### EPR, Bell's Inequality, and the E91 Protocol

Gaurav S. Kasbekar

Dept. of Electrical Engineering

IIT Bombay

#### References

- C. Bernhardt: Chapter 5
- M. Nielsen and I. Chuang: Section 2.6

#### Copenhagen Interpretation

Named after city where Neils Bohr lived According to Copenhagen interpretation: ■ Before a measurement is made, a quantum system exists in a superposition of multiple states o e.g., particle might be considered to be in multiple locations simultaneously ☐ The act of measuring forces the system to "collapse" into a single, definite state • Specific state it collapses into is probabilistic, i.e., it is governed by probabilities, not determined with certainty Observer or measuring apparatus plays crucial role in determining outcome of quantum event Observer not merely passively observing, but actively influencing system's state Some physicists, e.g., Einstein and Schrodinger did not subscribe to this model, in particular, to: ☐ interpretation of states jumping with given probabilities to basis states, and concept of action at a distance in entanglement They thought there should be a better model using: ☐ hidden variables and local realism (details later) John Stewart Bell devised an ingenious test, which could distinguish between above two models Later, several experiments based on Bell's result were conducted, which showed that:

☐ Einstein and Schrodinger's view was wrong and Copenhagen interpretation is correct

#### Local Realism

ı	A concept in physics that combines two ideas:
	☐ locality
	□ realism
١	Locality:
	an object's properties are determined only by its immediate environment
	$\square$ any influence on that object can only travel at or below speed of light
	distant objects cannot instantaneously affect each other
	Realism:
	objects have definite, pre-existing properties, regardless of whether they are being observed or measured
	Einstein, Boris Podolsky, and Nathan Rosen (EPR) published a paper, which stated that:
	special theory of relativity implied that information could not travel faster than speed of light,
	but instantaneous action at a distance (as in entangled particles) would mean that information could be sent from Alice to Bob instantaneously
	This problem known as <i>EPR paradox</i>

#### Hidden Variables

- Hidden variables:
  - □ properties of particles, to which we do not have access, which determine the outcomes of quantum measurements even before the measurements are actually made
- EPR proposed that:
  - I quantum mechanics might be incomplete because it does not account for these hidden variables,
  - ☐ quantum mechanics only provides probabilistic predictions, while a more complete theory (including hidden variables) could offer deterministic outcomes

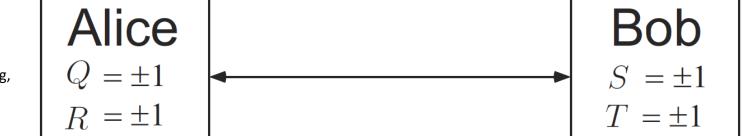
#### Versions of Bell's Inequality

- There are several versions of Bell's inequality
- We discuss two versions

### Bell's Inequality: Version 1 Clauser, Horne, Shimony, and Holt (CHSH Inequality)

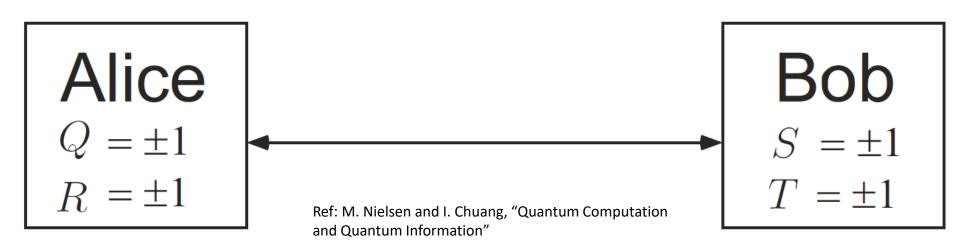
#### **CHSH** Inequality

- Imagine we perform following experiment, illustrated in Fig.
- Charlie prepares two particles:
  - ☐ sends one particle to Alice, and second particle to Bob
- Once Alice receives her particle, she performs a measurement on it
- She has available two different measurement apparatuses, so she could choose to do one of two different measurements
- These measurements are of physical properties which we shall label:
  - $\square$   $P_O$  and  $P_R$ , respectively
- Alice doesn't know in advance which measurement she will choose to perform:
  - ☐ when she receives particle, she randomly decides which measurement to perform
- For simplicity, assume that the measurements can each have one of two outcomes,  $+\ 1\ \text{or}\ -1$
- Suppose Alice's particle has a value:
  - $\square$  *Q* for the property  $P_Q$
  - $\square$  *R* for the property  $P_R$



