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# Lat-Net: Compressing Lattice Boltzmann Flow Simulations using Deep Neural Networks

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## Abstract

1 Computational Fluid Dynamics (CFD) is a hugely important subject with applica-  
2 tions in almost every engineering field, however, fluid simulations are extremely  
3 computationally and memory demanding. Towards this end, we present Lat-Net, a  
4 method for compressing both the computation time and memory usage of Lattice  
5 Boltzmann flow simulations using deep neural networks. Lat-Net employs convo-  
6 lutional autoencoders and residual connections in a fully differentiable scheme to  
7 compress the state size of a simulation and learn the dynamics on this compressed  
8 form. The result is a computationally and memory efficient neural network that  
9 can be iterated and queried to reproduce a fluid simulation. We show that once  
10 Lat-Net is trained, it can generalize to large grid sizes and complex geometries  
11 while maintaining accuracy. We also show that Lat-Net is a general method for  
12 compressing other Lattice Boltzmann based simulations such as Electromagnetism.

## 13 1 Introduction

14 Computational fluid dynamics (CFD) is a branch of fluid mechanics that deals with numerically  
15 solving and analyzing fluid flow problems such as those found in aerodynamics, geology, biology, etc.  
16 CFD simulations are known for their high computational requirements, memory usage, and run times.  
17 Because of this, there is an ever growing body of work using simulation data to create reduced order  
18 models or surrogate models that can be evaluated with significantly less resources. Towards this end,  
19 we develop a neural network approach that both compresses the computation time and memory usage  
20 of fluid simulations.

21 We investigate fluid simulations that contain complex time dependent turbulence. Simulations of  
22 this form are difficult because they require fluid solver to have high resolution and small times steps.  
23 Never the less, they frequently occur in nature and are an important area of study. Motivated by need  
24 for these simulations and the recent success of neural network based models in related areas (1) (2)  
25 (3), we choice this setting to test our model.

26 The most popular approach to modeling fluid flow is with the Navier stokes equation. The solution  
27 to this partial differential equation gives the flow velocity field for a given domain. Recently, a new  
28 method for simulating fluid flow has emerged named the Lattice Boltzmann Method (LBM). It is  
29 derived from the Boltzmann equation and grew out of Lattice Gas Automaton (LGA) in the late 80s  
30 (4). The main advantages of the LBM are its ability to run on complex geometries, its scalability to  
31 parallel architectures (particularly GPUs) and applicability to complex flows that contain phenomena  
32 such as heat transfer and chemical reactions. Our method is centered around this method of simulating  
33 flow.

34 Lat-Net works by compressing the state of a simulation while learning the dynamics of the simulation  
35 on this compressed form. The model can be broken up into three pieces, an encoder, compression

36 mapping, and decoder. The encoder compresses both the state of the simulation as well as the given  
37 boundary conditions. The compression mapping learns the dynamics on the compressed state that  
38 correspond to the dynamics in the fluid simulation. The decoder decompresses the compressed state  
39 allowing either the whole simulation state or desired pieces to be extracted.

40 We focus the content of this paper on LBM fluid simulations because this is the most popular use of the  
41 LBM, however, this method of simulation is known to be able to solve a large set of partial differential  
42 equations (5). In fact, LBM can simulate many physical systems of interest such as Electromagnetism,  
43 Plasma, Multiphase flow, Schrödinger equation etc. (6) (7) (8) (9). With this in mind, we keep our  
44 method general and show evidence our method works equally well on Electromagnetic simulations.  
45 However, because the dominate use of LBM is on fluid flow problems we center discussion on this  
46 subject.

47 Our work has the following contributions.

- 48 • It allows for simulations to be generated with less memory then the original flow solver.  
49 There is a crucial need for such methods because memory requirements grow cubic to grid  
50 size in 3D simulations. In practice, this quickly results in the need for large GPU clusters  
51 (10) (11).
- 52 • Once our model is trained, it can be used to generate significantly larger simulations.  
53 This allows the model to learn from a training set of small simulations and then generate  
54 simulations as much as 16 times bigger with little effect in accuracy.
- 55 • Our method is directly applicable to a variety of physics simulations, not just fluid flow. We  
56 show this with our electromagnetic example and note that the changes to our model are  
57 trivial.

