

$$a) \quad \lambda = - (1 + 2\vec{L} \cdot \vec{S}) - \frac{i\gamma}{r} (\vec{\alpha} \cdot \vec{r}), \quad \gamma = -\omega$$

$$\lambda^2 = + (1 + 2\vec{L} \cdot \vec{S})^2 + \left[\frac{i\gamma}{r} (\vec{\alpha} \cdot \vec{r}) \right]^2 + (1 + 2\vec{L} \cdot \vec{S}) \frac{i\gamma}{r} (\vec{\alpha} \cdot \vec{r})$$

$$+ \frac{i\gamma}{r} (\vec{\alpha} \cdot \vec{r}) (1 + 2\vec{L} \cdot \vec{S})$$

$$= 1 + 4\vec{L} \cdot \vec{S} + 4(\vec{L} \cdot \vec{S})^2 - \frac{r^2}{r^2} (\vec{\alpha} \cdot \vec{r})^2 + \frac{2i\gamma}{r} \vec{\alpha} \cdot \vec{r}$$

$$+ 2i\gamma \left[(\vec{L} \cdot \vec{S}) \cdot \frac{\vec{\alpha} \cdot \vec{r}}{r} + \frac{\vec{\alpha} \cdot \vec{r}}{r} \cdot \vec{L} \cdot \vec{S} \right]$$

$$\left(\frac{\vec{\alpha} \cdot \vec{r}}{r} \right)^2 = \frac{1}{r^2} (\vec{\alpha} \cdot \vec{r})^2 = \frac{1}{r^2} \sum_i r^i \alpha_j r^j = \frac{1}{r^2} \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix} r^i \begin{pmatrix} 0 & \sigma_j \\ \sigma_j & 0 \end{pmatrix} r^j$$

$$= \frac{1}{r^2} \begin{pmatrix} \sigma_i \sigma_j & 0 \\ 0 & \sigma_i \sigma_j \end{pmatrix} r^i r^j = \frac{1}{r^2} r^i r^j (\delta_{ij} \mathbb{1} + i \epsilon_{ijk} \begin{pmatrix} \sigma_k & 0 \\ 0 & \sigma_k \end{pmatrix})$$

$$= \mathbb{1} + i \frac{1}{r^2} r^i r^j \epsilon_{ijk} 2\sigma^k = 1 + \frac{2i}{r^2} (\vec{r} \times \vec{r}) \cdot \vec{S} = \mathbb{1}$$

$$4(\vec{L} \cdot \vec{S})^2 = 4\vec{L} \cdot \vec{S} \cdot \vec{L} \cdot \vec{S} = 4\vec{L}^2 \cdot \vec{S}^2 = 4\vec{L}^2 \cdot \frac{1}{4} \mathbb{1} = \vec{L}^2 \quad \left. \begin{array}{l} L_i S^i L_j S^j \\ = L_i L_j S^i S^j \\ \neq S^2 \end{array} \right\}$$

$$= (\vec{S} \cdot \vec{\alpha})(\vec{S} \cdot \vec{r})$$

$$= 1 + 4\vec{L} \cdot \vec{S} + \vec{L}^2 - \gamma^2 + \text{something}$$

$$(\vec{L} \cdot \vec{S}) \cdot \frac{\vec{\alpha} \cdot \vec{r}}{r} = \frac{1}{2} \begin{pmatrix} \vec{\alpha} \cdot \vec{L} & 0 \\ 0 & \vec{r} \cdot \vec{L} \end{pmatrix} \begin{pmatrix} 0 & \vec{\alpha} \cdot \vec{r}/r \\ \vec{\alpha} \cdot \vec{r}/r & 0 \end{pmatrix}$$

$$= \frac{1}{2} \left(\begin{pmatrix} 0 & \vec{L} \cdot \vec{r} + i(\vec{L} \times \vec{r}) \cdot \vec{\alpha} \\ \vec{L} \cdot \vec{r} + i(\vec{L} \times \vec{r}) \cdot \vec{\alpha} & 0 \end{pmatrix} \right) = \frac{1}{2} \begin{pmatrix} 0 & \vec{L} \cdot \vec{r} \\ \vec{L} \cdot \vec{r} & 0 \end{pmatrix}$$

