UNIVERSIDAD DE SEVILLA

Probing pairing correlations in nuclei with (t,p) reactions

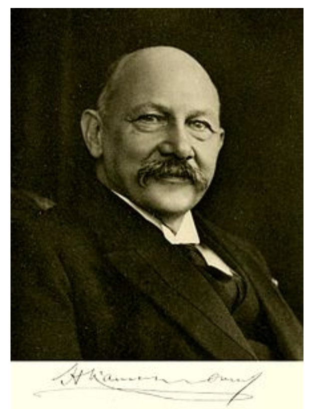
Gregory Potel Aguilar

December 17, 2025

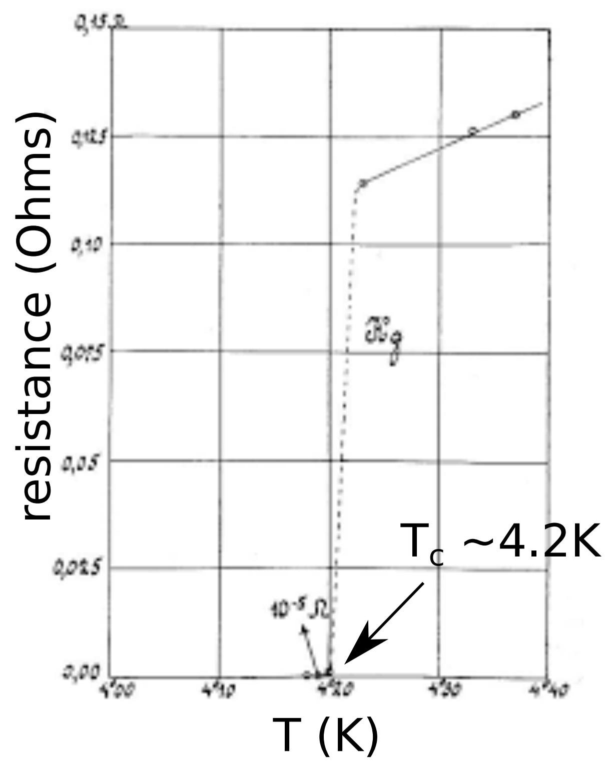
UNIVERSIDAD DE SEVILLA 1505

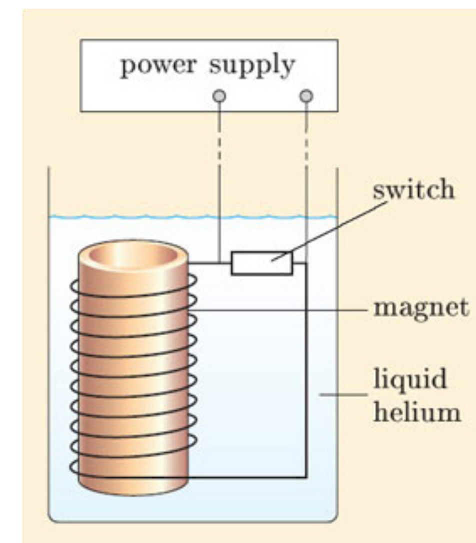
Escuela Técnica Superior de INGENIERÍA DE SEVILLA

