

H.S a)

\swarrow P is operator in Hilbert space
acts only on a_s, a_s^\dagger

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$$\begin{aligned} P \psi(x,t) P^{-1} & \quad (P = P^\dagger = P^{-1}) \\ &= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_s [P a_s(p) P u_s(p) e^{-ipx} + P b_s^\dagger(p) P v_s(p) e^{+ipx}] \\ &= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_s [\eta_a a_s(-p) \underbrace{u_s(p)}_{=\gamma^0 u_s(-p)} e^{-ipx} - \eta_b^* b_s^\dagger(-p) \underbrace{v_s(p)}_{=\gamma^0 v_s(-p)} e^{+ipx}] \end{aligned}$$

$$= \eta_a \gamma^0 \psi(t, -x) |0\rangle, \text{ if } \eta_a = -\eta_b^*$$

$\eta_a^2 = \eta_b^2 = \pm 1$, since its fermionic field, only even number of a_p or a_q

$$\Rightarrow \eta_a \eta_b = -\eta_a \eta_a^* = -1, \text{ i.e. } \eta_a \eta_a^* = 1$$

$$\eta_a = x + iy : (x + iy)(x - iy) = x^2 + y^2 = 1$$

$$(x + iy)(x + iy) = \pm 1, \quad 2ixy = 0, \quad x=0 \text{ or } y=0$$

$$\Rightarrow \eta_a = \pm 1, \pm i$$

$$\eta_b = \mp 1, \pm i$$

b) $T \psi(x,t) T^{-1} \quad (T = T^\dagger = T^{-1})$

$$\begin{aligned} &= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_s [T a_s(p) u_s(p) e^{-ipx} T + T b_s^\dagger(p) v_s(p) e^{+ipx} T] \\ &= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_s [(-)^{\frac{1}{2}+s} a_{-s}(-p) u_s^*(p) e^{+ipx} + (-)^{\frac{1}{2}+s} b_{-s}^\dagger(-p) v_s^*(p) e^{-ipx}] \end{aligned}$$

$$\left[\begin{aligned} i\gamma^5 \gamma^2 \gamma^0 u_s(p) &= (-)^{\frac{1}{2}-s} [u_{-s}(-p)]^* \\ \text{LHS} &= -\gamma^0 \gamma^1 \gamma^2 \gamma^3 \gamma^2 \gamma^0 = +\gamma^0 \gamma^1 (\gamma^2)^2 \gamma^3 \gamma^0 = -\gamma^0 \gamma^1 \gamma^3 \gamma^0 \\ &= -(\gamma^0)^2 \gamma^1 \gamma^3 = -\gamma^1 \gamma^3 \end{aligned} \right.$$

$$\Rightarrow -\gamma^1 \gamma^3 u_s(p) = (-)^{\frac{1}{2}-s} [u_{-s}(-p)]^*$$

$$\Leftrightarrow -(-)^{s-\frac{1}{2}} \gamma^1 \gamma^3 u_s(p) = [u_{-s}(-p)]^*$$

$$\Leftrightarrow -(-)^{-s-\frac{1}{2}} \gamma^1 \gamma^3 u_{-s}(-p) = [u_s(p)]^*$$

analogically $-(-)^{-s-\frac{1}{2}} \gamma^1 \gamma^3 v_{-s}(-p) = [v_s(p)]^*$

$$= -\gamma^1 \gamma^3 \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_s [a_{-s}(-p) u_{-s}(-p) e^{+ipx} + b_{-s}^\dagger(-p) v_{-s}(-p) e^{-ipx}] |0\rangle$$

inside integral $-p \rightarrow \tilde{p}$, $p = (E_p, p) \rightarrow \tilde{p} = (E_p, -p)$

$$= -\gamma^1 \gamma^3 \int \frac{d^3 \tilde{p}}{(2\pi)^3} \frac{1}{\sqrt{2E_{\tilde{p}}}} \sum_s [a_{-s}(\tilde{p}) u_{-s}(\tilde{p}) e^{-i\tilde{p}(-t, \underline{x})} + b_{-s}^\dagger(\tilde{p}) v_{-s}(\tilde{p}) e^{i\tilde{p}(-t, \underline{x})}] |0\rangle$$

$$= -\gamma^1 \gamma^3 \psi(-t, \underline{x}) |0\rangle$$

c) $C \psi(\underline{x}, t) C^{-1}$ $(C^\dagger = C = C^{-1} = -i\gamma^2)$

$$= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_s \left[\underbrace{C a_s(p) C}_{b_s(p)} \underbrace{u_s(p) e^{-ipx}}_{= -i\gamma^2 v_s^*(p)} + \underbrace{C b_s^\dagger(p) C}_{a_s^\dagger(p)} \underbrace{v_s(p) e^{-ipx}}_{= -i\gamma^2 u_s^*(p)} \right]$$

$$= -i\gamma^2 \psi^*(\underline{x}, t)$$

H.6

How do we get expression (19)?

