VE 320 Fall 2021

Introduction to Semiconductor Devices

Instructor: Rui Yang (杨睿)

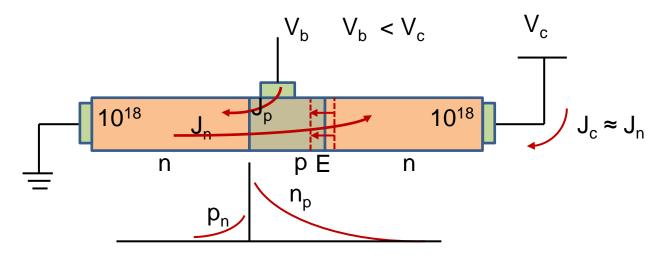
Office: JI Building 434

rui.yang@sjtu.edu.cn

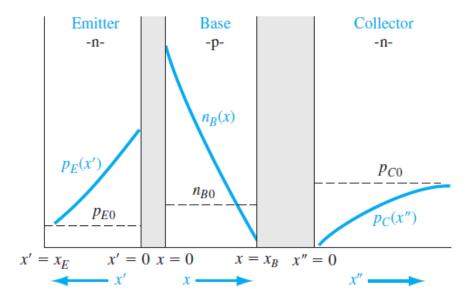


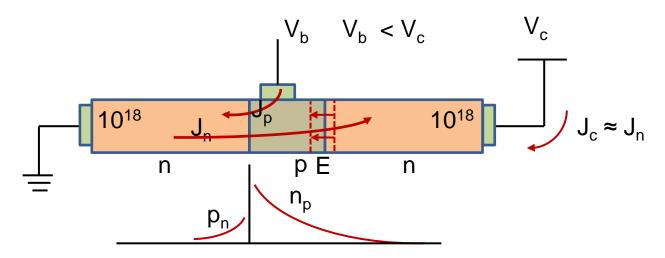
Lecture 9

BJT (Chapter 12)



Now consider actual npn device: forward active mode





Now consider actual npn device: forward active mode

In the neutral base region

$$D_B \frac{\partial^2 (\delta n_B(x))}{\partial x^2} - \frac{\delta n_B(x)}{\tau_{B0}} = 0$$

Excess electron concentration $\delta n_B(x) = n_B(x) - n_{B0}$

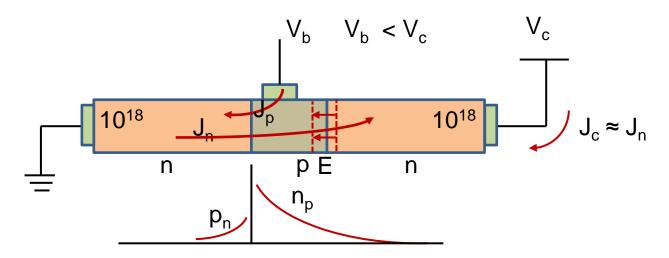
$$\delta n_B(x) = n_B(x) - n_{B0}$$

$$\delta n_B(x) = A \exp\left(\frac{+x}{L_B}\right) + B \exp\left(\frac{-x}{L_B}\right)$$

$$L_B = \sqrt{D_B \tau_{B0}}$$

$$\delta n_B(x=0) \equiv \delta n_B(0) = A + B$$

$$\delta n_B(x = x_B) \equiv \delta n_B(x_B) = A \exp\left(\frac{+x_B}{L_B}\right) + B \exp\left(\frac{-x_B}{L_B}\right)$$

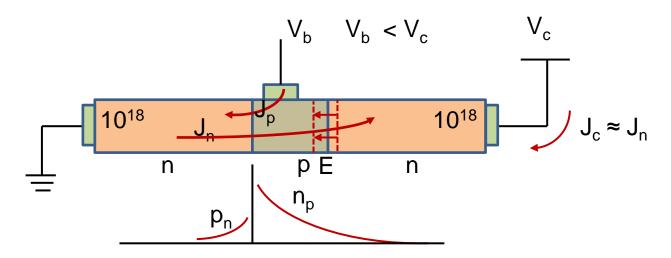


Now consider actual npn device: forward active mode In the neutral base region BE junction forward biased:

$$\delta n_B(0) = n_B(x=0) - n_{B0} = n_{B0} \left[\exp\left(\frac{eV_{BE}}{kT}\right) - 1 \right]$$

BC junction reverse biased:

$$\delta n_B(x_B) = n_B(x = x_B) - n_{B0} = 0 - n_{B0} = -n_{B0}$$

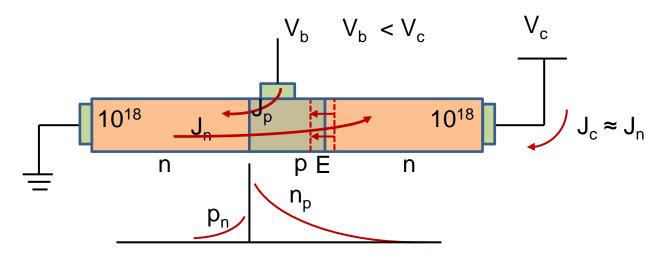


Now consider actual npn device: forward active mode In the neutral base region

We can get

$$A = \frac{-n_{B0} - n_{B0} \left[\exp\left(\frac{eV_{BE}}{kT}\right) - 1 \right] \exp\left(\frac{-x_B}{L_B}\right)}{2 \sinh\left(\frac{x_B}{L_B}\right)}$$

