

Higher order singular-value decomposition for strongly correlated systems

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Virginia Tech

Our group's efforts to model strong correlation

1/21



1. Spin-flip-EA-IP methods

- Strong correlation for spatial/spin near-degeneracies
 - Shannon Houck's talk: COMP 145



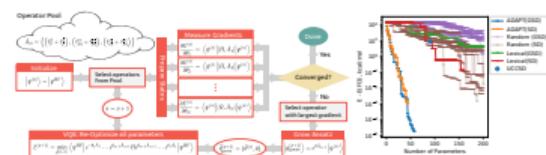
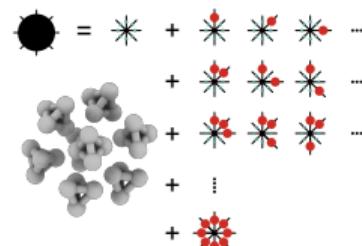
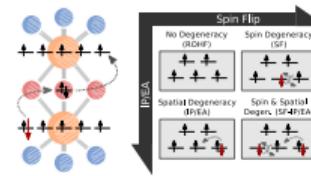
2. *n*-body Tucker approximation

- Tensor decomposition
 - Strong correlation for “quasi-clusterable” systems



3. ADAPT-VQE

- Strong correlation for all(?) systems
 - Quantum computation



Any two algorithms are equivalent when their performance is averaged across all possible problems

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 - PT, CC, ...
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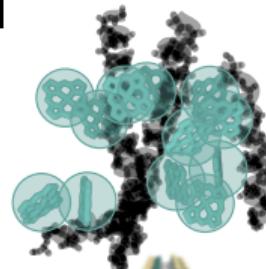
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- Low-dimensionality
 - DMRG, PEPS, ...
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- “**Quasi-Clusterable**”
 - Our focus here

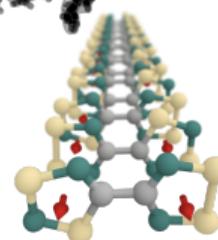
Artificial Photosynthesis



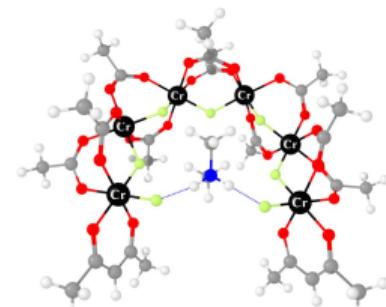
Photovoltaics
(Singlet Fission)



Molecule-based Magnets



Transition-Metal Chemistry



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Artificial
Photosynthesis

Photovoltaics
(Solar Cells)

Recent methodological advances:

- Block-correlated CC (BCCC): Li, JCP, 120, 5017, (2004)
- Renormalized Exciton Method (REM): Al Hajj et al. PRB, 72, 224412, (2005)
- Active-space Decomposition (ASD): Parker, et al. JCP. 139(2), 21108, (2013)
- Ab Initio Frenkel/Davydov Exciton (AIFDEM): Morrison & Herbert. JCTC, 10, 5366, (2014)
- Cluster mean-field (cMF): Jiménez-Hoyos & Scuseria. PRB, 92(8), 85101, (2015)
- excitonic CC (XCC): Lui & Dutoi, arXiv:1709.01966 (2017)
- “Quasi-Clusterable”
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- Imagine a “clusterable” molecular system comprised of 3 clusters,

$$\begin{aligned} |\psi\rangle &= C_{ijk} |A_i B_j C_k\rangle \\ &= C_{\alpha\beta\gamma} U_{i\alpha} U_{j\beta} U_{k\gamma} |A_i B_j C_k\rangle \quad \text{Tucker Decomposition} \\ &= C_{\alpha\beta\gamma} |\tilde{A}_\alpha \tilde{B}_\beta \tilde{C}_\gamma\rangle \end{aligned}$$

- Working in this compressed space rarely solves the problem - “curse of dimensionality”



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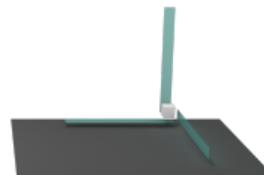
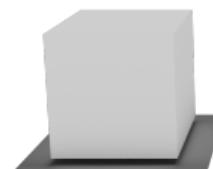
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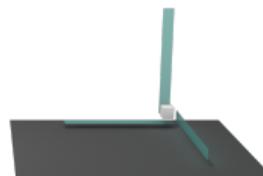
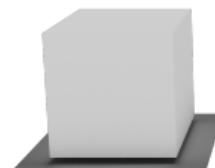




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- Define projector onto *important* and *unimportant* vectors for cluster, A :

$$\hat{P}_A = U_{i\alpha} U_{j\alpha}$$

$$\hat{Q}_A = \hat{I} - \hat{P}_A$$

- Assume that \hat{Q}_A is “small”

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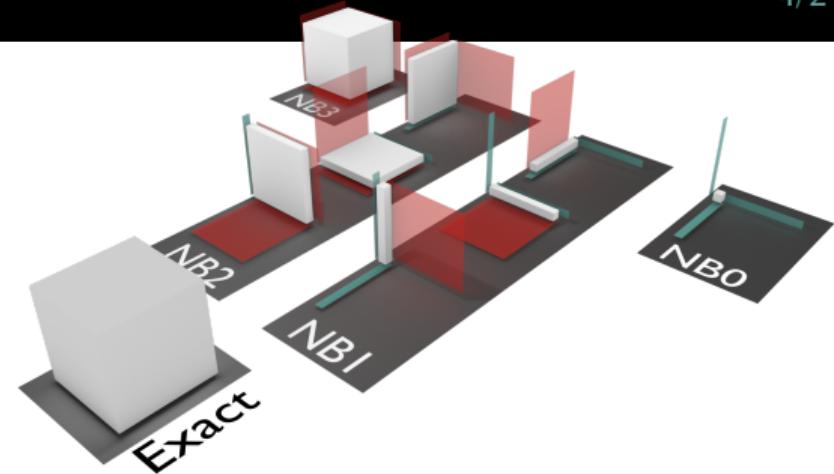
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0th Order

1th Order

2nd Order

3rd Order

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n-body Tucker

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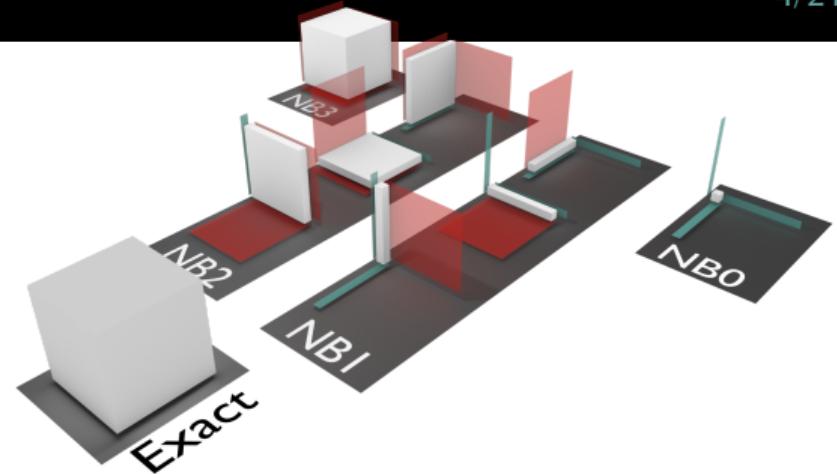
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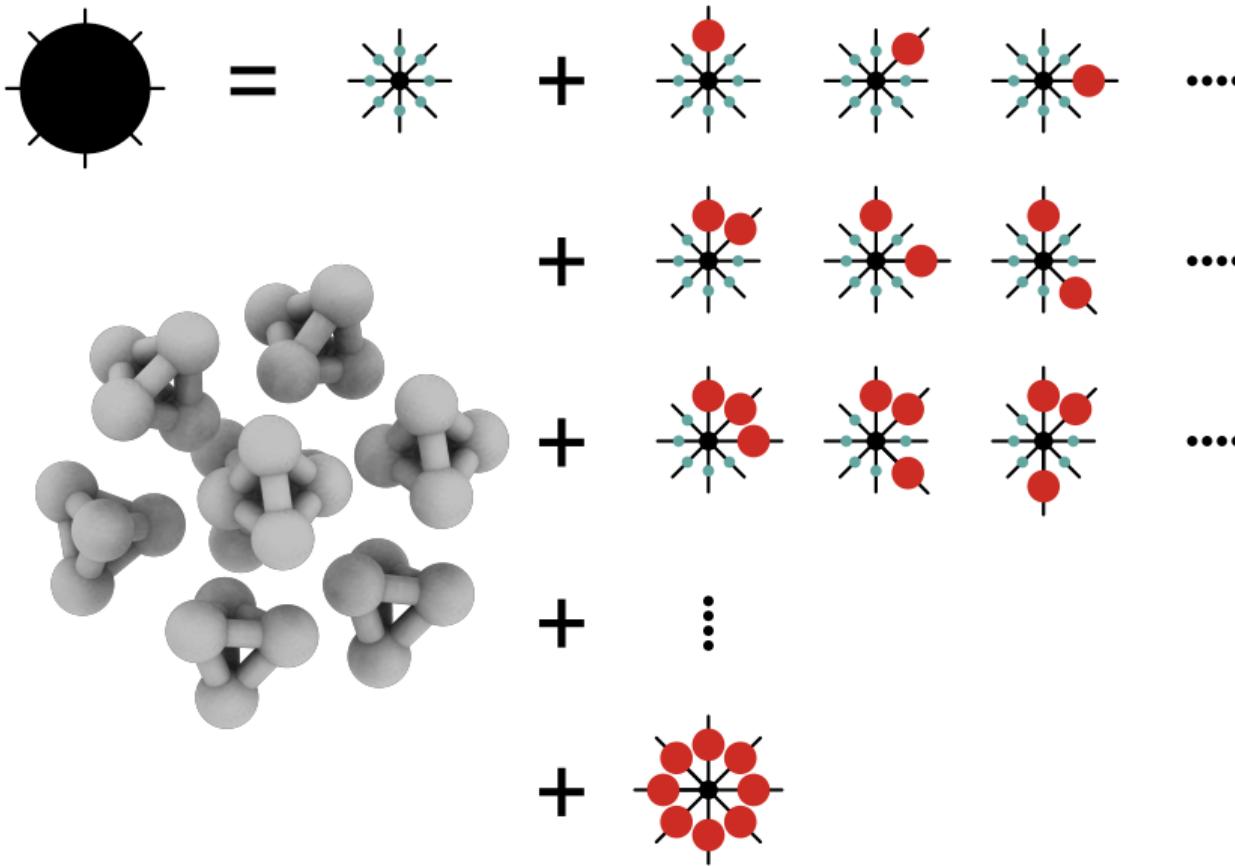
2nd Order