Introduction: superconductivity in metals



H. Kamerlingh Onnes



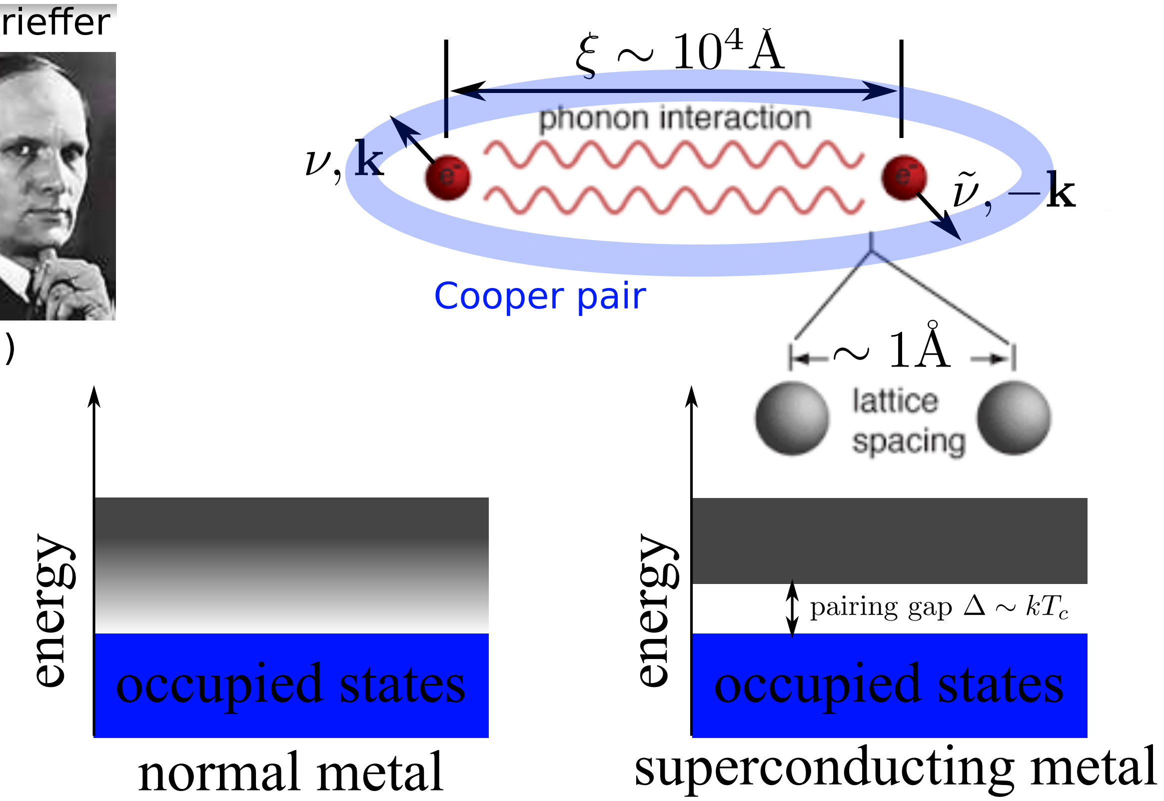


- In 1911, H. K. Onnes liquefies Helium and discovers superconductivity in mercury.
- When cooled below a critical temperature (e.g. $T_c = 7.26$ K for lead, $T_c = 3.69$ K for tin), many metals become superconductors.
- Persistent supercurrents can be induced in superconducting coils.

BCS theory and Cooper pairs

BCS theory (1957)

More than 40 years later arrives the theoretical breakthrough!

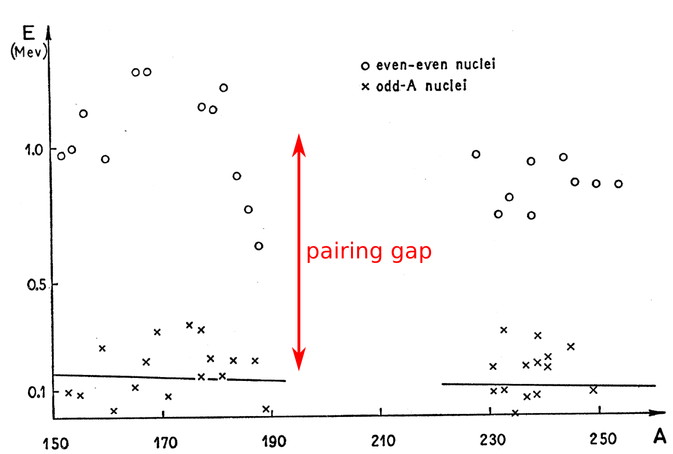


- Below T_c, electrons form enormous (correlation length $\xi \sim 10^4$Å) quasi-bosons (Cooper pairs).
- The binding interaction results from the screening of the Coulomb force and the exchange of lattice phonons.
- An energy gap develops in the low-lying spectrum.
- The Cooper pairs form a condensate $|BCS\rangle = \prod_\nu \left(U_\nu + V_\nu e^{i\phi} a_\nu^\dagger a_{\tilde{\nu}}^\dagger \right) |0\rangle$