$$S_{th} = 2i\gamma \begin{pmatrix} 0 & (\vec{L} + \vec{\alpha}) \cdot \vec{r} \\ (\vec{L} + \vec{\alpha}) \cdot \vec{r} & 0 \end{pmatrix}$$

$$a) (\vec{L} \cdot \vec{S})^2 = \frac{1}{4} \left((\vec{L} \cdot \vec{r})^2 - (\vec{L} \cdot \vec{r})^2 \right) = \frac{1}{2} (\vec{L} - 2\vec{L} \cdot \vec{S})^2$$

$$\begin{aligned}
 (\vec{L} \cdot \vec{r})^2 &= (L_i \tau_i)(L_j \tau_j) = \frac{1}{2} (L_i L_j \tau_i \tau_j + \underbrace{L_i L_j \tau_i \tau_j}_{2S_{ij}}) \\
 &= \frac{1}{2} (L_i L_j [\tau_i \tau_j] - i \epsilon_{ijk} L_k \tau_j \tau_i) \\
 &= L^2 + \frac{1}{4} \epsilon_{ijk} L_k \tau_i \tau_j - \frac{i}{4} (\epsilon_{ijk} L_k \tau_j \tau_i) \\
 &= L^2 + \frac{i}{4} \epsilon_{ijk} L_k (\tau_i \tau_j - \tau_j \tau_i) \quad \xrightarrow{\text{Or } (\vec{L} \cdot \vec{r})^2 = L_i \tau_i L_j \tau_j} \\
 &= L^2 + \frac{i}{4} \epsilon_{ijk} L_k [\tau_i, \tau_j] \\
 &= \vec{L}^2 - 2\vec{L} \cdot \vec{S}
 \end{aligned}$$

$$(\vec{L} \cdot \vec{S})(\vec{L} \cdot \vec{r}) \cdot \frac{1}{r} + \frac{1}{r} (\vec{L} \cdot \vec{r})(\vec{L} \cdot \vec{S})$$

$$\begin{aligned}
 &= \frac{1}{2} \begin{pmatrix} 0 & \frac{\tau_i \tau_j L_i \tau_j}{r} + \tau_j \tau_i \frac{\tau_j}{r} L_i \\ \frac{\tau_i \tau_j L_i \tau_j}{r} + \tau_j \tau_i \frac{\tau_j}{r} L_i & 0 \end{pmatrix} \\
 &= \frac{1}{2} \begin{pmatrix} 0 & -2 \frac{\vec{r} \cdot \vec{r}}{r} \\ -2 \frac{\vec{r} \cdot \vec{r}}{r} & 0 \end{pmatrix}
 \end{aligned}$$

$$(1+2\vec{L} \cdot \vec{S}) \frac{i\tau}{r} (\vec{L} \cdot \vec{r}) + \frac{i\tau}{r} (\vec{L} \cdot \vec{r})(1+2\vec{L} \cdot \vec{S}) = 2 \frac{i\tau}{r} (\vec{L} \cdot \vec{r}) - 2 \frac{i\tau}{r} (\vec{L} \cdot \vec{r})$$

$$= 0$$

$$c) [-\frac{1}{r} \partial_r^2 r, 1]$$

$$\begin{aligned}
 &= (1+2\vec{L} \cdot \vec{S}) \frac{1}{r} \partial_r^2 r + \frac{i\tau}{r} (\vec{L} \cdot \vec{r}) \frac{1}{r} \partial_r^2 r - \frac{1}{r} \partial_r^2 r (1+2\vec{L} \cdot \vec{S}) - \frac{1}{r} \partial_r^2 r \frac{i\tau}{r} (\vec{L} \cdot \vec{r})
 \end{aligned}$$

$$= \frac{i\tau}{r} [(\vec{L} \cdot \vec{r}) \frac{1}{r} \partial_r^2 r - \frac{\partial^2}{\partial r^2} r (\frac{1}{r} (\vec{L} \cdot \vec{r}))]$$