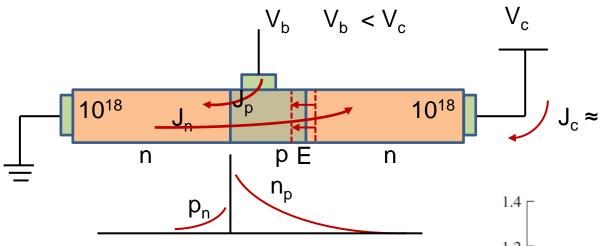
$$B = \frac{n_{B0} \left[\exp\left(\frac{eV_{BE}}{kT}\right) - 1 \right] \exp\left(\frac{x_B}{L_B}\right) + n_{B0}}{2 \sinh\left(\frac{x_B}{L_B}\right)}$$



Now consider actual npn device: forward active mode In the neutral base region

Excess minority carrier electron concentration in the base:

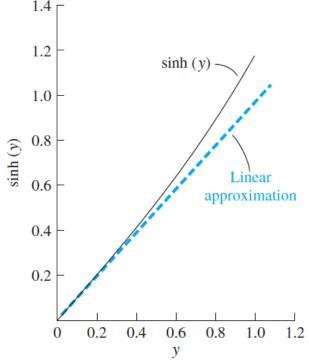
$$\delta n_B(x) = \frac{n_{B0} \left\{ \left[\exp\left(\frac{eV_{BE}}{kT}\right) - 1 \right] \sinh\left(\frac{x_B - x}{L_B}\right) - \sinh\left(\frac{x}{L_B}\right) \right\}}{\sinh\left(\frac{x_B}{L_B}\right)}$$



Now consider actual npn device: forward active mode In the neutral base region

Simplify: $x_B < L_B$

Function	Taylor expansion
sinh (x)	$x + \frac{x^3}{3!} + \frac{x^5}{5!} + \cdots$
$\cosh(x)$	$1 + \frac{x^2}{2!} + \frac{x^4}{4!} + \cdots$
tanh (x)	$x-\frac{x^3}{3}+\frac{2x^5}{15}+\cdots$



$$sinh(x) = \frac{1}{2} \left[\exp(x) - exp(-x) \right]$$

$$\exp(x) = 1 + x + \frac{1}{2}(x)^2 + \cdots$$

$$\exp(-x) = 1 - x + \frac{1}{2}(x)^2 - \dots$$

$$sinh(x) = \frac{1}{2} \left[\exp(x) - \exp(-x) \right] = x$$

$$\exp(x) = 1 + x + \frac{1}{2}(x)^2 + \dots$$

$$\exp(-x) = 1 - x + \frac{1}{2}(x)^2 - \dots$$

$$if x < 1$$

$$cosh(x) - 1 = \frac{1}{2} [exp(x) + exp(-x)] - 1$$

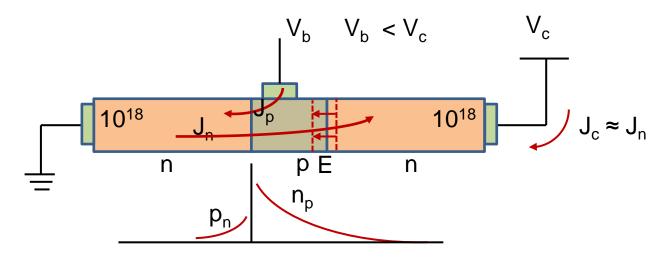
$$exp(x) = 1 + x + \frac{1}{2}(x)^{2} + \cdots$$

$$exp(-x) = 1 - x + \frac{1}{2}(x)^{2} - \cdots$$

$$cosh(x) - 1 = \frac{1}{2} \left[\exp(x) + \exp(-x) \right] - 1 = \frac{1}{2} x^2$$

$$\exp(x) = 1 + x + \frac{1}{2}(x)^{2} + \dots$$

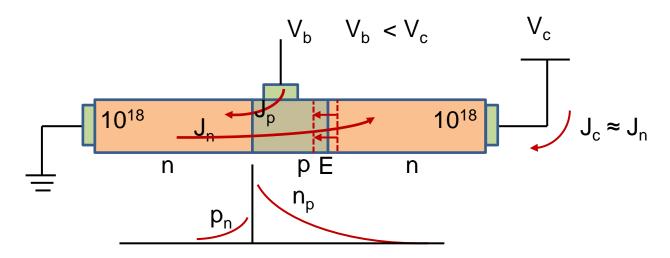
$$\exp(-x) = 1 - x + \frac{1}{2}(x)^{2} - \dots$$
if $x < 1$



Now consider actual npn device: forward active mode In the neutral base region

$$\delta n_B(x) = \frac{n_{B0} \left\{ \left[\exp\left(\frac{eV_{BE}}{kT}\right) - 1 \right] \sinh\left(\frac{x_B - x}{L_B}\right) - \sinh\left(\frac{x}{L_B}\right) \right\}}{\sinh\left(\frac{x_B}{L_B}\right)}$$

$$\delta n_B(x) \approx \frac{n_{B0}}{x_B} \left\{ \left[\exp\left(\frac{eV_{BE}}{kT}\right) - 1 \right] (x_B - x) - x \right\}$$



Now consider actual npn device: forward active mode In the neutral emitter region

