

H8.1

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$$a) \quad i\hbar \gamma^\mu \partial_\mu \bar{\Psi} = mc \bar{\Psi} \quad \bar{\Psi}(x^\mu) = u(\epsilon, \vec{p}) e^{i/\hbar(\vec{p} \cdot \vec{x} - \epsilon t)}$$

$$\Rightarrow i\hbar (\beta \partial_0 + \alpha^i \partial_i) \bar{\Psi} = mc \bar{\Psi}$$

$$\Rightarrow i\hbar \partial_t \bar{\Psi} = (c \vec{\alpha} \cdot \vec{p} + \beta mc^2) \bar{\Psi}$$

Plug in $\bar{\Psi}(x^\mu)$

$$i\hbar \partial_t u(\epsilon, \vec{p}) e^{i/\hbar(\vec{p} \cdot \vec{x} - \epsilon t)} = \begin{bmatrix} mc^2 & c \vec{\sigma} \cdot \vec{p} \\ c \vec{\sigma} \cdot \vec{p} & -mc^2 \end{bmatrix} u(\epsilon, \vec{p}) e^{i/\hbar(\vec{p} \cdot \vec{x} - \epsilon t)}$$

$$i\hbar \cdot \frac{i}{\hbar} \cdot (-\epsilon) u(\epsilon, \vec{p}) e^{i/\hbar(\vec{p} \cdot \vec{x} - \epsilon t)} = \begin{bmatrix} mc^2 & c \vec{\sigma} \cdot \vec{p} \\ c \vec{\sigma} \cdot \vec{p} & -mc^2 \end{bmatrix} u(\epsilon, \vec{p}) e^{i/\hbar(\vec{p} \cdot \vec{x} - \epsilon t)}$$

$$\Rightarrow \epsilon u(\epsilon, \vec{p}) = \begin{bmatrix} mc^2 & c \vec{\sigma} \cdot \vec{p} \\ c \vec{\sigma} \cdot \vec{p} & -mc^2 \end{bmatrix} u(\epsilon, \vec{p})$$

← not sure why we can eliminate $e^{i/\hbar(\vec{p} \cdot \vec{x} - \epsilon t)}$ here ...

$$\Rightarrow \begin{bmatrix} mc^2 - \epsilon & c \vec{\sigma} \cdot \vec{p} \\ c \vec{\sigma} \cdot \vec{p} & -mc^2 - \epsilon \end{bmatrix} u(\epsilon, \vec{p}) = 0$$

$$\text{determinant: } -(m^2 c^4 - \epsilon^2) - c^2 (\vec{\sigma} \cdot \vec{p})^2 = 0$$

$$\epsilon^2 - m^2 c^4 - c^2 \vec{p}^2 = 0$$

$$\left((\vec{\sigma} \cdot \vec{p})(\vec{\sigma} \cdot \vec{p}) = 1(\vec{p} \cdot \vec{p}) + i \vec{\sigma}(\vec{p} \times \vec{p}) = \vec{p} \cdot \vec{p} \right)$$

$$\Rightarrow \epsilon = \pm \sqrt{m^2 c^4 + c^2 \vec{p}^2} = \pm E$$

Return to the equation:

$$\begin{bmatrix} mc^2 - E & c \vec{\sigma} \cdot \vec{p} \\ c \vec{\sigma} \cdot \vec{p} & -mc^2 - E \end{bmatrix} u(E, \vec{p}) = 0$$

$$E = +E, \text{ particle: } u(E = E) = \begin{pmatrix} \frac{c \vec{\sigma} \cdot \vec{p}}{E + mc^2} \varphi \\ \varphi \end{pmatrix}$$

$$(mc^2 - E) \varphi + c \vec{\sigma} \cdot \vec{p} \left(\frac{c \vec{\sigma} \cdot \vec{p}}{E + mc^2} \right) \varphi$$

$$= \left(mc^2 - E + \frac{c^2 (\vec{p} \cdot \vec{p})}{E + mc^2} \right) \varphi = \frac{(E + mc^2)(-E + mc^2) + c^2 p^2}{E + mc^2} \varphi = 0 \Rightarrow \varphi$$

