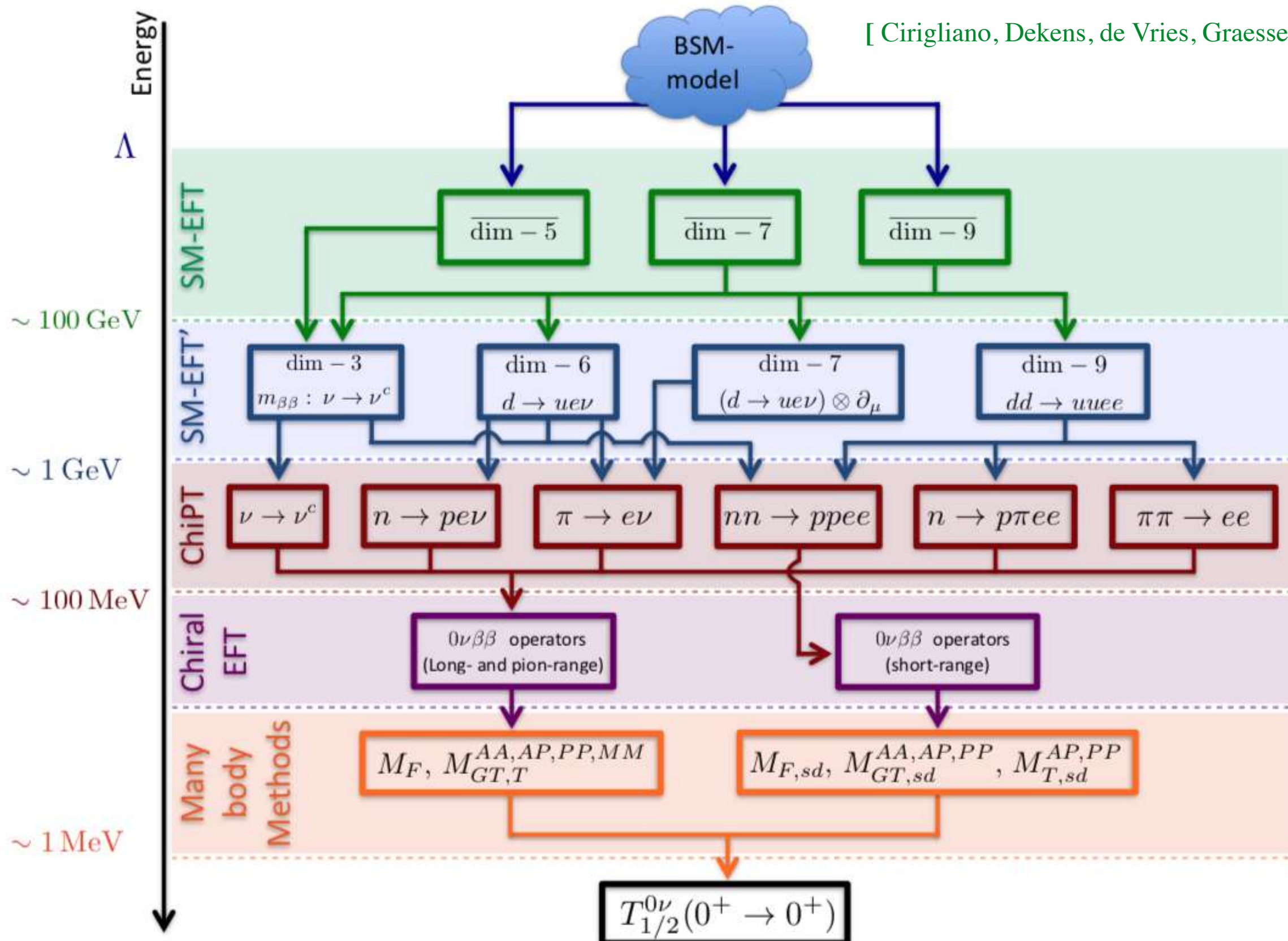
EFT Framework

Model independent systematical parametrization of new physics



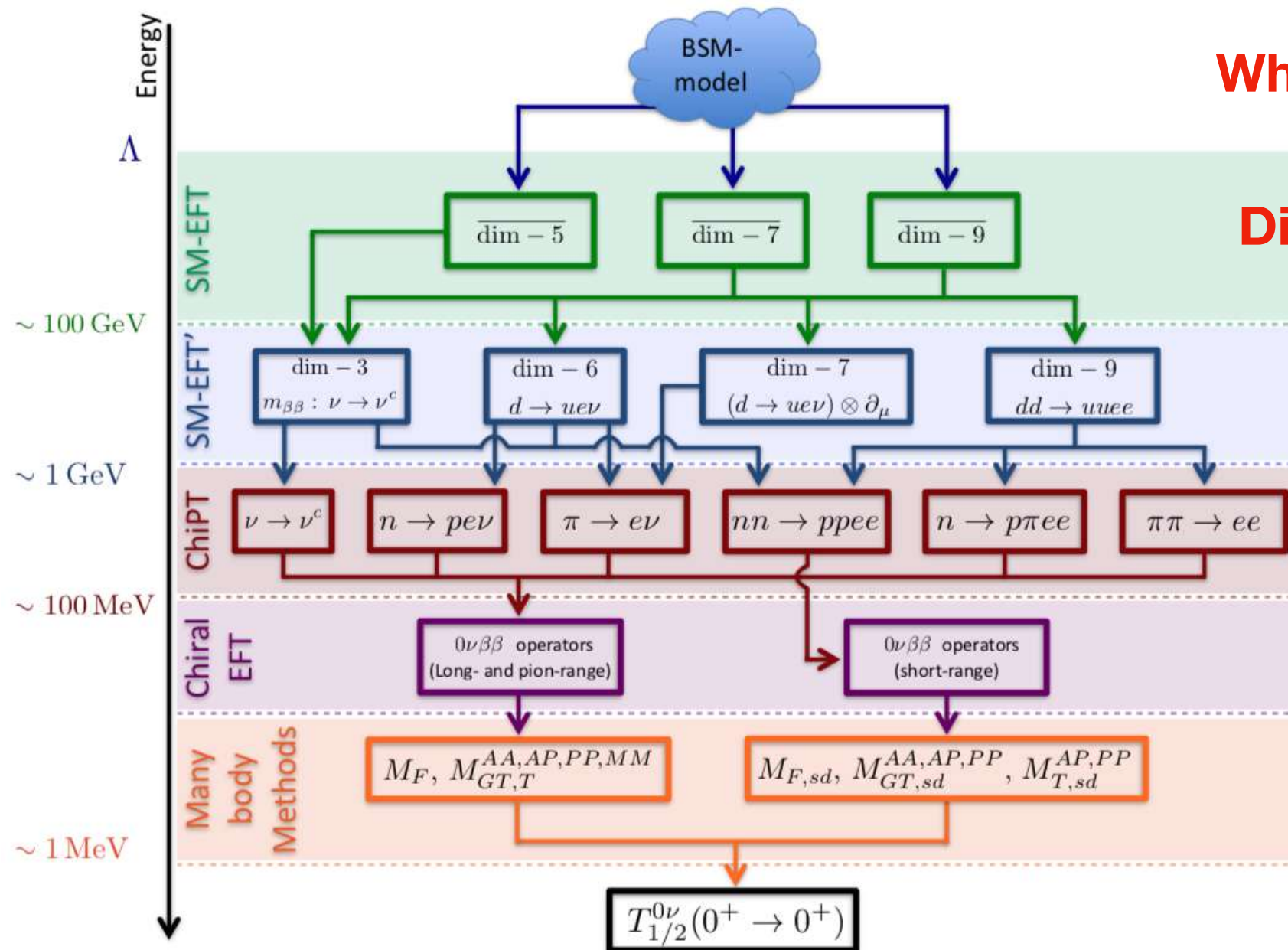
EFT for $0\nu\beta\beta$

[Cirigliano, Dekens, de Vries, Graesser, 2018]



Why Not Enough?

[Cirigliano, Dekens, de Vries, Graesser, 2018]



What is in the UV?

Dim-9 SMEFT not known!

Dim-8/9 LEFT?

More PiN ChiPT?

More Chiral EFT?

More nuclear ME?

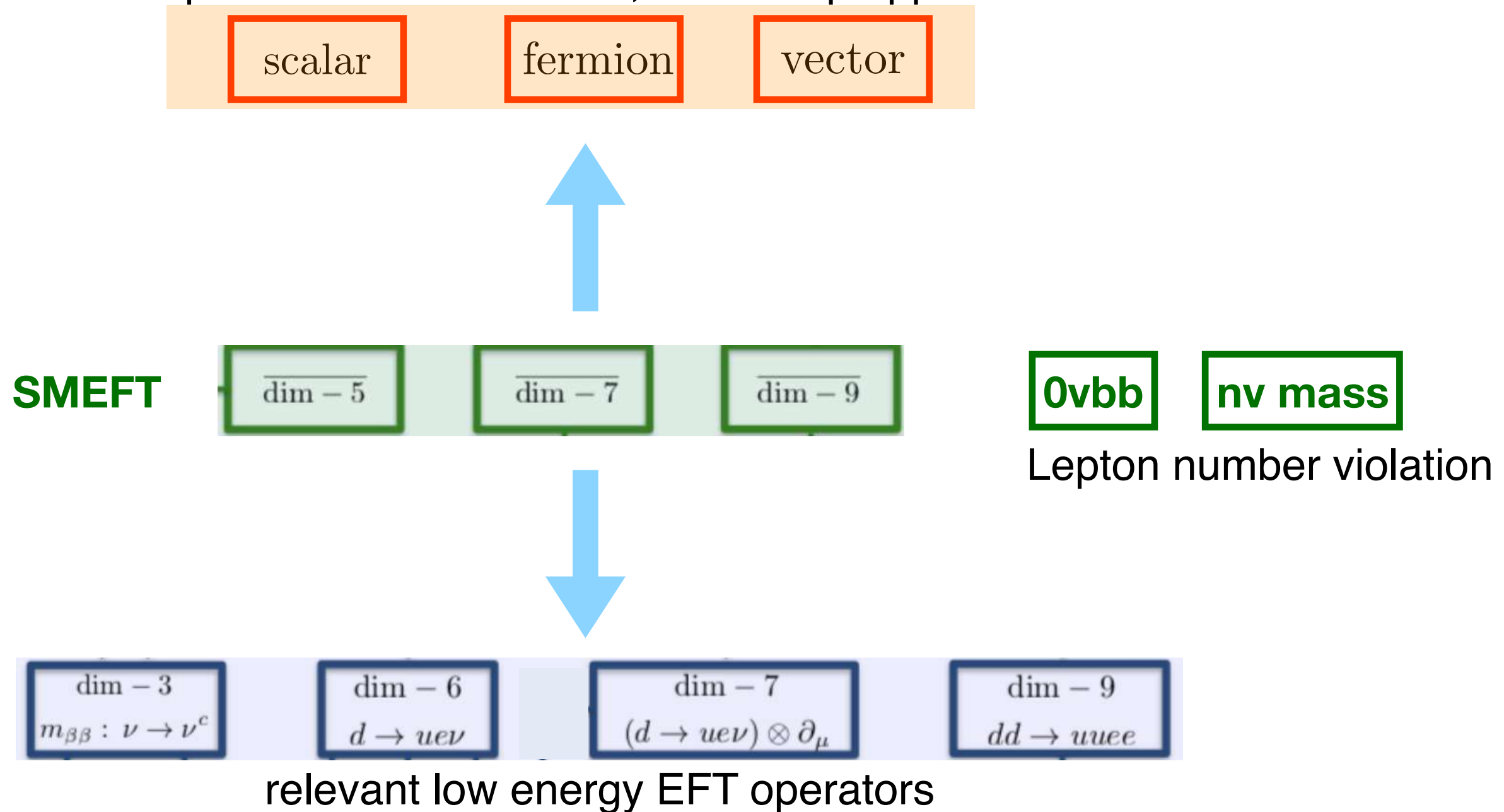
Still top-down: Integrate out and matching/running

SMEFT

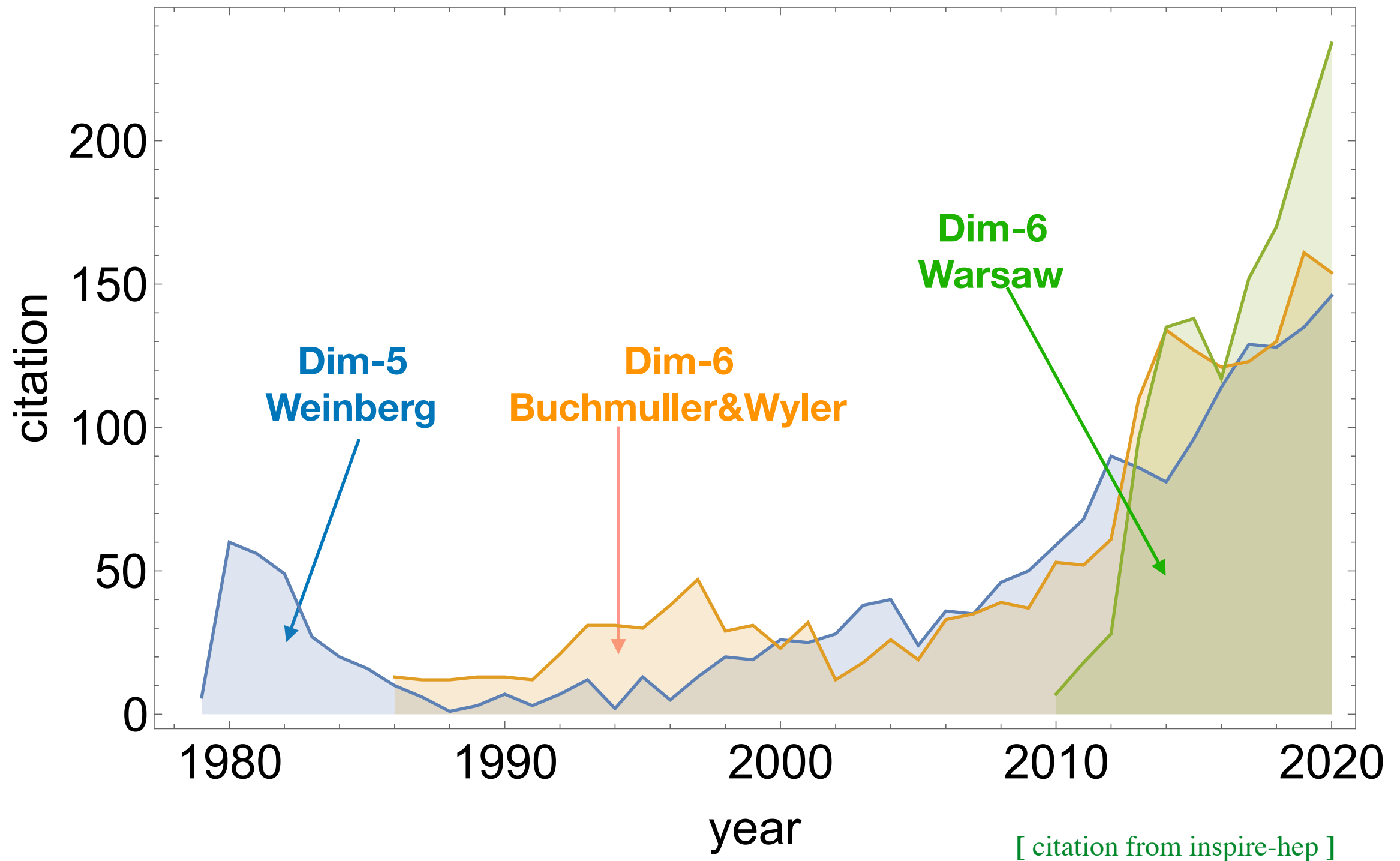
0vbb/nv operators and UV Resonances

SMEFT

Various possible UV realization, bottom-up approach

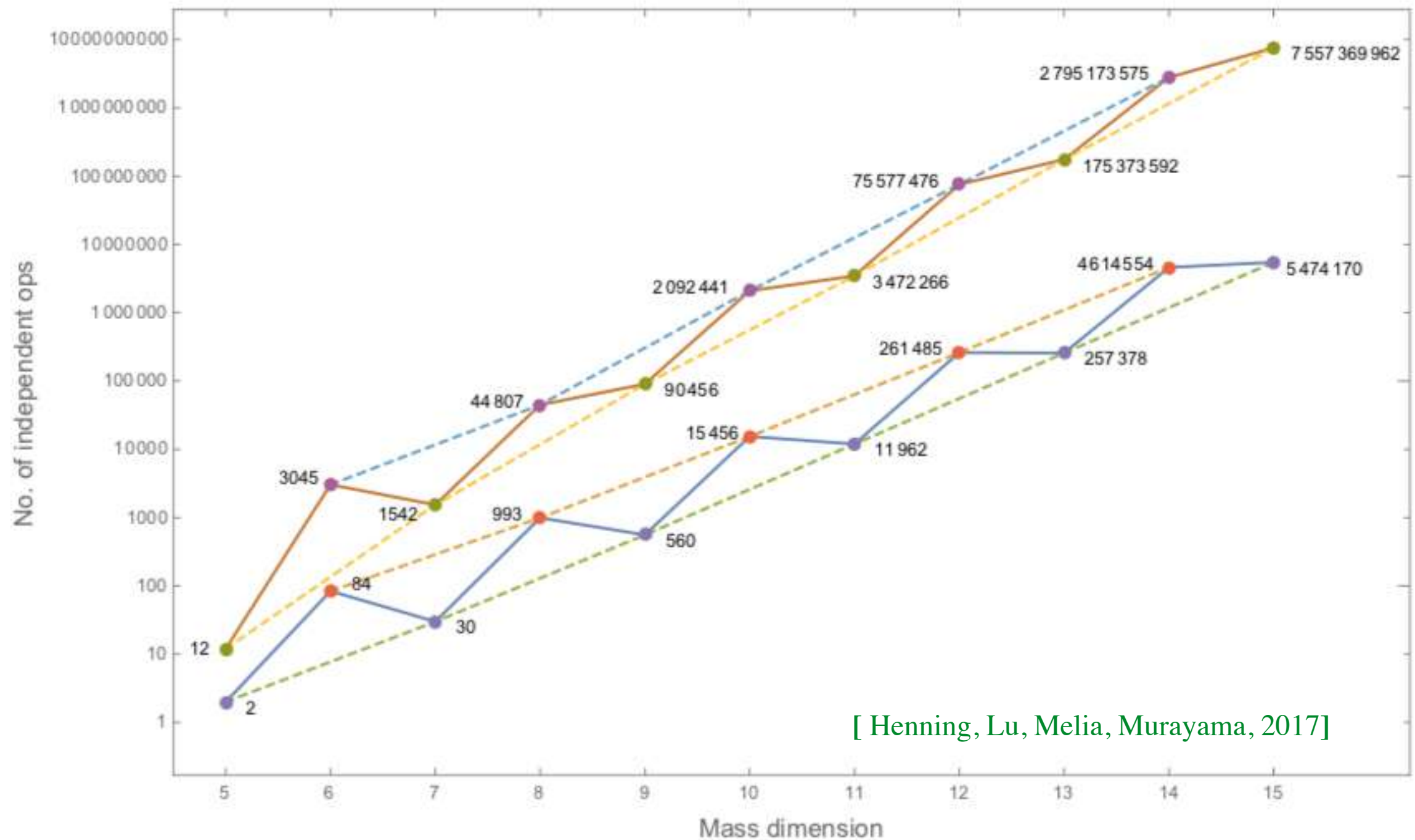


SMEFT Operators



Hilbert Series Counting

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \mathcal{L}_5 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \frac{1}{\Lambda^3} \mathcal{L}_7 + \frac{1}{\Lambda^4} \mathcal{L}_8 + \frac{1}{\Lambda^5} \mathcal{L}_9 + \dots,$$



[Henning, Lu, Melia, Murayama, 2017]

Higher Dim Operators!?

Given numbers of independent operators

does not mean we know explicit form of operators!

Derivatives

$BWHH^\dagger D^2$

2

Repeated fields

$QQQL$

57

$$\begin{aligned}
 & (D^2 H^\dagger) H B_{L\mu\nu} W_L^{\mu\nu}, (D^\mu D_\nu H^\dagger) H B_{L\mu\rho} W_L^{\nu\rho}, (D_\nu D^\mu H^\dagger) H B_{L\mu\rho} W_L^{\nu\rho}, (D_\mu H^\dagger) (D^\mu H) B_{L\nu\rho} W_L^{\nu\rho}, \\
 & (D_\mu H^\dagger) (D^\nu H) B_{L\nu\rho} W_L^{\mu\rho}, (D^\nu H^\dagger) (D_\mu H) B_{L\nu\rho} W_L^{\mu\rho}, (D_\mu H^\dagger) H (D^\mu B_{L\nu\rho}) W_L^{\nu\rho}, (D_\mu H^\dagger) H (D^\nu B_{L\nu\rho}) W_L^{\mu\rho}, \\
 & (D^\nu H^\dagger) H (D_\mu B_{L\nu\rho}) W_L^{\mu\rho}, (D_\mu H^\dagger) H B_{L\nu\rho} (D^\mu W_L^{\nu\rho}), (D_\mu H^\dagger) H B_{L\nu\rho} (D^\nu W_L^{\mu\rho}), (D^\nu H^\dagger) H B_{L\nu\rho} (D_\mu W_L^{\mu\rho}), \\
 & H^\dagger (D^2 H) B_{L\mu\nu} W_L^{\mu\nu}, H^\dagger (D^\mu D_\nu H) B_{L\mu\rho} W_L^{\nu\rho}, H^\dagger (D_\nu D^\mu H) B_{L\mu\rho} W_L^{\nu\rho}, H^\dagger (D^\mu H) (D_\mu B_{L\nu\rho}) W_L^{\nu\rho}, \\
 & H^\dagger (D^\nu H) (D_\mu B_{L\nu\rho}) W_L^{\mu\rho}, H^\dagger (D_\mu H) (D^\nu B_{L\nu\rho}) W_L^{\mu\rho}, H^\dagger (D^\mu H) B_{L\nu\rho} (D_\mu W_L^{\nu\rho}), H^\dagger (D^\nu H) B_{L\nu\rho} (D_\mu W_L^{\mu\rho}), \\
 & H^\dagger (D_\mu H) B_{L\nu\rho} (D^\nu W_L^{\mu\rho}), H^\dagger H (D^2 B_{L\mu\nu}) W_L^{\mu\nu}, H^\dagger H (D^\mu D_\nu B_{L\mu\rho}) W_L^{\nu\rho}, H^\dagger H (D_\nu D^\mu B_{L\mu\rho}) W_L^{\nu\rho}, \\
 & H^\dagger H (D^\mu B_{L\nu\rho}) (D_\mu W_L^{\nu\rho}), H^\dagger H (D^\nu B_{L\nu\rho}) (D_\mu W_L^{\mu\rho}), H^\dagger H (D_\mu B_{L\nu\rho}) (D^\nu W_L^{\mu\rho}), H^\dagger H B_{L\mu\nu} (D^2 W_L^{\mu\nu}), \\
 & H^\dagger H B_{L\mu\rho} (D^\mu D_\nu W_L^{\nu\rho}), H^\dagger H B_{L\mu\rho} (D_\nu D^\mu W_L^{\nu\rho}).
 \end{aligned}
 \tag{14}$$

$$Q_{prst}^{qqql} = C^{prst} \left(\begin{aligned} & \epsilon^{abc} \epsilon^{ik} \epsilon^{jl} (L_{pi} Q_{raj}) (Q_{sbk} Q_{tcl}) \\ & \epsilon^{abc} \epsilon^{ik} \epsilon^{jl} (L_{pi} Q_{sbk}) (Q_{raj} Q_{tcl}) \\ & \epsilon^{abc} \epsilon^{ij} \epsilon^{kl} (L_{pi} Q_{sbk}) (Q_{raj} Q_{tcl}) \\ & \epsilon^{abc} \epsilon^{ij} \epsilon^{kl} (L_{pi} Q_{raj}) (Q_{sbk} Q_{tcl}) \end{aligned} \right) \quad p, r, s, t = 1, 2, 3$$

Which 2 should be picked up?

