#### **VE 320 Fall 2021**

#### Introduction to Semiconductor Devices

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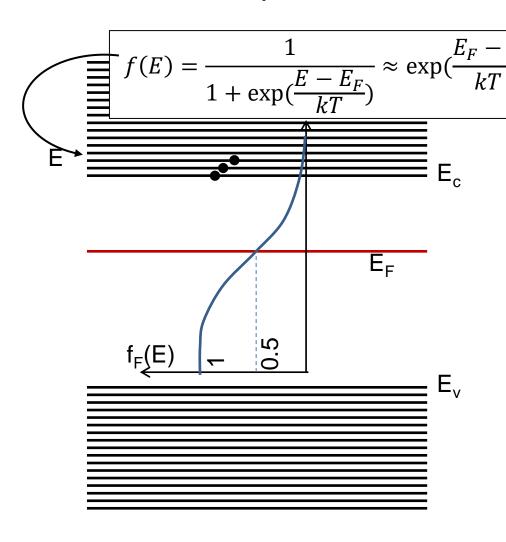
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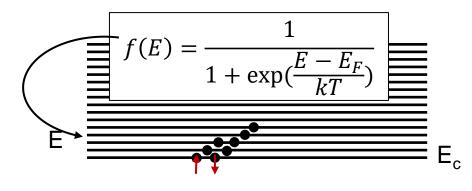
# Lecture 4

# The semiconductor in equilibrium (Chapter 4)



Non-degenerate Semiconductors:

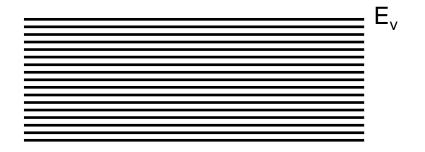
- Small number of impurity atoms
- Impurities introduce discrete, noninteracting donor energy states in n-type semiconductor
- Boltzmann approximation





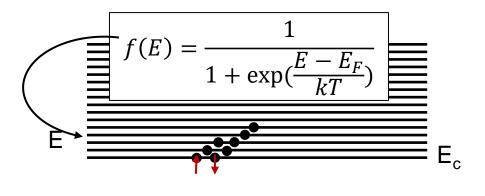
- High impurity concentration
- Donor electrons interact
- Split into a band of energies
- When n<sub>0</sub>>N<sub>c</sub>, then E<sub>F</sub>>E<sub>c</sub>:
   Degenerate n-type semiconductor

$$n_0 = N_c \exp\left[\frac{-(E_c - E_F)}{kT}\right]$$

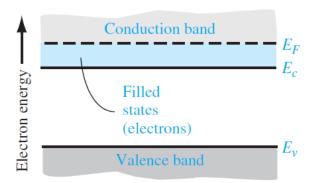


- When  $p_0 > N_v$ , the  $E_F < E_v$ :

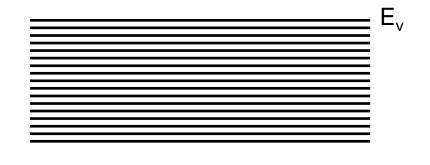
  Degenerate p-type semiconductor
- Fermi distribution, cannot use
   Boltzmann approximation