$$b) \quad \vec{p} = (p, 0, 0)^T$$

$$(\vec{\sigma} \cdot \vec{p}) \psi = \pm |\vec{p}| \psi \quad \Rightarrow \quad \begin{pmatrix} 0 & p \\ p & 0 \end{pmatrix} \psi = \pm p \psi$$

$$\Rightarrow \begin{pmatrix} 0 & p \\ p & 0 \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} = \pm p \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$$

$$\begin{cases} p\psi_2 = \pm p\psi_1 \\ p\psi_1 = \pm p\psi_2 \end{cases}$$

$$\Rightarrow \psi_1 = \pm \psi_2 \quad \psi = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} \quad \text{or} \quad \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$c) \quad N^2 = (E + mc^2) / 2E$$

$$u^\dagger u = \left[N \left(\frac{c \vec{\sigma} \cdot \vec{p}}{E + mc^2} \psi \right) \right]^\dagger \left[N \left(\frac{c \vec{\sigma} \cdot \vec{p}}{E + mc^2} \psi \right) \right]$$

$$= \frac{E + mc^2}{2E} \left[\psi^\dagger \psi + \left(\frac{c \vec{\sigma} \cdot \vec{p}}{E + mc^2} \psi \right)^\dagger \left(\frac{c \vec{\sigma} \cdot \vec{p}}{E + mc^2} \psi \right) \right]$$

$$\left[(\vec{\sigma} \cdot \vec{p})^\dagger = \begin{pmatrix} p_z & p_x - ip_y \\ p_x + ip_y & p_z \end{pmatrix}^\dagger = (\vec{\sigma} \cdot \vec{p}) \right]$$

$$= \frac{E + mc^2}{2E} \psi^\dagger \psi \left(1 + \frac{c^2 |\vec{p}|^2}{(E + mc^2)^2} \right)$$

$$= \psi^\dagger \psi \frac{E + mc^2}{2E} \cdot \frac{(E + mc^2)^2 + p^2 c^2}{(E + mc^2)^2}$$

$$= \psi^\dagger \psi \frac{1}{2E} \frac{E^2 + 2Emc^2 + E^2}{E + mc^2}$$

$$= \frac{1}{2E} \frac{2E(E + mc^2)}{E + mc^2} = 1$$

$$= \psi^\dagger \psi$$

$$a) \quad \rho = \bar{\Psi}^+ \bar{\Psi} \quad \vec{J} = \bar{\Psi}^+ (c\vec{\alpha}) \bar{\Psi}$$

$$\vec{J} = U(E)^+ (c\vec{\alpha}) U(E)$$

$$\vec{J}^i = cN^2 \begin{pmatrix} \psi \\ \frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi \end{pmatrix}^+ \begin{pmatrix} 0 & \sigma^i \\ \sigma^i & 0 \end{pmatrix} \begin{pmatrix} \psi \\ \frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi \end{pmatrix}$$

$$= cN^2 \begin{pmatrix} \frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi^+ \cdot \sigma^i, & \psi^+ \sigma^i \end{pmatrix} \begin{pmatrix} \psi \\ \frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi \end{pmatrix}$$

$$= cN^2 \left(\frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi^+ \cdot \sigma^i \psi + \psi^+ \sigma^i \frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi \right)$$

$$= cN^2 \left(\psi^+ \sigma^i \frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi + \psi^+ \sigma^i \frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi \right)$$

$$= N^2 \cdot 2c \psi^+ \sigma^i \frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi$$

$$= \frac{E+mc^2}{2E} \cdot 2c \psi^+ \sigma^i \frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi$$

$$\Rightarrow \vec{J} = \frac{c^2}{E} \psi^+ \vec{\sigma} (\vec{\sigma} \cdot \vec{p}) \psi = \frac{c^2}{E} \psi^+ \vec{p} \psi = \frac{c^2 \vec{p}}{E} \psi^+ \psi$$