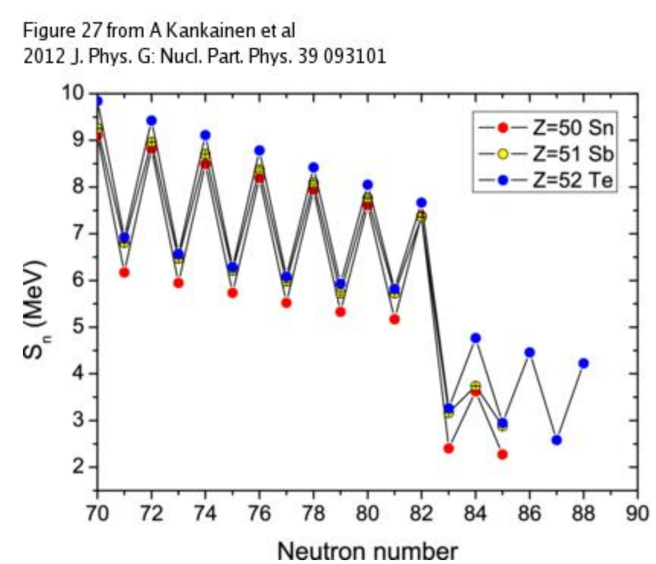
Some experimental evidence of nuclear superfluidity

- Gap in the spectrum of even–even nuclei associated with the breaking of a Cooper pair.
- Odd–even mass staggering: enhanced binding for even number of nucleons.
- Enhanced two–nucleon transfer reactions due to the coherence of the Cooper pair wave function.

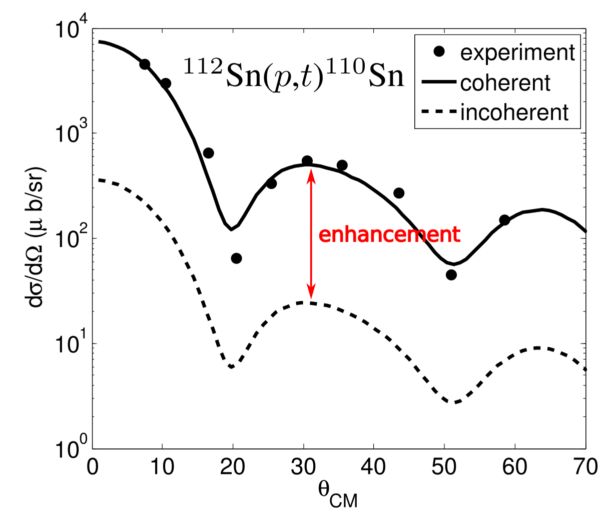
Bohr, Mottelson and Pines (1958)



