Bessel Function

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The Bessel's differential equation is given by

$$\frac{1}{z}\frac{\mathrm{d}}{\mathrm{d}z}\left(z\frac{\mathrm{d}u}{\mathrm{d}z}\right) + \left(1 - \frac{\nu^2}{z^2}\right)u = \frac{\mathrm{d}^2u}{\mathrm{d}z^2} + \frac{1}{z}\frac{\mathrm{d}u}{\mathrm{d}z} + \left(1 - \frac{\nu^2}{z^2}\right)u = 0. \tag{1}$$

The Bessel functions are the solution to this equation. The definitions of the Bessel functions are given by

Bessel
$$J_{\nu}(z) = \left(\frac{z}{2}\right)^{\nu} \sum_{n=0}^{\infty} \frac{(-1)^n (z/2)^{2n}}{n! \Gamma(\nu+n+1)} \quad (z \neq \text{negative real})$$
 (2)

Neumann
$$N_{\nu}(z) = Y_{\nu}(z) = \frac{1}{\sin \nu \pi} [\cos \nu \pi J_{\nu}(z) - J_{-\nu}(z)]$$
 ($\nu \neq \text{integer}, z \neq \text{negative real}$) (3)

$$N_{n}(z) = Y_{n}(z) = \frac{1}{\pi} \left[\frac{\partial J_{\nu}(z)}{\partial \nu} - (-1)^{n} \frac{\partial J_{-\nu}(z)}{\partial \nu} \right]$$

$$= \frac{2}{\pi} J_{n}(z) \left(\gamma + \log \frac{z}{2} \right)$$

$$- \frac{1}{\pi} \left(\frac{z}{2} \right)^{n} \sum_{k=0}^{\infty} \frac{(-1)^{k}}{k!(n+k)!} \left(\frac{z}{2} \right)^{2k} \left[\sum_{m=1}^{k} \frac{1}{m} + \sum_{m=1}^{n+k} \frac{1}{m} \right]$$

$$- \frac{1}{\pi} \left(\frac{z}{2} \right)^{-n} \sum_{n=0}^{n-1} \frac{(n-r-1)!}{r!} \left(\frac{z}{2} \right)^{2r} \quad (n=0,1,2,\cdots,z \neq \text{negative real})$$

where $\gamma = 0.57721...$ is the Euler's gamma,

and the last term is replaced by 0 for
$$n = 0$$
. (4)

Hankel (first kind)
$$H_{\nu}^{(1)} = J_{\nu}(z) + iN_{\nu}(z) \tag{5}$$

Hankel (second kind)
$$H_{\nu}^{(2)} = J_{\nu}(z) - iN_{\nu}(z). \tag{6}$$

 $J_{\nu},~N_{\nu},~H_{\nu}^{(1,2)}$ are also called the first, second, and the third kind of the cylindrical function, and they are collectively referred as the ν -th cylindrical or Bessel (in a broad sense) function.

Another expression for the Bessel function (of the first kind) having an integer parameter is

$$J_n(z) = \frac{\mathrm{i}^{-n}}{2\pi} \int_0^{2\pi} \cos n\theta e^{\mathrm{i}z\cos\theta} \mathrm{d}\theta. \tag{7}$$

Now we show this satisfies the Bessel's equation. J_n satisfies the following recurrence relation,

$$J_{n-1} + J_{n+1} = \frac{2n}{z} J_n. (8)$$

Proof.

$$z(J_{n-1} + J_{n+1}) = z \frac{\mathrm{i}^{-n+1}}{2\pi} \int_0^{2\pi} \cos(n-1)\theta e^{\mathrm{i}z\cos\theta} d\theta + z \frac{\mathrm{i}^{-n-1}}{2\pi} \int_0^{2\pi} \cos(n+1)\theta e^{\mathrm{i}z\cos\theta} d\theta$$

$$= \frac{\mathrm{i}^{-n}}{2\pi} \int_0^{2\pi} \mathrm{i}z(\cos(n-1)\theta - \cos(n+1)\theta) e^{\mathrm{i}z\cos\theta} d\theta$$

$$= \frac{\mathrm{i}^{-n}}{2\pi} \int_0^{2\pi} 2\mathrm{i}z\sin n\theta \sin \theta e^{\mathrm{i}z\cos\theta} d\theta$$

$$= -2\frac{\mathrm{i}^{-n}}{2\pi} \int_0^{2\pi} \sin n\theta \frac{\partial}{\partial \theta} \left(e^{\mathrm{i}z\cos\theta} \right) d\theta$$

$$= -2\frac{\mathrm{i}^{-n}}{2\pi} \left(\left[\sin n\theta e^{\mathrm{i}z\cos\theta} \right]_0^{2\pi} - \int_0^{2\pi} n\cos n\theta e^{\mathrm{i}z\cos\theta} d\theta \right)$$

$$= 2nJ_n.$$

Next, we evaluate the derivatives with respect to z.

