

PONTIFICIA UNIVERSIDAD CATÓLICA DE CHILE
DEPTO. DE ING. HIDRÁULICA Y AMBIENTAL
ICH1104. MECÁNICA DE FLUIDOS

PRIMER SEMESTRE DE 2012
Jueves 21 de junio de 2012.

INTERROGACIÓN N°3
Sin Apuntes. Tiempo total: 2:00 hrs.

NOMBRE..... PAUTA

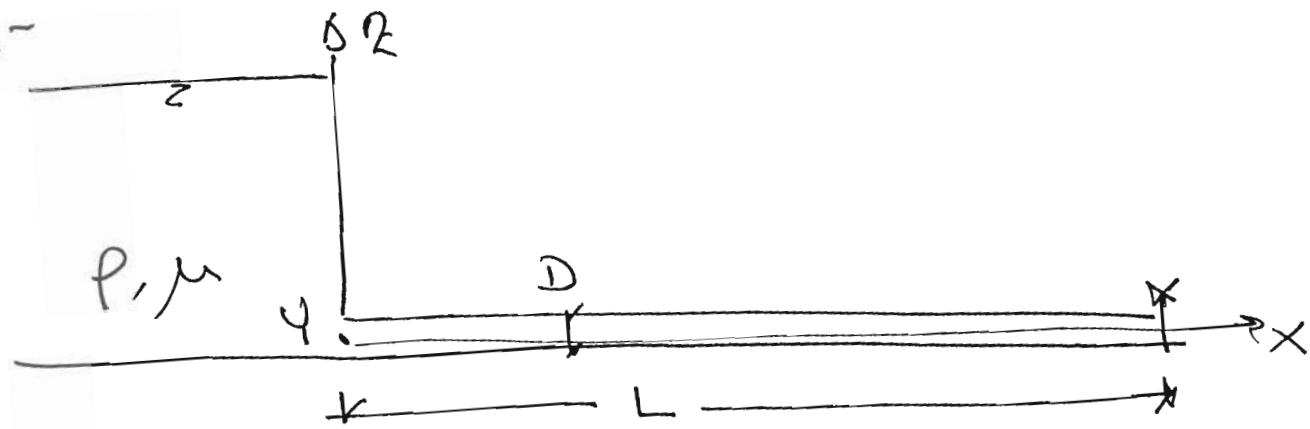
INSTRUCCIONES

Coloque su nombre en esta página. **No separe las hojas de este cuadernillo.**

Entregue en este cuadernillo sus respuestas en forma ordenada. Ocupe una hoja nueva para cada problema e indique claramente el problema que está contestando. Escriba con tinta o lápiz pasta. Explique su solución con las frases necesarias para poder seguirla. Ud. es responsable de que se entienda.

Sólo debe entregar este cuadernillo al final de la interrogación para ser corregido, de manera que debe asegurarse que sus respuestas sean completas, estén bien escritas y se pueda identificar claramente lo que contesta.

Problema	1	2	3	Final
Corrector				
Nota				



Q.-

$$\left. \begin{aligned} \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} &= f_{mx} - \frac{1}{\rho} \frac{\partial P}{\partial x} \\ \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} &= f_{my} - \frac{1}{\rho} \frac{\partial P}{\partial y} \\ \frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} &= f_{mz} - \frac{1}{\rho} \frac{\partial P}{\partial z} \\ \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} &= 0 \end{aligned} \right\}$$

Llamo $\vec{V} = u \hat{i}$ con $v = w = 0$, ademas

Fluido unidimensional $\frac{\partial v}{\partial x} = \frac{\partial v}{\partial y} = \frac{\partial v}{\partial z} = 0$

en la deriva $\frac{\partial w}{\partial x} = \frac{\partial w}{\partial y} = \frac{\partial w}{\partial z} = 0$

$$\frac{\partial v}{\partial t} = \frac{\partial w}{\partial t} = 0$$

La gravedad es $f_m = -g \hat{h} \Rightarrow f_{mx} = f_{my} = 0$

Si se incluye la ecuación de continuidad $\frac{\partial u}{\partial x} = 0$

$$\frac{\partial u}{\partial t} = -\frac{1}{\rho} \frac{\partial P}{\partial x}$$

$$0 = \frac{1}{\rho} \frac{\partial P}{\partial y}$$

Sistema de ecuaciones para este caso.