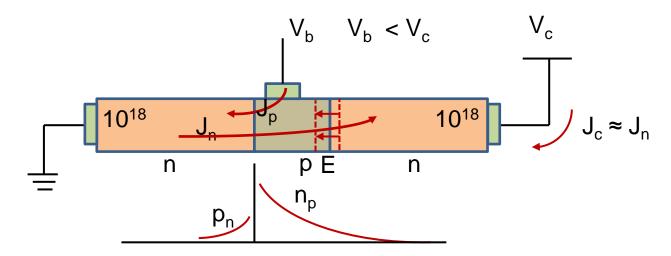
$$D_E \frac{\partial^2 [\delta p_E(x')]}{\partial x'^2} - \frac{\delta p_E(x')}{\tau_{E0}} = 0$$

$$\delta p_E(x') = p_E(x') - p_{E0}$$

General solution:
$$\delta p_E(x') = C \exp\left(\frac{+x'}{L_E}\right) + D \exp\left(\frac{-x'}{L_E}\right)$$

$$L_E = \sqrt{D_E \tau_{E0}}$$





Now consider actual npn device: forward active mode In the neutral emitter region

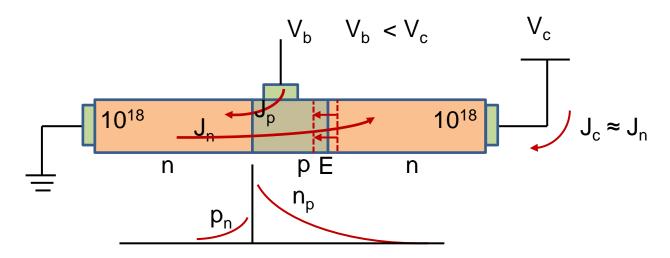
Boundary conditions:
$$\delta p_E(x'=0) \equiv \delta p_E(0) = C + D$$

$$\delta p_E(x' = x_E) \equiv \delta p_E(x_E) = C \exp\left(\frac{x_E}{L_E}\right) + D \exp\left(\frac{-x_E}{L_E}\right)$$

BE junction forward biased:
$$\delta p_E(0) = p_E(x'=0) - p_{E0} = p_{E0} \left[\exp \left(\frac{eV_{BE}}{kT} \right) - 1 \right]$$

Assume infinite surface recombination velocity at $x' = x_F$

$$\delta p_E(x_E) = 0$$



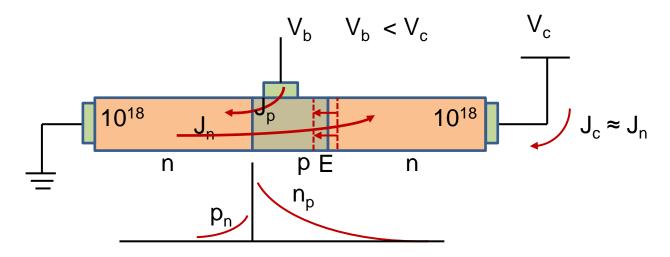
Now consider actual npn device: forward active mode In the neutral emitter region

Boundary conditions:

$$\delta p_E(x') = \frac{p_{E0} \left[\exp \left(\frac{eV_{BE}}{kT} \right) - 1 \right] \sinh \left(\frac{x_E - x'}{L_E} \right)}{\sinh \left(\frac{x_E}{L_E} \right)}$$

Again, when x_E is small:

$$\delta p_E(x') \approx \frac{p_{E0}}{x_E} \left[\exp\left(\frac{eV_{BE}}{kT}\right) - 1 \right] (x_E - x')$$



Now consider actual npn device: forward active mode

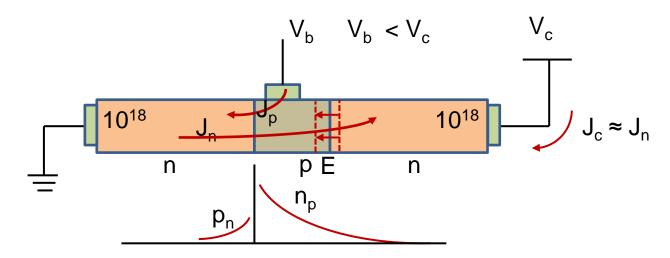
In the neutral collector region

$$D_C \frac{\partial^2 [\delta p_C(x'')]}{\partial x''^2} - \frac{\delta p_C(x'')}{\tau_{C0}} = 0$$

$$\delta p_C(x'') = p_C(x'') - p_{C0}$$

$$\delta p_C(x'') = G \exp\left(\frac{x''}{L_C}\right) + H \exp\left(\frac{-x''}{L_C}\right)$$

$$L_C = \sqrt{D_C \tau_{C0}}.$$



Now consider actual npn device: forward active mode In the neutral collector region

Boundary conditions:

Collector is long, and excess carrier concentration must be finite, so *G*=0

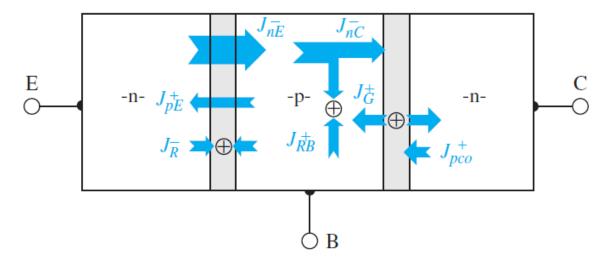
$$\delta p_C(x''=0) \equiv \delta p_C(0) = p_C(x''=0) - p_{C0} = 0 - p_{C0} = -p_{C0}$$

