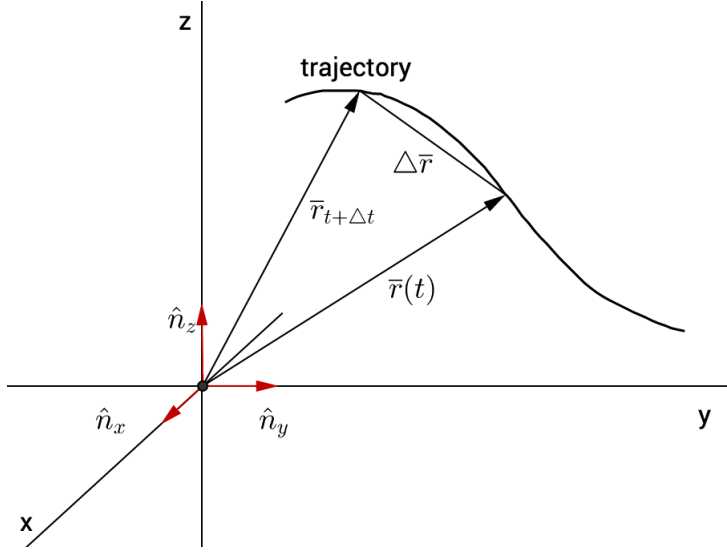


# 1 Cartesian coordinate system



$\Delta \bar{r}(t)$   
displacement over  
time interval  $(t, t + \Delta t)$

parametric equations  
of trajectory  $\begin{cases} x = x(t) \\ y = y(t) \\ z = z(t) \end{cases}$

Position

$$\bar{r}(t) = x(t)\hat{n}_x + y(t)\hat{n}_y + z(t)\hat{n}_z = (x(t), y(t), z(t)) \quad (1)$$

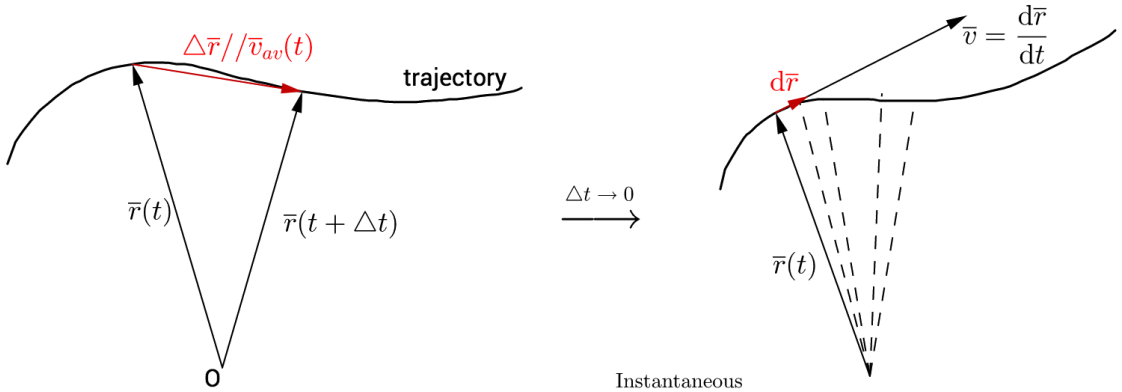
$\hat{n}_x, \hat{n}_y, \hat{n}_z$  - fixed unit vectors (time independent) i.e.  $\dot{\hat{n}}_x = \dot{\hat{n}}_y = \dot{\hat{n}}_z = 0$

Velocity

$$\bar{v}_{av} = \frac{\Delta \bar{r}}{\Delta t} \quad (\text{average}) \quad (2)$$

$$\begin{aligned} \bar{v}(t) &= \lim_{\Delta t \rightarrow 0} \frac{\Delta \bar{r}}{\Delta t} = \lim_{\Delta t \rightarrow 0} \frac{\bar{r}(t + \Delta t) - \bar{r}(t)}{\Delta t} \\ &= \lim_{\Delta t \rightarrow 0} \left[ \frac{x(t + \Delta t) - x(t)}{\Delta t} \hat{n}_x + \frac{y(t + \Delta t) - y(t)}{\Delta t} \hat{n}_y + \frac{z(t + \Delta t) - z(t)}{\Delta t} \hat{n}_z \right] \\ &= \underbrace{\dot{x}(t)}_{=v_x(t)} \hat{n}_x + \underbrace{\dot{y}(t)}_{=v_y(t)} \hat{n}_y + \underbrace{\dot{z}(t)}_{=v_z(t)} \hat{n}_z = (\dot{x}(t), \dot{y}(t), \dot{z}(t)) \quad (\text{instantaneous}) \end{aligned} \quad (3)$$

speed  $v(t) = |\bar{v}(t)| = \sqrt{[\dot{x}(t)]^2 + [\dot{y}(t)]^2 + [\dot{z}(t)]^2} \quad (4)$