What flavor relations should be imposed?

Operator as Young Tensor

算符的基元为Lorentz群的不可约表示：取最高权（无需运动方程）

$$H_i \in (0,0) \quad \psi_\alpha \in (1/2,0) \quad F_{L\alpha\beta} = \frac{i}{2} F_{\mu\nu} \sigma_{\alpha\beta}^{\mu\nu} \in (1,0) \quad D_{\alpha\dot{\alpha}} = D_\mu \sigma_{\alpha\dot{\alpha}}^\mu \in (1/2,1/2),$$

$$\partial^2 \phi = (0,0) + (0,1) + (1,0) + (1,1)$$

$$D_{\mu_1} D_{\mu_2} \phi = (D^2 \phi)_{\alpha\beta\dot{\alpha}\dot{\beta}} = \frac{1}{2} \epsilon_{\alpha\beta} \epsilon_{\dot{\alpha}\dot{\beta}} D^\mu D_\mu \phi - \frac{i}{4} \epsilon_{\dot{\alpha}\dot{\beta}} \sigma_{\alpha\beta}^{\mu\nu} [D_\mu, D_\nu] \phi - \frac{i}{4} \epsilon_{\alpha\beta} \bar{\sigma}_{\dot{\alpha}\dot{\beta}}^{\mu\nu} [D_\mu, D_\nu] \phi + \frac{1}{4} (D^2 \phi)_{(\alpha\beta)(\dot{\alpha}\dot{\beta})}.$$

$$\partial\psi = \left(1, \frac{1}{2}\right) + \left(0, \frac{1}{2}\right) \quad (D\psi)_{\alpha\beta\dot{\alpha}} = -\frac{1}{2} \epsilon_{\alpha\beta} (\not{D}\psi)_{\dot{\alpha}} + \frac{1}{2} (D\psi)_{(\alpha\beta)\dot{\alpha}}$$

$$\partial F_L = \left(\frac{1}{2}, \frac{1}{2}\right) \oplus \left(\frac{3}{2}, \frac{1}{2}\right)$$

$$\begin{array}{|c|} \hline \\ \hline \end{array} = [1^2]$$

算符在总动量的小群变换下为：U(N)表示（无需动量积分） $\epsilon^{\alpha_i \alpha_j} \rightarrow \sum_{k,l} \mathcal{U}_k^i \mathcal{U}_l^j \epsilon^{\alpha_k \alpha_l}$

$$\mathcal{M} = (\epsilon^{\alpha_i \alpha_j})^{\otimes n} (\tilde{\epsilon}_{\dot{\alpha}_i \dot{\alpha}_j})^{\otimes \tilde{n}} \prod_{i=1}^N (D^{r_i - |h_i|} \Phi_i)_{\alpha_i^{r_i - h_i}}^{\dot{\alpha}_i^{r_i + h_i}} = \begin{array}{c} [\tilde{n}^{N-2}] \\ N-2 \left\{ \begin{array}{c} \begin{array}{|c|} \hline \\ \hline \end{array} \dots \begin{array}{|c|} \hline \\ \hline \end{array} \\ \vdots \vdots \\ \begin{array}{|c|} \hline \\ \hline \end{array} \dots \begin{array}{|c|} \hline \\ \hline \end{array} \end{array} \right\}_{\tilde{n}} \otimes \underbrace{\begin{array}{|c|} \hline \\ \hline \end{array} \dots \begin{array}{|c|} \hline \\ \hline \end{array}}_n = \begin{array}{c} N-2 \left\{ \begin{array}{c} \begin{array}{|c|} \hline \\ \hline \end{array} \dots \begin{array}{|c|c|} \hline \hline \end{array} \dots \begin{array}{|c|} \hline \\ \hline \end{array} \\ \vdots \vdots \\ \begin{array}{|c|} \hline \\ \hline \end{array} \dots \begin{array}{|c|} \hline \\ \hline \end{array} \end{array} \right\}_{\tilde{n}} \overset{n}{\underbrace{\quad}} + \dots \end{array}$$

SSYT

Operator Construction

$$BWHH^\dagger D^2$$

Li, Ren, Shu, Xiao, **JHYu**, Zheng, arXiv: 2005.00008

Li, Ren, Xiao, **JHYu**, Zheng, arXiv: 2007.07899

$$\begin{aligned} & (D^2 H^\dagger) H B_{L\mu\nu} W_L^{\mu\nu}, (D^\mu D_\nu H^\dagger) H B_{L\mu\rho} W_L^{\nu\rho}, (D_\nu D^\mu H^\dagger) H B_{L\mu\rho} W_L^{\nu\rho}, (D_\mu H^\dagger) (D^\mu H) B_{L\nu\rho} W_L^{\nu\rho}, \\ & (D_\mu H^\dagger) (D^\nu H) B_{L\nu\rho} W_L^{\mu\rho}, (D^\nu H^\dagger) (D_\mu H) B_{L\nu\rho} W_L^{\mu\rho}, (D_\mu H^\dagger) H (D^\mu B_{L\nu\rho}) W_L^{\nu\rho}, (D_\mu H^\dagger) H (D^\nu B_{L\nu\rho}) W_L^{\mu\rho}, \\ & (D^\nu H^\dagger) H (D_\mu B_{L\nu\rho}) W_L^{\mu\rho}, (D_\mu H^\dagger) H B_{L\nu\rho} (D^\mu W_L^{\nu\rho}), (D_\mu H^\dagger) H B_{L\nu\rho} (D^\nu W_L^{\mu\rho}), (D^\nu H^\dagger) H B_{L\nu\rho} (D_\mu W_L^{\mu\rho}), \\ & H^\dagger (D^2 H) B_{L\mu\nu} W_L^{\mu\nu}, H^\dagger (D^\mu D_\nu H) B_{L\mu\rho} W_L^{\nu\rho}, H^\dagger (D_\nu D^\mu H) B_{L\mu\rho} W_L^{\nu\rho}, H^\dagger (D^\mu H) (D_\mu B_{L\nu\rho}) W_L^{\nu\rho}, \\ & H^\dagger (D^\nu H) (D_\mu B_{L\nu\rho}) W_L^{\mu\rho}, H^\dagger (D_\mu H) (D^\nu B_{L\nu\rho}) W_L^{\mu\rho}, H^\dagger (D^\mu H) B_{L\nu\rho} (D_\mu W_L^{\nu\rho}), H^\dagger (D^\nu H) B_{L\nu\rho} (D_\mu W_L^{\mu\rho}), \\ & H^\dagger (D_\mu H) B_{L\nu\rho} (D^\nu W_L^{\mu\rho}), H^\dagger H (D^2 B_{L\mu\nu}) W_L^{\mu\nu}, H^\dagger H (D^\mu D_\nu B_{L\mu\rho}) W_L^{\nu\rho}, H^\dagger H (D_\nu D^\mu B_{L\mu\rho}) W_L^{\nu\rho}, \\ & H^\dagger H (D^\mu B_{L\nu\rho}) (D_\mu W_L^{\nu\rho}), H^\dagger H (D^\nu B_{L\nu\rho}) (D_\mu W_L^{\mu\rho}), H^\dagger H (D_\mu B_{L\nu\rho}) (D^\nu W_L^{\mu\rho}), H^\dagger H B_{L\mu\nu} (D^2 W_L^{\mu\nu}), \\ & H^\dagger H B_{L\mu\rho} (D^\mu D_\nu W_L^{\nu\rho}), H^\dagger H B_{L\mu\rho} (D_\nu D^\mu W_L^{\nu\rho}). \end{aligned} \quad (14)$$

30

highest weight representation

$$(D^{r-|h|} \Phi)_{\alpha^{r-h}}^{\dot{\alpha}^{r+h}} \in \left(\frac{r-h}{2}, \frac{r+h}{2} \right)$$

$$\begin{aligned} & (DH^\dagger)_{\alpha\dot{\alpha}} (DH)_{\beta\dot{\beta}} B_{L\{\gamma\delta\}} W_{L\{\xi\eta\}} \epsilon^{\dot{\alpha}\dot{\beta}} \epsilon^{\alpha\beta} \epsilon^{\gamma\xi} \epsilon^{\delta\eta} \\ & (DH^\dagger)_{\alpha\dot{\alpha}} (DH)_{\beta\dot{\beta}} B_{L\{\gamma\delta\}} W_{L\{\xi\eta\}} \frac{1}{2} \epsilon^{\dot{\alpha}\dot{\beta}} \epsilon^{\delta\xi} (\epsilon^{\alpha\gamma} \epsilon^{\beta\eta} + \epsilon^{\beta\gamma} \epsilon^{\alpha\eta}) \\ & (DH^\dagger)_{\alpha\dot{\alpha}} H (DB_L)_{\{\beta\gamma\delta\},\dot{\beta}} W_{L\{\xi\eta\}} \epsilon^{\dot{\alpha}\dot{\beta}} \epsilon^{\alpha\beta} \epsilon^{\gamma\xi} \epsilon^{\delta\eta} \\ & (DH^\dagger)_{\alpha\dot{\alpha}} H B_{L\{\xi\eta\}} (DW_L)_{\{\beta\gamma\delta\},\dot{\beta}} \epsilon^{\dot{\alpha}\dot{\beta}} \epsilon^{\alpha\beta} \epsilon^{\gamma\xi} \epsilon^{\delta\eta} \\ & H^\dagger (DH)_{\alpha\dot{\alpha}} (DB_L)_{\{\beta\gamma\delta\},\dot{\beta}} W_{L\{\xi\eta\}} \epsilon^{\dot{\alpha}\dot{\beta}} \epsilon^{\alpha\beta} \epsilon^{\gamma\xi} \epsilon^{\delta\eta} \\ & H^\dagger (DH)_{\alpha\dot{\alpha}} B_{L\{\xi\eta\}} (DW_L)_{\{\beta\gamma\delta\},\dot{\beta}} \epsilon^{\dot{\alpha}\dot{\beta}} \epsilon^{\alpha\beta} \epsilon^{\gamma\xi} \epsilon^{\delta\eta} \\ & H^\dagger H (DB_L)_{\{\alpha\beta\gamma\},\dot{\alpha}} (DW_L)_{\{\xi\eta\delta\},\dot{\beta}} \epsilon^{\dot{\alpha}\dot{\beta}} \epsilon^{\alpha\xi} \epsilon^{\beta\eta} \epsilon^{\gamma\delta} \end{aligned}$$

7

$$\mathcal{M} = (\epsilon^{\alpha_i \alpha_j})^{\otimes n} (\tilde{\epsilon}_{\dot{\alpha}_i \dot{\alpha}_j})^{\otimes \tilde{n}} \prod_{i=1}^N (D^{r_i-|h_i|} \Phi_i)_{\alpha_i^{r_i-h_i}}^{\dot{\alpha}_i^{r_i+h_i}} \in [\mathcal{M}]_{N,n,\tilde{n}} = [\mathcal{A}]_{N,n,\tilde{n}} \oplus [\mathcal{B}]_{N,n,\tilde{n}}$$

1	1	1	3
2	2	2	4

1	1	1	2
2	2	3	4

$$B_L^{\alpha\beta} W_{L\alpha\beta} (DH^\dagger)^\gamma_{\dot{\alpha}} (DH)_\gamma^{\dot{\alpha}}, \quad B_L^{\alpha\beta} W_{L\alpha}{}^\gamma (DH^\dagger)_{\beta\dot{\alpha}} (DH)_\gamma^{\dot{\alpha}}$$

2

$$\begin{bmatrix} i & j \end{bmatrix} \times \begin{bmatrix} k \end{bmatrix} \times \begin{bmatrix} l \end{bmatrix} = \begin{bmatrix} i & j \\ k & l \end{bmatrix}$$

$$\epsilon^{ik} \epsilon^{jl} B_L^{\alpha\beta} W_{L\alpha\beta ij} (DH^\dagger)^{\gamma}_{\dot{\alpha}k} (DH)_\gamma^{\dot{\alpha}l}, \quad \epsilon^{ik} \epsilon^{jl} B_L^{\alpha\beta} W_{L\alpha}{}^\gamma{}_{ij} (DH^\dagger)_{\beta\dot{\alpha}k} (DH)_\gamma^{\dot{\alpha}l}.$$

SMEFT

Dimension-5

$$\epsilon_{ij}\epsilon_{mn}(L^iCL^m)H^jH^n$$

[Weinberg, 1979]

2

Dimension-6

Φ^6 and Φ^4D^2	$\psi^2\Phi^3$	X^3
$O_\Phi = (\Phi^\dagger\Phi)^3$	$O_{l\Phi} = (\Phi^\dagger\Phi)(\bar{L}_i l_j\Phi)$	$O_G = -f^{ABC}G_\mu^{AC}G_\nu^{BC}G_\rho^{CA}$
$O_{\Phi\Box} = (\Phi^\dagger\Phi)\Box(\Phi^\dagger\Phi)$	$O_{\sigma\Phi} = (\Phi^\dagger\Phi)(\bar{Q}_i u_j\Phi^c)$	$O_{\tilde{G}} = -f^{ABC}\tilde{G}_\mu^{AC}G_\nu^{BC}G_\rho^{CA}$
$O_{\Phi D} = (\Phi^\dagger D^\mu\Phi)^\dagger(\Phi^\dagger D_\mu\Phi)$	$O_{l\Phi} = (\Phi^\dagger\Phi)(\bar{Q}_i d_j\Phi)$	$O_W = -\epsilon^{abc}W_\mu^{ab}W_\nu^{bc}W_\rho^{ca}$
		$O_{\tilde{W}} = -\epsilon^{abc}\tilde{W}_\mu^{ab}W_\nu^{bc}W_\rho^{ca}$

$X^2\Phi^2$	ψ^2X	$(\bar{L}L)(\bar{L}L)$	$(\bar{R}R)(\bar{R}R)$	$(\bar{L}L)(\bar{R}R)$
$O_{\Phi G} = (\Phi^\dagger\Phi)G_\mu^{AC}G_\nu^{BC}G_\rho^{CA}$	O_{uG}	$O_{LL} = (\bar{L}_i\gamma_\mu L_j)(\bar{L}_k\gamma^\mu L_l)$	$O_{ll} = (\bar{l}_i\gamma_\mu l_j)(\bar{l}_k\gamma^\mu l_l)$	$O_{Ll} = (\bar{L}_i\gamma_\mu L_j)(\bar{l}_k\gamma^\mu l_l)$
$O_{\Phi\tilde{G}} = (\Phi^\dagger\Phi)\tilde{G}_\mu^{AC}G_\nu^{BC}G_\rho^{CA}$	O_{dG}	$O_{Q\bar{Q}}^{(1)} = (\bar{Q}_i\gamma_\mu Q_j)(\bar{Q}_k\gamma^\mu Q_l)$	$O_{uu} = (\bar{u}_i\gamma_\mu u_j)(\bar{u}_k\gamma^\mu u_l)$	$O_{L_u} = (\bar{L}_i\gamma_\mu L_j)(\bar{u}_k\gamma^\mu u_l)$
$O_{\Phi W} = (\Phi^\dagger\Phi)W_\mu^{ab}W_\nu^{bc}W_\rho^{ca}$	O_{lW}	$O_{Q\bar{Q}}^{(3)} = (\bar{Q}_i\gamma_\mu\tau^a Q_j)(\bar{Q}_k\gamma^\mu\tau^a Q_l)$	$O_{dd} = (\bar{d}_i\gamma_\mu d_j)(\bar{d}_k\gamma^\mu d_l)$	$O_{Ld} = (\bar{L}_i\gamma_\mu L_j)(\bar{d}_k\gamma^\mu d_l)$
$O_{\Phi\tilde{W}} = (\Phi^\dagger\Phi)\tilde{W}_\mu^{ab}W_\nu^{bc}W_\rho^{ca}$	O_{sW}	$O_{LQ}^{(1)} = (\bar{L}_i\gamma_\mu L_j)(\bar{Q}_k\gamma^\mu Q_l)$	$O_{ll} = (\bar{l}_i\gamma_\mu l_j)(\bar{l}_k\gamma^\mu l_l)$	$O_{Ql} = (\bar{Q}_i\gamma_\mu Q_j)(\bar{l}_k\gamma^\mu l_l)$
$O_{\Phi B} = (\Phi^\dagger\Phi)B_\mu B_\nu B_\rho$	O_{tW}	$O_{LQ}^{(3)} = (\bar{L}_i\gamma_\mu\tau^a L_j)(\bar{Q}_k\gamma^\mu\tau^a Q_l)$	$O_{ll} = (\bar{l}_i\gamma_\mu l_j)(\bar{d}_k\gamma^\mu d_l)$	$O_{Q_u}^{(1)} = (\bar{Q}_i\gamma_\mu Q_j)(\bar{u}_k\gamma^\mu u_l)$
$O_{\Phi\tilde{B}} = (\Phi^\dagger\Phi)\tilde{B}_\mu B_\nu B_\rho$	O_{bW}		$O_{ud}^{(1)} = (\bar{u}_i\gamma_\mu u_j)(\bar{d}_k\gamma^\mu d_l)$	$O_{Q_u}^{(3)} = (\bar{Q}_i\gamma_\mu\tau^a Q_j)(\bar{u}_k\gamma^\mu\tau^a u_l)$
$O_{\Phi\tilde{B}} = (\Phi^\dagger\Phi)\tilde{B}_\mu B_\nu B_\rho$	O_{tB}		$O_{ud}^{(3)} = (\bar{u}_i\gamma_\mu\tau^a u_j)(\bar{d}_k\gamma^\mu\tau^a d_l)$	$O_{Q_u}^{(1)} = (\bar{Q}_i\gamma_\mu Q_j)(\bar{d}_k\gamma^\mu d_l)$
$O_{\Phi WB} = -(\Phi^\dagger\tau^a\Phi)W_\mu^{ab}B_\nu B_\rho$	O_{tB}	$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$	B-violating	$O_{Ql}^{(3)} = (\bar{Q}_i\gamma_\mu\tau^a Q_j)(\bar{d}_k\gamma^\mu\tau^a d_l)$
$O_{\Phi\tilde{W}B} = -(\Phi^\dagger\tau^a\Phi)\tilde{W}_\mu^{ab}B_\nu B_\rho$	O_{tB}			
	O_{dB}			

[Buchmuller, Wyler, 1986]

[Grzadkowski, Iskrzynski, Misiak, Rosiek, 2010]

Dimension-7

$1 : \psi^2XH^2 + \text{h.c.}$	$2 : \psi^2H^4 + \text{h.c.}$
$Q_{l^2WH^2} \left \epsilon_{mn}(\tau^I\epsilon)_{jk}(l_p^m C i \sigma^{\mu\nu} l_r^n) H^\mu H^\nu W_{\mu\nu}^I \right $	$Q_{l^2H^4} \left \epsilon_{mn}\epsilon_{jk}(l_p^m C l_r^n) H^\mu H^\nu (H^\dagger H) \right $
$Q_{l^2BH^2} \left \epsilon_{mn}\epsilon_{jk}(l_p^m C i \sigma^{\mu\nu} l_r^n) H^\mu H^\nu B_{\mu\nu} \right $	

$3(B) : \psi^4H + \text{h.c.}$	$3(\tilde{B}) : \psi^4H + \text{h.c.}$
$Q_{l^3eH} \left \epsilon_{jk}\epsilon_{mn}(\bar{e}_p l_r^j)(l_s^k C l_t^m) H^n \right $	$Q_{lud^2H} \left \epsilon_{\alpha\beta\gamma}(\bar{l}_p d_r^\alpha)(u_s^\beta C d_t^\gamma) \tilde{H} \right $
$Q_{lewdH} \left \epsilon_{jk}(\bar{d}_p l_r^j)(u_s C e_t) H^k \right $	$Q_{lq^2dH} \left \epsilon_{\alpha\beta\gamma}\epsilon_{jk}(\bar{l}_p d_r^\alpha)(q_{sm}^\beta C q_t^\gamma) \tilde{H}^k \right $
$Q_{l^2qdH}^{(1)} \left \epsilon_{jk}\epsilon_{mn}(\bar{d}_p l_r^j)(q_s^k C l_t^m) H^n \right $	$Q_{ld^2H} \left \epsilon_{\alpha\beta\gamma}(\bar{l}_p d_r^\alpha)(d_s^\beta C d_t^\gamma) H \right $
$Q_{l^2qdH}^{(2)} \left \epsilon_{jm}\epsilon_{kn}(\bar{d}_p l_r^j)(q_s^k C l_t^m) H^n \right $	$Q_{eqd^2H} \left \epsilon_{\alpha\beta\gamma}\epsilon_{jk}(\bar{e}_p q_r^\alpha)(d_s^\beta C d_t^\gamma) \tilde{H}^k \right $
$Q_{l^2quH} \left \epsilon_{jk}(\bar{q}_p u_r)(l_{sm} C l_t^s) H^k \right $	
$4 : \psi^2H^3D + \text{h.c.}$	$5(B) : \psi^4D + \text{h.c.}$
$Q_{leH^3D} \left \epsilon_{mn}\epsilon_{jk}(l_p^m C \gamma^\mu e_r) H^\mu H^j i D_\mu H^k \right $	$Q_{l^2udD} \left \epsilon_{jk}(\bar{l}_p \gamma^\mu u_r)(l_s^j C i D_\mu l_t^k) \right $
$6 : \psi^2H^2D^2 + \text{h.c.}$	$5(\tilde{B}) : \psi^4D + \text{h.c.}$
$Q_{l^2H^2D^2}^{(1)} \left \epsilon_{jk}\epsilon_{mn}(l_p^j C D^\mu l_r^k) H^m (D_\mu H^n) \right $	$Q_{lq^2D} \left \epsilon_{\alpha\beta\gamma}(\bar{l}_p \gamma^\mu q_r^\alpha)(d_s^\beta C i D_\mu d_t^\gamma) \right $
$Q_{l^2H^2D^2}^{(2)} \left \epsilon_{jm}\epsilon_{kn}(l_p^j C D^\mu l_r^k) H^m (D_\mu H^n) \right $	$Q_{ed^2D} \left \epsilon_{\alpha\beta\gamma}(\bar{e}_p \gamma^\mu d_r^\alpha)(d_s^\beta C i D_\mu d_t^\gamma) \right $