## 58 2 Related Work

59 Recently, there have been several papers applying neural networks to fluid flow problems. Guo etc.  
60 (2) proposed to use a neural network to learn a mapping from boundary conditions to steady state  
61 flow. Most related to our own work, Yang etc. (3) and Tompson etc. (1) use a neural network to solve  
62 the Poisson equation in order to accelerate Eulerian fluid simulations. The key difference between  
63 this and Lat-Net is its ability to compress the memory usage and the generality of our method to other  
64 physics simulations.

65 There has also been an increasing body of work applying neural networks to other physics modeling  
66 problems. For example, neural networks have been readily adopted in many chemistry applications  
67 such as predicting molecular properties from descriptors, protein contact prediction and computational  
68 material design (12). Very recently, neural networks have been applied to quantum mechanics  
69 problems as seen in Mills etc. (13) and Giuseppe etc. (14) where neural networks are used to  
70 approximate solutions to the Schrödinger equation. In high energy Physics, Paganini etc. (15) uses a  
71 generative adversarial networks (GAN)(16) to model electromagnetic showers in a longitudinally  
72 segmented calorimeter. Many of these applications are relatively recent and indicate a resurgence of  
73 interest in applications of neural networks to modeling physics.

74 Reduced order Modeling is an area of research that focuses on techniques to reduce the dimensionality  
75 and computational complexity of mathematical models. A Reduced order model (ROM) is constructed  
76 from high-fidelity simulations and can subsequently be used to generate simulations for lower  
77 computation. The most popular ROM method for fluid dynamics is Galerkin projection (17) (18).  
78 This method uses Proper Orthogonal Decomposition to reduce the dimensionality of flow simulations  
79 and then finds the dynamics on this reduced space. There are other methods that build on this such  
80 as reduced basis methods and balanced truncation (19) (20). While these approaches are centered  
81 around the Navier stokes equation and thus not directly comparable to our own, we note that the  
82 compression mapping present in these methods is typically quite simple. Given the recent success  
83 neural networks have had in creating well structured encodings (such as Variational Autoencoders  
84 (21) (22)), we feel our approach is well justified.

## 85 3 Deep Neural Networks for Compressed Lattice Boltzmann

86 In this section, we present our model for compressing Lattice Boltzmann simulations.

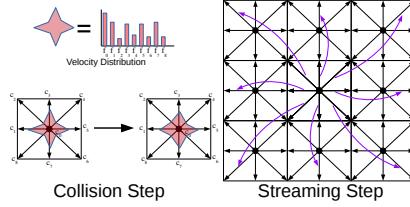


Figure 1: Illustration of the Lattice Boltzmann update steps

### 87 3.1 Review: The Lattice Boltzmann Method

88 In a Lattice Boltzmann simulation, the domain is discretized into an equal sized Cartesian grid. Each  
 89 cell of this grid contains a velocity distribution function  $f_i$  that describes the velocity of flow at that  
 90 point.  $f_i$  has values ranging over  $i$  that correspond to the  $\{\vec{c}_i\}$  directions of flow. In our 2 dimensional  
 91 simulations, there are 9 such directions (D2Q9 scheme). A figure showing this grid structure is seen  
 92 in 1. From the distribution function  $f_i$ , one can calculate the density ( $p$ ) and velocity ( $\vec{u}$ ) of the fluid  
 93 flow with the following equations.

$$p = \sum_i f_i \quad \text{and} \quad \vec{u} = \sum_i \vec{c}_i f_i \quad (1)$$

94 The lattice states are updated with two separate steps, the collision step and the streaming step. The  
 95 collision step mimics the flow interacting with itself and is updated in the following way,

$$f_i^t(x, t + \delta_t) = f_i(x, t) + \frac{1}{\tau}(f_i^{eq} - f_i) \quad (2)$$

96 where  $\tau$  is the relaxation constant and  $f_i^{eq}$  is the flow equilibrium. For our simulations, we use the  
 97  $f_i^{eq}$  from the Lattice Bhatnagar-Gros-Krook (LBGK) scheme (23). After the collision step is applied,  
 98 the flow propagates to adjacent cells following the streaming step.

$$f_i(x + c_i, t + \delta_t) = f_i^t(x, t + \delta_t) \quad (3)$$

99 This step will contain bounce back if one of the adjacent cells is a boundary. Figure 1 illustrates these  
 100 steps for the 2 dimensional case.

