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In what follows, we develop a variational ansatz for the ground state of quantum electrodynamical Hamiltonians. In particular, we develop an ansatz in which the ground state can be considered as a factorizable state of an effective matter and effective photon quasiparticle both in their respective ground states. This ansatz, analogous to the Hartree-Fock ansatz, leads to coupled eigen-equations describing the ground-state of the light-matter system. We benchmark our ansatz against a multi-level and multi-mode Rabi model parameterized by the strength of the characteristic momentum matrix elements describing the coupling of the matter and the electromagnetic fields. We find that when the momentum matrix elements are sufficiently weak, our theory finds the ground state energy to a remarkable accuracy of less than >99.9%, even in the non-perturbative coupling regime. Even in regimes where the correlation energy is 50% of the total energy, far outside of the validity domain of our ansatz, our ansatz predicts the energy up to a few percent. In regimes where our results are accurate, we claim we have found the effective quasiparticle description of the ground state. Moreover, in this regime, the ground state interaction energy manifests itself as the difference in zero-point energies between the quasiparticles and the bare particles. This generalizes the notion of the Casimir energy to a single atom in a cavity.

Our results service several aims. First, they suggest a novel variational approach for understanding ground states in ultrastrongly coupled light-matter systems. Second, they provide a general framework for understanding light-matter decoupling effects in the regime of strong light-matter interactions. Third, they prove a non-perturbative theory of Casimir-Polder forces. Finally, they provide a rigorous concept of correlation energy in quantum electrodynamics in a way analogous to electronic structure theory.

The outline of this work is as follows: in section I, we develop the variational ansatz for the ground state of the Hamiltonian of macroscopic quantum electrodynamics. In section II, we calculate the ground state energy according to this ansatz for the multi-level and multi-mode Rabi model and compare it to numerical diagonalization of the same problem. In section III, we study the accuracy of our ansatz as a function of the correlation energy in the system, and test several approaches to recover the correlation energy. We consider both perturbation theory formulated in the interacting modes as well as a self-consistent perturbation theory which allows the perturbative energy shift to influence the modes.

In this work, we consider the light-matter Hamiltonian corresponding to an emitter placed at position z = d in a

one-dimensional cavity oriented along the z-direction. As the cavity is considered for simplicity to be one-dimensional, the electric field is oriented along a single direction, denoted x, while the magnetic field is oriented along a direction transverse to both the electric field and the cavity length, denoted y. We consider the interaction in the velocity gauge. The Hamiltonian can then be written as:

$$H = H_{\text{matter}} + \frac{\epsilon_0}{2} \int dz \ (E^2 + c^2 B^2) + \frac{e}{m} A(d) p + \frac{q^2}{2m} A^2(d).$$
 (1)

The fields can be expressed as a mode expansion via:

$$A(z,t) = \sum_{n=1}^{\infty} \sqrt{\frac{\hbar}{2\epsilon_0 \omega_n}} (F_n(z) a_n e^{-i\omega_n t} + F_n^*(z) a_n^{\dagger} e^{i\omega_n t}),$$
(2)

with $E(z,t)=-\partial_t A(z,t)$ and $B(z,t)=\partial_z A(z,t)$. The $F_n(z)$ are the mode functions of the cavity, normalized such that $\int dz \ |F_n(z)|^2=1$. For a cavity of length L the modes are given by $F_n(z)=\sqrt{\frac{2}{L}}\sin\left(\frac{n\pi z}{L}\right)$ and the corresponding mode frequencies are $\omega_n=\frac{n\pi c}{L}$. The matter Hamiltonian we take to be a multilevel system with an arbitrary number of levels, N_a . The matter system we describe as an N_a site system, perhaps describing a molecule in an LCAO framework. Thus we parameterize the general family of matter Hamiltonians

$$H_{\text{matter}} = \sum_{i=1}^{N_a - 1} V_i |i\rangle\langle i| + t(|i\rangle\langle i + 1| + |i + 1\rangle\langle i|)$$
 (3)

I. VARIATIONAL ANSATZ

The concept behind the variational ansatz is shown in Figure (1). We seek a description of the ground state of the system as $\Psi \rangle \approx |\tilde{g}\rangle \otimes |\tilde{0}\rangle + |\delta\rangle$, where a \sim denotes an effective quantity. In other words, we seek a description in which the state is a factorizable state of matter and photon *quasi*particles, up to some correlations quantified by the state $|\delta\rangle$. These correlations essentially hybridize $|\tilde{g}\rangle \otimes |\tilde{0}\rangle$ with excitations of the effective emitter and virtual effective photons.

We start our analysis by considering a family of ansatze in which $\delta = 0$ (uncorrelated ansatz). The expectation value of

the Hamiltonian is

$$\langle \Psi | H | \Psi \rangle = \langle \tilde{g} | H_{\text{matter}} | \tilde{g} \rangle$$

$$+ \frac{\hbar}{4} \int dz \sum_{n=1}^{\infty} \left(\omega_n |F_n|^2 - \frac{c^2}{\omega_n} F_n^* \partial_z^2 F_n \right)$$

$$+ \frac{\hbar q^2}{4m\epsilon_0} \sum_{n=1}^{\infty} \int dz \, \delta(z - d) |F_n|^2$$
(4)

Notably, in this ansatz, the term in the Hamiltonian coupling the momentum of the matter to the photon makes no contribution. We impose constraints of matter normalization and photon mode normalization through Lagrange multipliers $\epsilon_a(\langle \tilde{g} | \tilde{g} \rangle - 1)$ and $\frac{\hbar \lambda_n}{2} \left(\int dz |F_n|^2 - 1 \right)$.

- A. Concept: Effectively non-interacting matter and photon systems
 - II. APPLICATION TO A MULTI-LEVEL SYSTEM COUPLED TO A MULTI-MODE CAVITY
 - III. STRATEGIES FOR RECOVERING THE CORRELATION ENERGY

IV. OUTLOOK

V. ACKNOWLEDGEMENTS

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$$|\Psi\rangle = \sum_{a=0}^{N_a} \sum_{n_1 \cdots n_{N_p}=0}^{\infty} c_{a;n_1 \cdots n_{N_p}} |a\rangle \otimes |n_1 \cdots n_N\rangle \qquad |\Psi\rangle \approx |\tilde{g}\rangle \otimes |\tilde{0}\rangle + \sum_{\tilde{a}=1}^{N_a} \sum_{\tilde{k}=1}^{N_p} c_{\tilde{a};\tilde{k}} |\tilde{a}\rangle \otimes |\tilde{1}_{\tilde{k}}\rangle$$

Fig. 1. Ground-state ansatz applied to matter in a cavity: effectively decoupled matter and photons. (Left) Bare description of the coupled light-matter ground state in terms of many virtual excitations of the emitter state and the bare cavity photons. (Right) Quasiparticle description of the coupled system as a factorable state an effective emitter in its ground state and the vacuum of an effective photonic degree of freedom.