[ Velocity vector  $\bar{v}$  is always tangent to the trajectory! ]

## Acceleration

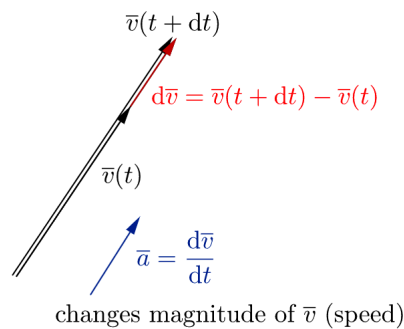
analogously  $\bar{a}_{av} = \frac{\Delta \bar{v}}{\Delta t}$  (average),  $\Delta \bar{v} = \bar{v}(t + \Delta t) - \bar{v}(t)$  (5)

$$\begin{aligned} \bar{a}(t) &= \lim_{\Delta t \rightarrow 0} \frac{\Delta \bar{v}}{\Delta t} = \dot{v}_x(t) \hat{n}_x + \dot{v}_y(t) \hat{n}_y + \dot{v}_z(t) \hat{n}_z \\ &= \underbrace{\ddot{x}(t)}_{=a_x(t)} \hat{n}_x + \underbrace{\ddot{y}(t)}_{=a_y(t)} \hat{n}_y + \underbrace{\ddot{z}(t)}_{=a_z(t)} \hat{n}_z \quad (\text{instantaneous}) \end{aligned} \quad (6)$$

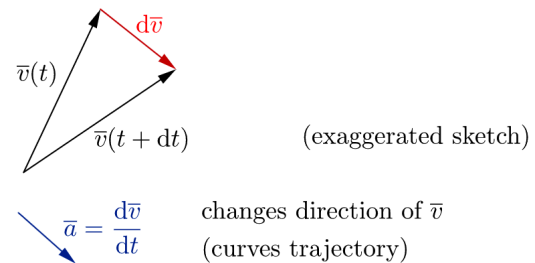
magnitude  $a(t) = |\bar{a}(t)| = \sqrt{[\ddot{x}(t)]^2 + [\ddot{y}(t)]^2 + [\ddot{z}(t)]^2}$  (7)

## 2 Tangential and normal components of acceleration

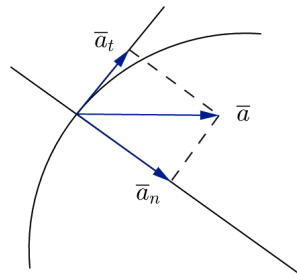
acceleration parallel to velocity  
(i.e. tangent to trajectory)



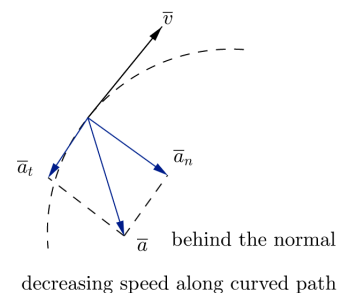
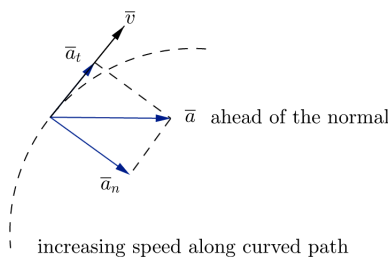
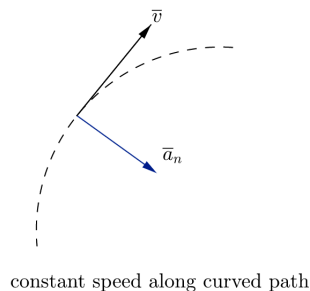
acceleration perpendicular to velocity  
(i.e. normal to trajectory)



