

$$\mu.12 \quad \mathcal{L} = -\frac{1}{2} H (\partial^2 + M^2) H + \bar{\Psi} (i\not{\partial} - m) \Psi - g \bar{\Psi} \Psi H$$

a)

Introduce sources j and η of H and Ψ

$$\mathcal{L} \rightarrow \mathcal{L} + H j + \bar{\Psi} \eta + \bar{\eta} \Psi,$$

$$\mathcal{L}_0 = -\frac{1}{2} H (\partial^2 + M^2) H + \bar{\Psi} (i\not{\partial} - m) \Psi + H j + \bar{\Psi} \eta + \bar{\eta} \Psi$$

the free Lagrangian

$$W[j] = \frac{\langle 0|0 \rangle_{j, \eta, \bar{\eta}}}{\langle 0|0 \rangle_{j=\eta=\bar{\eta}=0}}$$

$$\langle 0|0 \rangle_{j, \eta, \bar{\eta}} = \int \mathcal{D}\bar{\Psi} \mathcal{D}\Psi \mathcal{D}H \exp(i \int d^4x \mathcal{L}[\bar{\Psi}, \Psi, H, \bar{\eta}, \eta, j])$$

$$= \int \mathcal{D}\bar{\Psi} \mathcal{D}\Psi \mathcal{D}H \exp(i \int d^4x \mathcal{L}_0) \exp(-i \int d^4x g H \bar{\Psi} \Psi)$$

$$\langle 0|0 \rangle_{j=\eta=\bar{\eta}=0} = \int \mathcal{D}\bar{\Psi} \mathcal{D}\Psi \mathcal{D}H \exp(i \int d^4x \mathcal{L}[\bar{\Psi}, \Psi, H, \bar{\eta}=\eta=j=0])$$

$$= \int \mathcal{D}\bar{\Psi} \mathcal{D}\Psi \mathcal{D}H \exp(i \int d^4x \mathcal{L}_0|_{j=\eta=\bar{\eta}=0}) \exp(-i \int d^4x g H \bar{\Psi} \Psi)$$

$$H \rightarrow \frac{1}{i} \frac{\delta}{\delta j} \quad , \quad \Psi \rightarrow \frac{1}{i} \frac{\delta}{\delta \eta} \quad , \quad \bar{\Psi} \rightarrow -\frac{1}{i} \frac{\delta}{\delta \bar{\eta}}$$

$$\exp(-i \int d^4x g H \bar{\Psi} \Psi) \rightarrow \exp \left[-i \int d^4x \frac{1}{i} \frac{\delta}{\delta j(x)} \frac{1}{i} \frac{\delta}{\delta \eta(x)} \frac{1}{i} \frac{\delta}{\delta \bar{\eta}(x)} \right]$$

b) $F = -g \bar{\Psi} \Psi$, no external sources

$$\langle 0|0 \rangle_{j, \bar{\eta}, \eta} = \int \mathcal{D}\bar{\Psi} \mathcal{D}\Psi \mathcal{D}H \exp(i \int d^4x (-\frac{1}{2} H \hat{D}_H H + \bar{\Psi} \hat{D}_\Psi \Psi + F H))$$

$$\text{Shift} \quad H \rightarrow H' = H - i \int d^4y D_F(x-y) F(y)$$

$$F(y) = -g \bar{\Psi}(y) \Psi(y) H(y)$$

$$S_H = \int d^4x \left(-\frac{1}{2} H \hat{D}_H H + F H \right)$$

$$\begin{aligned} \rightarrow S_{H'} &= \int d^4x \left(-\frac{1}{2} H' \hat{D}_H H' + F H' \right) \\ &= \int d^4x \left\{ -\frac{1}{2} H' \hat{D}_H H' + \underbrace{F H}_{\text{Wick rotation}} \right\} \\ &\quad - i \int d^4y F(x) D_F(x-y) F(y) \end{aligned}$$

$$\Rightarrow \langle 0|0 \rangle_{j, n, \bar{n}} = \int \mathcal{D}H' \exp\left(-\frac{i}{2} \int d^4x H' \hat{D}_H H'\right) \times \int \mathcal{D}\bar{\Psi} \mathcal{D}\Psi \exp\left(i \int d^4x (\bar{\Psi} \hat{D}_\psi \Psi - i \int d^4y F(x) D_F(x-y) F(y))\right)$$

$$\left(\begin{array}{l} \text{First factor is Gaussian integral} \\ \int \mathcal{D}H' \exp\left(-\frac{i}{2} \int d^4x H' \hat{D}_H H'\right) = (\det \hat{D}_H)^{-\frac{1}{2}} \\ \text{It gets cancelled out in } W \end{array} \right) \leftarrow \text{Wick rotation}$$

$$= \tilde{N} \int \mathcal{D}\bar{\Psi} \mathcal{D}\Psi \exp\left(i \int d^4x (\bar{\Psi} \hat{D}_\psi \Psi - i \int d^4y F(x) D_F(x-y) F(y))\right)$$

$$\rightarrow W = \frac{1}{\langle 0|0 \rangle_{j, \bar{n}, n}} \int \mathcal{D}\bar{\Psi} \mathcal{D}\Psi \exp\left(i \int d^4x (\bar{\Psi} \hat{D}_\psi \Psi - i \int d^4y F(x) D_F(x-y) F(y))\right)$$

$$\mathcal{L}_{\text{eff}} = \bar{\Psi} \hat{D}_\psi \Psi - i \int d^4y F(x) D_F(x-y) F(y)$$

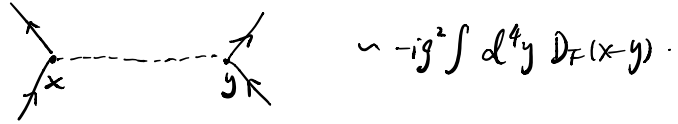
$$\begin{aligned} c) \quad p^2 \ll M^2, \quad D_F(x-y) &= \int \frac{d^4p}{(2\pi)^4} \frac{i}{p^2 - M^2 + i\epsilon} e^{-ip(x-y)} \\ &= \int \frac{d^4p}{(2\pi)^4} \underbrace{\frac{i}{p^2/M^2 - 1}}_{\substack{|p|^2 < M^2 \\ = i(-1 + \sigma(p^2/M^2))}} \cdot \frac{1}{M^2} e^{-ip(x-y)} \\ &= i(-1 + \sigma(p^2/M^2)) \end{aligned}$$

$$\begin{aligned} \rightarrow \mathcal{L}_{\text{eff}} &\cong \bar{\Psi} \hat{D}_\psi \Psi + \underbrace{i \int d^4y F(x) \frac{i}{M^2} \delta^{(4)}(x-y) F(y)}_{\substack{= -F^2(x) \frac{1}{M^2} \\ = -\frac{g^2}{M^2} (\bar{\Psi} \Psi)^2}} \\ &= -F^2(x) \frac{1}{M^2} \\ &= -\frac{g^2}{M^2} (\bar{\Psi} \Psi)^2 \end{aligned}$$

d) i)



ii)



iii)



e) We went from small scale physics to large scale by integrating field in generating function over.