$$= \frac{i\tau}{r} \left[(\vec{L} \cdot \vec{r}) \frac{1}{r} \partial_r^2 r - \frac{1}{r} (\vec{L} \cdot \vec{r}) \partial_r^2 r - \left(\frac{\partial^2}{\partial r^2} r - \frac{1}{r} \vec{L} \cdot \vec{r} \right) \right]$$

$$b) \quad \lambda(\lambda+1) = 1 + \vec{L}^2 + 2\vec{L} \cdot \vec{S} - \gamma^2 - (1 + 2\vec{L} \cdot \vec{S}) - \frac{i\gamma}{r} (\vec{\alpha} \cdot \vec{r})$$

$$= \vec{L}^2 - \gamma^2 - \frac{i\gamma}{r} \vec{\alpha} \cdot \vec{r}$$

$$\frac{1}{r^2} \vec{L}^2 - \left(E + \frac{r}{r}\right)^2 - \frac{i\gamma}{r^3} (\vec{\alpha} \cdot \vec{r})$$

$$= \frac{1}{r^2} \left\{ \vec{L}^2 - \gamma^2 - \frac{i\gamma}{r} \vec{\alpha} \cdot \vec{r} \right\} - E^2 - 2E \frac{r}{r}$$

$$= \frac{1}{r^2} \lambda(\lambda+1) - E^2 - 2E \frac{r}{r}$$

$\Rightarrow (2)$

$$c) \quad \lambda = -(1 + 2\vec{L} \cdot \vec{S}) - \frac{i\gamma}{r} (\vec{\alpha} \cdot \vec{r}) \propto r \rightarrow \partial_r^2 r = 0$$

$\Rightarrow \lambda$ and ∂_r^2 commute

$\Rightarrow \lambda$ and the whole terms in brackets commute

$$d) \quad \lambda^2 \psi = \lambda^2 \psi$$

$$\Rightarrow (1 - r^2 + \vec{L}^2 + 2\vec{L} \cdot \vec{s}) \psi = \lambda^2 \psi$$

$$(1 - r^2 + \cancel{\vec{L}^2} + \vec{s}^2 - \cancel{L^2} - \vec{s}^2) \psi = \lambda^2 \psi$$

$$\Rightarrow \lambda^2 = 1 - r^2 + s(s+1) - s(s+1) = -r^2 + j(j+1) + \frac{1}{4}$$
$$= -r^2 + (j + \frac{1}{2})^2 \quad , \quad \lambda = \pm |\lambda|$$

$$\lambda(\lambda+1)\psi = \lambda^2 \psi + \lambda \psi = \lambda^2 \psi \pm i\lambda \psi$$

$$= |\lambda| (|\lambda| \pm 1) \psi$$

$$= \begin{cases} |\lambda| (|\lambda| + 1) \psi & , \lambda > 0 \\ |\lambda| (|\lambda| - 1) \psi & , \lambda < 0 \end{cases}$$

$$\Rightarrow \text{define } \tilde{\ell} = \begin{cases} |\lambda| & , \lambda > 0 \\ |\lambda| - 1 & , \lambda < 0 \end{cases}$$

$$\Rightarrow \lambda(\lambda+1)\psi = \tilde{\ell}(\tilde{\ell}+1)\psi = \begin{cases} |\lambda| (|\lambda| + 1) \psi & , \lambda > 0 \\ (|\lambda| - 1) |\lambda| \psi & , \lambda < 0 \end{cases}$$

$$\Rightarrow (2) \rightarrow \left(-\frac{1}{r^2} \partial_r r + \frac{\tilde{\ell}(\tilde{\ell}+1)}{r^2} - \frac{2\epsilon r}{r} - \epsilon^2 + m^2 \right) \psi = 0$$

e) non-relativistic:

$$\left(-\frac{1}{r} \frac{\partial^2}{\partial r^2} r + \frac{l(l+1)}{r^2} - \frac{2mr}{r} - 2mE \right) \psi = 0$$

$$\Rightarrow \left\{ \frac{1}{2m} \left[-\frac{1}{r} \frac{\partial}{\partial r} r + \frac{l(l+1)}{r^2} \right] - \frac{r}{r} \right\} \psi = E \psi$$

