

Compound matrices for dispersion function

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1 P-SV problem settings in x-z plane

Equations satisfied by the P-wave potential scalar field ϕ and the SV-wave potential vector field $(0, \psi, 0)$.

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial z^2} + \frac{\omega^2}{\alpha^2} \phi = 0 \quad (1)$$

$$\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial z^2} + \frac{\omega^2}{\beta^2} \psi = 0 \quad (2)$$

The state variables of interest are expressed as following:

$$u_x = \frac{\partial \phi}{\partial x} - \frac{\partial \psi}{\partial z} \quad (3)$$

$$u_z = \frac{\partial \phi}{\partial z} + \frac{\partial \psi}{\partial x} \quad (4)$$

$$\sigma_{xz} = \rho \beta^2 \left(2 \frac{\partial^2 \phi}{\partial x \partial z} + \frac{\partial^2 \psi}{\partial x^2} - \frac{\partial^2 \psi}{\partial z^2} \right) \quad (5)$$

$$\sigma_{zz} = \rho \alpha^2 \frac{\partial^2 \phi}{\partial z^2} + \rho (\alpha^2 - 2\beta^2) \frac{\partial^2 \phi}{\partial x^2} + 2\rho \beta^2 \frac{\partial^2 \psi}{\partial x \partial z} \quad (6)$$

We consider layered medium with z-axis oriented positively downwards. We seek plane wave solution for the potential fields in a given homogeneous layer:

$$\phi = e^{ik(x-ct)} (A^+ e^{ikr_a z} + A^- e^{-ikr_a z}) \quad (7)$$

$$\psi = e^{ik(x-ct)} (B^+ e^{ikr_b z} + B^- e^{-ikr_b z}) \quad (8)$$

We use the following notation:

A^+, B^+ Amplitudes of downgoing waves

A^-, B^- Amplitudes of upgoing waves

ω circular frequency, common to all layers

k horizontal wavenumber, common to all layers

$c = \frac{\omega}{k}$ phase velocity, common to all layers

$r_a = \sqrt{\frac{c^2}{\alpha^2} - 1}$ can be complex number

$r_b = \sqrt{\frac{c^2}{\beta^2} - 1}$ can be complex number

$C_a = \cos(kr_a z)$ could become hyperbolic function for complex numbers

$C_b = \cos(kr_b z)$

$S_a = \sin(kr_a z)$

$S_b = \sin(kr_b z)$

$\mu = \rho\beta^2$

$\gamma = \frac{\beta^2}{c^2}$

$t = 2 - \frac{c^2}{\beta^2}$

After some algebraic manipulation, the state variables can be expressed as a function of the amplitudes vector:

$$\begin{bmatrix} -iku_x \\ -iku_z \\ \sigma_{xz} \\ \sigma_{zz} \end{bmatrix} = k^2 e^{ik(x-ct)} \cdot M \cdot P \cdot E \cdot \begin{bmatrix} A^+ \\ A^- \\ B^+ \\ B^- \end{bmatrix} \quad (9)$$

where

$$\begin{aligned}
M &= \begin{bmatrix} 1 & & & \\ & 1 & & \\ & & \mu & \\ & & & \mu \end{bmatrix} \\
P &= \begin{bmatrix} 1 & 1 & -r_b & r_b \\ r_a & -r_a & 1 & 1 \\ -2r_a & 2r_a & -t & -t \\ t & t & -2r_b & 2r_b \end{bmatrix} \\
E &= \begin{bmatrix} e^{ikr_a z} & & & \\ & e^{-ikr_a z} & & \\ & & e^{ikr_b z} & \\ & & & e^{-ikr_b z} \end{bmatrix}
\end{aligned}$$