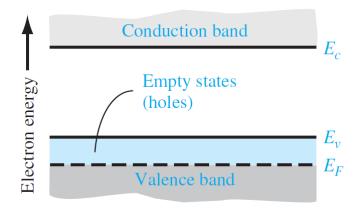


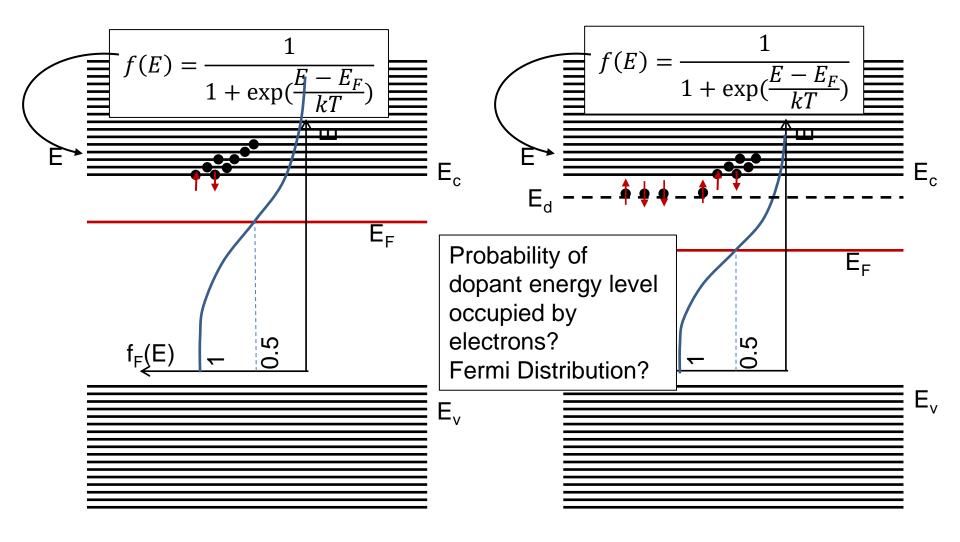
Degenerate n-type semiconductor: the states between  $E_{\rm F}$  and  $E_{\rm c}$  are mostly filled with electrons

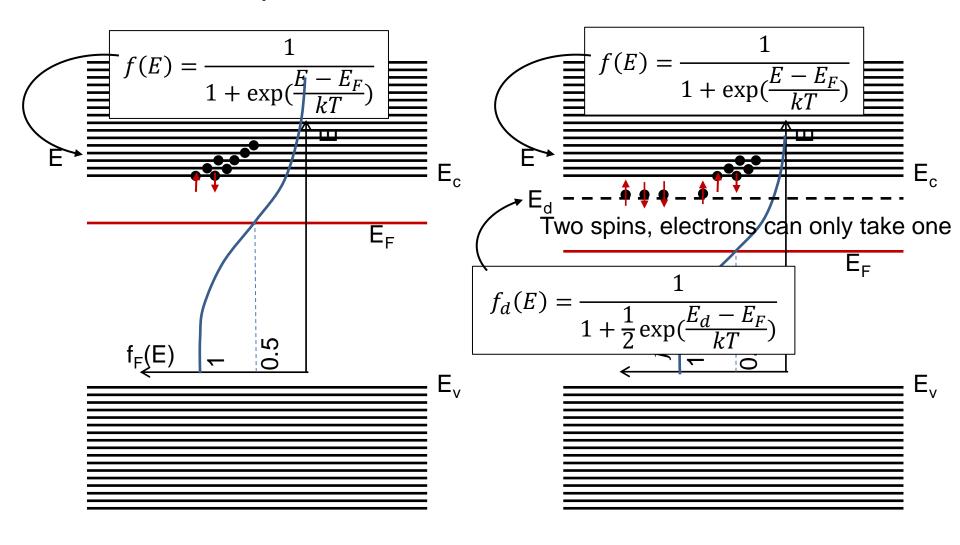


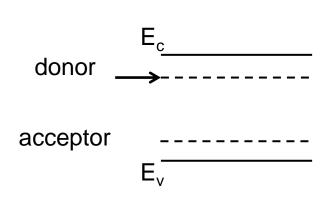
Degenerate p-type semiconductor: states between  $E_F$  and  $E_V$  are mostly empty











1 Probability of electrons occupying donor energy level

$$f_d(E) = \frac{1}{1 + \frac{1}{2} \exp(\frac{E_d - E_F}{kT})}$$

The density of electrons occupying the donor level:

$$1/2 = 1/g$$
, g is degeneracy factor

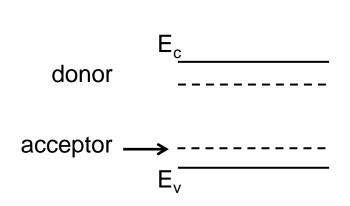
$$n_d = \frac{N_d}{1 + \frac{1}{2} \exp\left(\frac{E_d - E_F}{kT}\right)}$$

