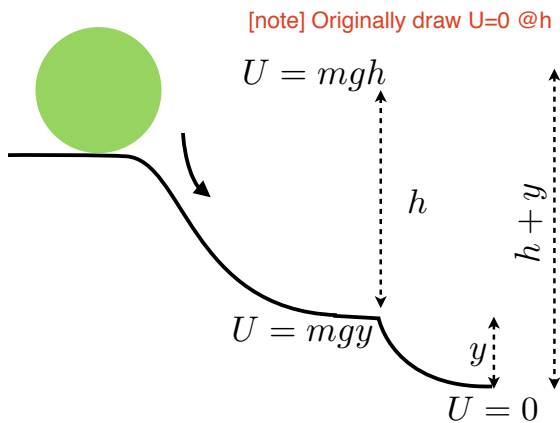


What is Potential?

- Recall from gravity, the change in potential energy of an object above the Earth is: $\Delta U_G = mgh$



- This is a measure of the amount of energy (work) that was required to push the ball up the hill to height h .
- This is also equal to the kinetic energy if the ball rolled back down the hill to $h=0$
- But what if h does not measure entire height of the hill?

$$\Delta U_G = mg(h + y) - mgy = mgh$$

- The mass doesn't matter, only the change in height matters, lets divide out the mass:

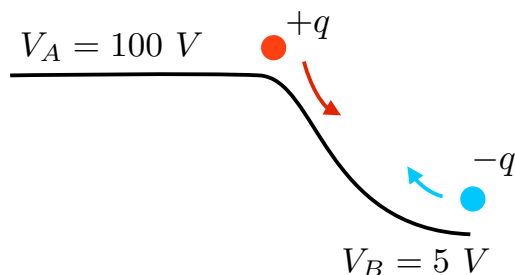
$$\frac{U_G}{m} = gh = V_G$$

This is the "**Gravitational Potential**": it is a measure of the amount of work required that does not require knowing the mass.

- The change in potential does not care where $U=0$ is, only the difference matters.
- If $\Delta h = 0$ then the change in potential is zero, for any mass.

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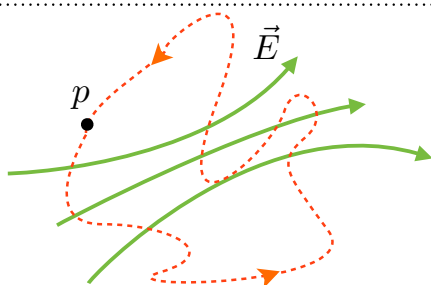
Electric potential works the exact same:



"Electrostatic Potential"

$$\frac{U_{el}}{e} = \frac{Q}{4\pi\epsilon_0 r} = V$$

- You can think of (+) charges as rolling down the potential hill to lower V . (-) charges go up the hill to higher potentials.



- Suppose I am given an E-field and a point p
- Now starting at p , I go along an arbitrary path and at the end, go back to point p
- Forces are conservative \rightarrow work from $p \rightarrow p$ is zero

$$V_p - V_p = \oint_A^A \vec{E} \cdot d\vec{l} = 0$$

integrate over closed loop

We now have 4 important equations for electrostatics

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Important Electrostatic Equations:

$$(1) \oint_S \vec{E} \cdot d\vec{A} = \frac{Q_{enc}}{\epsilon_0}$$

Gives E-field when charge distribution is known

$$(2) V_p = \int_p^\infty \vec{E} \cdot d\vec{l}$$

Relates potential to E-field along path from $p \rightarrow \infty$

$$(3) V_A - V_B = \int_A^B \vec{E} \cdot d\vec{l} \quad \text{or} \quad \Delta V = V_B - V_A = - \int_A^B \vec{E} \cdot d\vec{l}$$

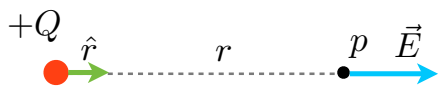
The change in potential from A \rightarrow B is related to the amount of E-field along the path from A \rightarrow B

$$(4) \oint \vec{E} \cdot d\vec{l} = 0 \quad \text{The work done along any closed path is zero.}$$

- If we know the E-field everywhere then (2) tells use that we know potential everywhere.
- But given the potential V everywhere, can we find the E-field?
- From (2) we see that the potential V is the integral of the E-field.
 - So E-field must be the derivative of the potential
 - Must worry about whether there is a minus sign or not.