We write:

$$D = Q \cdot E$$

$$Q = M \cdot P$$

$$Q^{-1} = P^{-1} \cdot M^{-1} = \begin{bmatrix} \gamma & -\frac{2\gamma-1}{2r_a} & -\frac{\gamma}{2\mu r_a} & -\frac{\gamma}{2\mu} \\ \gamma & \frac{2\gamma-1}{2r_a} & \frac{\gamma}{2\mu r_a} & -\frac{\gamma}{2\mu} \\ \frac{2\gamma-1}{2r_b} & \gamma & \frac{\gamma}{2\mu} & -\frac{\gamma}{2\mu r_b} \\ -\frac{2\gamma-1}{2r_b} & \gamma & \frac{\gamma}{2\mu} & \frac{\gamma}{2\mu r_b} \end{bmatrix}$$

and $T = Q \cdot E \cdot Q^{-1}$ with

$$\begin{aligned}
T_{11} &= \gamma(2C_a - tC_b) \\
T_{12} &= -i\gamma\left(\frac{t}{r_a}S_a + 2r_bS_b\right) & T_{31} &= -i\gamma\mu\left(4r_aS_a + \frac{t^2}{r_b}S_b\right) \\
T_{13} &= -i\frac{\gamma}{\mu}\left(\frac{1}{r_a}S_a + r_bS_b\right) & T_{32} &= 2\gamma\mu t(C_a - C_b) \\
T_{14} &= -\frac{\gamma}{\mu}(C_a - C_b) & T_{33} &= T_{11} \\
T_{21} &= i\gamma\left(2r_aS_a + \frac{t}{r_b}S_b\right) & T_{34} &= T_{21} \\
T_{22} &= -\gamma(tC_a - 2C_b) & T_{41} &= T_{32} \\
T_{23} &= T_{14} & T_{42} &= -i\gamma\mu\left(\frac{t^2}{r_a}S_a + 4r_bS_b\right) \\
T_{24} &= -i\frac{\gamma}{\mu}\left(r_aS_a + \frac{1}{r_b}S_b\right) & T_{43} &= T_{12} \\
& & T_{44} &= T_{22}
\end{aligned}$$

The matrix D is the field matrix relating the amplitudes to the state variables. The matrix T is the propagator matrix relating the state variables at two different depths in the same layer. This matrix satisfies $T^{-1}(z) = T(-z)$ and $T(z_1 + z_2) = T(z_1) \cdot T(z_2)$.

We introduce the delta matrices, or compound matrices of order 2, which are formed from a given matrix by taking the different combinations of order-2 determinant. If X is a 4×4 matrix, its compound matrix \bar{X} is 6×6 . A matrix entry \bar{X}_{ij} is itself an order-2 determinant obtained from the original matrix by selecting 2 rows and 2 columns. The order of pairs (of rows or columns) is $\{(12), (13), (14), (23), (24), (34)\}$.

With this we obtain the compound matrices for the needed matrices E , Q , Q^{-1} and T .

$$\bar{E} = \begin{bmatrix} 1 & & & & & \\ & e^{ik(r_a+r_b)z} & & & & \\ & & e^{ik(r_a-r_b)z} & & & \\ & & & e^{-ik(r_a-r_b)z} & & \\ & & & & e^{-ik(r_a+r_b)z} & \\ & & & & & 1 \end{bmatrix}$$

$\overline{Q}_{11} = -2r_a$	$\overline{Q}_{31} = 0$	$\overline{Q}_{51} = 2\mu tr_a$
$\overline{Q}_{12} = 1 + r_a r_b$	$\overline{Q}_{32} = -2\mu r_b + \mu tr_b$	$\overline{Q}_{52} = -2\mu r_a r_b - \mu t$
$\overline{Q}_{13} = 1 - r_a r_b$	$\overline{Q}_{33} = -\overline{Q}_{32}$	$\overline{Q}_{53} = 2\mu r_a r_b - \mu t$
$\overline{Q}_{14} = \overline{Q}_{13}$	$\overline{Q}_{34} = \overline{Q}_{32}$	$\overline{Q}_{54} = \overline{Q}_{53}$
$\overline{Q}_{15} = \overline{Q}_{12}$	$\overline{Q}_{35} = -\overline{Q}_{32}$	$\overline{Q}_{55} = \overline{Q}_{52}$
$\overline{Q}_{16} = -2r_b$	$\overline{Q}_{36} = 0$	$\overline{Q}_{56} = 4\mu r_b$
$\overline{Q}_{21} = 4\mu r_a$	$\overline{Q}_{41} = 0$	$\overline{Q}_{61} = -4\mu^2 tr_a$
$\overline{Q}_{22} = -\mu t - 2\mu r_a r_b$	$\overline{Q}_{42} = -\mu tr_a + 2\mu r_a$	$\overline{Q}_{62} = 4\mu^2 r_a r_b + (\mu t)^2$
$\overline{Q}_{23} = -\mu t + 2\mu r_a r_b$	$\overline{Q}_{43} = \overline{Q}_{42}$	$\overline{Q}_{63} = -4\mu^2 r_a r_b + (\mu t)^2$
$\overline{Q}_{24} = \overline{Q}_{23}$	$\overline{Q}_{44} = -\overline{Q}_{42}$	$\overline{Q}_{64} = \overline{Q}_{63}$
$\overline{Q}_{25} = \overline{Q}_{22}$	$\overline{Q}_{45} = -\overline{Q}_{42}$	$\overline{Q}_{65} = \overline{Q}_{62}$
$\overline{Q}_{26} = 2\mu tr_b$	$\overline{Q}_{46} = 0$	$\overline{Q}_{66} = -4\mu^2 tr_b$