gap in low-lying spectrum

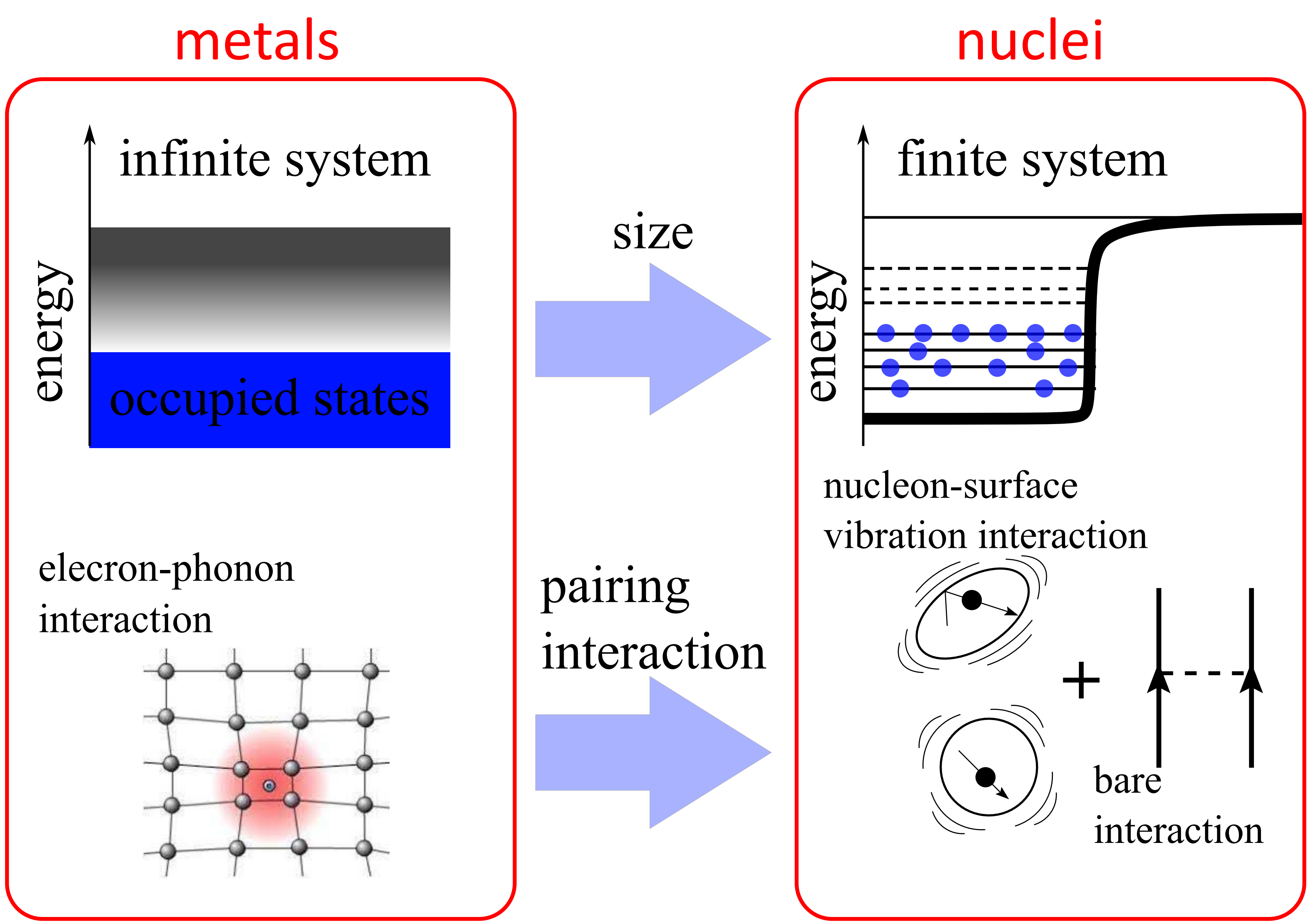


odd-even staggering



enhancement of 2-transfer