 $\frac{\mathrm{d}J_{n}}{\mathrm{d}z} = \frac{\mathrm{i}^{-n}}{2\pi} \int_{0}^{2\pi} \mathrm{i}\cos\theta\cos n\theta e^{\mathrm{i}z\cos\theta} \mathrm{d}\theta \\
= \frac{\mathrm{i}^{-n}}{2\pi} \int_{0}^{2\pi} \frac{\mathrm{i}}{2} (\cos(n+1)\theta + \cos(n-1)\theta) e^{\mathrm{i}z\cos\theta} \mathrm{d}\theta \\
= \frac{1}{2} \frac{\mathrm{i}^{-n+1}}{2\pi} \frac{\mathrm{i}^{2}}{\mathrm{i}^{2}} \int_{0}^{2\pi} \cos(n+1)\theta e^{\mathrm{i}z\cos\theta} \mathrm{d}\theta + \frac{1}{2} \frac{\mathrm{i}^{-n+1}}{2\pi} \int_{0}^{2\pi} \cos(n-1)\theta e^{\mathrm{i}z\cos\theta} \mathrm{d}\theta \\
= -\frac{1}{2} J_{n+1} + \frac{1}{2} J_{n-1} = -\frac{1}{2} (\frac{2n}{z} J_{n} - J_{n-1}) + \frac{1}{2} J_{n-1} = J_{n-1} - \frac{n}{z} J_{n}, \tag{9}$ $\frac{\mathrm{d}^{2} J_{n}}{\mathrm{d}z^{2}} = \frac{\mathrm{d} J_{n-1}}{\mathrm{d}z} - \frac{n}{z} \frac{\mathrm{d} J_{n}}{\mathrm{d}z} + \frac{n}{z^{2}} J_{n} \\
= \left(J_{n-2} - \frac{n-1}{z} J_{n-1} \right) - \frac{n}{z} \left(J_{n-1} - \frac{n}{z} J_{n} \right) + \frac{n}{z^{2}} J_{n} \\
= \left(-J_{n} + \frac{2(n-1)}{z} J_{n-1} \right) - \frac{2n-1}{2} J_{n-1} + \frac{n^{2}+n}{z^{2}} J_{n} \\
= -\frac{1}{z} J_{n-1} + \left(\frac{n^{2}+n}{z^{2}} - 1 \right) J_{n} \\
= -\frac{1}{z} \left(\frac{\mathrm{d} J_{n}}{\mathrm{d}z} + \frac{n}{z} J_{n} \right) + \left(\frac{n^{2}+n}{z^{2}} - 1 \right) J_{n} \\
= -\frac{1}{z} \frac{\mathrm{d} J_{n}}{z^{2}} + \left(\frac{n^{2}}{z^{2}} - 1 \right) J_{n}. \tag{10}$

Thus, J_n satisfies the Bessel's equation.

(7) can be derived from the generating function expression of the Bessel function

$$e^{iz\sin\theta} = \sum_{n=-\infty}^{\infty} J_n(z)e^{in\theta}.$$
 (11)

By replacing θ by $-\theta$, it is easily expanded to

$$e^{\pm iz\sin\theta} = \sum_{n=-\infty}^{\infty} J_n(z)e^{\pm in\theta}.$$
 (12)

We also use a transform $\theta' = \theta \pm \pi/2$ to obtain the expression for $\exp(\pm iz \cos \theta)$. Plugging the relations

$$\sin \theta = \sin(\theta' \mp \frac{\pi}{2}) = \mp \cos \theta' \sin \frac{\pi}{2} = \mp \cos \theta',$$

$$e^{in\theta} = e^{in(\theta' \mp \pi/2)} = e^{\mp in\pi/2} e^{in\theta'} = (\mp i)^n e^{in\theta'}$$
(13)

into (11), we get

$$e^{iz\sin\theta} = e^{\mp iz\cos\theta'} = \sum_{n=-\infty}^{\infty} J_n(z)(\mp i)^n e^{in\theta'}.$$
 (14)