101 It is interesting to note the simplicity of this method and its similarity to convolutional neural networks.  
 102 In fact, if we treat the lattice state  $f$  as a tensor of size  $(n_x, n_y, 9)$ , as we do for the remainder of this  
 103 paper, the streaming operator can be mimicked with a 3 by 3 convolution and the collision step can  
 104 be performed with a 1 by 1 convolution (D2Q9). This offers a unique way to interpret our method. In  
 105 a some sense, we are taking a large convolutional neural network and compressing it onto a much  
 106 smaller and more memory efficient network. With this mental picture in mind, we now describe our  
 107 approach.

### 108 3.2 Proposed Architecture

109 Figure 2 shows a sketch of the model. The figure can be understood by following the arrows starting  
 110 from the flow state  $f_t$  and the boundary  $b$ . We treat  $f_t$  as a tensor with shape  $(n_x, n_y, 9)$  for the  
 111 2D case and  $(n_x, n_y, n_z, 15)$  for the 3D case. The boundary is treated as a binary tensor of shape  
 112  $(n_x, n_y, 1)$  and  $(n_x, n_y, n_z, 1)$  with the value being 1 if the cell is solid. Bellow we walk through  
 113 each step of our method.

114 First, we compress both the state of the fluid simulation  $f_t$  and the boundary conditions  $b$  using  
 115 two separate neural networks  $\phi_{enc}$  and  $\phi'_{enc}$  respectively. The result from  $\phi_{enc}$  is a compressed  
 116 representation of the flow  $g_t$  and the result of  $\phi'_{enc}$  are two tensors  $b_{mul}$  and  $b_{add}$  of equal size to  $g_t$ .  
 117 These three tensors represent the entirety of the compressed state of the simulation.

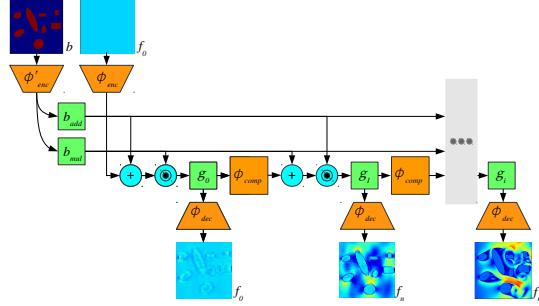


Figure 2: Illustration of the Lat-Net architecture

118 In a Lattice Boltzmann solver, the boundary conditions are used at each time-step to add bounce  
 119 back to the streaming step. In a similar way, our model applies the compressed boundary to the  
 120 compressed state every time-step. We do this in the following way,

$$g_t = (g_t \odot b_{mul}) + b_{add} \quad (4)$$

121 This method proved extremely successful at keeping the boundary information firmly planted through  
 122 the duration of the simulation. This method of applying boundary conditions was inspired by (24)  
 123 where they use a similar method to combine foreground and background information in video  
 124 prediction. After the boundary is applied to  $g_t$ , we can run the state through another neural network  
 125 to emulate the dynamics, i.e.  $\phi_{comp} : g_t \rightarrow g_{t+1}$ . Each step of  $\phi_{comp}$  is equivalent to  $n$  time-steps  
 126 of the Lattice Boltzmann solver. For example, in the 2 dimensional simulation, each step of  $\phi_{comp}$   
 127 mimics 120 steps of the Lattice Boltzmann solver. Once  $g_t$  is computed, we can extract out the  
 128 generated state of the simulation with a decoder network  $\phi_{dec}$ .

### 129 3.3 Network Implementation Details

130 We implement 2 networks trained on 2 dimensional and 3 dimensional lattice simulations. The  
 131 encoder, compression mapping and decoder pieces of the 2D network are each a series of 3 by 3  
 132 residual blocks with the sequences 4x(down res-res)-res, 4x(res), and 3x(transpose conv-res-res)-  
 133 transpose conv, respectively (25). For the 3D network, the sequences are 2x(down res)-res, 3x(res),  
 134 and (transpose conv-res-transpose conv) where 3 by 3 by 3 convolutions are used. The down residual  
 135 blocks are created by changing the first convolution to have kernel size 4 by 4, stride 2 and double the  
 136 filter size. The up sampling is achieved with transpose convolutions of kernel size 4 by 4, stride 2 and  
 137 half the filter size. For the last residual block on the 3D network, the filter size is halved once again.

