

We begin by defining the first law of thermodynamics

$$\frac{dU}{dt} = Q + W \quad (1)$$

If we define the rate of work in a closed system and rearrange eqn 1 for Q

$$Q = \frac{dU}{dt} + P \frac{dV}{dt} \quad (2)$$

$$W = -P \frac{dV}{dt}$$

Now redefine the internal energy change as...

$$\frac{dU}{dt} = \frac{\delta U}{\delta T} \frac{dT}{dt} + \frac{\delta U}{\delta V} \frac{dV}{dt} \quad (3)$$

Now sub in eqn 3 into 2 and rearrange

$$Q = \frac{\delta U}{\delta T} \frac{dT}{dt} + \frac{\delta U}{\delta V} \frac{dV}{dt} + P \frac{dV}{dt}$$

$$= C_v \frac{dT}{dt} + \left( \frac{\delta U}{\delta V} + P \right) \frac{dV}{dt} \quad (4)$$

$$C_v = \frac{\delta U}{\delta T} \quad (5)$$

Now in an adiabatic, constant volume process, internal energy is conserved. Additionally, under the adiabatic assumption no heat is added to the system. Therefore eqn 4 will become

$$0 = C_v \frac{dT}{dt} + P \frac{dV}{dt} \quad (6)$$

$$\frac{\delta U}{\delta V} = 0$$

Rearranging eqn 6

$$C_v \frac{dT}{dt} = -P \frac{dV}{dt} \quad (7)$$

Take the derivative of the ideal gas law with respect to time and rearrange

$$PV = NRT \quad P \frac{dV}{dt} + V \frac{dP}{dt} = NR \frac{dT}{dt}$$

$$P \frac{dV}{dt} = -V \frac{dP}{dt} + NR \frac{dT}{dt}$$

$$-P \frac{dV}{dt} = V \frac{dP}{dt} - NR \frac{dT}{dt} \quad (8)$$

Also assume you only have 1 molecule (i.e., N=1) and sub eqn 8 into eqn 7 and rearrange.

$$C_v \frac{dT}{dt} = V \frac{dP}{dt} - R \frac{dT}{dt}$$

$$C_v \frac{dT}{dt} + R \frac{dT}{dt} = V \frac{dP}{dt}$$

$$(C_v + R) \frac{dT}{dt} = V \frac{dP}{dt}$$

$$C_p \frac{dT}{dt} = V \frac{dP}{dt} \quad (9)$$

Where  $C_p$  is the specific heat at constant pressure Now again use the ideal gas law rearranged for volume  $V = \frac{RT}{P}$  in eqn 9. Then rearrange

$$\begin{aligned} C_p \frac{dT}{dt} &= \frac{RT}{P} \frac{dP}{dt} \\ \frac{1}{T} \frac{dT}{dt} &= \frac{\kappa}{P} \frac{dP}{dt} \\ \kappa &= \frac{R}{C_p} \end{aligned} \tag{10}$$

Regardless of the change with time, we find the following relationship from eqn 10

$$\frac{1}{T} dT = \frac{\kappa}{P} dP \tag{11}$$

We can now integrate these two sides between two temperatures and their corresponding pressures...

$$\begin{aligned} \int_{T_o}^{T_f} \frac{1}{T} dT &= \int_{P_o}^{P_f} \frac{\kappa}{P} dP \\ \ln(T) \Big|_{T_o}^{T_f} &= \kappa \ln(P) \Big|_{P_o}^{P_f} \\ \ln\left(\frac{T_f}{T_o}\right) &= \kappa \ln\left(\frac{P_f}{P_o}\right) \\ \ln\left(\frac{T_f}{T_o}\right) &= \ln\left(\frac{P_f}{P_o}\right)^\kappa \\ \frac{T_f}{T_o} &= \left(\frac{P_f}{P_o}\right)^\kappa \\ T_f &= T_o \left(\frac{P_f}{P_o}\right)^\kappa \end{aligned} \tag{12}$$

If we define the initial temperature to be the temperature of an air parcel at some pressure and the final temperature to be the temperature when brought down to a reference pressure (such as the surface then we have the definition of potential temperature

$$\theta = T \left( \frac{P_{sfc}}{P} \right)^\kappa \tag{13}$$