with energy: $E = -E_0 \left(\frac{z}{n} \right)^2 = -\frac{1}{2} m c^2 \omega^2 \left(\frac{z}{n} \right)^2 = -\frac{1}{2} m \left(\frac{r}{n} \right)^2$

with $n = k_0 + l + 1$

relativistic:

$$\left(-\frac{1}{r} \frac{\partial^2}{\partial r^2} r + \frac{\tilde{l}(\tilde{l}+1)}{r^2} - \frac{2Er}{r} - E^2 + m^2 \right) \psi = 0$$

$$\left\{ \frac{1}{2E} \left[\left(-\frac{1}{r} \frac{\partial}{\partial r} r + \frac{\tilde{l}(\tilde{l}+1)}{r^2} \right) - \frac{r}{r} \right] - \frac{r}{r} \right\} \psi = \frac{E^2 - m^2}{2E} \psi$$

$$m \rightarrow E, \quad l \rightarrow \tilde{l}, \quad n \rightarrow \tilde{n} = k_0 + \tilde{l} + 1 = n + \tilde{l} - l$$

$$\begin{cases} n + |\lambda| - l, & \lambda > 0 \\ n + |\lambda| - 1 - l, & \lambda < 0 \end{cases}$$

+ somehow get neglected

$$\frac{E^2 - m^2}{2E} = -\frac{1}{2} E \left(\frac{r}{n} \right)^2,$$

↑
Don't have to
 $n = N + 1 + j \pm \frac{1}{2}$

$$\Rightarrow E^2 \left(1 + \frac{r^2}{n^2} \right) = m^2$$

$$\Rightarrow E_n = m \sqrt{1 + r^2 \underbrace{\left(n + j - \frac{r^2}{n^2} + \left(j + \frac{1}{2} \right)^2 - \left(j + \frac{1}{2} \right) \right)}_{|\lambda|^{-2}}}$$

e)

$$\left(-\frac{1}{r} \partial_r^2 r + \frac{\ell(\ell+1)}{r^2} - 2 \underbrace{\frac{mr}{r}}_a - 2mE \right) \psi = 0$$

$$\left(-\frac{1}{r} \partial_r^2 r + \frac{\ell(\ell+1)}{r^2} - 2 \frac{mr}{r} - E^2 + m^2 \right) \psi = 0$$

$$f) \gamma = z \propto << 1$$

$$\frac{\partial |\lambda|}{\partial \gamma} = \frac{\partial}{\partial \gamma} \sqrt{-\gamma^2 + (j + \frac{1}{2})^2} = \frac{1}{2} \sqrt{-\gamma^2 + (j + \frac{1}{2})^2}^{-1} \cdot -2\gamma = -\gamma / |\lambda|$$

$$\frac{\partial n}{\partial \gamma} = \frac{\partial}{\partial \gamma} (n - j - \frac{1}{2} + |\lambda|) = -\frac{\gamma}{|\lambda|}$$

$$\frac{\partial E_{nj}}{\partial \gamma} = \frac{\partial}{\partial \gamma} \frac{m}{\sqrt{1 + \gamma^2 n^{-2}}} = -\frac{m}{2} (1 + \gamma^2 n^{-2})^{-3/2} \cdot \left\{ 2\gamma n^{-2} - 2\gamma^2 n^{-3} \cdot \frac{-\gamma}{|\lambda|} \right\}$$

$$= -\frac{m}{2} (1 + \gamma^2 n^{-2})^{-3/2} 2\gamma n^{-2} (1 + 2\gamma^2 / n|\lambda|)$$

$$\frac{\partial E_{nj}}{\partial \gamma} \Big|_{\gamma=0} = 0$$

$$\frac{\partial^2 E_{nj}}{\partial \gamma^2} = +\frac{3m}{4} (1 + \gamma^2 n^{-2})^{-5/2} \left[2\gamma n^{-2} (1 + 2\gamma^2 / n|\lambda|) \right]^2$$

