AST 301 Problem Set 8

Nathan Spilker

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a) We must prove that $\boldsymbol{\xi} \cdot \mathbf{U}$ is a constant of motion, where U is the tangent of a geodesic, or equivalently:

$$\nabla_{\mathbf{U}}(\boldsymbol{\xi} \cdot \mathbf{U}) = 0$$

This can be rewritten as:

$$\nabla_{\mathbf{U}}(\xi_{\mu}U^{\mu}) = 0$$

This is equivalent to (Schutz 6.48)

$$U^{\beta} \nabla_{\beta} (\xi_{\mu} U^{\mu}) = 0$$

Applying Leibniz's product rule:

$$U^{\beta}(U^{\mu}\xi_{\mu;\beta} + \xi_{\mu}U^{\mu}_{;\beta}) = 0$$

A geodesic parallel transports its own tangent vector, so

$$U^{\beta}U^{\mu}_{;\beta} = 0$$

We are left with:

$$U^{\beta}U^{\mu}\xi_{\mu;\beta}=0$$

From the killing equation, we have:

$$\xi_{\mu;\nu} + \xi_{\nu;\mu} = 0$$

In the sum two lines above, each component will subtract identically from another component. Equivalently $\xi_{\mu;\nu}$ is an anti-symmetric tensor while $U^{\mu}U^{\nu}$ is a symmetric tensor.

$$U^{01}U^{10}\xi_{0:1} + U^{10}U^{01}\xi_{1:0} = 0$$

due to killing's equation. Therefore it is satisfied that

$$\nabla_{\mathbf{U}}(\boldsymbol{\xi} \cdot \mathbf{U}) = 0$$

b) We have:

$$K^{\mu} = T^{\mu\nu}\xi_{\nu}$$

Differentiating both sides with respect to μ :

$$K^{\mu}_{\;\;;\mu} = (T^{\mu\nu}\xi_{\nu})_{;\mu}$$

By Leibniz's product rule:

$$=T^{\mu\nu}\xi_{\nu;\mu}+T^{\mu\nu}_{\ ;\mu}\xi_{\nu}$$

We have $T^{\mu\nu}_{\ ;\mu}=0$ by conservation of energy and momentum. Now our equation is:

$$=T^{\mu\nu}\xi_{\nu;\mu}$$

 $\xi_{\mu;\nu}$ is an anti-symmetric tensor while $T^{\mu\nu}$ is a symmetric tensor, so the sum will be 0. We have then proved:

$$K^{\mu} = 0$$

c) From the problem statement we have:

$$\xi_{\mu;\nu} + \xi_{\nu;\mu} = (g_{\mu\gamma}\delta_0^{\gamma})_{;\nu} + (g_{\nu\gamma}\delta_0^{\gamma})_{;\mu} = 0$$

$$=g_{\mu 0;\nu}+g_{\nu 0;\mu}$$

Expanding the covariant derivatives we get:

$$= g_{\mu 0,\nu} - g_{\lambda 0} \Gamma^{\lambda}_{\mu \nu} + g_{\nu 0,\mu} - g_{\lambda 0} \Gamma^{\lambda}_{\nu \mu}$$

From the definition of Christoffel symbols we get:

$$=g_{\mu 0,\nu}-\frac{1}{2}g_{\lambda 0}g^{\gamma \lambda}(g_{\gamma \mu,\nu}+g_{\gamma \mu,\nu}-g_{\nu \mu,\gamma})+g_{\nu 0,\mu}-\frac{1}{2}g_{\lambda 0}g^{\gamma \lambda}(g_{\gamma \nu,\mu}+g_{\gamma \mu,\nu}-g_{\nu \mu,\gamma})$$

Simplifying:

$$= g_{\mu 0,\nu} - \frac{1}{2} \delta_0^{\gamma} (g_{\gamma\mu,\nu} + g_{\gamma\nu,\mu} - g_{\nu\mu,\gamma}) + g_{\nu 0,\mu} - \frac{1}{2} \delta_0^{\gamma} (g_{\gamma\nu,\mu} + g_{\gamma\mu,\nu} - g_{\nu\mu,\gamma})$$

$$= g_{\mu 0,\nu} - \frac{1}{2} (g_{0\mu,\nu} + g_{0\nu,\mu} - g_{\nu\mu,0}) + g_{\nu 0,\mu} - \frac{1}{2} (g_{0\nu,\mu} + g_{0\mu,\nu} - g_{\nu\mu,0})$$

$$= g_{\nu\mu,0}$$

This proves that the metric does not depend on x^0 .

a) We have:

$$F_{\alpha\beta;\gamma} + F_{\gamma\alpha;\beta} + F_{\beta\gamma;\alpha} = 0$$

Expanding the covariant derivatives, we have:

$$F_{\alpha\beta,\gamma} - F_{\mu\beta}\Gamma^{\mu}_{\alpha\gamma} - F_{\alpha\mu}\Gamma^{\mu}_{\beta\gamma} + F_{\gamma\alpha,\beta} - F_{\mu\alpha}\Gamma^{\mu}_{\gamma\beta} - F_{\gamma\mu}\Gamma^{\mu}_{\alpha\beta} + F_{\beta\gamma,\alpha} - F_{\mu\gamma}\Gamma^{\mu}_{\beta\alpha} - F_{\beta\mu}\Gamma^{\mu}_{\gamma\alpha} = 0$$

 $F^{\mu\nu}$ is anti-symmetric and Christoffels are symmetric on their lower two indices:

$$F_{\alpha\beta,\gamma} + F_{\gamma\alpha,\beta} + F_{\beta\gamma,\alpha} (F_{\mu\beta} \Gamma^{\mu}_{\alpha\gamma} - F_{\mu\beta} \Gamma^{\mu}_{\alpha\gamma}) + (F_{\alpha\mu} \Gamma^{\mu}_{\beta\gamma} - F_{\alpha\mu} \Gamma^{\mu}_{\beta\gamma}) + (F_{\gamma\mu} \Gamma^{\mu}_{\alpha\beta} - F_{\gamma\mu} \Gamma^{\mu}_{\alpha\beta}) = 0$$