Ref: M. Nielsen and I. Chuang, "Quantum Computation and Quantum Information"

- Similarly, Bob is capable of measuring one of two properties,  $P_S$  or  $P_T$ , each taking value  $\pm 1$  or  $\pm 1$
- Bob waits until he has received the particle and then:
  - ☐ randomly decides which measurement to perform
- Timing of experiment arranged so that Alice and Bob do their measurements at the same time
- Therefore, according to classical model, measurement which Alice performs cannot disturb result of Bob's measurement (or vice versa):
  - ☐ since physical influences cannot propagate faster than light



Note that:

1) 
$$QS + RS + RT - QT = (Q + R)S + (R - Q)T$$

• Since  $R, Q \in \{-1, +1\}$ , it follows that:

$$\square (Q+R)S = 0 \text{ or } (R-Q)T = 0$$

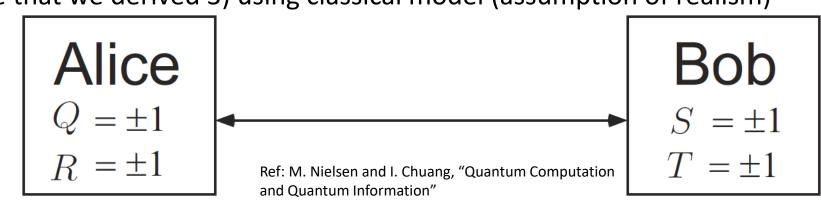
In either case, from 1), we get that:

$$2) \quad QS + RS + RT - QT = \pm 2$$

- Let p(q, r, s, t) be probability that, before the measurements are performed, system is in state where Q = q, R = r, S = s, T = t
  - ☐ these probabilities may depend on how Charlie performs his preparation, and on experimental noise
- By 2), we get:

3) 
$$E(QS) + E(RS) + E(RT) - E(QT) = \sum_{q,r,s,t} p(q,r,s,t) (qs + rs + rt - qt) \le 2$$

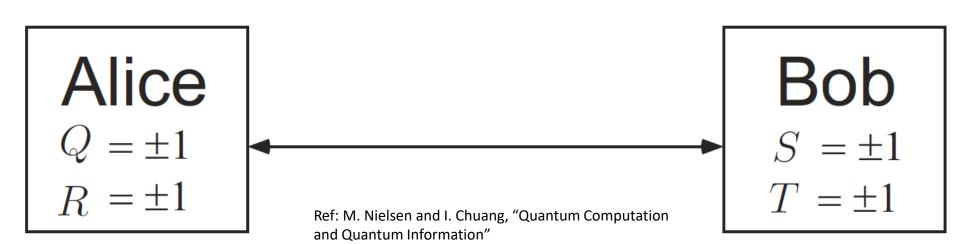
- Inequality in 3) known as CHSH inequality
- Note that we derived 3) using classical model (assumption of realism)



Recall: CHSH inequality:

3) 
$$E(QS) + E(RS) + E(RT) - E(QT) \le 2$$

- By repeating experiment many times, Alice and Bob can determine each quantity on LHS of 3)
- Thus, they can check to see whether it is obeyed in a real experiment
- Next, we calculate LHS of 3) using quantum mechanical model

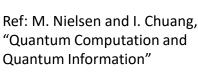


- Recall: CHSH inequality:

  3)  $E(QS) + E(RS) + E(RT) E(QT) \le 2$ Charlie prepares a quantum system of two qubits in states
- Suppose Charlie prepares a quantum system of two qubits in state:
- $|\psi\rangle = \frac{|01\rangle |10\rangle}{\sqrt{2}}$ He passes: (Contd.)
- He passes:

   ☐ first qubit to Alice, and
  - second qubit to BobThey perform measurements of following observables:
    - $Q = Z_1, S = \frac{-Z_2 X_2}{\sqrt{2}}$   $R = X_1, T = \frac{Z_2 X_2}{\sqrt{2}}$
    - Note: subscript 1 (respectively, 2) denotes qubit of Alice (respectively, Bob))