using $(\vec{\sigma} \cdot \vec{p}) \psi = \pm p \psi$

$$\left[\begin{aligned} \rho = \bar{\Psi}^+ \bar{\Psi} &= N^2 \left(\psi^+, \frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi^+ \right) \cdot \begin{pmatrix} \psi \\ \frac{c\vec{\sigma} \cdot \vec{p}}{E+mc^2} \psi \end{pmatrix} \\ &= \frac{E+mc^2}{2E} \psi^+ \psi \left(1 + \frac{c^2 p^2}{E+mc^2} \right) = \psi^+ \psi \quad \text{just like last part} \end{aligned} \right]$$

$$\Rightarrow \vec{J} = \frac{c^2 \vec{p}}{E} \psi^+ \psi = \frac{c^2 \vec{p}}{E} \rho = \pm \frac{c^2 \vec{p}}{E} \rho = \pm \vec{v} \rho$$

$$d) \quad J_i = \bar{\psi} (c \vec{\alpha}) \psi = \left(\psi^\dagger \quad \psi^\dagger \frac{c \vec{\sigma} \cdot \vec{p}}{E + mc^2} \right) \begin{pmatrix} 0 & c \sigma_i \\ c \sigma_i & 0 \end{pmatrix} \begin{pmatrix} \psi \\ \frac{c \vec{\sigma} \cdot \vec{p}}{E + mc^2} \psi \end{pmatrix}$$

$$= N^2 \left(\psi^\dagger \quad \psi^\dagger \frac{c \vec{\sigma} \cdot \vec{p}}{E + mc^2} \right) \begin{pmatrix} \frac{c^2 \sigma_i (\vec{\sigma} \cdot \vec{p})}{E + mc^2} \psi \\ c \sigma_i \psi \end{pmatrix}$$

$$= N^2 \left(\psi^\dagger \frac{c^2 \sigma_i (\vec{\sigma} \cdot \vec{p})}{E + mc^2} \psi + \psi^\dagger \frac{c^2 (\vec{\sigma} \cdot \vec{p}) \sigma_i}{E + mc^2} \psi \right)$$

$$\left[\{ \sigma_i, \sigma_j \} = 2 \delta_{ij} \quad \Rightarrow \quad \sigma_i \sigma_j = 2 \delta_{ij} - \sigma_j \sigma_i \right]$$

$$= N^2 \left(\frac{2 p_i c^2}{E + mc^2} \psi^\dagger \psi \right) = \frac{c^2 \vec{p}}{E} \psi^\dagger \psi$$

H8.2

$$a) \quad z < -\frac{a}{2}: \quad \bar{\Psi}_A(z) = A_+ e^{ipz/\hbar} \begin{pmatrix} 1 \\ 0 \\ \frac{pc}{E+mc^2} \\ 0 \end{pmatrix} + A_- e^{-ipz/\hbar} \begin{pmatrix} 1 \\ 0 \\ \frac{-pc}{E+mc^2} \\ 0 \end{pmatrix}$$

$$z \in [-\frac{a}{2}, \frac{a}{2}]: \quad \bar{\Psi}_B(z) = B_+ e^{ip'z/\hbar} \begin{pmatrix} 1 \\ 0 \\ \frac{p'c}{E-V_0+mc^2} \\ 0 \end{pmatrix} + B_- e^{-ip'z/\hbar} \begin{pmatrix} 1 \\ 0 \\ \frac{-p'c}{E-V_0+mc^2} \\ 0 \end{pmatrix}$$

$$z > \frac{a}{2}: \quad \bar{\Psi}_C(z) = C_+ e^{ipz/\hbar} \begin{pmatrix} 1 \\ 0 \\ \frac{pc}{E+mc^2} \\ 0 \end{pmatrix} + C_- e^{-ipz/\hbar} \begin{pmatrix} 1 \\ 0 \\ \frac{-pc}{E+mc^2} \\ 0 \end{pmatrix}$$

