

Quantum heat engine simulated on superconducting qubits

Nick Materise and Eliot Kapit

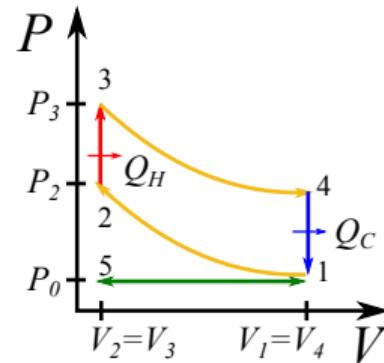
¹Colorado School of Mines, Department of Physics

March, 2 2020



Classical Otto Engine

1. Intake stroke
 2. Compression stroke (adiabatic),
 $V_1 \rightarrow V_2, P_0 \rightarrow P_2$
 3. Heating stroke (isochoric), $P_2 \rightarrow P_3$
 4. Expansion stroke (adiabatic),
 $V_2 \rightarrow V_1, P_3 \rightarrow P_2$
 5. Cooling stroke (isochoric), $P_3 \rightarrow P_0$
- » Efficiency $\eta_{\text{Otto}} = 1 - r^{1-\gamma}$,
 $r = V_1/V_2, \gamma = C_v/C_p$ ¹



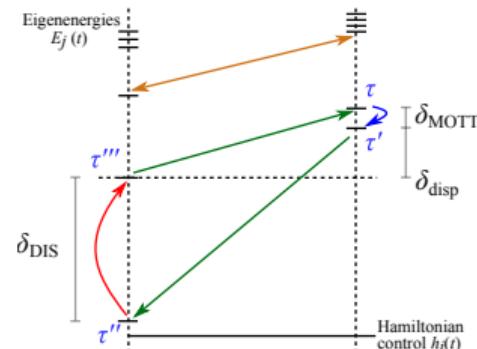
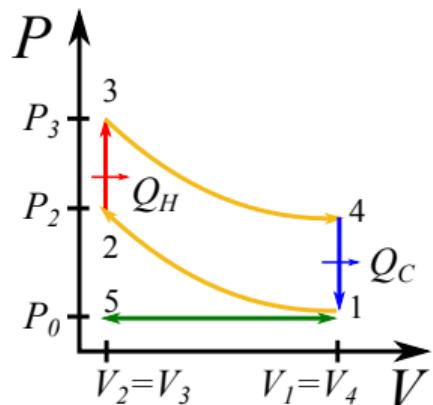
¹Quattroch 2006.

Near Ground State Disordered (NGSD) Engine

1. $\tau''' \rightarrow \tau$, Expansion
 2. $\tau \rightarrow \tau'$, Cooling
 3. $\tau' \rightarrow \tau''$, Compression
 4. $\tau'' \rightarrow \tau'''$, Heating
- » Efficiency

$$\eta_{\text{NGSD}} = 1 - \delta_{\text{MOTT}}/\delta_{\text{DIS}},$$

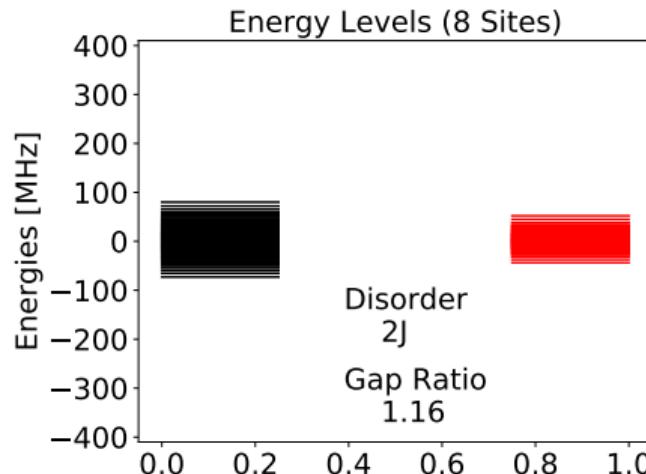
similar to²



²Yunger Halpern et al., Phys. Rev. B 99, 024203 (2019).

Near Ground State Disordered (NGSD) Engine

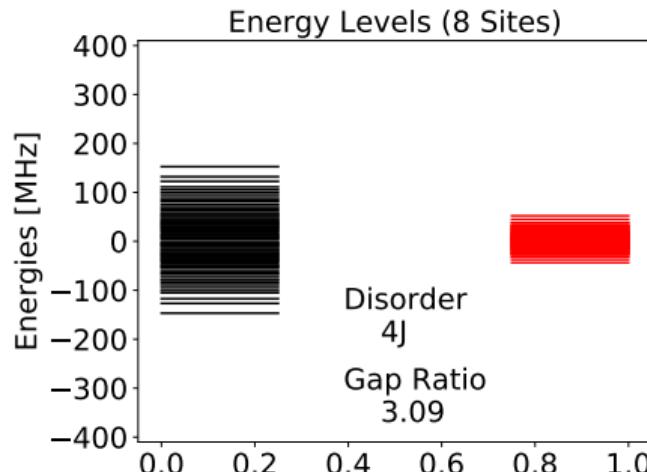
- ▶ Energy Spectrum of disordered (DIS) and Mott (MOTT) insulator phases



²Yunger Halpern et al., Phys. Rev. B **99**, 024203 (2019).