      Exercise: Show that eigenvalues of each of the above observables Q, R, S, and T are  $\pm 1$
    - **Exercise**: Show that eigenvalues of each of the above observables Q, R, S, and I are  $\pm 1$  recall: eigenvalues of observables are measurement outcomes; hence, the possible outcomes are  $\pm 1$  for each observable
- **Exercise:** Show that average values of above observables are:
- Thus:
- 4)  $\langle QS \rangle + \langle RS \rangle + \langle RT \rangle \langle QT \rangle = 2\sqrt{2}$
- Note that 3) and 4), which were obtained using classical and quantum mechanical model, respectively, contradict each other





 $|\psi\rangle = \frac{|01\rangle - |10\rangle}{\sqrt{2}}$   $Q = Z_1, S = \frac{-Z_2 - X_2}{\sqrt{2}}$   $R = X_1, T = \frac{Z_2 - X_2}{\sqrt{2}}$ 

Recall:

# CHSH Inequality (contd.)

• We now outline the proof of  $\langle QS \rangle = \frac{1}{\sqrt{2}}$ , which is part of above exercise

• 
$$Q|0\rangle = Z|0\rangle = |0\rangle$$
 and  $Q|1\rangle = Z|1\rangle = -|1\rangle$ 

- It can be shown that:  $S|0\rangle = \frac{-|0\rangle |1\rangle}{\sqrt{2}}$  and  $S|1\rangle = \frac{-|0\rangle + |1\rangle}{\sqrt{2}}$
- Hence:  $QS|\psi\rangle = QS\left(\frac{|01\rangle |10\rangle}{\sqrt{2}}\right) = \frac{Q|0\rangle \otimes S|1\rangle Q|1\rangle \otimes S|0\rangle}{\sqrt{2}}$ , which simplifies to:  $QS|\psi\rangle = \frac{-1}{2}[|00\rangle |01\rangle + |10\rangle + |11\rangle]$
- So  $\langle QS \rangle = \langle \psi | QS | \psi \rangle = -\left(\frac{\langle 01| \langle 10|}{2\sqrt{2}}\right) [|00\rangle |01\rangle + |10\rangle + |11\rangle] = \frac{1}{\sqrt{2}}$

Chuang, on and 
$$Q = \pm 1$$
  $R = \pm 1$   $T = \pm 1$ 

 $\square Q = Z_1, S = \frac{-Z_2 - X_2}{\sqrt{2}}$   $\square R = X_1, T = \frac{Z_2 - X_2}{\sqrt{2}}$ • Note that we need to

 $\square |\psi\rangle = \frac{|01\rangle - |10\rangle}{\sqrt{2}}$ 

Recall:

## CHSH Inequality (contd.)

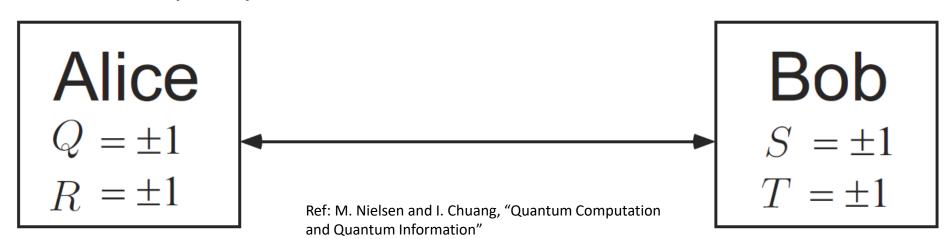
- Note that we need to calculate E(QS), which is the average value of the product of two scalars
- But we equated it to  $\langle QS \rangle = \langle \psi | QS | \psi \rangle = \langle \psi | Q \otimes S | \psi \rangle$
- Why does the tensor product  $Q \otimes S$  correspond to the measurement outcome qs, which is a product of two scalars?
- Recall the spectral decompositions:  $\square \ Q = \sum_{a} q P_{a} \text{ and }$ 
  - $\Box S = \sum_{S} S P_{S}$   $\Box \text{ where } P_{S} \text{ is projector onto eigenspace of } Q \text{ with eigenvalue } q$
  - lacksquare where  $P_q$  is projector onto eigenspace of Q with eigenvalue q
  - $\square Q \otimes S = (\sum_{a} q P_{a}) \otimes (\sum_{s} s P_{s}) = \sum_{a,s} (qs) (P_{a} \otimes P_{s})$
  - Hence,  $Q \otimes S$  is an observable with corresponding measurement outcomes qs

Ref: M. Nielsen and I. Chuang, "Quantum Computation and Quantum Information"