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- Single point charge:



$$\vec{E} = \frac{Q}{4\pi\epsilon_0 r^2} \hat{r}$$

← vector function

- We have already derived the potential: $V_p = \frac{Q}{4\pi\epsilon_0 r}$ ← scalar function

- Want to show that E is derivative of V:

$$\frac{dV}{dr} = -\frac{Q}{4\pi\epsilon_0 r^2}$$

- We want a vector equation, so multiply both sides by \hat{r} :

$$\frac{dV}{dr} \hat{r} = -\frac{Q}{4\pi\epsilon_0 r^2} \hat{r} = -\vec{E}$$

$$\vec{E} = -\frac{dV}{dr} \hat{r}$$



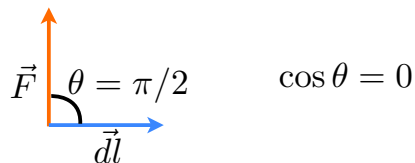
If you know the potential everywhere in space then you can find the E-field.

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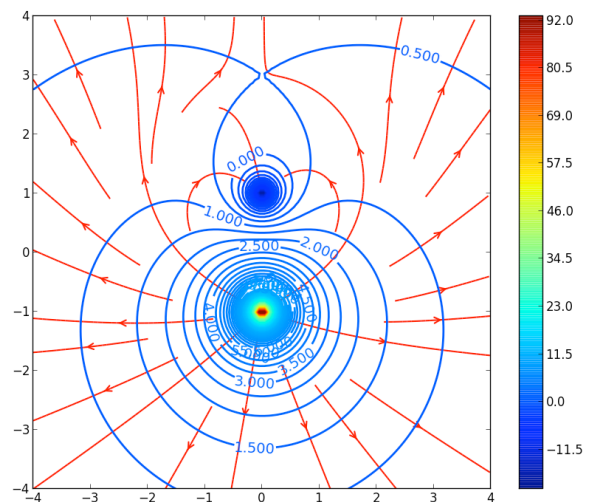
- Recall: E-field lines are always perpendicular to equipotential surfaces:

- Suppose you have a charge q and you always move perpendicular to E-field:

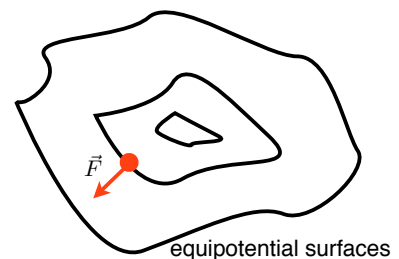
$$\int_A^B q \vec{E} \cdot d\vec{l} = \int_A^B \vec{F} \cdot d\vec{l} = \int_A^B F dl \cos \theta$$



$$\int_A^B \vec{F} \cdot d\vec{l} = W_{A \rightarrow B} = 0 \quad - \text{Work from A} \rightarrow \text{B is zero} \rightarrow \text{potential is constant}$$



A charge released on an equipotential surface will always move perpendicular to the surface, the the direction of the E-field if (+), or opposite direction if (-).



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Deeper Connection Between E-field and V:

- Suppose I am at point P in space with potential V_p and some E-field
- Now I take a step **only** in the x-direction:
 - If potential does not change \rightarrow no E-field in the x-direction.
 - E-field (if any) must be perpendicular since $V=\text{constant}$
 - If potential does change, then the electric field in the x-direction can be written as

$$|E_x| = \left| \frac{\Delta V}{\Delta x} \right|_{y,z=\text{const}} \left[\frac{V}{m} \right] \quad \text{Electric field is the change in potential over some distance}$$

- Same in other directions:

$$|E_y| = \left| \frac{\Delta V}{\Delta y} \right|_{x,z=\text{const}} \quad |E_z| = \left| \frac{\Delta V}{\Delta z} \right|_{x,y=\text{const}}$$