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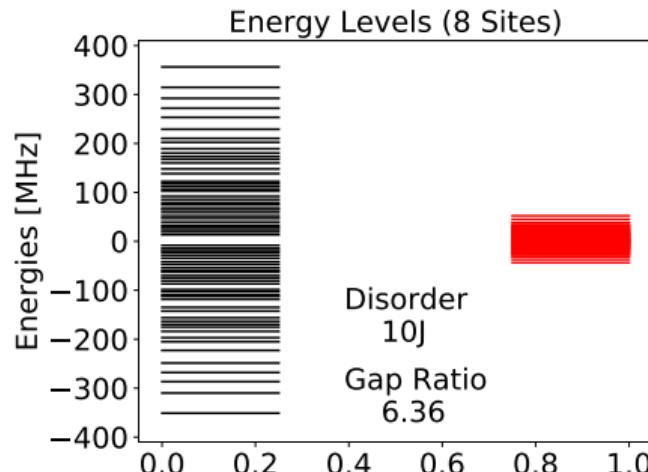
- ▶ Energy Spectrum of disordered (DIS) and Mott (MOTT) insulator phases
- ▶ Compared to the MBL engine, the NGSD engine does *not* operate at the middle of the spectrum



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Near Ground State Disordered (NGSD) Engine

- ▶ Energy Spectrum of disordered (DIS) and Mott (MOTT) insulator phases
- ▶ Compared to the MBL engine, the NGSD engine does *not* operate at the middle of the spectrum
- ▶ This allows us to achieve high efficiencies ($1 - 1/\text{Gap Ratio}$) in shorter times



²Yunger Halpern et al., Phys. Rev. B 99, 024203 (2019).

Bose Hubbard Model (Standard)

► Hamiltonian³

$$H = - \sum_{\langle ij \rangle} J_{ij} (a_i^\dagger a_j + \text{h.c.}) + \frac{U}{2} \sum_i n_i (n_i + 1) + \sum_i (h_i - \mu) n_i \quad (1)$$



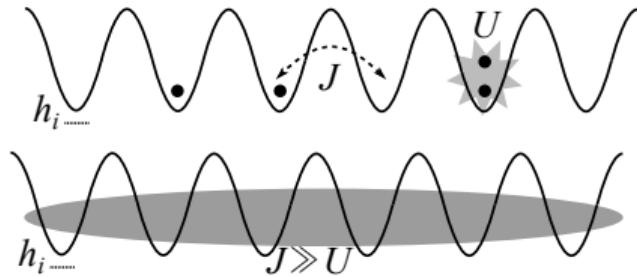
³Ma et al., Nature 566, 51 (2019).

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- ▶ Superfluid phase



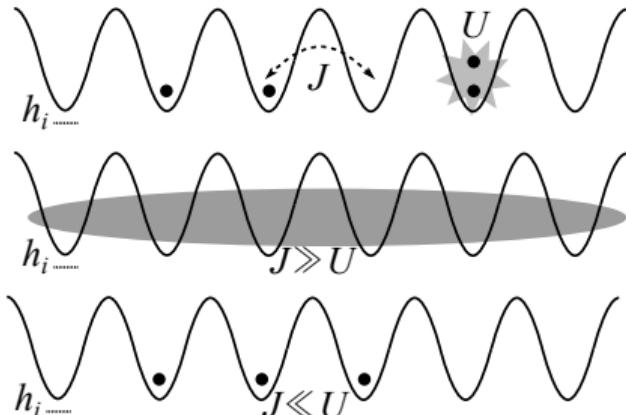
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- ▶ Superfluid phase
- ▶ Mott insulator phase (MOTT)



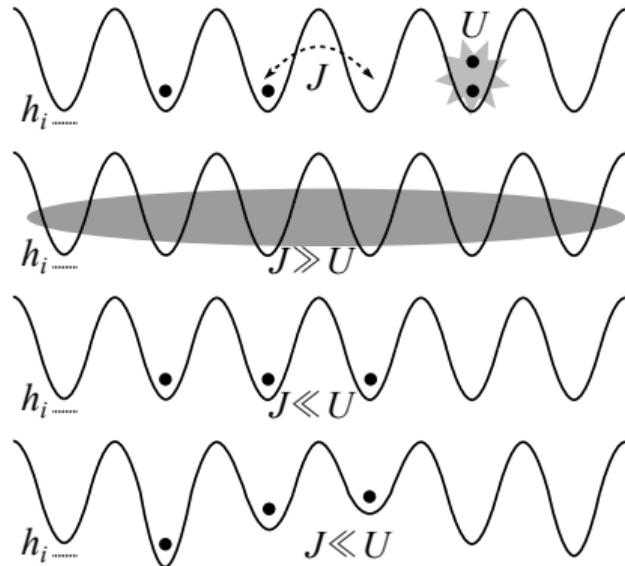
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- ▶ Superfluid phase
- ▶ Mott insulator phase (MOTT)
- ▶ Disordered phase (DIS)



