



CLASSMATE
BECAUSE YOU ARE ONE OF A KIND

Optics {concept}



Lake Como, the third largest and one of the most beautiful lakes in Italy, has been famously described as 'the looking glass of Venus'. The sprawling lake covers an area of 146 square kilometres.

Notebook

I N D E X

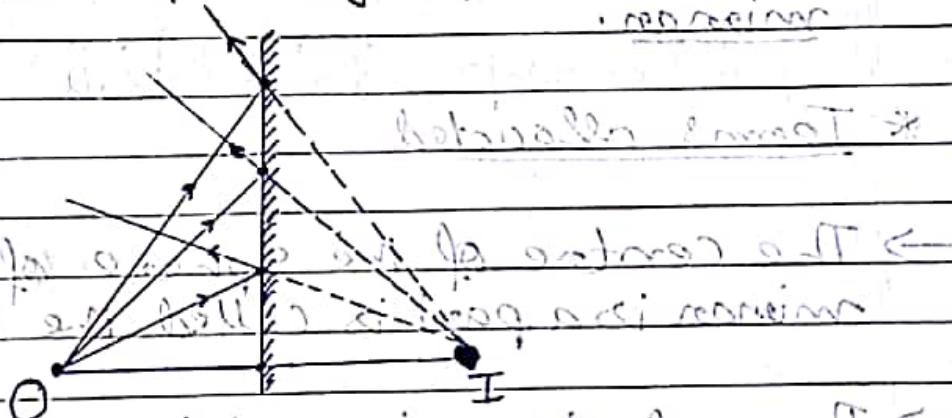
NAME: Aditya STD.: XII SEC.: A ROLL NO.: 01 SUB.: Physics Optics

Geometrical Optics

⇒ Laws of reflection

- ① The angle of incidence is equal to the angle of reflection.
- ② The incident ray, the reflected ray and the normal to the reflecting surface are coplanar.

* Formation of Image



→ All reflected rays meet at I when produced

→ Behind the mirror are no thing

→ An eye receiving the reflected rays feels that the rays are diverging from the point I.

→ The point I is called Image of Object O.

⇒ Spherical mirrors

→ A spherical mirror is a part cut from a hollow sphere with one surface silvered.

→ If reflection takes place at the convex surface; it is called a convex mirror.
and if reflection takes place at the concave surface; it is called a concave mirror.

* Terms associated

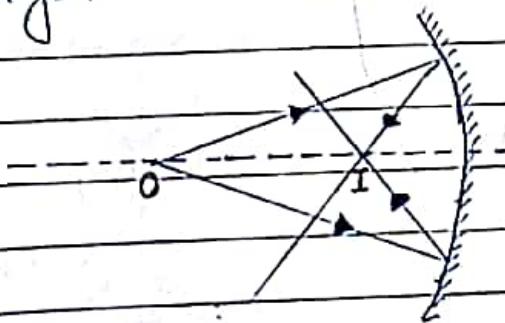
→ The centre of the sphere of which the mirror is a part is called the Centre of curvature.

→ The radius of sphere is called the radius of curvature.

→ The point on the mirror at the middle of the surface is called its pole.

→ The line joining the pole and the centre of curvature is called the Principal axis.

- * Focus \Rightarrow The point where reflected rays converge is called the focus of the mirror.
- * Focal plane \Rightarrow A plane through focus and perpendicular to principle axis is focal plane.
- * focal length \Rightarrow Distance between focus and pole is focal length.
- * Paraxial Rays \Rightarrow A ray close to the principle axis is called
- * Image tracing \rightarrow Draw Incident rays and reflected rays.
- \rightarrow The point of intersection of the incident rays is called the object & the point of intersection of the corresponding reflected rays is called its image.



Sign Convention

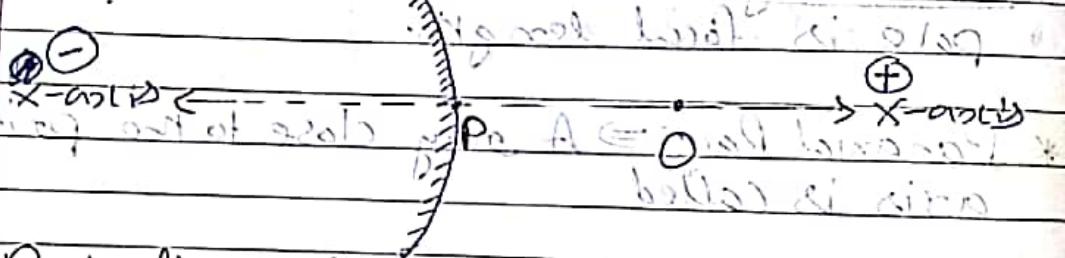
(Pole is taken as origin)

→ Pole is taken to be Origin.

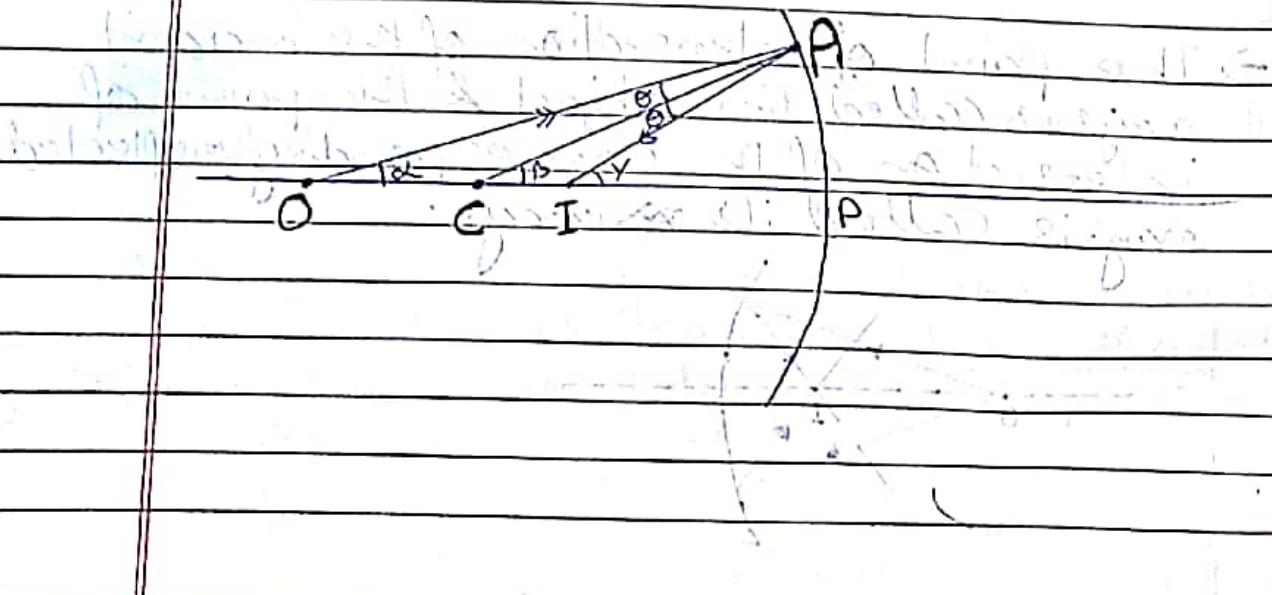
→ Principle axis is taken as X-axis.

→ The positive side of X-axis is taken along the direction of object to mirror.

→ Real inverted image will be formed.



Relation Between u, v & R for Spherical mirrors



$$\beta = \alpha + \theta \quad \left\{ \begin{array}{l} \text{Sum of exterior angle equals the sum} \\ \text{of two opposite interior angles} \end{array} \right.$$

$$\Rightarrow 2\beta = \alpha + \gamma \quad \left\{ \text{From above} \right\} \quad \text{--- (1)}$$

$$\Rightarrow \alpha \approx \frac{AP}{OP} \quad \beta = \frac{AP}{EP} \quad \gamma \approx \frac{AP}{PI} \quad \left\{ \text{From } \theta = \frac{l}{r} \right\}$$

$$\Rightarrow \frac{2AP}{CP} = \frac{AP}{PO} + \frac{AP}{PI} \quad \left\{ \text{From relation (1)} \right\}$$

$$\Rightarrow \frac{2}{R} = \frac{1}{u} + \frac{1}{v}$$

$$\Rightarrow \boxed{\frac{2}{R} = \frac{1}{u} + \frac{1}{v}}$$

* Relation between focal length and modulus of curvature

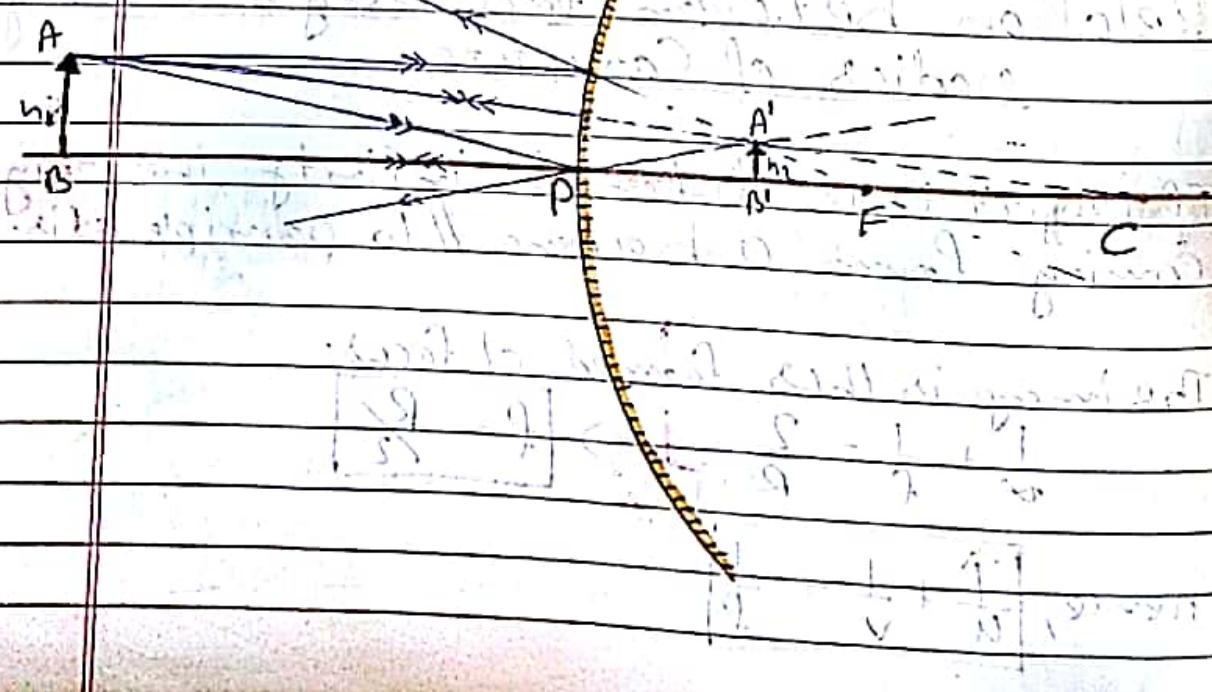
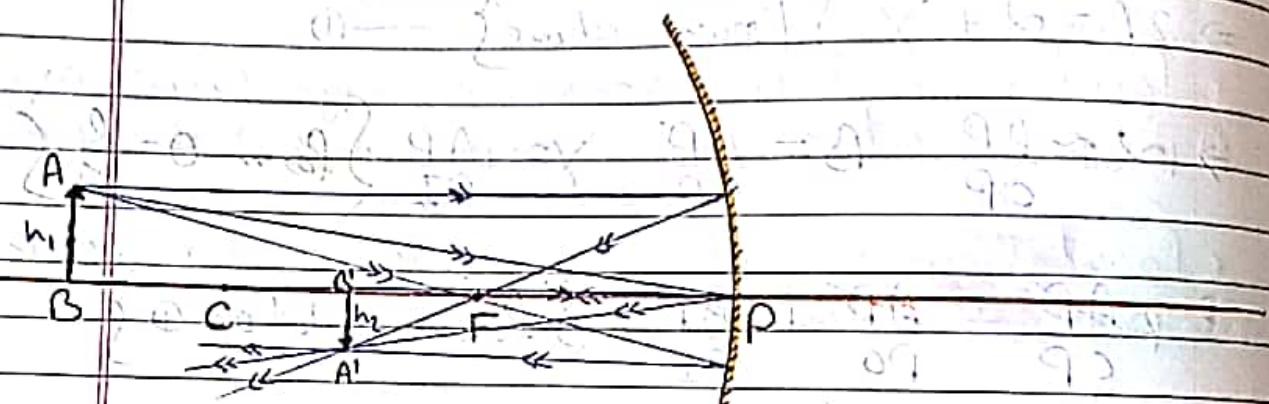
→ If object O is taken at infinity, the ray coming from O become // to principle axis.

→ The image is thus formed at focus.

$$\frac{1}{\infty} + \frac{1}{f} = \frac{2}{R} \Rightarrow f = \frac{R}{2}$$

Hence, $\boxed{\frac{1}{u} + \frac{1}{v} = \frac{1}{f}}$

Extended object & Magnification



* Lateral magnification \Rightarrow The ratio $\frac{\text{height of image}}{\text{height of object}}$

is called lateral or transverse magnification.

\Rightarrow Let h_1 & h_2 be height of object and image respectively.

$$\frac{h_1}{h_2} = -\frac{u}{v} \quad \left. \begin{array}{l} \text{as } \triangle ABP \sim \triangle A'P \\ u \quad v \end{array} \right\} \quad \text{--- (1)}$$

$$m = h_2 = -\frac{v}{u} \quad \left. \begin{array}{l} \text{from (1)} \\ h_1 \end{array} \right\}$$

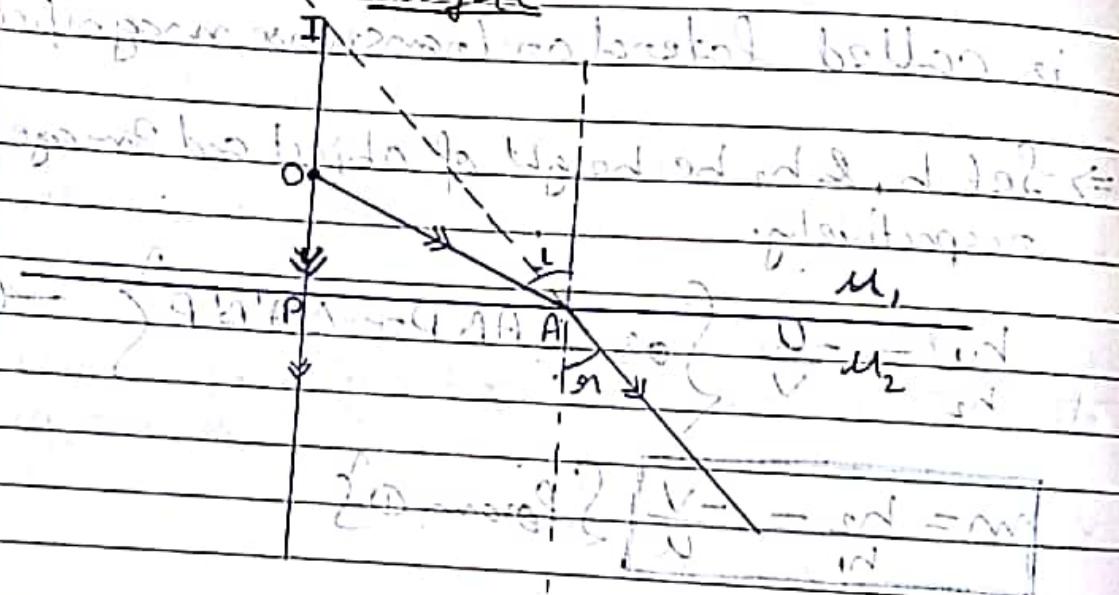
\Rightarrow Refraction at Plane Surface $\Rightarrow m$

When a light ray is incident on a surface separating two transparent media, the ray bends at the time of changing the medium. The angle of incidence is called angle of refraction & follow Snell's law.

$$\frac{\sin i}{\sin r} = \frac{v_1}{v_2} = \frac{n_2}{n_1} \quad \text{Speed}$$

Where v_1 & v_2 are speed of light in medium 1 & 2
and n_1 & n_2 are refractive index of medium 1 & 2 respectively.