b). - La ecuación para P_2 :

$$\frac{\partial u}{\partial t} = -\frac{1}{P} \frac{\partial P}{\partial x}$$

Si la velocidad es constante $Q = \frac{Q_0}{T} k$

$$\text{Pero } Q = u \frac{\pi D^2}{4}$$

$$\text{Entonces } u = \frac{4Q}{\pi D^2} = \frac{4Q_0}{\pi T D^2} k$$

$$\therefore \frac{\partial u}{\partial t} = \frac{4Q_0}{\pi T D^2}$$

La ecuación queda:

$$\frac{1}{P} \frac{\partial P}{\partial x} = -\frac{4Q_0}{\pi T D^2}$$

$$\int_{P_0}^P \frac{\partial P}{\partial x} = -\frac{4Q_0}{\pi T D^2} \int_0^x$$

condición de borde

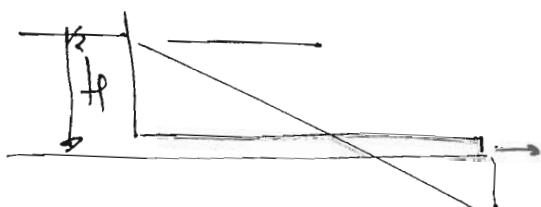
$$x=0 \Rightarrow P = gH$$

$$P - gH = -\frac{4Q_0 P}{\pi T D^2} x$$

$$\underline{\underline{P(x) = gH - \frac{4Q_0 P}{\pi T D^2} x}}$$

En particular en la sección $x=L$

$$\underline{\underline{P_L = gH - \frac{4Q_0 P}{\pi T D^2} L}}$$



L.-



base de Euler en la linea de corriente:

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = - \frac{1}{\rho} \frac{\partial P}{\partial x} \quad t > t_0$$

$$-\frac{\partial u}{\partial t} = \frac{1}{\rho} \frac{\partial P}{\partial x} + u \frac{\partial u}{\partial x}$$

$$-\int_1^2 \frac{\partial u}{\partial t} dx = \frac{1}{\rho} \int_1^2 \partial P + \int_1^2 u \partial u$$

$$-\underbrace{\frac{\partial u}{\partial t} (x_2 - x_1)}_{L} = \frac{1}{\rho} (P_2 - P_1) + \left(\frac{u_2^2}{2} - \frac{u_1^2}{2} \right)$$

$u_2 = u$

$$-L \frac{\partial u}{\partial t} = -\cancel{\frac{1}{\rho} SH} + \frac{u^2}{2}$$

$$-2L \frac{\partial u}{\partial t} = -2gH + u^2$$

$$\int_0^t \frac{\partial u}{2gH - u^2} = \int_{t_0}^t \frac{\partial u}{2L} = \frac{t - t_0}{2L}$$

↓ Esta integral no es obvia, miy pruebarronable.

$$\int \frac{du}{2gH - u^2} = \frac{1}{\sqrt{2gH}} \log \left| \frac{\sqrt{2gH} + u}{\sqrt{2gH} - u} \right| = \frac{t - t_0}{2L}$$

=====

d).- En régimen permanente, entre 1, 2:

$$Z_1 + \frac{P_1}{\gamma} + \frac{V_1^2}{2g} = Z_2 + \frac{P_2}{\gamma} + \frac{V_2^2}{2g}$$

donde $Z_1 = Z_2$

$$\frac{P_1}{\gamma} = H$$

$$V_1 = 0$$

$$P_2 = 0 \text{ (atm.)}$$

$$\therefore H = \frac{V_2^2}{2g} \Rightarrow V = \underline{\underline{\sqrt{2gH}}} = U_0$$