138 As mentioned above, each network is kept entirely convolutional. Fluid flow is inherently spatially  
 139 correlated so using convolutional layers allows this spatial information to be preserved. Keeping the  
 140 network convolutional also allows different input sizes to be used. This is how our model is able to  
 141 train on small simulations and then generate larger simulations.

142 Residual connections have been used in many deep learning architectures with much success. Adding  
 143 residual connections allows for much deeper networks to be trained, often resulting in improved  
 144 results (25). When training our model, it is necessary to unroll the compression network over several  
 145 time-steps. This has the same effect as making the network deeper. For this reason, it seems logical to  
 146 take advantage of this network architecture. We have seen that removing these residual connections  
 147 results in much slower convergence and worse accuracy.

### 148 3.4 Training Details

149 Lat-Net is trained by unrolling the network and comparing the generated flow with the true. Our  
 150 loss function is Mean Squared Error (MSE) with Image Gradient Difference Loss (GDL) (26). The  
 151 GDL is multiplied by  $\lambda_{GDL}$  and then added to the MES. In all our experiments,  $\lambda_{GDL}$  is set to 0.2.  
 152 Removing the GDL tended to produce less accurate models. Lat-Net is unrolled 5 time-steps and  
 153 then trained with the Adam optimizer (27).

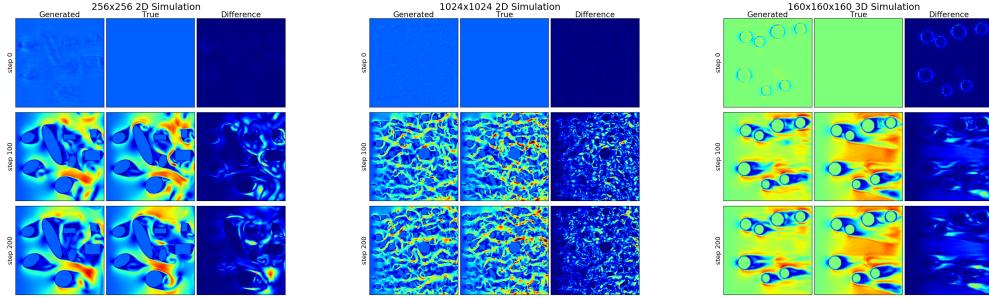


Figure 3: A visual comparison of flows generated by Lat-Net and the Lattice Boltzmann method. Each figure shows the Generated, True, and Difference of the flow for various time-steps.

## 154 4 Experiments

155 In this section, we describe our experiments testing Lat-Net on a variety of problems. Our experiments  
 156 are designed to test our models ability to generate large simulations as well as its transferability to  
 157 new boundary geometries. We also explore computation time and working memory usage. Finally,  
 158 we briefly show results applying this method to electromagnetic simulations.

### 159 4.1 Dataset Generation

160 In order to train and test our model, we generate sets of fluid and electromagnetic simulations. All  
 161 simulations were generated with the MechSys library (28).

162 The train set for the 2D fluid simulations consists of 50 runs of grid size 256 by 256 and 9 directional  
 163 flows in the lattice Boltzmann solver (D2Q9 scheme)(23). The simulations use periodic boundary  
 164 conditions on top and bottom as well as uniform inlet flow and outlet flow of 0.04 from the left and  
 165 right. 8 Objects are placed randomly with height and width sizes ranging from 140 to 20 cells. The  
 166 test set for the 2 dimensional simulations consists of 10 runs of size 256 by 256, and 5 runs of size  
 167 1024 by 1024 with the same boundary conditions and object densities. We also generate a test set of  
 168 size 256 by 512 with vehicle cross sections as objects. There are 28 cross sections used ranging from  
 169 trucks to minivans. For all 2 dimensional simulations, the ratio of network steps to Lattice Boltzmann  
 170 steps is 1 to 120.

