

# AST 301 Problem Set 7

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## 1

Must prove  $g_{\alpha\beta;\gamma} = 0$ . Definition of Christoffel symbols:

$$\Gamma_{\mu\nu}^{\lambda} = \frac{1}{2}g^{\lambda\kappa}(g_{\kappa\nu,\mu} + g_{\mu\kappa,\nu} - g_{\mu\nu,\kappa})$$

From the definition of the covariant derivative, we have:

$$g_{\alpha\beta;\gamma} = g_{\alpha\beta,\gamma} - \Gamma_{\alpha\gamma}^{\sigma}g_{\sigma\beta} - \Gamma_{\beta\gamma}^{\omega}g_{\alpha\omega}$$

using def. of Christoffel symbols:

$$= g_{\alpha\beta,\gamma} - \frac{1}{2}g^{\sigma\kappa}(g_{\kappa\gamma,\alpha} + g_{\kappa\alpha,\gamma} - g_{\alpha\gamma,\kappa})g_{\sigma\beta} - \frac{1}{2}g^{\omega\epsilon}(g_{\epsilon\beta,\gamma} + g_{\epsilon\gamma,\beta} - g_{\beta\gamma,\epsilon})g_{\alpha\omega}$$

Relationship between metric tensor and upper metric tensor:

$$g^{\sigma\kappa}g_{\sigma\beta} = \delta_{\beta}^{\kappa}$$

Using this property and the index-canceling property of the kronecker delta, we get:

$$\begin{aligned} &= g_{\alpha\beta,\gamma} - \frac{1}{2}(g_{\beta\gamma,\alpha} + g_{\beta\alpha,\gamma} - g_{\alpha\gamma,\beta}) - \frac{1}{2}(g_{\alpha\beta,\gamma} + g_{\alpha\gamma,\beta} - g_{\beta\gamma,\alpha}) \\ &= (g_{\alpha\beta,\gamma} - \frac{1}{2}g_{\alpha\beta,\gamma} - \frac{1}{2}g_{\alpha\beta,\gamma}) + (\frac{1}{2}g_{\beta\gamma,\alpha} - \frac{1}{2}g_{\beta\gamma,\alpha}) + (\frac{1}{2}g_{\alpha\gamma,\beta} - \frac{1}{2}g_{\alpha\gamma,\beta}) \\ &g_{\alpha\beta;\gamma} = 0 \end{aligned}$$

## 2

a)

$$\Gamma_{\mu\nu}^{\mu} = \frac{1}{2}(\ln|g|)_{,\nu}$$

expanding the derivative on the RHS:

$$= \frac{1}{2g}g_{,\nu}$$

Eqn 6.40 in Schutz is:

$$\Gamma_{\mu\alpha}^{\alpha} = (\sqrt{-g})_{,\mu}/\sqrt{-g}$$

Expanding the derivative on the RHS, we get

$$= \frac{1}{2g}g_{,\mu}$$

Because  $\Gamma_{\mu\alpha}^{\alpha} = \Gamma_{\alpha\mu}^{\alpha}$  and we may change the free variable as we choose, we have proved the identity:

$$\Gamma_{\mu\nu}^{\mu} = \frac{1}{2}(\ln|g|)_{,\mu}$$

b)

c)

$$\begin{aligned} F^{\mu\nu}_{;\nu} &= (\sqrt{-g}F^{\mu\nu})_{,\nu}/\sqrt{-g} \\ &= F^{\mu\nu}_{,\nu} + F^{\nu\mu}\frac{1}{2g}(g)_{,\nu} \end{aligned}$$

From the notes, we have the definition:

$$F^{\mu\nu}_{;\nu} = g^{-1/2}(g^{1/2}F^{\mu\nu})_{,\nu} + \Gamma_{\lambda\nu}^{\mu}F^{\lambda\nu}$$

expanding the derivative:

$$= F^{\nu\mu}_{,\nu} + F^{\nu\mu}\frac{1}{2g}(g)_{,\nu} + \Gamma_{\lambda\nu}^{\mu}F^{\lambda\nu}$$