- In x,y,z coordinates, the E-field can be written as: $\vec{E} = E_x \hat{x} + E_y \hat{y} + E_z \hat{z}$

$$\vec{E} = - \left(\frac{\partial V}{\partial x} \hat{x} + \frac{\partial V}{\partial y} \hat{y} + \frac{\partial V}{\partial z} \hat{z} \right) = -\nabla V$$

$$\frac{\partial}{\partial x} \rightarrow y, z = \text{const}$$

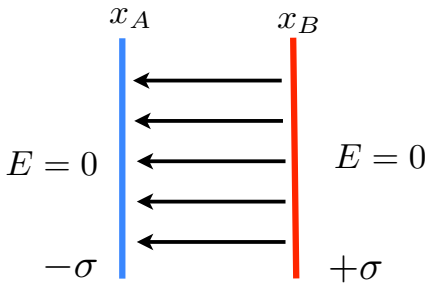
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Ex. $V(x) = 10^5 x$ $x \in [0, 0.01 \text{ m}]$ $V=0$ everywhere else

- What is E?

$$\vec{E} = -\frac{\partial V}{\partial x} = -10^5 \hat{x} \quad x \in [0, 0.01 \text{ m}] \quad E_y = E_z = 0$$

- This corresponds to the example of two charged conducting planes



- Lets write out what the E-field is

- Same as previous example $|E| = \frac{\sigma}{\epsilon_0} = 10^5 \frac{\text{V}}{\text{m}}$

- What is potential difference?

$$V_A - V_B = \int_{x_A}^{x_B} \vec{E} \cdot d\vec{x} = -10^5 \int_{x_A}^{x_B} \hat{x} \cdot d\vec{x}$$

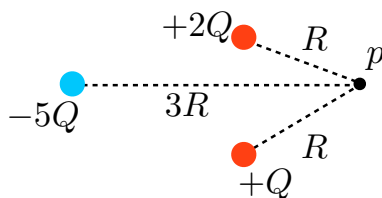
both in same direction

$$V_A - V_B = -10^5 \int_{x_A}^{x_B} dx = -10^5 (x_B - x_A)$$

- Since distance between A & B is 0.01m: $V_A - V_B = -1000 \text{ V}$ Potential grows linearly from A -> B

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Ex. Three point charges:



- What is the potential at point P?

- Recall: $V_p = \frac{Q}{4\pi\epsilon_0 r}$

- We can always use superposition to add up individual potential terms

$$V_p = V_p^{+Q} + V_p^{+2Q} + V_p^{-5Q}$$

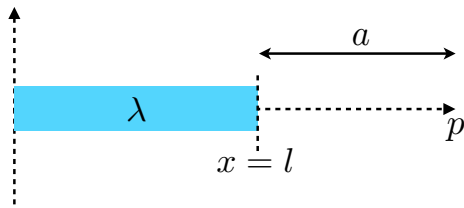
- Problem is simple now:

$$V_p^{+Q} = \frac{Q}{4\pi\epsilon_0 R} \quad V_p^{+2Q} = \frac{2Q}{4\pi\epsilon_0 R} \quad V_p^{-5Q} = \frac{-5Q}{4\pi\epsilon_0 (3R)}$$

$$V_p = \frac{4}{3} \frac{Q}{4\pi\epsilon_0 R}$$

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Ex. Charged rod:



- What is the potential at point P from rod with length l and charge density λ .

- We can calculate the E-field and use

$$V_p = \int_p^\infty \vec{E} \cdot d\vec{l}$$

but we are not given E and this is too much work.

- Instead we know that charged rod is just a collection of small individual charges dq .

- So lets add up the potential from each little charge dq to get answer:

$$dV_p = \frac{dq}{4\pi\epsilon_0 r} = \frac{\lambda dx}{4\pi\epsilon_0 [(a+l) - x]}$$

$$V_p = \frac{\lambda}{4\pi\epsilon_0} \int_0^l \frac{dx}{[(a+l) - x]} = \frac{\lambda}{4\pi\epsilon_0} [\log(a+l) - \log(a)]$$

$$V_p = \frac{\lambda}{4\pi\epsilon_0} \log\left(\frac{a+l}{a}\right)$$