

Theoretical Astroparticle Physik

Homework 5

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1 Time of recombination from photon last scattering

- a) Assume the expansion was only caused by non-relativistic matter with $\Omega_m = 0.27$. From previous homework, we know its energy density is proportional to its number density

$$\rho_M = m_M \cdot n_M \quad (1.1)$$

Per definition the ratios of Ω is the same as the ratio of energy density

$$m_M n_M = \frac{\Omega_M}{\Omega_B} m_p n_B^{tot}(T) \quad (1.2)$$

and we say the baryonic matter in this epoch is simply protons or hydrogen atoms, thus the mass $m_B = m_p$.

Again the baryonic matter density is directly related to photon density

$$n_B^{tot} = \eta_B \cdot n_\gamma(T) = \eta_B \frac{2\zeta(3)}{\pi^2} T^3 \quad (1.3)$$

Putting these three equations (1.1), (1.2), and (1.3) together

$$\rho_M = \frac{\Omega_M}{\Omega_B} m_p \eta_b \frac{2\zeta(3)}{\pi^2} T^3 \quad (1.4)$$

In matter-dominated epoch (the first equation is valid for all types of Universe with zero curvature), we have

$$\rho = \frac{3}{8\pi G} H^2 \quad (1.5)$$

$$H(t) = \frac{2}{3t} \quad (1.6)$$

Then the Hubble parameter as last scattering is

$$\begin{aligned}
H(T_r) &= \left[\frac{8\pi G}{3} \frac{\Omega_M}{\Omega_B} m_p \eta_b \frac{2\zeta(3)}{\pi^2} T^3 \right]^{1/2} \\
&= \left[\frac{8\pi}{3} \frac{0.27}{0.046} 6.2 \cdot 10^{-10} \frac{2 \cdot 1.20 \cdot 938.3 \text{ MeV} \cdot (0.26 \text{ eV})^3}{\pi^2 (1.22 \times 10^{19} \text{ GeV})^2} \right]^{1/2} \\
&= 2.87 \times 10^{-38} \text{ GeV} \\
&= 1.435 \times 10^{-24} \text{ cm}^{-1} \\
&= 4.45 \text{ Mpc}^{-1}
\end{aligned} \tag{1.7}$$

The time can be computed with (1.6)

$$t = \frac{2}{3H} = 2.32 \times 10^{37} \text{ GeV}^{-1} = 15.3 \times 10^{12} \text{ s} = 4.85 \times 10^5 \text{ year} \tag{1.8}$$

- b) Now we have matter and radiation contributions. The Friedmann equation becomes

$$H^2(t) = \frac{8\pi G}{3} \rho = \frac{8\pi G}{3} (\rho_{\text{rad}} + \rho_M)$$

Energy density of radiation, considering all relativistic species, follows

$$\rho_{\text{rad}} = g_* \frac{\pi^2}{30} T^4$$

where g_* is the effective degrees of freedom. Without exact knowledge of the Universe, in particular the relativistic particles, g_* is hard to determine. Thus we don't really want to plug it into Friedmann equation.

However, we also know that the energy density of radiation $\rho_{\text{rad}} \sim a^{-4}$. Then we can write

$$a = \frac{c}{T} \tag{1.9}$$

with c a numerical constant and its exact value will be irrelevant.

The Friedmann equation with both contributions reads

$$\begin{aligned}
\left(\frac{\dot{a}}{a} \right)^2 &= \frac{8\pi G}{3} \rho_c \left[\Omega_M \left(\frac{a_0}{a} \right)^3 + \Omega_{\text{rad}} \left(\frac{a_0}{a} \right)^4 \right] \\
\frac{da}{dt} &= H_0 \left(\Omega_M \frac{a_0^3}{a} + \Omega_{\text{rad}} \frac{a_0^4}{a^2} \right)^{1/2} \\
\int_0^{t_r} dt &= \frac{1}{H_0} \int_0^{a_r} da \left(\Omega_M \frac{a_0^3}{a} + \Omega_{\text{rad}} \frac{a_0^4}{a^2} \right)^{-1/2}
\end{aligned}$$

Do a simple substitution

$$da = -\frac{a}{T} dT$$

The integral on RHS gets flipped because of the negative sign

$$\begin{aligned} \int_0^{t_r} dt &= \frac{1}{H_0} \int_{T_r}^{\infty} \frac{dT}{T} \left(\Omega_M \frac{a_0^3}{a^3} + \Omega_{\text{rad}} \frac{a_0^4}{a^4} \right)^{-1/2} \\ t_r &= \frac{1}{H_0} \int_{T_r}^{\infty} \frac{dT}{T} \left(\Omega_M \frac{T^3}{T_0^3} + \Omega_{\text{rad}} \frac{T^4}{T_0^4} \right)^{-1/2} \end{aligned}$$

With Python and its package, the integral is numerically evaluated

$$t_r = 2.1 \times 10^6 \text{ year} \quad (1.10)$$

which cannot be right.

- c) The Universe transitioned from radiation dominated to matter dominated at 0.76 eV. Thus in order to get the correct time of last scattering, one need to include radiation.