3rd Order

0	A	B	C	AB	AC	BC	ABC
ABC	BC	AC	AB	C	B	A	0
ABC	BC	AC	AB	C	B	A	0

n -body Tucker

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- Heisenberg Hamiltonian implementation
- σ vector elements:

$$\sigma_{\alpha\beta\gamma}^k \Leftarrow \langle A_\alpha B_\beta C_\gamma | \hat{H}_{AB} | A_{\alpha'} B_{\beta'} C_{\gamma'} \rangle C_{\alpha'\beta'\gamma'}^k$$

	0	A	B	C	AB	AC	BC	ABC
ABC	BC	AC	AB	C	B	A	0	
BC	AC	AB	C	B	A	0		
AC	AB	C	B	A	0			
AB	C	B	A	0				
C	B	A	0					
B	A	0						
A	0							
0								

$$\begin{aligned}
 &= \sum_{i \in A} \sum_{j \in B} \mathcal{J}_{ij} \langle A_\alpha B_\beta | (\hat{S}_i^+ \hat{S}_j^- + \hat{S}_i^- \hat{S}_j^+ + 2\hat{S}_i^z \hat{S}_j^z) | A_{\alpha'} B_{\beta'} \rangle \langle C_\gamma | C'_\gamma \rangle C_{\alpha'\beta'\gamma'}^k \\
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 &= H_{\alpha\beta\alpha'\beta'} C_{\alpha'\beta'\gamma'}^k
 \end{aligned}$$

- The non-interacting Hamiltonian inverse seems to be an effective preconditioner:

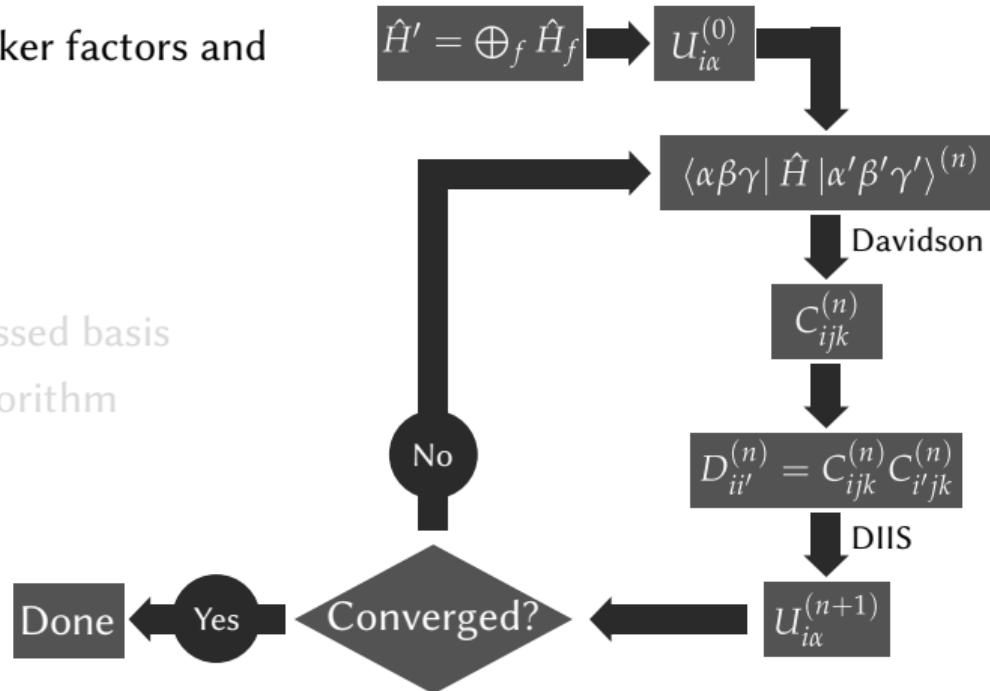
$$\hat{H} = \bigoplus_f \hat{H}_f$$

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- Two nested optimizations:
 1. Davidson: matrix diagonalization
 2. DIIS: tucker factor optimization
- All quantities computed in the compressed basis
- Davidson gets faster toward end of algorithm