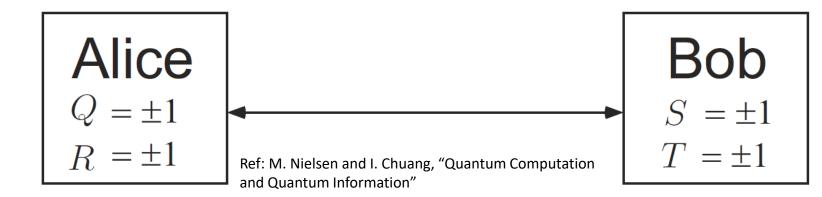
So:

 $\begin{array}{c|c} S & \longrightarrow & DOI \\ S & = \pm \\ T & = \pm \\ \end{array}$ 

- Recall:
  - 3)  $E(QS) + E(RS) + E(RT) E(QT) \le 2$ 
    - obtained using classical model
  - 4)  $\langle QS \rangle + \langle RS \rangle + \langle RT \rangle \langle QT \rangle = 2\sqrt{2}$ 
    - o obtained using quantum mechanical model
- Several experiments have been conducted to check as to which one of 3) and 4) is true
- Results were strongly in favor of quantum mechanical prediction
   4)
- So the CHSH inequality 3) is not obeyed by Nature
- It means that assumptions that went into the derivation of the CHSH inequality must be incorrect



- Recall:
  - 3)  $E(QS) + E(RS) + E(RT) E(QT) \le 2$ 
    - obtained using classical model
  - 4)  $\langle QS \rangle + \langle RS \rangle + \langle RT \rangle \langle QT \rangle = 2\sqrt{2}$ 
    - obtained using quantum mechanical model
- There are two assumptions made in various proofs of 3) which are questionable: realism and locality
- ☐ Realism:
  - assumption that the physical properties  $P_Q$ ,  $P_R$ ,  $P_S$ , and  $P_T$  have definite values, Q, R, S, and T, which exist independent of observation
  - We used this in above derivation to write  $E(QS) + E(RS) + E(RT) E(QT) = \sum_{q,r,s,t} p(q,r,s,t) (qs + rs + rt qt)$ , where p(q,r,s,t) is probability that before the measurements are performed, system is in state where Q = q, R = r, S = s, T = t
  - Note: If we try to define p(q,r,s,t) to be probability that *after* measurements are performed, the outcomes will be Q=q,R=r,S=s,T=t, then such a joint distribution does not exist in general since either  $P_Q$  or  $P_R$  (respectively,  $P_S$  or  $P_T$ ) is measured, but not both
- ☐ Locality:
  - assumption that Alice performing her measurement does not influence the result of Bob's measurement
  - ☐ If this assumption holds, then we can write:
    - $\square$  E(QS) = E(Q)E(S), E(RS) = E(R)E(S), E(RT) = E(R)E(T), and <math>E(QT) = E(Q)E(T)
    - $\Box \text{ So } E(QS) + E(RS) + E(RT) E(QT) = E(S) \big( E(Q) + E(R) \big) + E(T) \big( E(R) E(Q) \big) \leq |E(Q) + E(R)| + |E(Q) E(R)| \leq 2 \max(|E(Q)|, |E(R)|) \leq 2$
- These two assumptions together constitute assumption of local realism:
  - must be dropped



#### Bell's Inequality: Version 2

#### **Entangled Qubits in Different Bases**

Consider two entangled qubits in state:

$$\square_{\frac{1}{\sqrt{2}}} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \frac{1}{\sqrt{2}} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

- Now, consider a new orthonormal basis  $(|b_0\rangle, |b_1\rangle)$ :
  - $\square$  where the components of  $|b_0\rangle$  and  $|b_1\rangle$  are real numbers
- Claim:

$$\square_{\frac{1}{\sqrt{2}}} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \frac{1}{\sqrt{2}} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \frac{1}{\sqrt{2}} |b_0\rangle \otimes |b_0\rangle + \frac{1}{\sqrt{2}} |b_1\rangle \otimes |b_1\rangle$$

• Proof:

$$\square$$
 Let  $|b_0\rangle = \begin{vmatrix} a \\ b \end{vmatrix}$  and  $|b_1\rangle = \begin{vmatrix} c \\ d \end{vmatrix}$ , where  $a, b, c, d \in \mathbb{R}$ 

$$\Box \operatorname{Then} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = a \begin{bmatrix} a \\ b \end{bmatrix} + c \begin{bmatrix} c \\ d \end{bmatrix}$$