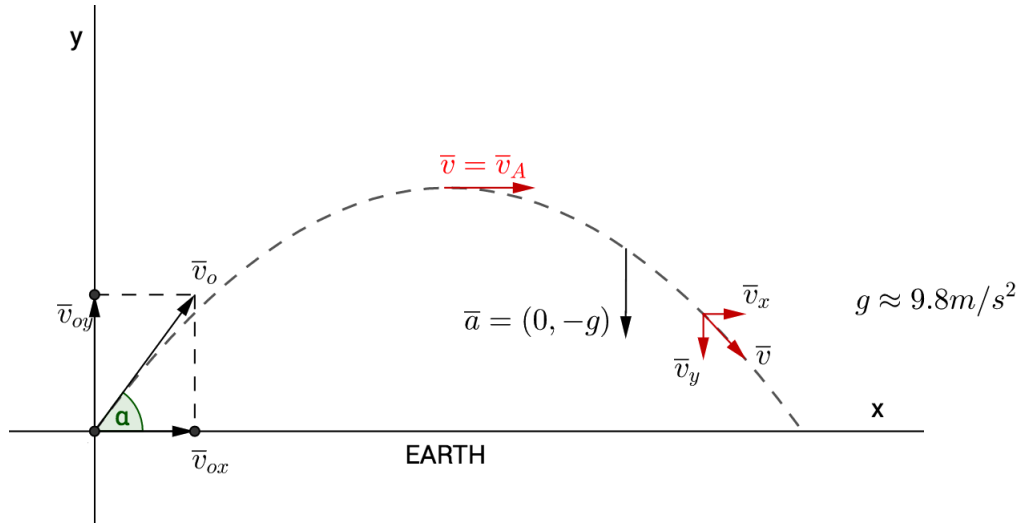
magnitude & direction of  $\bar{v}$  both change



special cases



Example (2D) projectile motion



Initial conditions (t=0)

$$\begin{cases} x(0) = 0 \\ y(0) = 0 \end{cases} \quad \begin{cases} v_x(0) = v_{0x} = v_0 \cos \alpha \\ v_y(0) = v_{0y} = v_0 \sin \alpha \end{cases}$$

Acceleration

$$\begin{aligned} a_x(t) &\equiv 0 \\ a_y(t) &= -g \end{aligned} \quad \text{that is} \quad \begin{aligned} \frac{dv_x}{dt} &= 0 \\ \frac{dv_y}{dt} &= -g \end{aligned}$$

Velocity

$$\int_{v_{0x}}^{v_x(t)} dv_x = \int_0^t 0 dt = 0 \quad \Rightarrow \quad v_x(t) = \text{const} = v_0 \cos \alpha \quad (8)$$

$$\int_{v_0 \sin \alpha}^{v_y(t)} dv_y = - \int_0^t g dt \quad \Rightarrow \quad v_y(t) = v_0 \sin \alpha - gt \quad (9)$$

Position

$$v_x(t) = \frac{dx}{dt} = v_0 \cos \alpha \quad \Rightarrow \quad \int_0^{x(t)} dx = \int_0^t v_0 \cos \alpha dt \quad \Rightarrow \quad x(t) = v_0 t \cos \alpha \quad (10)$$

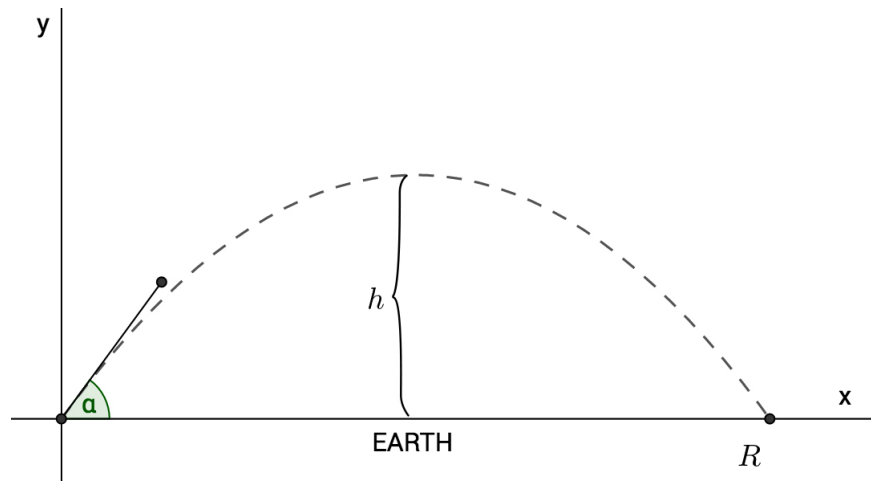
$$v_y(t) = \frac{dy}{dt} = v_0 \sin \alpha - gt \quad \Rightarrow \quad \int_0^{y(t)} dy = \int_0^t [v_0 \sin \alpha - gt] dt \quad \Rightarrow \quad y(t) = v_0 t \sin \alpha - \frac{1}{2}gt^2 \quad (11)$$