(This is called the Jacobi-Anger expansion.) Thus, we have the following relations,

$$e^{\pm iz\sin\theta} = \sum_{n=-\infty}^{\infty} J_n(z)e^{\pm in\theta}, \quad e^{\pm iz\cos\theta} = \sum_{n=-\infty}^{\infty} (\pm i)^n J_n(z)e^{in\theta}.$$
 (15)

Consider the integral

$$\frac{1}{2\pi} \int_0^{2\pi} \cos n\theta e^{\mathrm{i}z\cos\theta} \mathrm{d}\theta. \tag{16}$$

Substituting the above expression yields

$$\frac{1}{2\pi} \int_{0}^{2\pi} \cos n\theta e^{iz \cos \theta} d\theta = \frac{1}{2\pi} \int_{0}^{2\pi} \frac{e^{in\theta} + e^{-in\theta}}{2} \sum_{l=-\infty}^{\infty} i^{l} J_{l}(z) e^{il\theta} d\theta
= \sum_{l=\infty}^{\infty} \frac{i^{l}}{2} J_{l}(z) \frac{1}{2\pi} \int_{0}^{2\pi} \left(e^{i(n+l)\theta} + e^{i(-n+l)\theta} \right) d\theta
= \frac{i^{n}}{2} J_{n}(z) + \frac{i^{-n}}{2} J_{-n}(z)
= \frac{i^{n}}{2} J_{n}(z) \left(1 + (-1)^{n} i^{-2n} \right) = i^{n} J_{n}(z),$$
(17)

wheere $J_{-n}(z) = (-1)J_n(z)$ is used. Thus, we obtain

$$J_n(z) = \frac{\mathrm{i}^{-n}}{2\pi} \int_0^{2\pi} \cos n\theta e^{\mathrm{i}z\cos\theta} \mathrm{d}\theta \tag{18}$$

It is straightforwardly calculated that

$$\frac{1}{2\pi} \int_0^{2\pi} \cos n\theta e^{-iz\cos\theta} d\theta = (-i)^n J_n(z). \tag{19}$$

Thus,

$$J_n(z) = \frac{(\pm i)^{-n}}{2\pi} \int_0^{2\pi} \cos n\theta e^{\pm iz \cos \theta} d\theta.$$
 (20)

Next, we consider

$$\frac{1}{2\pi} \int_{0}^{2\pi} \sin n\theta e^{\pm iz \cos \theta} d\theta. \tag{21}$$

It is calculated as

$$\frac{1}{2\pi} \int_{0}^{2\pi} \sin n\theta e^{\pm iz \cos \theta} d\theta = \frac{1}{2\pi} \int_{0}^{2\pi} \frac{e^{in\theta} - e^{-in\theta}}{2i} \sum_{l=-\infty}^{\infty} (\pm i)^{l} J_{l}(z) e^{il\theta} d\theta
= \sum_{l=-\infty}^{\infty} \frac{(\pm i)^{l}}{2i} J_{l}(z) \frac{1}{2\pi} \int_{0}^{2\pi} \left(e^{i(n+l)\theta} - e^{i(-n+l)\theta} \right) d\theta
= -\frac{(\pm i)^{n}}{2i} J_{n}(z) + \frac{(\pm i)^{-n}}{2i} J_{-n}(z)
= \frac{(\pm i)^{n}}{2i} J_{n}(z) \left(-1 + (\pm i)^{-2n} (-1)^{n} \right) = 0.$$
(22)

Other combinations are

$$\frac{1}{2\pi} \int_0^{2\pi} \cos n\theta e^{\pm iz \sin \theta} d\theta, \quad \frac{1}{2\pi} \int_0^{2\pi} \sin n\theta e^{\pm iz \cos \theta} d\theta.$$
 (23)

