

QFT Introduction

Contact Information

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Homework

4-5 HW in total, due every 2-3 weeks
Take home final

Course Structure and Motivation

1. Classical Mech- point like massive particles, built off of action $S = \int dt L$ where $L = K - V$. Principle of least action is then applied, and variational method requires $\delta S = 0$. As a result, a trajectory that minimize the action is calculated. The resulting equation is the Euler- Lagrange equation:

$$\frac{\partial L}{\partial x_i} = \frac{d}{dt} \frac{\partial L}{\partial \dot{x}_i}$$

In addition to the Lagrangian formulation discussed above, another formulation is the Hamiltonian formulation where we construct Hamiltonian around the conjugate momentum $p_i \equiv \frac{\partial L}{\partial \dot{q}_i}$, and the Hamiltonian object $H = p_i \dot{q}_i - L$.

For quantity θ , we can calculate the time evolution of θ by using Poisson's bracket $\dot{\theta} = \{\theta, H\}$ where $\{A, B\} = \frac{\partial A}{\partial q_i} \frac{\partial B}{\partial p_i} - \frac{\partial A}{\partial p_i} \frac{\partial B}{\partial q_i}$.

Another important concept for classical mechanics is the phase space, where the phase space is the 2-d distribution of p vs q , and the trajectory is the phase-space orbit.

Note that the processes in classical mechanics are deterministic., meaning the time evolution has infinite precision of the position and momentum of the particle.

2. CFT: For objects that does not behave like point-like massive particles (light), the classical field theory is created to describe such systems. For those systems, we replace the concept of coordinates $q \rightarrow \Phi$ where the latter $\Phi(\vec{x}, t)$ is valued at every point in space.

The trajectory of the particle is then replaced by the time evolution of Φ . We still accomplish this by using the principle of least action $S = \int dt L$, where L is redefined as $L(\Phi, \dot{\Phi}, t)$. Effectively, this is a direct replacement from $q \rightarrow \Phi$. In the same light, we can re-define $p \rightarrow \Pi \equiv \frac{\partial L}{\partial \dot{\Phi}}$. We can re-write the E-L theorem:

$$\partial_\mu \left(\frac{\partial L}{\partial \partial_\mu \Phi} \right) - \frac{\partial L}{\partial \Phi} = 0$$

The difference here is that the degree of freedom for fields are more complicated as it is a continuous object of all spaces. However, **this formulation is still deterministic.**

3. Quantum Mech: QM still considers "massive" particles. It can be fairly easily obtained from the classical Hamiltonian formulation. However, the main difference is that the probability is not deterministic. Instead of a trajectory, we now consider a quantum state $|\phi\rangle$, which is a set of quantum numbers, and we consider the time evolution of the quantum states. We still define the Hamiltonian \hat{H} as an operator on vector space. This is given by the Schrodinger equation:

$$i\hbar \frac{d}{dt} |\phi\rangle = \hat{H} |\phi\rangle$$

And the expectation values of an observable θ , we simply take $\langle \phi | \theta | \phi \rangle$ of operator θ , and $[\hat{\theta}, \hat{H}]$ gives the time evolution of that operator.

Another feature of QM is that not all θ_i can be determined at the same time such as p_i, x_i to an infinite certainty. To determine if commensurability, we simply check if $[\theta_i, \theta_j] = 0$. Therefore, well-defined quantum numbers should be commensurable.

Additionally, one can compute transition probability that a particle go from one state to another state using the Hamiltonian as the time evolution operator by doing $\langle \phi_2 | e^{i\hat{H}t} | \phi_1 \rangle$.

4. QFT: The continuum limit of QM: Normal Quantum Mech does not allow for creation and annihilation of particles, which is required by taking the relativistic limit of the quantum mechanics. However, there are clear events in the universe where creation and annihilation processes are needed (e.g. electrons and positron collider resulting in W^+W^- pair.)

The only way to describe changing particle numbers, is to abandon the concept of particles. Instead, the particles are now excitation of a quantum field.

Under a Lorentz transformation, for a electron scattering from a nucleus, we could create the process of a nuclei emits a photon that creates an electron positron pair.

Similar to Hilbert Space in QM, we create a Fock Space, which is a direct sum of all of the n-particle Hilbert Spaces.

Another issue with the QM is the "causality". c is the ultimate speed in the universe. However, a simple QM computation of state transition can go beyond the speed of light. Therefore, to obey time-like causality, we need QFT. For instance, if we calculate the probability of a particle propagate from point $\vec{x} \rightarrow \vec{y}$, we formulate:

$$\langle y | e^{i\hat{H}t} | x \rangle$$

However, if x and y are so separated ($> ct$), one should get a 0 transition probability. However, in traditional QM, the transition probability is always non-zero.