 $E_{\rm d}$ : energy of the donor level

$$n_d = N_d - N_d^+$$

 $N_{d}^{+}$  is the concentration of ionized donors--Important





1 Probability of electrons occupying acceptor energy level

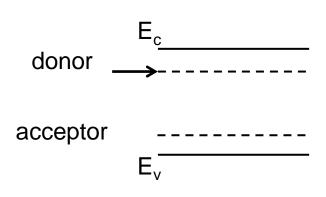
$$f_a(E) = \frac{1}{1 + \frac{1}{2} \exp(\frac{E_F - E_a}{kT})}$$

The density of holes occupying the acceptor level:

$$p_a = \frac{N_a}{1 + \frac{1}{g} \exp\left(\frac{E_F - E_a}{kT}\right)} = N_a - N_a^-$$

 $N_{\rm a}^{-}$  is the concentration of ionized acceptors--Important

The ground state degeneracy factor g is normally taken as 4 for the acceptor level in silicon and gallium arsenide because of the detailed band structure



1 Probability of electrons occupying dopant energy level

$$f_d(E) = \frac{1}{1 + \frac{1}{2} \exp(\frac{E_d - E_F}{kT})}$$

(2) Ionization rate

$$1 - f_d(E) = \frac{1}{1 + 2\exp(\frac{E_F - E_d}{kT})}$$

$$= \begin{cases} 1/3 & E_d = E_F \\ 1 & E_d \gg E_F + kT \\ 0 & E_d \ll E_F \end{cases}$$

Complete ionization

$$\begin{array}{c} E_{c} \\ \hline \\ \text{donor} \\ \hline \\ \text{acceptor} \\ \hline \\ E_{v} \\ \end{array}$$

Ratio of electrons that remains in dopant level to all electrons

$$(E_d - E_F) \gg kT$$

$$n_d \approx \frac{N_d}{\frac{1}{2} \exp\left(\frac{E_d - E_F}{kT}\right)} = 2N_d \exp\left[\frac{-(E_d - E_F)}{kT}\right]$$

Electrons in the conduction band:

$$n_0 = N_c \exp\left[\frac{-(E_c - E_F)}{kT}\right]$$

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$$\frac{n_d}{n_d+n_0} = \frac{2N_d \exp\left[\frac{-(E_d-E_F)}{kT}\right]}{2N_d \exp\left[\frac{-(E_d-E_F)}{kT}\right] + N_c \exp\left[\frac{-(E_c-E_F)}{kT}\right]} \quad \text{The donor states are essentially completely ionized for a typical do of 10$^{16} cm$^{-3}, almost all donor impurity atoms had donor donated an electron to$$

 $E_{\rm c}$ - $E_{\rm d}$ : ionization energy of the donor electrons

ionized for a typical doping of 10<sup>16</sup> cm<sup>-3</sup>, almost all donor impurity atoms have donated an electron to the conduction band.



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donor acceptor  $\mathsf{E}_\mathsf{v}$   $E_c$ - $E_d$ : ionization energy of the donor electrons

Ionization energy: the approximate energy required to elevate the donor electron into the conduction band

Bohr model:

Coulomb force = centripetal force

$$\frac{e^2}{4\pi \epsilon r_n^2} = \frac{m^* v^2}{r_n}$$

Angular momentum is also quantized  $m^* r_n v = n \hbar$ 

$$m^* r_n v = n \hbar$$

Radius of the orbit is quantized

$$r_n = \frac{n^2 \hbar^2 4\pi \epsilon}{m^* e^2}$$

Bohr radius

$$a_0 = \frac{4\pi\epsilon_0 \hbar^2}{m_0 e^2} = 0.53 \text{ Å} \qquad \frac{r_n}{a_0} = n^2 \epsilon_r \left(\frac{m_0}{m^*}\right)$$

*m*\* is the conductivity effective mass

acceptor

 $E_{c}$ - $E_{d}$ : ionization energy of the donor electrons

Total energy 
$$E = T + V$$

T: kinetic energy, V: potential energy

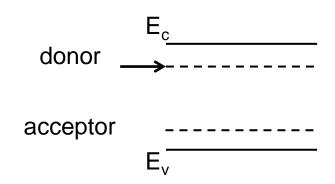
$$T = \frac{1}{2} m^* v^2 \qquad T = \frac{m^* e^4}{2(n\hbar)^2 (4\pi\epsilon)^2}$$

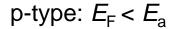
$$V = \frac{-e^2}{4\pi\epsilon r_n} = \frac{-m^* e^4}{(n\hbar)^2 (4\pi\epsilon)^2}$$