n -body Tucker: Cluster states via HOSVD

7/21

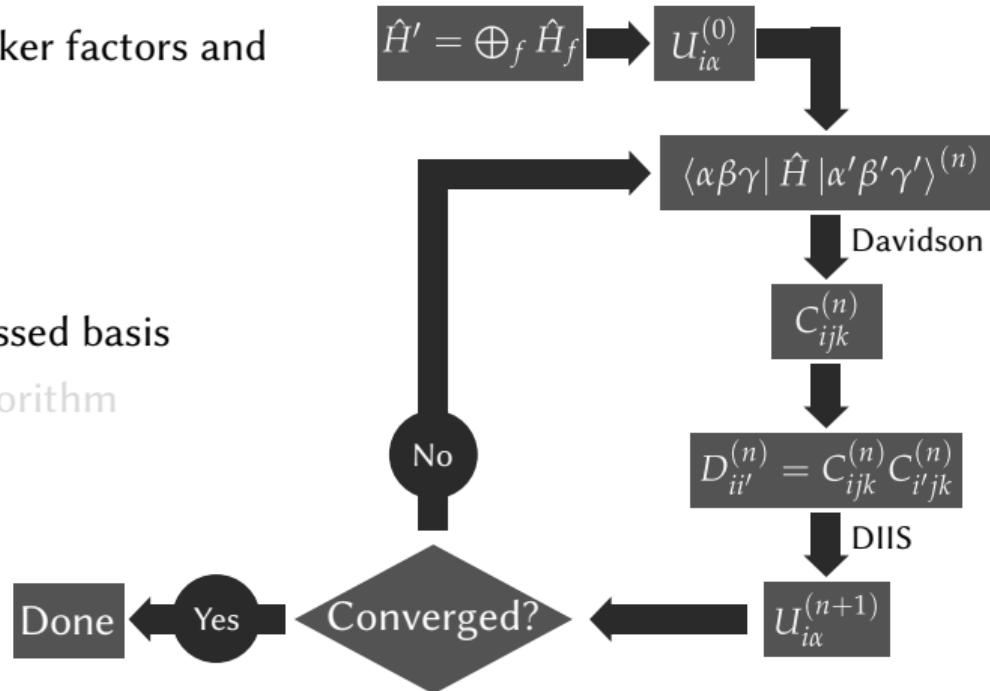
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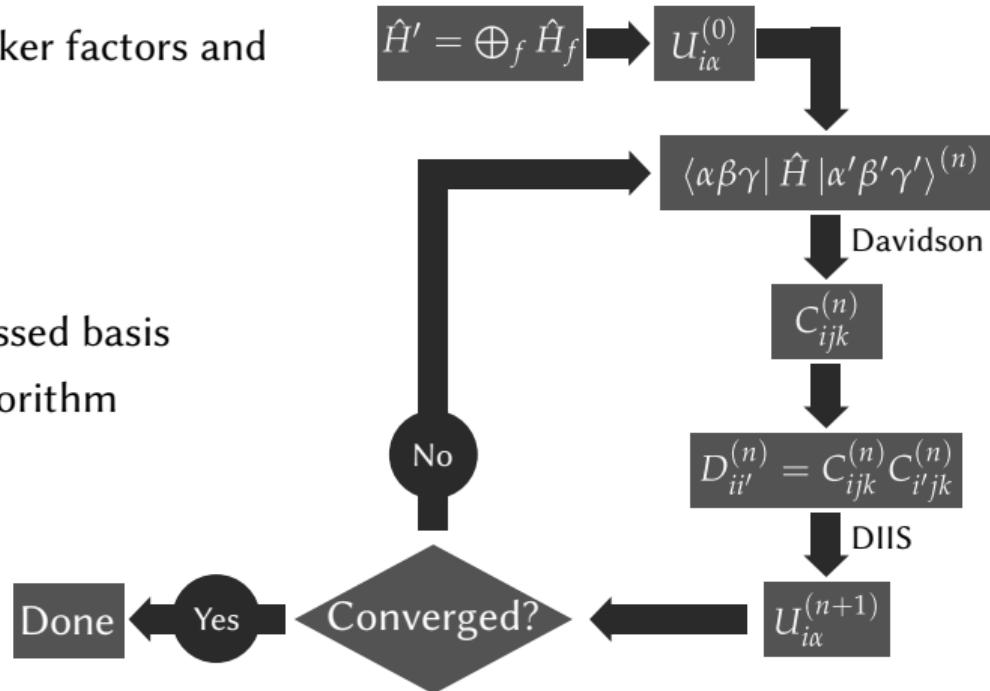
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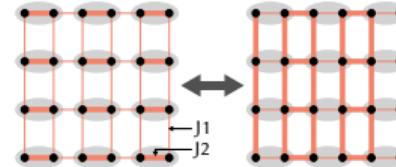
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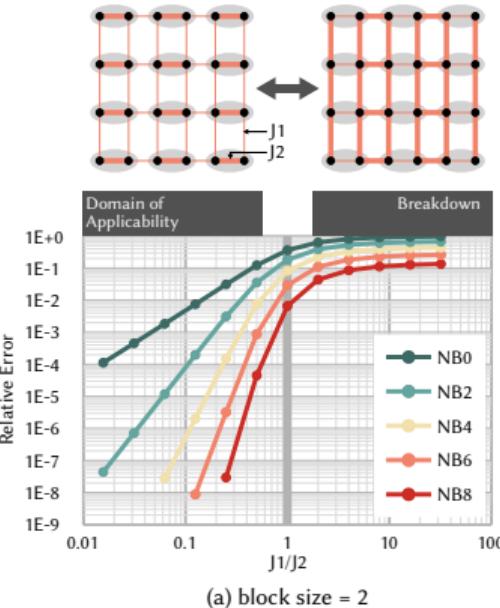
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- Clustered Hamiltonian
- Single P vector
- Exponential Convergence
- 4-site clusters
- For single P vector, Tucker optimization nearly negligible



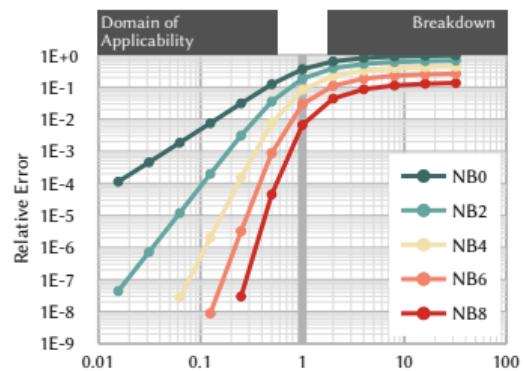
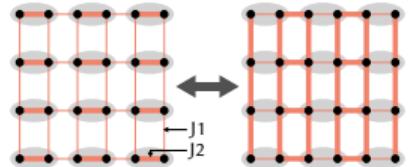
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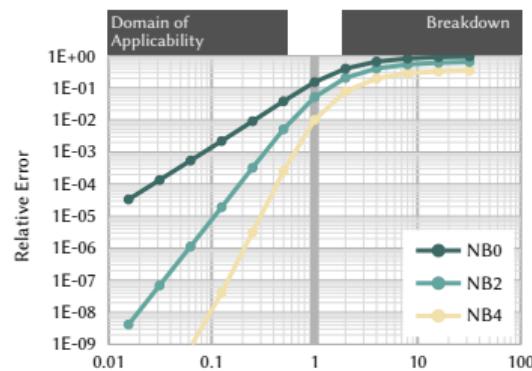
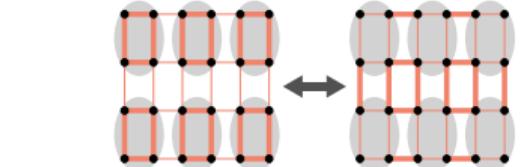
n -body Tucker: Variational optimization

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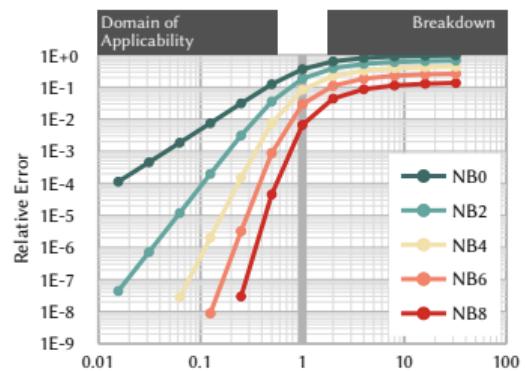
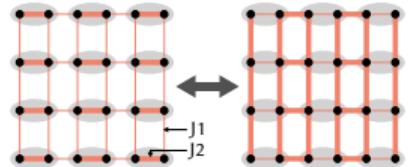


(a) block size = 2

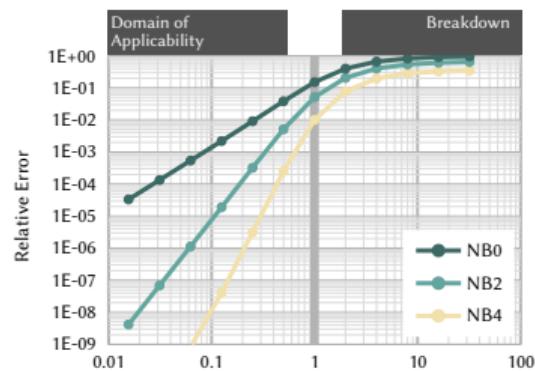
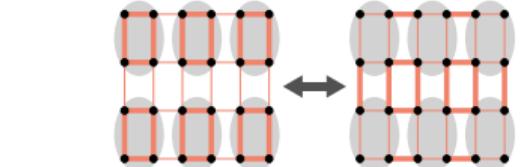


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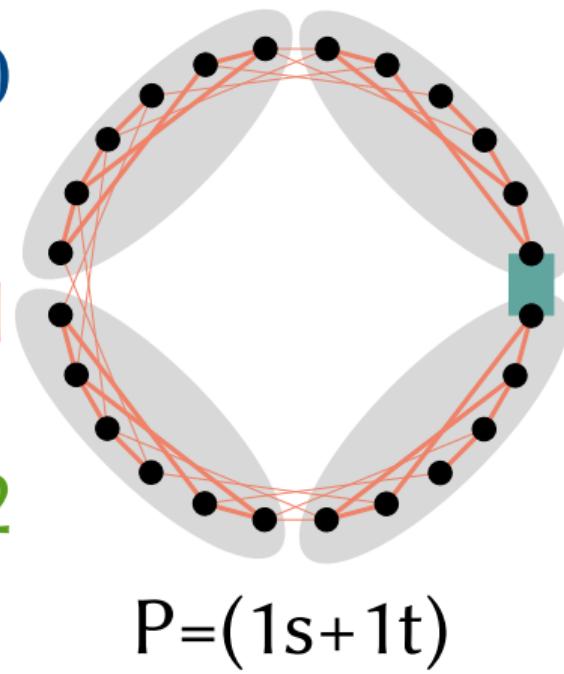
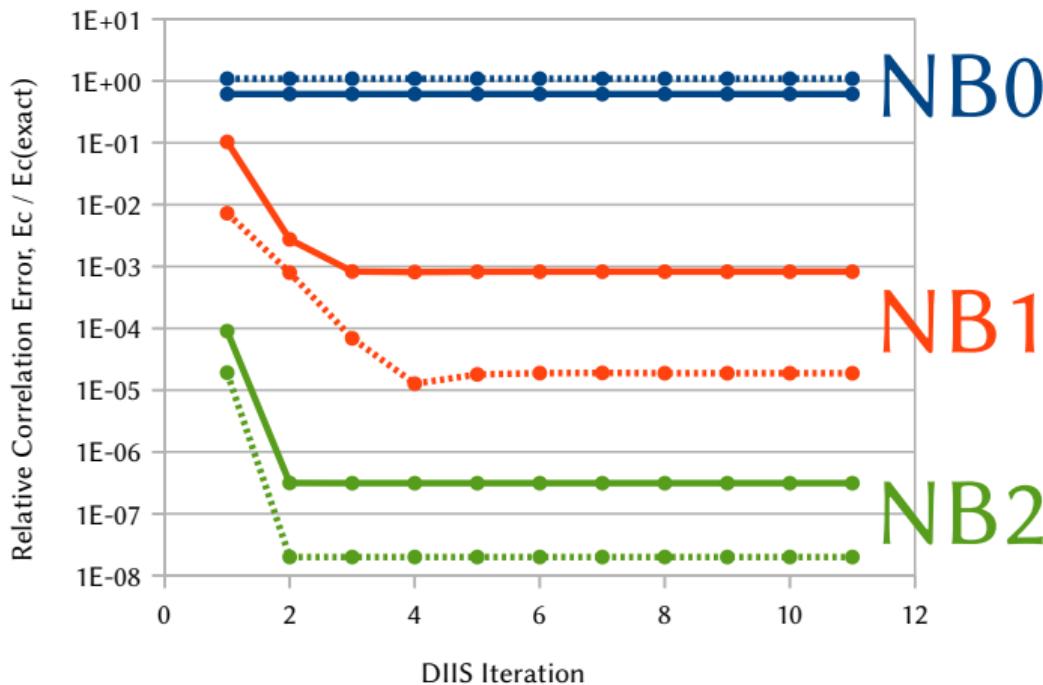