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2.- $f(z) = \ln z$ es un punto radial

a.- Poco o numeros: $f(z) = \frac{Q}{2\pi} \ln z$



Demolins $f(z) = iM \ln z$



La situación de una punto representarse con:

$$f(z) = \frac{Q}{2\pi} \ln r + iM \ln z$$

b.- $\phi = \phi_1 + \phi_2 ; \chi = \chi_1 + \chi_2$

Para ϕ_1, χ_1 :

$$\begin{aligned} f_1(z) &= \frac{Q}{2\pi} \ln z = \frac{Q}{2\pi} \ln r e^{i\theta} = \frac{Q}{2\pi} (\ln r + i\theta) \\ &= \frac{Q \ln r}{2\pi} + \frac{Q \theta}{2\pi} i \end{aligned}$$

$$\therefore \phi_1 = \frac{Q \ln r}{2\pi} ; \chi_1 = \frac{Q \theta}{2\pi}$$

Para ϕ_2, χ_2 :

$$\begin{aligned} f_2(z) &= iM \ln z = iM \ln r e^{i\theta} = iM (\ln r + i\theta) \\ &= iM \ln r - M\theta \end{aligned}$$

$$\therefore \phi_2 = -M\theta ; \chi_2 = M \ln r$$

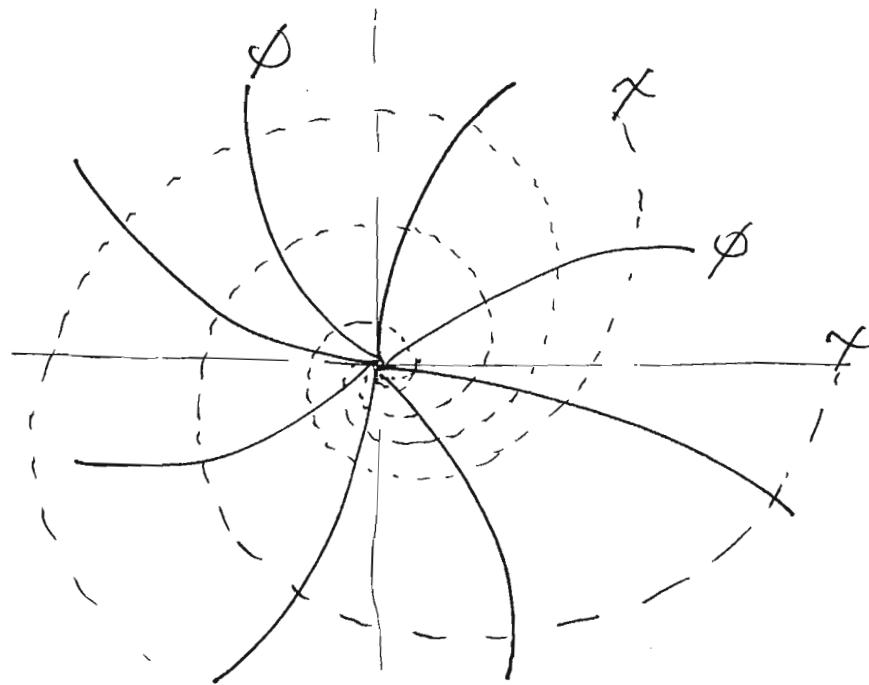
Entonces:

$$\phi = \frac{Q}{2\pi} \ln r - M\theta \quad \left. \right\}$$

$$\chi = \frac{Q\theta}{2\pi} + M \ln r \quad \left. \right\}$$

Achse für umbra limes von all tips:

$$\phi_x = \alpha \ln r + \beta \theta$$

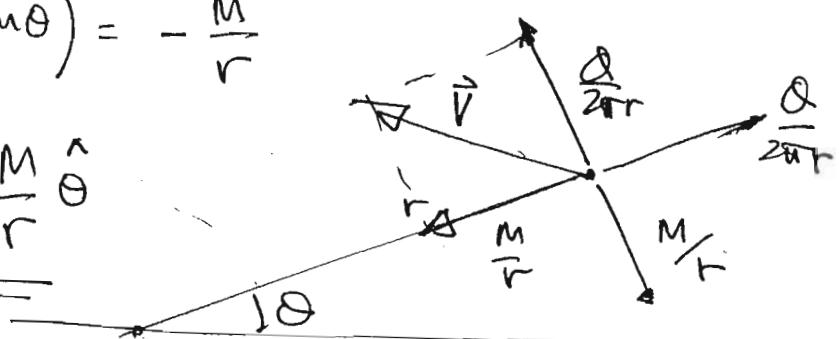


L.- La vektorfeld $\vec{V} = -\vec{\nabla}\phi = -\frac{\partial \phi}{\partial r} \hat{r} - \frac{1}{r} \frac{\partial \phi}{\partial \theta} \hat{\theta}$

$$\frac{\partial \phi}{\partial r} = \frac{\partial}{\partial r} \left(\frac{\partial}{2\pi} \ln r - M\theta \right) = \frac{\partial}{2\pi r} \star$$

$$\frac{1}{r} \frac{\partial \phi}{\partial \theta} = \frac{1}{r} \frac{\partial}{\partial \theta} \left(\frac{\partial}{2\pi} \ln r - M\theta \right) = -\frac{M}{r}$$

$$\therefore \vec{V} = -\frac{\partial}{2\pi r} \hat{r} + \frac{M}{r} \hat{\theta}$$



$$V_r = -\frac{\partial}{2\pi r} ; \quad V_\theta = \frac{M}{r}$$

d. Verifiziert:

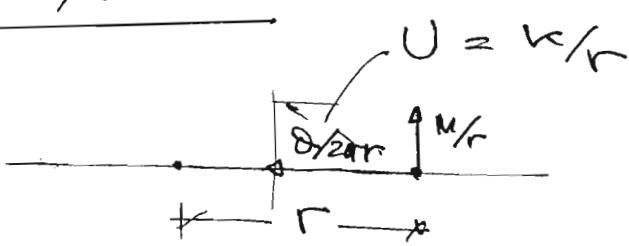
$$\begin{aligned}\vec{\nabla} \times \vec{V} &= \left(\frac{1}{r} \frac{\partial V_r}{\partial r} - \frac{1}{r} \frac{\partial V_\theta}{\partial \theta} \right) = \frac{1}{r} \frac{\partial \alpha}{\partial r} \hat{\phi} - \frac{1}{r} \frac{\partial}{\partial \theta} \left(-\frac{\alpha}{2\pi r} \right) \\ &= \frac{1}{r} \frac{\partial \alpha}{\partial r} \hat{m} + \frac{1}{r} \frac{\partial}{\partial \theta} \left(\alpha / 2\pi r \right) \\ &= 0\end{aligned}$$

$$\begin{aligned}\nabla \cdot \vec{V} &= \frac{1}{r} \frac{\partial}{\partial r} (r V_r) \hat{r} + \frac{1}{r} \frac{\partial V_\theta}{\partial \theta} \\ &= \frac{1}{r} \frac{\partial}{\partial r} \left(-\frac{\alpha \hat{\phi}}{2\pi r} \right) + \frac{1}{r} \frac{\partial}{\partial \theta} \left(\frac{m}{r} \right) \\ &= -\frac{1}{r} \frac{\partial}{\partial r} \left(\frac{\alpha}{2\pi} \right) + \frac{1}{r} \frac{\partial}{\partial \theta} \left(\frac{m}{r} \right) = 0\end{aligned}$$

$\therefore \vec{\nabla} \times \vec{V} = 0 ; \quad \nabla \cdot \vec{V} = 0$ anfangs
für $r=1 ; \theta = \pm \pi/2$

e.