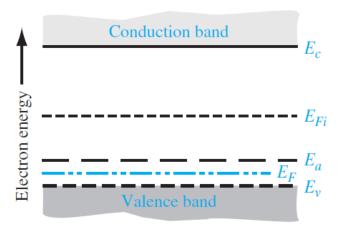
$$E = T + V = \frac{-m^* e^4}{2(n\hbar)^2 (4\pi\epsilon)^2}$$

Hydrogen atom in the lowest energy state:  $E_{\rm H}$  = -13.6 eV

Silicon ionization energy: is  $E_{\rm Si} = E_{\rm H} \times m^*/m_0/\epsilon_{\rm r}^2 = -25.8~{\rm meV} << E_{\rm g} = 1.12 {\rm eV}$ 







Freeze-out:

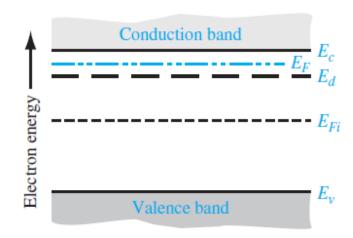
$$n_d = \frac{N_d}{1 + \frac{1}{2} \exp\left(\frac{E_d - E_F}{kT}\right)}$$

*T*=0K

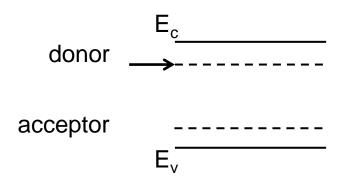
n-type semiconductor, each donor state must contain an electron,  $n_d = N_d$  or  $N_d^+ = 0$ 

$$\exp\left[(E_d - E_F)/kT\right] = 0$$

$$E_F > E_d$$



No electrons from the donor state are thermally elevated into the conduction band



Compensated semiconductor: contains both donor and acceptor impurity atoms in the same region

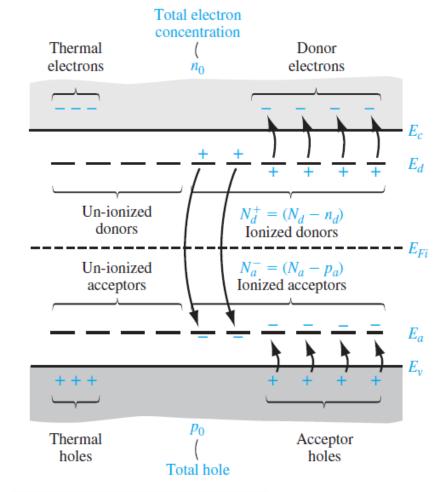
n-type compensated semiconductor:  $N_d > N_a$  p-type compensated semiconductor:  $N_a > N_d$   $N_a = N_d$ : completely compensated semiconductor, like intrinsic.

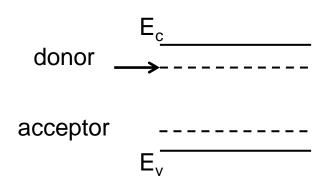
#### (3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

Compensated semiconductors at thermal equilibrium





Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 $n_d$ : concentration of electrons in the donor states  $p_a$ : concentration of holes in the acceptor states

Electron concentration in thermal equilibrium

Complete ionization:  $n_d=0$ ,  $p_a=0$ 

$$n_0 + N_a = N_d + p_0$$

$$n_0 + N_a = \frac{n_i^2}{n_0} + N_d$$

$$n_0^2 - (N_d - N_a)n_0 - n_i^2 = 0$$

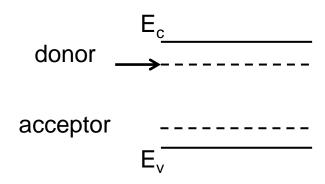
$$n_0 = \frac{(N_d - N_a)}{2} + \sqrt{\left(\frac{N_d - N_a}{2}\right)^2 + n_i^2}$$