## Trajectory

$$\text{parametric form} \quad \begin{cases} x(t) = v_0 t \cos \alpha \\ y(t) = v_0 t \sin \alpha - \frac{1}{2} g t^2 \end{cases} \Rightarrow t = \frac{x}{v_0 \cos \alpha} \quad (12)$$

$$\Rightarrow \text{trajectory (implicit form)} \quad y(x) = x \tan \alpha - \frac{1}{2} \frac{g}{v_0^2 \cos^2 \alpha} x^2 \quad (\text{parabola}) \quad (13)$$

## Height and range



## Height

$$\text{At the highest point} \quad v_y(t_h) = 0 \Rightarrow t_h = \frac{v_0 \sin \alpha}{g} \quad (14)$$

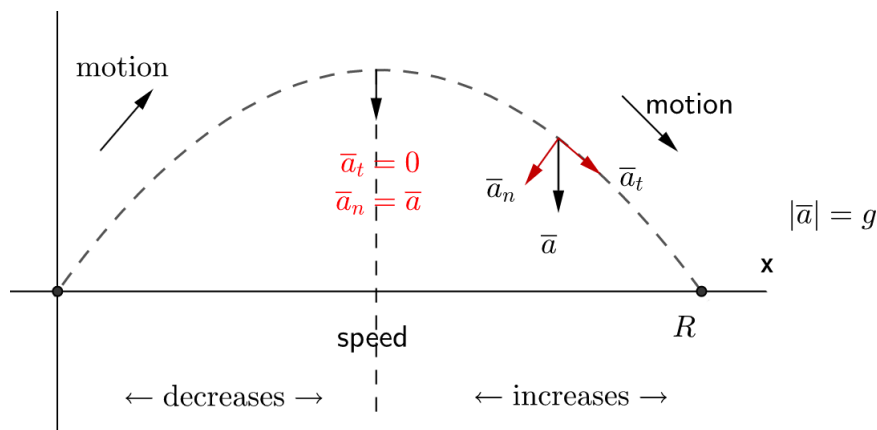
$$\text{and} \quad y(t_h) = \frac{v_0^2 \sin^2 \alpha}{2g} = h \quad \text{use (11)} \quad (15)$$

## Range

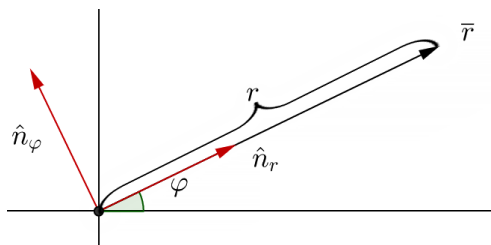
$$\begin{aligned} y(x_R) = 0 &\Rightarrow x_R \tan \alpha - \frac{1}{2} \frac{g}{v_0^2 \cos^2 \alpha} x_R^2 = 0 && \text{use (13)} \\ \Rightarrow x_R = 0 & \quad (\text{starting point}) \quad \text{or} \quad x_R = \frac{v_0^2 \sin 2\alpha}{g} && (16) \end{aligned}$$

What angle gives maximum (a) height?  $\alpha = 90^\circ$   
 (b) range?  $\alpha = 45^\circ$

## Acceleration in projectile motion



### 3 Polar coordinates (2D kinematics)



Position vector

$$\bar{r} = r\hat{n}_r$$

Trajectory

$$\text{parametric form} \quad \begin{cases} r = r(t) \\ \varphi = \varphi(t) \end{cases} \quad (17)$$

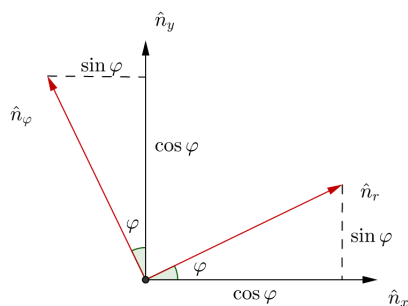
$$\text{or} \quad \text{implicit form} \quad r = r(\varphi) \text{ or } \varphi = \varphi(r) \quad (18)$$

Velocity

$$\bar{v} = \dot{\bar{r}} = \dot{r}\hat{n}_r + r\dot{\hat{n}}_r \quad (19)$$

Note that  $\dot{\hat{n}}_r$  is not zero, unlike in Cartesian coordinates, here  $\hat{n}_r, \hat{n}_\varphi$  are not fixed.