$$\vec{V} = -\frac{\alpha}{2\pi r} \hat{r} + \frac{m}{r} \hat{\phi}$$



$$|V| = \sqrt{\frac{\alpha^2}{4\pi^2 r^2} + \frac{m^2}{r^2}} = \underbrace{\frac{1}{r} \sqrt{\frac{\alpha^2}{4\pi^2} + \frac{m^2}{r^2}}}_{k} = \frac{k}{r}$$

4.- En el infinito: $B_\infty = \frac{P_{atm}}{\gamma} = \frac{P_0}{\gamma} = 0$

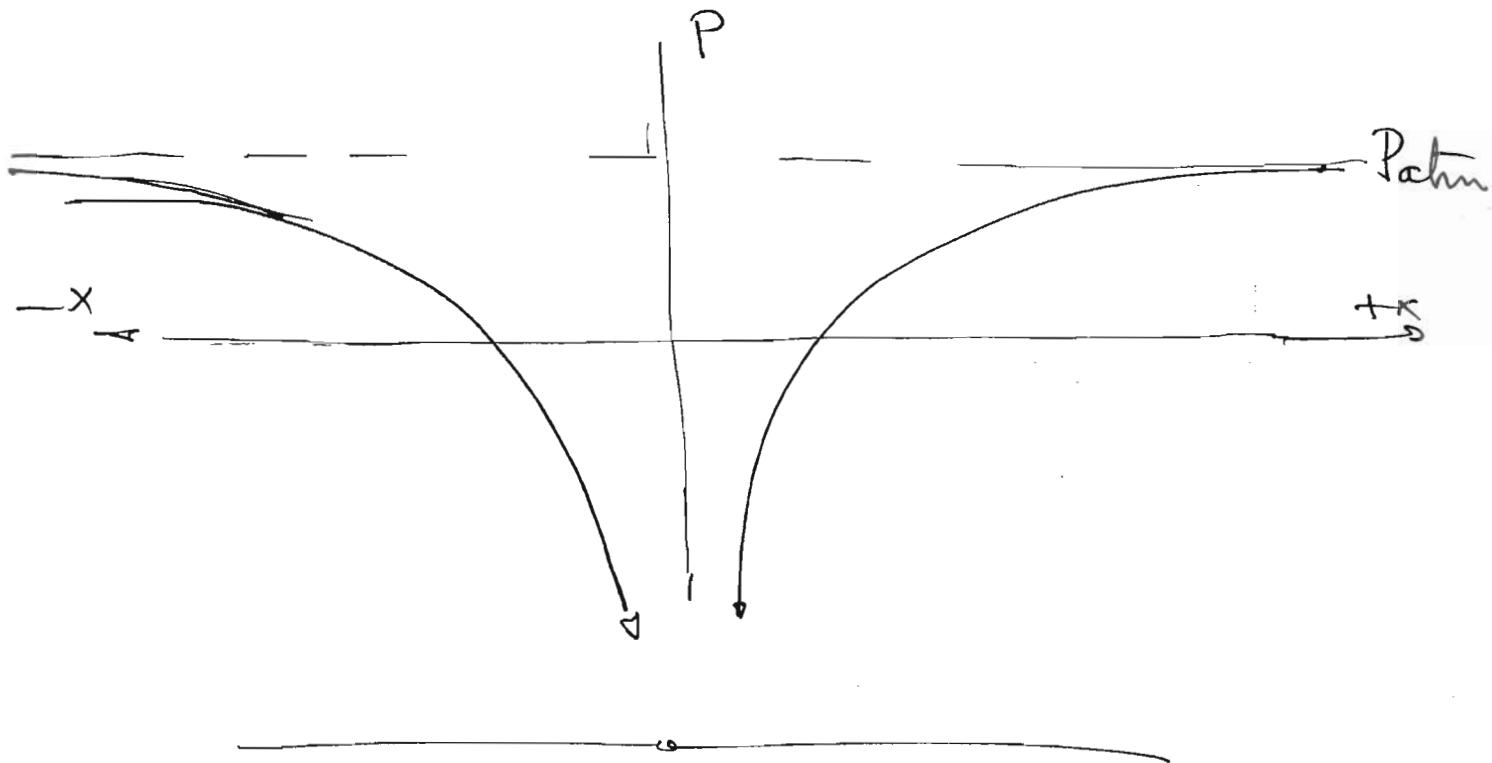
$$B_x = \frac{P}{\gamma} + \frac{V^2}{2g}$$