* Image due to Reflection at a Plane Surface



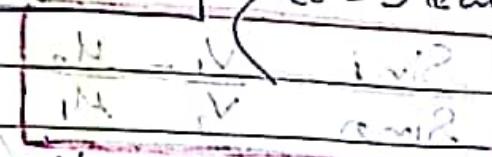
for small angle of incidence $\Rightarrow \sin i = \tan i = PA$

$$\frac{\sin i}{\sin r} = \frac{PA}{PO} = 1$$

$\therefore \sin i = \tan i = \frac{PA}{PO} = 1$ Point magnification

$$\frac{\mu_2}{\mu_1} = \frac{\sin i}{\sin r} = \frac{PI}{PO} = \frac{t}{t_0}$$

$t = \text{apparent depth}$
 $t_0 = \text{real depth}$



for $\mu_2 = 1$ & $\mu_1 = \mu$

$$\frac{1}{\mu} = \frac{t}{t_0} \Rightarrow \mu = \frac{t_0}{t}$$

$$\Delta t = t_0 - t = \left(1 - \frac{1}{n}\right) t_0 \quad \left\{ \begin{array}{l} \Delta t = \text{apparent shift} \\ \text{real shift} \end{array} \right.$$

\Rightarrow Critical angle $\{\theta_c\}$

For a particular refraction, the angle of incident at which angle of refraction become 90° is Critical angle (or RI of medium)

$$\theta_c = \sin^{-1} \left(\frac{n_2}{n_1} \right) \quad \left\{ \begin{array}{l} n_1 = \text{medium from where light is coming} \\ n_2 = \text{RI of medium where light is going} \end{array} \right.$$

For light coming from a medium to air

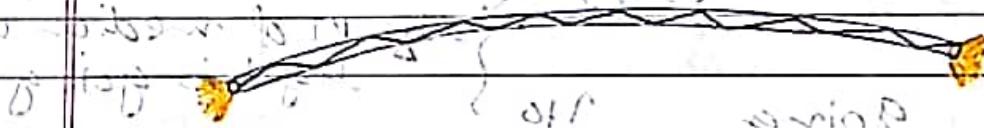
$$\theta_c = \sin^{-1} \left(\frac{1}{n} \right)$$

* Total internal reflection \Rightarrow It is a phenomenon that happens when a propagating wave strikes a medium boundary at an angle larger than Critical angle. The wave cannot pass through and is entirely reflected.

Optical Fibre

→ An optical fiber is a flexible, transparent fiber made of high quality extruded glass or plastic (slightly thicker than a human hair).

→ It can function as "light pipe" to transmit light between the two ends of the ~~the~~ fiber.



→ Because of the small radius of the fibre, light going into it makes a nearly glancing incidence on the wall.

$$i \rightarrow \theta_c$$

• Hence if total internal reflection takes place

then a signal propagates without loss of light

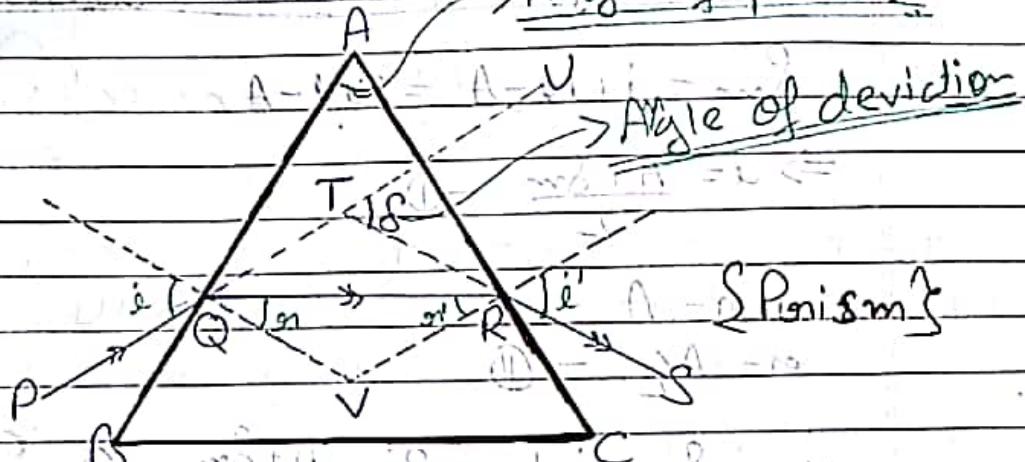
along with the polarization condition of light (horizontal and vertical) lost upon reflection.

∴ reflection may occur either horizontal or vertical

∴ reflection may occur either horizontal or vertical

Prism

Minimum deviation \Rightarrow Angle of prism



$$\delta = (i - r_1) + (i' - r'_1) = (i + i') - (r_1 + r'_1) \quad \text{--- (1)}$$

~~$$A = i - (r_1 + r'_1)$$~~

~~$$A = \pi - \frac{\pi}{2} + r_1 - \frac{\pi}{2} + r'_1 = r_1 + r'_1 \quad \text{--- (2)}$$~~

Using (1) and (2) we get

$$\Rightarrow \delta = i + i' - A$$

$$\delta = i - A + \sin^{-1} \left[\sin \left\{ A - \sin^{-1} \left(\frac{1}{n} \sin i \right) \right\} n \right]$$

$$\boxed{\sin \left\{ A - \sin^{-1} \left(\frac{1}{n} \sin i \right) \right\} n}$$

\rightarrow For angle of minimum deviation \therefore

$$\boxed{i = i'} \Rightarrow$$

$$\boxed{r_1 = r'_1}$$

* Relation between Refractive Index and the angle of Minimum deviation

$$\delta_m = i + A - A = 2i - A$$

$$\Rightarrow i = \frac{A + \delta_m}{2} \quad \text{--- (1)}$$

$$n + r = A$$

$$n = \frac{A - r}{2} \quad \text{--- (2)}$$

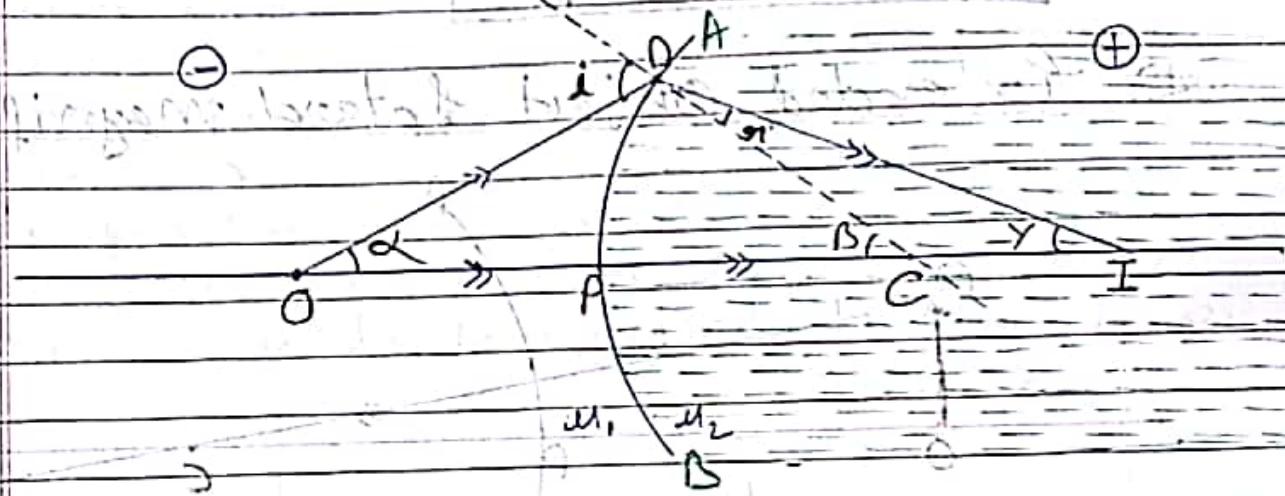
$$n = \frac{\sin i}{\sin \delta_m} = \frac{\sin \frac{A + \delta_m}{2}}{\sin \frac{A - r}{2}} \quad \left. \begin{array}{l} \text{from (1) & (2)} \\ \sin(r) = \sin(A - i) \end{array} \right\}$$

$$n = \frac{A + \delta_m}{A - r} \quad \left. \begin{array}{l} \text{for small } \delta_m \ll A \\ \text{and } A - r \approx A \end{array} \right\}$$

$$\delta_m = (n - 1) A$$

From the minimum deviation $A - \delta_m = 0$

⇒ Refraction at Spherical Surface



* Relation between u , v and R

$$i = \alpha + \beta \quad \text{--- (I)}$$

$$\beta = Y + \mu_1 \quad \text{--- (II)}$$

$$\frac{\sin i}{\sin r} = \mu_2 = \frac{Y}{\mu_1} \quad \left. \begin{array}{l} \text{for small angles} \\ \text{--- (III)} \end{array} \right.$$

Using (I), (II) and (III) we get $\frac{1}{v} - \frac{1}{u} = \frac{1}{R}$

$$\mu_1 \alpha + \mu_1 Y = (\mu_2 - \mu_1) \beta \quad \text{--- (4)}$$

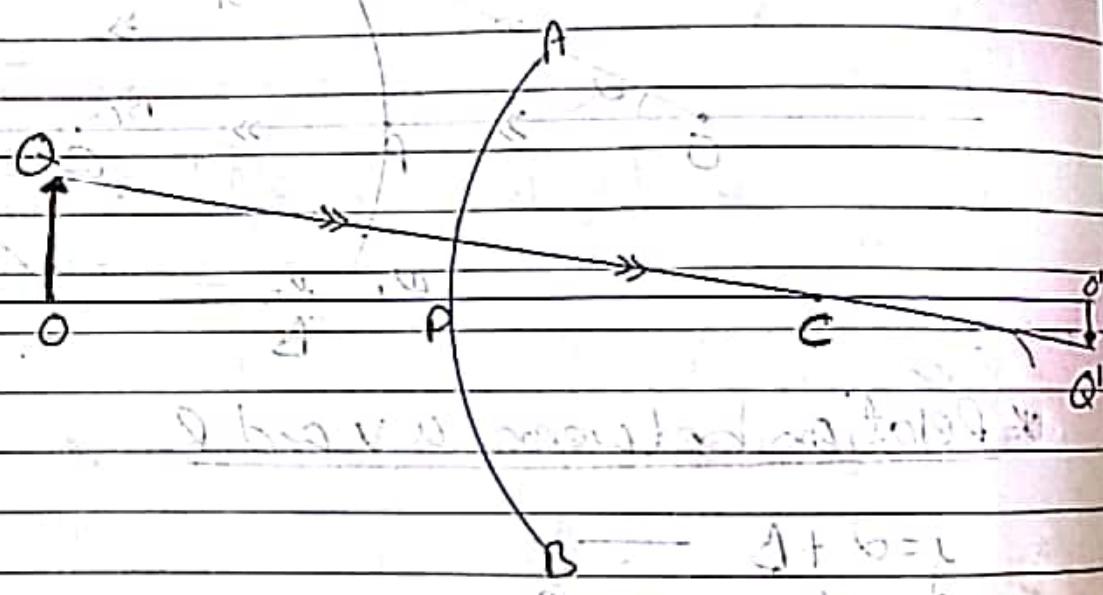
$$\alpha = \frac{DP}{PO}, \quad \beta = \frac{DP}{PC}, \quad Y = \frac{DP}{PI} \quad \text{--- (5)}$$

Using (4), (5) & (6) we get $\frac{1}{v} - \frac{1}{u} = \frac{1}{R}$

$$\Rightarrow \mu_1 \left(\frac{DP}{PO} \right) + \mu_2 \left(\frac{DP}{PI} \right) = (\mu_2 - \mu_1) \left(\frac{DP}{PC} \right)$$

$$\Rightarrow \frac{u_2 - u_1}{v} = \frac{u_2 - u_1}{R}$$

\Leftrightarrow Extended Object Lateral magnification



* Lateral magnification, $m = \frac{\text{height of Image}}{\text{height of Object}}$

$$m = -\frac{O'Q'}{OQ} = -\frac{PO' - PC}{PO + PC} = -\frac{R - R}{R + R} = -1$$

$$m = \frac{R - v}{R - u}$$

$$m = \frac{u_1 v}{u_2 u}$$

Eliminating R from
Equations

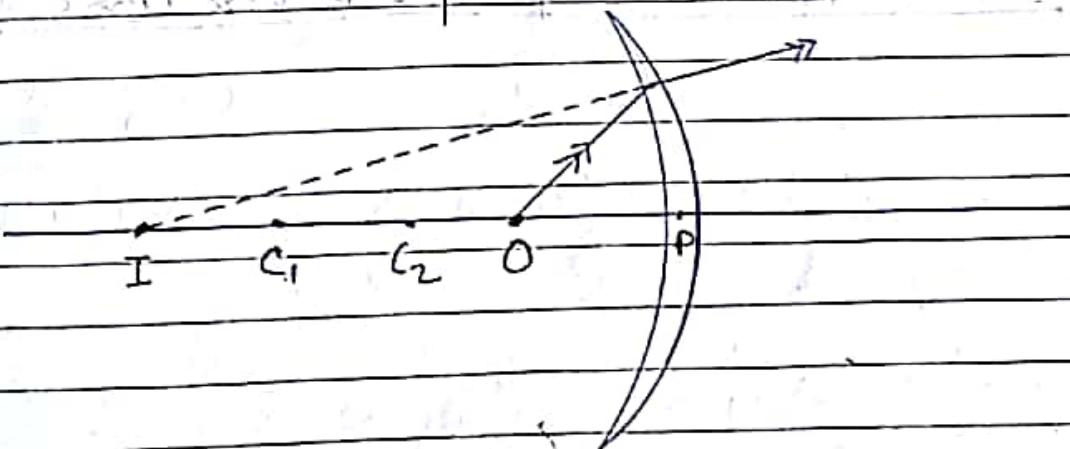
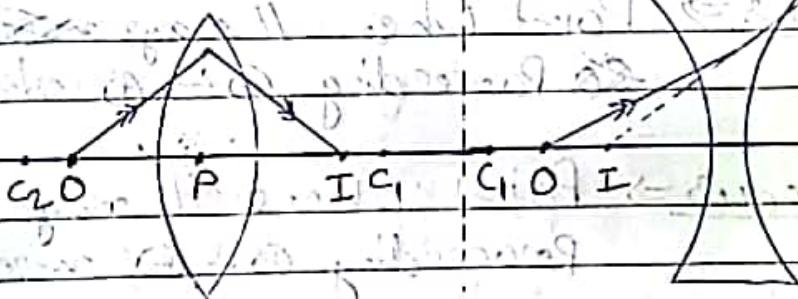
$$\frac{u_1}{v} \cdot \frac{u_1}{u} = \frac{u_2 - u_1}{R}$$

⇒ Refraction through thin lenses

using air-filled balloon

→ A lens is made of a transparent material bounded by two spherical surfaces.