How to find  $\dot{\hat{n}}_r$  and  $\dot{\hat{n}}_\varphi$ ?



$$\begin{cases} \hat{n}_r = \cos \varphi \hat{n}_x + \sin \varphi \hat{n}_y \\ \hat{n}_\varphi = -\sin \varphi \hat{n}_x + \cos \varphi \hat{n}_y \end{cases} \quad (20)$$

Now using (20)

$$\dot{\hat{n}}_r = -\dot{\varphi} \sin \varphi \hat{n}_x + \dot{\varphi} \cos \varphi \hat{n}_y = \dot{\varphi} (-\sin \varphi \hat{n}_x + \cos \varphi \hat{n}_y) = \dot{\varphi} \hat{n}_\varphi \quad (21)$$

$$\dot{\hat{n}}_\varphi = -\dot{\varphi} \cos \varphi \hat{n}_x - \dot{\varphi} \sin \varphi \hat{n}_y = -\dot{\varphi} (\cos \varphi \hat{n}_x + \sin \varphi \hat{n}_y) = -\dot{\varphi} \hat{n}_r \quad (22)$$

So velocity

$$\boxed{\bar{v} = \dot{r} \hat{n}_r + r \dot{\hat{n}}_r = \underbrace{\dot{r} \hat{n}_r}_{\text{radial component}} + \underbrace{r \dot{\varphi} \hat{n}_\varphi}_{\text{transversal component}}} \quad (23)$$

$$\text{speed} \quad v = |\bar{v}| = \sqrt{(\dot{r})^2 + (r\dot{\varphi})^2} \quad (24)$$

Acceleration

$$\begin{aligned} \bar{a} = \dot{\bar{v}} &= \ddot{r} \hat{n}_r + \dot{r} \dot{\hat{n}}_r + \dot{r} \dot{\varphi} \hat{n}_\varphi + r \ddot{\varphi} \hat{n}_\varphi + r \dot{\varphi} \dot{\hat{n}}_\varphi \\ &= \ddot{r} \hat{n}_r + \dot{r} \dot{\varphi} \hat{n}_\varphi + \dot{r} \dot{\varphi} \hat{n}_\varphi + r \ddot{\varphi} \hat{n}_\varphi - r (\dot{\varphi})^2 \hat{n}_r \\ &= (\ddot{r} - r \dot{\varphi}^2) \hat{n}_r + (r \ddot{\varphi} + 2\dot{r} \dot{\varphi}) \hat{n}_\varphi \end{aligned}$$

$$\boxed{\bar{a} = \underbrace{(\ddot{r} - r \dot{\varphi}^2) \hat{n}_r}_{\text{radial component}} + \underbrace{(r \ddot{\varphi} + 2\dot{r} \dot{\varphi}) \hat{n}_\varphi}_{\text{transversal component}}} \quad (25)$$

CAUTION!

in general, radial  $\neq$  normal transversal  $\neq$  tangential



Example

(kinematics in the polar coordinates)

Circular motion:

$$r = R = \text{const} \Rightarrow \dot{r} = \ddot{r} = 0$$

$\varphi = \varphi(t)$  - in general: any function of time

(a) uniform circular motion (particle travels at constant speed, assume counter-clockwise)

Velocity

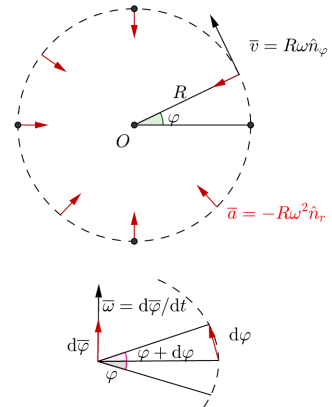
$$\bar{v} = \underbrace{\dot{r}}_{=0} \hat{n}_r + r \dot{\varphi} \hat{n}_\varphi = R \dot{\varphi} \hat{n}_\varphi \quad (26)$$

uniform motion, so  $|\bar{v}| = |v_\varphi| = v = \text{const}$

$$R \dot{\varphi} = v \Rightarrow \frac{d\varphi}{dt} = \frac{v}{R} \Rightarrow \int_0^{\varphi(t)} d\varphi = \int_0^t \frac{v}{R} dt$$