Now we must prove that

$$\Gamma_{\lambda\nu}^{\mu}F^{\lambda\nu} = 0$$

Both  $\nu$  and  $\lambda$  are dummy indices. There are  $4^2$  components for each value of  $\mu$ . 4 are null, when  $\lambda = \nu$  ( $F^{\lambda\nu}$  is anti-symmetric). Because  $\Gamma_{\lambda\nu}^{\mu} = \Gamma_{\nu\lambda}^{\mu}$

all components subtract each other out, for example:  $\Gamma_{01}^{\mu} F^{01} = -\Gamma_{10}^{\mu} F^{10}$   
Therefore,

$$F^{\mu\nu}_{;\nu} = (\sqrt{-g} F^{\mu\nu})_{,\nu} / \sqrt{-g}$$

d)

$$g^{\alpha\beta} g_{\beta\mu,\nu} + g^{\alpha\beta}_{,\nu} g_{\beta\mu} = 0$$

By the product rule, we get:

$$(g^{\alpha\beta} g_{\beta\mu})_{,\nu} = 0$$

Using  $g^{\alpha\beta} g_{\beta\mu} = \delta_{\mu}^{\alpha}$ , we can see that:

$$(\delta_{\mu}^{\alpha})_{,\nu} = 0$$

proving the identity:

$$g^{\alpha\beta} g_{\beta\mu,\nu} = -g^{\alpha\beta}_{,\nu} g_{\beta\mu}$$

e)

$$g^{\mu\nu}_{,\alpha} = -\Gamma_{\beta\alpha}^{\mu} g^{\beta\nu} - \Gamma_{\beta\alpha}^{\nu} g^{\mu\beta}$$

From definition of the covariant derivative of the metric tensor, we have:

$$g^{\mu\nu}_{;\alpha} = g^{\mu\nu}_{,\alpha} + \Gamma_{\beta\alpha}^{\mu} g^{\beta\nu} + \Gamma_{\beta\alpha}^{\nu} g^{\mu\beta} = 0$$

It follows directly that:

$$g^{\mu\nu}_{,\alpha} = -\Gamma_{\beta\alpha}^{\mu} g^{\beta\nu} - \Gamma_{\beta\alpha}^{\nu} g^{\mu\beta}$$

### 3

from the metric, we have:

$$g_{tt} = -e^{2\Phi(r)}$$

$$g_{rr} = e^{2\Lambda(r)}$$

$$g_{\theta\theta} = r^2$$

$$g_{\phi\phi} = r^2 \sin^2 \theta$$

The  $g^{\mu\mu} = 1/g_{\mu\mu}$ , because  $g_{\mu\nu}$  is diagonal.

Definition of Christoffel symbols:

$$\Gamma_{\mu\nu}^{\lambda} = \frac{1}{2} g^{\lambda\kappa} (g_{\kappa\nu,\mu} + g_{\mu\kappa,\nu} - g_{\mu\nu,\kappa})$$

We have the relation  $\lambda = \kappa$  because the metric is a diagonal matrix. We can only have nonzero components when  $\lambda = \nu$ ,  $\lambda = \mu$ , or  $\mu = \nu$  also due to the metric's diagonality. Using the equation for the Christoffel symbols:

$$\Gamma_{\mu\nu}^{\lambda} = \frac{1}{2}g^{\lambda\kappa}(g_{\kappa\nu,\mu} + g_{\mu\kappa,\nu} - g_{\mu\nu,\kappa})$$

We find:

$$\Gamma_{tr}^t = \Gamma_{rt}^t = \Phi'$$

$$\Gamma_{\theta r}^{\theta} = \Gamma_{r\theta}^{\theta} = \frac{1}{r}$$

$$\Gamma_{\phi r}^{\phi} = \Gamma_{r\phi}^{\phi} = \frac{1}{r}$$

$$\Gamma_{\phi\theta}^{\phi} = \Gamma_{\theta\phi}^{\phi} = \cot(\theta)$$

$$\Gamma_{rr}^r = \Lambda'$$

$$\Gamma_{tt}^r = -e^{2(\Phi-\Lambda)}$$

$$\Gamma_{\theta\theta}^r = -re^{-2\Lambda}$$

$$\Gamma_{\phi\phi}^r = -r\sin^2(\theta)e^{-2\Lambda}$$

All other Christoffel symbols are zero.