→ When the thickness of the lens is small compared to the other dimensions like object distance etc. we call it a thin lens.



* Principal axis \Rightarrow line joining C_1 & C_2 is called principle axis.

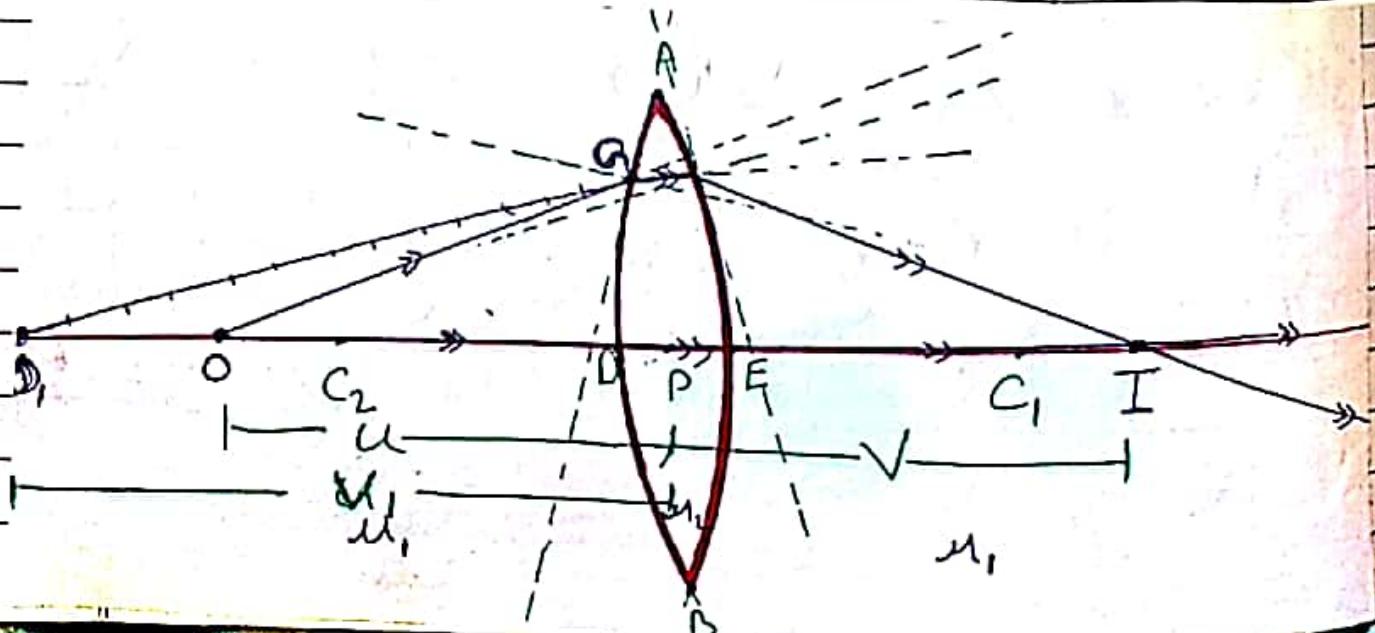
* Optical Centre \Rightarrow The centre of the thin lens which lies on the Principal axis is called the optical centre.

\rightarrow Real Focus

* First focus \Rightarrow Point where 11 rays proceeding $\ominus \rightarrow \oplus$ meets

* Second Focus \Rightarrow Point where 11 rays proceeding $\oplus \rightarrow \ominus$ meets.

\Rightarrow Lens maker's formula and lens formula



General equation for refraction at a Spherical Surface is :-

$$\frac{u_2}{v} - \frac{u_1}{u} = \frac{u_2 - u_1}{R}$$

Verified

Eulerian

\Rightarrow For refraction at surface (AOB) :-

$$\frac{u_2}{v_1} - \frac{u_1}{u} = \frac{u_2 - u_1}{R_1} \quad \text{--- (1)}$$

\Rightarrow For refraction at surface (AFB) :-

$$\frac{u_1}{v} - \frac{u_2}{v_1} = \left(\frac{u_1 - u_2}{R_2} \right) \quad \text{--- (2)}$$

\approx adding Eq. (1) & (2) we get :-

$$\frac{1}{v} - \frac{1}{u} = \left(\frac{u_2 - u_1}{R_1 + R_2} \right) \quad \text{--- (3)}$$

Considering Ideal Condition

$v = \infty$ {far away}

$u = f$ {Focus}

$u_2 = u$ {SRI of lens}

$u_1 = 1$ {ideal air}

$$\frac{1}{f} = (u-1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right) \quad \text{--- (4)}$$

* This is called Lens makers formula.

↳ Lateral magnification and magnification

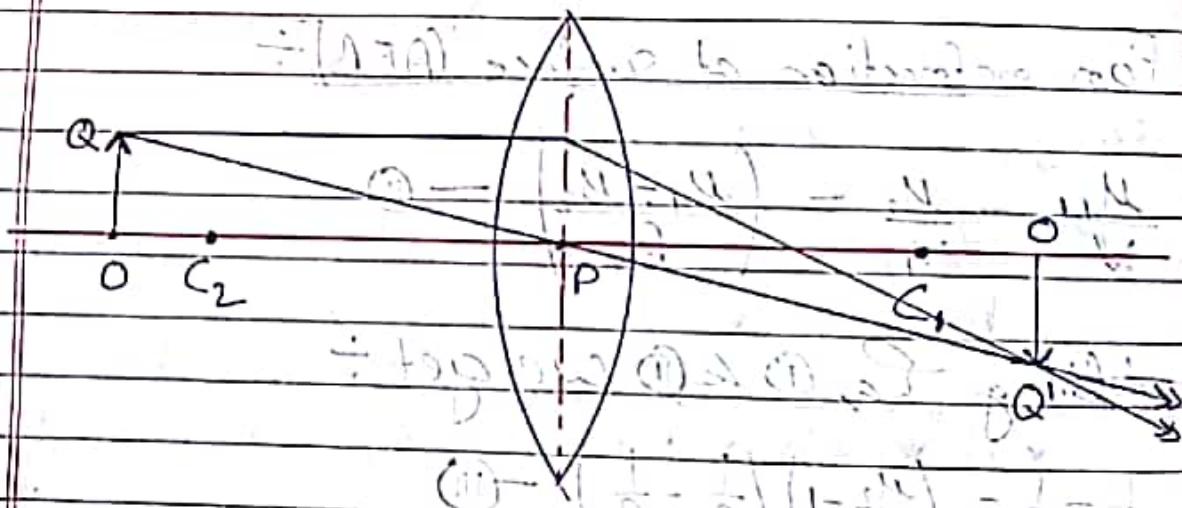
\Rightarrow Comparing Eq. ⑩ & ⑪ we get info?

$$\frac{1}{f} = \frac{1}{v} - \frac{1}{u}$$

$$\left(\frac{1}{M+1} - \frac{1}{M} = \frac{1}{v} - \frac{1}{u} \right)$$

* This is known as lens formula

\Rightarrow Extended Objects : Lateral magnification



$$m = \frac{h_2}{h_1} = \frac{-Q'P'}{QO} = \frac{-PO'}{QO} \quad \left\{ \text{as } \Delta QOP \approx \Delta Q'P' \right.$$

$$m = -\frac{v}{u} \Rightarrow m = \frac{v}{u}$$

Show that $m = \frac{v}{u} = \frac{h_2}{h_1}$

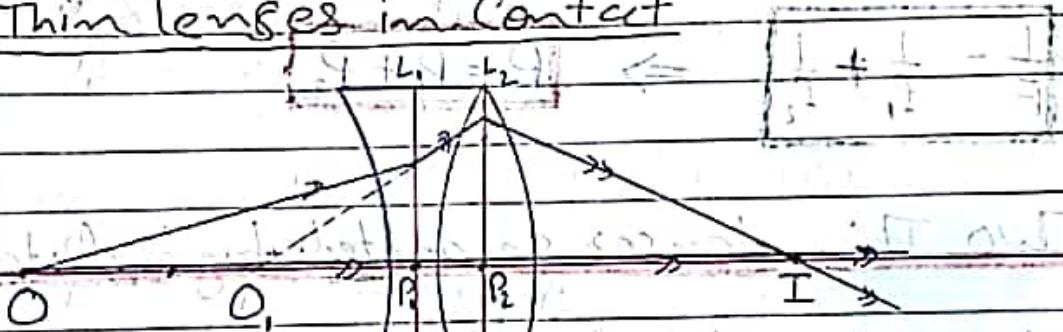
\Rightarrow Power of lens (P)

$$P = \frac{1}{f}$$

definition, where, P=focal length

\rightarrow SI unit of power of a lens is dioptre.

\Rightarrow Thin lenses in Contact



\Rightarrow Let, ~~u~~ u = object distance for the first lens

v = final image distance for the second lens.

v_1 = image distance of the first image O_1 for

$$*\frac{1}{v_1} - \frac{1}{u} = \frac{1}{f_1} \quad \{ \text{for First lens} \} \quad ①$$

$$*\frac{1}{v} - \frac{1}{v_1} = \frac{1}{f_2} \quad \{ \text{for Second lens} \} \quad ②$$

\therefore Adding ① & ② we get \therefore

$$\left[\frac{1}{v} - \frac{1}{u} = \frac{1}{f_1} + \frac{1}{f_2} \right] \quad ③$$

\Rightarrow If the combination is replaced by a single lens of focal length F such that it forms the image of O at the same position T .

$$\frac{1}{v} - \frac{1}{u} = \frac{1}{F} \quad \text{--- (ii)}$$

\Leftrightarrow Comparing (i) & (ii) we get \therefore

$$\frac{1}{F} = \frac{1}{f_1} + \frac{1}{f_2}$$

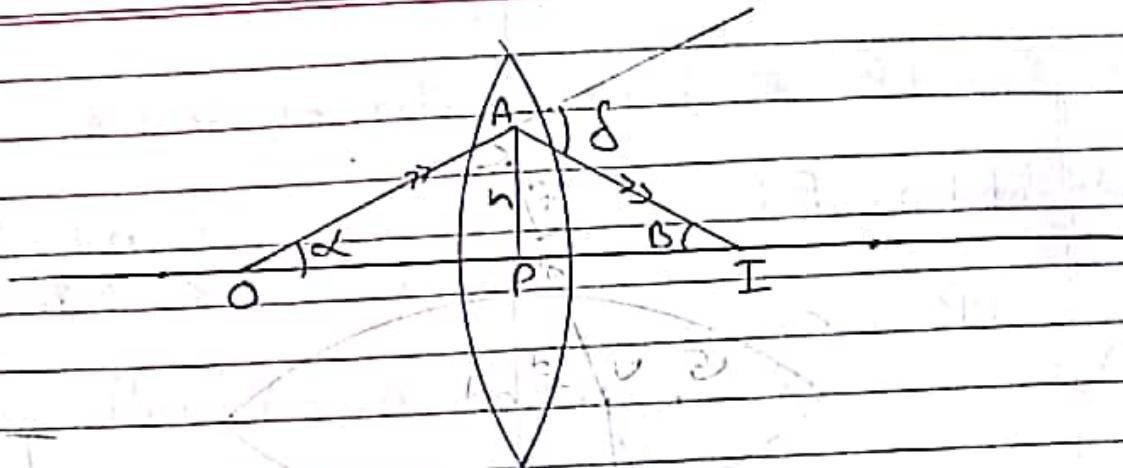
$$P = P_1 + P_2$$

\Rightarrow Two thin lenses separated by a distance

\Rightarrow When two thin lenses are separated by a distance, it is not equivalent to a single thin lens. In fact such a combination can only be equivalent to a thick lens.

\Rightarrow In a special case when the object is placed at infinity, the combination may be replaced by a single thin lens.

Expression for angle of deviation



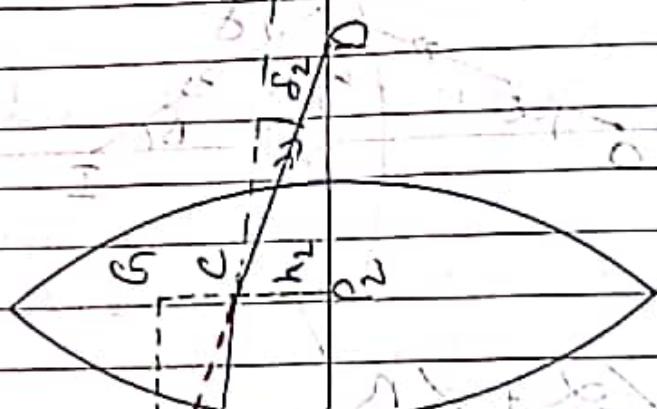
$$\Rightarrow \delta = \alpha + \beta.$$

$$\Rightarrow \delta = \tan \alpha + \tan \beta \quad \{ \text{For small } \alpha \text{ & } \beta \}$$

$$\Rightarrow \delta = \frac{h}{OP} + \frac{h}{OI} = h \left\{ \frac{1}{OP} + \frac{1}{OI} \right\}$$

$$\Rightarrow \delta = h \left\{ \frac{1}{\sqrt{h^2 + a^2}} \right\} \quad \{ \text{As } \frac{1}{\sqrt{h^2 + a^2}} = \frac{1}{P} \}$$

$$\Rightarrow \boxed{\delta = \frac{h}{P}} \quad \text{--- (1)}$$



$$|s_1 - s_2| = 2$$

$$|c_1 + c_2| = 2$$

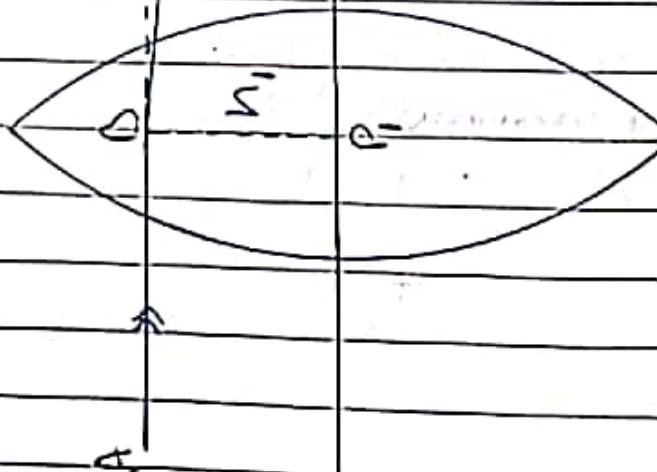
$$|c_1 - c_2| = 2$$

(To 10)

To 90

$$\{s_1 - s_2\}^2 = 8$$

$$\frac{d^2}{2} = 8$$



$$\Rightarrow S = S_1 + S_2 \quad \{ \text{from } \triangle BFC \} \quad - \textcircled{a}$$

\Rightarrow Equivalent thin lens should be placed along EP with its pole at P.