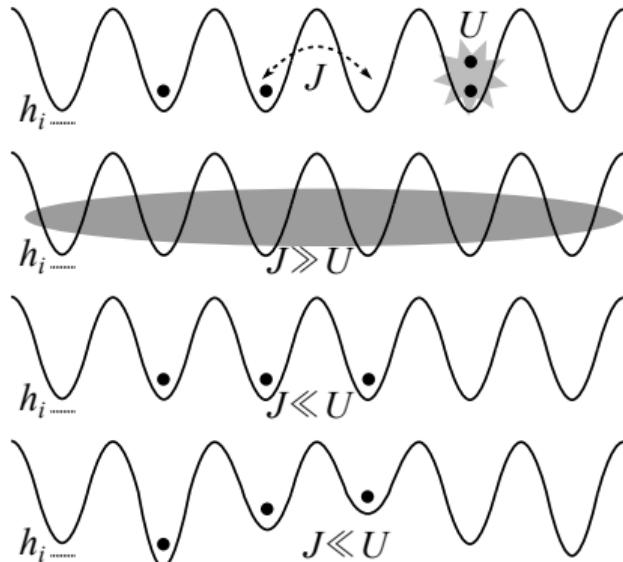
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Quantum Annealing

- ▶ Adiabatic evolution of $H_0 \rightarrow H_p$ ⁴

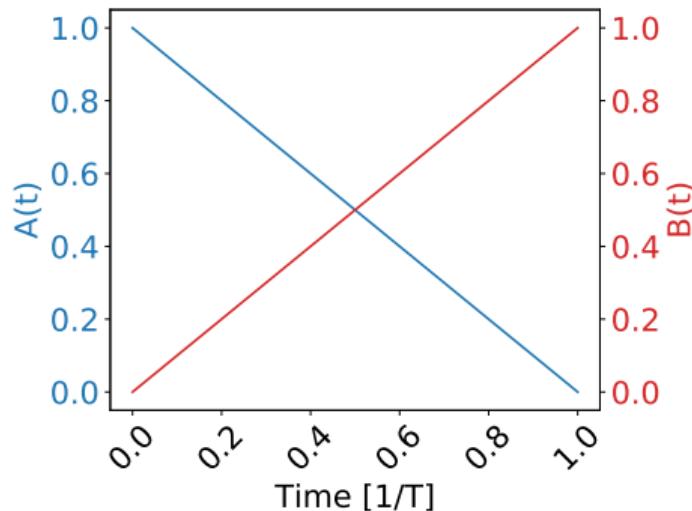
$$H(t) = H_0 A(t) + H_p B(t) \quad (2)$$

⁴Farhi et al. 2000.

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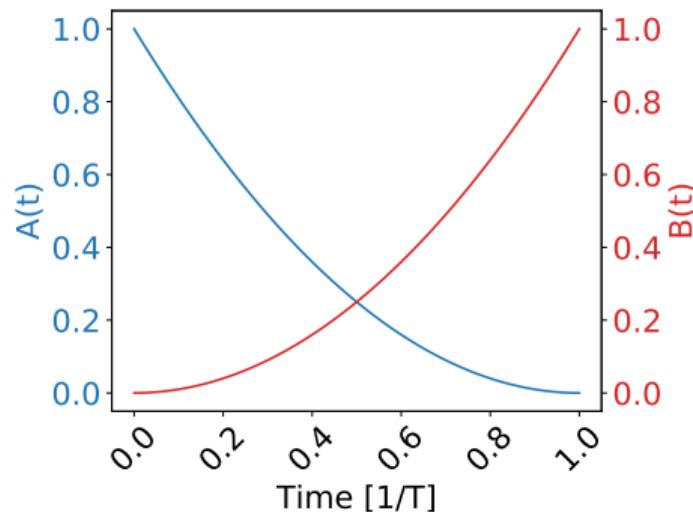
- ▶ Adiabatic evolution of $H_0 \rightarrow H_p$ ⁴

$$H(t) = H_0 A(t) + H_p B(t) \quad (2)$$

- ▶ Bose Hubbard as annealing problem

$$H(t) = H_0 + C(t) \sum_i h_i n_i, \quad (3)$$

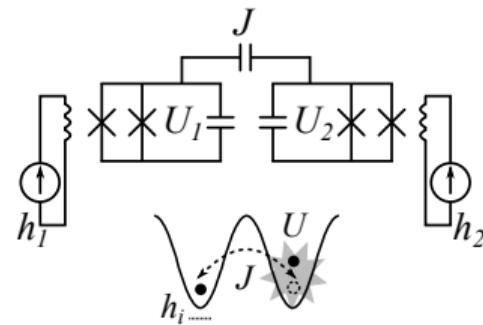
$$C(t) = \begin{cases} A(t), & \text{DIS} \rightarrow \text{MOTT} \\ B(t), & \text{MOTT} \rightarrow \text{DIS} \end{cases} \quad (4)$$



⁴Farhi et al. 2000.

