QFT Introduction

Contact Information

Prof. Linda Carpenter Office: M2054 PRB, office hour TBD lmc@physics.osu.edu

Homework

4-5 HW in total, due every 2-3 weeks Take home final

Textbook

Recommended: Srednicki Back-up: Peskin and Schroeder.

Course Structure and Motivation

1. Classical Mech- point like massive particles, built off of action $S = \int dt L$ where L - K - V. Principle of least action is then applied, and variational method requires $\delta S = 0$. As a result, a trajectory that minimize the action is calculated. The resulting equation is the Euler- Lagrange equation:

$$\frac{\partial L}{\partial x_i} = \frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial L}{\partial \dot{x}_i}$$

In addition to the Lagrangian formulation discussed above, another formulation is the Hamiltonian formulation where we construct Hamiltonian around the conjugate commentum $p_i \equiv \frac{\partial L}{\partial \dot{q}_i}$, and the Hamiltonian object $H = p_i \dot{q}_i - L.$

For quantity θ , we can calculate the time evolution of θ by using Possion's bracker $\dot{\theta} = \{\theta, H\}$ where $\{A, B\} = \frac{\partial A}{\partial q_i} \frac{\partial B}{\partial p_i} - \frac{\partial A}{\partial p_i} - \frac{\partial B}{\partial q_i}$.

Anothe important concept for classical mechanics is the phase space, where the phase space is the 2-d distri-

bution of p vs q, and the trajectory is the phase-space orbit.

Note that the processes in classical mechanics are deterministic., meaning the time evolutoiin has infinite precision of the position and momentum of the particle.

2. CFT: For objects that does not behave like point-like massive particles (light), the classical field theory is created to describe such systems. For those systems, we replace the concept of coordiates $q \to \Phi$ where the latter $\Phi(\vec{x},t)$ is valued at every point in space.

The trajectory of the particle is then replaced by the time evolution of Φ . We still accomplish is by using the principle of least action $S = \int dt L$, where L is redefined as $L(\Phi, \dot{\Phi}, t)$. Effectively, this is a direct replacement from $q \to \Phi$. In the same light, we can re-define $p \to \Pi \equiv \frac{\partial L}{\partial \dot{\Phi}}$. We can re-write the E-L theorem:

$$\partial_{\mu} \left(\frac{\partial L}{\partial \partial_{\mu} \Phi} \right) - \frac{\partial L}{\partial \Phi} = 0$$

The difference here is that the degree of freedom for fields are more complicated as it is a continuous object of all spaces. However, this formulation is stil deterministic.

3. Quantum Mech: QM still considers "massive" particles. It can be fairly easily obtained from the classical Hamiltonian formulation. However, the main difference is that the proabblity is not deterministic. Instead of a trajectory, we know consider a quantum state $|\phi\rangle$, which is a set of quantum numbers, and we consider the time evolution of the quantum states. We still define the Hamiltonian \hat{H} as an operator on vector space. This is given by the Schrodinger equation:

$$i\hbar \frac{\mathrm{d}}{\mathrm{d}t} |\phi\rangle = \hat{H} |\phi\rangle$$

And the expectation values of an observable θ , we simply take $\langle \phi | \theta | \phi \rangle$ of operator θ , and $|\hat{\theta}, \hat{H}|$ gives the time evolution of that operator.

1

Anothere feature of QM is that not all θ_i can be determined at the same time such as p_i, x_i to an infinite certainty. To determine if comensurability, we simply check if $[\theta_i, \theta_j] = 0$. Therefore, well-defined quantum numbers should be comensurable.

Additionally, one can compute transition probability that a particle go from one state to another state using the Hamiltonian as the time evolution operator by doing $\langle \phi_2 | e^{i\hat{H}t} | \phi_i \rangle$.

4. QFT: The continuum limit of QM: Normal Quantum Mech does not allow for creation and annaliation of particles, which is required by taking the relativistic limit of the quantum mechanics. However, there are clear events in the universe where creation and annaliation processes are needed (e.g. electrons and position collider resulting in W^+W^- pair.)

The only way to describe changing particle numbers, is to abandon the concept of particles. Instead, the particles are now excitation of a quantum field.

Under a Lorentz transforation, for a electron scattering from a nucleus, we could create the process of a nuclei emits a photon that creates an electron positron pair.

Similiar to Hilbert Space in QM, we create a Fock Space, which is a direct sum of all of the n-particle Hilbert Spaces.

Another issue with the QM is the "causality". c is the ultimate speed in the universe. However, a simple QM computation of state transition can go beyond the speed of light. Therefore, to obey time-like causality, we need QFT. For instance, if we calculate the probablity of a particle propergate from point $\vec{x} \to \vec{y}$, we formulate:

$$\langle y|\,e^{i\hat{H}t}\,|x\rangle$$

However, if x and y are so separated (> ct), one should get a 0 transition probablity. However, in traditional QM, the transition probablity is always none-zero.