\Rightarrow The focal length of the equivalent lens, $F = PD$.

$$S_1 = \frac{h_1}{f_1} \quad \& \quad S_2 = \frac{h_2}{f_2} \quad \& \quad S = \frac{h_1}{F} \quad \{ \text{from eqn } \textcircled{a} \}$$

$$\Rightarrow S = S_1 + S_2$$

$$\Rightarrow \frac{h_1}{F} = \frac{h_1}{f_1} + \frac{h_2}{f_2} \quad - \textcircled{b}$$

$$h_1 - h_2 = GC = BC \tan \delta_1 = d \tan \delta_1 = \frac{dh_1}{f_1}$$

$$h_2 = h_1 - dh_1 \quad - \textcircled{c}$$

Using \textcircled{a} & \textcircled{b} we get:

$$\frac{1}{F} = \frac{1}{f_1} + \frac{1}{f_2} - \frac{d}{f_1 f_2}$$

Position of Equivalent lens:

$$PP_2 = EG \Rightarrow GC \cot \theta = \frac{h_1 + h_2}{f} - \frac{h_1}{f_1}$$

$$PP_2 = \frac{h_1 - h_2}{f} \quad \textcircled{a}$$

using eq \textcircled{a}, \textcircled{b} & \textcircled{c} we get :

$$PP_2 = \left(\frac{dh_1}{f_1} \right) \left(\frac{F}{h_1} \right)$$

$$\boxed{PP_2 = \frac{dF}{f_1}}$$

$$\frac{1}{f} = \frac{1}{f_1} + \frac{1}{f_2} = \frac{1}{f_1} - \frac{1}{n_2} = \frac{n_2 - n_1}{n_1 n_2}$$

* Chromatic aberration \Rightarrow It is a type of distortion in which there is a failure of a lens to focus all colours to the same convergence point.

⇒ Defects of Image

→ The simple theory of image formation developed for mirrors and lenses suffers from various approximations.

* As a result, the actual images formed contain several defects.

Defects of Image

Monochromatic Aberration

Chromatic Aberration

Spherical Aberration

{ defect due to which 1/R rays may not meet at a point }

Coma

{ defect due to which 1/R rays do not meet at a point }

Astigmatism

{ defect due to which Point object image spread on principle axis }

Curvature

{ defect due to which Perfect image is obtained on curved surface not plane }

Distortion

{ defect due to which magnification of extended Object are different at different position }

2

Wave Optics

CLASSMATE

Date _____

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⇒ Huygen's Principle

⇒ Huygen's considered light to be a mechanical wave moving in a hypothetical medium which was named as ether.

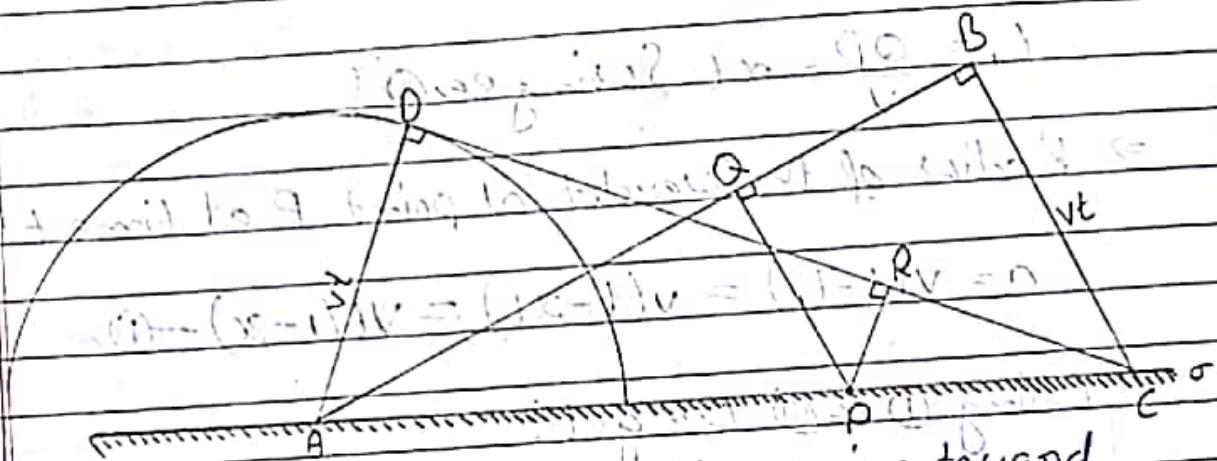
⇒ Huygen's principle may be stated in its most general form as follows :-

"Various points of an arbitrary surface, when reached by a wavefront, become secondary sources of light emitting secondary wavelets. The disturbance beyond the surface results from the superposition of these secondary wavelets."

* Huygen's Construction

"Various points of an arbitrary surface, as they are reached by a wavefront, becomes the sources of secondary wavelets. The geometrical envelope of these wavelets at any given later instant represents the new position of the wavefront at that instant"

* Reflection of Light



\Rightarrow Let us consider a wavelet AB moving toward reflecting surface σ . Let at time t point D come in contact with σ at point C.

∴ Using the formula of travel distance

\Rightarrow Let P be a point of σ between AC and C.

$$\frac{PR}{AD} = \frac{PC}{AC} \quad \text{as } \angle ACD = \angle PCR \quad \left\{ \begin{array}{l} \text{Let } x = \frac{AP}{AC} \\ \text{D is a virtual point on } \sigma \end{array} \right\}$$

$$\Rightarrow PR = AD(1-x) = vt(1-x) \quad \text{--- (1) At v = vt}$$

$$\text{also } \Rightarrow \frac{QP}{BC} = \frac{AP}{AC} = x \quad \left\{ \text{as } \triangle ADQ \sim \triangle ABC \right\}$$

$$\Rightarrow QP = BCx = vt \quad \text{--- (2) At v = vt}$$

$$\therefore QP = PR \quad \text{--- (3)}$$

\Rightarrow Let t_1 be the time taken by wavefront to reach point P.

$$t_1 = \frac{QP}{v} = xt \quad (\text{using eqn } ⑪)$$

\Rightarrow Radius of the wavelet at point P at time t'

$$a = v(t - t_1) = v(t - xt) = vt(1 - x) \quad ⑬$$

\Rightarrow Using ⑪ & ⑬ we get:

↳ ~~length of wavelet PR is also a function of time~~

$\Rightarrow a = PR$ ~~length of wavelet is constant~~

↳ ~~length of wavelet is constant~~ ~~length of wavelet is constant~~

* CD is tangent to ~~wavelet~~ Wavelet at point at time t'

* As P is an arbitrary point, hence ~~any~~ ~~any~~ ~~any~~ is tangent to every wavelet formed from point within AC.

* Hence CD is new wave front at time t'

\Rightarrow Hence incident ray is perpendicular to AC & reflected ray is 1 to CD.

$$\textcircled{2} \quad i = \angle BAC \quad & \quad r = \angle ACD$$

$$\Rightarrow \angle BAC = \angle ACD \quad \{ \text{as } \triangle ACD \cong \triangle ABC \}$$

$\Rightarrow i = r$
As Incident ray, Reflected ray & Normal lies
in same plane. {By Geometry}

Hence Law of Reflection verified.

$$\frac{\sin i}{\sin r} = \frac{\sin 90^\circ - \theta}{\sin (\theta - 90^\circ)} = \frac{\sin 90^\circ}{\sin (\theta - 90^\circ)} = \frac{1}{\tan (\theta - 90^\circ)}$$

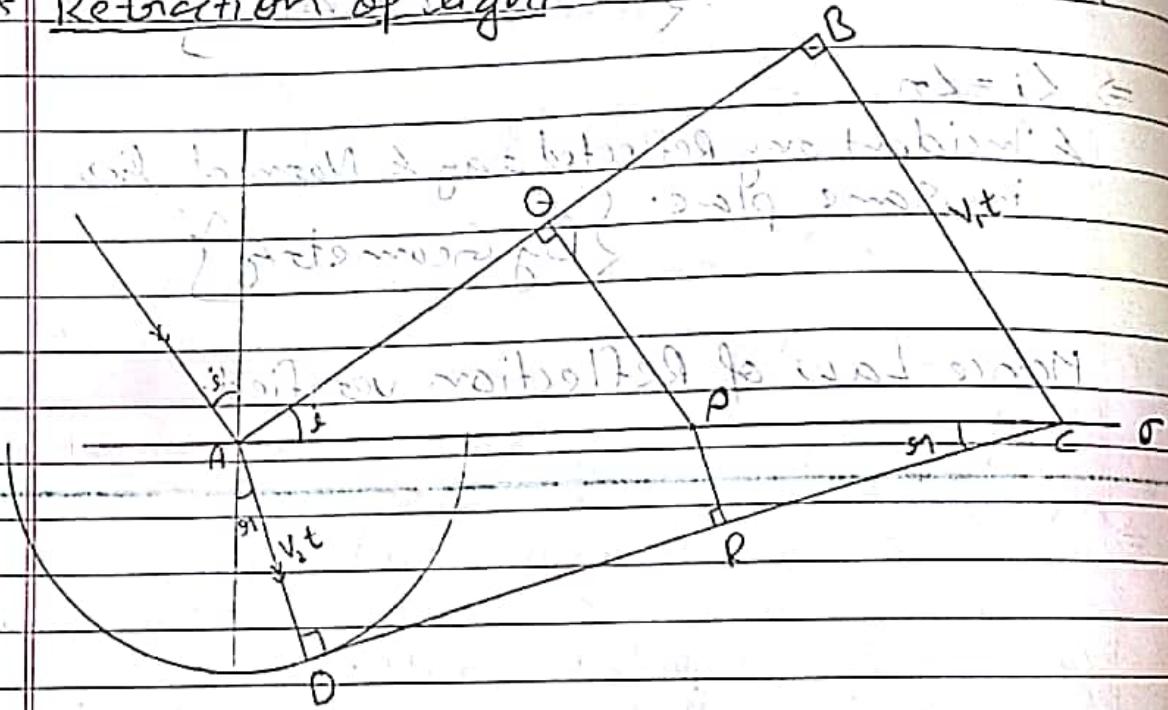
$$\Rightarrow \frac{\sin i}{\sin r} = \frac{1}{\tan (\theta - 90^\circ)} = \frac{1}{\cot (\theta - 90^\circ)} = \frac{1}{\operatorname{cosec} (\theta - 90^\circ)}$$

$$\Rightarrow \frac{\sin i}{\sin r} = \frac{1}{\operatorname{cosec} (\theta - 90^\circ)} = \sin (\theta - 90^\circ) = \sin (\theta - 90^\circ)$$

Now if $i = 90^\circ$ and $r = 90^\circ$ then $\theta = 90^\circ$ and $\theta - 90^\circ = 0^\circ$
 $\therefore \operatorname{cosec} 0^\circ = \infty$ and $\sin 0^\circ = 0$

$$\therefore \frac{\sin i}{\sin r} = \frac{1}{\operatorname{cosec} (\theta - 90^\circ)} = \frac{1}{\infty} = 0$$

* Refraction of light



$$\Rightarrow \frac{PR}{AD} = \frac{PC}{AC} = \frac{AC - AP}{AC} = 1 - \frac{AP}{AC} = 1 - x \quad \left\{ \text{Let } x = \frac{AP}{AC} \right\}$$

$$\Rightarrow PR = (AP)(1-x) = v_1 t (1-x) - \textcircled{1}$$

$$\text{Also } \Rightarrow \frac{PQ}{BC} = x \Rightarrow PQ = (BC)x = v_2 t x - \textcircled{2}$$

Let t_1 be the time taken by the wave front to reach point P.

$$t_1 = \frac{PQ}{v_1} = x t \quad \left\{ \text{Using } \textcircled{1} \right\}$$

// Let a be the radius of wavelet originated at point P.

$$a = v_1(t-t_i) = v_1 t (1-x) \quad \text{--- (III)}$$

// Using eq. (I) & (II) we get -

$$\Rightarrow a = PR$$

Medium CD is the new wave-front at time t .

$$\sin i = \frac{v_1 t}{AC} \quad \& \quad \sin r = \frac{v_2 t}{AC}$$

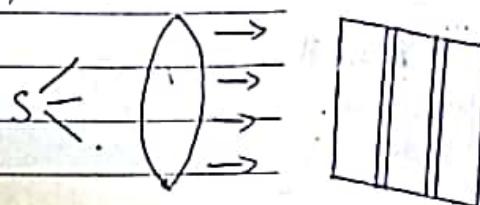
$$\Rightarrow \frac{\sin i}{\sin r} = \frac{v_1}{v_2} \quad \left. \begin{array}{l} \text{Hence Snell's Law is Law} \\ \text{of refraction is verified} \end{array} \right\}$$

\Rightarrow The ratio v_1/v_2 is called refractive index of medium 2 with respect to medium 1 and is denoted by μ_{21} .

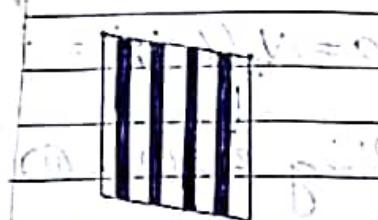
\Rightarrow If the medium 1 is vacuum, μ_{21} is simply the refractive index of the medium 2 and is denoted by μ_2 .

$$\Rightarrow \mu_{21} = \frac{v_1}{v_2} = \frac{c/v_1}{c/v_2} = \frac{v_2}{v_1} = \frac{\mu_2}{\mu_1}$$

Young's double slit Experiment



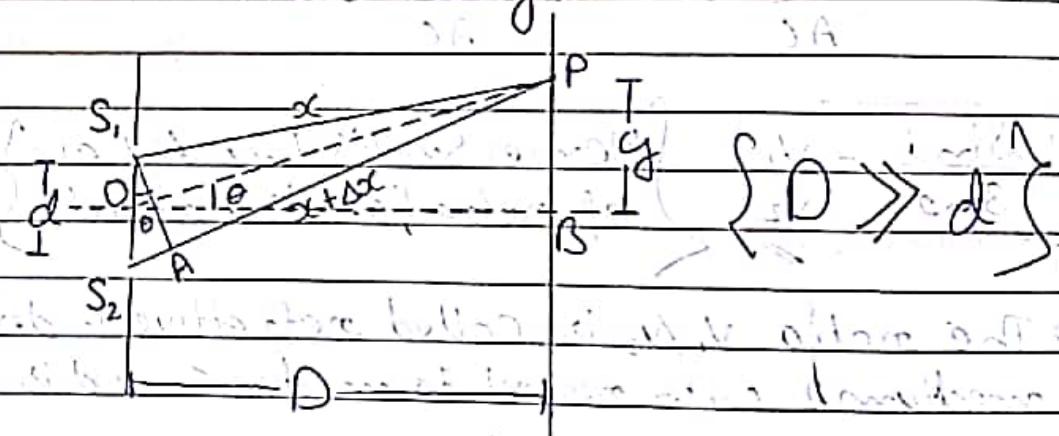
Source



Slits

Fringes on Screen

- * Fringes \Rightarrow Series of dark and bright strips obtained on screen is called fringes.