171 The train set for the 3D fluid simulations consists of 50 runs of grid size 40 by 40 by 160 and 15  
 172 directional flows in the lattice Boltzmann solver (D3Q15 scheme)(23). Similar to the 2D simulations,  
 173 periodic boundary conditions are used with same inlet and outlet flow. 4 spheres are randomly placed  
 174 with height and width 24. The reason different object geometries and sizes were not explored was  
 175 due to the fact that smaller objects or objects with complex geometries tended to have too course  
 176 a resolution for the lattice Boltzmann solver and larger objects required too large a simulation size.  
 177 The test set comprises 10 runs of 40 by 40 by 160 and 5 runs of 160 by 160 by 160 simulations with  
 178 the same object density. The ratio of network steps to Lattice Boltzmann steps is 1 to 60.

179 The train set for the electromagnetic simulations are grid size 256 by 256 with periodic boundaries.  
 180 An electromagnetic wave is initialized in the top of the simulations and proceeds to interact with  
 181 randomly placed objects of different dielectric constants. When the wave hits these objects, the  
 182 reflection and refraction phenomenon is seen. The test set consists of simulations of size 512 by 512  
 183 with the same object density.

### 184 4.2 Generating Simulations

185 A key component of our model is its ability to generate larger simulations than those trained on. To  
 186 test its effectiveness in doing so, we compare the accuracy of generated 2D and 3D simulations to  
 187 ground truth simulations.

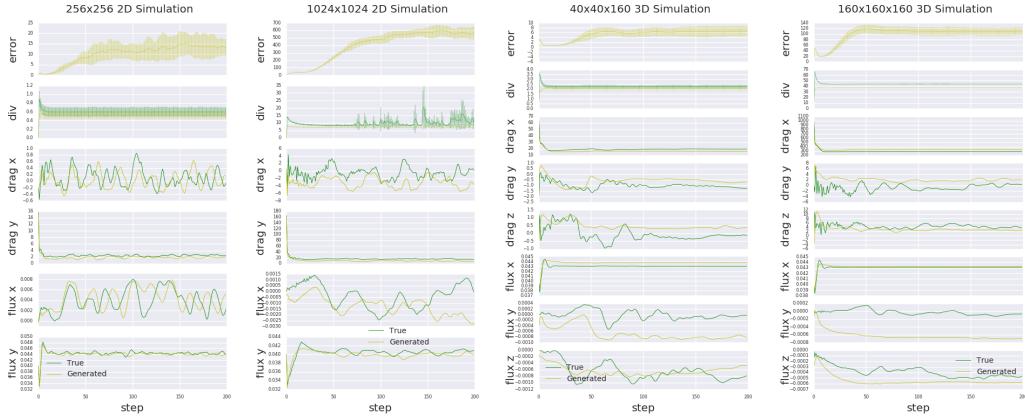


Figure 4: Comparison plot of the flows generated by Lat-Net and the Lattice Boltzmann method. Each plot shows the average mean squared error of the true and generated generated along with the average divergence of the velocity vector field for both simulations. The standard deviation is also displayed. In addition, the calculated values for drag and flux are displayed for a single simulation run. For the 1024 by 1024 simulation, the flow produced by the Lattice Boltzmann solver tended to produce instabilities resulting in the chaotic divergence observed.

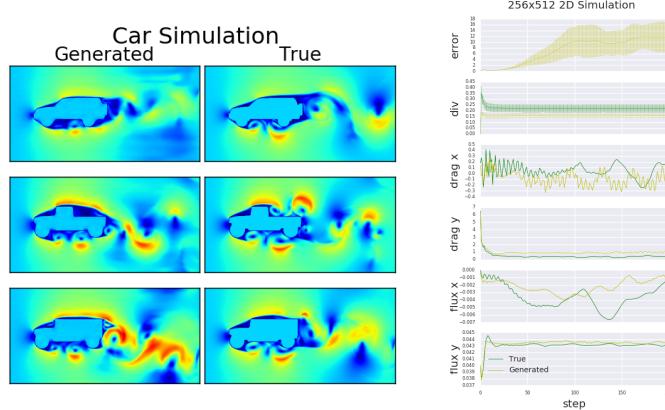


Figure 5: Comparison of generated flows for the vehicle cross section dataset. The images show the generated and true flow at step 100 for 3 different cars in the dataset. The plot shows the same values as described in 4.