BC junction reverse biased:

$$\delta p_C(x'') = -p_{C0} \exp\left(\frac{-x''}{L_C}\right)$$

Now consider actual npn device: forward active mode

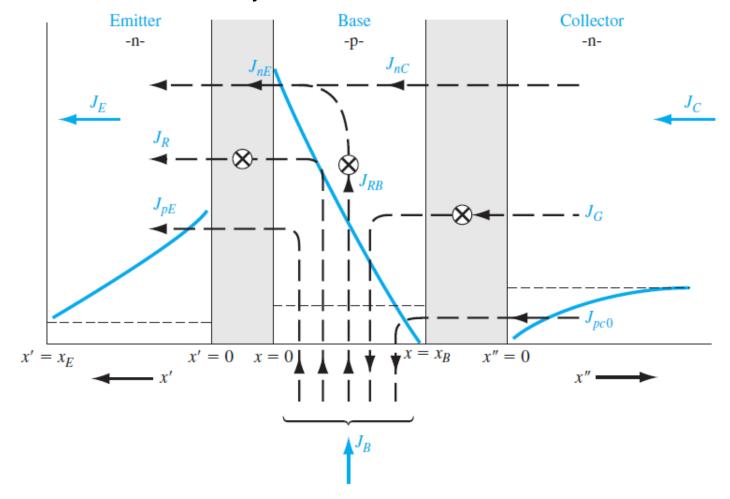
Electron and hole flux



 J_{RB}^{-} is the electron flux injected from the emitter into the base J_{RB}^{+} is the replacement hole flux due to the recombination in the base J_{nC}^{-} is the electron flux that reaches the collector J_{pE}^{+} is the holes from the base that are injected back into the emitter J_{R}^{-} is the recombination in the space charge region of the forward-biased BE junction J_{gC0}^{+} is the generation of electrons and holes in the reverse-biased BC junction J_{pC0}^{-} is the ideal reverse-saturation current in the BC junction

Now consider actual npn device: forward active mode

Electron and hole current density



Now consider actual npn device: forward active mode

Common-base current gain in DC

$$\alpha_0 = \frac{I_C}{I_E}$$

$$lpha_0 = rac{J_C}{J_E} = rac{J_{nC} + J_G + J_{pc0}}{J_{nE} + J_R + J_{pE}}$$

Small signal

$$lpha = rac{\partial J_C}{\partial J_E} = rac{J_{nC}}{J_{nE} + J_R + J_{pE}}$$

We can write

$$lpha = \left(rac{J_{nE}}{J_{nE}+J_{pE}}
ight)\left(rac{J_{nC}}{J_{nE}}
ight)\left(rac{J_{nE}+J_{pE}}{J_{nE}+J_{R}+J_{pE}}
ight)$$

Now consider actual npn device: forward active mode

We can write

$$\alpha = \gamma \alpha_T \delta$$

$$\gamma = \left(rac{J_{nE}}{J_{nE}+J_{pE}}
ight)$$

 $\gamma = \left(rac{J_{nE}}{J_{nE} + J_{pE}}
ight)$ Emitter injection efficiency factor. Diffusion current J_{pE} is not part of the collector current

$$lpha_T = \left(rac{J_{nC}}{J_{nE}}
ight)$$

Base transport factor. Recombination of excess minority carrier electrons in the base

$$\delta = rac{J_{nE} + J_{pE}}{J_{nE} + J_R + J_{pE}}$$

 $\delta = rac{J_{nE} + J_{pE}}{J_{nE} + J_{R} + J_{nF}}$ Recombination factor. Recombination in the forward-biased BE region

We want $\alpha=1$, so every term should be close to 1

Now consider actual npn device: forward active mode

$$\alpha = \gamma \alpha_T \delta$$

We can write

$$\gamma = \left(\frac{J_{nE}}{J_{nE} + J_{pE}}\right) = \frac{1}{\left(1 + \frac{J_{pE}}{J_{nE}}\right)}$$

$$J_{pE} = -eD_E \frac{d[\delta p_E(x')]}{dx'}\bigg|_{x'=0}$$

$$J_{nE} = (-)eD_B \frac{d[\delta n_B(x)]}{dx}\bigg|_{x=0}$$

$$J_{pE} = \frac{eD_{E}p_{E0}}{L_{E}} \left[\exp\left(\frac{eV_{BE}}{kT}\right) - 1 \right] \cdot \frac{1}{\tanh\left(x_{E}/L_{E}\right)}$$

$$J_{nE} = \frac{eD_B n_{B0}}{L_B} \left\{ \frac{1}{\sinh(x_B/L_B)} + \frac{\left[\exp(eV_{BE}/kT) - 1\right]}{\tanh(x_B/L_B)} \right\}$$

Now consider actual npn device: forward active mode

$$\alpha = \gamma \alpha_T \delta$$

We can write

$$\gamma = \left(rac{J_{nE}}{J_{nE}+J_{pE}}
ight) = rac{1}{\left(1+rac{J_{pE}}{J_{nE}}
ight)}$$

Assume

$$\exp\left(\frac{eV_{BE}}{kT}\right) \gg 1$$

$$\frac{\exp(eV_{BE}/kT)}{\tanh(x_B/L_B)} \gg \frac{1}{\sinh(x_B/L_B)}$$

$$\gamma = \frac{1}{1 + \frac{p_{E0}D_EL_B}{n_{B0}D_BL_E} \cdot \frac{\tanh(x_B/L_B)}{\tanh(x_E/L_E)}}$$

Now consider actual npn device: forward active mode

$$\alpha = \gamma \alpha_T \delta$$

$$\gamma = \frac{1}{1 + \frac{p_{E0}D_E L_B}{n_{B0}D_B L_E} \cdot \frac{\tanh(x_B/L_B)}{\tanh(x_E/L_E)}}$$