Let us consider two wave one from S_1 & another from S_2 . Interfinge at a general point P .

$$S_1P = x ; S_2P = x + \Delta x$$

* $E_1 = E_0 \sin(kx - \omega t)$ Electric Field at point P due to S_1

$$E_2 = E_{02} \sin [K(x + \Delta x) - \omega t]$$

$$\Rightarrow E_{02} \sin [Kx - \omega t + \delta] \quad \left. \begin{array}{l} \text{Electric field at point} \\ \text{P due to S}_1 \end{array} \right\}$$

$$\left. \begin{array}{l} S = K\Delta x = \frac{2\pi}{\lambda} \Delta x \end{array} \right\}$$

$$E = E_0 \sin (Kx - \omega t + \epsilon) \quad \left. \begin{array}{l} \text{Resultant electric} \\ \text{field at point P} \end{array} \right\}$$

$$\text{Where, } E_0^2 = E_{01}^2 + E_{02}^2 + 2E_{01}E_{02} \cos \delta$$

$$\tan \epsilon = \frac{E_{02} \sin \delta}{E_{01} + E_{02} \cos \delta}$$

The condition for constructive & destructive interference:

$$(i) \delta = 2\pi n \quad \text{for bright fringes}$$

$$\Delta x = n\lambda$$

$\forall n \in \text{Integers}$

$$(ii) \delta = (2n+1)\pi \quad \text{for dark fringes}$$

$$\Delta x = (n + \frac{1}{2})\lambda$$

* Intensity Variation

If the two slits are identical $E_{o1} = E_{o2} = E_0$

$$E_0^2 = 2E_0^2(1 + \cos\delta)$$

$$\Rightarrow I = 4I'(1 + \cos\delta) = 4I' \cos^2\delta$$

Note

→ Intensity is proportional to Square of Amplitude.

→ $I = \text{Resultant Intensity}$

→ $I' = \text{Intensity due to Single Slit}$

* $I = 4I'$ for Constructive Interference.

* $I = 0$ for ~~dark~~ destructive interference.

* Fringe-width and determination of wavelength

* Fringe-width \Rightarrow The separation on the screen between the centres of the two consecutive bright fringes or two consecutive dark fringes is called the fringe-width.

$\Rightarrow S.P \& S_2 P$ are nearly \parallel and hence $S.A$ is very nearly \perp to $S.P. S_2 P \& O.P.H$ i.e. $\angle = 90^\circ$

{as $\theta \gg d$ }

$$\Rightarrow \angle S_2 S.A = \angle P_0 B = \theta$$

$$\Rightarrow \Delta x = PS_2 - PS_1 \approx PS_2 - PA = S_2 A \quad \text{as } \theta \ll 1 \Rightarrow \theta = \frac{\Delta x}{D}$$

$$\Rightarrow d \sin \theta \approx d \Delta x \quad \text{[for small } \theta\text{]}$$

$$\Delta x = \frac{dy}{D}$$

// for bright fringes

$$\Delta x = \frac{dy}{D} = m\lambda$$

Hence fringe-width is $\frac{\Delta x}{d}$

$$\Rightarrow y = m D \lambda$$

$$w = \frac{\Delta x}{d}$$

// for dark fringes

$$\Delta x = \left(m + \frac{1}{2}\right) \lambda$$

$$\Rightarrow y = \left(m + \frac{1}{2}\right) \frac{D \lambda}{d}$$

Optical Path

Let us consider a light wave travelling in a medium of refractive index n . Then if it travels a distance Δx it can be said that

its equation may be written as :-

$$E = E_0 \sin \omega(t - \frac{\Delta x}{n})$$

$$E = E_0 \sin \omega(t - \frac{n\Delta x}{c})$$

In the light wave travels a distance Δx , the phase changes by :-

$$\delta_1 = n \frac{c}{c} \Delta x - \text{initial phase} = n \Delta x - \Delta x = (n-1) \Delta x$$

If the light travels through a distance Δx in vacuum then phase changes by :-

$$\delta_2 = \frac{c}{c} \Delta x - \text{initial phase}$$

By (i) & (ii), we see that a wave travelling through a distance Δx in a medium of refractive index n suffers the same phase change as when it travels a distance $n\Delta x$ in ~~Vacuum~~ Vacuum.

⇒ The quantity nsx is called the optical path of the light.

→ The geometrical path and the optical path are equal only when light travels in vacuum or in air where the refractive index is close to 1.

⇒ The concept of optical path may also be introduced in terms of change in wavelength as the wave changes its medium.

⇒ If the wavelength of light in vacuum is λ_0 and that in the medium is λ_m then :-

$$\lambda_0 = \frac{c}{f}$$

$$\lambda_m = \frac{v}{f} = \frac{c}{nf}$$

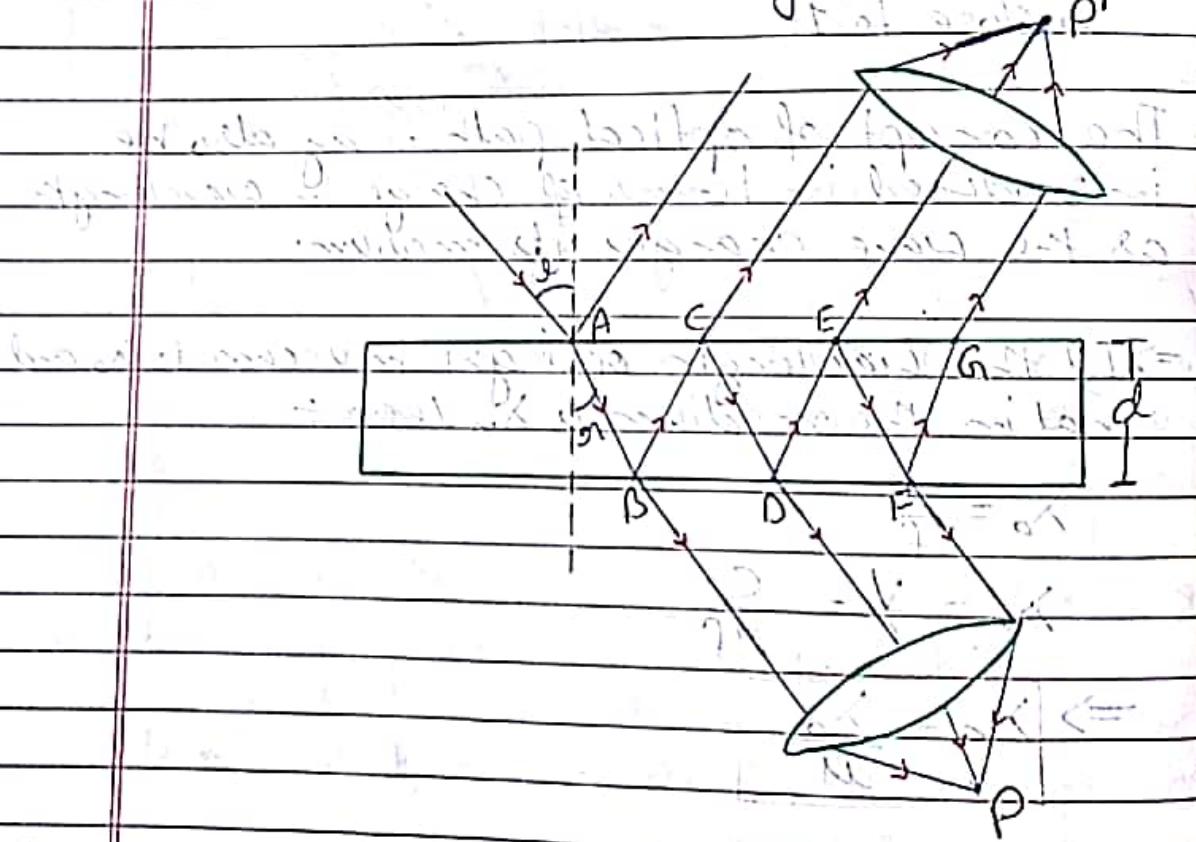
$$\Rightarrow \lambda_m = \frac{\lambda_0}{n}$$

Hence the Path λ_m in a medium is equivalent to path $\lambda_0 = n\lambda_m$ in vacuum.

\Rightarrow Interference from thin films

Consider a thin film made of a transparent material with plane parallel faces separated by a distance d .

\Rightarrow Suppose a $\perp\!\!\!/\!$ beam of light is incident on the film at an angle θ as shown.



⇒ The amplitude of the individual transmitted wave is different for different waves, it gradually decreases as more reflection's are involved.

Let us consider the phase difference between BP & DP through a special case when $j=0$.

⇒ Geometrical path difference = $2d$

⇒ Optical path difference = $2nd$

$$\Rightarrow \Delta x = 2nd \Rightarrow \delta = \frac{2\pi}{\lambda} \Delta x = \frac{2\pi}{\lambda} \times 2nd$$

⇒ All these waves are in phase if

$$\delta = 2n\pi$$

$$\Rightarrow \frac{2\pi}{\lambda} \times 2nd = 2n\pi \quad \left. \begin{array}{l} \text{Condition for maximum} \\ \text{Illumination in refraction} \end{array} \right\}$$

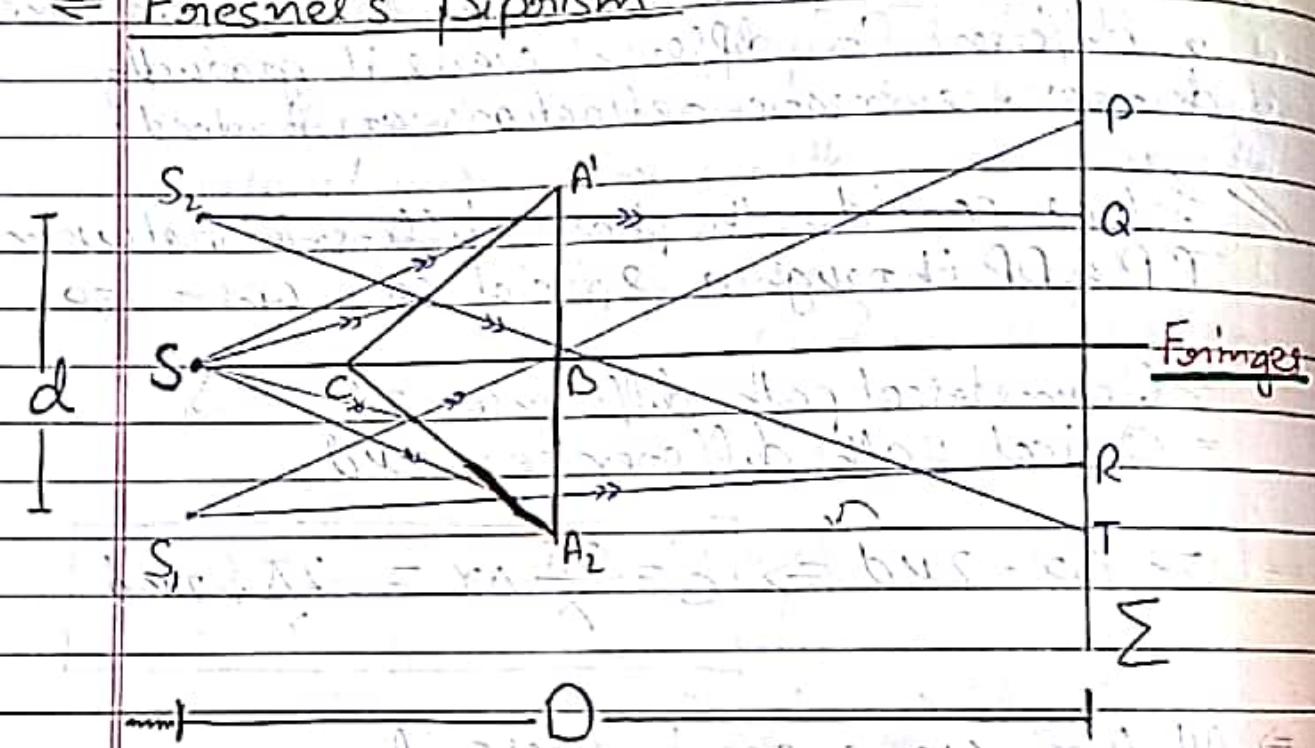
$$2nd = n\lambda$$

⇒ All these waves are out of phase if

$$\delta = (2n+1)\pi$$

$$2nd = \left(n + \frac{1}{2}\right)\lambda \quad \left. \begin{array}{l} \text{Condition for minimum} \\ \text{Illumination in refraction} \end{array} \right\}$$

⇒ Fresnel's Bipism



⇒ One can treat the points S_1 & S_2 as two coherent sources sending light to the screen.

⇒ The arrangement is then equivalent to a Young's double slit experiment with S_1 & S_2 acting as the two slits.

⇒ The fringe-width obtained on the screen is $\frac{\lambda}{d}$

$$\text{cw} = \frac{\lambda}{d}$$

⇒ Coherent & Incoherent sources

* Coherent Sources ⇒ Two sources of light waves are said to be coherent if initial phase difference $\Delta\phi$ between the waves remains constant in time.

Fringes

* Incoherent Sources ⇒ Two sources of light waves are said to be incoherent if initial phase difference $\Delta\phi$ between the waves changes randomly with time.

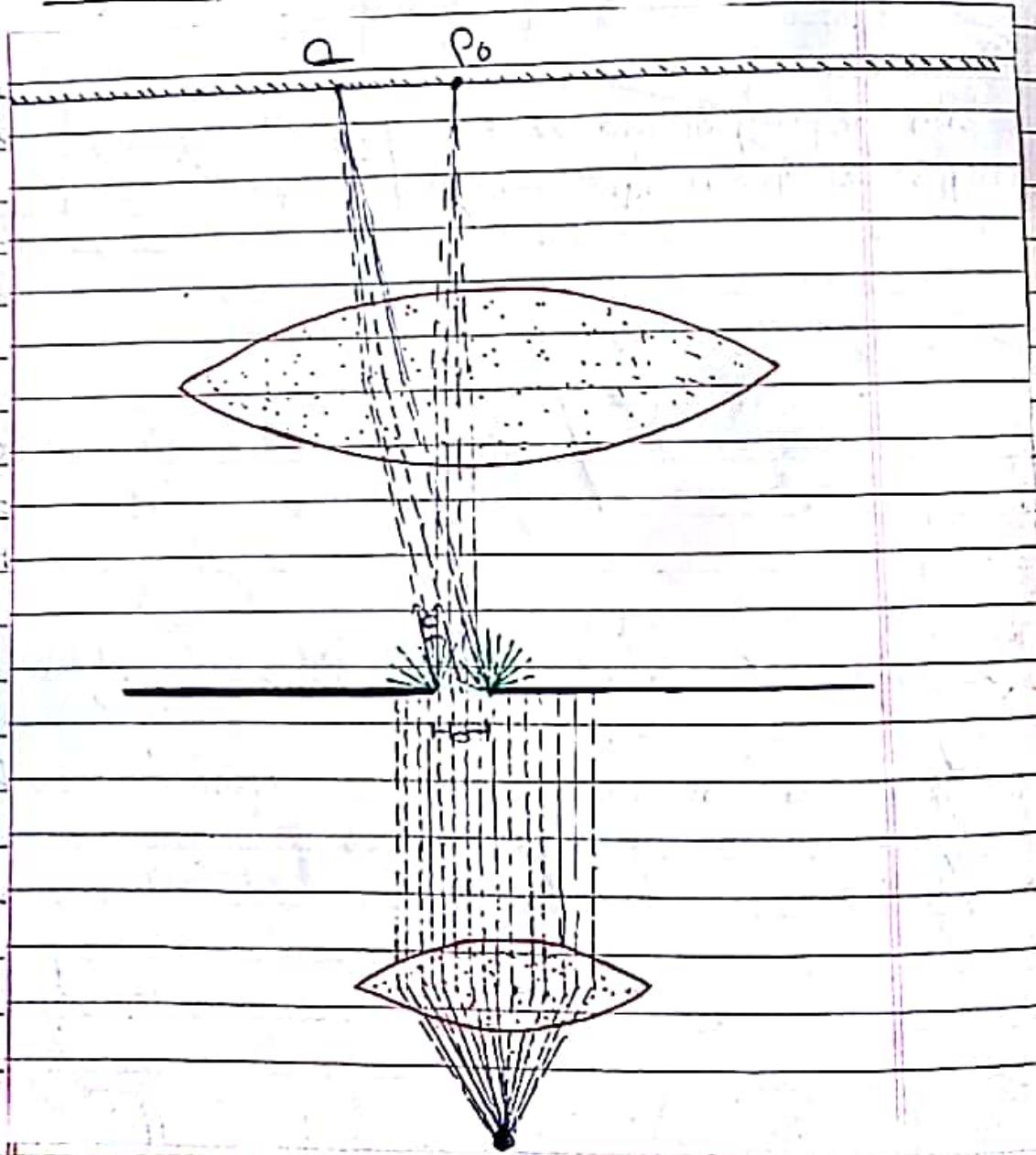
⇒ Diffraction of Light

* Diffraction ⇒ When a wave is obstructed by an obstacle, the rays bend round the corner. This phenomenon is known as diffraction.