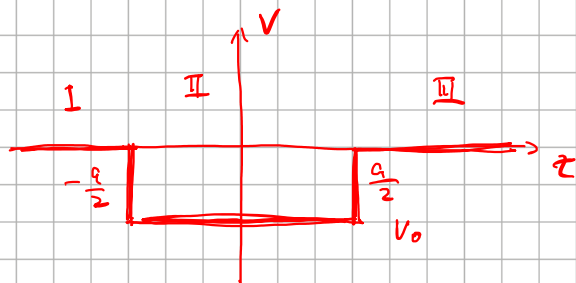
taking: $c = \hbar = 1$

Continuity:

$$\begin{cases} \bar{\Psi}_A(-\frac{a}{2}) = \bar{\Psi}_B(-\frac{a}{2}) \\ \bar{\Psi}_B(\frac{a}{2}) = \bar{\Psi}_C(\frac{a}{2}) \end{cases}$$

$$\begin{aligned} \rightarrow \quad & A_+ e^{-ipa/2} \begin{pmatrix} 1 \\ 0 \\ \frac{pc}{E+mc^2} \\ 0 \end{pmatrix} + A_- e^{+ipa/2} \begin{pmatrix} 1 \\ 0 \\ \frac{-pc}{E+mc^2} \\ 0 \end{pmatrix} \\ & = B_+ e^{-ip'a/2} \begin{pmatrix} 1 \\ 0 \\ \frac{p'c}{E-V_0+mc^2} \\ 0 \end{pmatrix} + B_- e^{+ip'a/2} \begin{pmatrix} 1 \\ 0 \\ \frac{-p'c}{E-V_0+mc^2} \\ 0 \end{pmatrix} \end{aligned}$$

a)



Greiner C.9

for I, III

$$(\vec{\alpha} \cdot \vec{p} + \beta m) \psi = E \psi$$

for II

$$(\vec{\alpha} \cdot \vec{p} + \beta m) \psi = (E - V_0) \psi$$

Take Spin up:

$$\psi_I(z) = A e^{ip_1 z} \begin{pmatrix} 1 \\ p_1 \\ E+m \\ 0 \end{pmatrix} + A' e^{-ip_1 z} \begin{pmatrix} 1 \\ -p_1 \\ E+m \\ 0 \end{pmatrix}$$

$$\psi_{II}(z) = B e^{ip_2 z} \begin{pmatrix} 1 \\ p_2 \\ E-V_0+m \\ 0 \end{pmatrix} + B' e^{-ip_2 z} \begin{pmatrix} 1 \\ -p_2 \\ E-V_0+m \\ 0 \end{pmatrix} \quad \leftarrow \begin{matrix} \text{omitting} \\ 0 \end{matrix}$$

$$\psi_{II}(z) = C e^{ip_1 z} \begin{pmatrix} 1 \\ p_1 \\ E+m \\ 0 \end{pmatrix} + C' e^{-ip_1 z} \begin{pmatrix} 1 \\ -p_1 \\ E+m \\ 0 \end{pmatrix}$$

"Current" comes from the other side \rightarrow anti-particle!

continuity at borders:

$$\psi_I(-\frac{a}{2}) = \psi_{II}(-\frac{a}{2})$$

$$\psi_{II}(\frac{a}{2}) = \psi_{III}(\frac{a}{2})$$

$$\Rightarrow z = -\frac{a}{2} :$$

$$\textcircled{1} A e^{-ip_1 a/2} + A' e^{ip_1 a/2} = B e^{-ip_2 a/2} + B' e^{ip_2 a/2}$$

$$\textcircled{2} (A e^{-ip_1 a/2} - A' e^{ip_1 a/2}) \frac{p_1}{E+m} = (B e^{-ip_2 a/2} - B' e^{ip_2 a/2}) \frac{p_2}{E-V_0+m}$$

$$z = \frac{a}{2} :$$

$$\textcircled{3} \quad B e^{i p_2 a/2} + B' e^{-i p_2 a/2} = C e^{i p_1 a/2} + C' e^{-i p_1 a/2}$$

$$\textcircled{4} \quad (B e^{i p_2 a/2} - B' e^{-i p_2 a/2}) \frac{p_2}{E - V_0 + m} = (C e^{i p_1 a/2} - C' e^{-i p_1 a/2}) \frac{p_1}{E + m}$$

