Quantum Field Theory

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3 Quantization of the Dirac field

3.1 Dirac equation

$$\left(i\gamma^{\mu}\partial_{\mu}-m\right)\phi(x)=0\tag{3.1.1}$$

Standard representation (Dirac's)

$$\gamma_0 = \begin{pmatrix} \mathbb{1}_2 & 0 \\ 0 & -\mathbb{1}_2 \end{pmatrix} \quad \boldsymbol{\gamma} = \begin{pmatrix} 0 & \boldsymbol{\sigma} \\ -\boldsymbol{\sigma} & 0 \end{pmatrix}$$
 (3.1.2)

Lorentz transformation

$$\Lambda = \exp\left(\frac{1}{2}\omega_{\mu\nu}M^{\mu\nu}\right) \tag{3.1.3}$$

with ω set of parameters and M the generator of Lie algebra.

Spinor representation

$$S^{\rho\sigma} = \frac{1}{4} \left[\gamma^{\rho}, \gamma^{\sigma} \right] = \frac{1}{2i} \sigma^{\rho\sigma} \tag{3.1.4}$$

(3.1.5)

Spinor transformation

$$S(\Lambda) = \exp\left(\frac{1}{2}\omega_{\mu\nu}S^{\mu\nu}\right) \tag{3.1.6}$$

$$\psi_a'(x) = S_{ab}(\Lambda)\psi_b(\Lambda^{-1}x) \tag{3.1.7}$$

adjoint spinor

$$\bar{\psi} = \psi^{\dagger} \gamma^0 \tag{3.1.8}$$

Fifth gamma matrix

$$\gamma^5 := i\gamma^0 \gamma^1 \gamma^2 \gamma^3 \tag{3.1.9}$$

$$\left\{\gamma^{\mu}, \gamma^{5}\right\} = 0 \tag{3.1.10}$$

$$(\gamma^5)^2 = \mathbb{1}_4 \tag{3.1.11}$$

Plane wavesolutions

$$\psi(x) = \begin{cases} u(p)e^{-ipx} & \text{positive frequency} \\ v(p)e^{ipx} & \text{negative frequency} \end{cases}$$
(3.1.12)

$$u_s(p) = \sqrt{E_p + m} \left(\frac{\chi_s}{\frac{u \cdot p}{E_p + m} \chi_s} \right) e^{-ipx} v_s(p) = \sqrt{E_p + m} \left(\frac{\frac{u \cdot p}{E_p + m} \tilde{\chi}_s}{\tilde{\chi}_s} \right) e^{ipx}$$
(3.1.13)

with

$$\chi_{\frac{1}{2}} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad x_{-\frac{1}{2}} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$
$$s = \pm \frac{1}{2} \quad \tilde{\chi}_s = \chi_{-s}$$

Orthogonality of spinor

$$\bar{u}_s(p)u_{s'}(p) = -\bar{v}_s(p)v_{s'}(p) = 2m\delta_{ss'}$$
(3.1.14)

$$\bar{u}_s(p)v_{s'}(p) = 0 (3.1.15)$$

Spin sums

$$\sum_{s} u_{s}(p)\bar{u}_{s}(p) = p + m \tag{3.1.16}$$

$$\sum_{s} v_s(p)\bar{v}_s(p) = p - m \tag{3.1.17}$$

3.2 Dirac Lagrangian and quantization

$$\mathcal{L} = \bar{\psi}(x)(i\partial \!\!\!/ - m)\psi(x) \tag{3.2.1}$$

Quantization

$$\left\{\psi_a(\mathbf{x}), \psi_b^{\dagger}(\mathbf{x}')\right\} = \delta_{ab}\delta^{(3)}(\mathbf{x} - \mathbf{x}') \tag{3.2.2}$$

$$\{\psi_a(\mathbf{x}), \psi_b(\mathbf{x}')\} = \{\psi_a^{\dagger}(\mathbf{x}), \psi_b^{\dagger}(\mathbf{x}')\} = 0$$
(3.2.3)

Field operators

$$\psi(\mathbf{x}) = \int \frac{\mathrm{d}^3 p}{(2\pi)^3 \sqrt{2E_p}} \sum_s (a_{\mathbf{p}}^s u_s(\mathbf{p}) e^{i\mathbf{p}\cdot\mathbf{x}} + b_{\mathbf{p}}^{s\dagger} v_s(\mathbf{p}) e^{-i\mathbf{p}\cdot\mathbf{x}})$$
(3.2.4)

thus the anticommutations of ladder operators:

$$\begin{cases}
a_{\mathbf{p}}^{s}, a_{\mathbf{p'}}^{s'\dagger} \\
\end{cases} = \begin{cases}
b_{\mathbf{p}}^{s}, b_{\mathbf{p'}}^{s'\dagger} \\
\end{cases} = (2\pi)^{3} \delta_{ss'} \delta^{(3)}(\mathbf{p} - \mathbf{p'})$$

$$\{a, a\} = \begin{cases}
a^{\dagger}, a^{\dagger} \\
\end{cases} = \dots = 0$$

Hamiltonian in terms of Fourier modes (with normal ordering)

$$H = \int \frac{\mathrm{d}^{3} p}{(2\pi)^{3}} \sum_{s} E_{p} (a_{\mathbf{p}}^{s\dagger} a_{\mathbf{p}}^{s} - b_{\mathbf{p}}^{s\dagger} b_{\mathbf{p}}^{s})$$
 (3.2.5)

3.3 Particles and antiparticles

$$Q = e \int d^3x \psi^{\dagger}(x)\psi(x)$$
 (3.3.1)

$$: Q := e \int \frac{\mathrm{d}^{3} p}{(2\pi)^{3}} \sum_{s} (a_{\mathbf{p}}^{s\dagger} a_{\mathbf{p}}^{s} - b_{\mathbf{p}}^{s\dagger} b_{\mathbf{p}}^{s})$$
 (3.3.2)

3.4 Dirac propagator and anticommutators

$$S_{ab}(x - y) = \{ \psi_a(x), \bar{\psi}_b(y) \}$$

= $(i\partial + m) [D(x - y) - D(y - x)]$ (3.4.1)

Time ordering of Dirac fields

$$T(\phi_a(x)\bar{\psi}_b(y)) = \Theta(x^0 - y^0)\psi_a(x)\bar{\psi}_b(y) - \Theta(y^0 - x^0)\bar{\psi}_b(y)\psi_a(x)$$
(3.4.2)

Feynman propogator for the Dirac field

$$S_F(x-y) = \langle 0|T\psi(x)\bar{\psi}(y)|0\rangle = \int \frac{d^4p}{(2\pi)^4} \frac{i(\not p+m)}{p^2 - m^2 + i\epsilon} e^{-ip\cdot(x-y)}$$
(3.4.3)

3.5 Discrete symmetries of the Dirac Field

	orientation perserving	orientation not perserving
(ortho)chronous	$\mathcal{L}_{\scriptscriptstyle{+}}^{\uparrow}$	$\mathcal{L}_{-}^{\uparrow}=\mathcal{P}\mathcal{L}_{+}^{\uparrow}$
non-orthochronous	$\mathcal{L}_{-}^{\downarrow}=\mathcal{T}\mathcal{L}_{+}^{\uparrow}$	$\mathcal{L}_{+}^{\downarrow} = \mathcal{PTL}_{+}^{\uparrow}$