# Useful Identities in Atomic Physics

### Nicolás Quesada

The quantum numbers of the energy eigenstates of the non-relativistic hydrogenoid atom are: n (energy), l (angular momentum) and m (z projection of angular momentum).  $\langle \cdot \rangle$  means average over *energy* eigenstates.

### Scales, Constants and Special Values of Hydrogenoid Wave Functions [1]

Energies for the Coulomb Potential 
$$(V(\mathbf{r}) = -\frac{1}{4\pi\epsilon_0} \frac{Ze^2}{r})$$
 are  $E_n = -\frac{\mu}{2n^2} \left(\frac{Ze^2}{4\pi\epsilon_0\hbar}\right)^2 = -\frac{e^2}{4\pi\epsilon_0a_0} \frac{Z^2}{2n^2} = -\frac{1}{2}\mu c^2 \frac{(Z\alpha)^2}{n^2}$ .

Fine structure constant:  $\alpha = \frac{e^2}{4\pi\epsilon_0\hbar c}$ . Bohr radius:  $a_0 = 4\pi\epsilon_0\hbar^2/(\mu e^2)$ .  $|\psi_{nlm}(0)|^2 = \frac{Z^3}{\pi a_0^3 n^3} \delta_l^0 \delta_m^0$ .

## Expected Values, the Virial Theorem and the Gamma Function [2,3]

Virial Theorem (valid for *any* potential): If  $H = T(\mathbf{p}) + V(\mathbf{r})$  and  $T(\mathbf{p}) = \frac{\mathbf{p}^2}{2\mu}$  then  $2\langle T \rangle = \langle \mathbf{r} \cdot \nabla V \rangle$ .

Expectation values for the Coulomb potential:  $\left\langle \frac{1}{r} \right\rangle = \frac{Z}{a_0 n^2}$ ,  $\left\langle \frac{1}{r^2} \right\rangle = \frac{Z^2}{a_0^2 n^3 (l+1/2)}$ .

 $\text{Recursion Relation: } 0 = \frac{s}{4} \left[ (2l+1)^2 - s^2 \right] \left( \frac{a_0}{Z} \right)^2 \left\langle r^{s-2} \right\rangle - (2s+1) \left( \frac{a_0}{Z} \right) \left\langle r^{s-1} \right\rangle + \frac{s+1}{n^2} \left\langle r^s \right\rangle \; .$ 

$$\Gamma(z) = \int_0^\infty t^{z-1} e^{-t} dt \ , \ \Gamma(n+1) = n! \ , \ \Gamma(1-z) \ \Gamma(z) = \frac{\pi}{\sin(\pi z)} \ , \ \Gamma(z) \ \Gamma\left(z + \frac{1}{2}\right) = 2^{1-2z} \ \sqrt{\pi} \ \Gamma(2z).$$

# Spherical Harmonics, Wigner 3j Symbols and Clebsch – Gordan Coefficients [4]

$$\int_0^{2\pi} d\phi \int_0^{\pi} \sin\theta d\theta \, \mathcal{Y}_{l_1}^{m_1}(\theta,\phi) \mathcal{Y}_{l_2}^{m_2}(\theta,\phi) \mathcal{Y}_{l_3}^{m_3}(\theta,\phi) = \sqrt{\frac{(2l_1+1)(2l_2+1)(2l_3+1)}{4\pi}} \begin{pmatrix} l_1 & l_2 & l_3 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} l_1 & l_2 & l_3 \\ m_1 & m_2 & m_3 \end{pmatrix}.$$

Wigner 3
$$j$$
 — Clebsch–Gordan (CG) relation: 
$$\begin{pmatrix} j_1 & j_2 & j_3 \\ m_1 & m_2 & m_3 \end{pmatrix} \equiv \frac{(-1)^{j_1-j_2-m_3}}{\sqrt{2j_3+1}} \langle j_1 m_1 j_2 m_2 | j_3 - m_3 \rangle.$$

Selection rules for Wigner 3j Symbol  $\begin{pmatrix} l_1 & l_2 & L \\ m_1 & m_2 & -M \end{pmatrix}$  (they are identical to CG Selection Rules):  $-l_i \le m_i \le l_i, \quad m_1 + m_2 = M, \quad |l_1 - l_2| \le L \le l_1 + l_2 \quad l_1 + l_2 + L \in \mathbb{Z}.$ 

Spherical components of a cartesian vector  $\vec{e}=(e_x,e_y,e_z)$ :  $e_{\pm 1}=\mp\frac{1}{\sqrt{2}}\left(e_x\pm ie_y\right)$  and  $e_0=e_z$ .

 $\text{Ladder Operators: } \hat{L}_{\pm} \equiv \hat{L}_x \pm i \hat{L}_y. \hspace{0.5cm} \hat{L}_{\pm} |l,m\rangle = \hbar \sqrt{l(l+1) - m(m\pm 1)} |l,m\pm 1\rangle.$ 

#### References

- [1] P. Ewart, Atomic Physics Lecture notes https://users.physics.ox.ac.uk/~ewart/
- [2] S. Jeon, Lecture Notes for Quantum Physics II. http://www.physics.mcgill.ca/~jeon/Phys457/
- [3] E. W. Weisstein, "Gamma Function." http://mathworld.wolfram.com/GammaFunction.html
- [4] E. W. Weisstein, "Wigner 3j-Symbol." http://mathworld.wolfram.com/Wigner3j-Symbol.html