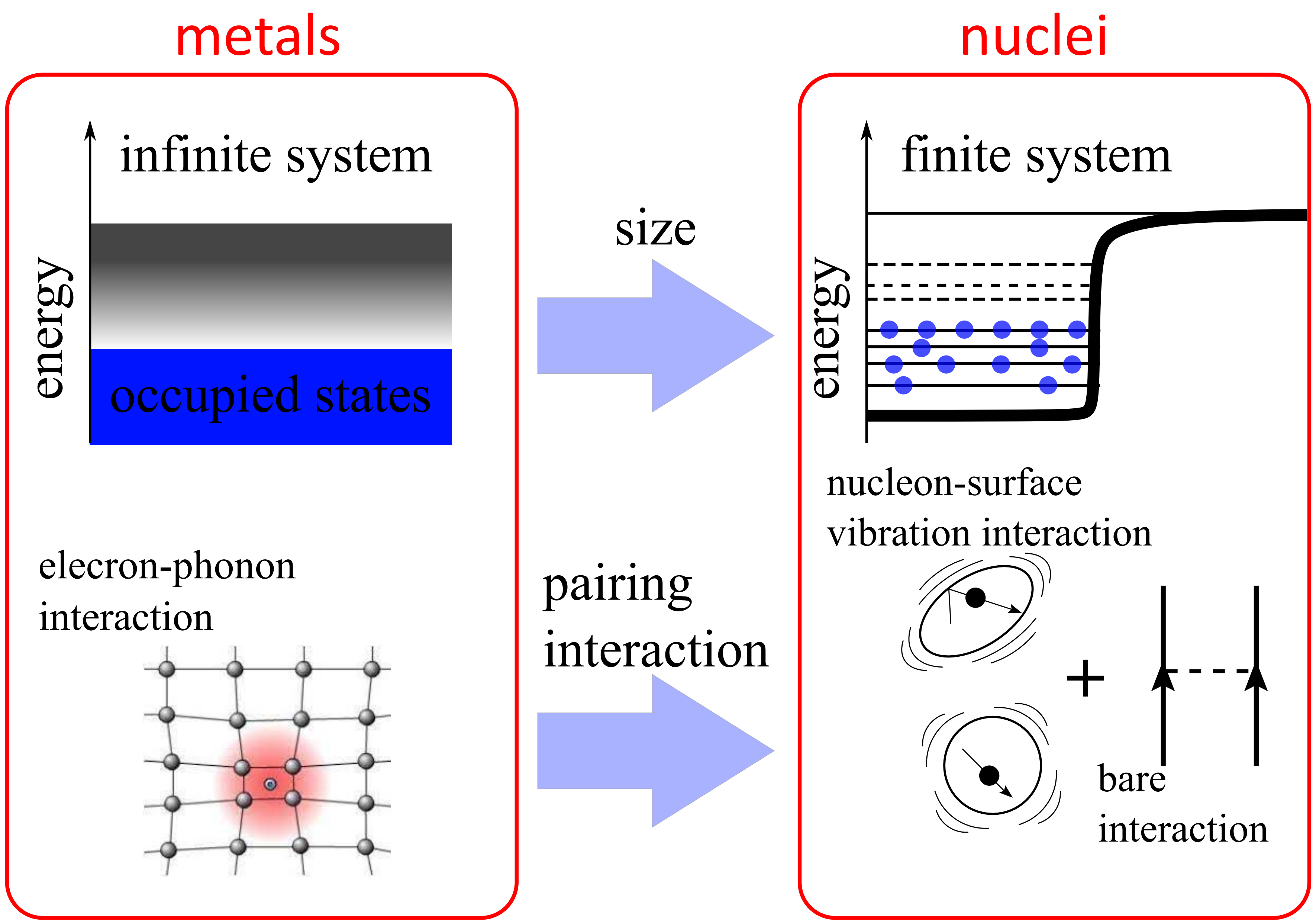
But, metals and nuclei are quite different, aren't they?