Hence

$$\varphi = \varphi(t) = \frac{v}{R} t = \omega t \quad (27)$$



Note  $\omega$  is angular velocity (here constant) Angular velocity is a vector, in general,

$$\bar{\omega} = \frac{d\bar{\varphi}}{dt}$$

Acceleration

$$\bar{a} = (\ddot{r} - r\dot{\varphi}^2)\hat{n}_r + (r\ddot{\varphi} + 2\dot{r}\dot{\varphi})\hat{n}_\varphi = -R\omega^2\hat{n}_r \quad (28)$$

Summary

$$\begin{aligned} \bar{v} &= R\omega\hat{n}_\varphi \\ \bar{a} &= -R\omega^2\hat{n}_\varphi \end{aligned}$$

Here:  $\hat{n}_\varphi \parallel$  tangential direction  
 $\hat{n}_r \parallel$  normal direction  
both  $|\bar{v}|$  and  $|\bar{a}|$  are constant in time

(b) non-uniform circular motion:  $\varphi = \varphi(t)$  - arbitrary function of time  
Now

$$\dot{\varphi} = \dot{\varphi}(t) = \omega(t)$$

angular velocity

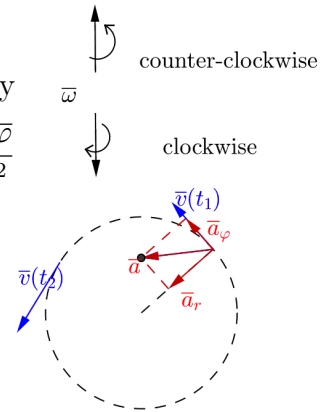
$$\ddot{\varphi} = \ddot{\varphi}(t) = \dot{\omega}(t) = \varepsilon(t) \quad \text{angular acceleration (in general)} \quad \bar{\varepsilon} = \frac{d^2\bar{\varphi}}{dt^2}$$

$$\bar{v} = R\omega(t)\hat{n}_\varphi$$

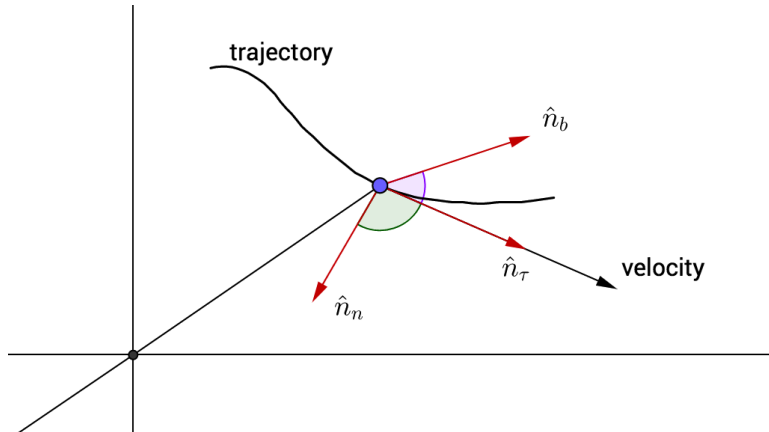
$$(29)$$

$$\bar{a} = \underbrace{-R\omega^2(t)\hat{n}_r}_{\text{curves the trajectory}} + \underbrace{R\varepsilon(t)\hat{n}_\varphi}_{\text{changes the magnitude of } \bar{v} \text{ (i.e. } v = R|\omega(t)| \neq \text{const})}$$

$$(30)$$



## 4 Natural coordinate system



unit vectors (versors)

$\hat{n}_\tau$  - tangent (along  $\bar{v}$ )

$\hat{n}_n$  - normal

$\hat{n}_b$  - binormal

$$\text{Velocity} \quad \bar{v}(t) = v\hat{n}_\tau \quad \text{or} \quad \hat{n}_\tau = \frac{\text{velocity(vector)} \quad \bar{v}}{\text{speed(scalar)} \quad v} = \frac{\dot{\vec{r}}}{|\dot{\vec{r}}|}$$

Note: Normal unit vector - perpendicular (orthogonal) to  $\hat{n}_\tau$  - many choices possible

! (in 3D)

Unique choice:

$$\hat{n}_n = \frac{\frac{d\hat{n}_\tau}{dt}}{\left| \frac{d\hat{n}_\tau}{dt} \right|} \quad \text{because } \frac{d\hat{n}_\tau}{dt} \text{ in general is not a unit}$$

vector.