→ Fraunhofer diffraction ⇒ It deals with the limiting cases where the light approaching the diffracting object is parallel and monochromatic, and where the image plane is at a distance large compared to the size of the diffracting object.

* Fresnel diffraction \Rightarrow Erasmeel diffraction
is a general form of Fraunhofer diffraction
in which all the restrictions are relaxed.

\Rightarrow Fraunhofer diffraction by a single Slit



\Rightarrow Suppose a parallel beam of light is incident normally on a slit of width b .

\Rightarrow The light is received by a screen placed at a large distance. In practice, this condition is achieved by placing the screen at the focal plane of a converging lens placed just after the slit.

\Rightarrow Let us consider a point P which collects the waves originating from different points of the slit at an angle θ .

\Rightarrow The optical path difference between the waves sent by the upper edge A of the slit and the wave sent by the centre of the slit is $\frac{b}{2} \sin \theta$.

\Rightarrow Consider the angle for which $\frac{b}{2} \sin \theta = \frac{\lambda}{2}$

$$\Rightarrow \delta = \frac{2\pi}{\lambda} \times \frac{\lambda}{2} = \pi$$

\Rightarrow The two wave will cancel each other.

\Rightarrow The whole slit can be divided into such pairs and hence, the intensity at P will be zero.

\Rightarrow Condition for minima, $b \sin \theta = n \lambda$ {dark fringe}

Condition for maxima,

$$b \sin \theta = (n + \frac{1}{2}) \lambda \quad \{n \rightarrow 1, 2, \dots\}$$

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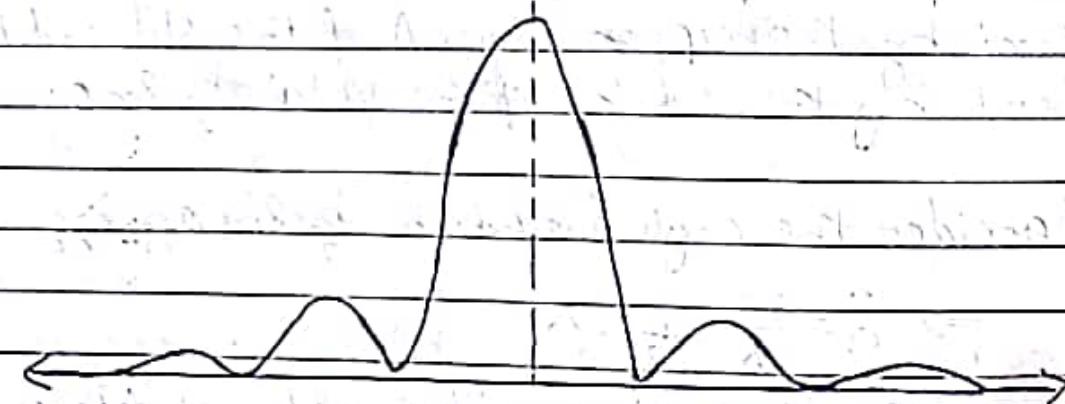
⇒ A detailed mathematical analysis shows that the amplitude E'_0 of the electric field at a general point P is :-

$$E'_0 = E_0 \frac{\sin \beta}{\beta}$$

$$\beta = \frac{1 - \frac{c}{v}}{2} b \sin \theta = \frac{\lambda}{\lambda} b \sin \theta$$

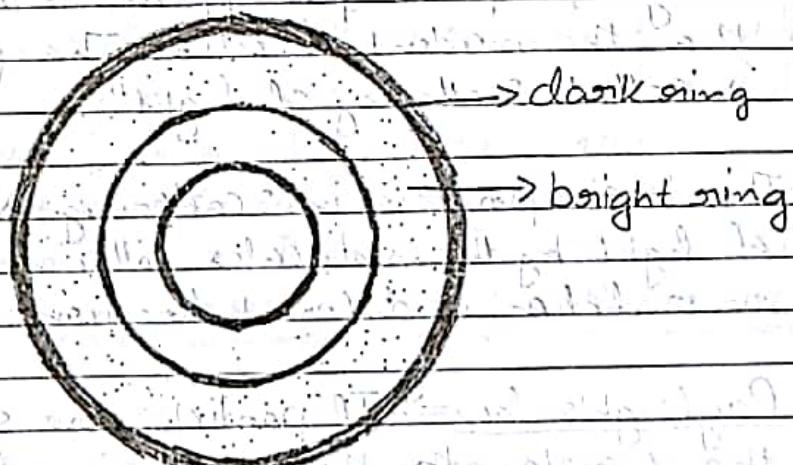
$$I = I_0 \frac{\sin^2 \beta}{\beta^2}$$

} Intensity is proportional to
the square of the amplitude

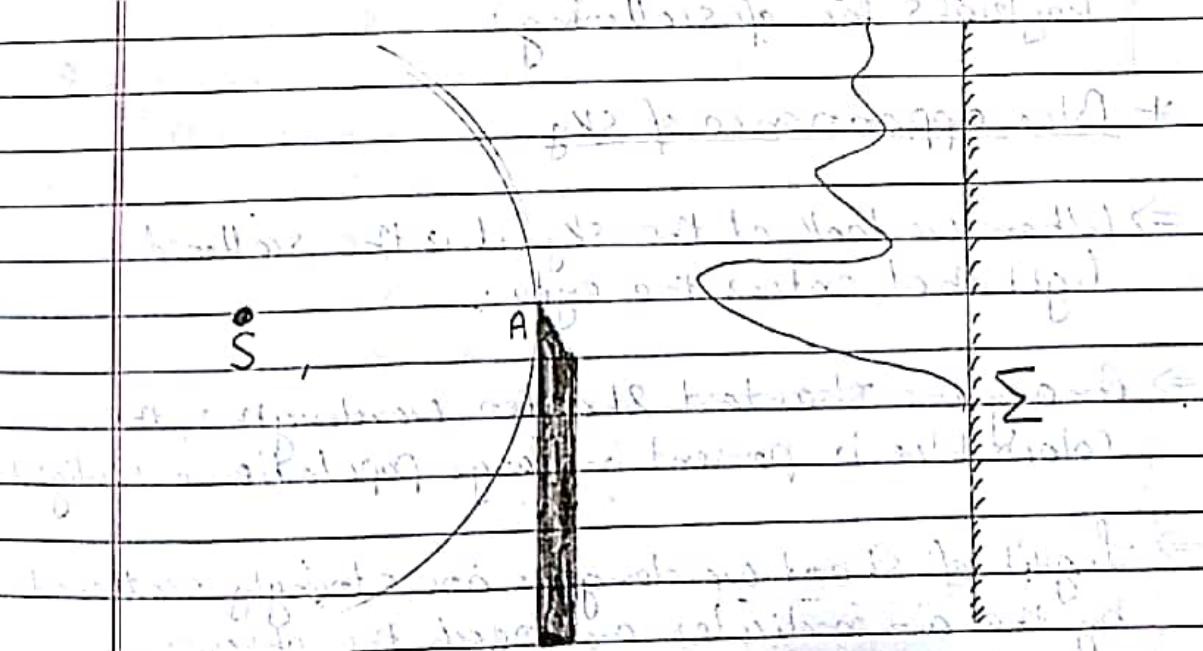


Rough Intensity Variation

⇒ Fraunhofer diffraction by a circular aperture



⇒ Fresnel diffraction at a straight edge



⇒ Scattering of light

"When a parallel beam of light passes through a gas, a part of it appears in directions other than the incident direction. This phenomenon is called scattering of light"

⇒ The basic process in scattering is absorption of light by the molecules followed by its re-emission in different direction.

* Rayleigh's law ⇒ If particle are smaller than the wavelength, the scattering of light is proportional to $1/\lambda^4$. ~~This is known as~~ This is known as Rayleigh's law of scattering.

* Blue appearance of sky

⇒ When we look at the sky, it is the scattered light that enters the eyes.

⇒ Among the shortest shorter wavelengths, the colour blue is present in large proportion in sunlight.

⇒ Light of short wavelengths are strongly scattered by the air molecules and reach the observer.

⇒ This explains blue colour of sky and violet colour of rainbow.

1. * Red appearance of sun at the sunset and at the sunrise.

⇒ At these times, the sunlight has to travel a large distance through the atmosphere. The blue and neighbouring colours are scattered away in the path and the light reaching the observer is predominantly red.

D

$$(21 \times 4 - 10) \times 12.7 = 3$$

$$(16 \times 4 - 10) \times 12.7 = 23$$

Polarization of light

* Plane polarized light \Rightarrow If electric field at a point always remains parallel to a fixed direction as time passes, the light is called linearly polarized along that direction. Also called plane polarized light.

* Unpolarized light \Rightarrow If electric field at a point changes direction randomly, then the light is called unpolarized light.

\Rightarrow Suppose an unpolarized light wave travels along the x-axis. The electric field at any instant is in the Y-Z plane, ~~so~~ we can break the field into its components E_y & E_z .

\Rightarrow The fact that the resultant electric field changes its direction randomly may be mathematically expressed by saying that E_y and E_z have a phase difference δ that changes randomly with time.

$$E_y = E_1 \sin(\omega t - k_x x + \delta)$$

$$E_z = E_2 \sin(\omega t - k_x x)$$

\Rightarrow If $\delta = \pi/2$ and $E_1 = E_2$ then,

$$\tan\theta = \frac{E_x}{E_y} = \frac{E_x \sin(\omega t - kx)}{E_y \sin(\omega t - kx + \pi/2)} = \tan(\omega t - kx)$$

$$\Rightarrow \theta = \omega t - kx$$

$$\text{And } E^2 = E_{xy}^2 + E_x^2 = E_1^2 = E_2^2$$

\Rightarrow The tip of the electric field thus, goes in a circle at a uniform angular speed. Such light is called Circularly polarized light.

\Rightarrow If $\delta = \pi/2$ but $E_1 \neq E_2$, the tip of electric field traces out an ellipse. Such a light wave is called Elliptically polarized light.

★ Polaroids

\Rightarrow Plane sheets in the shape of circular discs called Polaroids are commercially available which transmit light with E-vector parallel to a special direction in the sheet.

\Rightarrow These polaroids have long chain of hydrocarbon which become conducting at optical frequencies.

⇒ When light falls perpendicularly on the sheet, the electric field parallel to the chains is absorbed in setting up electric currents in the chains but the field perpendicular to the chain gets transmitted.

⇒ The direction perpendicular to the chains is called the transmission axis of the polaroid.

★ Polarization by Reflection and Refraction:

⇒ If the light is incident on the surface with an angle of incidence i such that:

$$\tan i = n$$

⇒ The reflected light is completely polarized with the electric field along the z-direction.

⇒ The refracted ray is never completely polarized.

⇒ The angle i is called Brewster's angle and equation itself is known as the Brewster's Law.

★ Polarization by Scattering

⇒ When unpolarized light is scattered by small particles, the scattered light is partially polarized.

"Austrian Nobel Laureate Karl Von Frisch performed experiments for several years on bees and concluded that the bees can not only distinguish unpolarized light from polarized light but can also determine the direction of polarization"

★ Malus law

⇒ If the E-Vector is at an angle θ with the transmission axis, light is partially transmitted. The intensity of the transmitted light is:

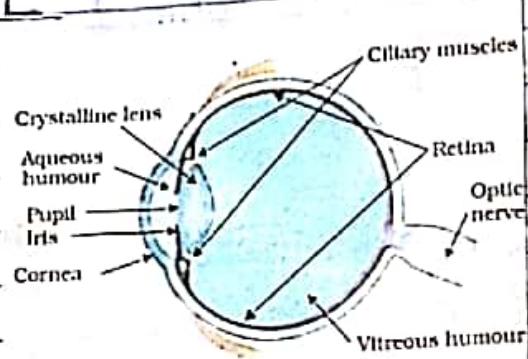
$$I = I_0 \cos^2 \theta$$

3

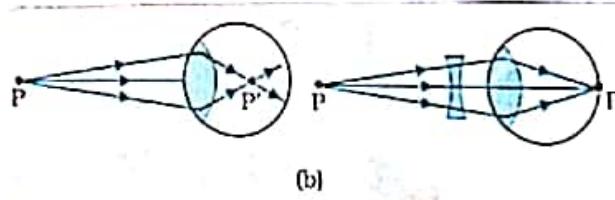
Optical Instruments



EYE

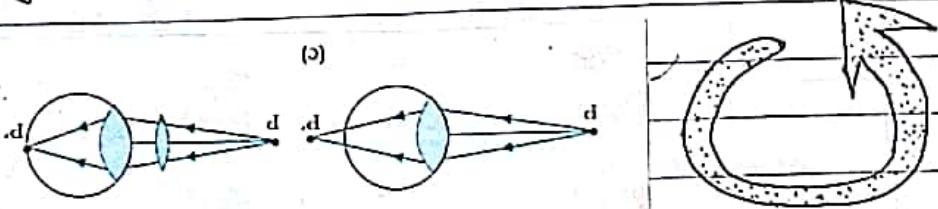


- * Accommodation \Rightarrow Accommodation is the process by which the vertebrate eye changes optical power to maintain a clear image.
- * Presbyopia \Rightarrow Presbyopia is a condition where with age, the eye exhibits a progressively diminished ability to focus on near object.
- * Myopia (nearsightedness) \Rightarrow Myopia is a condition of the eye where the light that comes in does not directly focus on the retina, but in front of it.

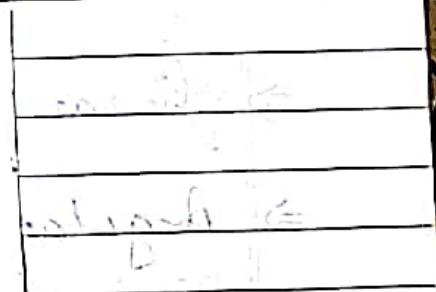
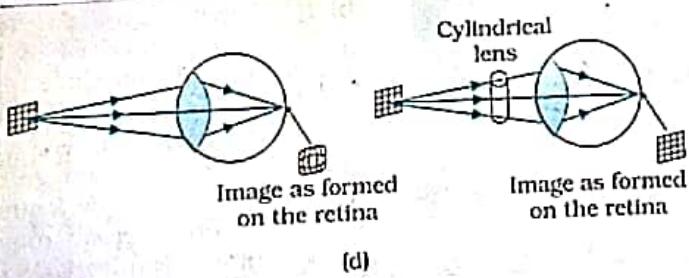


(b)

- * Hypermetropia (Sightlessness) \Rightarrow Hypermetropia is a condition of the eye where the light that comes in does not directly focus on the retina but \rightarrow away from it.

