Quantum Field Theory ST 2017

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1st Examination

1 Short questions

20p

(a) Derive the Euler-Lagrange equations for the Dirac and photon fields from the following Lagrangian

$$\mathcal{L} = \bar{\psi} \left(i \rlap{/}{D} - m \right) \psi - rac{1}{4} \epsilon^{\mu
u lpha \beta} F_{lpha eta} F_{\mu
u} \, .$$

4p

with $i\not D \psi = (i\not \partial - eA) \psi$ and $F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu}$.

(b) Consider a Yukawa theory in 6 dimensions with

$$\mathcal{L} = \bar{\psi} \left(i\partial \!\!\!/ - m \right) \psi + \frac{1}{2} \left(\partial_{\mu} \phi \, \partial^{\mu} \phi - M^2 \phi^2 \right) - g \bar{\psi} i \gamma^5 \psi \phi$$

Is this theory super-renormalizable, renormalizable, or non-renormalizable?

- Write down the Fourier decomposition of a massive fermion field $\psi(x)$ in the Heisenberg momentum p. Calculate $\langle 0|\psi(x)|\mathbf{p},r\rangle$, where $|\mathbf{p},r\rangle$ is a free particle with spin r and three
- (d) What are the propagators and vertex Feynman rules for the Lagrangian density

$$\mathcal{L} = \partial_{\mu}\phi^*\partial^{\mu}\phi - m^2\phi^*\phi + \frac{1}{2}\partial_{\mu}\pi\partial^{\mu}\pi - \sqrt{\lambda}\left(\phi^*\phi\right)^2 + \lambda\left(\phi^*\phi\right)\pi^2 - g\left(\phi^*\phi\right)\pi \quad ? \quad \mathbf{4p}$$

<u>Hint</u>: No proof is required. Be careful about the prefactors.

(e) Consider ϕ^4 theory with $\mathcal{L}_{int} = -\frac{\lambda}{4!}\phi^4$. Give the symmetry factors of the two diagrams



In addition prove them using Wick contractions.

 $\mathbf{q}_{\mathbf{0}}$

2 Proton-antiproton annihilation

20p

Consider a Yukawa theory given by

$$\mathcal{L} = \bar{\psi} \left(i \partial \!\!\!/ - m \right) \psi + \bar{\chi} \left(i \partial \!\!\!/ - \mu \right) \chi + \frac{1}{2} \left(\partial^{\mu} \!\!\!/ \!\!\!/ \!\!\!/ \!\!\!/ \!\!\!/ \!\!\!/ \!\!\!/ - M^2 \sigma^2 \right) - g \bar{\psi} i \gamma^5 \psi \dot{\phi} - \bar{g} \bar{\chi} i \gamma^5 \chi \dot{\phi} \,.$$

where $\psi(x)$ denotes the proton field. $\chi(x)$ the neutron field and $\phi(x)$ a real scalar field. The coupling constants are given by g and \bar{g} . Furthermore, note $\mu > m$. We now assign the spins s, \bar{s}, r, \bar{r} and the momenta p, \bar{p}, k, \bar{k} such that the process is given by

$$\psi^{(s)}(p) + \bar{\psi}^{(\bar{s})}(\bar{p}) \rightarrow \chi^{(r)}(k) + \bar{\chi}^{(\bar{r})}(\bar{k}) \,.$$

It is useful to define the Mandelstam variables $s=(p+\bar{p})^2$. $t=(p-k)^2$ and $u=(p-\bar{k})^2$

(a) What are the Feynman rules for the interaction vertices? <u>Hint</u>: Be careful about the factor i.

 \mathbf{q}

- (b) Which Feynman diagram(s) contribute to the process proton-antiproton to neutron-antineutron to the lowest order? Draw them, label the lines with momenta and spins. and determine the invariant matrix element $i\mathcal{M}$. 2 p
- Calculate $|\mathcal{M}|^2$ by averaging over the initial and summing over the final spin states. Show that it can be written in terms of Mandelstam variables as

$$\overline{|\mathcal{M}|^2} = \left(\frac{g\bar{g}}{s - M^2}\right)^2 \cdot s^2.$$
 10p

<u>Hint</u>: Use the relation $(\gamma^5)^{\dagger} = \gamma^5$.