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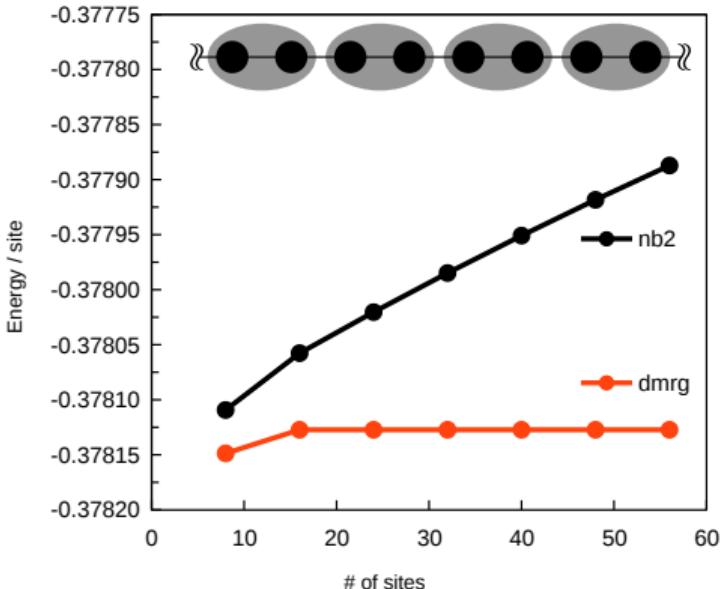


(b) block size = 4

- For multidimensional P spaces, significant benefit from HOSVD



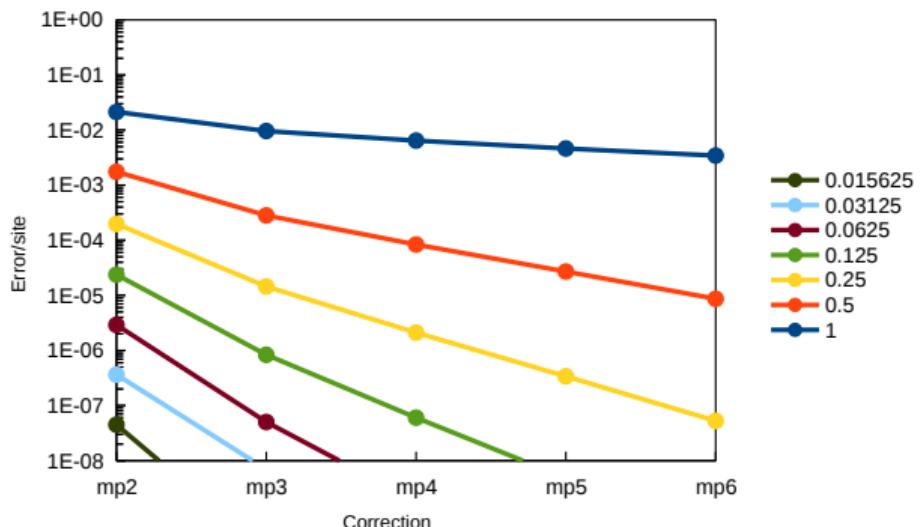
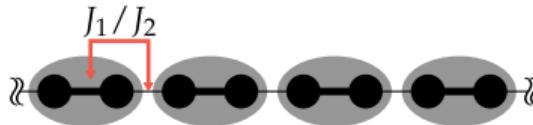
- Compare energy per site to DMRG calculations (ITensor)
- Variational solution not size-extensive



n -body Tucker: High order Perturbation theory

11/21

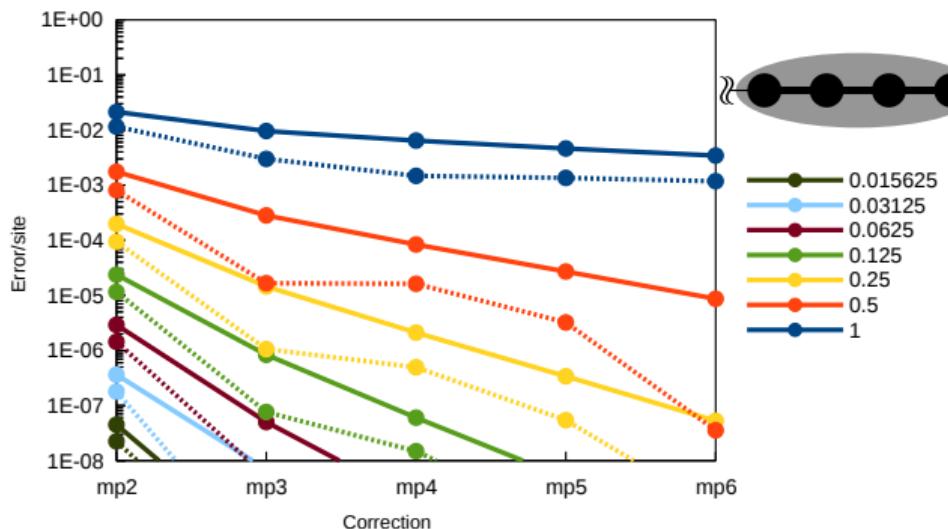
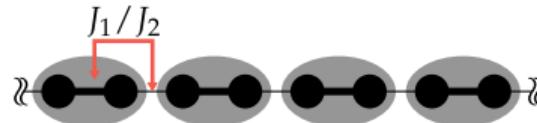
- Higher-order perturbations
- 1D system PBC, thermodynamic limit