Proposed Experimental Realization

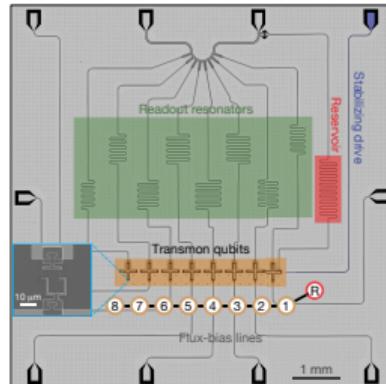
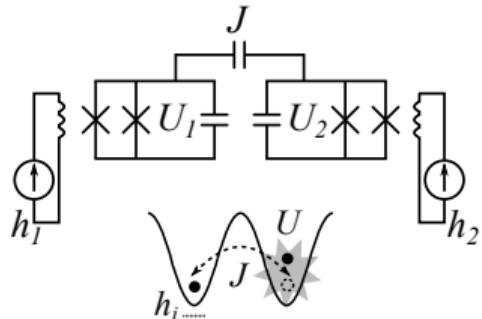
1. Superconducting circuit elements³



³Ma et al., Nature 566, 51 (2019).

Proposed Experimental Realization

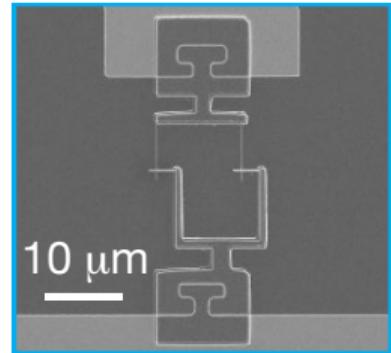
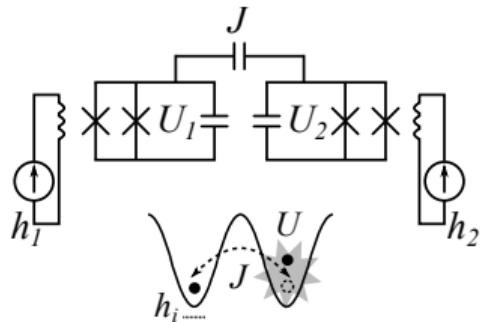
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2. Device layout



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Proposed Experimental Realization

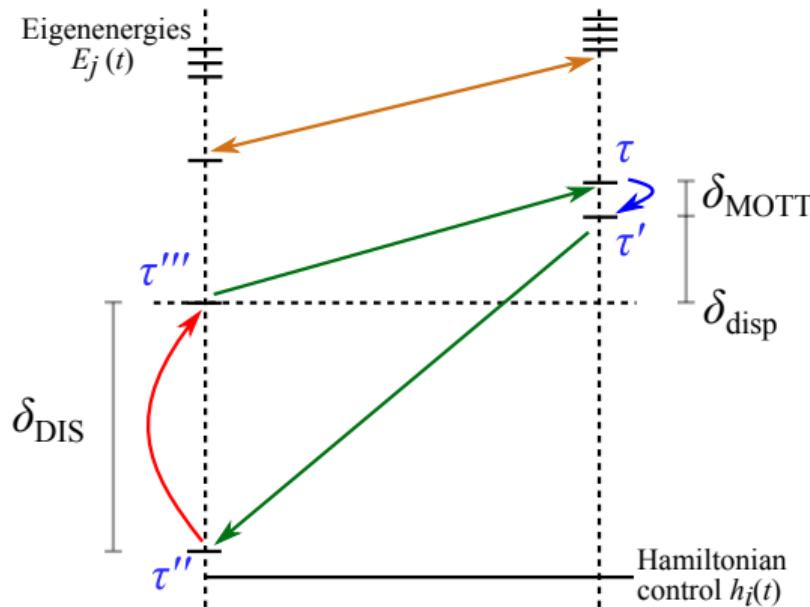
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2. Device layout
3. SQUIDs forming tunable transmons



³Ma et al., Nature 566, 51 (2019).