For γ to be close to 1, we need $p_{\rm E0} << n_{\rm B0}$

$$p_{E0} = \frac{n_i^2}{N_E} \qquad n_{B0} = \frac{n_i^2}{N_B}$$

Then $N_{\rm E} >> N_{\rm B}$

Many more electrons from the n-type emitter than holes from the p-type base will be injected across the BE space charge region

Assume
$$x_B \ll L_B$$
 and $x_E \ll L_E$
$$\gamma \approx \frac{1}{1 + \frac{N_B}{N_E} \cdot \frac{D_E}{D_B} \cdot \frac{x_B}{x_E}}$$

Now consider actual npn device: forward active mode

$$lpha = \gamma lpha_T \delta$$
 $lpha_T = J_{nC}/J_{nE}$

$$J_{nC} = (-)eD_B \frac{d[\delta n_B(x)]}{dx}\Big|_{x=x_R} \qquad J_{nE} = (-)eD_B \frac{d[\delta n_B(x)]}{dx}\Big|_{x=0}$$

Since

$$J_{nC} = (-)eD_B \frac{a[on_B(x)]}{dx} \Big|_{x=x_B}$$

$$/kT$$
) - 1] 1

Then

$$J_{nC} = \frac{eD_B n_{B0}}{L_B} \left\{ \frac{\left[\exp(eV_{BE}/kT) - 1 \right]}{\sinh(x_B/L_B)} + \frac{1}{\tanh(x_B/L_B)} \right\}$$

$$\delta n_B(x) = \frac{n_{B0} \left\{ \left[\exp\left(\frac{eV_{BE}}{kT}\right) - 1 \right] \sinh\left(\frac{x_B - x}{L_B}\right) - \sinh\left(\frac{x}{L_B}\right) \right\}}{\sinh\left(\frac{x_B}{L_B}\right)}$$

Assume

$$V_{BE} \gg kT/e$$

$$\alpha_T = \frac{J_{nC}}{J_{nE}} \approx \frac{\exp(eV_{BE}/kT) + \cosh(x_B/L_B)}{1 + \exp(eV_{BE}/kT)\cosh(x_B/L_B)}$$

Now consider actual npn device: forward active mode

$$\alpha = \gamma \alpha_T \delta$$

$$\alpha_T = \frac{J_{nC}}{J_{nE}} \approx \frac{\exp(eV_{BE}/kT) + \cosh(x_B/L_B)}{1 + \exp(eV_{BE}/kT)\cosh(x_B/L_B)}$$

We want α_{T} close to 1, so

$$x_B \ll L_B$$

$$\alpha_T \approx \frac{1}{\cosh(x_B/L_B)}$$

Taylor expansion:

$$\alpha_T \approx \frac{1}{\cosh(x_B/L_B)} \approx \frac{1}{1 + \frac{1}{2}(x_B/L_B)^2} \approx 1 - \frac{1}{2}(x_B/L_B)^2$$

Now consider actual npn device: forward active mode

$$\alpha = \gamma \alpha_T \delta$$

Assume
$$J_{\it pE} \ll J_{\it nE}$$

$$\delta = \frac{J_{nE} + J_{pE}}{J_{nE} + J_{R} + J_{pE}} \approx \frac{J_{nE}}{J_{nE} + J_{R}} = \frac{1}{1 + J_{R}/J_{nE}}$$

Recombination current density in a forward-biased pn junction

$$J_{R} = \frac{ex_{BE}n_{i}}{2\tau_{0}} \exp\left(\frac{eV_{BE}}{2kT}\right) = J_{r0} \exp\left(\frac{eV_{BE}}{2kT}\right)$$

$$J_{nE} = J_{s0} \exp\left(\frac{eV_{BE}}{kT}\right) \qquad J_{s0} = \frac{eD_{B}n_{B0}}{L_{B} \tanh(x_{B}/L_{B})}$$

$$\delta = \frac{1}{1 + \frac{J_{r0}}{J_{s0}} \exp\left(\frac{-eV_{BE}}{2kT}\right)}$$

We want δ close to 1, so $V_{\rm BE}$ should be large enough

Now consider actual npn device: forward active mode

Current
$$B_0=I_C/I_B$$
 $I_E=I_B+I_C$
 $rac{I_E}{I_C}=rac{I_B}{I_C}+1$
Gain $rac{1}{lpha_0}=rac{1}{eta_0}+1$
 $eta=rac{lpha}{1-lpha}$

Now consider actual npn device: forward active mode In summary Emitter injection efficiency

 $\gamma \approx \frac{1}{1 + \frac{N_B}{N_E} \cdot \frac{D_E}{D_B} \cdot \frac{x_B}{x_E}}$ $(x_B \ll L_B), (x_E \ll L_E)$

Base transport factor

$$\alpha_T \approx \frac{1}{1 + \frac{1}{2} \left(\frac{x_B}{L_B}\right)^2} \qquad (x_B \ll L_B)$$

Recombination factor

$$\delta = \frac{1}{1 + \frac{J_{r0}}{J_{s0}} \exp\left(\frac{-eV_{BE}}{2kT}\right)}$$