Can calculate the electron concentration in an n-type semiconductor,  $N_d > N_a$ 

$$p_0 = \frac{n_i^2}{n_0}$$

 $N_{\rm a}$  can be 0: only n doping





Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

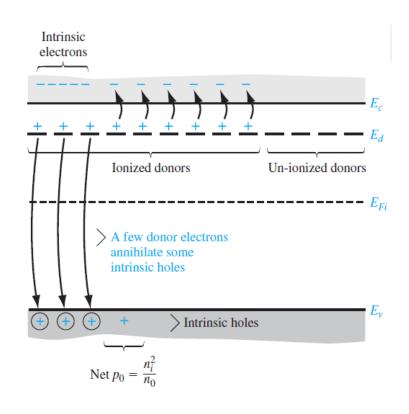
 $n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$ 

$$n_d$$
: concentration of electrons in the donor states  $p_a$ : concentration of holes in the acceptor states

#### Electron concentration in thermal equilibrium

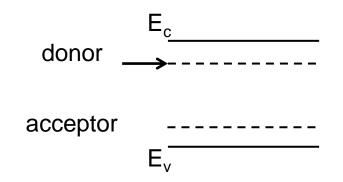
Complete ionization:  $n_d$ =0,  $p_a$ =0

$$n_0 > n_i$$
,  $p_0 < n_i$ 



Electron concentration in thermal equilibrium

High temperature: like intrinsic! – Usually bad for semiconductor devices



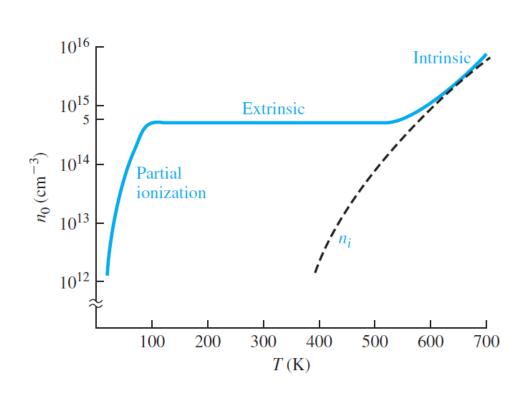
Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$



 $n_d$ : concentration of electrons in the donor states  $p_a$ : concentration of holes in the acceptor states

# Fermi level position:

Special case:  $N_a=0$ , Only n-type doping

Boltzmann approximation

$$n_0 = N_d^+ + p_0$$

$$n_0 p_0 = n_i^2$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)$$



 $n_i >> N_d^+ \Rightarrow T \text{ very high}$ 

$$n_0 = p_0 = n_i = \sqrt{N_c N_v} \exp(-\frac{E_g}{2kT})$$

$$\begin{array}{c} & & E_c \\ \hline \text{donor} & & \\ \hline \end{array}$$

Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 $n_d$ : concentration of electrons in the donor states  $p_a$ : concentration of holes in the acceptor states

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Compensated semiconductors

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$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 $n_d$ : concentration of electrons in the donor states  $p_a$ : concentration of holes in the acceptor states

$$n_0 = N_d^+ + p_0$$

$$n_0 p_0 = n_i^2$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)$$



$$n_i \leq N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 \approx N_d^+$$

(1) Assuming complete ionization:

$$n_0 \approx N_d$$

$$E_c = kT \ln \left( N_c \right)$$

$$E_c - E_F = kT \ln \left(\frac{N_c}{n_0}\right) = kT \ln \left(\frac{N_c}{N_d}\right)$$

 $N_{d}\uparrow$ ,  $E_{c}$ - $E_{F}\downarrow$ , Fermi level closer to the conduction band

# Fermi level position:

Special case:  $N_a=0$ , Only n-type doping

Boltzmann approximation

donor

 $\mathsf{E}_\mathsf{v}$ 

Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 $n_d$ : concentration of electrons in the donor states  $p_a$ : concentration of holes in the acceptor states

$$n_0 = N_d^+ + p_0$$

$$n_0 p_0 = n_i^2$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)$$



$$n_i \leq N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 \approx N_d^+$$

(1) Assuming complete ionization:

$$n_0 \approx N_d$$

$$E_c - E_F = kT \ln \left(\frac{N_c}{n_0}\right) = kT \ln \left(\frac{N_c}{N_d}\right)$$