$$\overline{Q}^{-1}_{11} = \frac{\gamma(2\gamma - 1)}{r_a}$$

$$\overline{Q}^{-1}_{12} = \frac{\gamma^2}{\mu r_a}$$

$$\overline{Q}^{-1}_{13} = 0$$

$$\overline{Q}^{-1}_{14} = 0$$

$$\overline{Q}^{-1}_{15} = \frac{\overline{Q}^{-1}_{11}}{2\mu}$$

$$\overline{Q}^{-1}_{16} = \frac{\overline{Q}^{-1}_{12}}{2\mu}$$

$$\overline{Q}^{-1}_{21} = \gamma^2 + \frac{(2\gamma - 1)^2}{4r_a r_b}$$

$$\overline{Q}^{-1}_{22} = \frac{\gamma^2}{2\mu} + \frac{\gamma(2\gamma - 1)}{4\mu r_a r_b}$$

$$\overline{Q}^{-1}_{23} = -\frac{\gamma}{4\mu r_b}$$

$$\overline{Q}^{-1}_{24} = \frac{\gamma}{4\mu r_a}$$

$$\overline{Q}^{-1}_{25} = \overline{Q}^{-1}_{22}$$

$$\overline{Q}^{-1}_{26} = \frac{\gamma^2}{4\mu^2} \left(1 + \frac{1}{r_a r_b}\right)$$

$$\overline{Q}^{-1}_{31} = 2\gamma^2 - \overline{Q}^{-1}_{21}$$

$$\overline{Q}^{-1}_{32} = \frac{\gamma^2}{\mu} - \overline{Q}^{-1}_{22}$$

$$\overline{Q}^{-1}_{33} = -\overline{Q}^{-1}_{23}$$

$$\overline{Q}^{-1}_{34} = \overline{Q}^{-1}_{24}$$

$$\overline{Q}^{-1}_{35} = \overline{Q}^{-1}_{32}$$

$$\overline{Q}^{-1}_{36} = \frac{\gamma^2}{4\mu^2} \left(1 - \frac{1}{r_a r_b}\right)$$