luego $B_\infty = B_x$

$$\frac{P}{\gamma} = \frac{P_0}{\gamma} - \frac{V^2}{2g}$$

$$\therefore P = P_0 - \frac{V^2 l}{2} \quad \text{donde } V^2 = \frac{K^2}{r^2}$$

$$P = P_0 - \frac{l}{2} \frac{K^2}{r^2} \quad \text{donde } K^2 = \frac{Q^2}{4\pi^2} + M^2$$



3.- a) Si $\vec{V} = w \hat{k}$ en $w = v = 0$
 Jergo unidimensional
 en m/s:

$$\text{en la ec. x: } \frac{\partial u}{\partial t} = u \frac{\partial u}{\partial x} = v \frac{\partial u}{\partial y} = w \frac{\partial u}{\partial z} = \frac{\partial^2 u}{\partial x^2} = \frac{\partial^2 u}{\partial y^2} = \frac{\partial^2 u}{\partial z^2} = 0$$

$$\text{“ “ “ y: } \frac{\partial v}{\partial t} = u \frac{\partial v}{\partial x} = v \frac{\partial v}{\partial y} = w \frac{\partial v}{\partial z} = \frac{\partial^2 v}{\partial x^2} = \frac{\partial^2 v}{\partial y^2} = \frac{\partial^2 v}{\partial z^2} = 0$$

$$\text{“ “ “ z: } u \frac{\partial w}{\partial x} = v \frac{\partial w}{\partial y} = 0$$

$$\text{Entimida: } \frac{\partial u}{\partial x} = \frac{\partial v}{\partial y} = 0$$

$$\text{Permanente: } \frac{\partial u}{\partial t} = \frac{\partial v}{\partial t} = \frac{\partial w}{\partial t} = 0$$

b) $f_m = -g \hat{h} \Rightarrow f_{mx} = f_{my} = 0$

$$\frac{\partial P}{\partial z} \approx \frac{\Delta P}{\Delta z} = -\frac{P_0 + P_{atm}}{h} \quad \text{pues } P_0 = \frac{F}{A}$$

$$\therefore \underline{\frac{\partial P}{\partial z} = -\frac{E}{A h}}$$

c) La ecuación de continuidad da $\frac{\partial w}{\partial z} = 0$

$$\text{Entonces } \frac{\partial^2 w}{\partial z^2} = 0$$

$$\text{Por simetría } \frac{\partial w}{\partial \theta} = 0 \Rightarrow \frac{\partial^2 w}{\partial \theta^2} = 0$$

Entonces las ecuaciones de NS se reducen a:

$$\left. \begin{aligned} 0 &= -\frac{1}{\rho} \frac{\partial P}{\partial x} \\ 0 &= -\frac{1}{\rho} \frac{\partial P}{\partial y} \\ 0 &= -g + \frac{1}{\rho} \frac{\partial P}{\partial z} + \frac{\mu}{\rho} \left(\frac{\partial^2 w}{\partial x^2} \right) \end{aligned} \right\}$$

lors $\frac{\partial P}{\partial z} = \frac{F}{A_h}$ el minima en:

$$\left. \begin{aligned} \frac{\partial P}{\partial x} &= \frac{\partial P}{\partial y} = 0 & (P = \text{constante en planes horizontales}) \\ \frac{\mu}{\rho} \frac{\partial^2 w}{\partial x^2} &= g + \frac{F}{\rho A_h} \end{aligned} \right\}$$

