12.3.5 Ion-atom scattering; general binary collision dynamics

A collision between two particles that have an interaction potential V(r) and which collide with an *impact parameter p* is shown in Figure 12.6. It is desired to find the orbit of two particles in an elastic collision and to relate the interaction potential to the differential cross section $\sigma(E, \theta)$.

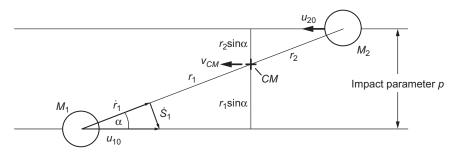


FIGURE 12.6: Geometry for derivation of elastic collision between an energetic ion and a stationary atom, interacting by a potential V(r).

In the system considered, a particle mass M_1 is moving initially with kinetic energy E toward an initially stationary particle mass M_2 . The center of mass (CM) is located on the line joining the two masses at a distance

$$r_1 = \frac{M_2}{M_1 + M_2} r$$
 and $r = r_1 + r_2$ (12.28)

from the mass M_1 . The initial velocity of particle 1 in the center-of-mass system, u_{10} , is decomposed into two perpendicular components \dot{r}_1 and \dot{S}_1 , such that $\vec{u}_{10} = \dot{\vec{S}}_1 + \dot{\vec{r}}_1$. The line between the particles makes an angle α with the initial direction of the particles in the CM system. Only the initial kinetic energy in the CM system is convertible to potential energy, and this is written

$$E_{CM} = \frac{1}{2}(M_1 + M_2)v_{CM}^2 , \qquad (12.29)$$

and using Equation (12.3):

$$E_{CM} = \frac{M_1}{M_1 + M_2} E = \frac{M_1}{M_1 + M_2} \left[\frac{M_1 v_{10}^2}{2} \right]$$
 (12.30)

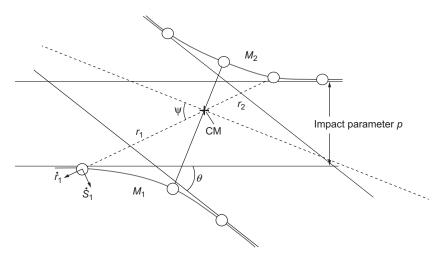


FIGURE 12.7: Geometry during the collision in center-of-mass coordinates.

where *E* is the initial kinetic energy of particle 1 in the laboratory frame. Now using conservation of energy and angular momentum during the collision, it is possible to derive a relationship between the scattering angle in the center of mass θ and the impact parameter p. The trajectory of the particles as they interact and are deflected by angle θ are shown in Figure 12.7.

Conservation of energy

As the two energetic particles approach each other, they convert kinetic energy into potential energy, V(r), so that at the distance of closest approach, the kinetic energy is minimized. Conservation of energy for the system is

$$\begin{split} E_{CM} &= V(r) + \frac{1}{2} M_1 u_{10}^2 + \frac{1}{2} M_2 u_{20}^2 = V(r) + \frac{1}{2} M_1 (\dot{r}_1^2 + \dot{S}_1^2) \\ &+ \frac{1}{2} M_2 (\dot{r}_2^2 + \dot{S}_2^2) \quad . \end{split} \tag{12.31}$$

The tangential speed \dot{S} is equal to $r\dot{\psi}$, so Equation (12.31) is

$$E_{CM} = V(r) + \frac{1}{2}M_1(\dot{r}_1^2 + r_1^2\dot{\psi}^2) + \frac{1}{2}M_2(\dot{r}_2^2 + r_2^2\dot{\psi}^2) , \qquad (12.32)$$

and using the definition of the energy of the center of mass,

$$E_{CM} = V(r) + \frac{1}{2} \left(\frac{M_1 M_2}{M_1 + M_2} \right) (\dot{r}^2 + r^2 \dot{\psi}^2)$$
 (12.33)

along with

$$\dot{r}^2 = \left(\frac{dr}{d\psi}\frac{d\psi}{dt}\right)^2 = \left(\frac{dr}{d\psi}\right)^2\dot{\psi}^2\tag{12.34}$$

yields

$$E_{CM} = V(r) + \frac{1}{2} \left(\frac{M_1 M_2}{M_1 + M_2} \right) \left(\left[\frac{dr}{d\psi} \right]^2 + r^2 \right) \dot{\psi}^2 \quad . \tag{12.35}$$

Conservation of angular momentum

The angular momentum of a mass M about an axis is $\vec{r} \times (M\vec{u}) = rMu \sin \alpha$ = $rM\dot{S}$, where r is the distance between the particle and the axis, u is the velocity, and α is the angle between r and u. Then, the total angular momentum in the CM system is equal to

$$L = r_1 M_1 u_{10} \sin \alpha + r_2 M_2 v_{CM} \sin \alpha . \qquad (12.36)$$

The tangential velocity \dot{S} is equal to $r\dot{\psi}$; thus,

$$L = L_1 + L_2 = M_1 \dot{S}_1 r_1 + M_2 \dot{S}_2 r_2 = M_1 r_1^2 \dot{\psi} + M_2 r_2^2 \dot{\psi} . \tag{12.37}$$

Using the definition of CM velocity in Equation (12.36) and equating to Equation (12.37)

$$L = \frac{M_1 M_2}{M_1 + M_2} v_{10} (r_1 \sin \alpha + r_2 \sin \alpha) = \frac{M_1 M_2}{M_1 + M_2} v_{10} p$$
$$= M_1 r_1^2 \dot{\psi} + M_2 r_2^2 \dot{\psi}$$
(12.38)

where p is the impact parameter. Using Equations (12.28) in Equation (12.38), we obtain an equation for $\dot{\psi}$

$$\dot{\psi} = v_{10} p / r^2 \quad , \tag{12.39}$$

which can be eliminated in Equation (12.38)

$$E_{CM} = V(r) + \frac{1}{2} \left(\frac{M_1 M_2}{M_1 + M_2} \right) \left(\left[\frac{dr}{d\psi} \right]^2 + r^2 \right) \left(\frac{v_{10} p}{r^2} \right)^2 . \tag{12.40}$$

Using Equation (12.30) to eliminate v_{10} ,

$$\frac{d\psi}{dr} = \frac{p}{r^2} \frac{1}{\left[1 - \frac{V(r)}{E_{CM}} - \frac{p^2}{r^2}\right]^{1/2}} . \tag{12.41}$$

When the particles are at their closest, $r = r_0$ and $\psi = \pi/2$, and when $r \to \infty$, $\psi \rightarrow \theta/2$. Integrating Equation (12.41) between these limits:

$$\int_{\theta/2}^{\pi/2} d\psi = \frac{\pi}{2} - \frac{\theta}{2} \quad , \tag{12.42}$$

produces the Classical Scattering Integral

$$\theta = \pi - 2 \int_{r_0}^{\infty} \frac{p dr}{r^2 \left[1 - \frac{V(r)}{E_{CM}} - \frac{p^2}{r^2} \right]^{1/2}} . \tag{12.43}$$

Equation (12.43) relates the impact parameter p to the scattering angle in the center of mass θ . At the distance of closest approach, $dr/d\psi = 0$ and

$$1 - \frac{V(r_o)}{E_{CM}} - \frac{p^2}{r_o^2} = 0 \quad . \tag{12.44}$$

Equation (12.44) can be solved for the distance of closest approach as a function of impact parameter. In a head-on collision, p = 0 and Equation (12.44) reduces to $V(r_o) = E_{CM}$, which, using Equation (12.27), yields the distance of closest approach in a head-on collision.