1) So 
$$\begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \left( a \begin{bmatrix} a \\ b \end{bmatrix} + c \begin{bmatrix} c \\ d \end{bmatrix} \right) \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} a \\ b \end{bmatrix} \otimes \begin{bmatrix} a \\ 0 \end{bmatrix} + \begin{bmatrix} c \\ d \end{bmatrix} \otimes \begin{bmatrix} c \\ 0 \end{bmatrix}$$

2) Similarly, 
$$\begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} a \\ b \end{bmatrix} \otimes \begin{bmatrix} 0 \\ b \end{bmatrix} + \begin{bmatrix} c \\ d \end{bmatrix} \otimes \begin{bmatrix} 0 \\ d \end{bmatrix}$$

Adding 1) and 2), we get: 
$$\begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} a \\ b \end{bmatrix} \otimes \begin{bmatrix} a \\ b \end{bmatrix} + \begin{bmatrix} c \\ d \end{bmatrix} \otimes \begin{bmatrix} a \\ b \end{bmatrix} + \begin{bmatrix} c \\ d \end{bmatrix} \otimes \begin{bmatrix} a \\ b \end{bmatrix} + \begin{bmatrix} c \\ d \end{bmatrix} \otimes \begin{bmatrix} a \\ b \end{bmatrix} + \begin{bmatrix} c \\ d \end{bmatrix} \otimes \begin{bmatrix} a \\ b \end{bmatrix} + \begin{bmatrix} c \\ d \end{bmatrix} \otimes \begin{bmatrix} a \\ b \end{bmatrix} + \begin{bmatrix} c \\ d \end{bmatrix} \otimes \begin{bmatrix} a \\ b \end{bmatrix} + \begin{bmatrix} c \\ d \end{bmatrix} \otimes \begin{bmatrix} a \\ b \end{bmatrix} + \begin{bmatrix} c \\ d \end{bmatrix} \otimes \begin{bmatrix} a \\ b \end{bmatrix} + \begin{bmatrix} c \\ d \end{bmatrix} \otimes \begin{bmatrix} a \\ b 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#### Bases Used for Bell's Inequality

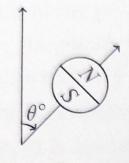
- For proving Bell's inequality, we use three different bases
- These correspond to rotating our measuring device through  $0^{\circ}$ ,  $120^{\circ}$ , and  $240^{\circ}$
- Above three bases denoted by:

$$\square$$
  $(|\uparrow\rangle, |\downarrow\rangle), (|\downarrow\rangle, |\uparrow\rangle), and  $(|\swarrow\rangle, |\nearrow\rangle)$ , respectively$ 

- Assume that qubits are encoded in spin of particles

• Recall: basis associated with rotating our apparatus by 
$$\theta$$
:  $\binom{\cos(\theta/2)}{-\sin(\theta/2)}$ ,  $\binom{\sin(\theta/2)}{\cos(\theta/2)}$ )

Hence:



$$\left( \left[ \begin{array}{c} \cos\left(\frac{\theta}{2}\right) \\ -\sin\left(\frac{\theta}{2}\right) \end{array} \right], \left[ \begin{array}{c} \sin\left(\frac{\theta}{2}\right) \\ \cos\left(\frac{\theta}{2}\right) \end{array} \right] \right)$$

(a) Measurement angle

(b) Basis

#### Classical Explanation of Entanglement

- Consider two entangled qubits:
  - one with Alice and other with Bob
  - $\square$  state of the qubits is  $\frac{1}{\sqrt{2}} | \uparrow \rangle | \uparrow \rangle + \frac{1}{\sqrt{2}} | \downarrow \rangle | \downarrow \rangle$
- Recall:
  - □ when Alice and Bob perform a measurement of their qubits in basis  $(|\uparrow\rangle, |\downarrow\rangle)$ , both get same answer (either 0 or 1)
- Recall: for proving Bell's inequality:
  - $\square$  we rotate our measuring device through  $0^{\circ}$ ,  $120^{\circ}$ , and  $240^{\circ}$
  - $\square$  we measure spin of particles in the bases  $(|\uparrow\rangle, |\downarrow\rangle), (|\downarrow\rangle, |\uparrow\rangle)$ , and  $(|\swarrow\rangle, |\nearrow\rangle)$
- Classical theory states:
  - ☐ there is a definite outcome of the measurement in each basis that is already determined before we perform the measurement (realism)
- Quantum theory states:
  - outcome of measurement not determined until the time we perform the measurement