Compensated semiconductor:

$$N_{\rm d} \rightarrow N_{\rm d} - N_{\rm a}$$

# Fermi level position:

Special case:  $N_a$ =0, Only n-type doping

Boltzmann approximation

$$n_0 = N_d^+ + p_0$$

$$n_0 p_0 = n_i^2$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)$$



$$n_i \leq N_d^+ \Rightarrow T \text{ not very high}$$

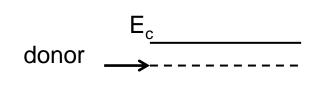
$$n_0 \approx N_d^+$$

(1) Assuming complete ionization:

$$n_0 \approx N_d$$

Also, 
$$n_0 = n_i \exp \left[ (E_F - E_{Fi})/kT \right]$$

$$E_F - E_{Fi} = kT \ln \left(\frac{n_0}{n_i}\right)$$



Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 $n_d$ : concentration of electrons in the donor states  $p_a$ : concentration of holes in the acceptor states

# Fermi level position:

Special case:  $N_d=0$ , Only p-type doping

$$n_i << N_a^+ \Rightarrow T \text{ not very high}$$

$$p_0 \approx N_a^+$$

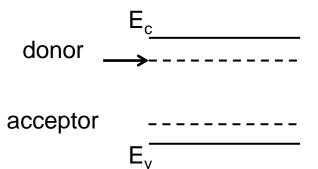
(1) Assuming complete ionization:

$$p_0 \approx N_a$$

Boltzmann approximation

$$p_0 = N_v \exp \left[ -(E_F - E_v)/kT \right]$$
$$E_F - E_v = kT \ln \left( \frac{N_v}{N_a} \right)$$

$$E_{Fi} - E_F = kT \ln \left(\frac{p_0}{n_i}\right)$$



Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

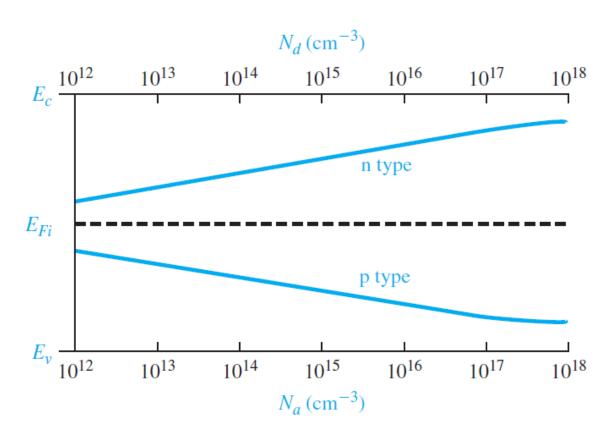
 $n_d$ : concentration of electrons in the donor states  $p_a$ : concentration of holes in the acceptor states

#### Fermi level position:

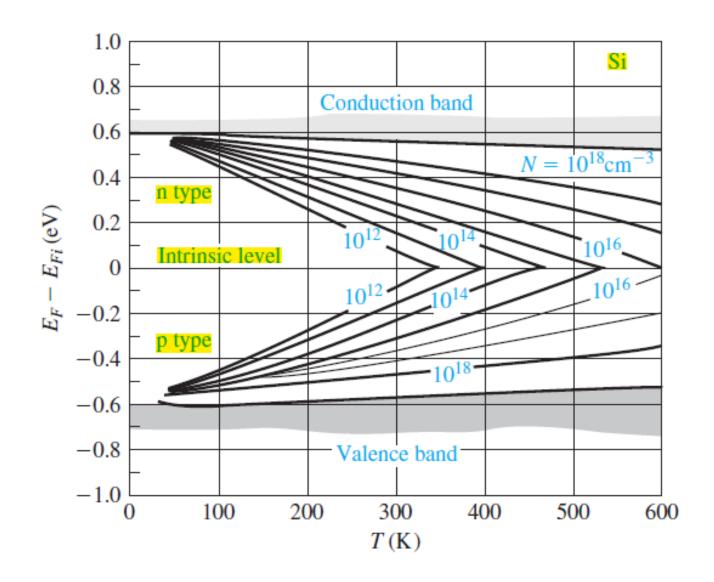
T not very high

#### Assuming complete ionization:

#### Boltzmann approximation



# Fermi level position:



# Fermi level position:

Special case:  $N_a$ =0, Only n-type doping

Boltzmann approximation

$$\begin{array}{c} & E_c \\ \hline \rightarrow & \end{array}$$

Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 $n_d$ : concentration of electrons in the donor states  $p_a$ : concentration of holes in the acceptor states