[Lehman, 2014]
[Liao, Ma, 2018]

Dimension-8

[Li, Ren, Shu, Xiao, **Yu**, Zheng, 2020]

N	(n, \bar{n})	Subclasses	N_{type}	N_{norm}	N_{operator}	Equations
4	(4, 0)	$F_L^4 + h.c.$	14	26	26	(4.19)
		$F_L^2\psi\psi^1D + h.c.$	22	22	$22n_f^2$	(4.51)
		$\psi^4D^2 + h.c.$	4+4	18+14	$12n_f^4 + n_f^3(5n_f - 1)$	(4.75, 4.78, 4.80)
		$F_L\psi^2\phi D^2 + h.c.$	16	32	$32n_f^2$	(4.44)
		$F_L^2\psi^2D^2 + h.c.$	8	12	12	(4.14)
	(2, 2)	$F_L^2F_R^2$	14	17	17	(4.19)
		$F_LF_R\psi\psi^1D$	27	35	$35n_f^2$	(4.50, 4.51)
		$\psi^2\psi^{12}D^2$	17+4	54+8	$\frac{1}{2}n_f^2(75n_f^2 + 11) + 6n_f^4$	(4.74, 4.79-4.81)
		$F_R\psi^2\phi D^2 + h.c.$	16	16	$16n_f^2$	(4.44)
		$F_LF_R\phi^2D^2$	5	6	6	(4.14)
5	(3, 0)	$\psi\psi^1\phi^2D^3$	7	16	$16n_f^2$	(4.31, 4.32)
		ϕ^4D^4	1	3	3	(4.8)
	(2, 1)	$F_L\psi^4 + h.c.$	12+10	66+54	$42n_f^4 + 2n_f^3(9n_f + 1)$	(4.86, 4.88, 4.89, 4.91)
		$F_L^2\psi^2\phi + h.c.$	32	60	$60n_f^2$	(4.47, 4.48)
		$F_L^3\phi^2 + h.c.$	6	6	6	(4.16)
		$F_L\psi^2\psi^{12} + h.c.$	84+24	172+32	$2n_f^2(59n_f^2 - 2) + 24n_f^4$	(4.84-4.85), (4.88-4.92)
		$F_R^2\psi^2\phi + h.c.$	32	36	$36n_f^2$	(4.47, 4.48)
		$\psi^3\psi^1\phi D + h.c.$	32+14	180+56	$n_f^3(135n_f - 1) + n_f^2(29n_f + 3)$	(4.66, 4.69-4.72)
		$F_L\psi\psi^1\phi^2D + h.c.$	38	92	$92n_f^2$	(4.39, 4.40)
6	(2, 0)	$\psi^2\phi^3D^2 + h.c.$	6	36	$36n_f^2$	(4.28)
		$F_L\phi^4D^2 + h.c.$	4	6	6	(4.10)
	(1, 1)	$\psi^4\phi^2 + h.c.$	12+4	48+18	$5(5n_f^4 + n_f^2) + \frac{3}{2}(8n_f^2 + n_f^2)$	(4.55, 4.59, 4.62, 4.64)
		$F_L\psi^2\phi^3 + h.c.$	16	22	$22n_f^2$	(4.36)
		$F_L^2\phi^4 + h.c.$	8	10	10	(4.12)
		$\psi^2\psi^{12}\phi^2$	23+10	57+14	$n_f^2(42n_f^2 + n_f + 2) + 3n_f^3(3n_f - 1)$	(4.54, 4.55, 4.59-4.63)
		$\psi\psi^1\phi^4D$	7	13	$13n_f^2$	(4.24, 4.25)
		ϕ^6D^2	1	2	2	(4.8)
	(1, 0)	$\psi^4\phi^3 + h.c.$	6	6	$6n_f^2$	(4.21)
7	(0, 0)	ϕ^8	1	1	1	(4.8)
Total		48	471+70	1070+196	993($n_f = 1$), 44807($n_f = 3$)	

[Murphy, 2020]

993

Dimension-9

[Li, Ren, Shu, Xiao, **Yu**, Zheng, 2020]

N	(n, \bar{n})	Classes	N_{type}	N_{norm}	N_{operator}	Equations
4	(3, 2)	$\psi^3\psi^1D^3 + h.c.$	0 + 4 + 2 + 0	10	$\frac{2}{3}n_f^2(7n_f^2 - 1)$	(5.50)(5.51)
		$\psi^2\phi^2D^4 + h.c.$	0 + 0 + 2 + 0	6	$3n_f(n_f + 1)$	(5.21)
5	(3, 1)	$F_L\psi^3\psi^1D + h.c.$	0 + 10 + 6 + 0	72	$32n_f^4$	(5.59)(5.60)
		$\psi^4\phi D^2 + h.c.$	0 + 4 + 4 + 0	100	$40n_f^4$	(5.45-5.48)
		$F_L\psi^2\phi^2D^2 + h.c.$	0 + 0 + 4 + 0	34	$17n_f^2 - n_f$	(5.28)(5.29)
	(2, 2)	$F_R\psi^3\psi^1D + h.c.$	0 + 10 + 6 + 0	54	$4n_f^3(6n_f + 1)$	(5.59)(5.60)
		$\psi^2\psi^{12}\phi D^2$	0 + 4 + 4 + 0	84	$n_f^3(49n_f + 1)$	(5.45-5.48)
		$F_R\psi^2\phi^2D^2 + h.c.$	0 + 0 + 4 + 0	20	$2n_f(5n_f - 1)$	(5.28)(5.29)
		$\psi\psi^1\phi^3D^3$	0 + 0 + 2 + 0	6	$6n_f^2$	(5.19)
	(3, 0)	$\psi^6 + h.c.$	2 + 4 + 6 + 0	116	$\frac{1}{24}n_f^2(415n_f^4 + 53n_f^3 + 59n_f^2 + 139n_f + 6)$	(5.63-5.70)
		$F_L\psi^4\phi + h.c.$	0 + 12 + 10 + 0	102	$2n_f^3(21n_f + 1)$	(5.54-5.56)
		$F_L^2\psi^2\phi^2 + h.c.$	0 + 0 + 8 + 0	20	$2n_f(5n_f + 2)$	(5.32)
6	(2, 1)	$\psi^4\psi^{12} + h.c.$	4 + 26 + 20 + 4	244	$\frac{1}{6}n_f^3(382n_f^2 - 9n_f^2 + 2n_f + 21)$	(5.63-5.69)
		$F_L\psi^2\psi^{12}\phi + h.c.$	0 + 24 + 24 + 0	92	$52n_f^4$	(5.54-5.56)
		$F_L^2\psi^{12}\phi^2 + h.c.$	0 + 0 + 8 + 0	12	$2n_f(3n_f + 2)$	(5.32)
		$\psi^3\psi^1\phi^2D + h.c.$	0 + 12 + 18 + 0	186	$\frac{2}{3}n_f^2(146n_f^2 + 1)$	(5.39-5.42)
		$F_L\psi\psi^1\phi^3D + h.c.$	0 + 0 + 8 + 0	12	$12n_f^2$	(5.25)
		$\psi^2\phi^4D^2 + h.c.$	0 + 0 + 4 + 0	24	$2n_f(6n_f + 1)$	(5.17)
	(2, 0)	$\psi^4\phi^3 + h.c.$	0 + 6 + 6 + 0	32	$\frac{4}{3}n_f^2(10n_f^2 - 1)$	(5.35-5.37)
		$F_L\psi^2\phi^4 + h.c.$	0 + 0 + 4 + 0	8	$2n_f(2n_f - 1)$	(5.23)
7	(1, 1)	$\psi^2\psi^{12}\phi^3$	0 + 6 + 10 + 0	24	$14n_f^4$	(5.35-5.37)
		$\psi\psi^1\phi^5D$	0 + 0 + 2 + 0	2	$2n_f^2$	(5.12)
8	(1, 0)	$\psi^2\phi^6 + h.c.$	0 + 0 + 2 + 0	2	$n_f^2 + n_f$	(5.9)
Total		42	6+122+164+4	1262	8 + 204 + 348 + 0 ($n_f = 1$) 2862 + 42234 + 44874 + 486 ($n_f = 3$)	

[Liao, Ma, 2020]

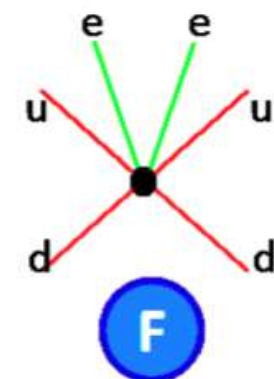
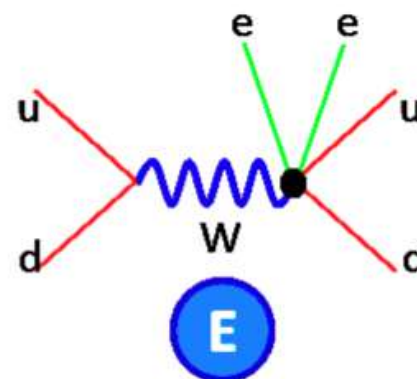
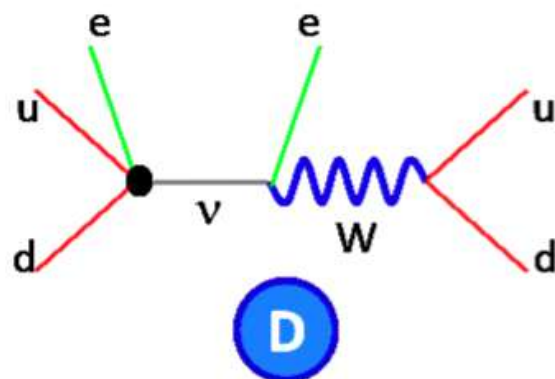
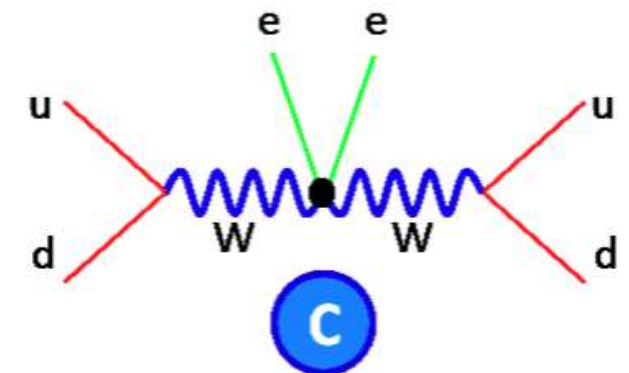
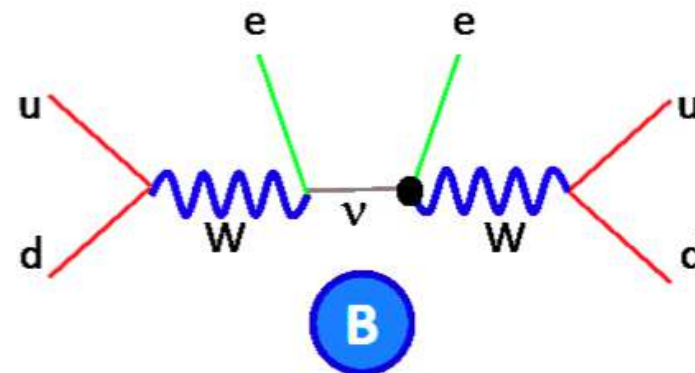
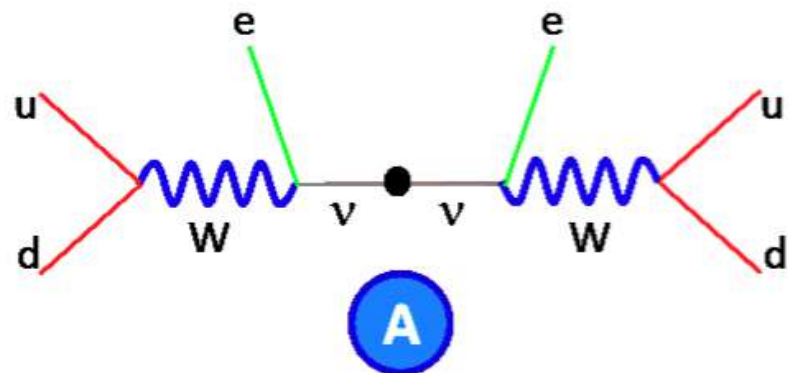
560

Jiang-Hao Yu

18

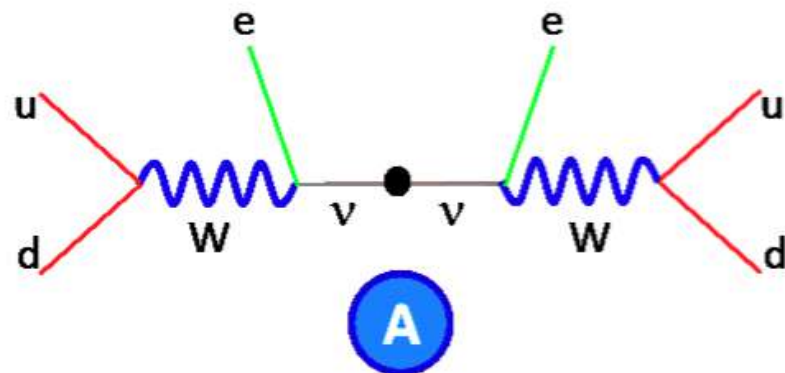
0vbb Related Operators

SMEFT broken phase:



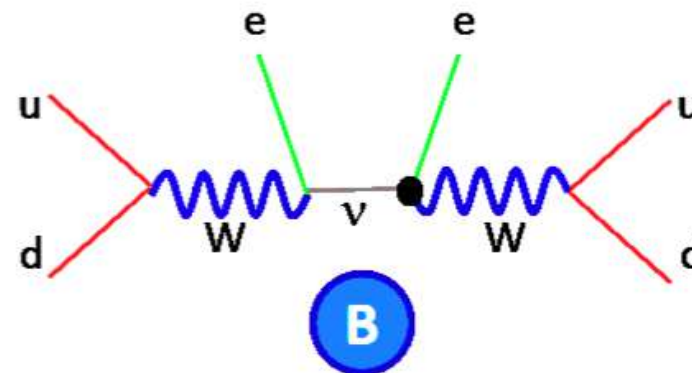
0vbb Related Operators

Relate to SMEFT unbroken operators:



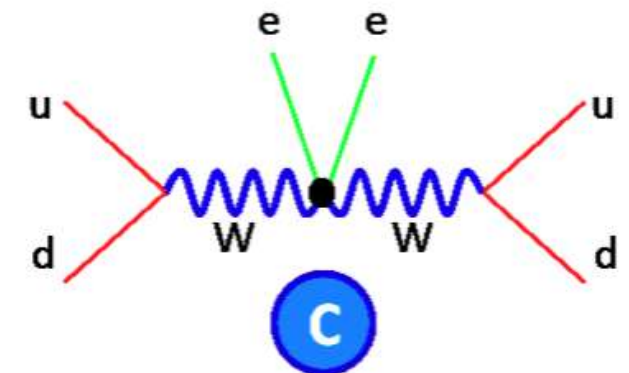
$$(\bar{\ell}_\alpha \phi) (\tilde{\phi}^\dagger \ell_\beta)$$

Dim-5



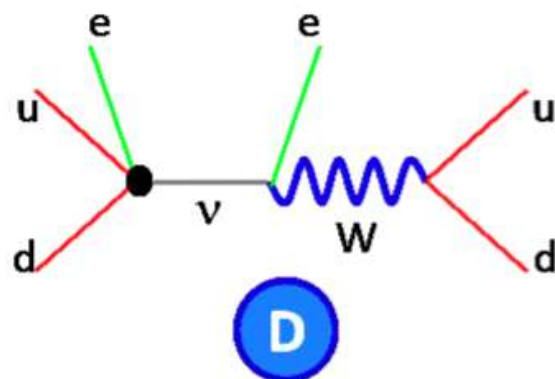
$$(\phi^\dagger D_\mu \tilde{\phi}) (\phi^\dagger \bar{e}_{\alpha R} \gamma^\mu \tilde{\ell}_\beta)$$

Dim-7, 9

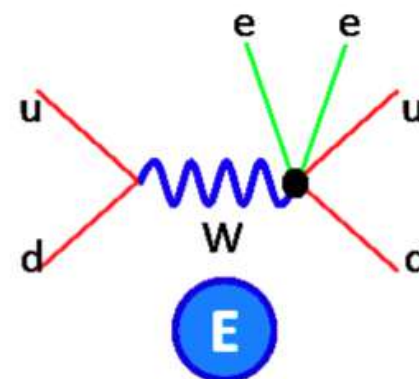


$$\bar{e}_{\alpha R} e_{\beta R}^c (\phi^\dagger D \tilde{\phi})^2$$

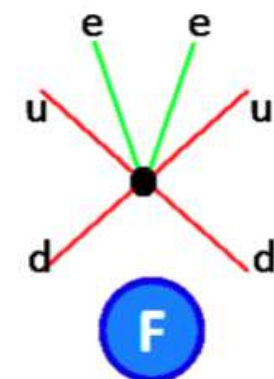
Dim-7, 9



Dim-7, 9



Dim-9

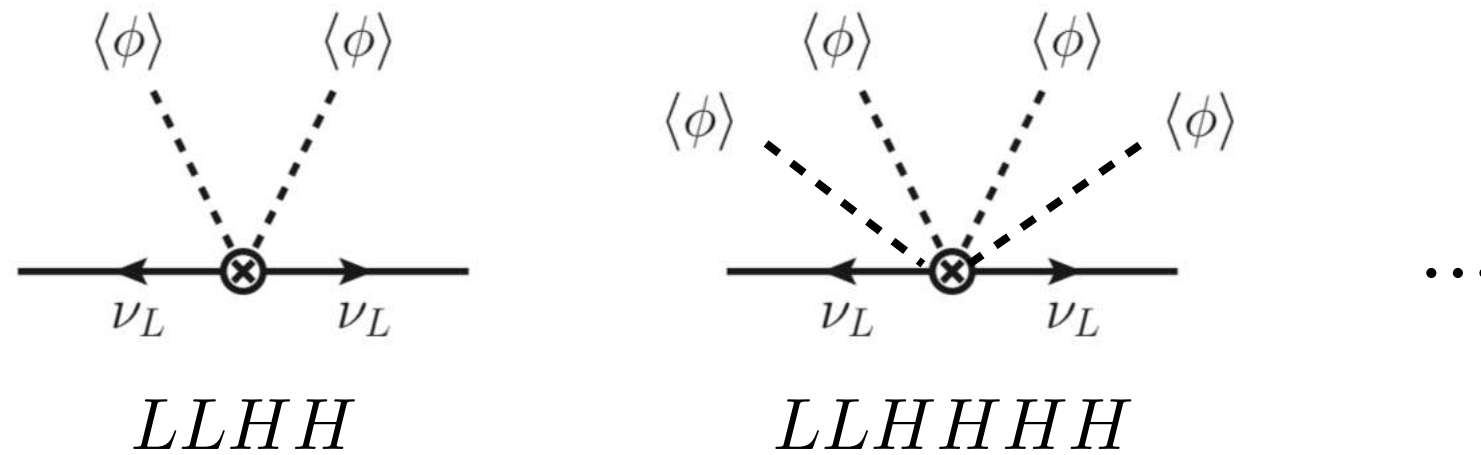


Dim-9

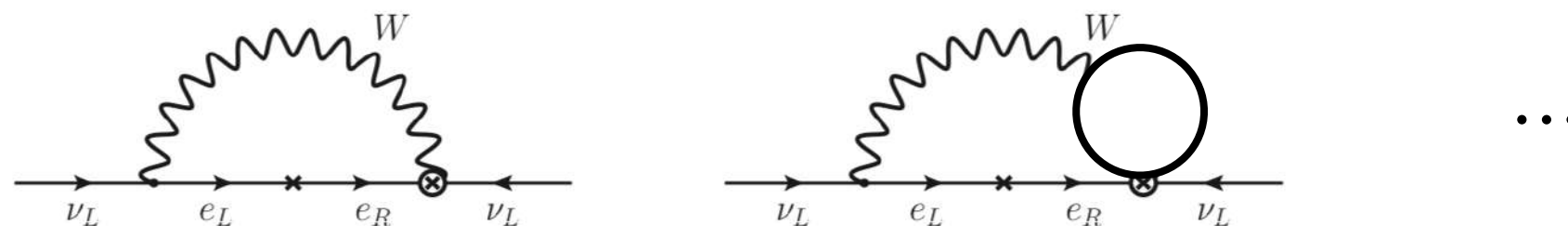
Not complete

Nv Mass Related Operators

Higgs taking VEV:

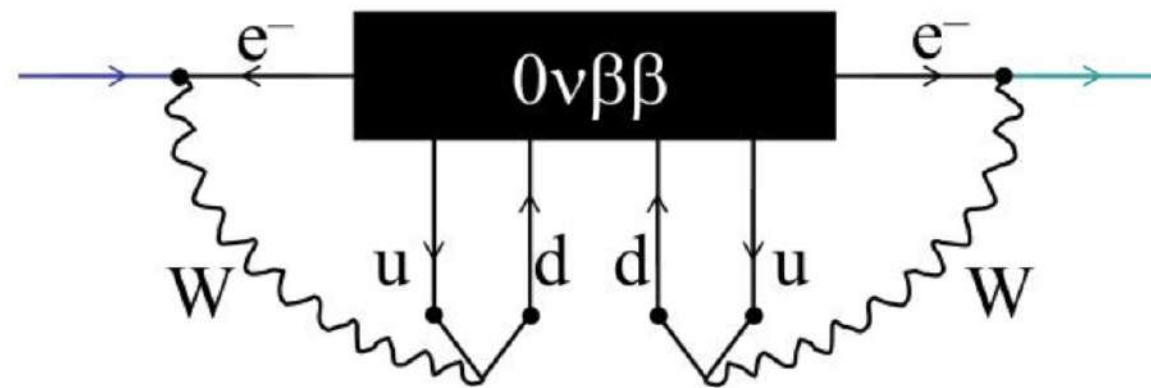


Could from other lepton number violation operators: anomalous RG

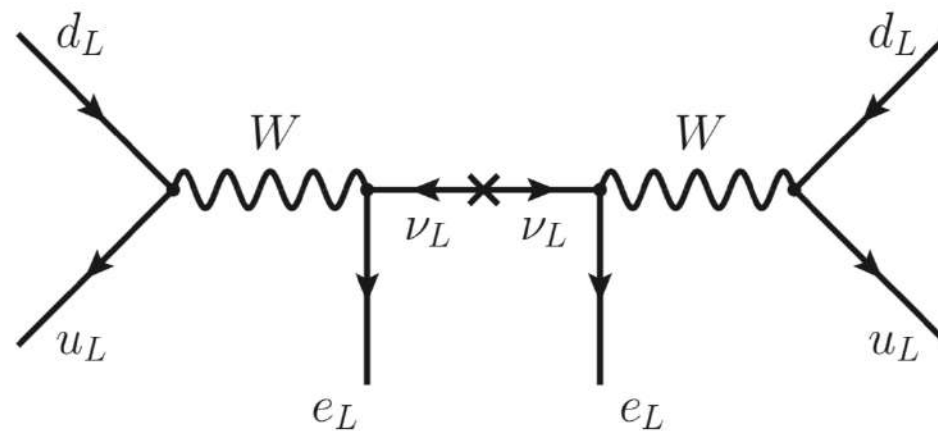


Neutrino Masses and $0\nu\beta\beta$

Schechter-Valle Theorem: whatever processes cause $0\nu\beta\beta$, its observation would imply existence of Majorana mass term

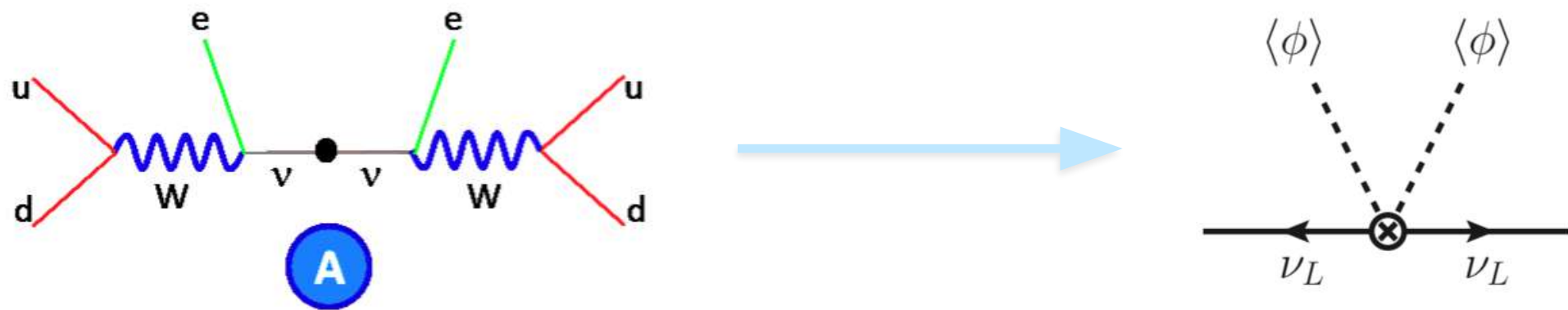


[Schechter-Valle, 1982]



Strong Correlation

Standard mechanism: origin of $0\nu\beta\beta$ = origin of neutrino masses

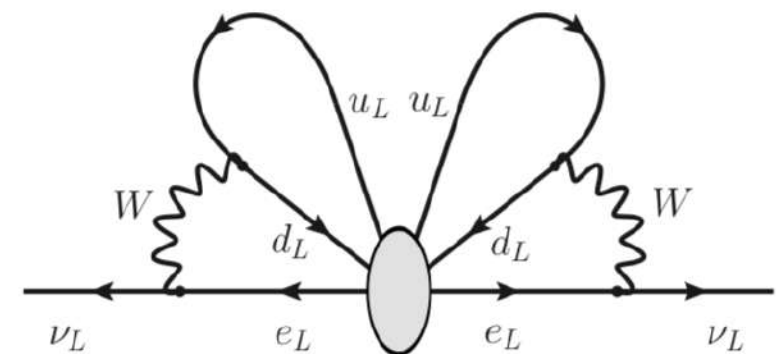
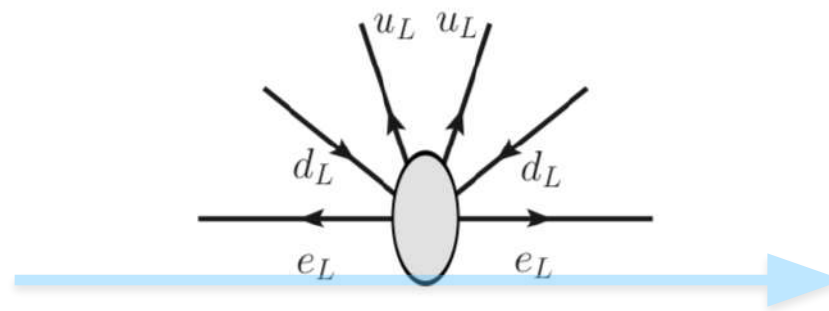
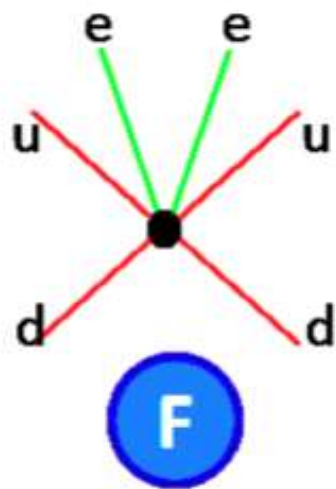


$$L^{\mu\nu} = - \int \int dx_2 dx_1 \sum_i \bar{e}(x_1) \gamma^\mu (1 - \gamma_5) U_{ei} \underbrace{\nu_{iL}(x_1) \bar{\nu}_{iL}^c(x_2)} \gamma_\nu (1 + \gamma_5) U_{ei} e_L^c(x_2) \quad m_{\beta\beta} = \sum_{i=1}^3 U_{ei}^2 m_{\nu_i}.$$

$0\nu\beta\beta$ has direct connection to neutrino physics

Strong Correlation???

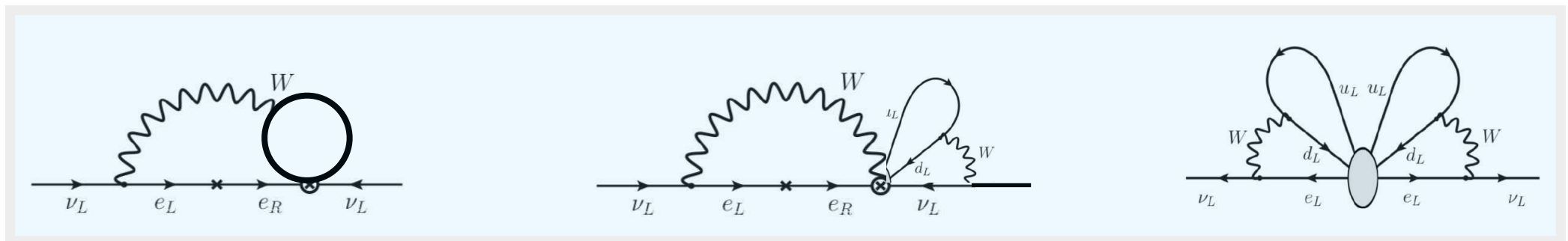
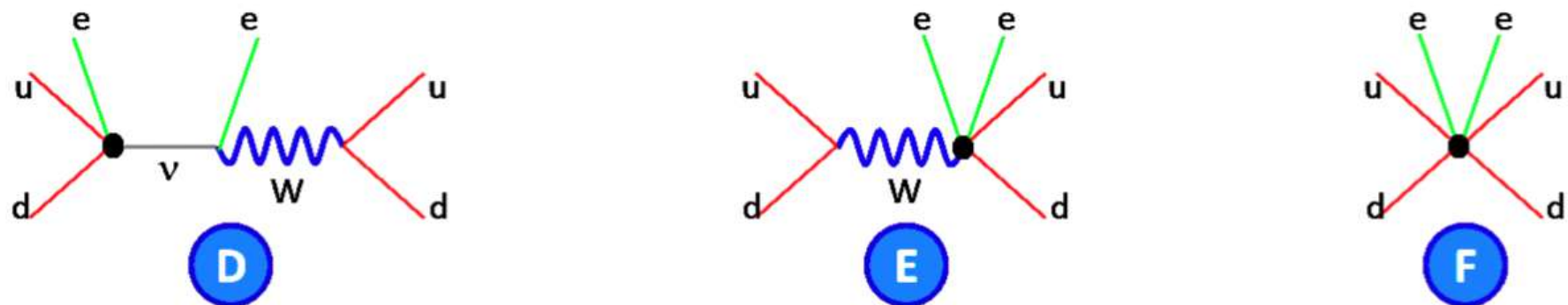
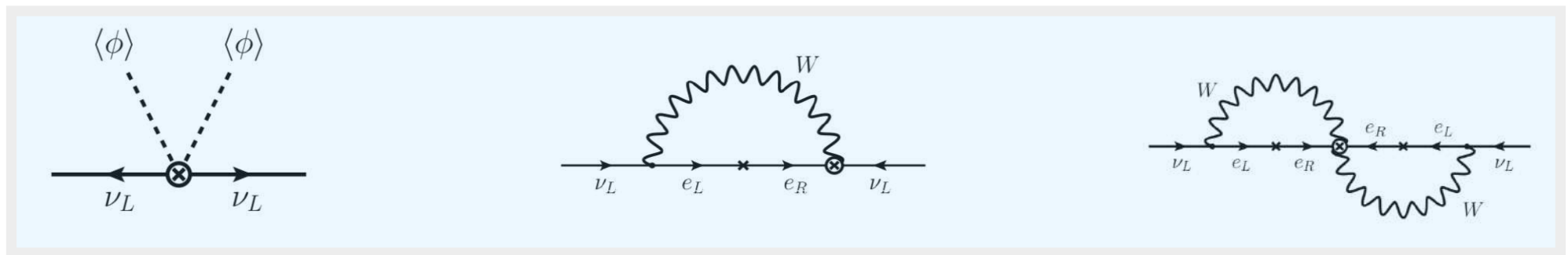
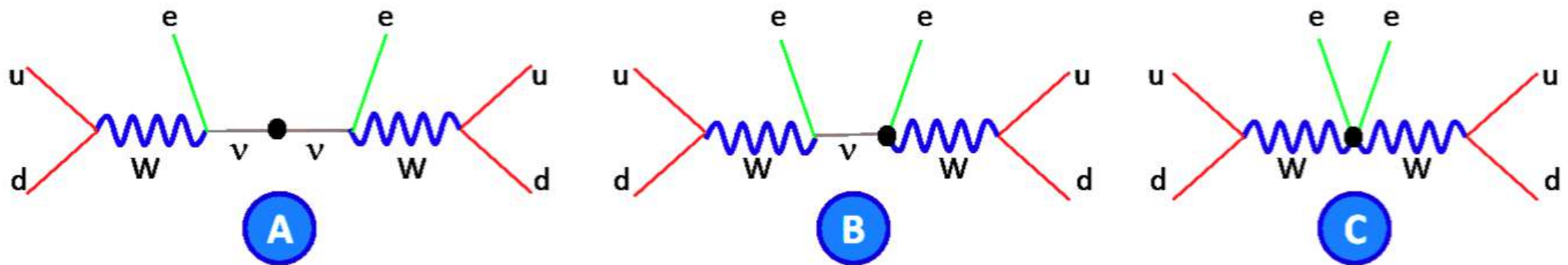
Lepton number violation operator: origin of $0\nu\beta\beta$ = small part of ν mass



Very tiny neutrino mass

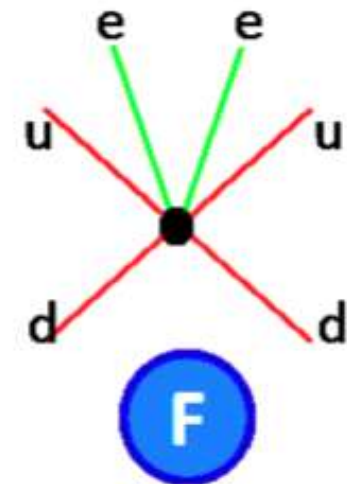
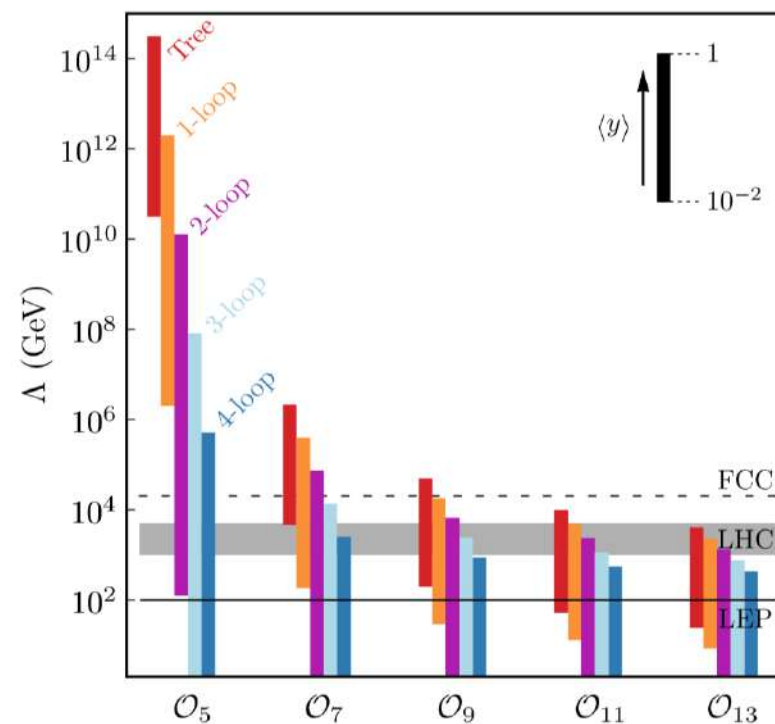
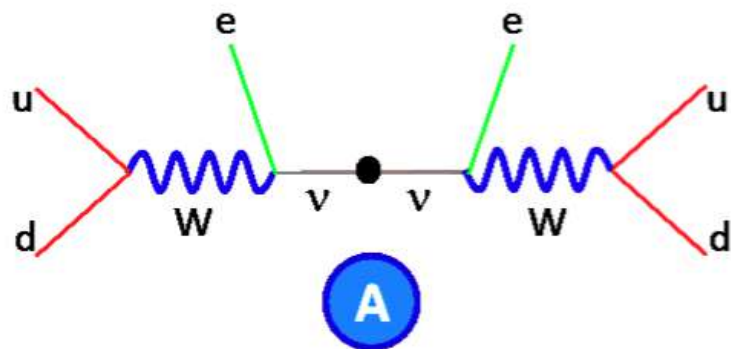
$0\nu\beta\beta$ does not need to connect to current neutrino exp.

Neutrino Masses and $0\nu\beta\beta$



Which One Dominate $0\nu b\bar{b}$?

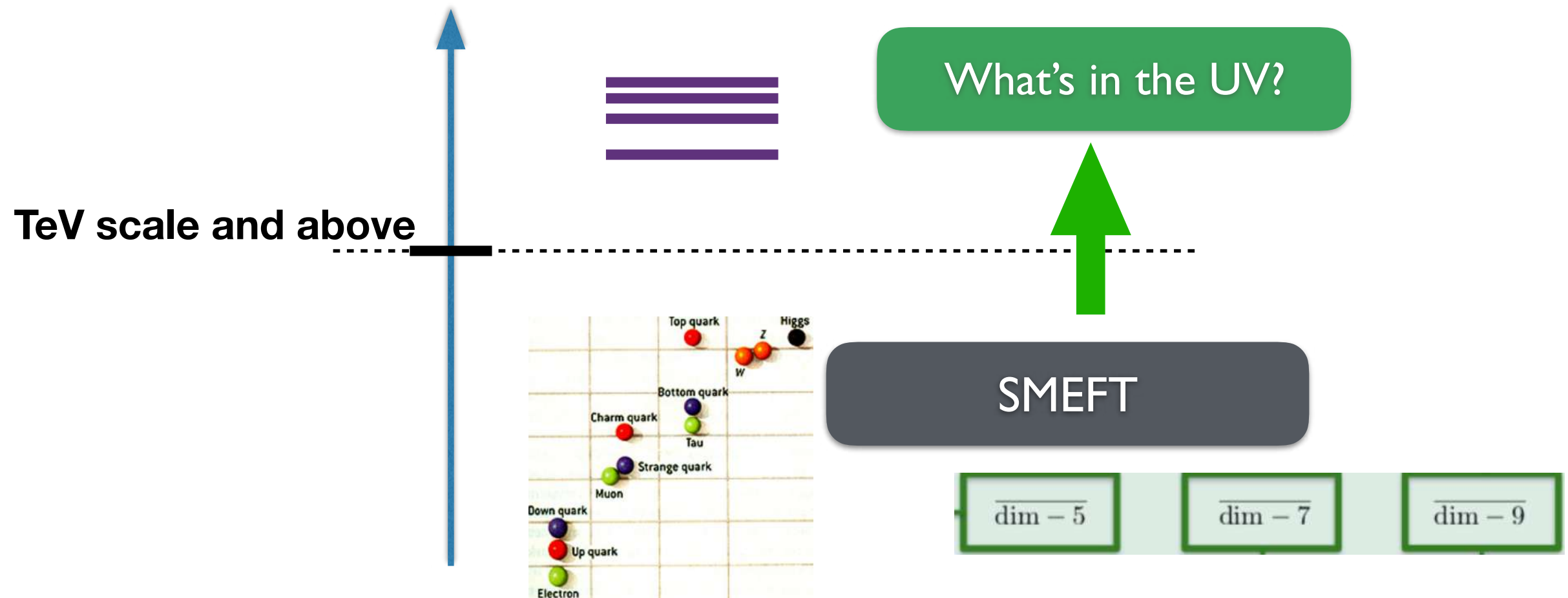
The higher dim operator, the lower cutoff scale



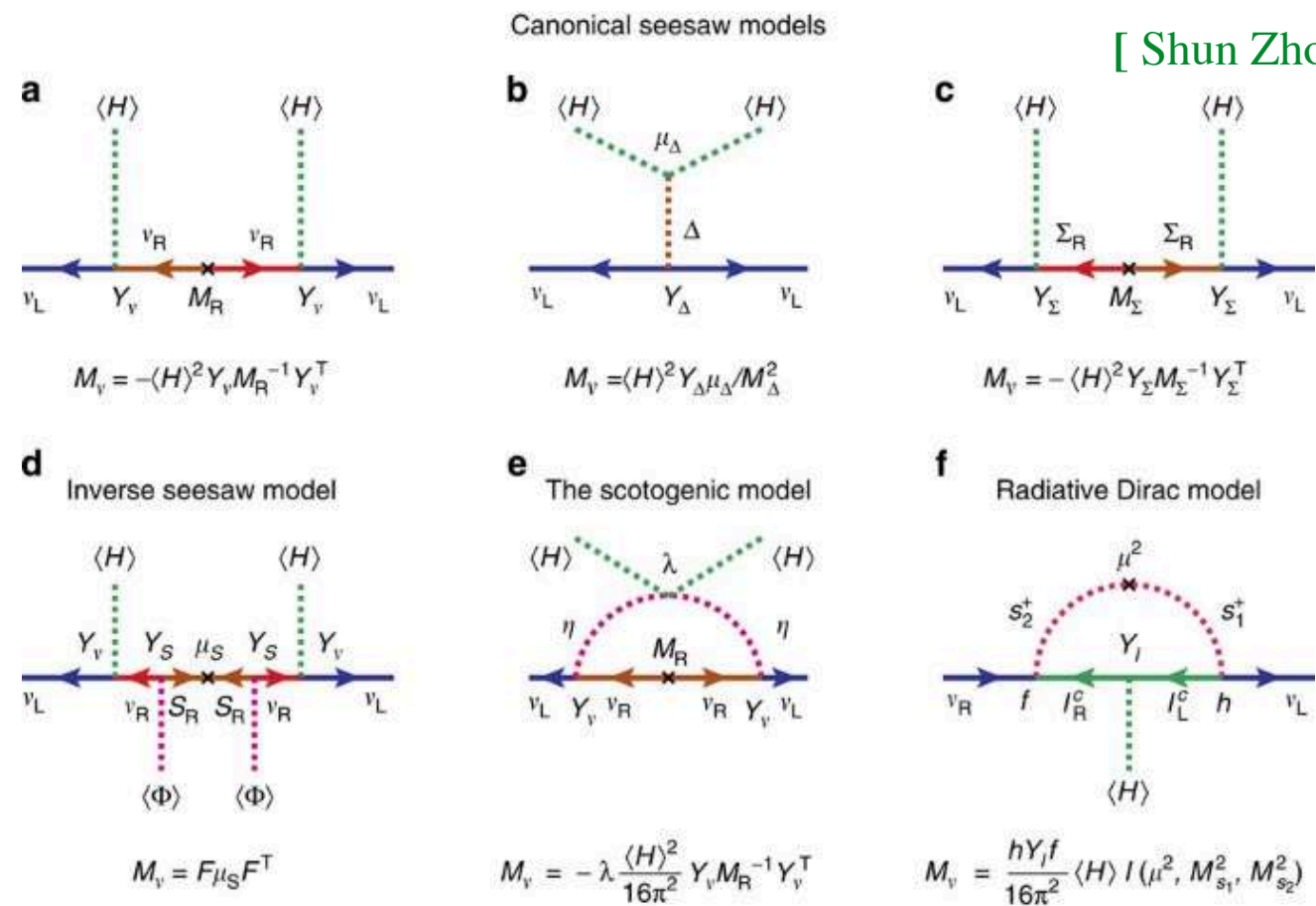
Could be comparable!

