#### Homework 10 Solution

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This is my GitHub link: Nana Ama's GitHub Link

#### 1 Introduction

In one-dimension(1D), the Schrodinger equation of a particle is expressed as

$$-\frac{\hbar^2}{2M}\frac{\partial^2 \psi}{\partial x^2} = i\hbar \frac{\partial \psi}{\partial t} \tag{1}$$

This particle, an electron, was then placed in a box with impenetrable walls in one dimension was solved in questions one and two using the Crank-Nicolson method and the spectral method. Boundary conditions of  $\psi(x=0)=\psi(x=L)=0$  were applied in both cases. The mass of the electron is  $M=9.109\times 10^{-31}kg$ , in abox of length  $L=1\times 10^{-8}m$ . At time t=0, the wavefunction of the electron is of the form

$$\psi(x,0) = \exp\left[-\frac{(x-x_0)^2}{2\sigma^2}\right] \exp(ikx)$$
 (2)

where  $x_0 = \frac{L}{2}, \sigma = 1 \times 10^{-10} m$ , and  $k = 5 \times 10^{10} m^{-1}$ 

#### 2 Crank-Nicolson Method

The Crank-Nicolson equations can be written in the form

$$A\psi(t+h) = B\psi(t) \tag{3}$$

with

$$\psi(t) = \begin{pmatrix} \psi(a,t) \\ \psi(2a,t) \\ \psi(3a,t) \\ \vdots \end{pmatrix}$$

$$(4)$$

where a is the spacing of the spacial grid points and h is the size of the time step and the matrices A and B are tridiagonal matrices

$$A = \begin{pmatrix} a_1 & a_2 & & & \\ a_2 & a_1 & a_2 & & \\ & a_2 & a_1 & a_2 & \\ & & a_2 & a_1 & \\ & & & \ddots & \end{pmatrix}, B = \begin{pmatrix} b_1 & b_2 & & \\ b_2 & b_1 & b_2 & & \\ & b_2 & b_1 & b_2 & \\ & & b_2 & b_1 & \\ & & & \ddots & \end{pmatrix}$$
(5)

and

$$a_1 = 1 + h \frac{i\hbar}{2ma^2}, a_2 - h \frac{i\hbar}{4ma^2}, b_1 = 1 - h \frac{i\hbar}{2ma^2}, b_2 = h \frac{i\hbar}{4ma^2}$$
 (6)

The vector  $\psi(t)$  was calculated using the initial wavefunction above Equation .. numerically by

$$v = B\psi \tag{7}$$

with N=1000 spatial slices and a=L/N for a single step. where

$$v_i = b_1 \psi_i + b_2 (\psi_{i+1} + \psi_{i-1}) \tag{8}$$

For the next step, the scipy, banded module was used to solve the linear system Ax = v, which gives the updated  $\psi$  It was then extended to perform repeated steps at a separation  $h = 1 \times 10^{-18} s$ . An animation of the solution was made at each time step.

### 3 Spectral Method

Th spectral method involves using a Fourier transform to solve for the k components. With the same 1D Schrödinger equation as in Equation ..., the potential solution to this equation is

$$\psi_k = \sin(\frac{\pi kx}{L} \exp(\frac{iEt}{\hbar})) \tag{9}$$

with

$$E = \frac{\pi^2 \hbar^2 k^2}{2ML^2} \tag{10}$$

with a full solution expressed as a linear combination of these individual solutions on grid points  $x_n = \frac{nL}{N}$ 

$$\psi(x_n, t) = \frac{1}{N} \sum_{k=1}^{N-1} \sin(\frac{\pi k n}{N}) \exp(i\frac{\pi^2 \hbar k^2}{2ML^2})$$
(11)

and  $b_k = \alpha_k + i\eta_k$  with  $\alpha_k$  being the real part and  $\eta_k$ , the imaginary part.

A programme was made to calculate these  $b_k$  coefficients with the same parameters as those in the Crank-Nicolson method. Discrete sine transformations were performed on each initial wavefunction real and imaginary array at each grid point for all  $k = 1, \dots, N-1$ . The real part of the wavefunction was taken and extended at an arbitrary time using the inverse discrete sine transformation. This programme was tested by making a graph of the wavefunction at time  $t = 10^{-16}s$  as seen in Figure ... The programme was also extended to make an animation.

## 4 Discussion of Results

These 2 methods produced the same results as seen in Figure ....

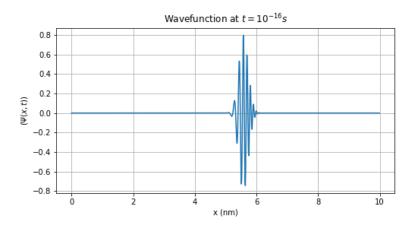


Figure 1: Waveform at  $t = 1 \times 10^{-16} s$ 

The test at time  $t = 1 \times 10^{-16} s$  yielded the plot in Figures 1 The progress of propagation of the wave can be seen in Figure 2 where it loses energy as it travels along the positive x direction and meets a boundary (that is the potential barrier), bounces back and travels along the negative x direction effectively changing its waveform.

# Propagation of Wave at Different Times

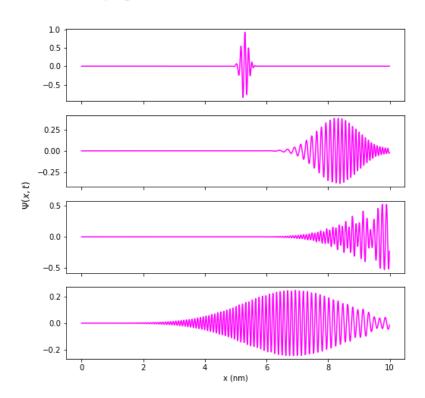


Figure 2: Crank-Nicolson Method