Common-base current gain

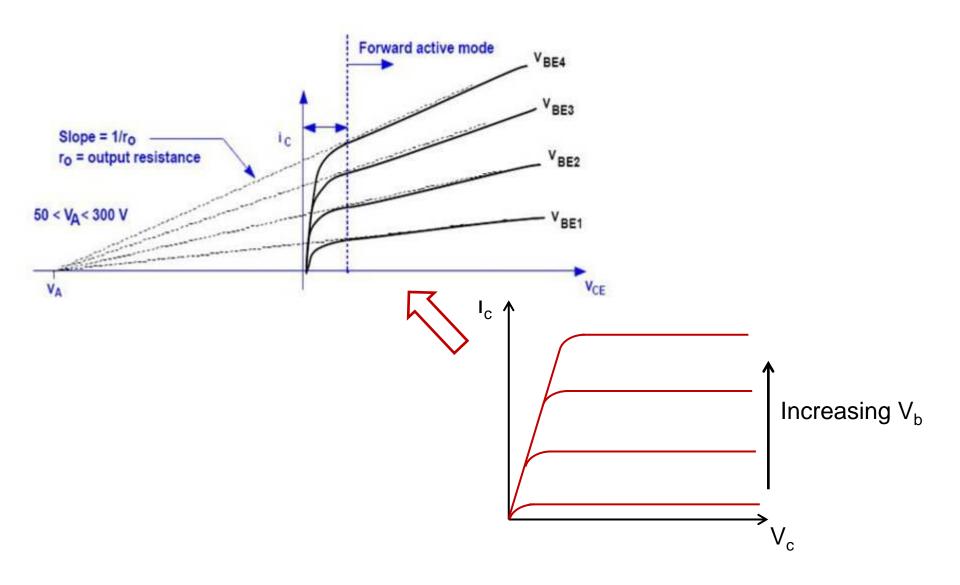
$$\alpha = \gamma \alpha_T \delta \approx \frac{1}{1 + \frac{N_B}{N_E} \cdot \frac{D_E}{D_B} \cdot \frac{x_B}{x_E} + \frac{1}{2} \left(\frac{x_B}{L_B}\right)^2 + \frac{J_{r0}}{J_{s0}} \exp\left(\frac{-eV_{BE}}{2kT}\right)}$$

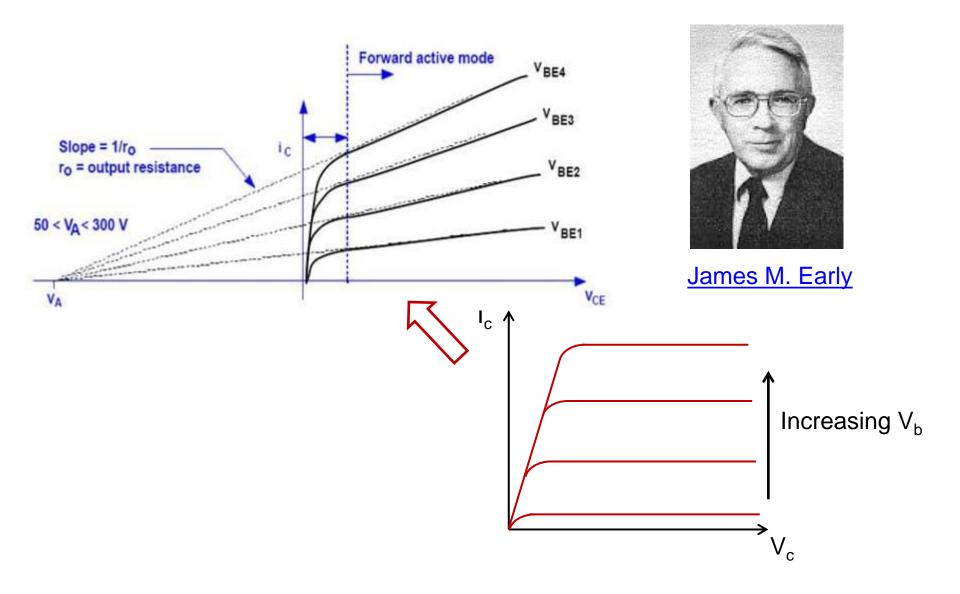
Common-emitter current gain

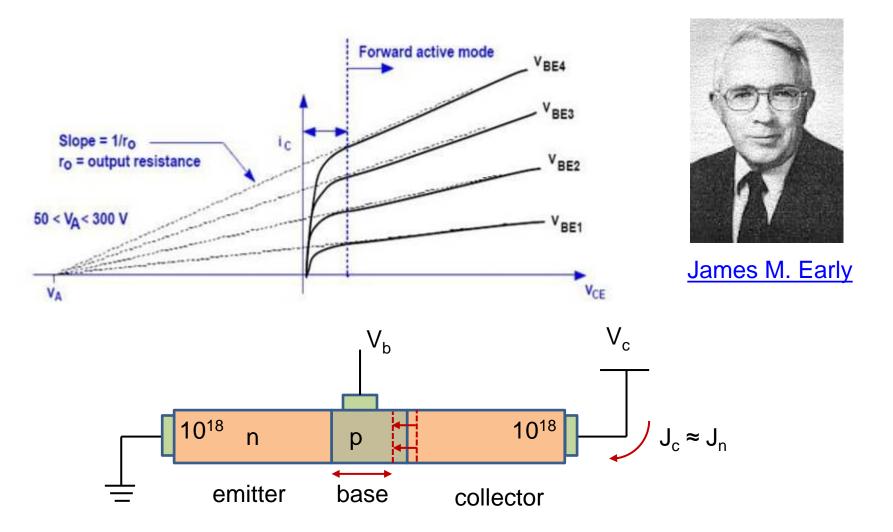
$$\beta = \frac{\alpha}{1 - \alpha} \approx \frac{1}{\frac{N_B}{N_E} \cdot \frac{D_E}{D_B} \cdot \frac{x_B}{x_E} + \frac{1}{2} \left(\frac{x_B}{L_B}\right)^2 + \frac{J_{r0}}{J_{s0}} \exp\left(\frac{-eV_{BE}}{2kT}\right)}$$