Not necessarily related to neutrino physics

What is in the UV?



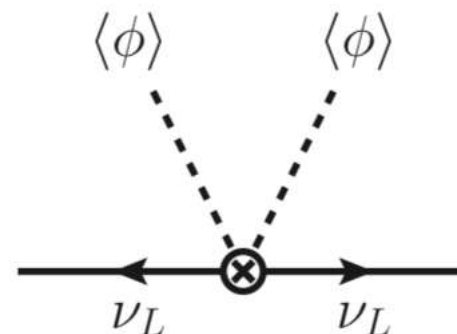
UV Realization of $N\nu$ Masses



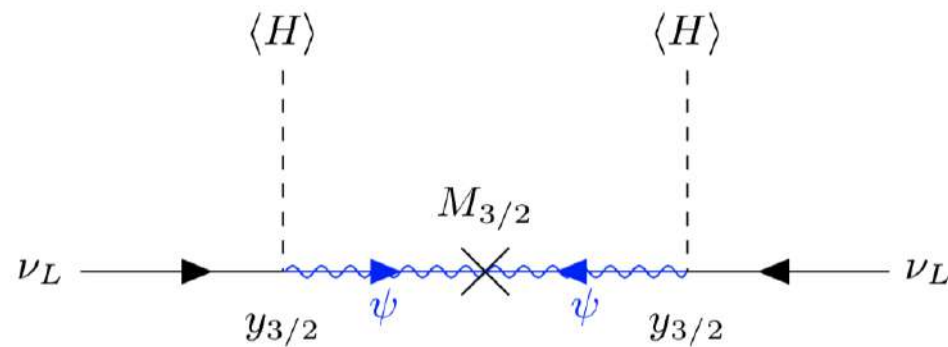
[Shun Zhou, 2016]

Integrate out heavy particles

Top-down approach



More UV realization?



Type-3/2 Seesaw Mechanism

Durmuş Demir,¹ Canan Karahan,^{2,*} and Ozan Sargin³

¹*Sabancı University, Faculty of Engineering and Natural Sciences, 34956 Tuzla İstanbul, Turkey*

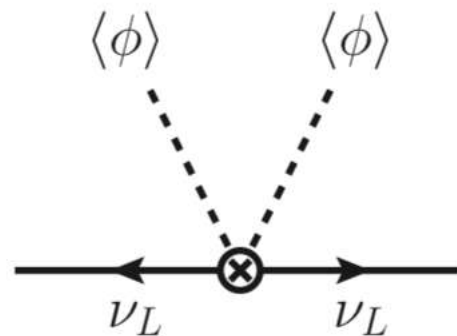
²*Physics Engineering Department, İstanbul Technical University, 34469 Maslak İstanbul, Turkey*

³*İzmir Institute of Technology, Department of Physics, 35430, İzmir, Turkey*

(Dated: May 17, 2021)



Bottom-up Approach



J-Basis Operator: Partial Wave

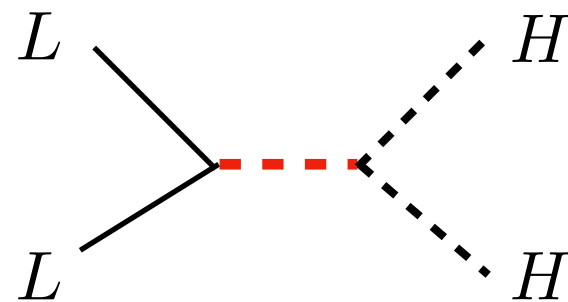
$$\mathcal{Y} [\boxed{p \ r}] \epsilon_{ik} \epsilon_{jl} \epsilon_{\alpha\beta} L_p^{\alpha i} L_r^{\beta j} H^k H^l$$

Partial wave expansion on operator

$$\mathbf{W}^2 \mathcal{B}^J = -s J(J+1) \mathcal{B}^J$$

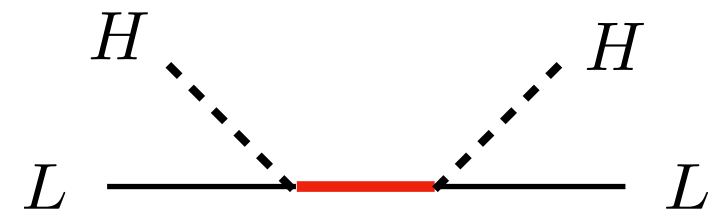
$$\mathbf{W}^2 = \frac{s}{8} \sum_{i,j=1}^N (\langle i, \partial_j \rangle \langle j, \partial_i \rangle + [i, \partial_j][j, \partial_i]) - \frac{1}{4} \sum_{i,j,k,l} [i, j] \langle j, \partial_k \rangle \langle k, l \rangle [l, \partial_i]$$

$LL \rightarrow HH$ channel



$J = 0$

$LH \rightarrow LH$ channel



$J = \frac{1}{2}$

Type-II: **SU(2) triplet**, or singlet (excluded by repeated field)

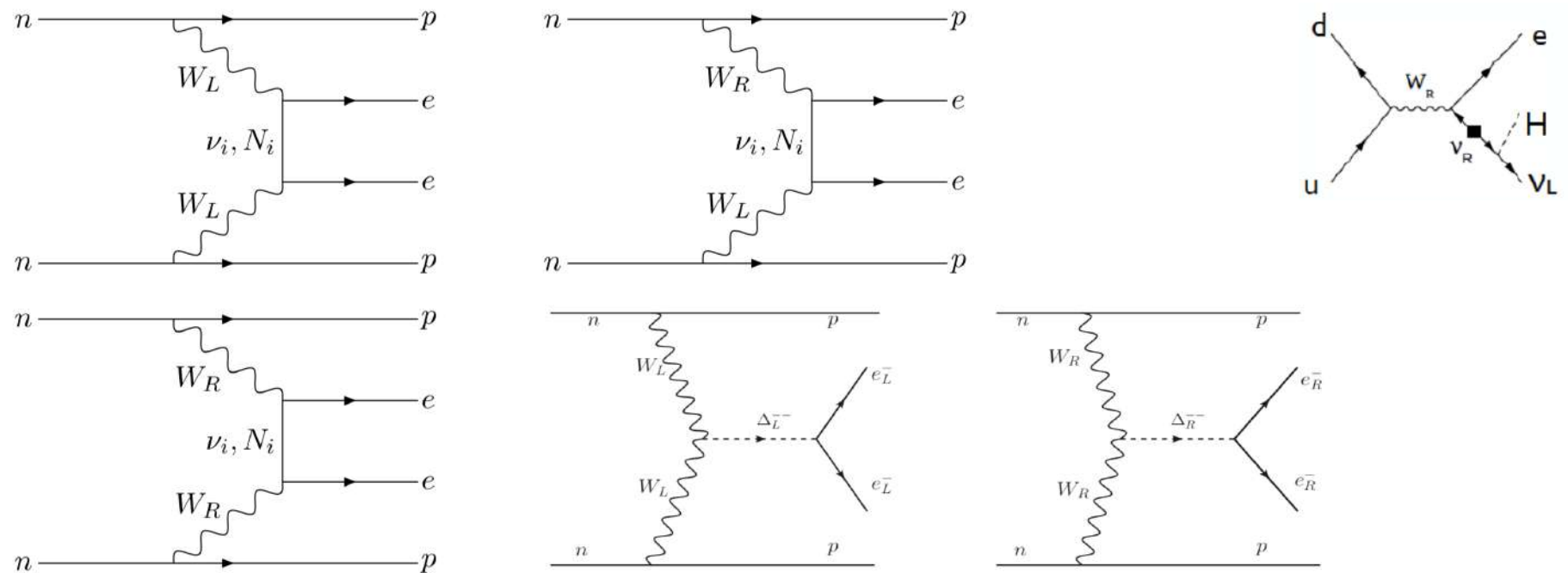
Type-I and III: **SU(2) single and triplet**

j-basis	Model
$\mathcal{O}_{HH \rightarrow LL}^{(0,1)} = \mathcal{O}^A$	N/A
$\mathcal{O}_{HH \rightarrow LL}^{(0,3)} = \mathcal{O}^S$	type II

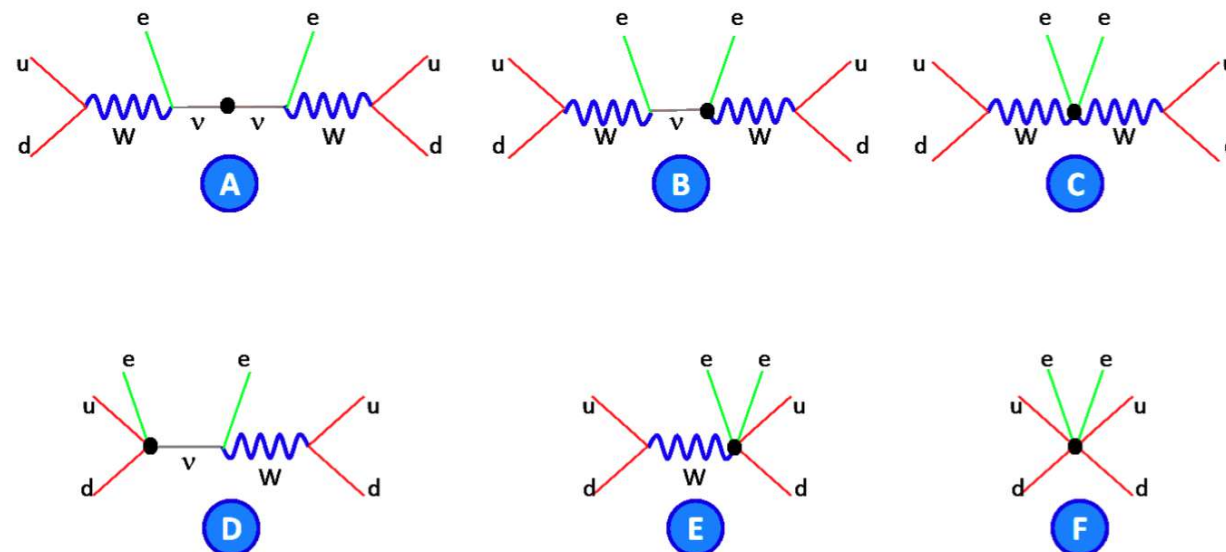
j-basis	Model
$\mathcal{O}_{HL \rightarrow HL}^{(1/2,1)} = \mathcal{O}^S + \mathcal{O}^A$	type I
$\mathcal{O}_{HL \rightarrow HL}^{(1/2,3)} = \mathcal{O}^S - 3\mathcal{O}^A$	type III

$$\mathcal{O}^S = (HL)(HL), \quad \mathcal{O}^A = (HH)(LL)$$

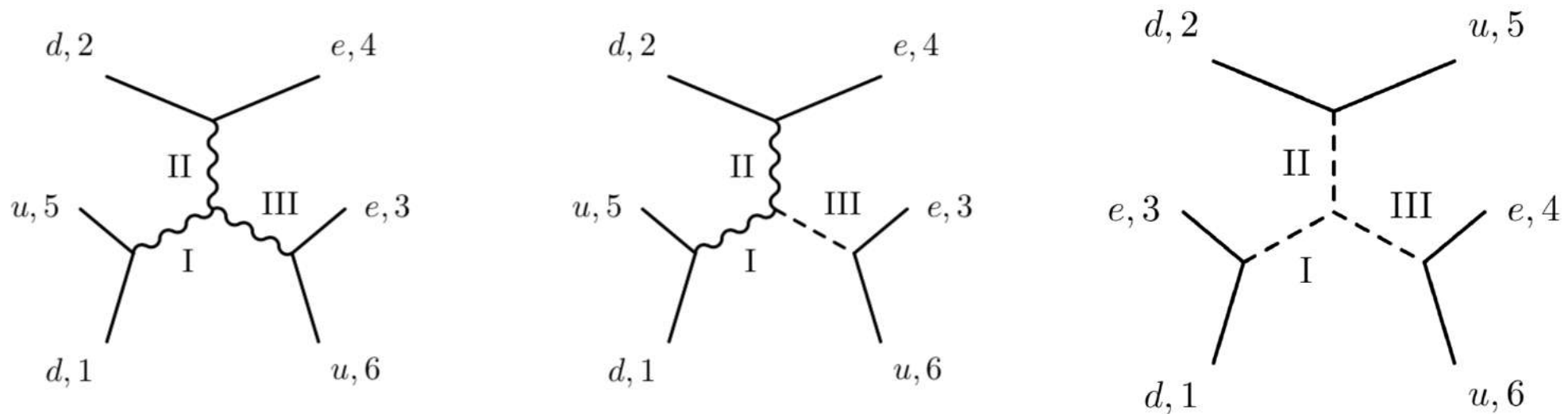
UV Realization of $0\nu b\bar{b}$



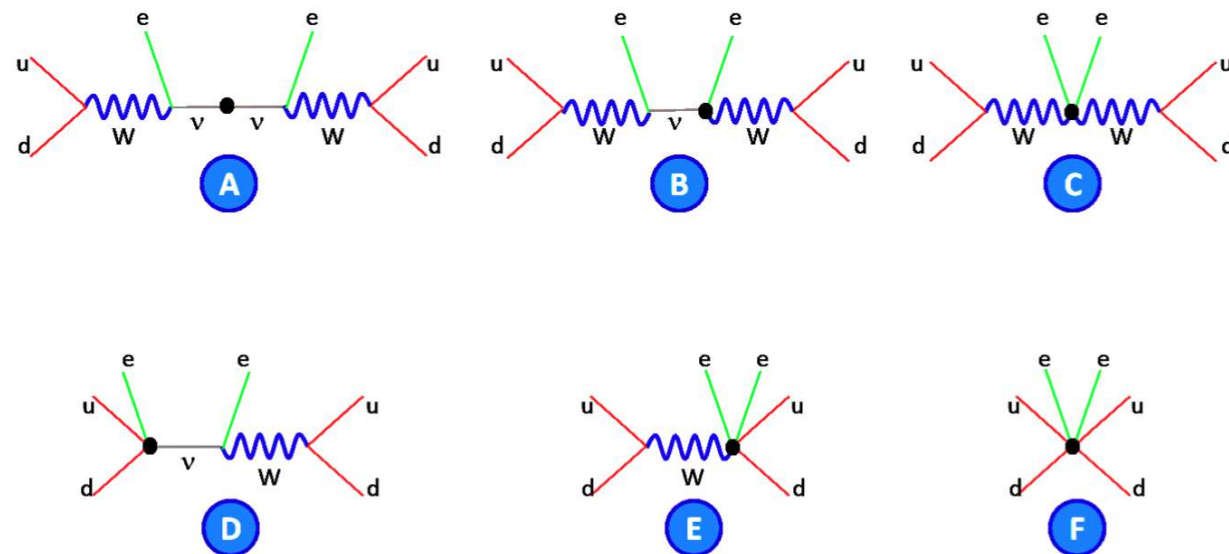
Integrate out heavy particles



More UV Realizations



$$W^2 \mathcal{B}^J = -sJ(J+1)\mathcal{B}^J$$

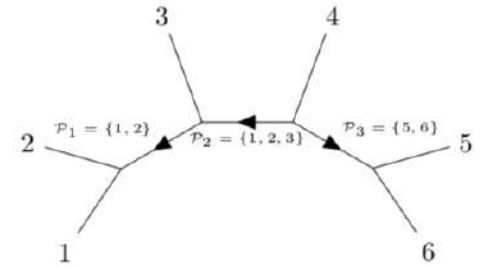
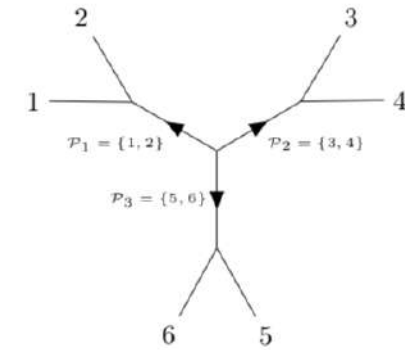


Example

Dim-9 operators:

$$T_{SU(3)}^{abcdef} \psi^{\dagger 6}$$

$$T_{SU(2)}^{ijkl} \psi^4 \psi^{\dagger 2}$$

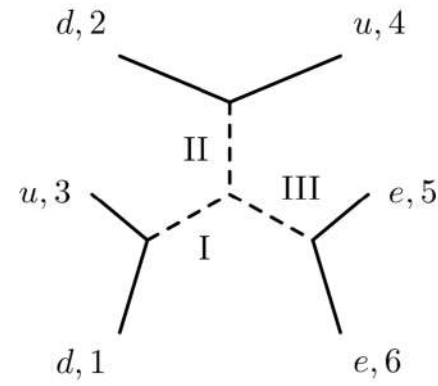
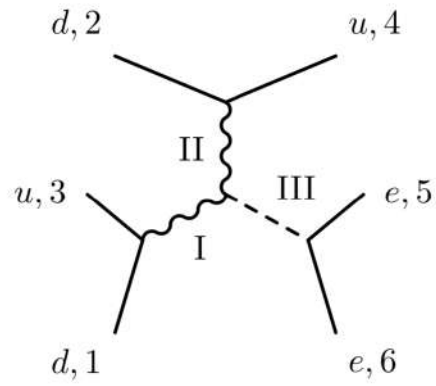


Lorentz	y-basis
ψ^6	$\mathcal{B}_1 = \langle 12 \rangle \langle 34 \rangle \langle 56 \rangle$ $\mathcal{B}_2 = \langle 12 \rangle \langle 35 \rangle \langle 46 \rangle$ $\mathcal{B}_3 = \langle 13 \rangle \langle 24 \rangle \langle 56 \rangle$ $\mathcal{B}_4 = \langle 13 \rangle \langle 25 \rangle \langle 46 \rangle$ $\mathcal{B}_5 = \langle 14 \rangle \langle 25 \rangle \langle 36 \rangle$
$\psi^4 \psi^{\dagger 2}$	$\mathcal{B}_1 = \langle 12 \rangle \langle 34 \rangle [56]$ $\mathcal{B}_2 = \langle 13 \rangle \langle 24 \rangle [56]$

gauge classes	y-basis
$T_{SU(3)}^{abcdef}$	$T_1 = \epsilon^{ace} \epsilon^{bdf}$ $T_2 = \epsilon^{acd} \epsilon^{bef}$ $T_3 = \epsilon^{abe} \epsilon^{cdf}$ $T_4 = \epsilon^{abd} \epsilon^{cef}$ $T_5 = \epsilon^{abc} \epsilon^{def}$
$T_{SU(2)}^{ijkl}$	$T'_1 = \epsilon^{ij} \epsilon^{kl}$ $T'_2 = \epsilon^{ik} \epsilon^{jl}$
$T_{SU(2)}^{ij}$	$T'_1 = \epsilon^{ij}$

type	$\bigoplus_{[\lambda]} n_{[\lambda]} \{[\lambda_1], [\lambda_2], \dots\}$
$d_c^{\dagger 4} u_c^{\dagger 2}$	$2\{\square_u, \square_d\} \oplus \{\square_u, \square_d\} \oplus$ $2\{\square_u, \square_d\} \oplus 2\{\square_u, \square_d\} \oplus$ $\{\square_u, \square_d\} \oplus 2\{\square_u, \square_d\} \oplus \{\square_u, \square_d\}$
$Q^4 d_c^{\dagger 2}$	$\{\square_Q, \square_{d^\dagger}\} \oplus 3\{\square_Q, \square_{d^\dagger}\} \oplus$ $2\{\square_Q, \square_{d^\dagger}\} \oplus 2\{\square_Q, \square_{d^\dagger}\} \oplus \{\square_Q, \square_{d^\dagger}\}$
$Q^2 d_c^{\dagger 3} u_c^{\dagger}$	$\{\square_Q, \square_{d^\dagger}\} \oplus 2\{\square_Q, \square_{d^\dagger}\} \oplus$ $\{\square_Q, \square_{d^\dagger}\} \oplus \{\square_Q, \square_{d^\dagger}\} \oplus$ $\{\square_Q, \square_{d^\dagger}\} \oplus \{\square_Q, \square_{d^\dagger}\}$

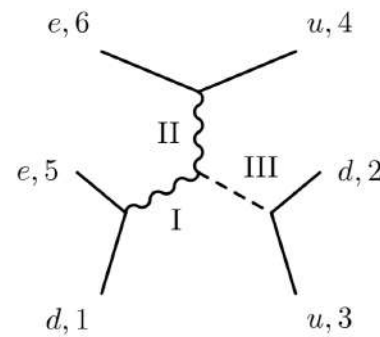
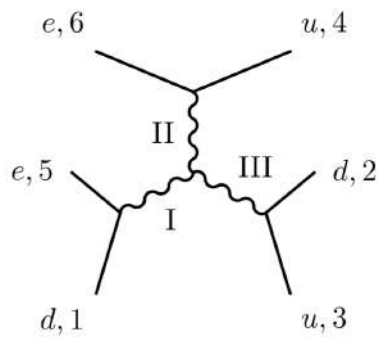
Example



$$\mathcal{O}_1 = \frac{1}{4}(L^\dagger_u{}^j L^\dagger_v{}^i)(Q_{pai}Q_{rbj})(u_{cs}{}^b u_{ct}{}^a)$$

$$\mathcal{O}_2 = -\frac{1}{4}(L^\dagger_u{}^j L^\dagger_v{}^i)(Q_{pai}u_{cs}{}^b)(Q_{rbj}u_{ct}{}^a).$$

(\mathbf{r}_i, J_i)	$(1, 1, 0)$	$(0, 0, 0)$
$(\mathbf{8}_2, \mathbf{8}_2, \mathbf{1}_3)$	$\frac{4}{3}\mathcal{O}_1 + 2\mathcal{O}_2$	$-\frac{2}{3}\mathcal{O}_2$
$(\mathbf{8}_2, \mathbf{8}_2, \mathbf{1}_1)$	$-2\mathcal{O}_2$	$\frac{4}{3}\mathcal{O}_1 + \frac{2}{3}\mathcal{O}_2$
$(\mathbf{1}_2, \mathbf{1}_2, \mathbf{1}_3)$	$-4\mathcal{O}_1 - 6\mathcal{O}_2$	$2\mathcal{O}_2$
$(\mathbf{1}_2, \mathbf{1}_2, \mathbf{1}_1)$	$6\mathcal{O}_2$	$-4\mathcal{O}_1 - 2\mathcal{O}_2$

