

H.4

$$\mathcal{L} = \bar{\psi} (i \not{\partial} - m) \psi$$

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a) EL-eq. $\partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \psi)} \right) - \frac{\partial \mathcal{L}}{\partial \psi} = 0$

$$\partial_\mu (\bar{\psi} i \gamma^\mu) + m \bar{\psi} = 0$$

$$\bar{\psi} (i \gamma^\mu \overleftarrow{\partial}_\mu + m) = 0$$

$$\partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \bar{\psi})} \right) - \frac{\partial \mathcal{L}}{\partial \bar{\psi}} = 0$$

$(i \not{\partial} - m) \psi = 0$ think of $\not{\partial}$ acting on ψ

hermitian conjugation:

$$\left(\begin{array}{l} \psi^\dagger (i \gamma^\mu \partial_\mu - m)^\dagger = 0 \\ \psi^\dagger (-i \gamma^0 \gamma^\mu \gamma^0 \partial_\mu - m) = 0 \\ \psi^\dagger (-i \gamma^0 \gamma^\mu \partial_\mu - m \gamma^0) = 0 \\ \rightarrow \bar{\psi} (i \gamma^\mu \overleftarrow{\partial}_\mu + m) = 0 \end{array} \right)$$

b) $T^{\mu\nu} = \frac{\partial \mathcal{L}}{\partial (\partial_\mu \psi)} \partial^\nu \psi + (\partial^\nu \bar{\psi}) \frac{\partial \mathcal{L}}{\partial (\partial_\mu \bar{\psi})} - g^{\mu\nu} \mathcal{L}$

$$= \bar{\psi} i \gamma^\mu \partial^\nu \psi - g^{\mu\nu} \bar{\psi} (i \not{\partial} - m) \psi$$

if one uses EOM here, $T^{\mu\nu} = i \bar{\psi} \gamma^\mu \partial^\nu \psi$

(why are we allowed to use EOM here?)

PS: one can get symmetric $T^{\mu\nu}$ using symmetric \mathcal{L}



c) $(i\not{\partial} - m)\psi = 0$, ψ is a solution

$m=0$: $i\not{\partial}\psi = 0$

$$i\not{\partial}\gamma^5\psi = i\gamma^\mu\gamma^5\partial_\mu\psi = -i\gamma^5\gamma^\mu\partial_\mu\psi = 0$$

$$(\{\gamma^\mu, \gamma^5\} = 0)$$

$\Rightarrow \gamma^5\psi$ is also a solution

d) $P_{L/R} = \frac{1}{2}(\mathbb{1} \mp \gamma^5)$

i) $P_{L/R}^2 = \frac{1}{4}(\mathbb{1} \mp \gamma^5)^2 = \frac{1}{4}(\mathbb{1} + (\gamma^5)^2 \mp 2\mathbb{1}\gamma^5)$

$$= \frac{1}{4}(2\mathbb{1} \mp 2\mathbb{1}\gamma^5) = \frac{1}{2}(\mathbb{1} \mp \gamma^5) = P_{L/R}$$

✓

$$P_L P_R = P_R P_L = 0$$

$$P_L + P_R = \mathbb{1}$$

ii) $\psi_{L/R} := P_{L/R}\psi$

$$\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix},$$

$$P_{L/R}\psi = \frac{1}{2} \begin{pmatrix} \mathbb{1} \pm \mathbb{1} & 0 \\ 0 & \mathbb{1} \mp \mathbb{1} \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} \quad (\text{using chiral rep.})$$

$$= \begin{cases} \begin{pmatrix} \psi_1 \\ 0 \end{pmatrix} = \psi_L, L \\ \begin{pmatrix} 0 \\ \psi_2 \end{pmatrix} = \psi_R, R \end{cases}$$

$$\Rightarrow \psi = \psi_L + \psi_R$$

$$\mathcal{L} = \bar{\psi}(i\not{\partial} - m)\psi = (\bar{\psi}_L + \bar{\psi}_R)(i\not{\partial} - m)(\psi_L + \psi_R)$$

$$= \bar{\psi}_L(i\not{\partial} - m)\psi_L + \underbrace{\bar{\psi}_L(i\not{\partial} - m)\psi_R + \bar{\psi}_R(i\not{\partial} - m)\psi_L}_{\text{diagonal and}} + \bar{\psi}_R(i\not{\partial} - m)\psi_R$$

$$(1, 0) \begin{pmatrix} 0 \\ 1 \end{pmatrix} = 0$$

$$= \bar{\psi}_L(i\not{\partial} - m)\psi_L + \bar{\psi}_R(i\not{\partial} - m)\psi_R = \mathcal{L}_L + \mathcal{L}_R$$