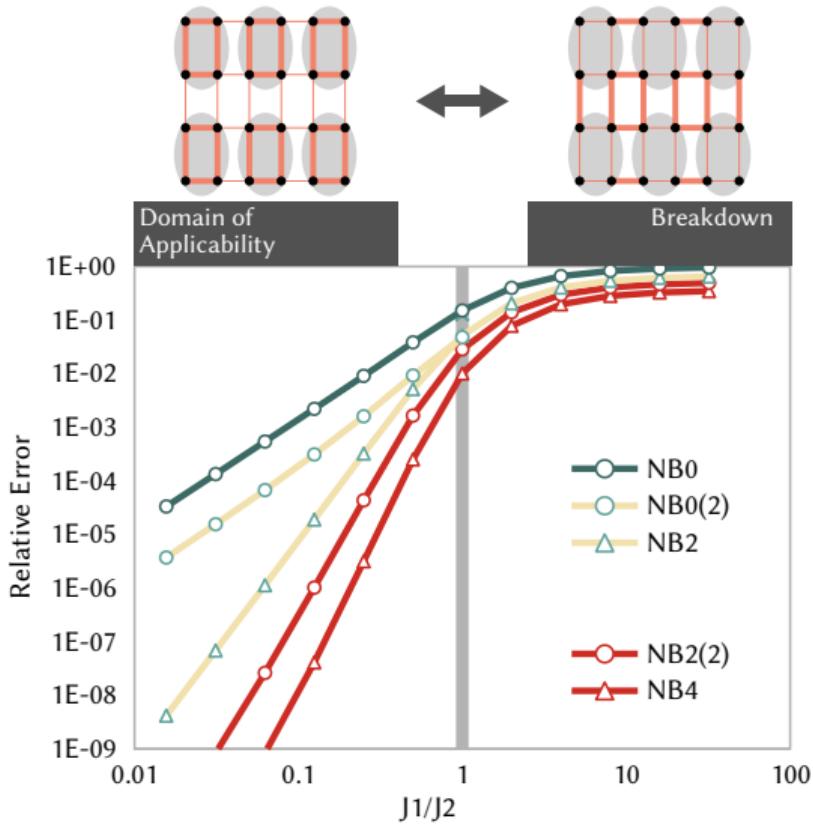
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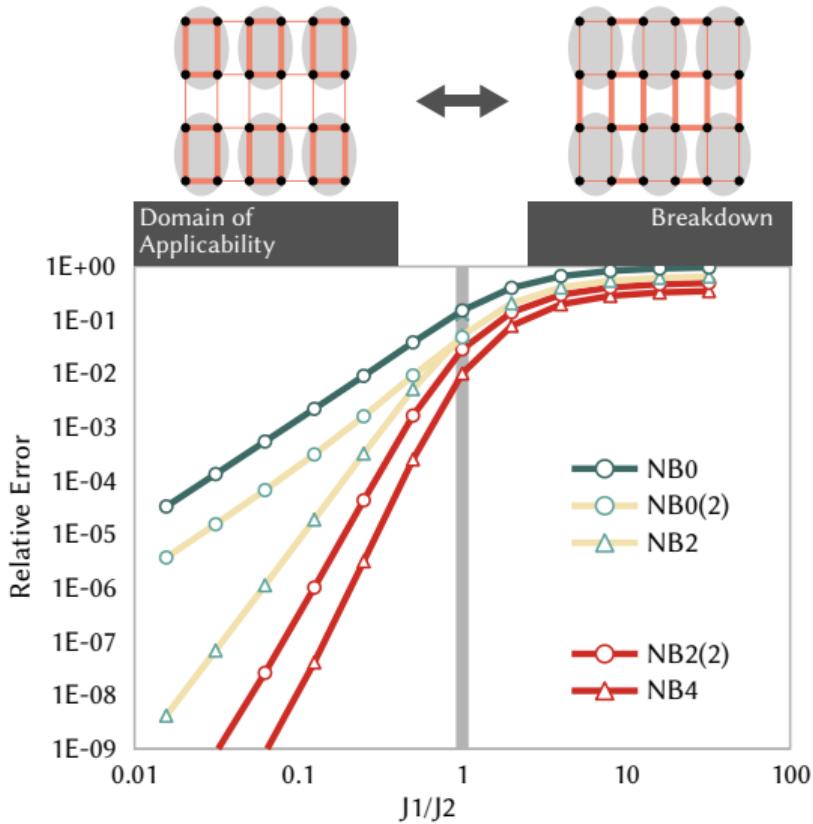
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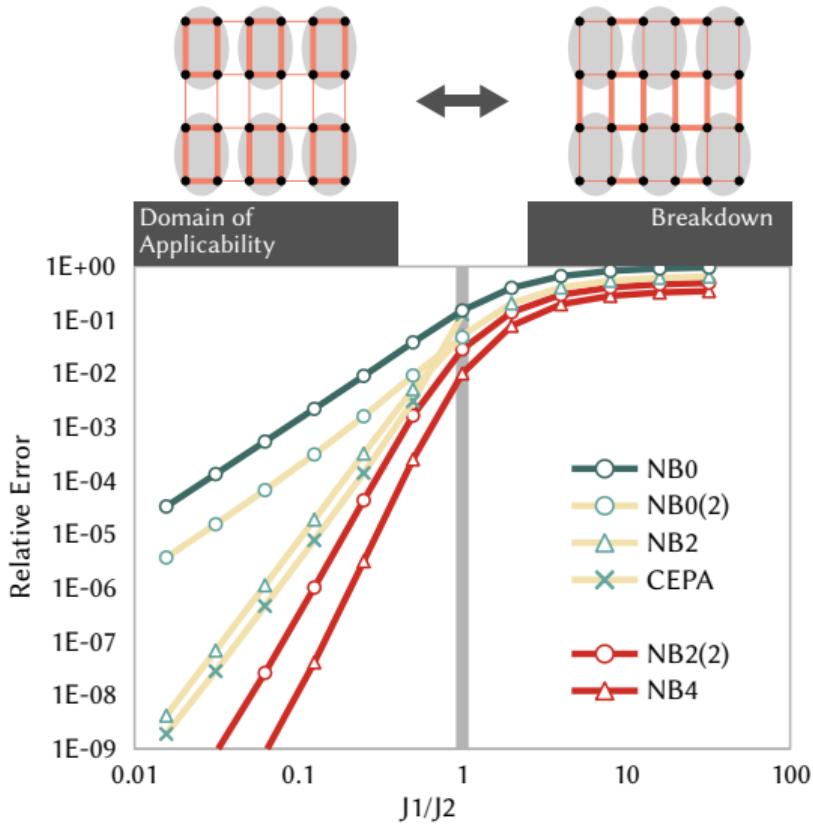
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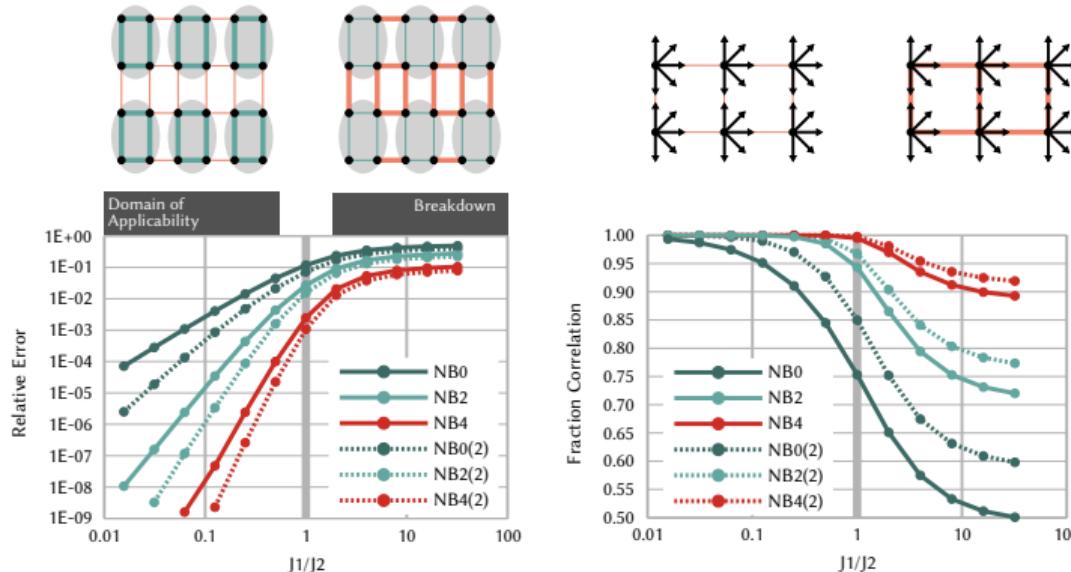
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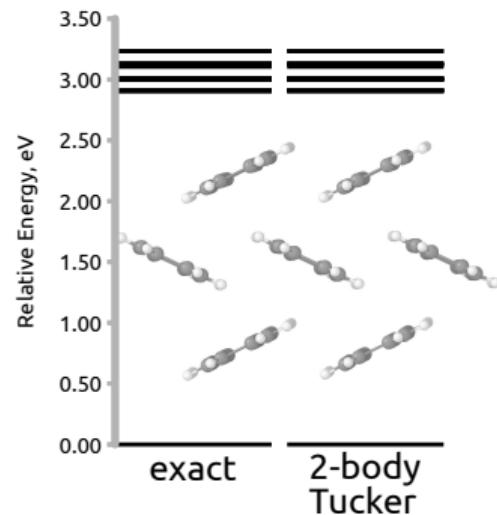


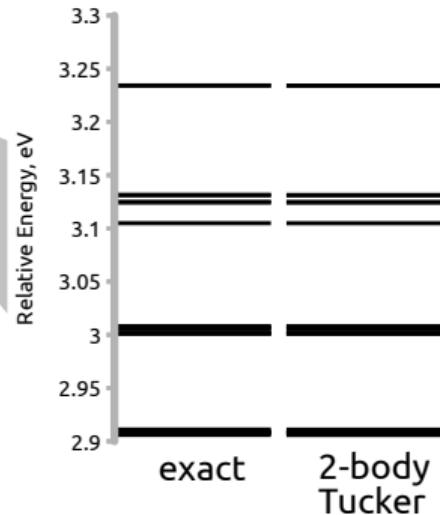
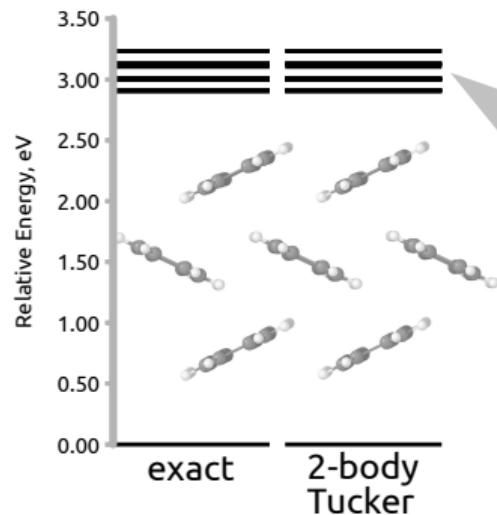
n -body Tucker: Ferromagnetic clusters