$$\overline{Q}^{-1}_{41} = \overline{Q}^{-1}_{31}$$

$$\overline{Q}^{-1}_{42} = \overline{Q}^{-1}_{32}$$

$$\overline{Q}^{-1}_{43} = \overline{Q}^{-1}_{23}$$

$$\overline{Q}^{-1}_{44} = -\overline{Q}^{-1}_{24}$$

$$\overline{Q}^{-1}_{45} = \overline{Q}^{-1}_{32}$$

$$\overline{Q}^{-1}_{46} = \overline{Q}^{-1}_{36}$$

$$\overline{Q}^{-1}_{51} = \overline{Q}^{-1}_{21}$$

$$\overline{Q}^{-1}_{52} = \overline{Q}^{-1}_{22}$$

$$\overline{Q}^{-1}_{53} = -\overline{Q}^{-1}_{23}$$

$$\overline{Q}^{-1}_{54} = -\overline{Q}^{-1}_{24}$$

$$\overline{Q}^{-1}_{55} = \overline{Q}^{-1}_{22}$$

$$\overline{Q}^{-1}_{56} = \overline{Q}^{-1}_{26}$$

$$\overline{Q}^{-1}_{61} = \frac{\gamma(2\gamma - 1)}{r_b}$$

$$\overline{Q}^{-1}_{62} = \frac{\overline{Q}^{-1}_{61}}{2\mu}$$

$$\overline{Q}^{-1}_{63} = 0$$

$$\overline{Q}^{-1}_{64} = 0$$

$$\overline{Q}^{-1}_{65} = \frac{\gamma^2}{\mu r_b}$$

$$\overline{Q}^{-1}_{66} = \frac{\overline{Q}^{-1}_{65}}{2\mu}$$

$$\begin{aligned}
\bar{T}_{11} &= -\gamma^2 \left(4t - (t^2 + 4)C_a C_b + (4r_a r_b + \frac{t^2}{r_a r_b})S_a S_b \right) \\
\bar{T}_{12} &= -\frac{\gamma^2}{\mu} \left(2 + t - (2 + t)C_a C_b + (2r_a r_b + \frac{t}{r_a r_b})S_a S_b \right) \\
\bar{T}_{13} &= -i\frac{\gamma}{\mu} \left(\frac{C_a S_b}{r_b} + r_a S_a C_b \right) \\
\bar{T}_{14} &= i\frac{\gamma}{\mu} \left(\frac{S_a C_b}{r_a} + r_b C_a S_b \right) \\
\bar{T}_{15} &= \bar{T}_{12} \\
\bar{T}_{16} &= -\frac{\gamma^2}{\mu^2} \left(2 + (\frac{1}{r_a r_b} + r_a r_b)S_a S_b - 2C_a C_b \right) \\
\bar{T}_{21} &= \gamma^2 \mu \left(2t(2 + t)(1 - C_a C_b) + (8r_a r_b + \frac{t^3}{r_a r_b})S_a S_b \right) \\
\bar{T}_{22} &= 1 + C_a C_b - \bar{T}_{11} \\
\bar{T}_{23} &= i\gamma \left(\frac{t}{r_b} C_a S_b + 2r_a S_a C_b \right) \\
\bar{T}_{24} &= -i\gamma \left(\frac{t}{r_a} S_a C_b + 2r_b C_a S_b \right) \\
\bar{T}_{25} &= \bar{T}_{22} - 1 \\
\bar{T}_{26} &= -\bar{T}_{12} \\
\bar{T}_{31} &= -i\gamma\mu \left(4r_b C_a S_b + \frac{t^2}{r_a} S_a C_b \right) \\
\bar{T}_{32} &= \bar{T}_{24} \\
\bar{T}_{33} &= C_a C_b \\
\bar{T}_{34} &= \frac{r_b}{r_a} S_a S_b \\
\bar{T}_{35} &= \bar{T}_{24} \\
\bar{T}_{36} &= -\bar{T}_{14} \\
\bar{T}_{41} &= i\gamma\mu \left(4r_a S_a C_b + \frac{t^2}{r_b} C_a S_b \right) \\
\bar{T}_{42} &= \bar{T}_{23} \\
\bar{T}_{43} &= \frac{r_a}{r_b} S_a S_b \\
\bar{T}_{44} &= \bar{T}_{33} \\
\bar{T}_{45} &= \bar{T}_{23} \\
\bar{T}_{46} &= -\bar{T}_{13} \\
\bar{T}_{51} &= \bar{T}_{21} \\
\bar{T}_{52} &= \bar{T}_{25} \\
\bar{T}_{53} &= \bar{T}_{23} \\
\bar{T}_{54} &= \bar{T}_{24} \\
\bar{T}_{55} &= \bar{T}_{22} \\
\bar{T}_{56} &= -\bar{T}_{12} \\
\bar{T}_{61} &= -\gamma^2 \mu^2 \left(8t^2(1 - C_a C_b) + (16r_a r_b + \frac{t^4}{r_a r_b})S_a S_b \right) \\
\bar{T}_{62} &= -\bar{T}_{21} \\
\bar{T}_{63} &= -\bar{T}_{41} \\
\bar{T}_{64} &= -\bar{T}_{31} \\
\bar{T}_{65} &= -\bar{T}_{21} \\
\bar{T}_{66} &= \bar{T}_{11}
\end{aligned}$$