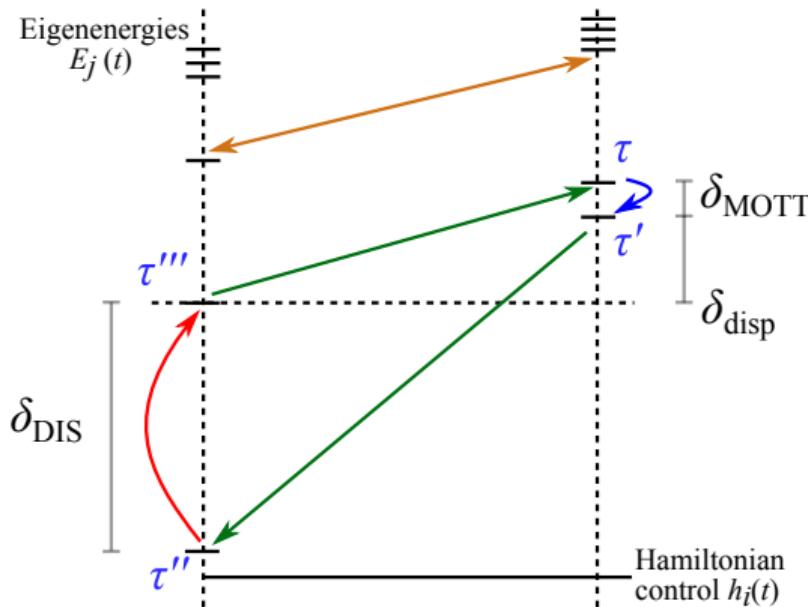
Results – *Compression & Expansion Strokes*

1. Evolve adiabatically between MOTT & DIS eigenstates, $(\tau' \rightarrow \tau'')$, $(\tau''' \rightarrow \tau)$



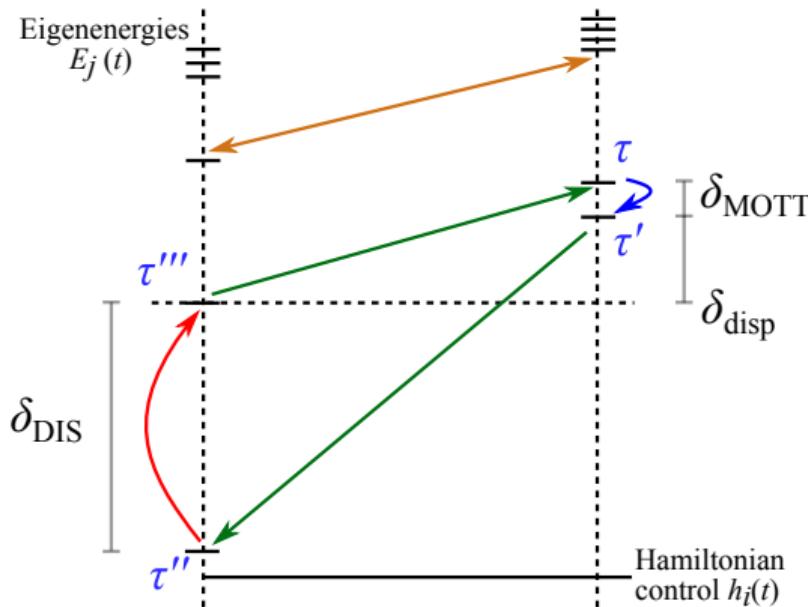
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1. Evolve adiabatically between MOTT & DIS eigenstates, $(\tau' \rightarrow \tau'')$, $(\tau''' \rightarrow \tau)$
2. Calculate fidelity $|\langle \psi_f | \psi_t \rangle|^2$ and residual energy $|E_t - \langle H \rangle|$



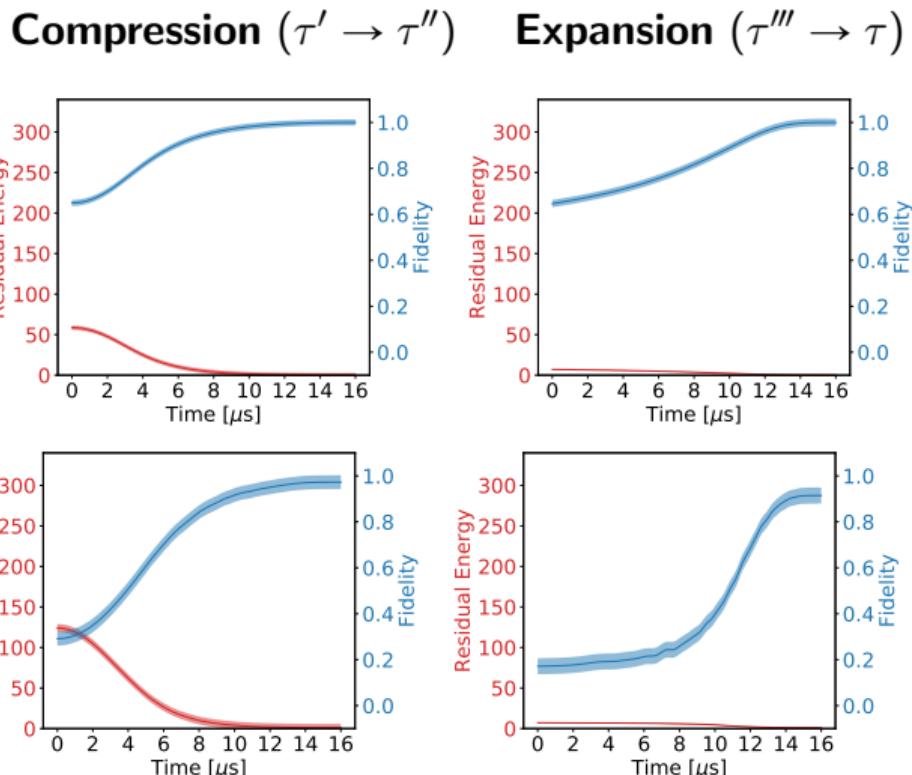
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3. Repeat, average over 100 different $h_i \in [-10J, 10J]$ realizations