#### Bell's Inequality

- We generate a stream of *n* pairs of qubits:
  - ☐ from each pair, one sent to Alice and other to Bob
- Each pair of qubits is in state:

$$\Box \frac{1}{\sqrt{2}} |\uparrow\rangle |\uparrow\rangle + \frac{1}{\sqrt{2}} |\downarrow\rangle |\downarrow\rangle$$

- Alice:
  - $\square$  measures each qubit in direction  $0^{\circ}$ ,  $120^{\circ}$ , or  $240^{\circ}$ 
    - for each qubit, direction of measurement selected uniformly at random, independently of other qubits
  - does not keep track of direction she chose for a qubit
  - $\square$  records the outcome (0 or 1) of each measurement
- Subsequently, Bob does the same for his qubits:
  - ☐ his directions of measurements selected independently of Alice
- In this way, Alice and Bob generate strings of length n each of 0s and 1s
- What is the fraction of bits of these strings in which Alice and Bob are in agreement?
- Bell realized that quantum mechanics model and classical model gave different numbers for the answer

#### **Answer of Quantum Mechanics**

- Consider a pair of qubits, which is in state:
  - $\Box \quad \tfrac{1}{\sqrt{2}} \, | \, \! \uparrow \rangle | \, \! \uparrow \rangle + \tfrac{1}{\sqrt{2}} \, | \, \! \downarrow \rangle | \, \! \downarrow \rangle$
- Recall: if Alice and Bob select same direction of measurement (i.e., same basis), their outcomes are same
  - Consider case where they select different directions of measurement
- E.g.: suppose:
  - $\Box$  Alice chooses  $(|\searrow\rangle, |\nwarrow\rangle)$
  - $\square$  Bob chooses  $(|\checkmark\rangle, |\nearrow\rangle)$
- Entangled state can be written as:
- When Alice performs measurement:
  - she gets 0 and state jumps to  $| \Sigma \rangle | \Sigma \rangle$  w.p.  $\frac{1}{2}$
  - $\square$  she gets 1 and state jumps to  $| \nwarrow \rangle | \nwarrow \rangle$  w.p.  $\frac{1}{2}$
- Suppose Alice gets 0 and state jumps to  $|\downarrow\rangle|\downarrow\rangle$
- Now:
- So when Bob performs measurement:
  - $\Box$  gets 0 w.p.  $\frac{1}{4}$  and 1 w.p.  $\frac{3}{4}$
- Thus, when Alice gets 0, states of Alice and Bob agree w.p.  $\frac{1}{4}$
- Exercise: Show that even when Alice gets 1, states of Alice and Bob agree w.p.  $\frac{1}{4}$
- In summary:
  - when Alice chooses  $(|\Sigma\rangle, |\mathbb{N}\rangle)$  and Bob chooses  $(|\mathcal{L}\rangle, |\mathcal{P}\rangle)$ , their bits agree w.p.  $\frac{1}{4}$  and disagree w.p.  $\frac{3}{4}$
- Exercise: Show that in all other cases in which Alice and Bob choose different bases, their bits agree w.p.  $\frac{1}{4}$  and disagree w.p.  $\frac{3}{4}$
- In summary, for a given pair of qubits:
  - $\square$  w.p.  $\frac{1}{3}$ , Alice and Bob choose same basis and their bits agree w.p. 1
  - $\square$  w.p.  $\frac{2}{3}$ , Alice and Bob choose different bases and their bits agree w.p.  $\frac{1}{4}$
- According to quantum mechanics model, what is the fraction of bits of the strings of Alice and Bob for which the bits are in agreement?
  - $\mathbf{J} = \frac{1}{2}$

#### Classical Answer

- Classical view:
  - ☐ Measurements in all directions determined right from the start
- There are eight possible configurations:
  - **O**000, 001, 010, 011, 100, 101, 110, 111
  - where the three bits in each configuration give the answer (0 or 1) if we measure a qubit in bases  $(|\uparrow\rangle, |\downarrow\rangle), (|\downarrow\rangle, |\uparrow\rangle)$ , and  $(|\swarrow\rangle, |\nearrow\rangle)$ , respectively
- Due to the entanglement between qubits of Alice and Bob:
  - for each pair of qubits, configurations of Alice and Bob are identical
- Table shows all possible outcomes of measurement, where:
  - $\square$  a, b, and c denote  $(|\uparrow\rangle, |\downarrow\rangle)$ ,  $(|\searrow\rangle, |\nwarrow\rangle)$ , and  $(|\swarrow\rangle, |\nearrow\rangle)$ , respectively
  - $\square$  (x, y), where  $x, y \in \{a, b, c\}$ , denotes that Alice (respectively, Bob) measures qubit in basis x (respectively, y)
  - $\square$  Entries in table show whether measurements agree (A) or disagree (D)
- Probabilities that should be assigned to different configurations unknown