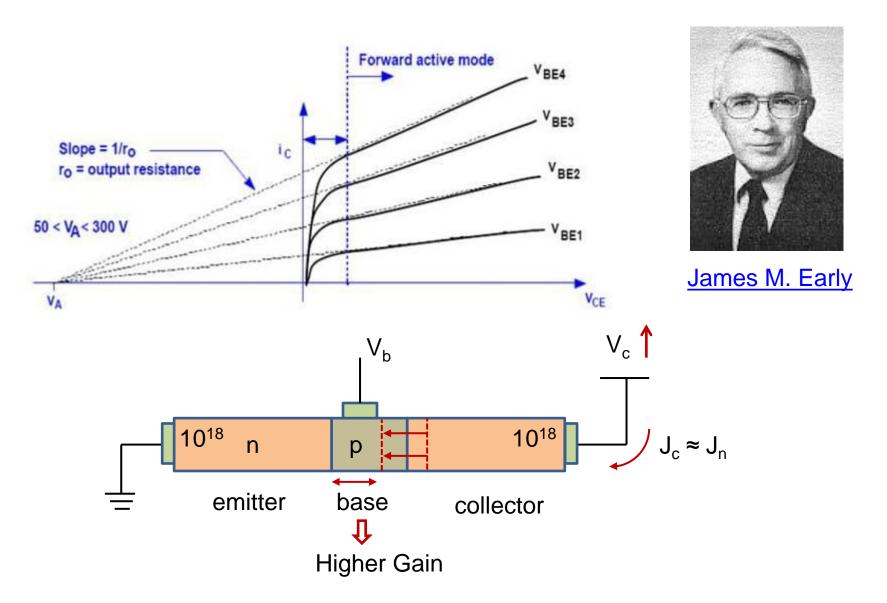
Example:

If we assume a typical value of β to be 100, then $\alpha = 0.99$. If we also assume that $\gamma = \alpha_T = \delta$, then each factor would have to be equal to 0.9967 in order that $\beta = 100$. This calculation gives an indication of how close to unity each factor must be in order to achieve a reasonable current gain.





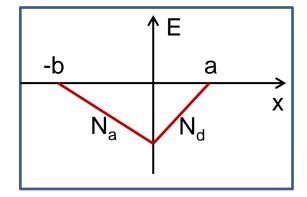


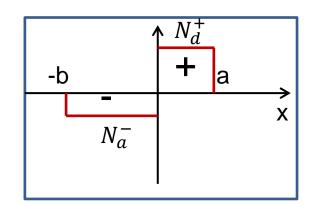


Previously...

$$a = \sqrt{\frac{2\varepsilon(V_{bi} - V_R)}{e} \frac{N_a}{N_d} \frac{1}{N_a + N_d}}$$

$$a = \sqrt{\frac{2\varepsilon(V_{bi} - V_R)}{e} \frac{N_a}{N_d} \frac{1}{N_a + N_d}} \qquad N_a^- b = N_d^+ a \Rightarrow b = \sqrt{\frac{2\varepsilon(V_{bi} - V_R)}{e} \frac{N_d}{N_a} \frac{1}{N_a + N_d}}$$





Previously...

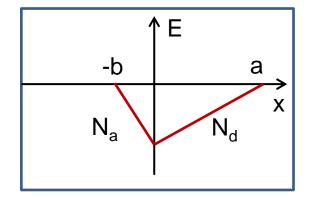
$$a = \sqrt{\frac{2\varepsilon(V_{bi} - V_R)}{e} \frac{N_a}{N_d} \frac{1}{N_a + N_d}}$$

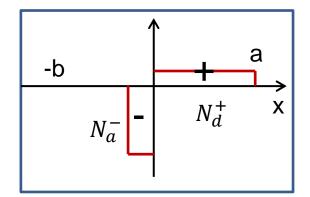
$$a = \sqrt{\frac{2\varepsilon(V_{bi} - V_R)}{e} \frac{N_a}{N_d} \frac{1}{N_a + N_d}} \qquad N_a^- b = N_d^+ a \Rightarrow b = \sqrt{\frac{2\varepsilon(V_{bi} - V_R)}{e} \frac{N_d}{N_a} \frac{1}{N_a + N_d}}$$