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$$\gamma^\mu P_{RL} = P_{LR} \gamma^\mu, \quad \bar{\psi} P_{LR} = \psi^\dagger \gamma^0 P_{LR} = \psi^\dagger P_{RL} \gamma^0 = (P_{RL} \psi)^\dagger \gamma^0 = \bar{\psi}_{RL}$$

$$\text{kinetic term} = i\bar{\psi} \not{\partial} (P_L + P_R) \psi = i\bar{\psi} \not{\partial} P_L \psi + i\bar{\psi} \not{\partial} P_R \psi \quad | \text{ and use } P_L P_R = P_R P_L = 0$$

$$= i\bar{\psi} P_R \not{\partial} P_L \psi + i\bar{\psi} P_L \not{\partial} P_R \psi = i\bar{\psi}_L \not{\partial} \psi_L + i\bar{\psi}_R \not{\partial} \psi_R$$

$$\text{mass term} \quad m\bar{\psi} \psi = m\bar{\psi}_L \psi_R + m\bar{\psi}_R \psi_L \quad \leftarrow \text{Dirac mass term is not chiral symmetric!}$$

e) the trafo $\begin{Bmatrix} \psi \\ \psi^\dagger \end{Bmatrix} \longrightarrow \begin{Bmatrix} \gamma^5 \psi \\ \psi^\dagger \gamma^5 \end{Bmatrix}$ If one wants to preserve CS, no Dirac mass, Mass from Higgs

i.e. $\begin{Bmatrix} \psi \\ \bar{\psi} \end{Bmatrix} = \begin{Bmatrix} \psi \\ \psi^\dagger \gamma^0 \end{Bmatrix} \longrightarrow \begin{Bmatrix} \psi \\ \psi^\dagger \gamma^5 \gamma^0 \end{Bmatrix} = \begin{Bmatrix} \psi \\ -\bar{\psi} \gamma^5 \end{Bmatrix}$ \uparrow Standard Model

$$\bar{\psi} \psi \longrightarrow -\bar{\psi} \gamma^5 \gamma^5 \psi = -\bar{\psi} \psi$$

$$\bar{\psi} i \not{\partial} \psi \longrightarrow -\bar{\psi} \gamma^5 i \not{\partial} \gamma^5 \psi = -\bar{\psi} i \not{\partial} \psi$$

$$\bar{\psi} \gamma^\mu \psi \longrightarrow -\bar{\psi} \gamma^5 \gamma^\mu \gamma^5 \psi = +\bar{\psi} \gamma^\mu \psi$$

$\{\gamma^5, \gamma^\mu\} = 0$

$$\bar{\psi} \gamma^\mu \gamma^5 \psi \longrightarrow -\bar{\psi} \gamma^5 \gamma^\mu \gamma^5 \gamma^5 \psi = +\bar{\psi} \gamma^\mu \gamma^5 \psi$$

$$\bar{\psi} \sigma^{\mu\nu} \psi \longrightarrow -\bar{\psi} \gamma^5 \sigma^{\mu\nu} \gamma^5 \psi = \bar{\psi} \sigma^{\mu\nu} \psi$$

$$\left(\begin{array}{l} \sigma^{\mu\nu} = \frac{i}{2} [\gamma^\mu, \gamma^\nu] \\ \{\sigma^{\mu\nu}, \gamma^5\} = \frac{i}{2} \{\gamma^\mu \gamma^\nu - \gamma^\nu \gamma^\mu, \gamma^5\} = 0 \end{array} \right)$$

f) $\partial_\mu V^\mu(x) = \partial_\mu (\bar{\psi} \gamma^\mu \psi) = \partial_\mu \bar{\psi} \gamma^\mu \psi + \bar{\psi} \gamma^\mu \partial_\mu \psi$