(d) In the center-of-mass system (CMS) the differential cross section is given by

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = \frac{1}{64\pi^2 s} \frac{|\mathbf{p}_f|}{|\mathbf{p}_i|} \overline{|\mathcal{M}|^2}.$$

where \mathbf{p}_{ι} (\mathbf{p}_{f}) are the CMS particle three momenta of the initial (final) state. Express the differential cross section in terms of the Mandelstam variables. What happens for

(e) Consider now proton–neutron scattering

$$\psi^{(s)}(p) + \chi^{(\bar{r})}(\bar{k}) \rightarrow \psi^{(\bar{s})}(\bar{p}) + \chi^{(r)}(k) \,.$$

Calculate $|\mathcal{M}|^2$ for this process and express it in terms of the new Mandelstam variables. We define our new Mandelstam variables to be $\tilde{s}=(p+\bar{k})^2$. $\tilde{t}=(p-\bar{p})^2$. and $\tilde{u}=(p-k)^2$

3 Wave function renormalization

20p

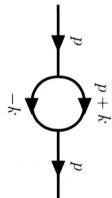
In the following we are going to study the self energy in o^3 -theory with renormalization. 4 dimensions) at next-to-leading order and determine its contribution to the wave function $\mathcal{L}_{\mathrm{int}}$ $\frac{\lambda}{3!}O^3$ (in

(a) Briefly show that the two-point function is given by

$$\int d^4x \, e^{ipx} \, \langle \Omega \, | T\phi(x)\phi(0) | \, \Omega \rangle = \frac{i}{p^2 - m_0^2 - \Sigma(p^2)} \, .$$

bare mass. Show that the correction to the wave function renormalization is given by $\frac{\mathrm{d}\Sigma(p^2)}{\mathrm{d}p^2}|_{p^2=m^2}$. where $-i\Sigma(p^2)$ denotes the sum of all one-particle-irreducible diagrams and m_0 the $-|_{p^2=m^2}$.

(b) Consider the following second-order contribution to $\Sigma(p^2)$ denoted by $\Sigma_2(p^2)$:



propagators using Feynman parameters, and simplify as far as possible. Write down the amplitude corresponding to the diagram in d dimensions, combine the

$$\frac{1}{AB} = \int_0^1 dx \frac{1}{[xA + (1-x)B]^2}$$

- (c) Perform a Wick rotation of the self energy contribution. What can you deduce about the degree of divergence of the loop integral for $\Sigma_2(p^2)$ and $\frac{\mathrm{d}\Sigma_2(p^2)}{\mathrm{d}p^2}$ for $d \to 4$? <u>Hint</u>: Perform the differentiation $\frac{d}{dp^2}$ under the integral.
- <u>a</u> We now want to calculate the contribution to the wave function renormalization. Perform a real number and should not be carried out. the Euclidean integral. The remaining integral over the Feynman parameter evaluates to

Hints:

1.
$$\int \frac{\mathrm{d}^{d} l_{E}}{(2\pi)^{d}} \frac{1}{(l_{E}^{2} + \Delta)^{n}} = \frac{1}{(4\pi)^{\frac{d}{2}}} \frac{\Gamma(n - \frac{d}{2})}{\Gamma(n)} \left(\frac{1}{\Delta}\right)^{n - \frac{d}{2}}$$

2.
$$\Gamma(x+1) = x \cdot \Gamma(x)$$
 and $\Gamma(1) = 1$

3. Note that
$$m^2 - m_0^2 = \mathcal{O}(\lambda^2)$$
.

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2nd Examination

1 Short questions

20p

Derive the Euler-Lagrange equations for the Dirac field from the following Lagrangian:

$$\mathcal{L}=ar{\psi}\left(iD\!\!\!/-m
ight)\psi-rac{1}{4}F_{\mu
u}F^{\mu
u}$$

with $\not D \psi = (\not \partial - i e \not A) \psi$ and $F_{\mu\nu} = \partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu}$.