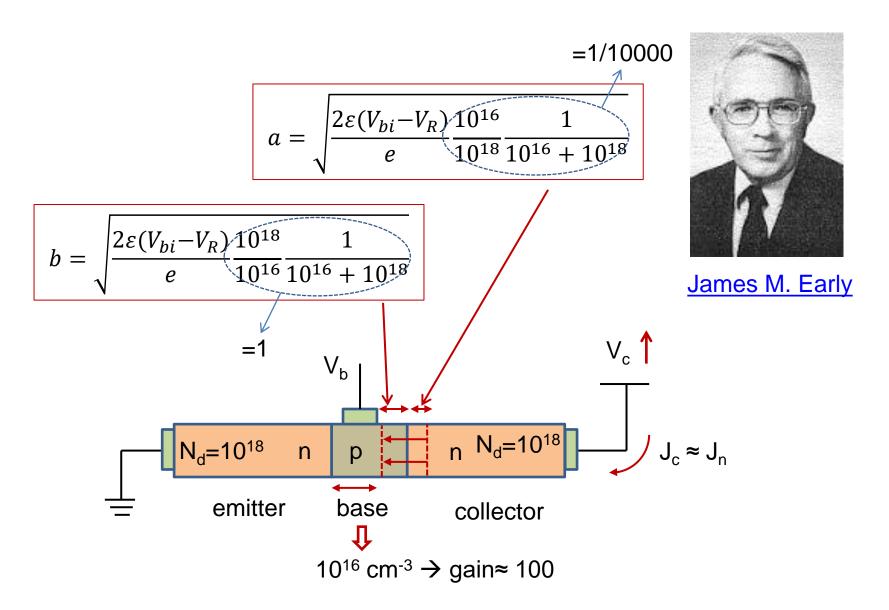
$$N_{a} = 100 N_{d}$$

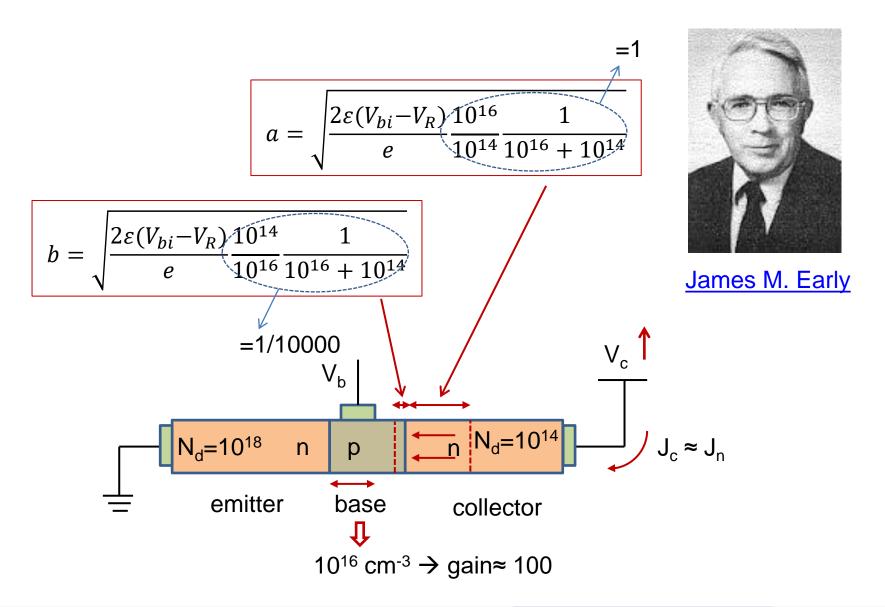
$$a = \sqrt{\frac{2\varepsilon(V_{bi} - V_R)}{e} \frac{100N_{\vec{a}}}{N_d} \frac{1}{100N_{\vec{a}}}}$$

$$b = \sqrt{\frac{2\varepsilon(V_{bi} - V_R)}{e} \frac{N_d}{100N_d}} \frac{1}{100N_d}$$

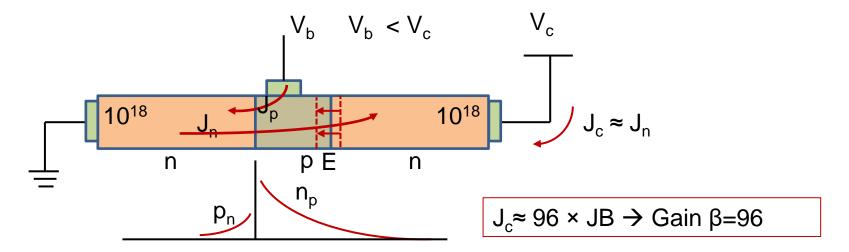








Summary



- 1. highest doping concentration is limited by solubility (<10²⁰)
- 2. Lowest doping concentration is limited by n_i and fabrication process

Basic facts:

- Narrower base → larger gain
- 2. $\beta \approx N_d/N_a$, higher emitter-to-base doping ratio \rightarrow higher gain
- 3. Trade-off for base doping concentration (gain and Early effect)

Breakdown

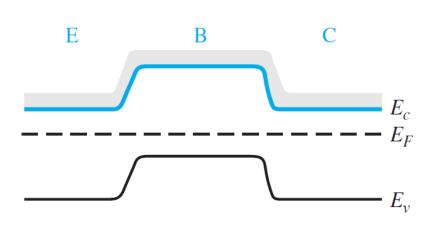
2 breakdown mechanisms:

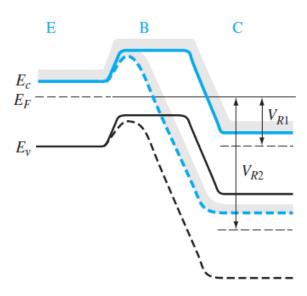
- 1. Punch through
- 2. Avalanche breakdown: with transistor gain

Punch through: large BC junction reverse bias voltage

Depletion region extends through the base region

BE potential barrier lowered, and produce a large current



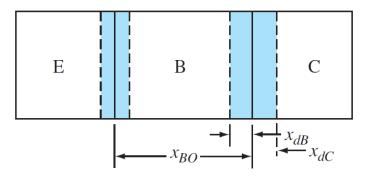


Breakdown

Punch-through: large BC junction reverse bias voltage

Depletion region extends through the base region

BE potential barrier lowered, and produce a large current



$$x_{dB} = x_{BO} = \left\{ \frac{2\epsilon_s(V_{bi} + V_{pt})}{e} \cdot \frac{N_C}{N_B} \cdot \frac{1}{N_C + N_B} \right\}^{1/2}$$

Ignore $V_{\rm bi}$

$$V_{pt} = \frac{ex_{BO}^2}{2\epsilon_s} \cdot \frac{N_B(N_C + N_B)}{N_C}$$

Large breakdown voltage: low N_c