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- Magnetic clusters have degenerate groundstates
- Requires larger P space
- n -body terms correct a renormalized spin model

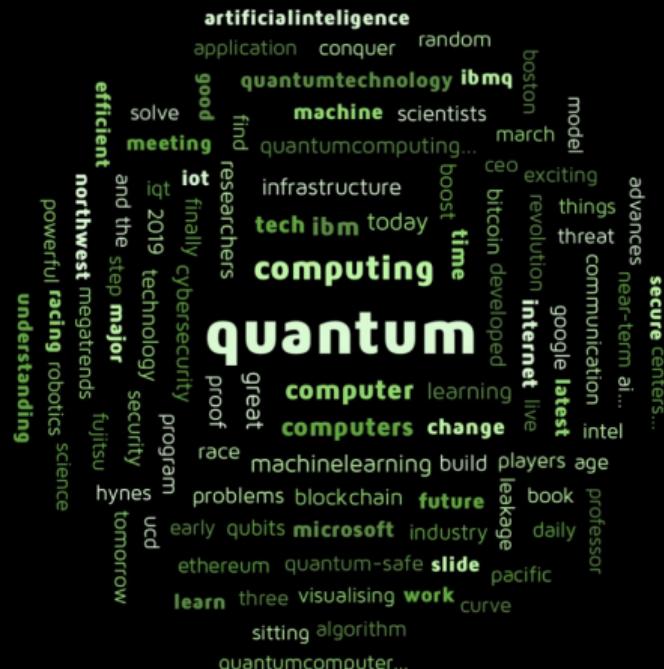






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 - *Quantum Computation* provide a possible solution

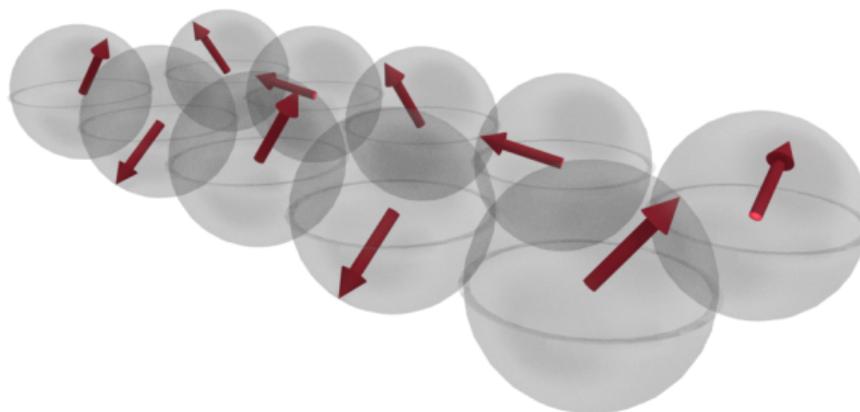


Pulled from #quantumcomputing twitter hashtag

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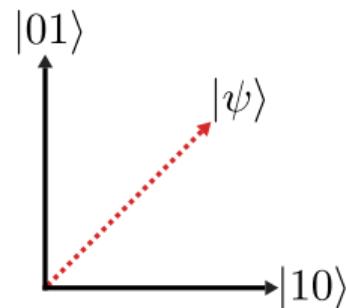
Pulled from #quantumcomputing twitter hashtag

- Computational framework comprised of *entangled* quantum bits (qubits)
- **Offensive simplification:** Qubits are “controllable two-level systems”
- Universal quantum computation is quite general
- **Our focus:** simulating quantum systems



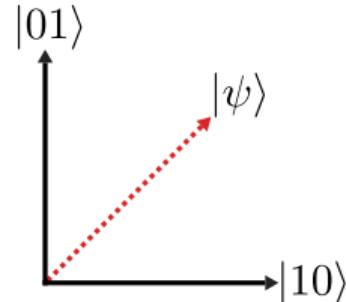
- Natural representation of quantum states → Computational speedup
- **Classical Computer:** Only represents “classical” basis vectors
 - Store (manipulate) quantum state as vector of numbers in memory
- **Quantum Computer:** Directly represents arbitrary vector in Hilbert space
 - Create (measure) quantum state on hardware to mimic simulated quantum state

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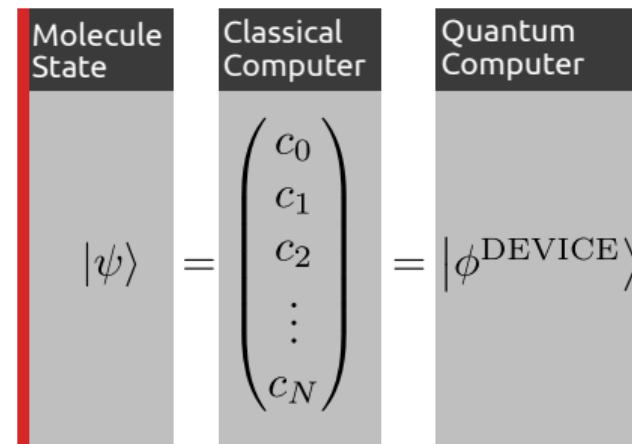
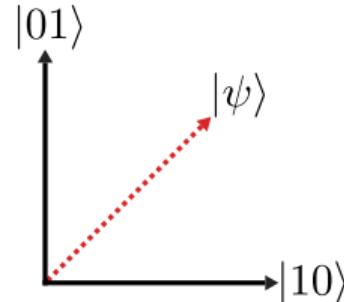
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Molecule State	$ \psi\rangle$	=	$\begin{pmatrix} c_0 \\ c_1 \\ c_2 \\ \vdots \\ c_N \end{pmatrix}$
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- Quantum Computers implement unitary rotations: $|\psi\rangle = \hat{U} |0\rangle$
- Unitary Coupled-Cluster

$$|\psi\rangle = \exp\left\{\sum_i \theta_i \hat{\tau}_i\right\} |0\rangle = \hat{U}(\vec{\theta}) |0\rangle$$

- For an arbitrary choice of parameters, $\vec{\theta}$, \hat{H} can be measured

$$E(\theta) = \langle \hat{H} \rangle_\theta = \sum_{pq} h_{pq} \left\langle \hat{p}^\dagger \hat{q} \right\rangle_\theta + \sum_{pqrs} \langle pq | rs \rangle \left\langle \hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r} \right\rangle_\theta$$

- Variationally minimize to choose optimal $\vec{\theta}$: **Variational Quantum Eigensolver (VQE)**

$$\frac{\partial E(\theta)}{\partial \theta} = 0$$

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 - only 1 or 2 qubit rotations generally available – Trotterization:

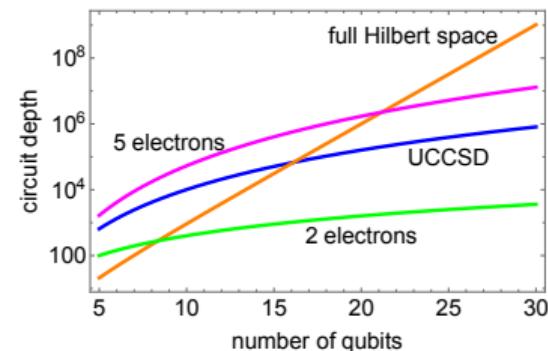
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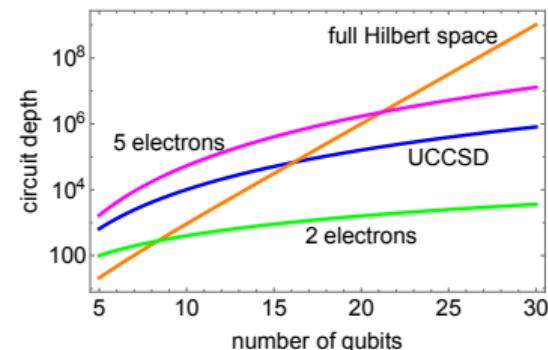
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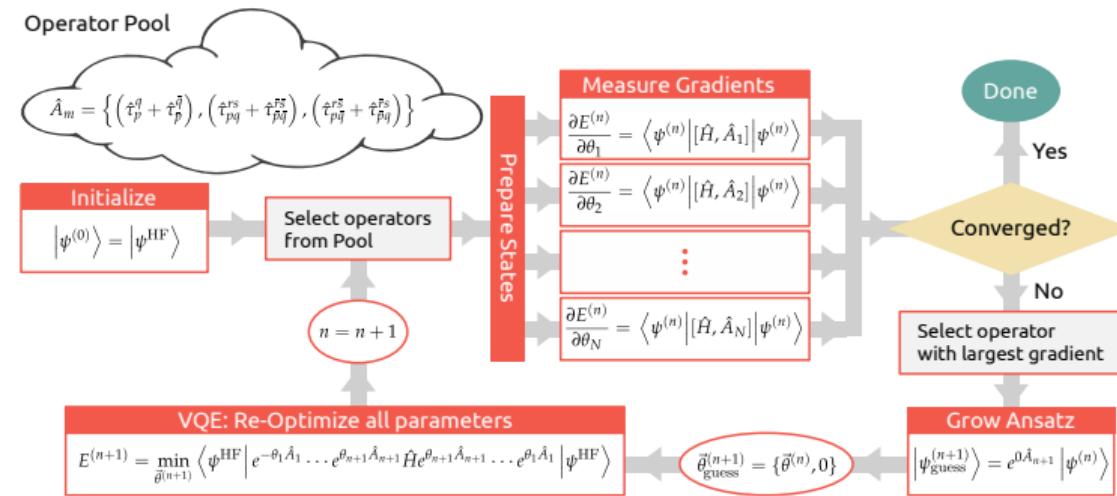
- Grow the ansatz by sequentially choosing the rotation with the largest energy derivative at each iteration
- Adaptive Derivative-Assembled Pseudo-Trotter ansatz: ([ADAPT-VQE](#))
 - ✓ Converges to exact solution: FCI
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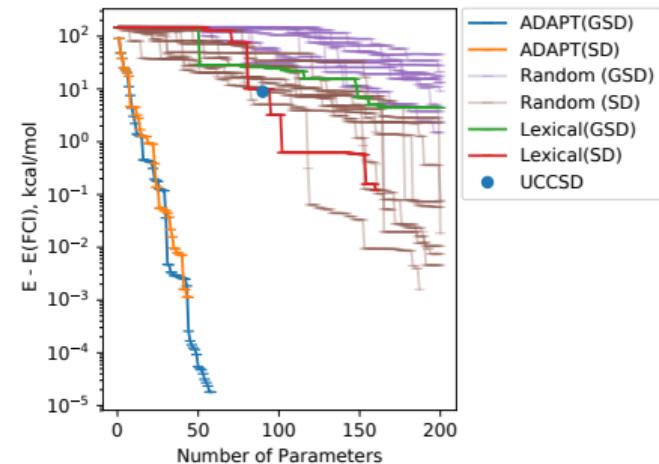
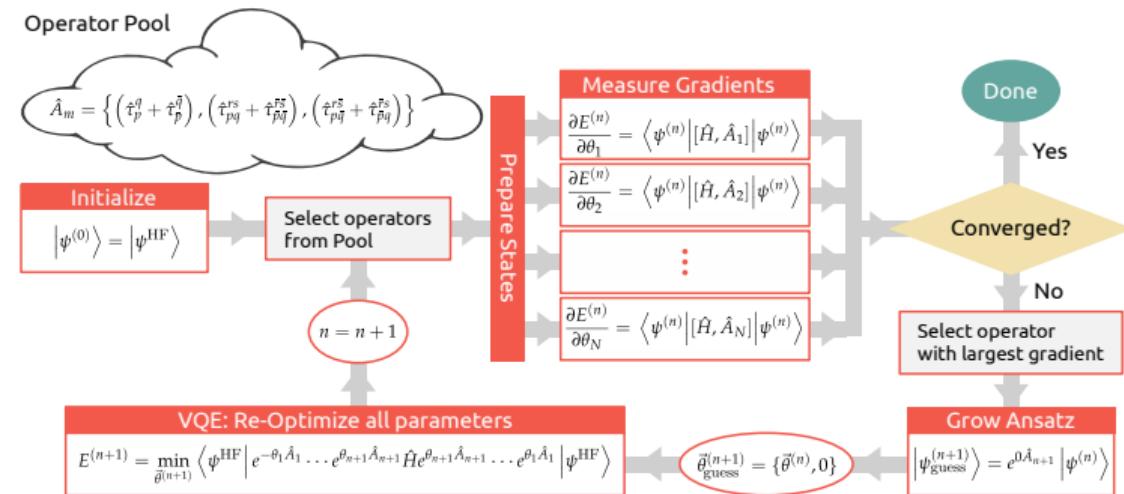
ADAPT-VQE: Algorithm

20/21



- BeH₂ at 2R_{eq} with STO-3G
- Extremely fast convergence with number of parameters!

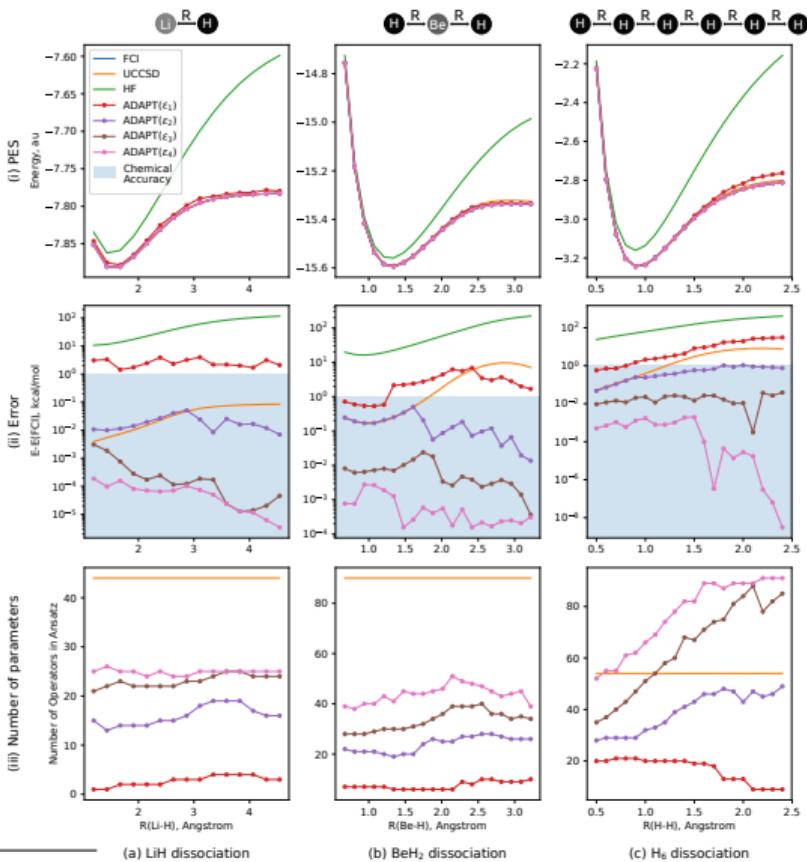
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ADAPT-VQE: Potential Energy Curves

21/21





Students

Shannon Houck
Vibin Abraham
Harper Grimsley
Robert Smith

Postdocs

Daniel Claudino
Oinam Meitei

Collaborators

Sophia Economou (VT Physics)
Ed Barnes (VT Physics)
Kyungwha Park (VT Physics)
Dave Pappas (NIST/UC Boulder)



Funding

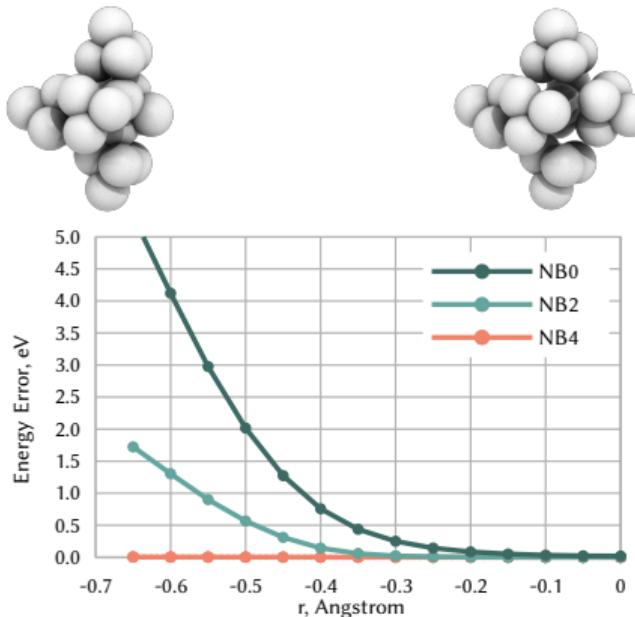
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NSF: 1839136



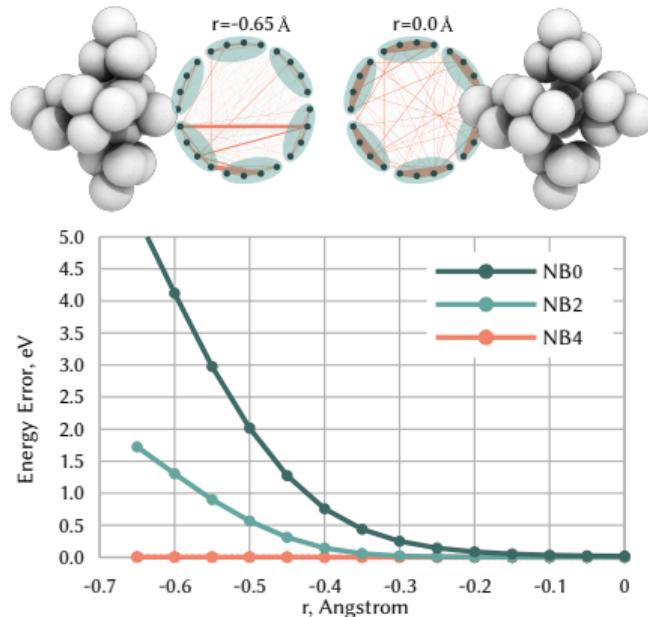
Postdoc and graduate student positions available!