Which simplifies to:

$$F_{\alpha\beta,\gamma} + F_{\gamma\alpha,\beta} + F_{\beta\gamma,\alpha} = 0$$

b) We have:

$$F^{\mu\nu}_{,\nu} = 4\pi J^{\mu}$$

Differentiating both sides with respect to μ (minkowskian metric):

$$F^{\mu\nu}_{,\nu\mu} = 4\pi J^{\mu}_{,\mu}$$

 $F^{\mu\nu}$ is anti-symmetric:

$$-F^{\nu\mu}_{,\nu\mu}=4\pi J^{\mu}_{,\mu}$$

Because partial derivatives commute we have:

$$-F^{\nu\mu}_{,\mu\nu} = 4\pi J^{\mu}_{,\mu}$$

We also have:

$$F^{\mu\nu}_{,\nu\mu}=4\pi J^{\mu}_{,\mu}$$

We are free to switch dummy indices:

$$F^{\nu\mu}_{,\mu\nu} = 4\pi J^{\mu}_{,\mu}$$

Which proves that

$$J^{\mu}_{,\mu} = 0$$

c) We start with:

$$F^{\mu\nu}_{\ ;\nu} = 4\pi J^{\mu}$$

Differentiating with respect to μ :

$$F^{\mu\nu}_{;\nu\mu} = 4\pi J^{\mu}_{;\mu}$$

We can write:

$$\frac{1}{2}(F^{\mu\nu}_{\ ;\nu\mu}-F^{\nu\mu}_{\ ;\nu\mu})=4\pi J^{\mu}_{\ ;\mu}$$

We are free to change dummy indices

$$\frac{1}{2}(F^{\mu\nu}_{\;\;;\nu\mu}-F^{\mu\nu}_{\;\;;\mu\nu})=4\pi J^{\mu}_{\;\;;\mu}$$

Multiplying both sides by a negative 1

$$\frac{1}{2}(F^{\mu\nu}_{;\mu\nu} - F^{\mu\nu}_{;\nu\mu}) = -4\pi J^{\mu}_{;\mu}$$

We have the identity for commutators of 2,0 tensors:

$$F^{\mu\nu}_{;\mu\nu} - F^{\mu\nu}_{;\nu\mu} = F^{\sigma\nu}R^{\mu}_{\sigma\nu\mu} + F^{\mu\sigma}R^{\nu}_{\sigma\nu\mu}$$

After substituting this identity into the above equation, we are left with:

$$8\pi J^{\mu}_{\;\;;\mu} = -F^{\sigma\nu}R^{\mu}_{\sigma\nu\mu} - F^{\mu\sigma}R^{\nu}_{\sigma\nu\mu}$$

Because Riemann tensor is anti-symmetric on the last two indices:

$$= -F^{\sigma\nu}R^{\mu}_{\sigma\nu\mu} + F^{\mu\sigma}R^{\nu}_{\sigma\mu\nu}$$

Relabeling indices:

$$= -F^{\sigma\mu}R^{\nu}_{\sigma\mu\nu} + F^{\mu\sigma}R^{\nu}_{\sigma\mu\nu}$$

Because $F^{\mu\nu}$ is anti-symmetric:

$$= F^{\mu\sigma}R^{\nu}_{\sigma\mu\nu} + F^{\mu\sigma}R^{\nu}_{\sigma\mu\nu}$$
$$= 2F^{\mu\sigma}R^{\nu}_{\sigma\mu\nu}$$

Riemann tensor is anti-symmetric on the last two indices:

$$=-2F^{\mu\sigma}R^{\nu}_{\sigma\nu\nu}$$

We recognize this as the Ricci tensor

$$=-2F^{\mu\sigma}R_{\sigma\mu}$$

This sum is equal to zero because $F^{\mu\nu}$ is an anti-symmetric tensor and the ricci tensor is symmetric. Therefore,

$$J^{\mu}_{:\mu} = 0$$

a) Calculated nonzero Christoffels with Mathematica script, (code attached)

$$\begin{split} \Gamma^{t}_{yy} &= -e^{2f(x-t)}f^{'}(x-t) \\ \Gamma^{t}_{zz} &= -e^{2g(x-t)}g^{'}(x-t) \\ \Gamma^{x}_{yy} &= -e^{2f(x-t)}f^{'}(x-t) \\ \Gamma^{x}_{zz} &= -e^{2g(x-t)}g^{'}(x-t) \\ \Gamma^{y}_{yt} &= \Gamma^{y}_{ty} &= -f^{'}(x-t) \\ \Gamma^{z}_{zt} &= \Gamma^{z}_{tz} &= -g^{'}(x-t) \\ \Gamma^{y}_{yx} &= \Gamma^{y}_{xy} &= f^{'}(x-t) \\ \Gamma^{z}_{zx} &= \Gamma^{z}_{xz} &= g^{'}(x-t) \\ \end{split}$$

- b) Components of the Riemann tensor calculated through the mathematica script as well.
 - c)
 - d) Definition of the Ricci tensor:

$$R_{\alpha\beta} = R^{\mu}_{\alpha\mu\beta}$$

Vacuum field equations state that:

$$R_{\alpha\beta} = 0$$

From $R_{zxz}^x = 0$ we get:

$$2g'(x-t)^2 + g''(x-t) = 0$$

From $R_{tyt}^y = 0$ we get:

$$2f'(x-t)^2 + f''(x-t) = 0$$