questions still arise

- Can we observe Cooper pairs in nuclei?
- How do we make a quantitative assessment of pair correlations in nuclei?
- How do we export our knowledge of nuclear superfluidity to nuclear matter?

But, metals and nuclei are quite different, aren't they?

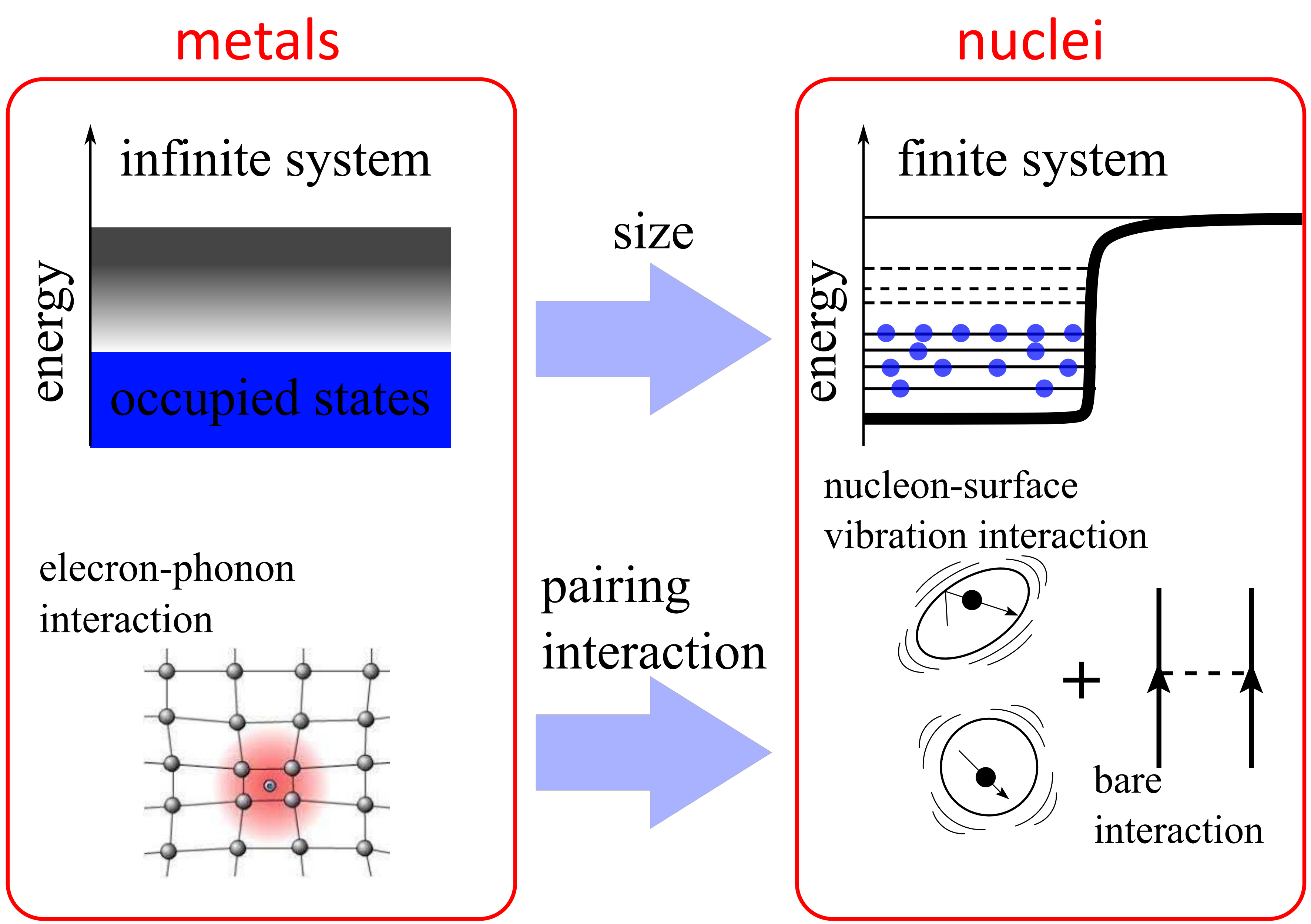


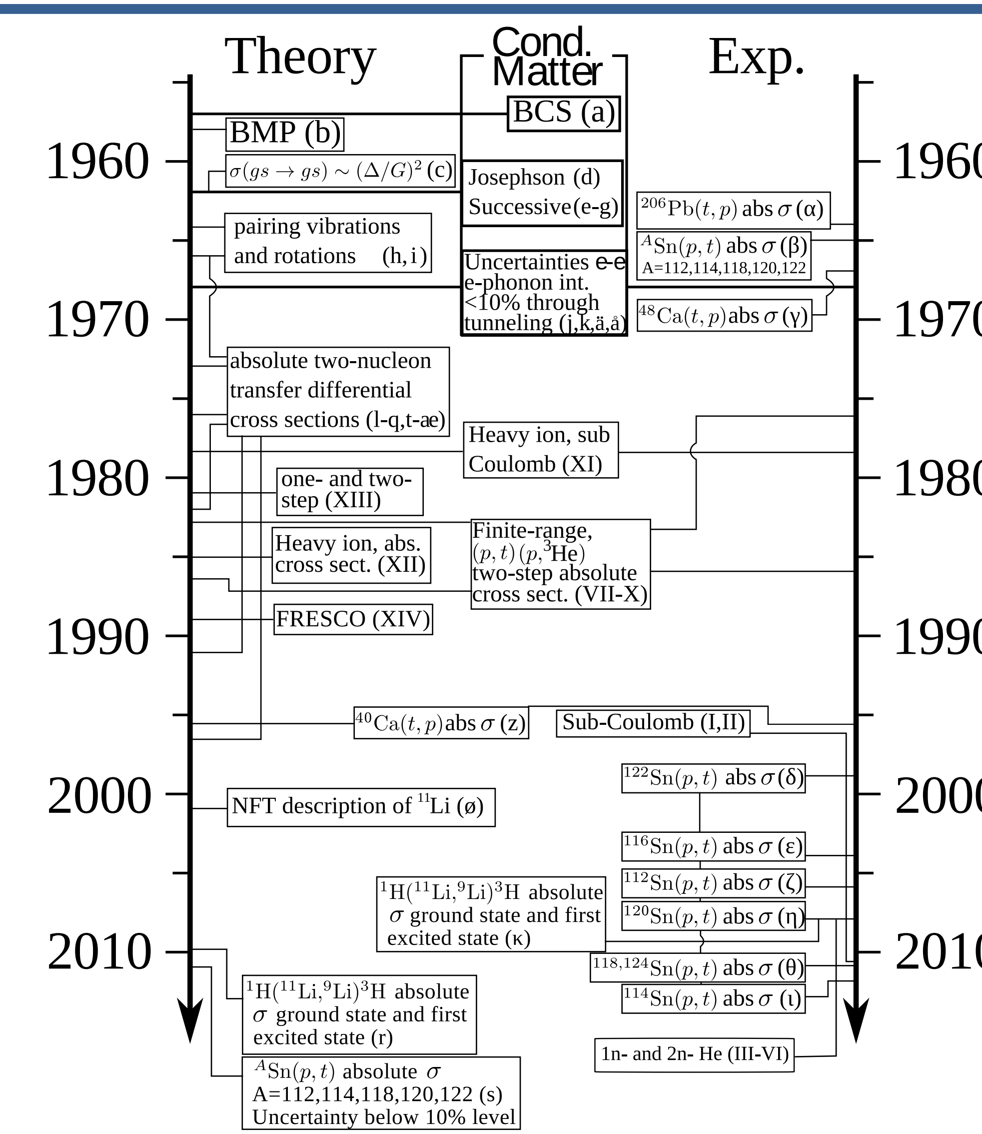
questions still arise