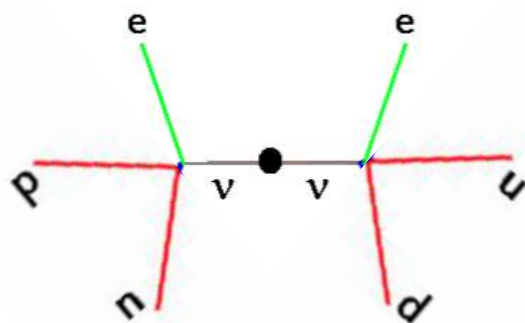
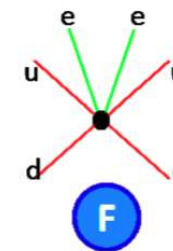
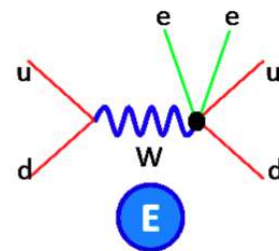
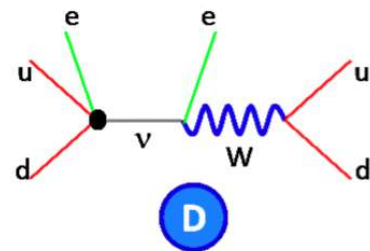
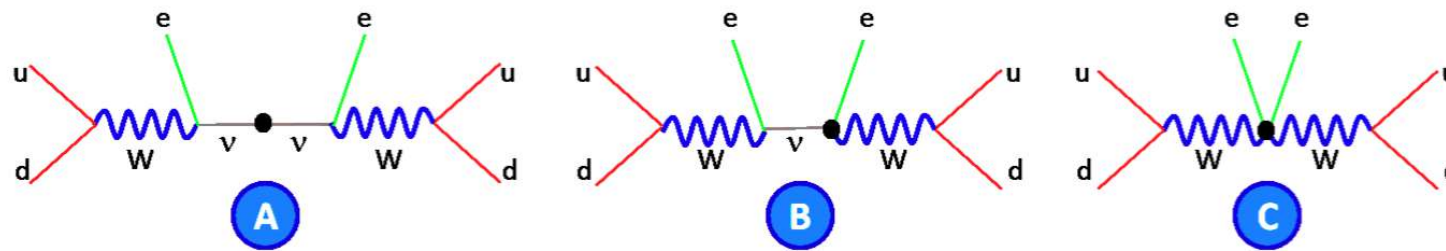


(\mathbf{r}_i, J_i)	$(1, 1, 1)$	$(1, 1, 0)$
$(\mathbf{3}_3, \mathbf{8}_2, \bar{\mathbf{3}}_2)$	$-\frac{8}{3}\mathcal{O}_1$	$\frac{4}{3}\mathcal{O}_1 + \frac{8}{3}\mathcal{O}_2$
$(\mathbf{3}_1, \mathbf{8}_2, \bar{\mathbf{3}}_2)$	$-\frac{8}{3}\mathcal{O}_1$	0
$(\mathbf{3}_3, \mathbf{1}_2, \bar{\mathbf{3}}_2)$	$-\frac{8}{3}\mathcal{O}_1$	$\frac{4}{3}\mathcal{O}_1 + \frac{8}{3}\mathcal{O}_2$
$(\mathbf{3}_1, \mathbf{1}_2, \bar{\mathbf{3}}_2)$	$-\frac{8}{3}\mathcal{O}_1$	0

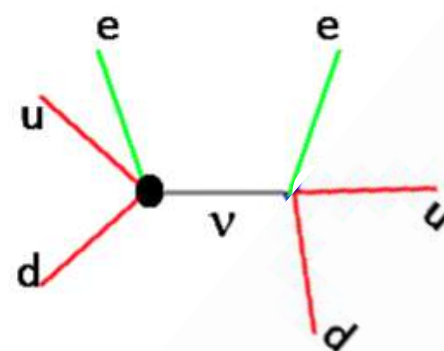
LEFT

Nucleon currents and weak sources

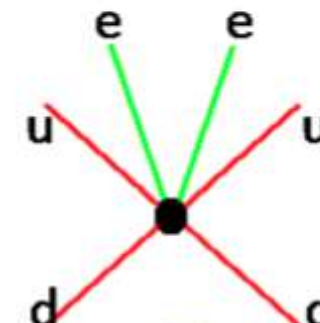
LEFT Related Operators



Dim-3, 5

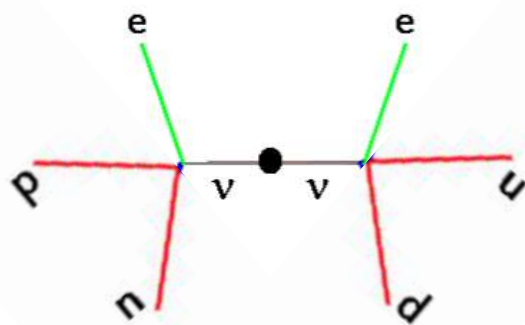
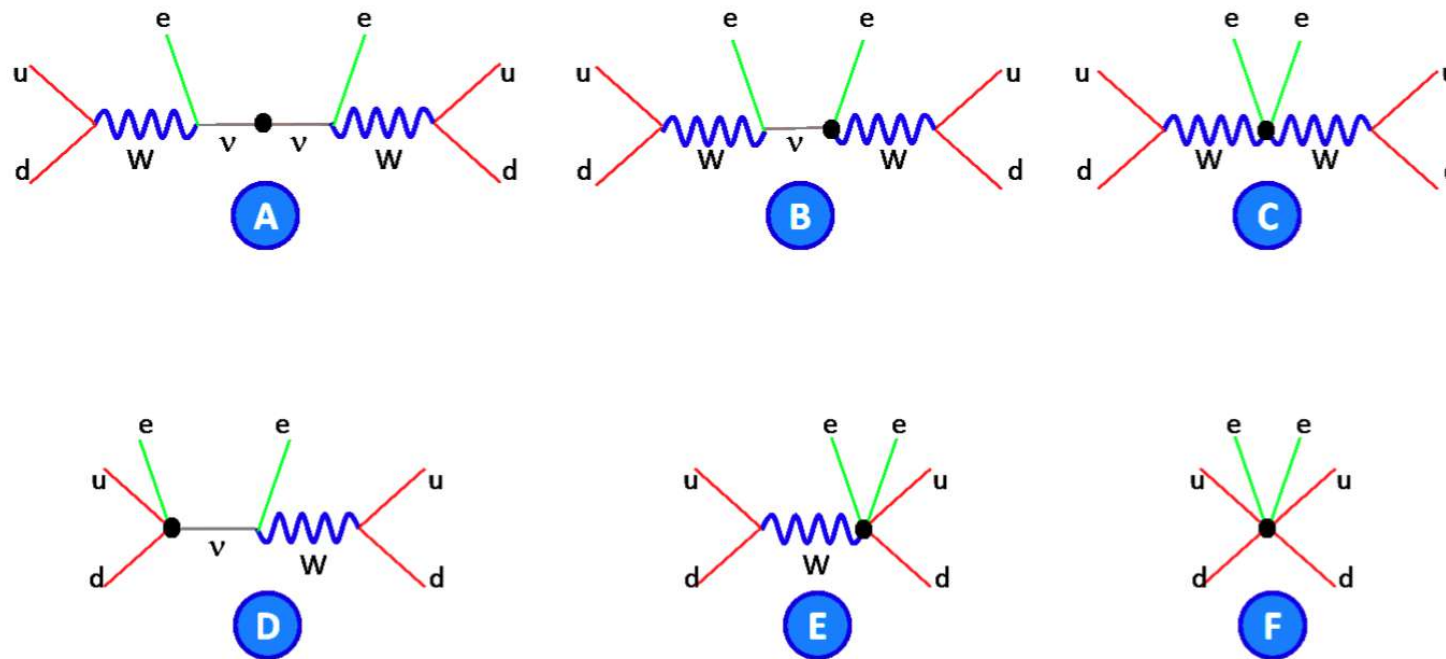


Dim-6, 7, 8, 9

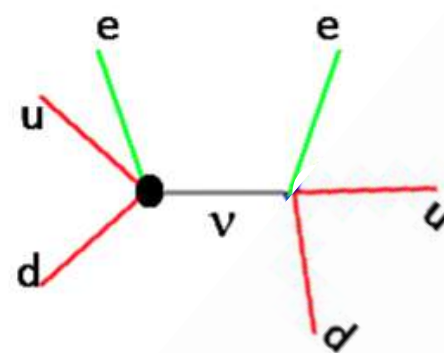


Dim-9, 11

LEFT Related Operators



Long-range interaction

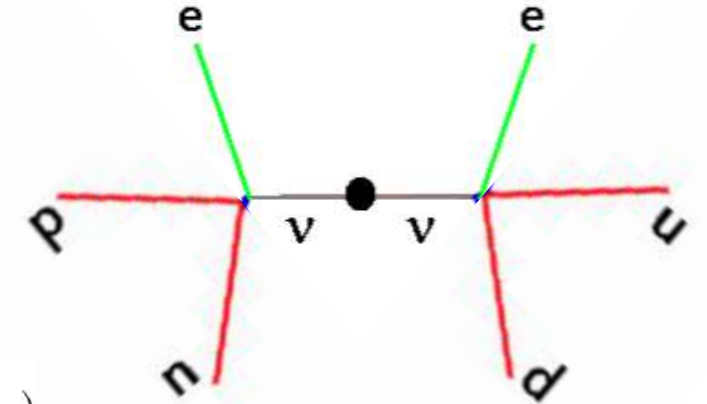


Short-range interaction

Long-Range Interaction

Standard mechanism: long-range neutrino potential

$$\mathcal{L} = \frac{G_F \cos \theta_C}{\sqrt{2}} [\bar{u} \gamma^\mu (1 - \gamma_5) d] \sum_{i=1}^3 U_{ei} [\bar{e} \gamma^\mu (1 - \gamma_5) \nu_i] + \text{h.c.}.$$



$$L^{\mu\nu} = - \int \int dx_2 dx_1 \sum_i \bar{e}(x_1) \gamma^\mu (1 - \gamma_5) U_{ei} \underbrace{\nu_{iL}(x_1) \bar{\nu}_{iL}^c(x_2)} \gamma_\nu (1 + \gamma_5) U_{ei} e_L^c(x_2)$$

$$\begin{aligned} \frac{m_i}{q^2 - m_i^2} &\propto \frac{m_i}{q^2} \quad \text{if } m_i^2 \ll q^2 \\ &\propto -\frac{1}{m_i} \quad \text{if } m_i^2 \gg q^2. \end{aligned}$$

$$m_{\beta\beta} = \sum_{i=1}^3 U_{ei}^2 m_{\nu_i}.$$

$$m_{\beta\beta} \rightarrow m_{\beta\beta} + \sum_{i=1}^{n_N} V_{eN_i}^2 m_{N_i}, \quad (m_{N_i} \ll 100 \text{ MeV}).$$

$$J_{\mu\nu}^{fi} = \sum_n \langle f | J_{\mu L}(\vec{x}_1) | n \rangle \langle n | J_{\nu L}(\vec{x}_2) | i \rangle e^{-i(E_n - E_f)x_{10}} e^{-i(E_n - E_i)x_{20}} + (\mu \rightarrow \nu, x_{10} \rightarrow x_{20}).$$

completeness relation

$$J_{0L}(\vec{x}) \simeq \sum_i \delta(\vec{x} - \vec{x}_i) f_1(0) \tau_i^+$$

$$\vec{J}_L(\vec{x}) \simeq \sum_i \delta(\vec{x} - \vec{x}_i) g_1(0) \vec{\sigma}_i \tau_i^+$$

$$J_{\mu L}(\vec{x}_1) J_{\nu L}(\vec{x}_2) = \sum_{i,j} \tau_i^+ \tau_j^+ \delta(\vec{x}_1 - \vec{x}_i) \delta(\vec{x}_2 - \vec{x}_j) [f_1^2(0) - g_1^2(0) \vec{\sigma}_i \cdot \vec{\sigma}_j].$$

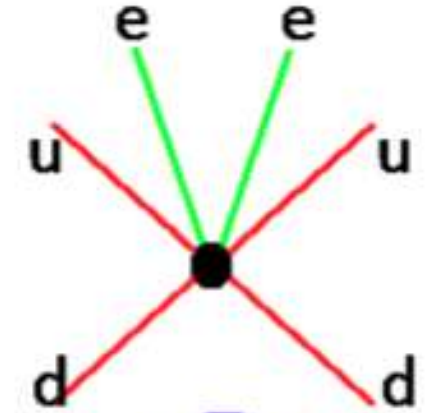
Short-Range Interaction

General quark currents = dim-9 LEFT operators

$$\mathcal{L}_{\text{SR}} = \frac{G_F^2}{2m_p} \sum_{\text{chiralities}} [\epsilon_1^\bullet J_\circ J_\circ j_\circ + \epsilon_2^\bullet J_\circ^{\mu\nu} J_{\circ\mu\nu} j_\circ + \epsilon_3^\bullet J_\circ^\mu J_{\circ\mu} j_\circ + \epsilon_4^\bullet J_\circ^\mu J_{\circ\mu\nu} j^\nu + \epsilon_5^\bullet J_\circ^\mu J_\circ j_\mu]$$

$$J_{R,L} = \bar{u}_a(1 \pm \gamma_5)d_a, \quad J_{R,L}^\mu = \bar{u}_a\gamma^\mu(1 \pm \gamma_5)d_a, \quad J_{R,L}^{\mu\nu} = \bar{u}_a\sigma_{\mu\nu}(1 \pm \gamma_5)d_a,$$

$$j_{R,L} = \bar{e}(1 \mp \gamma_5)e^c, \quad j^\mu = \bar{e}\gamma^\mu\gamma_5e^c.$$



Complete dim-9 LEFT 6-fermion operator basis

[No need to know dim-9 SMEFT 6-fermion operators]



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A neutrinoless double beta decay master formula from effective field theory

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Journal of High Energy Physics **2018**, Article number: 97 (2018) | [Cite this article](#)

\mathcal{O}_1^{RRR}	$[\bar{u}^i(1+\gamma_5)d_i] [\bar{u}^j(1+\gamma_5)d_j] [\bar{e}(1+\gamma_5)e^c]$
\mathcal{O}_1^{RRL}	$[\bar{u}^i(1+\gamma_5)d_i] [\bar{u}^j(1+\gamma_5)d_j] [\bar{e}(1-\gamma_5)e^c]$
$\mathcal{O}_1^{LRR} \equiv \mathcal{O}_1^{RLR}$	$[\bar{u}^i(1-\gamma_5)d_i] [\bar{u}^j(1+\gamma_5)d_j] [\bar{e}(1+\gamma_5)e^c]$
$\mathcal{O}_1^{LRL} \equiv \mathcal{O}_1^{RLL}$	$[\bar{u}^i(1-\gamma_5)d_i] [\bar{u}^j(1+\gamma_5)d_j] [\bar{e}(1-\gamma_5)e^c]$
\mathcal{O}_1^{LLR}	$[\bar{u}^i(1-\gamma_5)d_i] [\bar{u}^j(1-\gamma_5)d_j] [\bar{e}(1+\gamma_5)e^c]$
\mathcal{O}_1^{LLL}	$[\bar{u}^i(1-\gamma_5)d_i] [\bar{u}^j(1-\gamma_5)d_j] [\bar{e}(1-\gamma_5)e^c]$
\mathcal{O}_2^{RRR}	$[\bar{u}^i\sigma^{\mu\nu}(1+\gamma_5)d_i] [\bar{u}^j\sigma_{\mu\nu}(1+\gamma_5)d_j] [\bar{e}(1+\gamma_5)e^c]$
\mathcal{O}_2^{RRL}	$[\bar{u}^i\sigma^{\mu\nu}(1+\gamma_5)d_i] [\bar{u}^j\sigma_{\mu\nu}(1+\gamma_5)d_j] [\bar{e}(1-\gamma_5)e^c]$
$\mathcal{O}_2^{LRR} \equiv \mathcal{O}_2^{RLR}$	$[\bar{u}^i\sigma^{\mu\nu}(1-\gamma_5)d_i] [\bar{u}^j\sigma_{\mu\nu}(1-\gamma_5)d_j] [\bar{e}(1+\gamma_5)e^c]$
$\mathcal{O}_2^{LRL} \equiv \mathcal{O}_2^{RLL}$	$[\bar{u}^i\sigma^{\mu\nu}(1-\gamma_5)d_i] [\bar{u}^j\sigma_{\mu\nu}(1-\gamma_5)d_j] [\bar{e}(1-\gamma_5)e^c]$
\mathcal{O}_3^{RRR}	$[\bar{u}^i\gamma^\mu(1+\gamma_5)d_i] [\bar{u}^j\gamma_\mu(1+\gamma_5)d_j] [\bar{e}(1+\gamma_5)e^c]$
\mathcal{O}_3^{RRL}	$[\bar{u}^i\gamma^\mu(1+\gamma_5)d_i] [\bar{u}^j\gamma_\mu(1+\gamma_5)d_j] [\bar{e}(1-\gamma_5)e^c]$
$\mathcal{O}_3^{LRR} \equiv \mathcal{O}_3^{RLR}$	$[\bar{u}^i\gamma^\mu(1-\gamma_5)d_i] [\bar{u}^j\gamma_\mu(1+\gamma_5)d_j] [\bar{e}(1+\gamma_5)e^c]$
$\mathcal{O}_3^{LRL} \equiv \mathcal{O}_3^{RLL}$	$[\bar{u}^i\gamma^\mu(1-\gamma_5)d_i] [\bar{u}^j\gamma_\mu(1+\gamma_5)d_j] [\bar{e}(1-\gamma_5)e^c]$
\mathcal{O}_3^{LLR}	$[\bar{u}^i\gamma^\mu(1-\gamma_5)d_i] [\bar{u}^j\gamma_\mu(1-\gamma_5)d_j] [\bar{e}(1+\gamma_5)e^c]$
\mathcal{O}_3^{LLL}	$[\bar{u}^i\gamma^\mu(1-\gamma_5)d_i] [\bar{u}^j\gamma_\mu(1-\gamma_5)d_j] [\bar{e}(1-\gamma_5)e^c]$
\mathcal{O}_4^{RR}	$[\bar{u}^i\gamma^\mu(1+\gamma_5)d_i] [\bar{u}^j\sigma_{\mu\nu}(1+\gamma_5)d_j] [\bar{e}\gamma^\nu\gamma_5e^c]$
\mathcal{O}_4^{RL}	$[\bar{u}^i\gamma^\mu(1+\gamma_5)d_i] [\bar{u}^j\sigma_{\mu\nu}(1-\gamma_5)d_j] [\bar{e}\gamma^\nu\gamma_5e^c]$
\mathcal{O}_4^{LR}	$[\bar{u}^i\gamma^\mu(1-\gamma_5)d_i] [\bar{u}^j\sigma_{\mu\nu}(1+\gamma_5)d_j] [\bar{e}\gamma^\nu\gamma_5e^c]$
\mathcal{O}_4^{LL}	$[\bar{u}^i\gamma^\mu(1-\gamma_5)d_i] [\bar{u}^j\sigma_{\mu\nu}(1-\gamma_5)d_j] [\bar{e}\gamma^\nu\gamma_5e^c]$
\mathcal{O}_5^{RR}	$[\bar{u}^i\gamma^\mu(1+\gamma_5)d_i] [\bar{u}^j(1+\gamma_5)d_j] [\bar{e}\gamma_\mu\gamma_5e^c]$
\mathcal{O}_5^{RL}	$[\bar{u}^i\gamma^\mu(1+\gamma_5)d_i] [\bar{u}^j(1-\gamma_5)d_j] [\bar{e}\gamma_\mu\gamma_5e^c]$
\mathcal{O}_5^{LR}	$[\bar{u}^i\gamma^\mu(1-\gamma_5)d_i] [\bar{u}^j(1+\gamma_5)d_j] [\bar{e}\gamma_\mu\gamma_5e^c]$
\mathcal{O}_5^{LL}	$[\bar{u}^i\gamma^\mu(1-\gamma_5)d_i] [\bar{u}^j(1-\gamma_5)d_j] [\bar{e}\gamma_\mu\gamma_5e^c]$

Low Energy EFT

Dimension-5

Dim-5 operators			
N	(n, \tilde{n})	Classes	$\mathcal{N}_{\text{type}}$
3	(2, 0)	$F_L \psi_L^2 + h.c.$	10 + 0 + 2 + 0

10

[Jenkins, Manohar, Stoffer, 2017]

Dimension-6

Dim-6 operators

N	(n, \tilde{n})	Classes	$\mathcal{N}_{\text{type}}$	$\mathcal{N}_{\text{term}}$
3	(3, 0)	$F_L^3 + h.c.$	2 + 0 + 0 + 0	2
4	(2, 0)	$\psi_L^4 + h.c.$	14 + 12 + 8 + 2	78
	(1, 1)	$\psi_L^2 \psi_R^2$	40 + 20 + 12 + 0	84
Total		5	56 + 32 + 20 + 2	164

Dimension-7

Dim-7 operators

N	(n, \tilde{n})	Classes	$\mathcal{N}_{\text{type}}$	$\mathcal{N}_{\text{term}}$
4	(3, 0)	$F_L^2 \psi_L^2 + h.c.$	16 + 0 + 4 + 0	32
	(2, 1)	$F_L^2 \psi_L^2 \psi_R + h.c.$	16 + 0 + 4 + 0	24
		$\psi_L^3 \psi_R D + h.c.$	50 + 32 + 22 +	
Total		6	82 + 32 + 30 +	

120

166

[Liao, Ma, Wang, 2020]

Dimension-8

[Li, Ren, Xiao, Yu, Zheng, 2020]