If m=0 the Lagrangian is invariant under the global transformation $\psi \to e^{i\alpha\gamma^5}\psi$ $\bar{\psi} \to \psi^{\dagger} e^{-i\alpha\gamma^5}\gamma^0$. Deduce the Noether current.

(b) Consider the following interaction terms:

$$\mathcal{L}_{\mathrm{int}} = c_1 \left(ilde{F}_{\mu
u} ilde{F}^{\mu
u}
ight)^2 + c_2 \left(ilde{\psi} i \sigma^{\mu
u} \psi
ight) F_{\mu
u} + c_3 \left(ilde{\psi} i \gamma^5 A \psi
ight) \phi^2 + c_4 \left(\phi(\partial^\mu \phi) A_\mu
ight) .$$

with the Klein-Gordon field ϕ . Dirac field ψ , electromagnetic field A_{μ} , its corresponding

field strength tensor $F_{\mu\nu}$, its dual $\tilde{F}_{\mu\nu} = \frac{1}{2} \epsilon_{\mu\nu\alpha\beta} F^{\alpha\beta}$, and $\sigma^{\mu\nu} = \frac{i}{2} [\gamma^{\mu}, \gamma^{\nu}]$. Determine the mass dimension of the coupling constants c_i , $i = 1, \ldots, 4$, in d = 4 direnormalizable, or non-renormalizable. mensions and deduce whether the individual interaction terms are super-renormalizable.

(c) Why do the creation and annihilation operators of Dirac particles fulfill anticommutation instead of commutation relations? What is the result of $\left(a_{\mathbf{p}}^{(s)}\right)^{\dagger}\left(a_{\mathbf{q}}^{(r)}\right)^{\dagger}|0\rangle$? **2p**

Show that the superficial degree of divergence for an arbitrary loop-diagram in ϕ^4 theory can be written as

$$D = d + (d - 1) \cdot V - \left(\frac{d}{2} - 1\right) \cdot E.$$

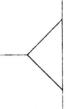
the number of external lines. where d denotes the space-time-dimension, V stands for the number of vertices and E for

Discuss the renormalizability of the theory for d < 4, d = 4, and d > 4.

> ±. 61

You may want to use the number of loops L and the number of internal lines I in the course of the derivation.

(e) Consider $\mathcal{L}_{int} = -\frac{\lambda}{3!}\phi^3$. What is the symmetry factor of the following diagram?



Proof your result using Wick contractions.

2 Electron-positron annihilation

20p

grangian for this process is then given by: We consider electron-positron annihilation into a muon-antimuon pair. The interaction La-

$$\mathcal{L}_{\mathrm{int}} = -\epsilon \bar{\psi} \gamma^{\mu} A_{\mu} \psi - \epsilon \bar{\xi} \gamma^{\mu} A_{\mu} \xi.$$

where ψ is the electron field. ξ the muon field, and $\epsilon = |\epsilon|$ is the elementary charge. We assign momenta p, p', k, k' and spins s, s', r, r' in the following way:

$$\psi^{(s)}(p) + \bar{\psi^{(s')}}(p') \rightarrow \xi^{(r)}(k) + \bar{\xi}^{(r')}(k')$$
.

but <u>not</u> the muon mass m_{μ} . Assume that we are in an energy region where we can neglect the electron mass m_e $(m_e \approx 0)$

It is useful to define the Mandelstam variables $s = (p + p')^2$, $t = (p - k)^2$, and $u = (p - k')^2$

(a) Show

1)
$$s - 2pp' - 2m_{\mu}^2 + 2kk'$$

2) $t = m_{\mu}^2 - 2pk = m_{\mu}^2 - 2p'k'$
3) $u = m_{\mu}^2 - 2pk' - m_{\mu}^2 - 2p'k$
4) $s + t + u = 2m_{\mu}^2$