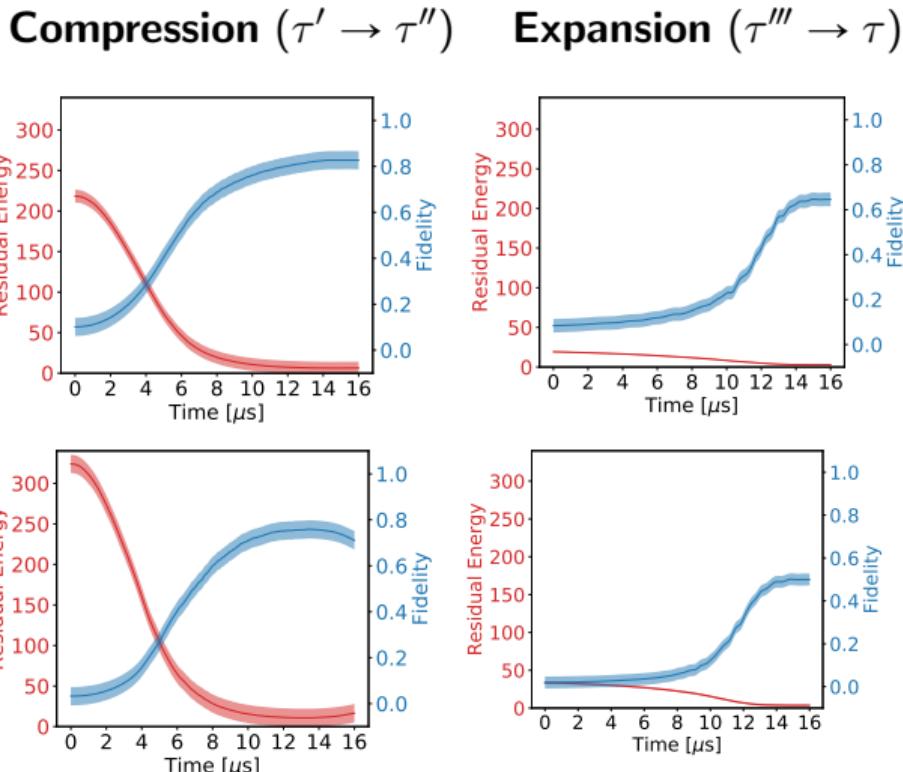
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4. 2,4 sites



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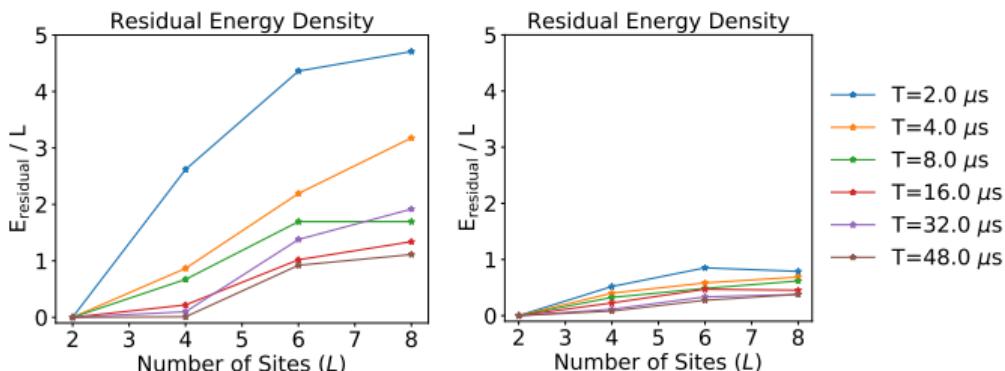
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3. Repeat, average over 100 different $h_i \in [-10J, 10J]$ realizations
4. 6,8 sites