a) $P |\bar{\Psi}, S=0, 1\rangle$ → wave function to bind two particles

$$= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \phi_s(|p|) P \left\{ \underbrace{a_{\frac{1}{2}}^{\dagger}(p)}_{PP} \underbrace{b_{-\frac{1}{2}}^{\dagger}(p)}_{PP} + (-1)^{S-1} \underbrace{a_{-\frac{1}{2}}^{\dagger}(p)}_{PP} \underbrace{b_{\frac{1}{2}}^{\dagger}(-p)}_{PP} \right\} |0\rangle$$

one for singlet
three for triplet

$$= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \left\{ \eta_a^* \eta_b^* a_{\frac{1}{2}}^{\dagger}(-p) b_{-\frac{1}{2}}^{\dagger}(+p) + (-1)^{S-1} \eta_a^* \eta_b^* a_{-\frac{1}{2}}^{\dagger}(-p) b_{\frac{1}{2}}^{\dagger}(+p) \right\} |0\rangle$$

$$= \eta_a^* \eta_b^* \int \frac{d^3 \tilde{p}}{(2\pi)^3} \frac{1}{\sqrt{2E_{\tilde{p}}}} \left\{ a_{\frac{1}{2}}^{\dagger}(\tilde{p}) b_{-\frac{1}{2}}^{\dagger}(-\tilde{p}) + (-1)^{S-1} a_{-\frac{1}{2}}^{\dagger}(\tilde{p}) b_{\frac{1}{2}}^{\dagger}(-\tilde{p}) \right\} |0\rangle$$

$$= -|\eta_a|^2 |\bar{\Psi}, S=0, 1\rangle$$

$$= -|\bar{\Psi}, S=0, 1\rangle$$

positive parity!

b) $C |\bar{\Psi}, S=0, 1\rangle$

$$= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \phi_s(|p|) \left\{ C a_{\frac{1}{2}}^{\dagger}(p) C C b_{-\frac{1}{2}}^{\dagger}(-p) + (-1)^{S-1} C a_{-\frac{1}{2}}^{\dagger}(p) C C b_{\frac{1}{2}}^{\dagger}(-p) \right\} C C |0\rangle$$

$$= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \phi_s(|p|) \left\{ b_{\frac{1}{2}}^{\dagger}(p) a_{-\frac{1}{2}}^{\dagger}(-p) + (-1)^{S-1} b_{-\frac{1}{2}}^{\dagger}(p) a_{\frac{1}{2}}^{\dagger}(-p) \right\} |0\rangle$$

$$= (-1)(-1)^{S-1} \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \phi_s(|p|) \left\{ (-1)^{S-1} a_{-\frac{1}{2}}^{\dagger}(-p) b_{\frac{1}{2}}^{\dagger}(p) + a_{\frac{1}{2}}^{\dagger}(-p) b_{-\frac{1}{2}}^{\dagger}(p) \right\} |0\rangle$$

$$\left[((-1)^{S-1})^2 \equiv 1, \quad \{a_s, b_s\} = \{a_s^{\dagger}, b_s^{\dagger}\} = 0 \right]$$

changing $-p \rightarrow \tilde{p}$

$$= (-1)^S |\bar{\Psi}, S=0, 1\rangle$$

c) $C A_{\mu} C^{-1} = -A_{\mu} \Rightarrow$ photo is its own antiparticle
eigenvalue -1

$$C |n\rangle = C |1\rangle \otimes \underbrace{\dots}_{n} \otimes |1\rangle$$

$$= (-1)^n |n\rangle$$

$$C|\bar{\Psi}, S=0\rangle = |\bar{\Psi}, S=0\rangle \Rightarrow \text{decay into even numbers of photons} \\ 2, 4, 6, \dots$$

$$C|\bar{\Psi}, S=1\rangle = -|\bar{\Psi}, S=1\rangle \Rightarrow \text{decay into odd numbers of photons} \\ 3, 5, 7, \dots$$

For these two decays, same initial state, $S=1$ decay has more possibilities for final states.

$$\Rightarrow P_{S=1} > P_{S=0}$$

$$\Rightarrow \tau_{S=1} < \tau_{S=0}, \text{ not consistent...}$$

2x bigger phase space \rightarrow decay easier
3x