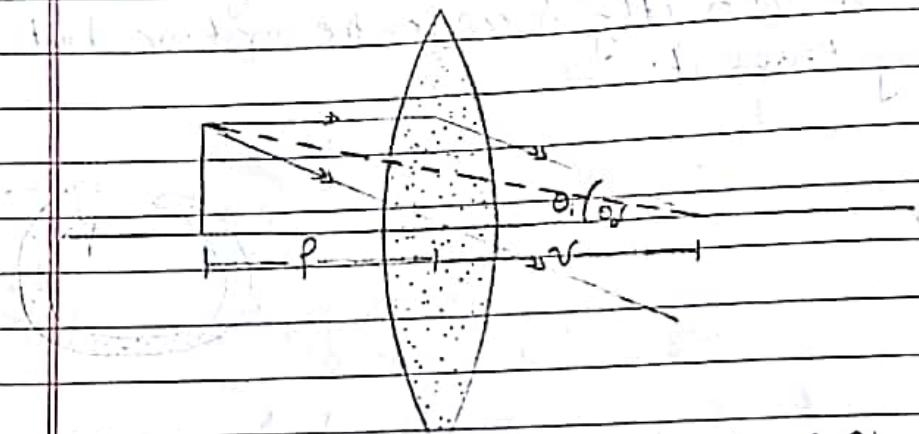


- * Astigmatism \Rightarrow Astigmatism is an optical defect in which vision is blurred due to the inability of the optics of the eye to focus a point object into a sharp focused image on the retina.



Defects in eye accommodation
1. (a) Myopia
2. (b) Hypermetropia
3. (c) Presbyopia

* MICROSCOPE



* Linear magnification = $\frac{\text{height of Image}}{\text{height of Object}}$

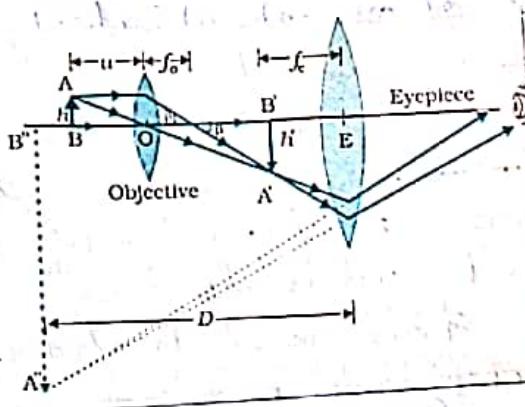
$$\Rightarrow \text{Angular magnification} = \frac{\Theta_i}{\Theta_o}$$

$$\Rightarrow \text{linear magnification} = \frac{v}{u} = \left(1 + \frac{D}{f}\right)$$

$$\Rightarrow \text{Angular magnification} = D/f$$

\Rightarrow A simple microscope has a limited maximum magnification (≤ 9)

* Compound microscope \Rightarrow light microscope that uses two converging lens system: the objective and the eyepiece.



$$m_o = \frac{h'}{h} = \frac{L}{f_o} \quad m_e = 1 + \frac{D}{f_e}$$

$$\Rightarrow m = m_o m_e = \frac{L}{f_o} \left(1 + \frac{D}{f_e}\right) \quad \begin{cases} m = \text{Total} \\ \text{magnification} \end{cases}$$

$$m = \frac{L}{f_o} \left(1 + \frac{D}{f_e}\right)$$

★ TELESCOPE

"Telescope is an instrument that aids in the observation of remote objects by collecting electromagnetic radiation"

⇒ Magnifying power of a telescope is the ratio of the angle β subtended at the eye by the image to the angle α subtended at the eye by the object.

$$m = \frac{\beta}{\alpha} = \frac{P_o}{P_e}$$

Where, P_o and P_e are the focal lengths of the objective and ~~the~~ eyepiece respectively.

★ Dop

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★ Doppler Effect

"Doppler effect is change in frequency of wave for an observer moving relative to its source".

$$\frac{\Delta f}{f_0} = \frac{V_{\text{radial}}}{C}$$

V_{radial} is considered positive
when the source moves away from the observer.

* Red Shift \Rightarrow Source moving away from observer results in Redshift.

* Blue Shift \Rightarrow Source moving toward observer results in Blueshift.

~~$\Delta f = (v_r - v_A) + (v_o - v_A) + A$~~

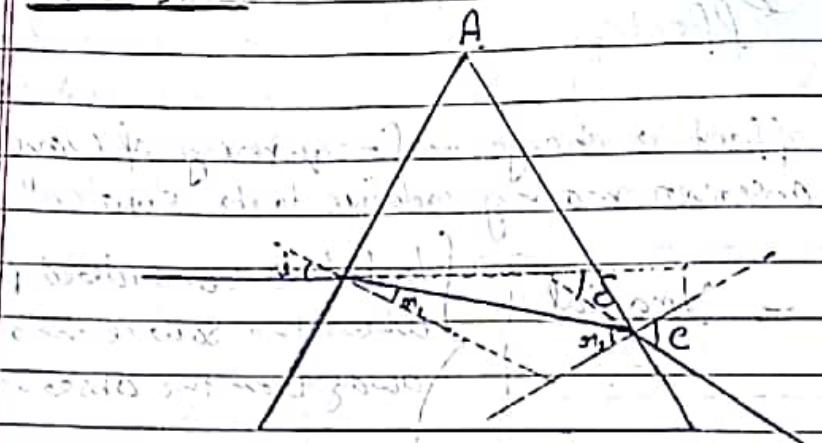
~~$v_r - v_o = A \Leftarrow$~~

~~$(v_r - v) + (v_o - v) = ?$~~

~~$(v_r + v_o) - 2v = ? \Leftarrow$~~

~~$\{ \text{It's always zero!} \} A - 2v = ? \Leftarrow$~~

~~$v_r - v = \text{distance between } v_r \text{ and } v_o$~~
 ~~$v_o - v = \text{distance between } v_o \text{ and } v_r$~~

Poism

$\angle A$ = Angle of poism
 i = Angle of incidence
 e = Angle of emergence

$$A + (90^\circ - \alpha_1) + (90^\circ - \alpha_2) = 180$$

$$\Rightarrow A = \alpha_1 + \alpha_2 - 10$$

$$\delta = (i - \alpha_1) + (e - \alpha_2)$$

$$\Rightarrow \delta = i + e - (\alpha_1 + \alpha_2)$$

$$\Rightarrow \delta = i + e - A \quad \{ \text{from equation 1} \}$$

for minimum deviation $i = e$ ~~$\alpha_1 = \alpha_2$~~
 Hence, $\alpha_1 = \alpha_2$ ~~$\{ \text{from Symmetry} \}$~~

Waves & Vibrations

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$$\Rightarrow \delta = 2i - A \quad \Rightarrow i = \frac{\delta + A}{2}$$

$$\Rightarrow n_1 + n_2 = A$$

$$\Rightarrow 2n_1 = A \quad (\text{as } n_1 = n_2)$$

$$\Rightarrow n_1 = A/2$$

$$\mu = \frac{8 \sin i}{8 \sin n_1} = \frac{8m \left(\frac{\delta + A}{2} \right)}{8m \left(A/h \right)}$$

$$\Rightarrow \mu = \frac{\sin \left(\frac{\delta + A}{2} \right)}$$

$$\mu = \mu_0 + \frac{A}{\lambda^2} \quad \text{Law of dispersion}$$



Understanding Concepts

★ Simple Microscope

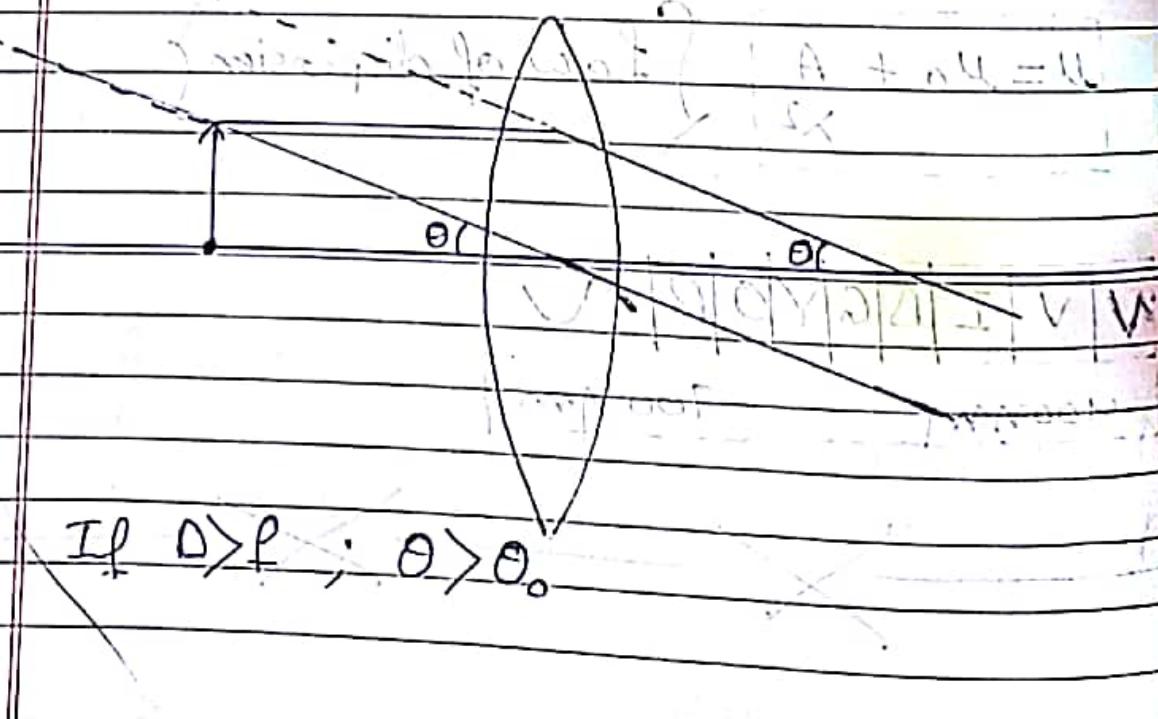
⇒ When we view an object with naked eyes, the object must be placed somewhere between infinity and the near point.

⇒ Maximum angle is subtended on the eye when the object is placed at the near point.

$$\theta_0 = \frac{h}{D} \quad \left. \begin{array}{l} h = \text{size of object} \\ D = \text{least distance} \end{array} \right\}$$

// Suppose, the lens has a focal length f ; $D > f$

$$\theta = \frac{h}{f} \quad \left. \begin{array}{l} f = \text{focal length} \end{array} \right\}$$



$$\text{If } D > f ; \theta > \theta_0$$

Hence, the eye perceives a larger image than it could have had without the microscope.

$$\text{Angular magnification} = \frac{\theta}{\theta_0} = \frac{h/f - D}{h/D}$$

$$\begin{aligned}\text{Lateral magnification} &= \frac{h_i}{h_o} = \frac{V}{u} = V \left(\frac{f}{f+D} \right) \\ &= \left(1 + \frac{D}{f} \right)^{-1}\end{aligned}$$

$$\text{So } V = -D$$

~~in diagram below~~ \Rightarrow ~~with diagram~~

~~a converging lens forms a real image at D~~

~~lens $-D$ forms a virtual image at D~~

~~enlarged image at D formed by $-D$~~

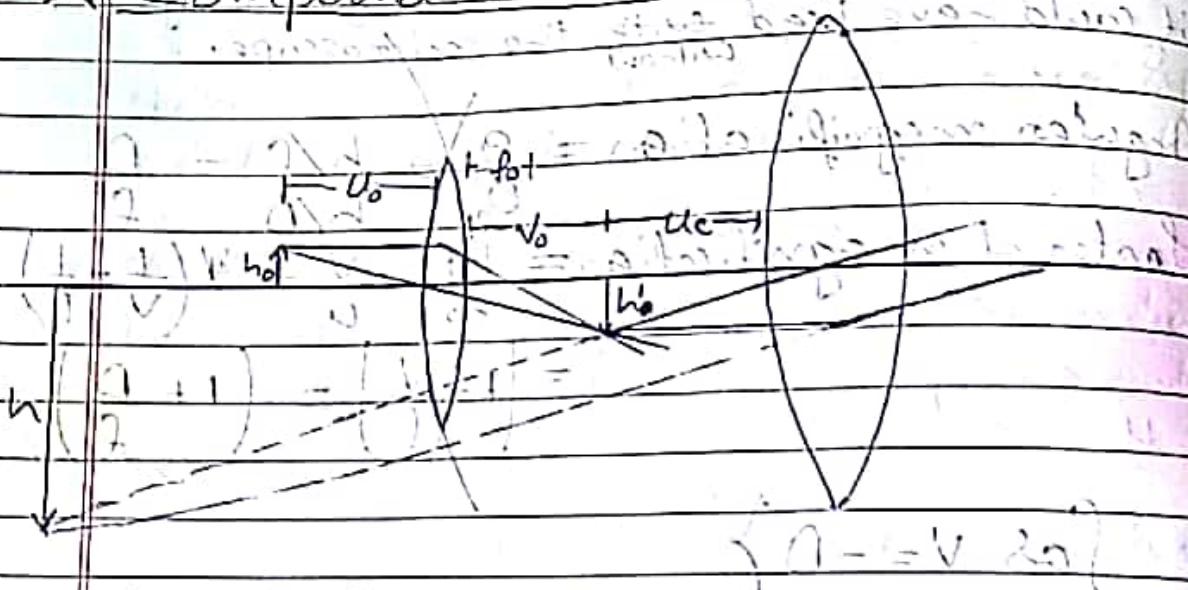
$$M = \frac{V}{u} = A$$

~~lens $-D$ forms a virtual image at D~~

$$M = \frac{V}{u} = A$$



Compound microscope



\Rightarrow Magnification by compound microscope is in two steps :-

- ① In first step the objective produces a magnified image of the given object.
- ② In second step the eyepiece produces an angular magnification

$$\theta_o = h_o / D \quad \text{--- (1)}$$

\Rightarrow When compound microscope is used, the final image subtend an angle θ' on the eye piece

$$\theta' = \frac{h'}{D_e} \quad \text{--- (1)}$$

\Rightarrow Magnifying power of the compound microscope
is therefore :-

As magnification = $\frac{h'}{h} \times D$ where D = distance between

$$m = \frac{h'}{h} = \frac{f_o}{D} \times \frac{D}{f_e} \quad \text{Using } h' = f_o u \text{ and } h = f_e u$$

$$\therefore m = \left(\frac{f_o}{f_e} \right) \left(\frac{D}{D} \right) \quad \text{Hence } A =$$

$$\Rightarrow \frac{f_o}{f_e} \quad \text{and } A = \frac{f_o}{f_e}$$

$$\Rightarrow \left(\frac{L}{f_o} \right) \left(\frac{D}{f_e} \right)$$

Hence,
$$m = \left[\frac{L}{f_o} \left(\frac{D}{f_e} \right) \right]$$

* Limit of Resolution

The fact that a lens forms a disc image of a point source, puts a limit on resolving two neighbouring points imaged by lens.

⇒ A detailed analysis shows that the radius of the central bright region is approximately given by :-

$$\text{or } r_b \approx \frac{0.61\lambda f}{b} \quad \left. \begin{array}{l} b = \text{radius of lens} \\ f = \text{focal length} \end{array} \right\}$$

$$\left(\frac{0.61}{0.7} \right) \left(\frac{1}{0.2} \right) = \text{not small}$$

* Validity of Ray Optics

Ray optics is good appx approximation at a distance Z_F where

$$Z_F = \frac{a^2}{\lambda}$$

{ } a = Size of aperture
 { } aperture \Rightarrow Slit or hole
 { } λ = Wave-length

The quantity Z_F is called Fresnel distance.