- Can we observe Cooper pairs in nuclei?
- How do we make a quantitative assessment of pair correlations in nuclei?
- How do we export our knowledge of nuclear superfluidity to nuclear matter?

(t,p) reactions are a specific probe of nuclear pairing correlations

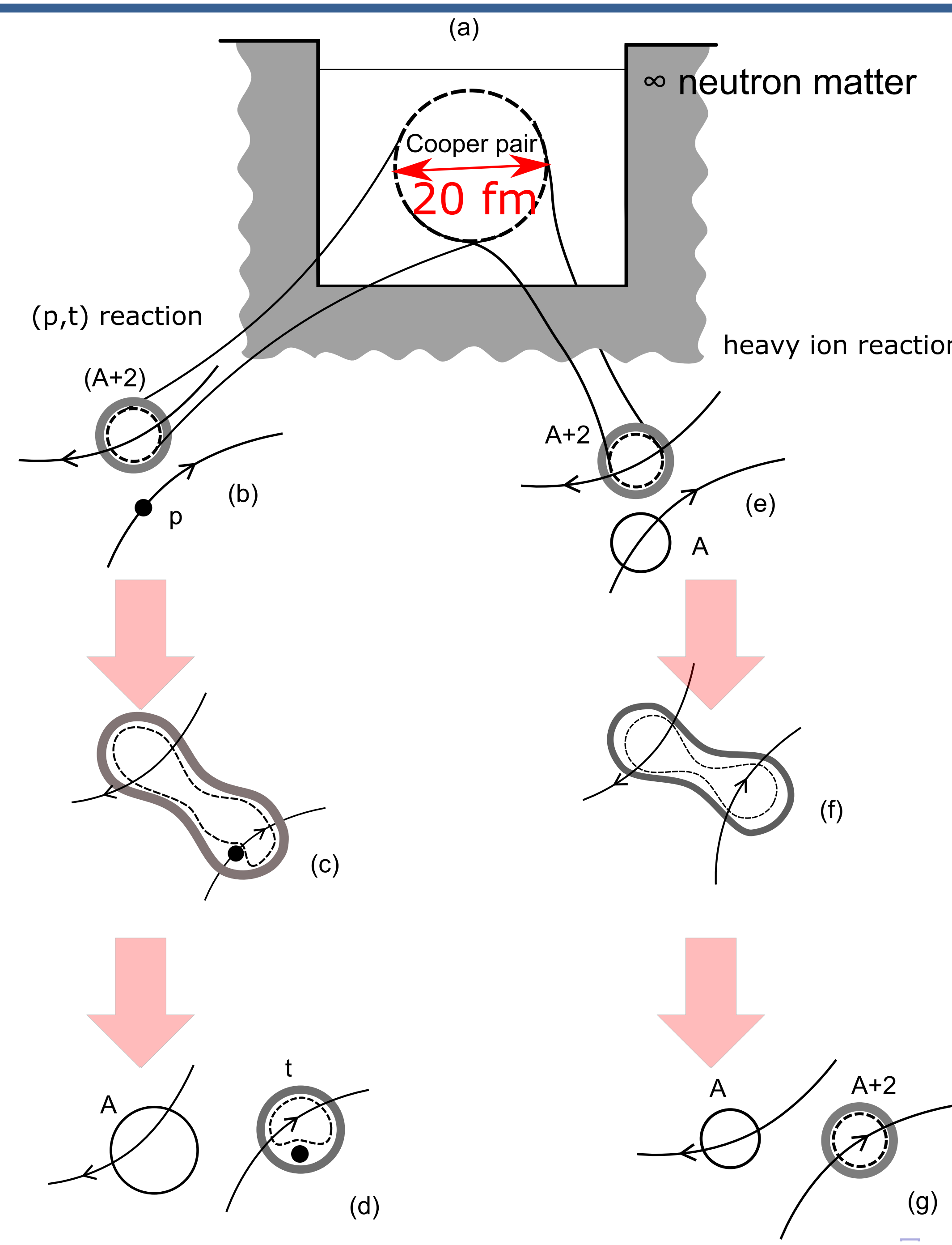
But, metals and nuclei are quite different, aren’t they?