$$n_0 = N_d^+ + p_0 \qquad \qquad n_0 p_0 = n_i^2$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)$$



$$n_i << N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 \approx N_d^+$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right) = N_d^+ = \frac{N_d}{1 + 2\exp(\frac{E_F - E_d}{kT})}$$

$$n_i << N_d^+ \Rightarrow T \ not \ very \ high$$

$$n_0 = N_d^+$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right) = \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)}$$

$$n_i << N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 = N_d^+$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right) = \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)}$$

$$\exp\left(\frac{E_F - E_c}{kT}\right) = \frac{n_0}{N_c}$$

$$n_i << N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 = N_d^+$$

$$(n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)) = \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)}$$

$$\exp\left(\frac{E_F - E_c}{kT}\right) = \frac{n_0}{N_c}$$

$$\frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)} = \frac{N_d}{1 + 2\exp\left(\frac{E_C - E_d}{kT}\right)} \exp\left(\frac{E_C - E_d}{kT}\right)$$

$$n_i << N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 = N_d^+$$

$$(n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)) = \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)}$$

$$\exp\left(\frac{E_F - E_c}{kT}\right) = \frac{n_0}{N_c}$$

$$\frac{N_d}{1 + 2\exp(\frac{E_F - E_d}{kT})} = \frac{N_d}{1 + 2\exp(\frac{E_c - E_d}{kT})\exp(\frac{E_F - E_c}{kT})}$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right) = N_d^+ = \frac{N_d}{1 + 2\exp(\frac{E_A}{kT})\frac{n_0}{N_c}}$$

$$n_i << N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 = N_d^+$$

$$(n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)) = \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)}$$

$$\exp\left(\frac{E_F - E_c}{kT}\right) = \frac{n_0}{N_c}$$

$$\exp\left(\frac{E_F - E_c}{kT}\right) = \frac{n_0}{N_c} \qquad \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)} = \frac{N_d}{1 + 2\exp\left(\frac{E_c - E_d}{kT}\right)\exp\left(\frac{E_F - E_c}{kT}\right)}$$

$$2\exp\left(\frac{E_A}{kT}\right)n_0^2 + N_c n_0 - N_d N_c = 0$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right) = N_d^+ = \frac{N_d}{1 + 2\exp(\frac{E_A}{kT})\frac{n_0}{N_c}}$$

$$2\exp\left(\frac{E_A}{kT}\right)n_0^2 + N_c n_0 - N_d N_c = 0$$

$$n_0 = N_c \times \frac{-1 + \sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})}}{4\exp(\frac{E_A}{kT})}$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right) \Rightarrow \exp\left(\frac{E_F - E_c}{kT}\right) = \frac{-1 + \sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})}}{4\exp(\frac{E_A}{kT})}$$

$$E_F = E_c + kT ln(\frac{\sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})} - 1}{4\exp(\frac{E_A}{kT})})$$

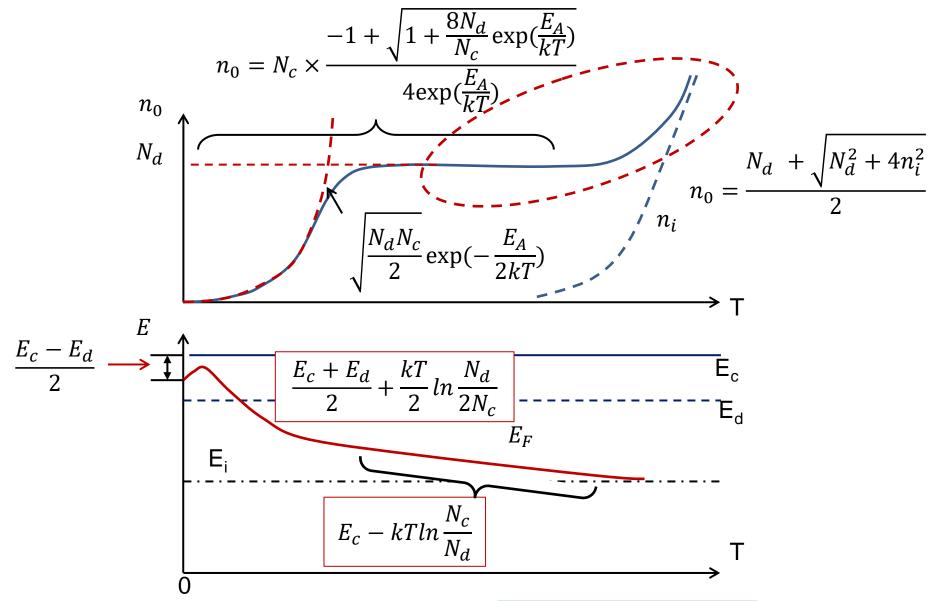
(2) Incomplete ionization, Boltzmann approximation, only n doping:

$$n_0 = N_c \times \frac{-1 + \sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})}}{4\exp(\frac{E_A}{kT})} = \begin{cases} \sqrt{\frac{N_d N_c}{2}} \exp(-\frac{E_A}{2kT}) & T \text{ small} \\ N_d & T \text{ big} \end{cases}$$

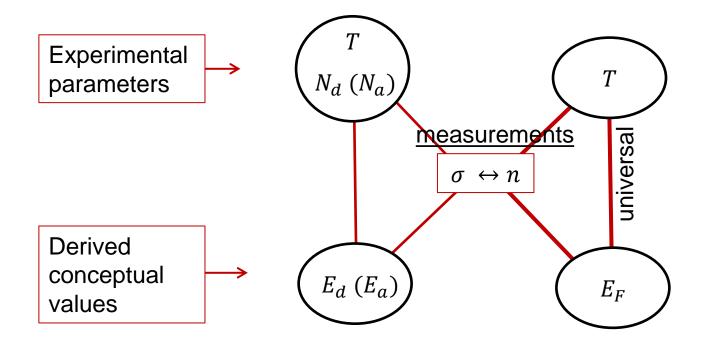
$$E_F = E_c + kT ln(\frac{\sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})} - 1}{4\exp(\frac{E_A}{kT})}) = \begin{cases} \frac{E_c + E_d}{2} + \frac{kT}{2} ln \frac{N_d}{2N_c} & T \text{ small} \\ E_c - kT ln \frac{N_c}{N_d} & T \text{ big} \end{cases}$$

$$n_0 = N_c \times \frac{-1 + \sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})}}{4\exp(\frac{E_A}{kT})} = \begin{cases} \sqrt{\frac{N_d N_c}{2}} \exp(-\frac{E_A}{2kT}) & T \text{ small} \\ N_d & T \text{ big} \end{cases}$$

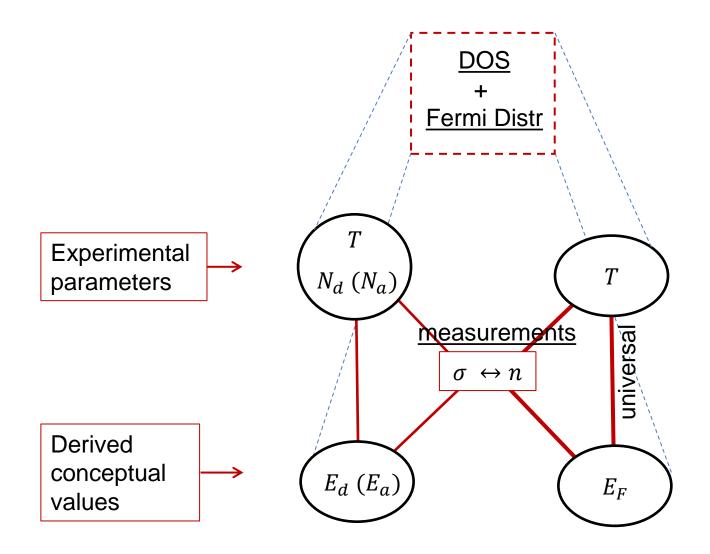
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# Summary of Lecture 4



# Summary of Lecture 4



# Summary of Lecture 4

#### Why is $E_F$ useful?

In thermal equilibrium, the Fermi energy level is a constant throughout a system