Huygens Principle

According to Huygens principle, each point of the wavefront is the source of secondary disturbance. Wavelet emerging from these points spreads out in all directions with speed of the wave. A common tangent to all these wavelets gives the new position of the wavefront at a later time.

Electromagnetic wave

Type	Wavelength Range
① Radio	\rightarrow or μm
② Microwave	$0.1\text{m} - 1\text{m}$
③ Infrared	$1\text{mm} - 700\text{nm}$
④ Light	$780\text{nm} - 400\text{nm}$
⑤ Ultraviolet	$400\text{nm} - 1\text{nm}$
⑥ X-rays	$1\text{nm} - 10^{-3}\text{nm}$
⑦ Gamma rays	$< 10^{-3}\text{nm}$

Optics Revision

① Rectilinear propagation of light.

"In a medium whose refractive index is not changing, light travels in straight path independent of any other parameter"

{ Classical prospective }

Note:

→ In relativistic prospective light bends in a medium with constant refractive index in presence of gravity field. But that deviation is insignificant relative to classical prospective"

② Reflection and Refraction at Plane and Spherical Surface

A) Reflection at Plane Surface

[i] Newton Law of Reflection

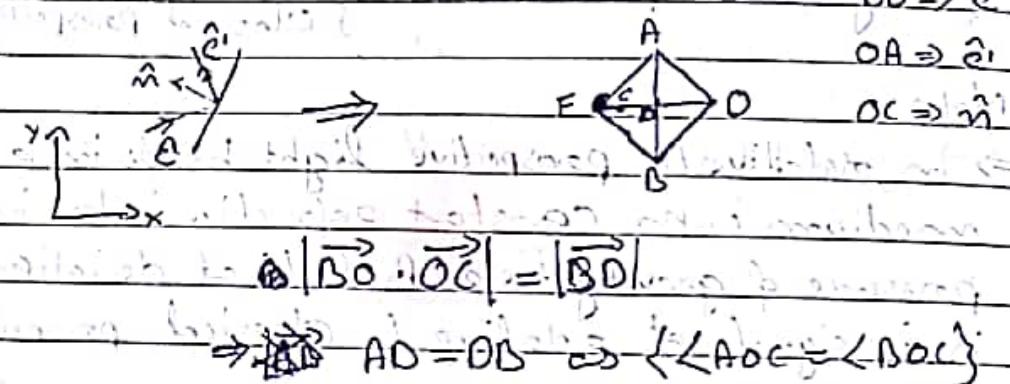
- ① The angle of incidence is equal to the angle of reflection.
- ② The incident ray, the reflected ray and the normal to the reflecting surface are coplanar.

$$\text{Angle of incidence} = \text{Angle of reflection}$$

Diagram

[II] Newton Law of reflection in vector form

Let us consider \hat{e} , \hat{e}' and \hat{n} be unit vectors parallel to incident ray, reflected ray and normal to the plane.



$$|\vec{BO} \cdot \vec{OC}| = |\vec{BD}|$$

$$\Rightarrow \vec{AD} = \vec{OB} - \vec{OA} \quad (\angle AOC = \angle BOD)$$

$$|\vec{AD}| = 2|\vec{BD}| = 2|\vec{BO} \cdot \vec{OC}|$$

~~$$\vec{AO} + \vec{OB} = \vec{OA} + \vec{BD}$$~~

~~$$\vec{AO} + \vec{OB} + \vec{OA} = \vec{OE}$$~~

~~$$|\vec{OE}| = |\vec{AO}| \quad \text{as } AOB\text{ is Square}$$~~

~~$$\Rightarrow \vec{OE} = 2|\vec{BO} \cdot \vec{OC}| \vec{OE}$$~~

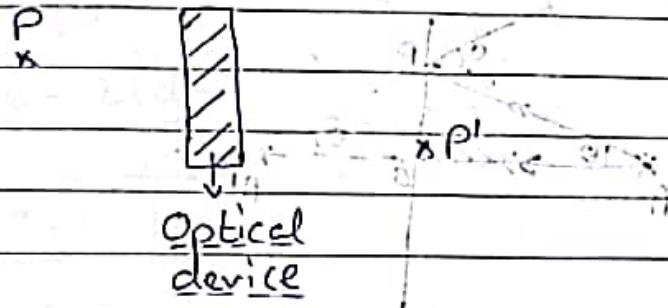
~~$$\Rightarrow \vec{OB} + \vec{OA} = 2|\vec{BO} \cdot \vec{OC}| \vec{OC}$$~~

$$\Rightarrow -\hat{e} + \hat{e}' = 2(\hat{n}\hat{e})\hat{n}$$

~~$$\Rightarrow \hat{e}' = \hat{e} - 2(\hat{n}\hat{e})\hat{n}$$~~

$$\boxed{\hat{e}' = \hat{e} - 2(\hat{n}\hat{e})\hat{n}}$$

III] Image



→ A point is visible if light comes from it or spreads from it or seems to spread from it in all directions.

→ The point P_1 is the image of point P if light rays from P intersect at P_1 .

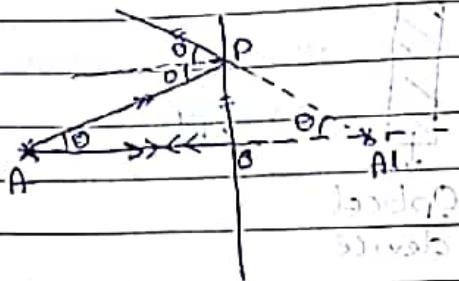
→ Intensity or Brightness of Image will be proportional to amount of light intersect at point P_1 .

- ★ Real Image ⇒ Image in which light rays coming from object actually intersect.

- ★ Virtual Image ⇒ Image in which light rays coming from object only seems to intersect but not actually.

$$(H_{obj} - H_{img}) \theta =$$

IV] Image formation due to plane mirror

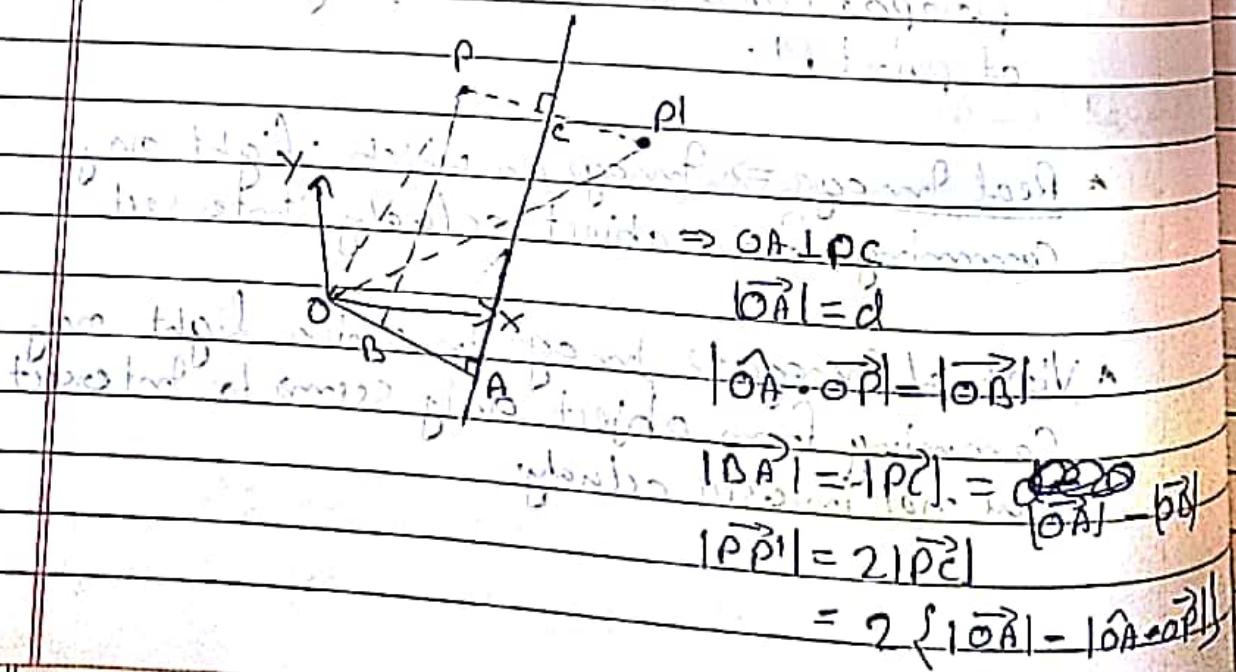


$\triangle APO \sim \triangle A'PO$ [as $\angle PAO = \angle PA'O$ & $\angle POA = \angle POA' = 90^\circ$]
 $\Rightarrow OA = O'A'$ {independent of O }

V] Image formation due to plane mirror in vector form

Let \vec{OA} be the position vector of point object and its image.

Let $\vec{n} = \hat{n}d$ be the normal of plane mirror



$$\vec{OP} + \vec{P'P} = \vec{OP}$$

$$\vec{OP} = \vec{OP} - \vec{PP'}$$

$$\vec{m}_i = \vec{m}_o - 2(d - (om))\hat{n}$$

$$\Rightarrow \boxed{\vec{m}_i = \vec{m}_o - 2[d - (om)]\hat{n}}$$

(B) Reflection at Curved Surface

[I] Mirror formula

$$\frac{1}{V} + \frac{1}{U} = \frac{2}{R} = \frac{1}{f}$$

Mirror formula is symmetric about Image & Object

$$m = -\frac{V}{U}$$

$$\frac{U-1}{U} = \frac{U-1}{V}$$

[II] Coordinate of Image with given Coordinate of object with pole at origin and minor and principle axis oriented along y & x-axis respectively

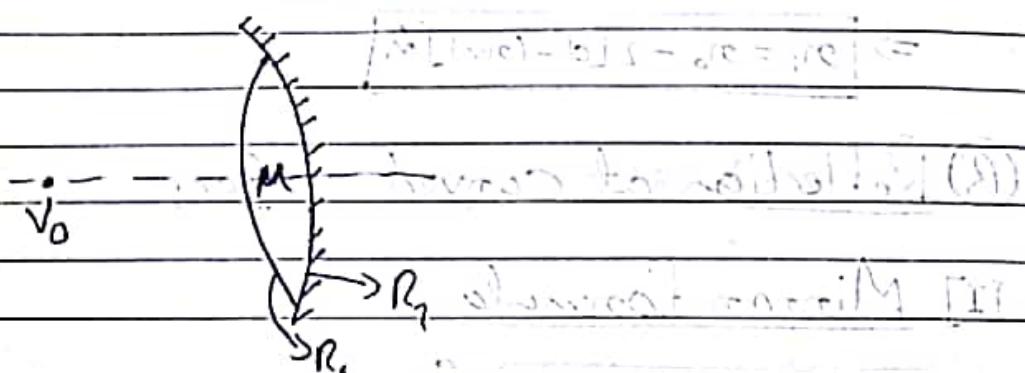
$$[U_x, U_y]$$

$$\Rightarrow \left[\frac{U_x f}{f + b_x}, \frac{-U_y f}{f + b_y} \right]$$

$$\left[\frac{U_x f}{U_x - f}, \frac{-U_y f}{U_y - f} \right]$$

② Refraction at plane Surface

Thin lens Combined With Mirror



$$\mu \times \frac{1}{v_1} - \frac{1}{v_0} = \frac{\mu-1}{R_1}$$

$$\mu \times \frac{1}{v_2} + \frac{1}{v_1} = \frac{2}{R_2}$$

$$\frac{1}{v_2} - \frac{1}{v_0} = \frac{1-\mu}{R_1}$$

$$\textcircled{+} \quad \frac{1}{v_3} - \frac{1}{v_0} = \frac{2\mu + 2(1-\mu)}{R_1 + R_2} = 2(\mu-1) \left[\frac{1}{R_1} - \frac{1}{R_2} \right] + \frac{2}{R_2}$$

$$\Rightarrow \frac{1}{v_3} - \frac{1}{v_0} = \frac{1}{f_m} - \frac{2}{f_l}$$

\Rightarrow So above combination of mirror and lens can be thought as mirror of focal length f .

$$f = \frac{1}{f_m} - \frac{2}{f_l}$$

11 Appendix: Formula Optics

① Spherical mirror

* ~~using $\frac{1}{u} + \frac{1}{v} = \frac{1}{f}$~~ $\Rightarrow \frac{1}{u} + \frac{1}{v} = \frac{1}{f}$

$$\Rightarrow \frac{1}{u} + \frac{1}{v} = \frac{1}{f} \quad \{ \text{Mirror formula} \} \checkmark$$

$$\Rightarrow f = \frac{R}{2} \quad \{ \text{Radius of curvature} \}$$

$$\Rightarrow m = -\frac{v}{u} \checkmark$$

② Apparent depth

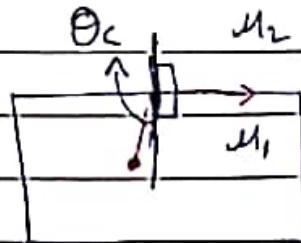
$$\Delta t = \left(1 - \frac{1}{n_2}\right) t \quad \checkmark$$

$$\frac{\text{Real depth}}{\text{Apparent depth}} = \frac{n_1}{n_2} \quad \{ \text{Can be worked out from refractor at spherical surface} \}$$

$$\Rightarrow \frac{t}{t_0} = \frac{n_1}{n_2} \quad \{ \text{Shift formula} \}$$

③ Critical angle

$$\Rightarrow \theta_c = \sin^{-1} \left(\frac{n_2}{n_1} \right) \checkmark$$



④ Prism

$$\Rightarrow n = \frac{\sin \left(\frac{A + \delta_m}{2} \right)}{\sin \left(\frac{A}{2} \right)} \quad \{ \begin{array}{l} \delta_m = \text{angle of minimum deviation} \\ \text{ } \end{array}$$

⑤ Refraction at spherical surface

$$\Rightarrow \frac{n_2 - n_1}{u} = \frac{n_2 - n_1}{R} \quad \{ \text{Refraction formula} \}$$

$$\Rightarrow m = \frac{R-V}{R-U} = \frac{U_1 V}{U_2 U} \checkmark = \frac{V/U_2}{U/U_1}$$

⑥ Refraction through lenses

$$\Rightarrow \frac{1}{f} = (n-1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right) \quad \left\{ \text{lenses makers formula} \right\} \checkmark$$

$$\Rightarrow \frac{1}{f} = \frac{1}{V} - \frac{1}{U} \quad \left\{ \text{lenses formula} \right\} \checkmark$$

$$\Rightarrow m = \frac{V}{U} \checkmark$$

$$\Rightarrow P = \frac{1}{f} \checkmark$$

⑦ Thin lenses in contact

$$\Rightarrow P_{\text{eff}} = \sum P \checkmark$$

⑧ Dispersion through Prism

$$M = M_o + \frac{A}{\lambda} \checkmark$$

⑨ Two thin lenses separated by a distance

$$\Rightarrow \frac{1}{F} = \frac{1}{f_1} + \frac{1}{f_2} - \frac{d}{f_1 f_2} \quad \left\{ \text{effective focal length} \right\}$$

$$\Rightarrow S_2 = \frac{df}{f_1} \quad \left\{ \text{distance from second lens} \right\}$$

~~X~~

~~X~~