Is  $\hat{n}_n \perp \hat{n}_\tau$ ? Yes!

$$\begin{aligned} \hat{n}_\tau \cdot \hat{n}_\tau = 1 & \xrightarrow[\text{with respect to } t]{\text{differentiate}} \frac{d\hat{n}_\tau}{dt} \cdot \hat{n}_\tau + \hat{n}_\tau \cdot \frac{d\hat{n}_\tau}{dt} = 0 \\ \frac{d\hat{n}_\tau}{dt} \cdot \hat{n}_\tau = 0 & \implies \frac{d\hat{n}_\tau}{dt} \perp \hat{n}_\tau \quad \text{and} \quad \hat{n}_n \perp \hat{n}_\tau \end{aligned}$$

The normal versor  $\hat{n}_n$  points along the radius of curvature.

The binormal unit vector determined by

$$\hat{n}_b = \hat{n}_\tau \times \hat{n}_n \quad (\text{right-handed system})$$

$\hat{n}_\tau$ ,  $\hat{n}_n$  and  $\hat{n}_b$  are the three unit vectors of the natural coordinate system, "sliding" along the particle's trajectory.

### Velocity

$$\bar{v} = v\hat{n}_\tau \quad (31)$$

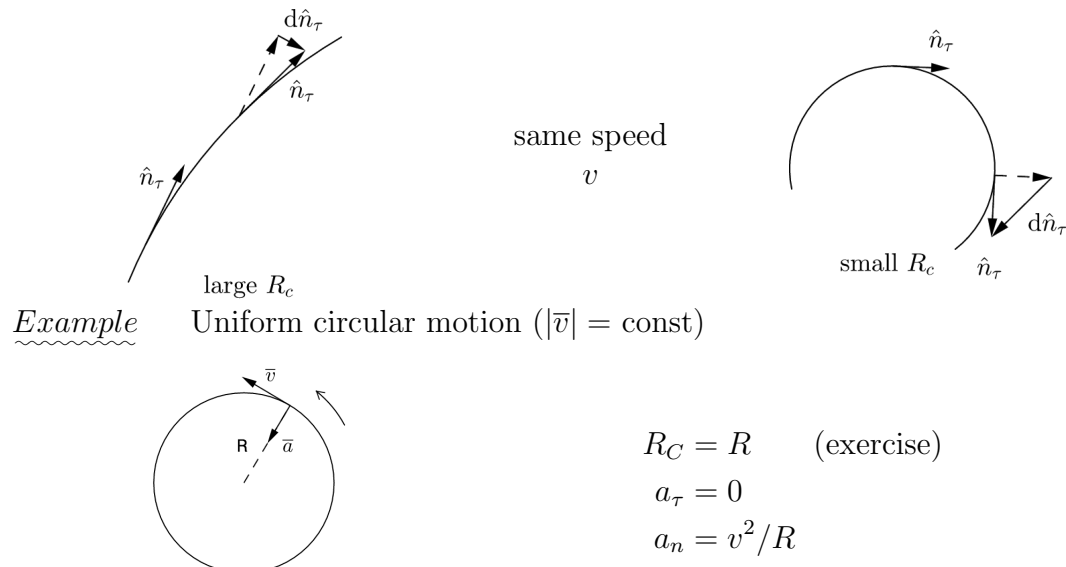
### Acceleration

$$\bar{a} = \dot{\bar{v}} = \dot{v}\hat{n}_\tau + v\dot{\hat{n}}_\tau \stackrel{\text{use definition of } \hat{n}_n}{=} \dot{v}\hat{n}_\tau + v \left| \frac{d\hat{n}_\tau}{dt} \right| \hat{n}_n$$