Comparing the accuracy of our simulations require some consideration. The naive approach is to compare the MSE between the generated and true simulation at various time-steps. The problem with this approach is that fluid flow is a chaotic dynamical system and small perturbations in flow quickly compound leading to dramatic differences at latter times. For this reason, we compare a variety of metrics in evaluating the generated flows accuracy. Similar to (1), we compare the divergence of the generated and true velocity vector field to test our models stability. We also compare the computed values of drag and flux. Flow simulations are often run to calculate such values so comparing this is a strong indicator of our models real world applicability. The drag is calculated directly from the lattice state via the momentum transfer method (23). The flux value is the average of the flux in each non-boundary cell. These values can be used to calculate important quantities such as the drag coefficient and Reynolds number. Lastly, we visually inspect the produced flow to check for instabilities and blurring effects.

Table 1: Computation Time of Network

Simulation	Comp. Size	Comp. Mapping	Full State	Plane	Line	Point
(1024, 1024, 9)	(64, 64, 128)	2.7 ms	36.2 ms	NA	6.7 ms	6.6 ms
(160, 160, 160, 15)	(40, 40, 40, 64)	23.1 ms	272.1 ms	38.2 ms	25.6 ms	24.1 ms

200 In figure 4, we see the predicted values for different grid sizes in 2D and 3D simulations. In the 2D  
 201 simulations, Lat-Net is able to effectively transfer to larger domain sizes with very similar calculated  
 202 values of drag and flux. The generated flow also maintains its stability even after hundreds of steps. In  
 203 the 3D simulation we see that, while our model predicts realistic values for the 40x40x160 simulation,  
 204 it tends to have a slight bias in the direction of flow that manifests itself in the 160x160x160  
 205 simulation.

206 When visually inspecting our produced flows (figure 3), we see a slight blurring effect but overall  
 207 similar structure in the 2D flows. We attribute this blur effect to the dimensionality reduction and use  
 208 of MSE. This can possibly be overcome with the use of generative adversarial network (16) where  
 209 the loss is derived from a discriminator network. Another solution may be to craft a loss function that  
 210 takes advantage of the statistical properties of flow (7). We leave these pursuits to future work.

211 There is a distinct difference in the generated and true flow for the 3D 160x160x160 simulation.  
 212 While the generated flow appears accurate close to the objects, in regions far between objects the  
 213 network tends to underestimate the flow velocity. We believe this is due to these types of regions not  
 214 being present in the train set and is the probable cause for the biases seen in the drag and flux. As  
 215 mentioned above, our 3D train set is limited due to memory constraints and so developing a diverse  
 216 train set to overcome this proved difficult.

217 The boundaries used in the above evaluation are drawn from the same distribution as the train set.  
 218 This motivates the question of how our model performs on drastically different geometries. To test  
 219 this, we apply our model to predicting flow around vehicle cross sections. Surprisingly, even though  
 220 our model is only trained on flows around simple shapes (ovals and rectangles) it can effectively  
 221 generalize to this distinctly different domain. In figure 5, we see the predicted flows are quite similar  
 222 but with the same blurring effect. Calculating the same values as above, we see the flow is stable and  
 223 produces similar drag and flux.

### 224 4.3 Computation and Memory Compression

225 In this section, we investigate the computational speed-up of our model. The standard performance  
 226 metric for Lattice Boltzmann Codes is Million Lattice Updates per Second (MLUPS). This metric is  
 227 calculated by the following equation,

$$MLUP = \frac{n_x \times n_y \times n_z \times 10^{-6}}{\text{Compute Time}} \quad (5)$$

228 where  $n_x$ ,  $n_y$ , and  $n_z$  are the dimensions of the simulation. For 3 dimensional simulations like the  
 229 ones seen in this paper, a speed of 1,200 MLUPS can be achieved with a Nvidia K20 GPU and single  
 230 precision floats (29). We use this as our benchmark value to compare against.

231 The computation time and memory usage of the encoder can be neglected because this is a one time  
 232 cost for the simulation. In addition, if the simulation is started with uniformly initialized flow as seen  
 233 in our experiments, the computation to compress the flow is extremely redundant and can easily be  
 234 optimized.