Results – *Compression & Expansion Strokes*

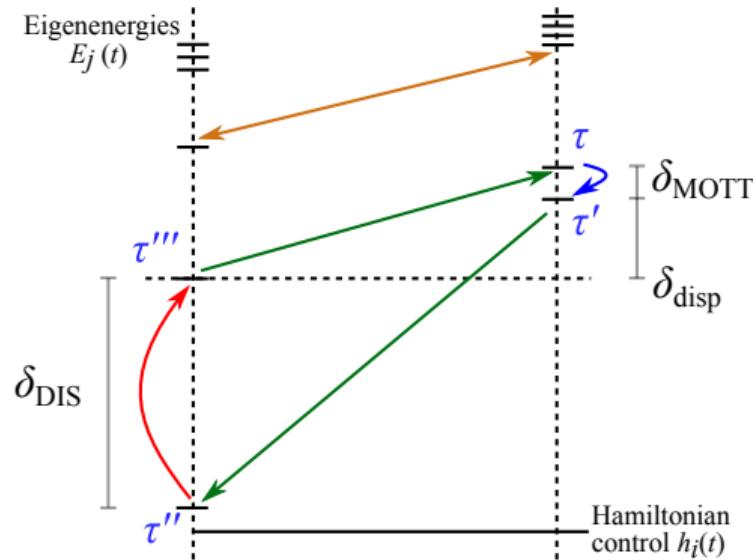
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4. Residual Energy Density

Compression ($\tau' \rightarrow \tau''$) **Expansion** ($\tau''' \rightarrow \tau$)



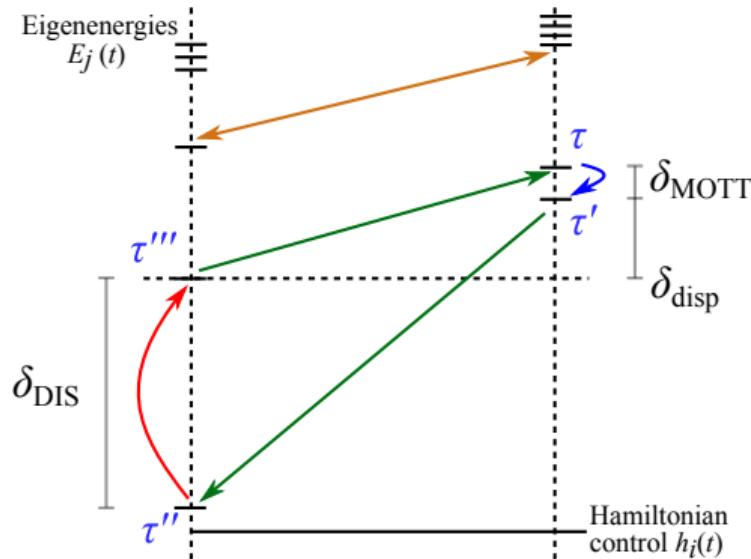
Outlook and Conclusions

- ▶ NGSD engine projected to have efficiency as function of ratio of MOTT / DIS gaps



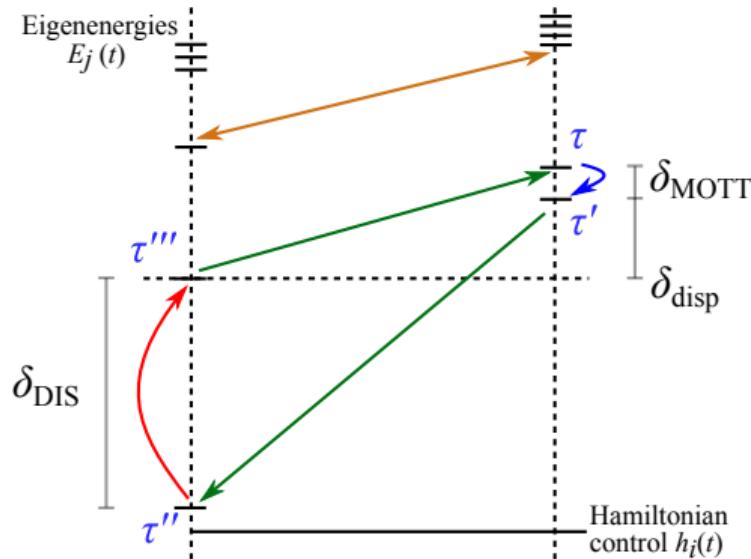
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Outlook and Conclusions

- ▶ NGSD engine projected to have efficiency as function of ratio of MOTT / DIS gaps
- ▶ Adiabatic time scales for *compression and expansion* strokes look promising for $N_s \leq 6$
- ▶ Repeat calculations on the DWAVE machine, explicit modeling of the cold bath



Acknowledgements

- ▶ Kapit Group: Eric Jones, David Rodríguez Pérez, Zhijie Tang, Mallory Zabrocky
 - F07.00006** Variational Preparation of Quantum Hall States on a Lattice
 - L16.00004** Is Fault Welcoming Quantum Computing Realistic?
 - L16.00010** Oscillatory quantum optimization methods applied to problems with large ground state bands
 - M09.00005** Considerations for incorporating small logical qubits in digital error correction codes
- ▶ Graduate Fellowship for Science, Technology, Engineering, and Mathematics Diversity (formerly, National Physical Sciences Consortium Fellowship)
- ▶ NSF Grant PHY-1653820