N	(n, \tilde{n})	Subclasses	$\mathcal{N}_{\text{type}}$	$\mathcal{N}_{\text{term}}$	$\mathcal{N}_{\text{operator}}$	Equations
4	(4, 0)	$F_L^4 + h.c.$	14	26	26	(4.19)
	(3, 1)	$F_L^3 \psi_L^1 D + h.c.$	22	22	$22n_f^2$	(4.51)
		$\psi^4 D^2 + h.c.$	4+4	18+14	$12n_f^4 + n_f^2(5n_f - 1)$	(4.75, 4.78, 4.80)
		$F_L \psi^2 \phi D^2 + h.c.$	16	32	$32n_f^2$	(4.44)
		$F_L^2 \phi^2 D^2 + h.c.$	8	12	12	(4.14)
	(2, 2)	$F_L^2 F_R^2$	14	17	17	(4.19)
		$F_L F_R \psi \psi^1 D$	27	35	$35n_f^2$	(4.50, 4.51)
		$\psi^2 \psi^{12} D^2$	17+4	54+8	$\frac{1}{2}n_f^2(75n_f^2 + 11) + 6n_f^4$	(4.74, 4.79-4.81)
		$F_R \psi^2 \phi D^2 + h.c.$	16	16	$16n_f^2$	(4.44)
		$F_L F_R \phi^2 D^2$	5	6	6	(4.14)
		$\psi \psi^1 \phi^2 D^3$	7	16	$16n_f^2$	(4.31, 4.32)
		$\phi^4 D^4$	1	3	3	(4.8)
5	(3, 0)	$F_L \psi^4 + h.c.$	12+10	66+54	$42n_f^4 + 2n_f^2(9n_f + 1)$	(4.86, 4.88, 4.89, 4.91)
		$F_L^2 \psi^2 \phi + h.c.$	32	60	$60n_f^2$	(4.47, 4.48)
		$F_L^3 \phi^2 + h.c.$	6	6	6	(4.16)
	(2, 1)	$F_L \psi^2 \psi^{12} + h.c.$	84+24	172+32	$2n_f^2(59n_f^2 - 2) + 24n_f^4$	(4.84-4.85), (4.88-4.92)
		$F_R^2 \psi^2 \phi + h.c.$	32	36	$36n_f^2$	(4.47, 4.48)
		$\psi^3 \psi^1 \phi D + h.c.$	32+14	180+56	$n_f^2(135n_f - 1) + n_f^3(29n_f + 3)$	(4.66, 4.69-4.72)
		$F_L \psi \psi^1 \phi^2 D + h.c.$	38	92	$92n_f^2$	(4.39, 4.40)
		$\psi^2 \phi^3 D^2 + h.c.$	6	36	$36n_f^2$	(4.28)
		$F_L \phi^4 D^2 + h.c.$	4	6	6	(4.10)
6	(2, 0)	$\psi^4 \phi^2 + h.c.$	12+4	48+18	$5(5n_f^4 + n_f^2) + \frac{2}{3}(8n_f^4 + n_f^2)$	(4.55, 4.59, 4.62, 4.64)
		$F_L \psi^2 \phi^3 + h.c.$	16	22	$22n_f^2$	(4.36)
		$F_L^2 \phi^4 + h.c.$	8	10	10	(4.12)
	(1, 1)	$\psi^2 \psi^{12} \phi^2$	23+10	57+14	$n_f^2(42n_f^2 + n_f + 2) + 3n_f^3(3n_f - 1)$	(4.54, 4.55, 4.59-4.63)
		$\psi \psi^1 \phi^4 D$	7	13	$13n_f^2$	(4.24, 4.25)
		$\phi^6 D^2$	1	2	2	(4.8)
7	(1, 0)	$\psi^2 \phi^5 + h.c.$	6	6	$6n_f^2$	(4.21)
8	(0, 0)	ϕ^8	1	1	1	(4.8)
Total		48	471+70	1070+196	$993(n_f = 1), 44807(n_f = 3)$	

783

[Murphy, 2020]

Jiang-Hao Yu

Dimension-9

[Li, Ren, Xiao, Yu, Zheng, 2020]

N	(n, \tilde{n})	Classes	$\mathcal{N}_{\text{type}}$	$\mathcal{N}_{\text{term}}$	$\mathcal{N}_{\text{operator}}$	Equations
4	(3, 2)	$\psi^3 \psi^1 D^3 + h.c.$	0 + 4 + 2 + 0	10	$\frac{2}{3}n_f^2(7n_f^2 - 1)$	(5.50)(5.51)
		$\psi^2 \phi^2 D^4 + h.c.$	0 + 0 + 2 + 0	6	$3n_f(n_f + 1)$	(5.21)
5	(3, 1)	$F_L \psi^3 \psi^1 D + h.c.$	0 + 10 + 6 + 0	72	$32n_f^4$	(5.59)(5.60)
		$\psi^4 \phi D^2 + h.c.$	0 + 4 + 4 + 0	100	$40n_f^4$	(5.45-5.48)
		$F_L \psi^2 \phi^2 D^2 + h.c.$	0 + 0 + 4 + 0	34	$17n_f^2 - n_f$	(5.28)(5.29)
	(2, 2)	$F_R \psi^3 \psi^1 D + h.c.$	0 + 10 + 6 + 0	54	$4n_f^3(6n_f + 1)$	(5.59)(5.60)
		$\psi^2 \psi^{12} \phi D^2$	0 + 4 + 4 + 0	84	$n_f^3(49n_f + 1)$	(5.45-5.48)
		$F_R \psi^2 \phi^2 D^2 + h.c.$	0 + 0 + 4 + 0	20	$2n_f(5n_f - 1)$	(5.28)(5.29)
		$\psi \psi^1 \phi^3 D^3$	0 + 0 + 2 + 0	6	$6n_f^2$	(5.19)
6	(3, 0)	$\psi^6 + h.c.$	2 + 4 + 6 + 0	116	$\frac{1}{24}n_f^2(415n_f^4 + 53n_f^3 + 59n_f^2 + 139n_f + 6)$	(5.63-5.70)
		$F_L \psi^4 \phi + h.c.$	0 + 12 + 10 + 0	102	$2n_f^3(21n_f + 1)$	(5.54-5.56)
		$F_L^2 \psi^2 \phi^2 + h.c.$	0 + 0 + 8 + 0	20	$2n_f(5n_f + 2)$	(5.32)
	(2, 1)	$\psi^4 \psi^{12} + h.c.$	4 + 26 + 20 + 4	244	$\frac{1}{6}n_f^3(382n_f^3 - 9n_f^2 + 2n_f + 21)$	(5.63-5.69)
		$F_L \psi^2 \psi^{12} \phi + h.c.$	0 + 24 + 24 + 0	92	$52n_f^4$	(5.54-5.56)
		$F_L^2 \psi^2 \phi^2 + h.c.$	0 + 0 + 8 + 0	12	$2n_f(3n_f + 2)$	(5.32)
		$\psi^3 \psi^1 \phi^2 D + h.c.$	0 + 12 + 18 + 0	186	$\frac{2}{3}n_f^2(146n_f^2 + 1)$	(5.39-5.42)
		$F_L \psi \psi^1 \phi^2 D + h.c.$	0 + 0 + 8 + 0	12	$12n_f^2$	(5.25)
		$\psi^2 \phi^4 D^2 + h.c.$	0 + 0 + 4 + 0	24	$2n_f(6n_f + 1)$	(5.17)
7	(2, 0)	$\psi^4 \phi^3 + h.c.$	0 + 6 + 6 + 0	32	$\frac{4}{3}n_f^2(10n_f^2 - 1)$	(5.35-5.37)
		$F_L \psi^2 \phi^3 + h.c.$	0 + 0 + 4 + 0	8	$2n_f(2n_f - 1)$	(5.23)
	(1, 1)	$\psi^2 \psi^{12} \phi^3$	0 + 6 + 10 + 0	24	$14n_f^4$	(5.35-5.37)
		$\psi \psi^1 \phi^5 D$	0 + 0 + 2 + 0	2	$2n_f^2$	(5.12)
8	(1, 0)	$\psi^2 \phi^6 + h.c.$	0 + 0 + 2 + 0	2	$n_f^2 + n_f$	(5.9)
Total		42	6+122+164+4	1262	$8 + 204 + 348 + 0 (n_f = 1)$ $2862 + 42234 + 44874 + 486 (n_f = 3)$	

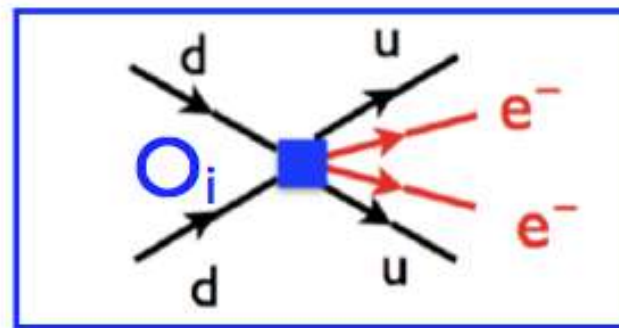
3774

40

ChiPT: Quark to Nucleon

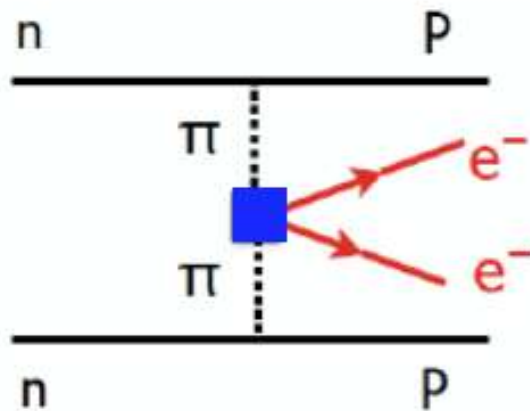
Chiral perturbation theory + Heavy baryon EFT + LNV external source

Pion-range
effects

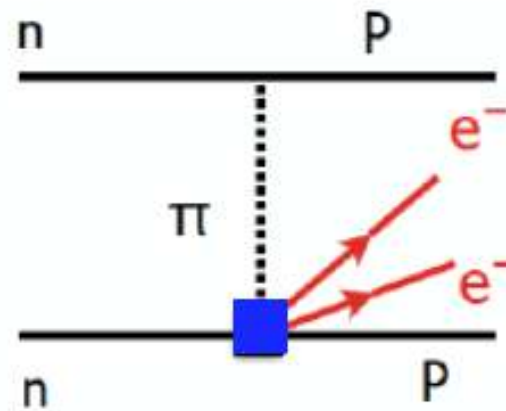


[Cirigliano, Dekens, de Vries, Graesser, 2018]

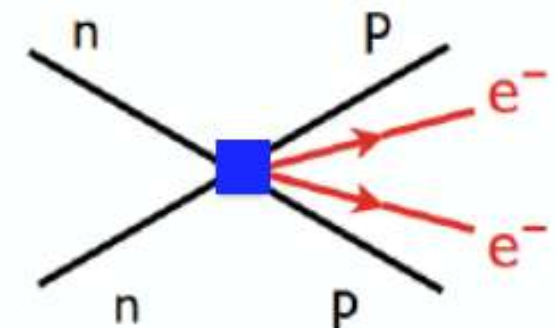
Short-range
effects



$\mathcal{L}_{\pi\pi}$



$\mathcal{L}_{\pi N}$



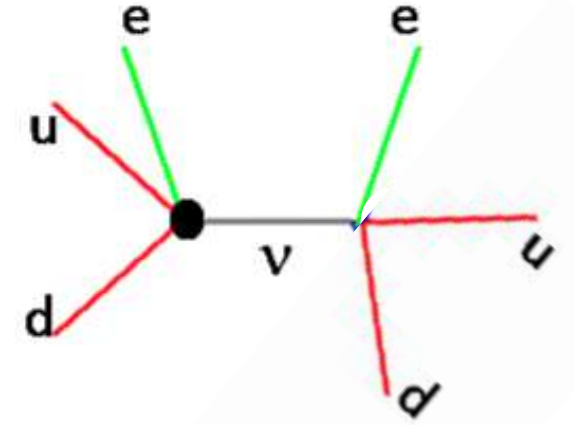
\mathcal{L}_{NN}

Long-Range from LNV Operators

Long-range neutrino potential: no ν mass dependence

$$\mathcal{L}^{4\text{-Fermi}} = \mathcal{L}^{\text{SM}} + \mathcal{L}^{\text{LNV}} = \frac{G_F}{\sqrt{2}} \left[j_{V-A}^\mu J_{V-A,\mu} + \sum_{\alpha, \beta \neq V-A} \epsilon_\alpha^\beta j_\beta J_\alpha \right]$$

$$\begin{aligned} J_{V\pm A}^\mu &= (J_{R/L})^\mu \equiv \bar{u}\gamma^\mu(1 \pm \gamma_5)d, & j_{V\pm A}^\mu &\equiv \bar{e}\gamma^\mu(1 \pm \gamma_5)\nu, \\ J_{S\pm P} &= J_{R/L} \equiv \bar{u}(1 \pm \gamma_5)d, & j_{S\pm P} &\equiv \bar{e}(1 \pm \gamma_5)\nu, \\ J_{T_{R/L}}^{\mu\nu} &= (J_{R/L})^{\mu\nu} \equiv \bar{u}\gamma^{\mu\nu}(1 \pm \gamma_5)d, & j_{T_{R/L}}^{\mu\nu} &\equiv \bar{e}\gamma^{\mu\nu}(1 \pm \gamma_5)\nu, \end{aligned}$$



Complete dim-6 LEFT 4-fermion operator basis

$$\begin{aligned} \mathcal{L}_{\Delta L=2}^{(6)} &= \frac{2G_F}{\sqrt{2}} \left\{ C_{\text{VL},ij}^{(6)} \bar{u}_L \gamma^\mu d_L \bar{e}_{R,i} \gamma_\mu C \bar{\nu}_{L,j}^T + C_{\text{VR},ij}^{(6)} \bar{u}_R \gamma^\mu d_R \bar{e}_{R,i} \gamma_\mu C \bar{\nu}_{L,j}^T \right. \\ &\quad \left. + C_{\text{SR},ij}^{(6)} \bar{u}_L d_R \bar{e}_{L,i} C \bar{\nu}_{L,j}^T + C_{\text{SL},ij}^{(6)} \bar{u}_R d_L \bar{e}_{L,i} C \bar{\nu}_{L,j}^T + C_{\text{T},ij}^{(6)} \bar{u}_L \sigma^{\mu\nu} d_R \bar{e}_{L,i} \sigma_{\mu\nu} C \bar{\nu}_{L,j}^T \right\} + \text{h.c.} \end{aligned} \quad (7)$$

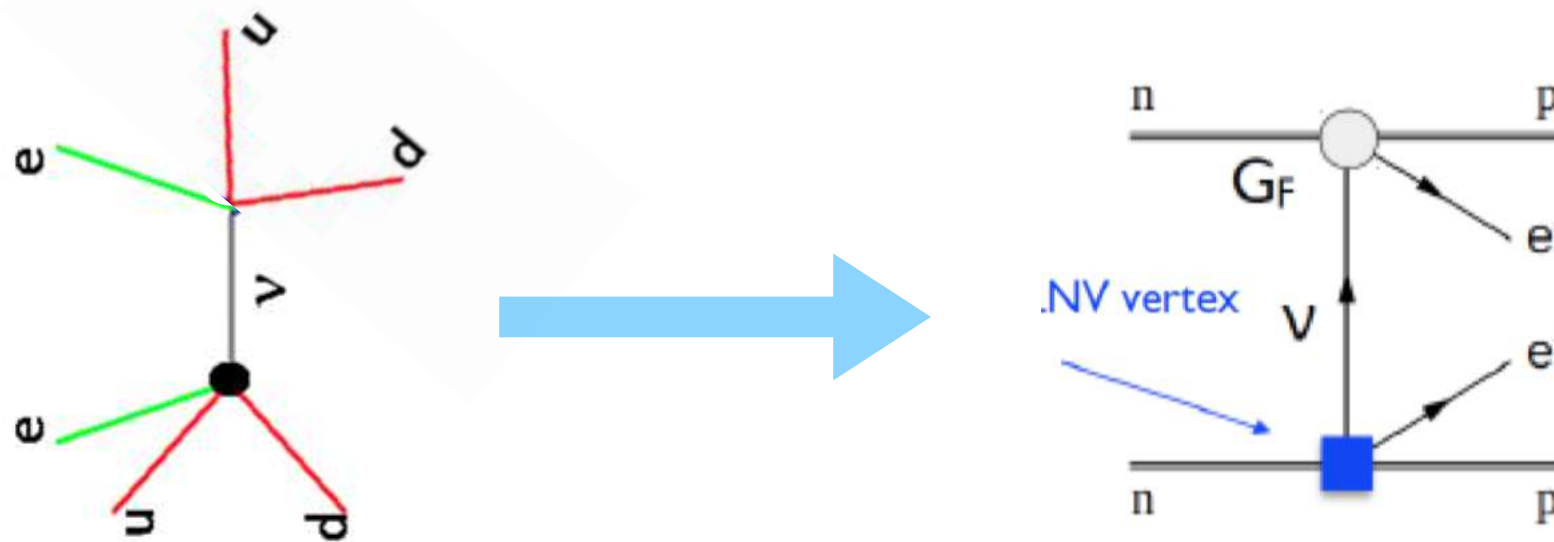
$$\mathcal{L}_{\Delta L=2}^{(7)} = \frac{2G_F}{\sqrt{2}v} \left\{ C_{\text{VL},ij}^{(7)} \bar{u}_L \gamma^\mu d_L \bar{e}_{L,i} C i \overleftrightarrow{\partial}_\mu \bar{\nu}_{L,j}^T + C_{\text{VR},ij}^{(7)} \bar{u}_R \gamma^\mu d_R \bar{e}_{L,i} C i \overleftrightarrow{\partial}_\mu \bar{\nu}_{L,j}^T \right\} + \text{h.c.} \quad (8)$$

dim-7/8/9 LEFT 4-fermion operator?

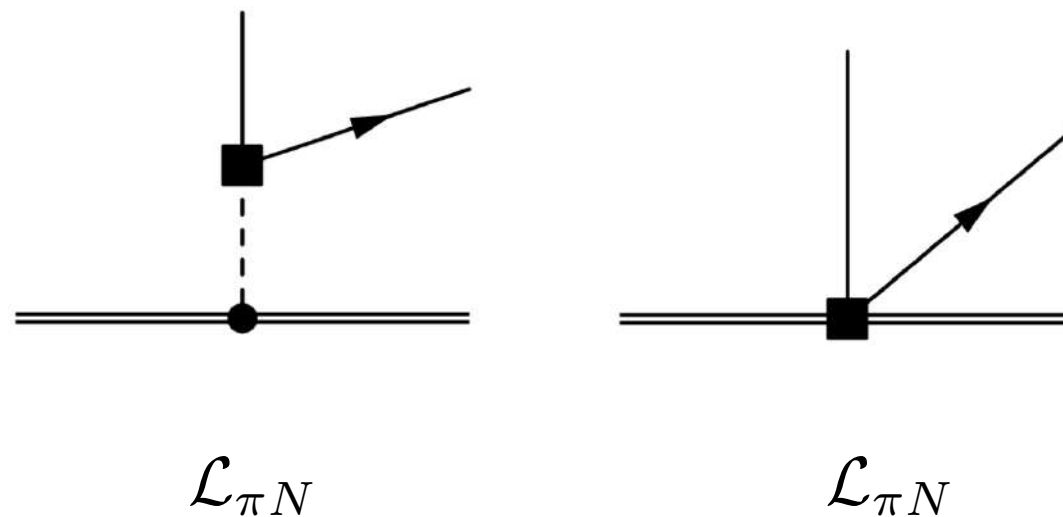
No dim-9 SMEFT 4-fermion operator!