235 As seen in table 1, the computation time of the compression mapping is 23.1 ms for a 3D simulation of  
 236 grid size 160 by 160 by 160. Because each step of the compression mapping is equivalent to 60 Lattice  
 237 Boltzmann steps, this equates to 10,600 MLUPS and a roughly 9x speed increase (a similar speed-up  
 238 is seen with the 2D simulation). This does not give a complete picture though. Once the compressed  
 239 states have been generated, the flow must be extracted with the decoder. Unfortunately, this requires  
 240 considerable amounts of computation and memory because it involves applying convolutions to  
 241 the full state size. Fortunately, there are ways around this. In many applications of CFD, it is not  
 242 necessary to have the full state information of the flow at each time-step. For example, calculating  
 243 the drag only requires integrating over the surface of the object. By using the convolutional nature  
 244 of the decoder, we can extract specific pieces of the flow without needing to compute the full state.

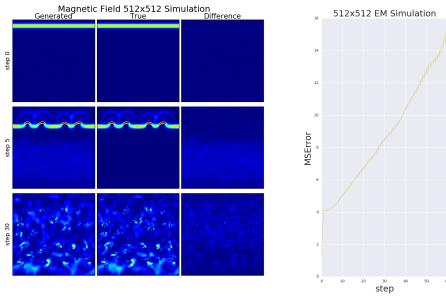


Figure 6: Comparison of generated Electromagnetic fields. The images show the true and generated magnetic field at various time-steps. The reflection and refraction phenomena can clearly be seen in both. The plot shows the mean squared error of the true and generated simulation.

245 In table 1, we show computation times for extracting flow information of a plane, line and single  
 246 point. While these computations can still be somewhat expensive, they do not necessarily need to be  
 247 performed at every time-step and require very little working memory.

248 Unfortunately, there are some measurements that do require the full state information to compute  
 249 such as the average flux seen in our tests. Our method is currently unable to handle these without  
 250 requiring high run-times and large working memory. A possible solution is training a separate neural  
 251 network that takes in the compressed state and predicts the desired measurement. This would negate  
 252 the need to extract out the full state and keep memory usage low. We leave this and similar ideas for  
 253 future work.

254 While Lat-Net compresses the simulation state size by more than an order of magnitude, the working  
 255 memory requirements for the compression network must be considered. A typical GPU based Lattice  
 256 Boltzmann solver requires around 1.5 times as much working GPU memory as the memory size  
 257 of the lattice (29). For example, the maximum sized D3Q15 lattice that can fit on a GPU with 8  
 258 Gigabytes is  $446^3$ . We have observed that in our implementation of Lat-Net, the maximum 3D  
 259 compression network we can run with an 8 Gigabyte GPU corresponds to a lattice size of  $672^3$ . This  
 260 represents a 3.4x efficiency gain in working memory usage. While this is certainly a nontrivial gain,  
 261 we feel that further improvements can be realized with a more memory efficient implementation of  
 262 the compression mapping.

#### 263 4.4 Electromagnetic Results

264 Finally, we illustrate the generality of our method by applying it to electromagnetic simulations. The  
 265 same neural network architecture is used as in the 2D flow simulations with the only difference being  
 266 the filter size on the compression is half of that in the flow network. The loss is kept identical however  
 267 the lattice values are scaled up by a factor of 10 so they are on the same range as the flow lattice  
 268 values. In figure 6, we see very similar waves formed with the same reflection and refraction.

## 269 5 Conclusion

270 Fluid Simulations are incredibly important for a variety of tasks however they are extremely computa-  
 271 tion and memory intensive. In this work, we have developed a unique method to overcome this using  
 272 deep neural networks. We have demonstrated it is capable of accurately reconstructing a variety of  
 273 simulations under different conditions with significantly less computation and memory. We have also  
 274 shown that our method can be readily applied to other physics simulations such as electromagnetic  
 275 simulations. While our method has proved successful on the problems in this paper, there is still  
 276 significant room for improvement. A loss function that either takes into account the statistical nature  
 277 of the flow or uses recent advances in GANs could produce sharper, more realistic flow. Training  
 278 a network to extract desired measurements from the compressed state such as average flux would  
 279 overcome the current memory limitation for such a task. We leave these and other improvements for  
 280 future work.

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