Entropy stable notes

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1 Curved meshes

Non-affine geometric factors

Let F(U) denote the flux matrix whose rows are

$$m{F}(m{U}) = \left(egin{array}{c} m{F}_x \ m{F}_y \end{array}
ight).$$

The conservation law we're interested in is the following

$$\frac{\partial \boldsymbol{U}}{\partial t} + \nabla \cdot \boldsymbol{F}(\boldsymbol{U}) = 0,$$

where the divergence is taken over each column of F(U).

Let G denote the Jacobian of the geometric mapping

$$oldsymbol{G}_{ij} = rac{\partial \widehat{oldsymbol{x}}_j}{\partial oldsymbol{x}_i},$$

and let J denote the determinant of G. One can show (Kopriva?) that

$$\widehat{\nabla} \cdot (J \mathbf{G}^T) = 0.$$

We also have that the mapped normals obey the property

$$JG\widehat{n}\widehat{J}^f = nJ^f$$
.

At the continuous level, the physical gradient and divergence satisfy

$$(\nabla u, \boldsymbol{v})_{D^k} = \left(J\boldsymbol{G}\widehat{\nabla} u, \boldsymbol{v} \right)_{\widehat{D}}, \qquad (\nabla \cdot \boldsymbol{u}, v)_{D^k} = \left(\widehat{\nabla} \cdot \left(J\boldsymbol{G}^T u \right), \boldsymbol{v} \right)_{\widehat{D}},$$

as well as a corresponding integration by parts property.

2 Flux differencing

Replace the flux derivative with

$$(\nabla \cdot \boldsymbol{F}_{S}(\boldsymbol{U}(\boldsymbol{x}), \boldsymbol{U}(\boldsymbol{x}'))|_{\boldsymbol{x}'=\boldsymbol{x}})_{Dk}$$
.

It is possible to remove the effect of geometric aliasing by evaluating the above term via

$$\left(\left.\widehat{
abla}\cdot oldsymbol{F}_{S}^{k}(oldsymbol{U}(oldsymbol{x}),oldsymbol{U}(oldsymbol{x}'))\right|_{oldsymbol{x}'=oldsymbol{x}}
ight)_{\widehat{D}}, \qquad oldsymbol{F}_{S}^{k}=\left\{\!\left\{Joldsymbol{G}
ight\}\!\right\}oldsymbol{F}_{S}\left(oldsymbol{U}(oldsymbol{x}),oldsymbol{U}(oldsymbol{x}')
ight)$$

We can evaluate the new geometrically de-aliased flux using the discrete gradient of Chen and Shu

$$(\nabla_{h} \cdot \boldsymbol{u}, v)_{\Omega} = \sum_{k} \left(\widehat{\nabla} \cdot J \boldsymbol{G}^{T} \Pi_{N} \boldsymbol{u}, v \right)_{\widehat{D}} + \frac{1}{2} \left\langle J \boldsymbol{G}^{T} \boldsymbol{u}^{+} - \Pi_{N} J \boldsymbol{G}^{T} \boldsymbol{u}, v \widehat{\boldsymbol{n}} \right\rangle_{\partial \widehat{D}} + \frac{1}{2} \left\langle J \boldsymbol{G}^{T} \boldsymbol{u} - \Pi_{N} J \boldsymbol{G}^{T} \boldsymbol{u}, \Pi_{N} \left(v \widehat{\boldsymbol{n}} \right) \right\rangle_{\partial \widehat{D}}$$

Another alternative is to treat curvilinear grids using a direct skew-symmetric form.

3 Limiting

Ignoring high order terms in the Taylor expansion, we have

$$u(\Pi_N v) - u = \frac{\partial u}{\partial v} (\Pi_N v - v).$$

We want to limit u with a compression towards the mean $u = \bar{u} + \theta(u - \bar{u})$ such that $|u(\Pi_N v)| \le C|u|$.

3.1 Option 1

One option: ensure $|u(\Pi_N v) - u(\bar{v})| \approx |u - u(\bar{v})|$ based on Taylor series.

$$u\left(\Pi_N v\right) - u(\bar{v})$$

Limit $\tilde{u} = \bar{u} + \theta(u - \bar{u})$ such that

$$\min \Pi_N v \ge \min v, \qquad \max \Pi_N v \le \max v.$$

Another option: estimate or filter $\Pi_N v$

3.2 Other options

- Adaptively choose time-step size?
- Expensive option: bisection-type algorithm for finding theta.