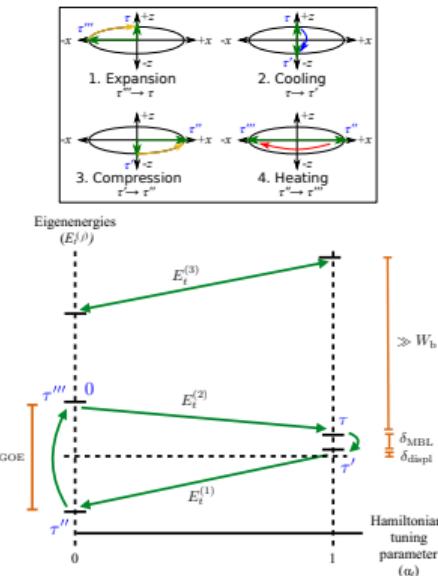
Additional Slides

Many Body Localized (MBL) Engine

- ▶ Single qubit engine cycle³

$$H(t) = (1 - \alpha_t) \frac{\delta_{\text{GOE}}}{2} \sigma^x + \alpha_t \frac{\delta_{\text{MBL}}}{2} \sigma^z \quad (5)$$

1. $\tau''' \rightarrow \tau$, Expansion
2. $\tau \rightarrow \tau'$, Cooling
3. $\tau' \rightarrow \tau''$, Compression
4. $\tau'' \rightarrow \tau'''$, Heating



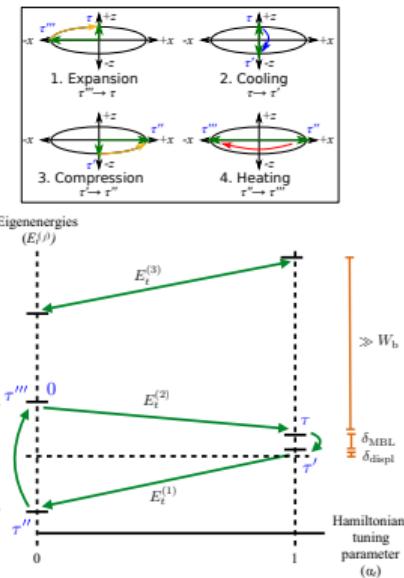
³Yunger Halpern et al., Phys. Rev. B 99, 024203 (2019).

Many Body Localized (MBL) Engine

- ▶ Heat and work absorbed³

$$Q := \int_0^\tau \text{Tr} \left(\rho \frac{dH}{dt} \right) dt \quad (6)$$

$$W := \int_0^\tau \text{Tr} \left(\frac{d\rho}{dt} H \right) dt \quad (7)$$



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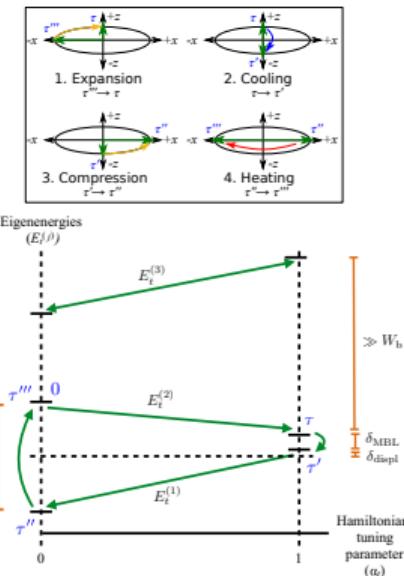
Many Body Localized (MBL) Engine

- ▶ Applied to qubit MBL engine³

$$\langle Q_{\text{in}} \rangle = \langle Q_4 \rangle = \frac{\delta_{\text{GOE}}}{2} \quad (5)$$

$$\langle W_{\text{tot}} \rangle = \frac{\delta_{\text{GOE}}}{2} + \frac{\delta_{\text{MBL}}}{2} \quad (6)$$

$$\eta_{\text{MBL}} = \frac{\langle W_{\text{tot}} \rangle}{\langle Q_{\text{in}} \rangle} = 1 - \frac{\delta_{\text{GOE}}}{\delta_{\text{MBL}}} \quad (7)$$



³Yunger Halpern et al., Phys. Rev. B 99, 024203 (2019).

Bose Hubbard Model (Two-Level Perturbative)

▶ Hamiltonian

$$\begin{aligned} H = & -J \sum_{i=1}^{N_s-1} (\sigma_i^+ \sigma_{i+1}^- + \text{h.c.}) \\ & + \frac{4J}{U^2} \sum_{i=1}^{N_s-1} \frac{1}{2} (\mathbb{1} + \sigma_i^z) \frac{1}{2} (\mathbb{1} + \sigma_{i+1}^z) \\ & + \frac{2J}{U^2} \sum_{i=1}^{N_s-2} (\sigma_{i+2}^+ \sigma_{i+1}^+ \sigma_{i+1}^- \sigma_i^- + \text{h.c.}) \quad (8) \end{aligned}$$



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▶ Note, all operators are now particle number conserving

$$\sigma_i^+ \sigma_j^- |ij\rangle = |ji\rangle \quad (9)$$

$$\sigma_i^z |ij\rangle = \text{sign}(i-1) |ij\rangle \quad (10)$$

References

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