ChiPT: Quark to Nucleon

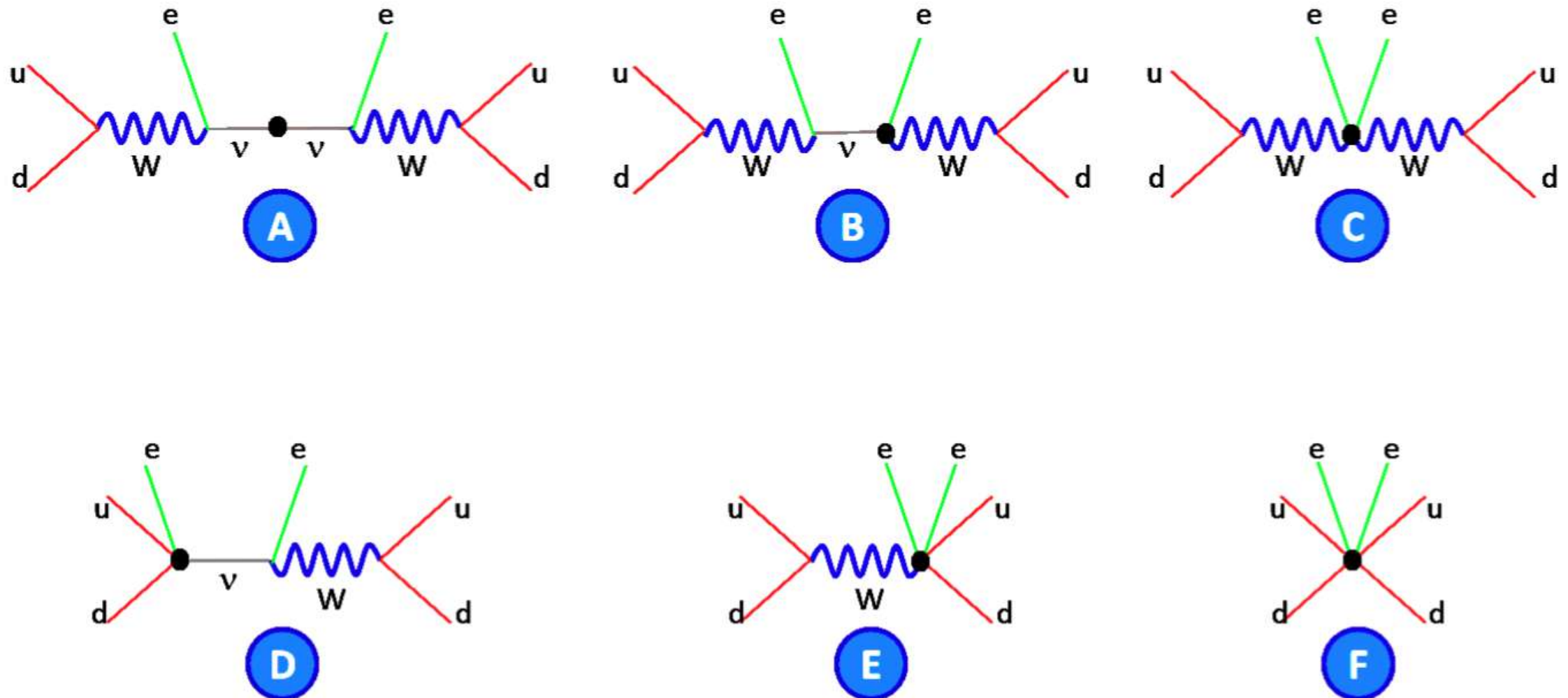
Chiral perturbation theory + Heavy baryon EFT + LNV external source



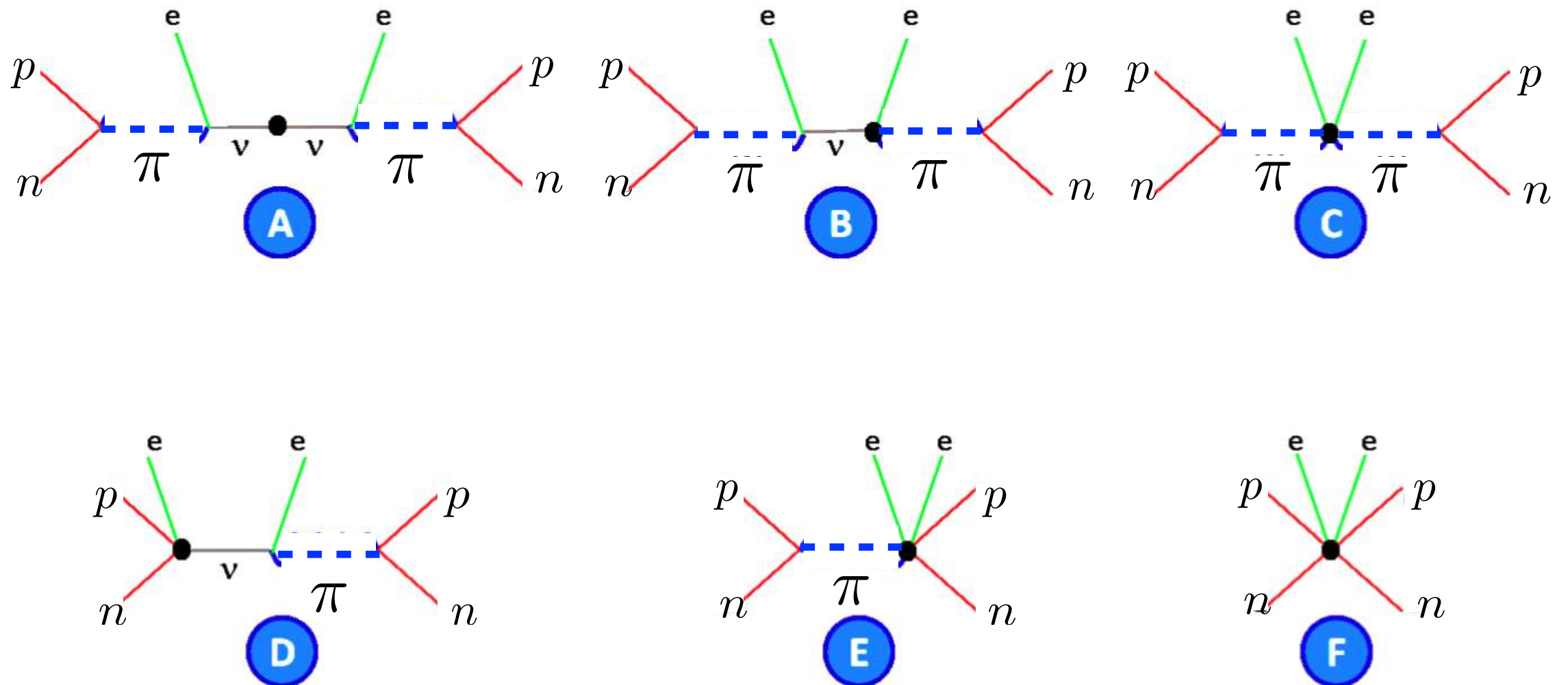
[Cirigliano, Dekens, de Vries, Graesser, 2018]



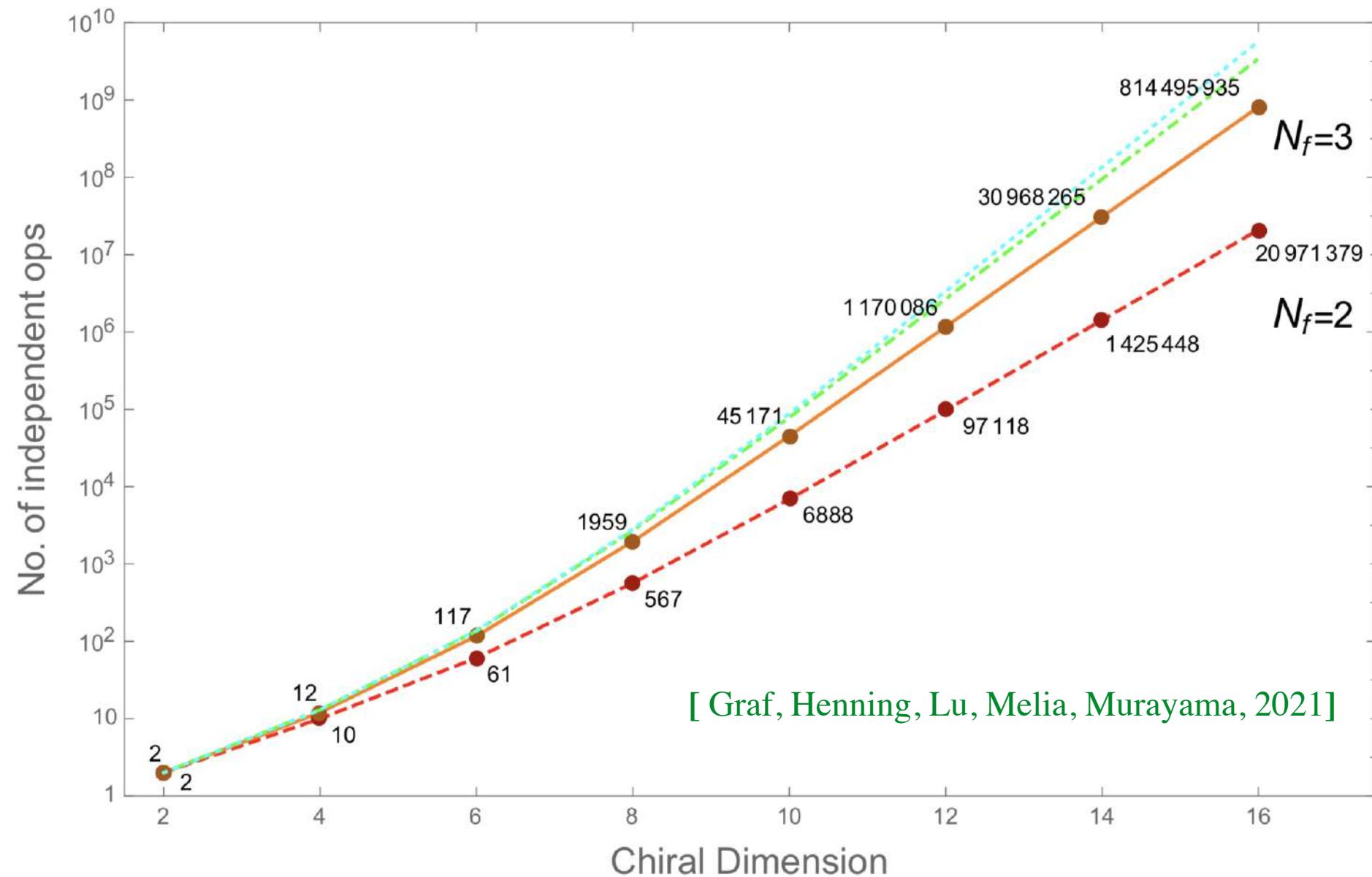
0vbb at SMEFT



0vbb at ChiPT



ChiPT Lagrangian



J. Bijnens, N. Hermansson-Truedsson, and S. Wang, 2019

Fettes, Meisner, Mojzis, Steininger, 2000

Girlanda, Pastore, Schiavilla, Viviani, 2010

Jiang-Hao Yu

Nv Potential Master Formula

$$\mathcal{A} = \langle 0^+ | \sum_{m,n} \int \frac{d^3 \mathbf{q}}{(2\pi)^3} e^{i\mathbf{q} \cdot \mathbf{r}} V(\mathbf{q}^2) | 0^+ \rangle \quad V(\mathbf{q}^2) = V_3(\mathbf{q}^2) + V_6(\mathbf{q}^2) + V_7(\mathbf{q}^2) + V_9(\mathbf{q}^2).$$

[Cirigliano, Dekens, de Vries, Graesser, 2018]

$$V_3(\mathbf{q}^2) = -(\tau^{(1)+}\tau^{(2)+})(4g_A^2 G_F^2 V_{ud}^2) m_{\beta\beta} \bar{u}(k_1) P_R C \bar{u}^T(k_2) \left\{ \frac{1}{\mathbf{q}^2} \left(-\frac{1}{g_A^2} h_F(\mathbf{q}^2) + \boldsymbol{\sigma}^{(1)} \cdot \boldsymbol{\sigma}^{(2)} h_{GT}(\mathbf{q}^2) + S^{(12)} h_T(\mathbf{q}^2) \right) + \frac{2g_\nu^{NN}}{g_A^2} h_F(\mathbf{q}^2) \right\}$$

$$V_a(\mathbf{q}^2) = \tau^{(1)+}\tau^{(2)+} 4g_A^2 G_F^2 V_{ud} \left(B \left(C_{\text{SL}}^{(6)} - C_{\text{SR}}^{(6)} \right) + \frac{m_\pi^2}{v} \left(C_{\text{VL}}^{(7)} - C_{\text{VR}}^{(7)} \right) \right) \frac{1}{\mathbf{q}^2} \bar{u}(k_1) P_R C \bar{u}^T(k_2) \left\{ \boldsymbol{\sigma}^{(1)} \cdot \boldsymbol{\sigma}^{(2)} \left(\frac{1}{2} h_{GT}^{AP}(\mathbf{q}^2) + h_{GT}^{PP}(\mathbf{q}^2) \right) + S^{(12)} \left(\frac{1}{2} h_T^{AP}(\mathbf{q}^2) + h_T^{PP}(\mathbf{q}^2) \right) \right\}. \quad (83)$$

More matrix element?

$$V_9(\mathbf{q}^2) = -(\tau^{(1)+}\tau^{(2)+}) g_A^2 \frac{4G_F^2}{v} \bar{u}(k_1) P_R C \bar{u}^T(k_2) \times \left[- \left(\boldsymbol{\sigma}^{(1)} \cdot \boldsymbol{\sigma}^{(2)} - S^{(12)} \right) \left(\frac{C_{\pi\pi\text{L}}^{(9)}}{6} \frac{\mathbf{q}^2}{(\mathbf{q}^2 + m_\pi^2)^2} - \frac{C_{\pi\text{NL}}^{(9)}}{3} \frac{\mathbf{q}^2}{\mathbf{q}^2 + m_\pi^2} \right) + \frac{2}{g_A^2} C_{\text{NNL}}^{(9)} \right]$$

Sterile Neutrino EFT

Dimension-5

Dim-5 operators			
N	(n, \tilde{n})	Classes	$\mathcal{N}_{\text{type}}$
3	(2, 0)	$F_L \psi^2 + h.c.$	0 + 0 + 2 + 0
4	(1, 0)	$\psi^2 \phi^2 + h.c.$	0 + 0 + 2 + 0
Total		4	0 + 0 + 4 + 0

2

[Aguila, Bar-Shalom, Soni, Wudka, 2009]

Dimension-6

Dim-6 operators			
N	(n, \tilde{n})	Classes	$\mathcal{N}_{\text{type}}$
4	(2, 0)	$\psi^4 + h.c.$	4 + 2 + 0 + 2
		$F_L \psi^2 \phi + h.c.$	4 + 0 + 0 + 0
	(1, 1)	$\psi^2 \psi^{\dagger 2}$	10 + 2 + 0 + 0
		$\psi \psi^{\dagger} \phi^2 D$	3 + 0 + 0 + 0
5	(1, 0)	$\psi^2 \phi^3 + h.c.$	2 + 0 + 0 + 0
Total		8	23 + 4 + 0 + 2

29

Dimension-7

Dim-7 operators			
N	(n, \tilde{n})	Classes	$\mathcal{N}_{\text{type}}$
4	(3, 0)	$F_L^2 \psi^2 + h.c.$	0 + 0 + 6 + 0
	(2, 1)	$F_L^2 \psi^{\dagger 2} + h.c.$	0 + 0 + 6 + 0
		$\psi^3 \psi^{\dagger} D + h.c.$	0 + 4 + 20 + 0
		$F_L \psi \psi^{\dagger} \phi D + h.c.$	0 + 0 + 8 + 0
5	(2, 0)	$\psi^4 \phi + h.c.$	0 + 2 + 10 + 0
		$F_L \psi^2 \phi^2 + h.c.$	0 + 0 + 6 + 0
	(1, 1)	$\psi^2 \psi^{\dagger 2} \phi$	0 + 4 + 22 + 0
		$\psi \psi^{\dagger} \phi^3 D$	0 + 0 + 2 + 0
6	(1, 0)	$\psi^2 \phi^4 + h.c.$	0 + 0 + 2 + 0
Total		18	0 + 10 + 86 + 0

80

[Bhattacharya, Wudka, 2016]

[Liao, Ma, 2017]

Dimension-8

[Li, Ren, Xiao, Yu, Zheng, 2021]

N	(n, \tilde{n})	Classes	$\mathcal{N}_{\text{type}}$
4	(3, 1)	$\psi^4 D^2 + h.c.$	4 + 0 + 2 + 2
		$F_L \psi^2 \phi D^2 + h.c.$	4 + 0 + 0 + 0
	(2, 2)	$F_L F_R \psi \psi^{\dagger} D$	3 + 0 + 0 + 0
		$\psi^2 \psi^{\dagger 2} D^2$	10 + 2 + 0 + 0
		$F_R \psi^2 \phi D^2 + h.c.$	4 + 0 + 0 + 0
		$\psi \psi^{\dagger} \phi^2 D^3$	3 + 0 + 0 + 0
5	(3, 0)	$F_L \psi^4 + h.c.$	10 + 4 + 0 + 2
		$F_L^2 \psi^2 \phi + h.c.$	8 + 0 + 0 + 0
	(2, 1)	$F_L \psi^2 \psi^{\dagger 2} + h.c.$	42 + 12 + 0 + 0
		$F_L^2 \psi^{\dagger 2} \phi + h.c.$	8 + 0 + 0 + 0
		$\psi^3 \psi^{\dagger} \phi D + h.c.$	24 + 6 + 0 + 2
		$F_L \psi \psi^{\dagger} \phi^2 D + h.c.$	12 + 0 + 0 + 0
6	(2, 0)	$\psi^4 \phi^2 + h.c.$	8 + 2 + 0 + 2
		$F_L \psi^2 \phi^3 + h.c.$	4 + 0 + 0 + 0
	(1, 1)	$\psi^2 \psi^{\dagger 2} \phi^2$	16 + 4 + 0 + 2
		$\psi \psi^{\dagger} \phi^4 D$	3 + 0 + 0 + 0
7	(1, 0)	$\psi^2 \phi^5 + h.c.$	2 + 0 + 0 + 0
Total		31	167 + 30 + 2 + 10

323

Dimension-9

[Li, Ren, Xiao, Yu, Zheng, 2021]

N	(n, \tilde{n})	Classes	$\mathcal{N}_{\text{type}}$
4	(4, 1)	$F_L^2 \psi^2 D^2 + h.c.$	0 + 6 + 0 + 0
		$F_L F_R \psi^2 D^2 + h.c.$	0 + 6 + 0 + 0
	(3, 2)	$F_L^2 \psi^{\dagger 2} D^2 + h.c.$	0 + 6 + 0 + 0
		$\psi^3 \psi^{\dagger} D^3 + h.c.$	4 + 20 + 0 + 0
		$F_L \psi \psi^{\dagger} \phi D^3 + h.c.$	0 + 8 + 0 + 0
		$\psi^2 \phi^2 D^4 + h.c.$	0 + 4 + 0 + 0
5	(4, 0)	$F_L^3 \psi^2 + h.c.$	0 + 10 + 0 + 0
		$F_L^3 \psi^{\dagger 2} + h.c.$	0 + 4 + 0 + 0
	(3, 1)	$F_L \psi^3 \psi^{\dagger} D + h.c.$	10 + 42 + 0 + 0
		$F_L^2 \psi \psi^{\dagger} \phi D + h.c.$	0 + 16 + 0 + 0
		$\psi^4 \phi D^2 + h.c.$	2 + 10 + 0 + 0
		$F_L \psi^2 \phi^2 D^2 + h.c.$	0 + 8 + 0 + 0
	(2, 2)	$F_L F_R^2 \psi^2 + h.c.$	0 + 12 + 0 + 0
		$F_R \psi^3 \psi^{\dagger} D + h.c.$	10 + 42 + 0 + 0
		$F_L F_R \psi \psi^{\dagger} \phi D$	0 + 10 + 0 + 0
		$\psi^2 \psi^{\dagger 2} \phi D^2$	4 + 22 + 0 + 0
	(1, 1)	$F_R \psi^2 \phi^2 D^2 + h.c.$	0 + 8 + 0 + 0
		$\psi \psi^{\dagger} \phi^3 D^3$	0 + 2 + 0 + 0
6	(3, 0)	$\psi^6 + h.c.$	6 + 10 + 6 + 2
		$F_L \psi^4 \phi + h.c.$	6 + 26 + 0 + 0
		$F_L^2 \psi^2 \phi^2 + h.c.$	0 + 12 + 0 + 0
	(2, 1)	$\psi^4 \psi^{\dagger 2} + h.c.$	40 + 106 + 14 + 0
		$F_L \psi^2 \psi^{\dagger 2} \phi + h.c.$	24 + 116 + 0 + 0
		$F_L^2 \psi^{\dagger 2} \phi^2 + h.c.$	0 + 10 + 0 + 0
		$\psi^3 \phi^1 \phi^2 D + h.c.$	10 + 44 + 0 + 0
	(1, 1)	$F_L \psi \psi^{\dagger} \phi^3 D + h.c.$	0 + 8 + 0 + 0
		$\psi^2 \phi^4 D^2 + h.c.$	0 + 4 + 0 + 0
7	(2, 0)	$\psi^4 \phi^3 + h.c.$	2 + 12 + 0 + 0
		$F_L \psi^2 \phi^4 + h.c.$	0 + 6 + 0 + 0
	(1, 1)	$\psi^2 \psi^{\dagger 2} \phi^3$	4 + 22 + 0 + 0
		$\psi \psi^{\dagger} \phi^5 D$	0 + 2 + 0 + 0

1358

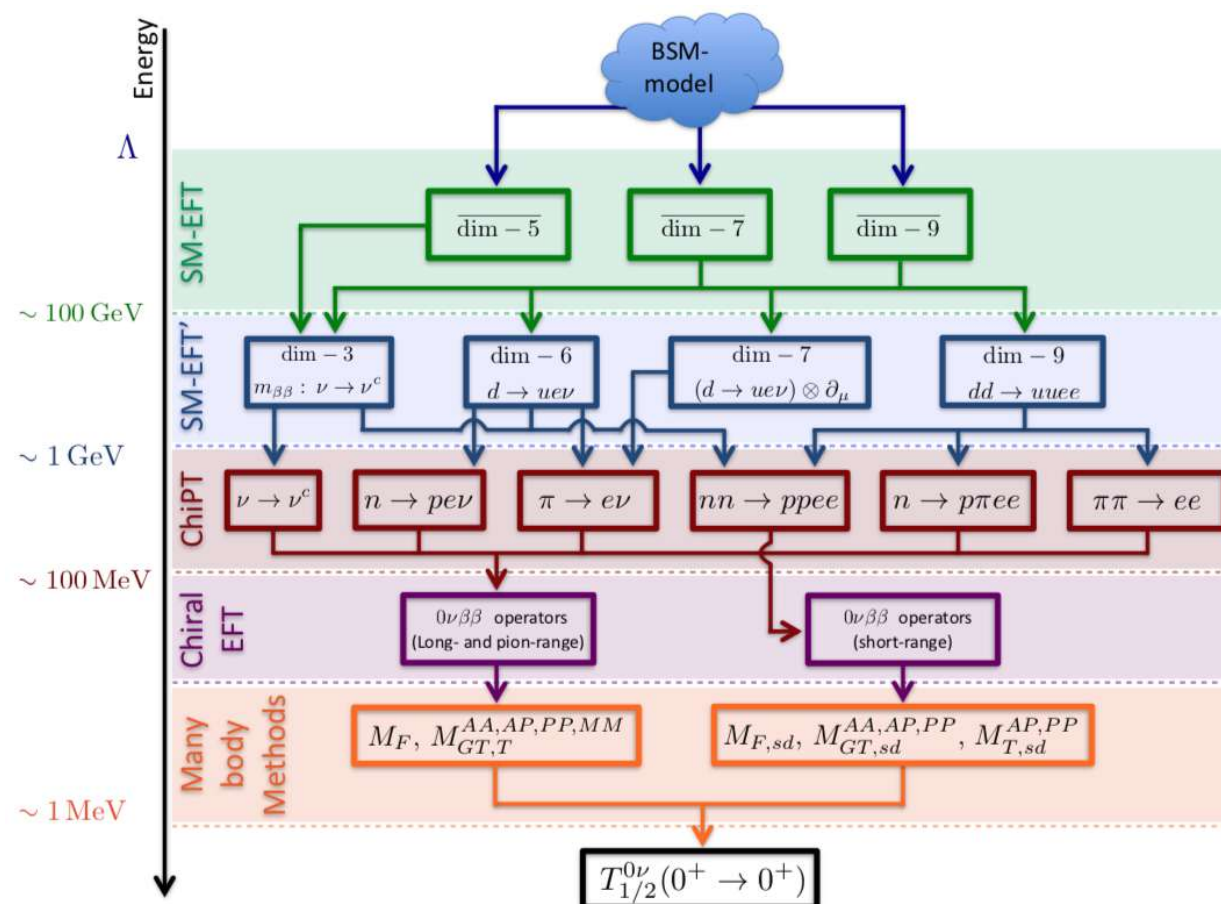
Jiang-Hao Yu

48

Summary

Summary

- $0\nu\beta\beta$ involves in many scales: SMEFT, LEFT, ChiEFT
- The complete bases just written down recently 2020 - 2021
- The formalism needs to be extended in each EFT levels



Thanks very much!