Since Alice and Bob choose each of their three bases with equal probabilities, each of the nine pairs of bases

occurs w.p.  $\frac{1}{9}$ 

- Note that each row contains at least five As
  - Hence, overall probability that measurements of Alice and Bob

result in same outcome is at least  $\frac{3}{9}$ 

According to classical model, what is the fraction of bits of the strings of Alice and Bob for which the bits

are in agreement?  $\Box \text{ at least } \frac{5}{2}$ 

Ref: "Quantum Computing fo
Everyone" by C. Bernhardt

	Measurement directions								
Config.	(a,a)	(a,b)	(a,c)	(b,a)	( <i>b</i> , <i>b</i> )	(b,c)	(c,a)	(c,b)	(c,c)
000	A	A	A	A	A	A	A	A	A
001	A	A	D	A	A	D	D	D	A
010	A	D	A	D	A	D	A	D	A
011	A	D	D	D	A	A	D	A	A
100	A	D	D	D	A	A	D	A	A
101	A	D	A	D	A	D	A	D	A
110	A	A	D	A	A	D	D	D	A
111	A	A	A	A	A	A	A	A	A

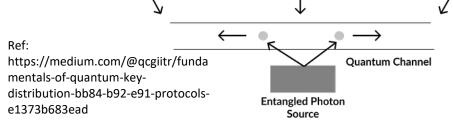
#### Summary of Bell's Result

- Recall above question:
  - ☐ What is the fraction of bits of the strings in which Alice and Bob are in agreement?
- Answer of quantum mechanics:  $\frac{1}{2}$
- Classical answer: at least  $\frac{5}{9}$
- Thus, Bell's test gives us a way to distinguish between the two theories
- Above experiment has been performed by several researchers
- Outcomes have always been in agreement with quantum mechanics

- E91 QKD Protocol
  In 1991, Artur Ekert proposed QKD protocol based on entangled qubits used in Bell's test
- We generate a stream of 3n pairs of qubits:
  - ☐ from each pair, one sent to Alice and other to Bob
- Each pair of qubits is in state:

$$\Box \frac{1}{\sqrt{2}} |\uparrow\rangle |\uparrow\rangle + \frac{1}{\sqrt{2}} |\downarrow\rangle |\downarrow\rangle$$

- Alice and Bob measure each qubit using a random choice of three bases:
  - as in Bell test
- As in BB84 protocol, for each measurement:
  - ☐ Alice and Bob record both outcome and basis that they chose
- After Alice and Bob have made 3n measurements:
  - ☐ They compare the sequences of bases that they chose
  - ☐ This can be done on insecure channel: ONote that they only reveal bases, but not measurement outcomes
  - $lue{}$  They will agree on approximately n of the sequence of bases



Classical Channel

Bob

#### E91 QKD Protocol (contd.)

- For the *n* pairs of qubits for which Alice and Bob selected same basis:
  - ☐ Alice and Bob get same measurement outcome (0 or 1)
- This string of *n* bits will be their key if an eavesdropper Eve is not listening in
- Now Alice and Bob test to find if Eve is present
- If Eve is eavesdropping, she will have to measure the qubits:
  - ☐ When she measures them, the entangled states of Alice and Bob become unentangled
- Alice and Bob look at their bits strings of length 2n each that correspond to times when they chose different bases
- From above Bell inequality calculation, they know that if their states are entangled, in each place, they should agree w.p.  $\frac{1}{4}$
- However, if Eve is measuring one of the qubits, the fraction of bits for which they agree changes
- E.g., if Eve measures a qubit before Alice and Bob have made their measurements:
- Hence, Alice and Bob can use following test to check for presence of Eve:
  - $\Box$ They calculate fraction of the 2n qubits for which their measurement outcomes agree
  - $\square$  If it is close to  $\frac{1}{4}$ , they can conclude that nobody has interfered and use the key