Define the (instantaneous) radius of trajectory's curvature  $R_c \stackrel{\text{def}}{=} \frac{v}{\left| \frac{d\hat{n}_\tau}{dt} \right|}$ . Then

$$\begin{aligned} \bar{a} = & \underbrace{\dot{v}\hat{n}_\tau}_{\text{tangential component } \bar{a}_\tau} + \underbrace{\frac{v^2}{R_c}\hat{n}_n}_{\text{normal component } \bar{a}_n} \\ & \text{mutually perpendicular } (|\bar{a}| = \sqrt{a_\tau^2 + a_n^2}) \end{aligned} \quad (32)$$

Note Interpretation of  $R_C$





Again:

$$a_\tau = \dot{v} \quad \text{changes the magnitude of } \bar{v}$$

$$a_n = v^2/R_C \quad \text{changes the direction of } \bar{v}$$

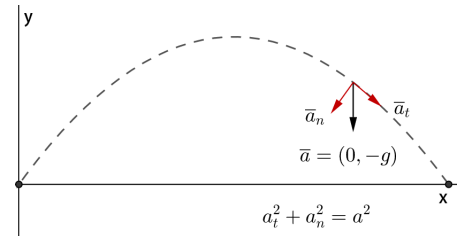
Example Projectile motion (also see hm2)

$$\bar{v} = (v_0 \cos \alpha) \hat{n}_x + (v_0 \sin \alpha - gt) \hat{n}_y$$

$$|\bar{v}| = v = \sqrt{[v_x(t)]^2 + [v_y(t)]^2}$$

and  $a_t = \dot{v}$

$$a_n = \sqrt{g^2 - (\dot{v})^2}$$



## 5 Concluding remarks

### 5.1 average speed vs average time

distance traveled over  
time interval  $(t_1, t_2)$

$$\text{average speed} = \frac{\int_{t_1}^{t_2} |\bar{v}(t)| dt}{t_2 - t_1}$$

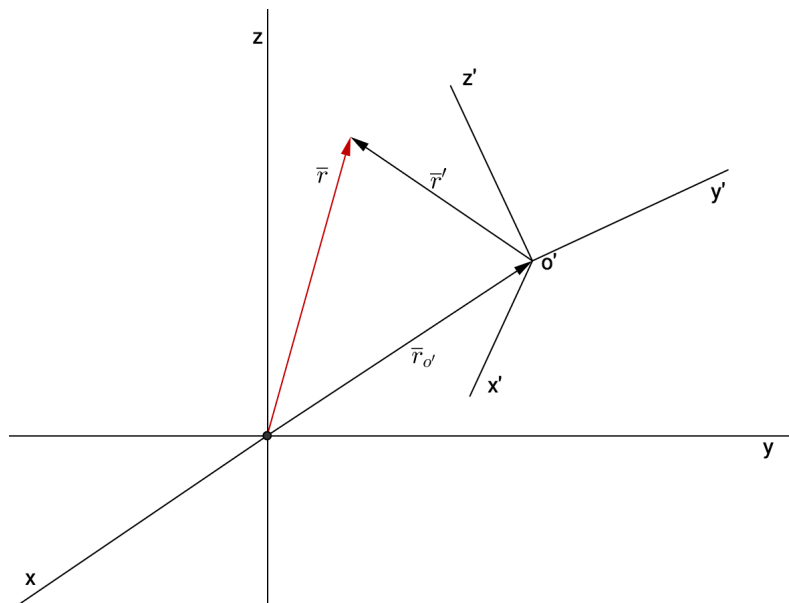
(number)

net displacement over  
time interval  $(t_1, t_2)$

$$\text{average velocity} = \frac{\int_{t_1}^{t_2} \bar{v}(t) dt}{t_2 - t_1}$$

(vector)

### 5.2 relative motion and Galilean transformation



$$\bar{r} = \bar{r}_{o'} + r'$$

Assume  $\dot{\bar{r}}_{o'} = \bar{v}_{o'} = \text{const}$ , that is  $x'y'z'$  moves along a straight line (no rotations, either)

$$\Rightarrow \bar{v} = \bar{v}_{o'} + \bar{v}'$$

Use  $\bar{r}_{o'} = \bar{v}_{o'}t + \bar{r}_{o',init}$  (choose to be 0) (i.e.  $O = O'$  at  $t = 0$ )

$$\Rightarrow \boxed{\bar{r} = \bar{v}_{o'}t + \bar{r}'} \quad \text{Galilean transformation} \quad (33)$$