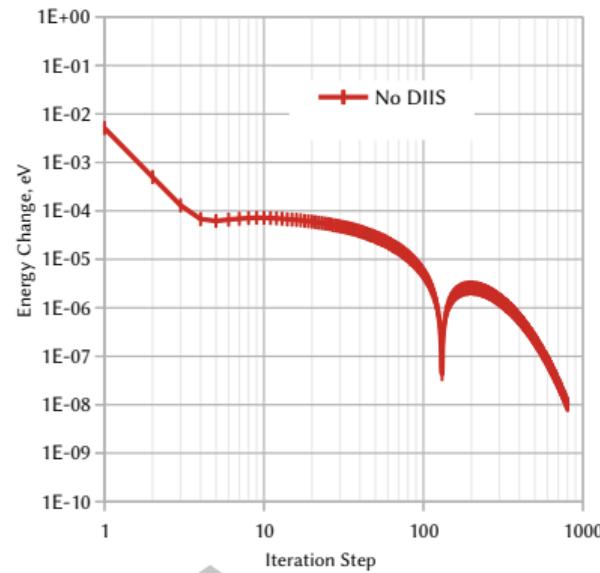
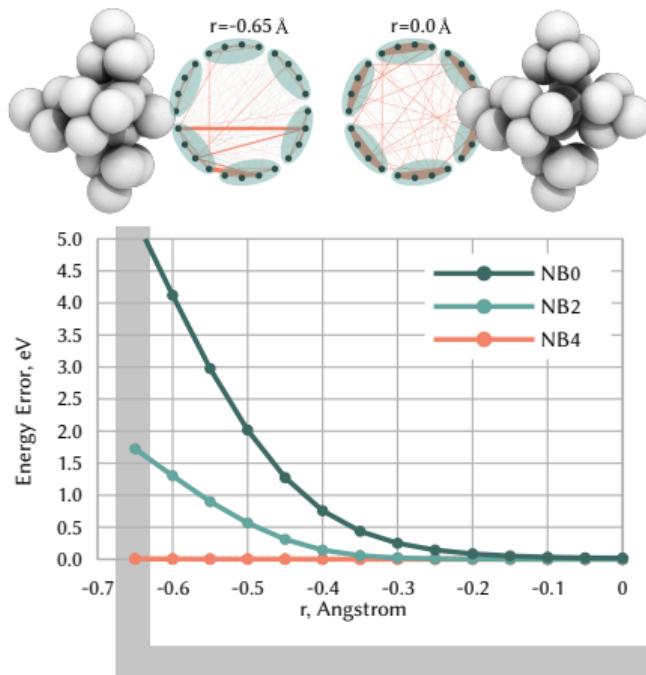
- Convergence of self-consistent “Tucker basis”



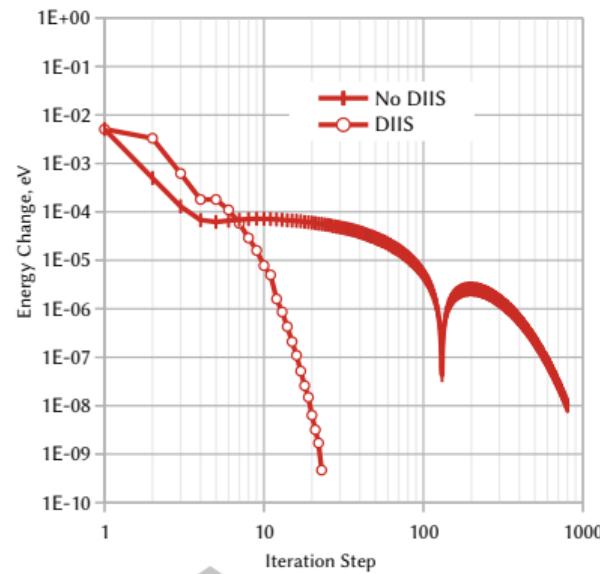
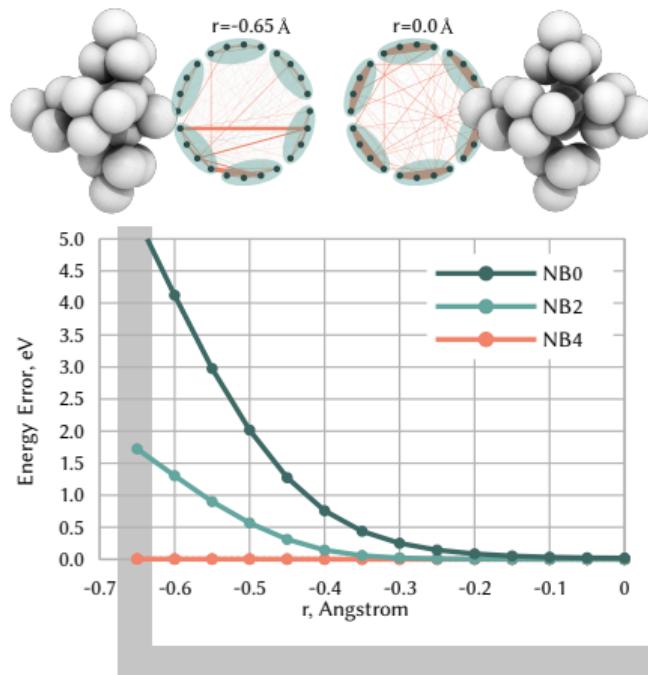
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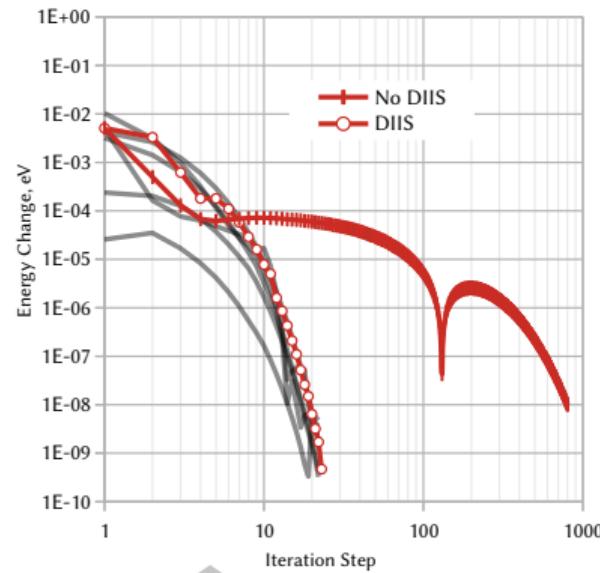
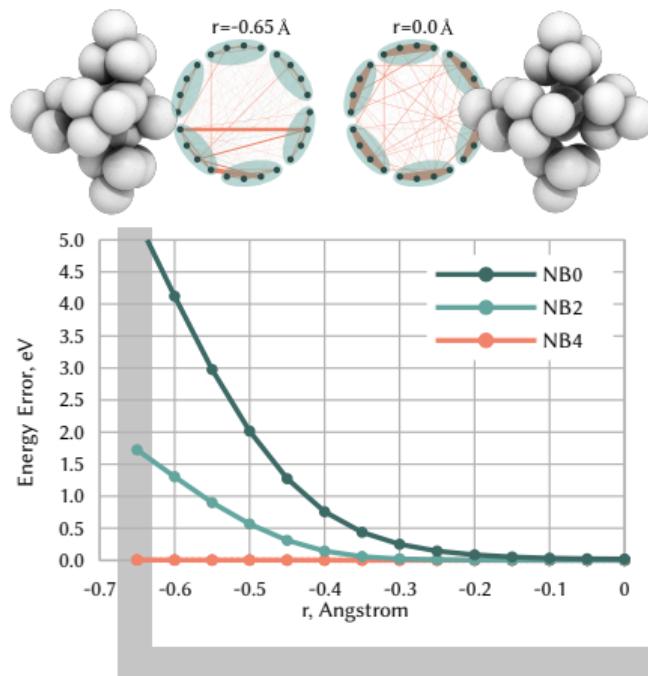
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- Perturbation theory via Löwdin Partitioning