$$- \frac{m}{2} (1 + \gamma^2 n^{-2})^{-3/2} \left\{ 2n^{-2} (1 + 2\gamma^2 / n|\lambda|) - 4\gamma n^{-3} \cdot \frac{-\gamma}{|\lambda|} (1 + 2\gamma^2 / n|\lambda|) \right\}$$

$$+ 2\gamma n^{-2} \left[4\gamma / n|\lambda| - 2\gamma^2 / n^2 |\lambda| \cdot \frac{-\gamma}{|\lambda|} - 2\gamma^2 / n|\lambda|^2 \cdot \frac{-\gamma}{|\lambda|} \right] \}$$

$$\frac{\partial^2}{\partial \gamma^2} E_{nj} \Big|_{\gamma=0} = -\frac{m}{2} \cdot 2n^{-2} = -\frac{m}{n^2} = -\frac{m}{n^2}$$

$$\Rightarrow E_{nj} = m \left(1 - \frac{\gamma^2}{2n^2} + O(\gamma^4) \right)$$

H9.2

$$a) [\hat{a}_i, \hat{a}_j] = 0, [\hat{a}_i^+, \hat{a}_j^+] = 0$$

$$[\hat{a}_i \hat{a}_j, \hat{a}_k] = \hat{a}_i [\hat{a}_j, \hat{a}_k] + [\hat{a}_i, \hat{a}_k] \hat{a}_j = 0$$

$$[\hat{a}_i^+ \hat{a}_j, \hat{a}_k] = 0$$

$$[\hat{a}_i \hat{a}_j, \hat{a}_k^+] = \hat{a}_i [\hat{a}_j, \hat{a}_k^+] + [\hat{a}_i, \hat{a}_k^+] \hat{a}_j \\ = \hat{a}_i \delta_{jk} + \delta_{ik} \hat{a}_j$$

$$[\hat{a}_i \hat{a}_j, \hat{a}_k] = \hat{a}_i [\hat{a}_j, \hat{a}_k] + [\hat{a}_i, \hat{a}_k] \hat{a}_j \\ = -\hat{a}_i \delta_{jk} + \delta_{ik} \hat{a}_j$$

$$b) \hat{H} = \sum_{i,j} t_{ij} \hat{a}_i^+ \hat{a}_j + \frac{1}{2} \sum_{i,j,k,l} V_{ijkl} \hat{a}_i^+ \hat{a}_j^+ \hat{a}_k \hat{a}_l$$

$$\left[\sum_{i,j} t_{ij} \hat{a}_i^+ \hat{a}_j, \sum_i \hat{a}_i \hat{a}_i \right]$$

$$= \sum_{i,j,k} t_{ij} [\hat{a}_i^+ \hat{a}_j, \hat{a}_k^+ \hat{a}_k] = \sum t_{ij} ([\hat{a}_i^+ \hat{a}_j, \hat{a}_k^+] \hat{a}_k + \hat{a}_k^+ [\hat{a}_i^+ \hat{a}_j, \hat{a}_k])$$

$$= \sum_{i,j,k} t_{ij} \left\{ (\hat{a}_i^+ [\hat{a}_j, \hat{a}_k^+] + [\hat{a}_i^+, \hat{a}_k^+] \hat{a}_j) \hat{a}_k + \hat{a}_k^+ (\hat{a}_i^+ [\hat{a}_j, \hat{a}_k] + [\hat{a}_i^+, \hat{a}_k] \hat{a}_j) \right\}$$

$$= \sum_{i,j,k} t_{ij} (\hat{a}_i^+ \delta_{jk} \hat{a}_k + \hat{a}_k^+ \delta_{ik} \hat{a}_j)$$

$$= \sum_{i,j} t_{ij} (\hat{a}_i^+ \hat{a}_j + \hat{a}_i^+ \hat{a}_j) = 0$$

$$\begin{aligned} & [AB, C] \\ &= A[C, B] + [A, C]B \\ & [A, BC] \\ &= [A, B]C + BCA, C \end{aligned}$$

$$[\hat{N}, \sum_{i,j,k} a_i^{\dagger} a_j^{\dagger} a_k]$$

$$= \sum V [a_k^{\dagger} a_k, a_i^{\dagger} a_j^{\dagger} a_k]$$

$$= \sum V \{ a_k^{\dagger} [a_k, a_i^{\dagger} a_j^{\dagger} a_k] + [a_k^{\dagger}, a_i^{\dagger} a_j^{\dagger} a_k] a_k \}$$

$$= \sum V \{ a_k^{\dagger} [a_k, a_i^{\dagger} a_j^{\dagger}] a_k + a_i^{\dagger} a_j^{\dagger} [a_k^{\dagger}, a_k] a_k \}$$