Each calculated as follows,

$$\frac{1}{2\pi} \int_{0}^{2\pi} \cos n\theta e^{\pm iz \sin \theta} d\theta = \frac{1}{2\pi} \int_{0}^{2\pi} \frac{e^{in\theta} + e^{-in\theta}}{2} \sum_{l=-\infty}^{\infty} J_{l}(z) e^{\pm il\theta} d\theta
= \sum_{l=-\infty}^{\infty} \frac{1}{2} J_{l}(z) \frac{1}{2\pi} \int_{0}^{2\pi} \left(e^{i(n\pm l)\theta} + e^{i(-n\pm l)\theta} \right) d\theta
= \frac{1}{2} J_{\mp n}(z) + \frac{1}{2} J_{\pm n}(z) = \frac{1}{2} (1 + (-1)^{n}) J_{n}(z)
= \begin{cases} J_{n}(z) & n = \text{even} \\ 0 & n = \text{odd} \end{cases},$$

$$\frac{1}{2\pi} \int_{0}^{2\pi} \sin n\theta e^{\pm iz \sin \theta} d\theta = \frac{1}{2\pi} \int_{0}^{2\pi} \frac{e^{in\theta} - e^{-in\theta}}{2i} \sum_{l=-\infty}^{\infty} J_{l}(z) e^{\pm il\theta} d\theta
= \sum_{l=-\infty}^{\infty} \frac{1}{2i} J_{l}(z) \frac{1}{2\pi} \int_{0}^{2\pi} \left(e^{i(n\pm l)\theta} - e^{i(-n\pm l)\theta} \right) d\theta
= -\frac{i}{2} J_{\mp n}(z) + \frac{i}{2} J_{\pm n}(z) = \pm \frac{i}{2} (1 - (-1)^{n}) J_{n}(z)
= \begin{cases} 0 & n = \text{even} \\ \pm i J_{n}(z) & n = \text{odd} \end{cases}.$$
(25)

We can combine these results in the following way,

$$\frac{1}{2\pi} \int_0^{2\pi} e^{in\theta} e^{\pm z \sin \theta} d\theta = \frac{1}{2\pi} \int_0^{2\pi} (\cos n\theta + i \sin n\theta) e^{\pm z \sin \theta} d\theta$$

$$= \begin{cases} J_n(z) & n = \text{even} \\ \mp J_n(z) & n = \text{odd} \end{cases}$$

$$= (\mp 1)^n J_n(z). \tag{26}$$

Or, we can write

$$J_n(z) = \frac{(\mp 1)^{-n}}{2\pi} \int_0^{2\pi} e^{\mathrm{i}n\theta} e^{\pm \mathrm{i}z \sin\theta} \mathrm{d}\theta. \tag{27}$$

Summary

Here, we summarize the relations.

$$e^{\pm iz\sin\theta} = \sum_{n} J_n e^{in\theta} \qquad e^{\pm iz\cos\theta} = \sum_{n} (\pm i)^n J_n e^{in\theta}$$
 (28)

$$\frac{1}{2\pi} \int_0^{2\pi} \cos n\theta e^{\pm iz\cos\theta} d\theta = (\pm i)^n J_n(z) \qquad J_n(z) = \frac{(\pm i)^{-n}}{2\pi} \int_0^{2\pi} \cos n\theta e^{\pm iz\cos\theta} d\theta \qquad (29)$$

$$\frac{1}{2\pi} \int_0^{2\pi} \sin n\theta e^{\pm iz \cos \theta} d\theta = 0 \tag{30}$$

$$\frac{1}{2\pi} \int_0^{2\pi} e^{\mathrm{i}n\theta} e^{\pm \mathrm{i}z\cos\theta} = (\pm \mathrm{i})^n J_n(z) \qquad J_n(z) = \frac{(\pm \mathrm{i})^{-n}}{2\pi} \int_0^{2\pi} e^{\mathrm{i}n\theta} e^{\pm \mathrm{i}z\cos\theta} d\theta \qquad (31)$$

$$\frac{1}{2\pi} \int_0^{2\pi} \cos n\theta e^{\pm iz \sin \theta} d\theta = \begin{cases} J_n(z) & n = \text{even} \\ 0 & n = \text{odd} \end{cases}$$
 (32)

$$\frac{1}{2\pi} \int_0^{2\pi} \sin n\theta e^{\pm iz \sin \theta} d\theta = \begin{cases} 0 & n = \text{even} \\ \pm iJ_n(z) & n = \text{odd} \end{cases}$$
 (33)

$$\frac{1}{2\pi} \int_0^{2\pi} e^{\mathrm{i}n\theta} e^{\pm \mathrm{i}z\sin\theta} \mathrm{d}\theta = (\mp 1)^n J_n(z) \qquad J_n(z) = \frac{(\mp 1)^{-n}}{2\pi} \int_0^{2\pi} e^{\mathrm{i}n\theta} e^{\pm \mathrm{i}z\sin\theta} \mathrm{d}\theta \qquad (34)$$

Modified Bessel functions

By giving pure imaginary argument to the Bessel functions, we obtain the modified Bessel functions.