d) Ilamando $\kappa = \rho g - \frac{F}{\rho A_h} \rho = \rho g - \frac{F}{A_h}$

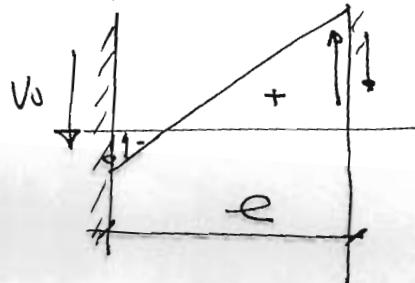
la ecuacion es:

$$\mu \frac{\partial^2 w}{\partial x^2} = \kappa$$

para los en flujos laminar: $\zeta = \mu \frac{\partial w}{\partial x}$

$$\therefore \frac{\partial \zeta}{\partial x} = \kappa \Rightarrow \underline{\zeta = Kx + C_0}$$

C_0 para determinar definir.



$$e) \quad \bar{F} = \mu \frac{\partial w}{\partial x} = kx + c_0$$

$$\partial w = \left(\frac{k}{\mu} x + \frac{c_0}{\mu} \right) \partial x$$

$$w = \frac{kx^2}{2\mu} + \frac{c_0 x}{\mu} + c_1$$

en la pared fija: $x=0 \Rightarrow w=0$
 $\therefore c_1 = 0$

en el punto en boga $x=e \rightarrow w=-V_0$

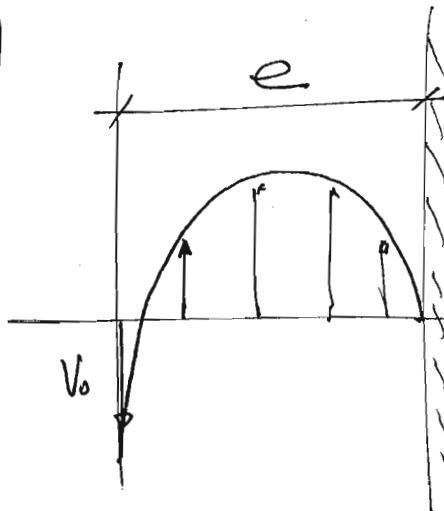
$$\therefore -V_0 = \frac{ke^2}{2\mu} + \frac{c_0 e}{\mu}$$

$$\begin{aligned} c_0 &= -\left(V_0 - \frac{ke^2}{2\mu}\right) \frac{\mu}{e} \\ &= -\frac{\mu V_0}{e} - \frac{ke}{2} \end{aligned}$$

Ari la rebuñada quide dada prr:

$$w = \frac{kx^2}{2\mu} + \frac{x}{\mu} \left(-\frac{\mu V_0}{e} - \frac{ke}{2} \right)$$

$$w = \frac{kx^2}{2\mu} - \frac{V_0}{e} x - \frac{ke}{2\mu} x$$



$$\begin{aligned}
 f: \quad \frac{Q}{4a} &= \int_0^e w dx = \int_0^e \left(\frac{kx^2}{2\mu} - \frac{v_0}{e} x - \frac{ke}{2\mu} x \right) dx \\
 &= \frac{k}{2\mu} \frac{x^3}{3} \Big|_0^e - \frac{v_0}{e} \frac{x^2}{2} \Big|_0^e - \frac{ke x^2}{2\mu} \Big|_0^e = \\
 &= \frac{ke^3}{6\mu} - \frac{v_0 e^2}{2} - \frac{ke^3}{4\mu} = \\
 &= \frac{ke^3}{\mu} \left(\frac{1}{6} - \frac{1}{4} \right) - \frac{v_0 e}{2} = \frac{ke^3}{12\mu} - \frac{v_0 e}{2} \\
 \therefore Q &= \underline{\left(\frac{ke^3}{12\mu} - \frac{v_0 e}{2} \right) 4a}
 \end{aligned}$$

Dato en igual a $v_0 4a$ para lo que el fricción interior