- (b) Which diagram contributes to the leading order? Draw it, label the lines with momenta and spins, and determine the invariant matrix element \mathcal{M} .
- (c) Calculate $|\mathcal{M}|^2$ by averaging $|\mathcal{M}|^2$ over the initial spins and summing over the final spin states. Simplify $|\mathcal{M}|^2$ as far as possible (using $m_e = 0$) and verify that it reduces to

$$\overline{|\mathcal{M}|^2} - \frac{2e^4}{8^2} \left(t^2 + u^2 + 4m_{\mu}^2 s - 2m_{\mu}^4\right) .$$

$$\text{Tr} \left(\gamma^{\mu}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}\right) = 4\left(g^{\mu\nu}g^{\rho\sigma} - g^{\mu\rho}g^{\nu\sigma} + g^{\mu\sigma}g^{\nu\rho}\right)$$
12p

(d) Assume now to be in the ultrarelativistic case in which $m_{\mu} \approx 0$. In the center-of-mass system (CMS) the differential cross section is given by

$$rac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = rac{1}{64\pi^2 \mathrm{s}} \cdot \overline{|\mathcal{M}|^2}.$$

Express the cross section in terms of the Mandelstam variable s and the angle Θ between the incoming electron (\mathbf{p}) and the outgoing muon three momentum (\mathbf{k}) .

3 The electron self energy

20p

In the following we are going to study the electron self energy in QED at next-to-leading order.

(a) Briefly show that the two-point function is given by

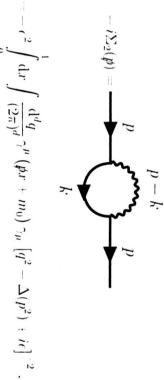
$$\int \mathrm{d}^{4}x\,\epsilon^{ipx}\left\langle \Omega\left|T\psi(x)\bar{\psi}(0)\right|\Omega\right\rangle =\frac{i^{-1}}{\not\!p-m_{0}-\Sigma\left(\not\!p\right)}\,.$$

where $-i\Sigma(p)$ denotes the sum of all one-particle-irreducible diagrams and m_0 the bare electron mass. What is the correction to the physical electron mass m?

4p.

Hint: Use the relation $\frac{i(p+m_0)}{p^2-m_0^2} = \frac{i}{p-m_0}$. Furthermore use that the geometric sum for a matrix valued operator A is given by $\sum_{n=0}^{\infty} A^n = [\mathbb{I} - A]^{-1}$.

(b) Show that the second-order self energy contribution $\Sigma_2(p)$ is defined by



with $\Delta(p^2) = (1-x)(m_0^2 - xp^2)$. <u>Hint</u>: Use the relation $\frac{1}{AB} = \int_0^1 \mathrm{d}x \frac{1}{[xA + (1-x)B]^2}$.

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(c) Show that in d dimensions the Dirac structure reduces to

$$\gamma^{\mu} (px + m_0) \gamma_{\mu} = (2 - d)xp + dm_0.$$
 2p

- (d) What is the mass dimension of the coupling constant ϵ in d dimensions? Rewrite it as $\epsilon = \mu^{\alpha} \epsilon_{0}$, where ϵ_{0} has the same mass dimension as ϵ in four dimensions. What is α ? 2 p
- (e) Perform a Wick rotation in order to evaluate the momentum integral.

 $^{2}\mathrm{p}$

Use the relation $\int rac{\mathrm{d}^d q_E}{(2\pi)^d} rac{1}{(q_E^2 + \Delta)^n} = rac{1}{(q_E^2 + \Delta)^n}$ $(4\pi)^{\frac{d}{2}}$ $rac{\Gamma\left(n-rac{d}{2}
ight)}{\Gamma(n)}\left(rac{1}{\Delta}
ight)^{n-rac{d}{2}}$

(f) Identify and calculate the divergent part of the mass renormalization in the variable

Hints: Use the relations:

1)
$$\Gamma(2) = 1$$
 and $\Gamma(x) = \frac{1}{x} - \Gamma_E + \mathcal{O}(x^2)$

$$2) \quad a' = 1 + \epsilon \ln a$$

3)
$$m_0 = m + \mathcal{O}(\epsilon^2)$$