# Working Title

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Modeling the liquid in a binary PFC model as regular solution instead of an ideal one leads to lots of interesting features. Look at all those features. We'll even discuss how it can be a contributor to non-classical nucleation pathways. Very sexy.

Keywords: nucleation, growth, phase field crystal

#### I. INTRODUCTION

# II. FREE ENERGY FUNCTIONAL FOR BINARY ALLOYS

To contruct a free energy functional for a binary alloy we begin by splitting the free energy into ideal and excess components.

$$\mathcal{F}[\rho_A, \rho_B] = \mathcal{F}_{id}[\rho_A, \rho_B] + \mathcal{F}_{ex}[\rho_A, \rho_B]$$
 (1)

The ideal component comes from the non-interacting, kinetic part of the Hamiltonian while the excess component comes from the interaction potential part of the Hamiltonian. The ideal component can be computed exactly as,

$$\beta \mathcal{F}_{id}[\rho_A, \rho_B] = \sum_{i=A,B} \int dr \rho_i(r) \left( \ln(\Lambda_i^3 \rho_i(r)) - 1 \right) \quad (2)$$

Where  $\rho_i(r)$  and  $\Lambda_i^3$  are the density field and thermal de Broglie volume of the ith species. The excess free energy cannot, in general, be computed exactly and thus approximate it by expanding the excess free energy about a uniform reference mixture with densities  $\rho_A^0$  and  $\rho_B^0$ .

$$\beta \mathcal{F}_{ex}[\rho_A, \rho_B] = \beta \mathcal{F}_{ex}^0 + \sum_i \int dr C_i^{(1)}(r) \Delta \rho_i(r)$$

$$+ \frac{1}{2} \sum_{i,j} \int dr dr' \Delta \rho_i(r') C_{ij}^{(2)}(r, r') \Delta \rho_j(r) + \dots$$
(3)

We now perform a change of variable to solute concentration, c and dimensionless total density, n:

$$n = \Delta \rho / \rho_0 = \frac{\Delta \rho_A + \Delta \rho_B}{\rho_A^0 + \rho_B^0} \tag{4}$$

$$c = \frac{\rho_B}{\rho_A + \rho_B} \tag{5}$$

Under this change of variables the change in ideal free energy from the reference becomes,

$$\frac{\beta \Delta \mathcal{F}_{id}[n,c]}{\rho_0} = \int dr (n+1) \ln(n+1) - n$$

$$+ (n+1) \left[ c \ln \left( \frac{c}{c_0} \right) + (1-c) \ln \left( \frac{1-c}{1-c_0} \right) \right]$$

Following Greenwood et al [cite] if we assume that the concentration field varies over a much longer length scale than the total density, the change in excess free energy from the reference can be written as,

$$\frac{\beta \Delta \mathcal{F}_{ex}[n,c]}{\rho_0} = -\frac{1}{2} \int dr n \left( \int dr' C_{eff}^n n' + C_{eff}^c c' \right) - \frac{1}{2} \int dr (c - c_0) \left( \int dr' C_{eff}^{c_0 n} n' + C_{eff}^{c c_0} c' \right)$$
(7)

Where the n-c effective correlation functions are,

$$C_{eff}^{n} = c^{2} C_{BB}^{(2)} + (1 - c)^{2} C_{AA}^{(2)} + 2c(1 - c) C_{AB}^{(2)}$$
 (8)

$$C_{eff}^{c} = cC_{BB}^{(2)} - (1 - c)C_{AA}^{(2)} + (1 - 2c)C_{AB}^{(2)}$$
 (9)

$$C_{eff}^{c_0n} = C_{eff}^c \tag{10}$$

$$C_{eff}^{c_0c} = C_{BB}^{(2)} + C_{AA}^{(2)} - 2C_{AB}^{(2)}$$
(11)

#### A. Simplified Regular Free Energy Functional

Following Greenwood  $et\ al\ [cite]$  we can now expand the ideal free energy about reference density in a Taylor expansion and coarse grain the excess contributions the the free energy. In contrast to Greenwood  $et\ al$ , we will not assume that the k=0 mode of the effective concentration-concentration correlation is zero. The effect of eliminating this assumption is to produce a free energy of mixing with a regular solution model instead of an ideal solution. That is to say, there is both an entropy and enthaply of mixing in the theory.

$$\frac{\beta \Delta \mathcal{F}[n,c]}{\rho_0} = \int dr \frac{n^2}{2} - \eta \frac{n^3}{6} + \chi \frac{n^4}{12} - \frac{1}{2} n C_{eff}^n * n + \int dr \frac{W_c}{2} |\nabla c|^2 + (n+1) f_{mix}(c) + \frac{1}{2} \epsilon c^2$$
(12)

## B. Model Dynamics

# III. EQUILIBRIUM PROPERTIES

- A. Syntectic Phase Diagram
- B. Precipitation from Solution

## IV. APPLICATION

- A. Non-classical nucleation pathways of precipation from solution
- B. Anomolous growth of nanoparticles