$$[(i\not{\partial} - m)\psi = 0; \bar{\psi}(i\not{\partial} + m) = 0]$$

$$= im\bar{\psi} \psi + \bar{\psi}(-im)\psi = 0$$

$$\partial_\mu A^\mu(x) = \partial_\mu (\bar{\psi} \gamma^\mu \gamma^5 \psi) = \partial_\mu \bar{\psi} \gamma^\mu \gamma^5 \psi + \bar{\psi} \gamma^\mu \gamma^5 \partial_\mu \psi$$

$$= \not{\partial} \bar{\psi} \gamma^5 \psi - \bar{\psi} \gamma^5 \not{\partial} \psi$$

$$= im\bar{\psi} \gamma^5 \psi + i\bar{\psi} \gamma^5 m \psi$$

$$\partial_\mu A^\mu = 0 \quad \text{iff} \quad m=0$$

$$\psi(x) \longrightarrow e^{i\alpha} \psi(x), \quad \alpha \Delta \psi = i\alpha \psi$$

$$\alpha \Delta \bar{\psi} = -i\alpha \bar{\psi}$$

$$\begin{aligned}
 j^\mu &= \frac{\partial \mathcal{L}}{\partial (\partial_\mu \psi)} \partial^\mu \psi + \frac{\partial \mathcal{L}}{\partial (\partial_\mu \bar{\psi})} \partial^\mu \bar{\psi} - \underbrace{X^\mu}_{=0} \\
 &= \bar{\psi} \cdot i \gamma^\mu \cdot i \partial^\mu \psi \\
 &= -i \bar{\psi} \gamma^\mu \psi \propto V(x)
 \end{aligned}$$

$$g) \quad \mathcal{L} = \bar{\psi} (i \not{\partial} - m) \psi$$

$$\rightarrow \tilde{\mathcal{L}} = \mathcal{L} + \partial_\mu X^\mu(x)$$

$$= \frac{i}{2} \bar{\psi} (\gamma^\mu \vec{\partial}_\mu - m) \psi - \frac{i}{2} \bar{\psi} (\overleftarrow{\partial}_\mu \gamma^\mu + m) \psi$$

$$\partial_\mu X^\mu(x) = -\frac{i}{2} \bar{\psi} (\gamma^\mu \vec{\partial}_\mu - \cancel{\gamma^\mu}) \psi - \frac{i}{2} \bar{\psi} (\overleftarrow{\partial}_\mu \gamma^\mu + \cancel{\gamma^\mu}) \psi$$

$$= -\frac{i}{2} \bar{\psi} \gamma^\mu (\vec{\partial}_\mu + \overleftarrow{\partial}_\mu) \psi = -\frac{i}{2} \partial_\mu (\bar{\psi} \gamma^\mu \psi)$$

the difference can be rewritten into total

divergence, thus the EOM from $\tilde{\mathcal{L}}$ is the same from \mathcal{L}

$$\Rightarrow \tilde{T}^{\mu\nu} = \frac{\partial \mathcal{L}}{\partial (\partial_\mu \psi)} \partial^\nu \psi + (\partial^\nu \bar{\psi}) \frac{\partial \mathcal{L}}{\partial (\partial_\mu \bar{\psi})} - \underbrace{g^{\mu\nu} \mathcal{L}}_{=0}$$

$$= \frac{i}{2} \bar{\psi} \gamma^\mu (\partial^\nu \psi) - (\partial^\nu \bar{\psi}) \frac{i}{2} \gamma^\mu \psi$$

$$= \frac{i}{2} \bar{\psi} \gamma^\mu (\vec{\partial}^\nu - \overleftarrow{\partial}^\nu) \psi$$

$$\mathcal{L} = \frac{i}{2} \bar{\psi} (\gamma^\mu \vec{\partial}_\mu - m) \psi - \frac{i}{2} \bar{\psi} (\overleftarrow{\partial}_\mu \gamma^\mu + m) \psi$$

$$\partial_\mu \frac{\partial \mathcal{L}}{\partial (\partial_\mu \psi)} - \frac{\partial \mathcal{L}}{\partial \psi} = 0 \Rightarrow \bar{\psi} (i \not{\partial} + m) = 0$$

$$\text{w.r.t. } \bar{\psi} \quad \Rightarrow \quad (i \not{\partial} - m) \psi = 0$$