$$\text{Write : } \gamma = \frac{\frac{p_1}{E + m}}{\frac{p_2}{E - V_0 + m}} = \frac{\sqrt{E^2 - m^2}}{E + m} \cdot \frac{E - V_0 + m}{\sqrt{(E - V_0)^2 - m^2}}$$

$$= \sqrt{\frac{E - m}{E + m} \cdot \frac{(E - V_0) - m}{(E - V_0) + m}}$$

in Matrix form:

$$z = \frac{a}{2} :$$

$$\textcircled{1} + \textcircled{2} : 2 A e^{-i p_2 a/2} = B e^{i p_1 a/2} \left(1 + \frac{1}{\gamma}\right) + B' e^{-i p_1 a/2} \left(1 - \frac{1}{\gamma}\right)$$

$$A e^{-i p_2 a/2} = \frac{1}{2} B e^{i p_1 a/2} \left(\frac{\gamma+1}{\gamma}\right) + \frac{1}{2} B' e^{-i p_1 a/2} \left(\frac{\gamma-1}{\gamma}\right) \quad \text{then multiply } e^{+i p_1 a/2}$$

→ do the same to A' : $\textcircled{1} - \textcircled{2}$

$$z = \frac{a}{2} :$$

$$\textcircled{3} + \textcircled{4} : 2 B e^{i p_2 a/2} = C e^{i p_1 a/2} (1 + \gamma) + C' e^{-i p_1 a/2} (1 - \gamma) \quad \text{multiply by } e^{i p_2 a/2}$$

Same to get B'

$$\textcircled{3} - \textcircled{4}$$

$$\Rightarrow \begin{pmatrix} A \\ A' \end{pmatrix} = \frac{1}{2} \begin{pmatrix} \frac{\gamma+1}{\gamma} e^{i(p_1 - p_2)a/2} & \frac{\gamma-1}{\gamma} e^{i(p_2 + p_1)a/2} \\ \frac{\gamma-1}{\gamma} e^{-i(p_2 + p_1)a/2} & \frac{\gamma+1}{\gamma} e^{i(p_2 - p_1)a/2} \end{pmatrix} \begin{pmatrix} B \\ B' \end{pmatrix}$$

$$\begin{pmatrix} B \\ B' \end{pmatrix} = \frac{1}{2} \begin{pmatrix} (1+\gamma) e^{i(p_2 - p_1)a/2} & (1-\gamma) e^{-i(p_1 + p_2)a/2} \\ (1-\gamma) e^{i(p_1 + p_2)a/2} & (1+\gamma) e^{i(p_2 - p_1)a/2} \end{pmatrix} \begin{pmatrix} C \\ C' \end{pmatrix}$$

$$\Rightarrow \begin{pmatrix} A \\ A' \end{pmatrix} = \frac{1}{4\gamma} \begin{pmatrix} (1+r)^2 e^{i(p_1-p_2)a} - (1-r)^2 e^{i(p_1+p_2)a}, (1-r^2)(e^{-ip_2a} - e^{ip_2a}) \\ -(1-r)^2(e^{ip_2a} - e^{-ip_2a}) \end{pmatrix}$$

$\underbrace{(1+r)^2 e^{-i(p_1-p_2)a} - (1-r)^2 e^{-i(p_1+p_2)a}}_{e^{-i(p_1+p_2)a}}$
 $\begin{pmatrix} C \\ C' \end{pmatrix}$

$$\int dz \psi^\dagger \psi = 1$$

$$b) \quad E^2 = m^2 + p_1^2 \quad (E - V_0)^2 = m^2 + p_2^2$$

p_i could be imaginary / real

$$\begin{cases} E^2 > m^2 & \longrightarrow \begin{cases} E > m \\ E < -m \end{cases} \\ E^2 < m^2 \text{ \& } (E - V_0)^2 > m^2 & \longrightarrow -m < E < m \text{ \& } \begin{cases} E - V_0 > m \rightarrow E > m + V_0 \\ E - V_0 < -m \rightarrow E < -m + V_0 \end{cases} \\ E^2, (E - V_0)^2 < m^2 & \longrightarrow \text{in } (-\infty, \infty) \quad \psi = 0 \end{cases}$$

oscillatory in II, damps in I, III
 ↗