First kind
$$I_{\nu}(z) = \left(\frac{z}{2}\right)^{\nu} \sum_{n=0}^{\infty} \frac{(z/2)^{2n}}{n!\Gamma(\nu+n+1)}$$
 (35)

$$= \begin{cases} e^{-\mathrm{i}\nu\pi/2} J_{\nu} \left(e^{\mathrm{i}\pi/2} z \right) & (-\pi < \arg z < \pi/2) \\ e^{3\mathrm{i}\nu\pi/2} J_{\nu} \left(e^{-3\mathrm{i}\pi/2} z \right) & (\pi/2 < \arg z < \pi) \end{cases}$$
 $(z \neq \text{negative real})$ (36)

Second kind
$$K_{\nu}(z) = \frac{\pi}{2} \frac{I_{-\nu}(z) - I_{\nu}(z)}{\sin \nu \pi} = \frac{i\pi}{2} e^{i\nu\pi/2} H_{\nu}^{(1)}(iz) = \frac{i\pi}{2} e^{-i\nu\pi/2} H_{-\nu}^{(1)}(iz)$$
 (37)

$$(\nu \neq \text{integer}, z \neq \text{negative real}),$$
 (38)

$$K_n(z) = K_{-n}(z) = \frac{(-1)^n}{2} \left[\frac{\partial I_{-\nu}(z)}{\partial \nu} - \frac{\partial I_{\nu}(z)}{\partial \nu} \right]_{\nu=n}$$
(39)

$$=(-1)^{n+1}I_n(z)\left(\gamma + \log\frac{z}{2}\right) \tag{40}$$

$$+\frac{(-1)^n}{2} \sum_{k=0}^{\infty} \frac{(z/2)^{n+2k}}{k!(n+k)!} \left[\sum_{m=1}^k \frac{1}{m} + \sum_{m=1}^{k+n} \frac{1}{m} \right]$$
 (41)

$$+\frac{1}{2}\sum_{r=0}^{n-1}(-1)^r\frac{(n-r-1)!}{r!}\left(\frac{r}{2}\right)^{2r-n}, \text{ the last term is replaced by 0 for } n=0.$$
 (42)

These modified Bessel functions satisfy the modified Bessel's equation,

$$\frac{\mathrm{d}^2 w}{\mathrm{d}z^2} + \frac{1}{z} \frac{\mathrm{d}w}{\mathrm{d}z} - \left(1 + \frac{\nu^2}{z^2}\right) w = 0. \tag{43}$$

The generating function expression of I_n is given by

$$\exp\left[\frac{x}{2}\left(t+\frac{1}{t}\right)\right] = \sum_{n=-\infty}^{\infty} I_n(z)t^n,\tag{44}$$

or by replacing t by $e^{i\theta}$

$$e^{z}\cos\theta = \sum_{n=-\infty}^{\infty} I_{n}(z)e^{\mathrm{i}n\theta}.$$
 (45)

Integrals

$$\int_0^\infty e^{-a^2 x^2} x J_0(bx) dx = \frac{e^{-b^2/(4a^2)}}{2a^2} \quad \text{(Weber)}$$

$$\int_{0}^{\infty} e^{-a^{2}x^{2}} x^{2} J_{1}(bx) dx = \frac{be^{-b^{2}/(4a^{2})}}{4a^{4}}$$
(47)

$$\int_{0}^{\infty} e^{-a^{2}x^{2}} x J_{0}(bx) dx = \frac{e^{-b^{2}/(4a^{2})}}{2a^{2}} \quad \text{(Weber)}$$

$$\int_{0}^{\infty} e^{-a^{2}x^{2}} x^{2} J_{1}(bx) dx = \frac{be^{-b^{2}/(4a^{2})}}{4a^{4}}$$

$$\int_{0}^{\infty} e^{-a^{2}x^{2}} x J_{n}^{2}(bx) dx = \frac{e^{-b^{2}/(2a^{2})}}{2a^{2}} I_{n} \left(\frac{b^{2}}{2a^{2}}\right)$$
(48)