H9.3

a) $\hat{T} = \sum_{i,j} \hat{t} = \sum_{i,j} t_{ij} \sum_{\alpha} |i\rangle_{\alpha} \langle j|_{\alpha}$ acting on an arbitrary state

$$\frac{1}{\sqrt{N}} C |n_1, n_2, \dots \rangle = \sum_{i,j} t_{ij} \sum_{\alpha} |i\rangle_{\alpha} \langle j|_{\alpha} |i_1\rangle_1 |i_2\rangle_2 \dots |i_N\rangle_N$$

for every particle
normalisation

$$= C \sum_{i,j} t_{ij} \sum_{\alpha} |i_1\rangle_1 |i_2\rangle_2 \dots |i\rangle_{\alpha} |i_N\rangle_N \delta_{ji\alpha}$$

$$\langle j|_{\alpha} |i\rangle_{\alpha} = \delta_{ji\alpha}$$

in the process a particle with state $j = i_2$
is replaced by state i , i.e. $\hat{a}_i^+ \hat{a}_j^-$

comes from $\sum_{\alpha} \delta_{ji\alpha}$, acts on n_j -particles

↑

$$= \sum_{i,j} t_{ij} n_j C |n_1, \dots, n_{i+1}, \dots, n_j - 1, \dots \rangle$$

$(C = \frac{1}{\sqrt{n_1! n_2! \dots}},$ since the occupation number of i, j states have changed
the C has to change to keep the wave function normalised)

$$= \sum_{i,j} t_{ij} n_j \sqrt{n_{i+1}} \frac{1}{\sqrt{n_j}} C |n_1, \dots, n_{i+1}, \dots, n_j - 1, \dots \rangle$$

$$= \sum_{i,j} t_{ij} \hat{a}_i^+ \hat{a}_j^- |i_1, i_2, \dots, i_N\rangle$$

→ the original ket vector(s)

$$\Rightarrow \hat{F} = \sum_{i,j} t_{ij} \hat{a}_i^+ \hat{a}_j^-$$

$$(2) \quad \hat{F} = \frac{1}{2} \sum_{\alpha \neq \beta} \hat{f}^{(2)}(\chi_{\alpha}, \chi_{\beta}) = \frac{1}{2} \sum_{\alpha, \beta, \gamma, \delta} \hat{f}_{\alpha \beta \gamma \delta}^{(2)} |\alpha\rangle |\beta\rangle \langle \gamma | \langle \delta|$$

$$= \frac{1}{2} \sum_{\alpha, \beta, \gamma, \delta} \hat{f}_{\alpha \beta \gamma \delta}^{(2)} \hat{a}_{\alpha}^+ \hat{a}_{\beta}^+ \hat{a}_{\gamma}^- \hat{a}_{\delta}^-$$

$$F = \frac{1}{2} \sum_{\alpha \neq \beta} f^{(2)}(\chi_{\alpha}, \chi_{\beta}) = \frac{1}{2} \sum_{\alpha \neq \beta} \sum_{i, j, k, l} |i\rangle_{\alpha} \langle j|_{\beta} \langle i | f^{(2)} | k l \rangle \langle l |_{\beta} \langle k |_{\alpha}$$

$$= \frac{1}{2} \sum_{i, k, j, l} f^{(2)}_{ijkl} \hat{a}_i^+ \hat{a}_j^+ \hat{a}_k^- \hat{a}_l^- \sum_{\alpha \neq \beta} |0\rangle_{\alpha} \langle 0|_{\beta} \langle 0 |_{\beta} \langle 0 |_{\alpha}$$