## 4

From the metric, we have:

$$g_{tt} = -e^{2f(t,x)}$$

$$g_{xx} = e^{2f(t,x)}$$

$g^{xx} = \frac{1}{g_{xx}}$  and  $g^{tt} = \frac{1}{g_{tt}}$  because the metric is diagonal. From the definition of the Riemann tensor

$$R_{xtx}^t = \Gamma_{xx,t}^t - \Gamma_{xt,x}^t + \Gamma_{\sigma t}^t \Gamma_{xx}^{\sigma} - \Gamma_{\sigma x}^t \Gamma_{xt}^{\sigma}$$

Calculating the Christoffel symbols:

$$\begin{aligned}
\Gamma_{xx}^t &= \frac{1}{2}g^{tt}(g_{tx,x} + g_{xt,x} - g_{xx,t}) \\
&= \frac{1}{2}e^{-2f}(2\frac{\partial f}{\partial t}e^{2f}) \\
&= \frac{\partial f}{\partial t}
\end{aligned}$$

$$\begin{aligned}
\Gamma_{xt}^t &= \frac{1}{2}g^{tt}(g_{tx,t} + g_{tt,x} - g_{xt,t}) \\
&= \frac{-1}{2}e^{-2f}(-2\frac{\partial f}{\partial x}e^{2f}) \\
&= \frac{\partial f}{\partial x}
\end{aligned}$$

$$\begin{aligned}
\Gamma_{tt}^t &= \frac{1}{2}g^{tt}(g_{tt,t} + g_{tt,t} - g_{tt,t}) \\
&= \frac{-1}{2}e^{-2f}(-2\frac{\partial f}{\partial t}e^{2f}) \\
&= \frac{\partial f}{\partial t}
\end{aligned}$$

$$\begin{aligned}
\Gamma_{xx}^x &= \frac{1}{2}g^{xx}(g_{xx,x} + g_{xx,x} - g_{xx,x}) \\
&= \frac{1}{2}e^{-2f}(2\frac{\partial f}{\partial x}e^{2f}) \\
&= \frac{\partial f}{\partial x}
\end{aligned}$$

$$\begin{aligned}
\Gamma_{xt}^x &= \frac{1}{2}g^{xx}(g_{xx,t} + g_{xt,x} - g_{xt,x}) \\
&= \frac{1}{2}e^{-2f}(2\frac{\partial f}{\partial t}e^{2f}) \\
&= \frac{\partial f}{\partial t}
\end{aligned}$$

Plugging these into our equation for the Riemann tensor, we get:

$$R_{xtx}^t = \frac{\partial^2 f}{\partial t^2} - \frac{\partial^2 f}{\partial x^2} - (\frac{\partial f}{\partial t})^2 + (\frac{\partial f}{\partial x})^2 + (\frac{\partial f}{\partial t})^2 - (\frac{\partial f}{\partial x})^2$$

$$= \frac{\partial^2 f}{\partial t^2} - \frac{\partial^2 f}{\partial x^2}$$

We are also looking for, using null coordinates:

$$R_{vuv}^v = \Gamma_{vv,u}^v - \Gamma_{vu,v}^v + \Gamma_{\sigma u}^v \Gamma_{vu}^\sigma - \Gamma_{\sigma v}^v \Gamma_{vu}^\sigma$$

Metric:

$$g_{uv} = -e^{2f}$$

upper metric obtained from taking inverse of the matrix  $g_{uv}$ ,

$$g^{uv} = -e^{-2f}$$

Finding Christoffels:

$$\begin{aligned} \Gamma_{vv}^v &= \frac{1}{2} g^{vu} (g_{uv,v} + g_{uv,v} - g_{vv,v}) \\ &= \frac{-1}{2} e^{-2f} \left( -4 \frac{\partial f}{\partial v} e^{2f} \right) \\ &= 2 \frac{\partial f}{\partial v} \end{aligned}$$

$$\begin{aligned} \Gamma_{vu}^v &= \frac{1}{2} g^{vu} (g_{uv,u} + g_{uu,v} - g_{vu,u}) \\ &= 0 \end{aligned}$$

$$\begin{aligned} \Gamma_{uu}^v &= \frac{1}{2} g^{vu} (g_{uu,u} + g_{uu,u} - g_{uu,u}) \\ &= 0 \end{aligned}$$

$$\begin{aligned} \Gamma_{vu}^u &= \frac{1}{2} g^{vu} (g_{vv,u} + g_{vu,v} - g_{vu,v}) \\ &= 0 \end{aligned}$$

Therefore,

$$R_{vuv}^v = 2 \frac{\partial^2 f}{\partial u \partial v}$$