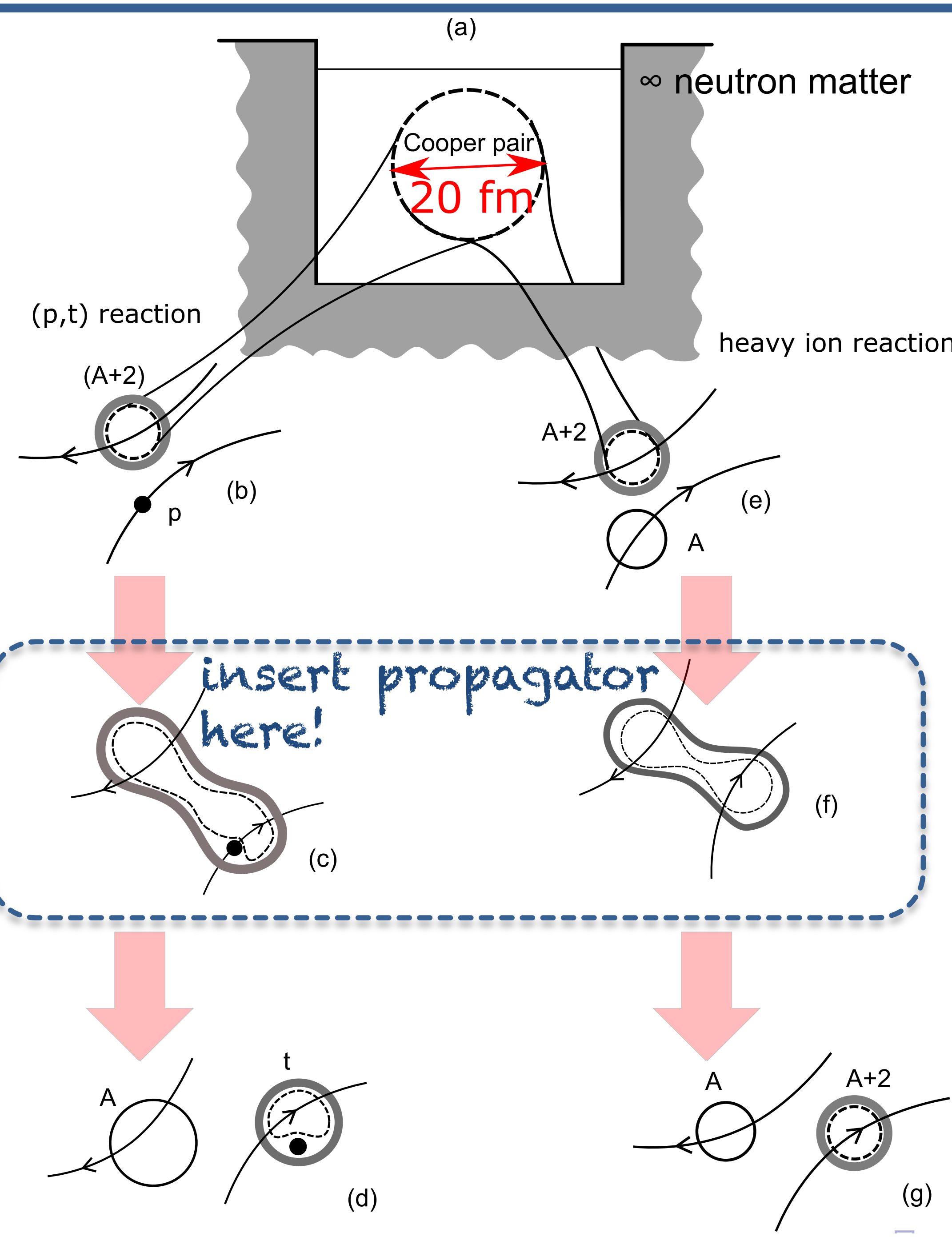


Rep. Prog. Phys. **76** (2013) 106301
GP, Idini, Barranco, Vigezzi, Broglia

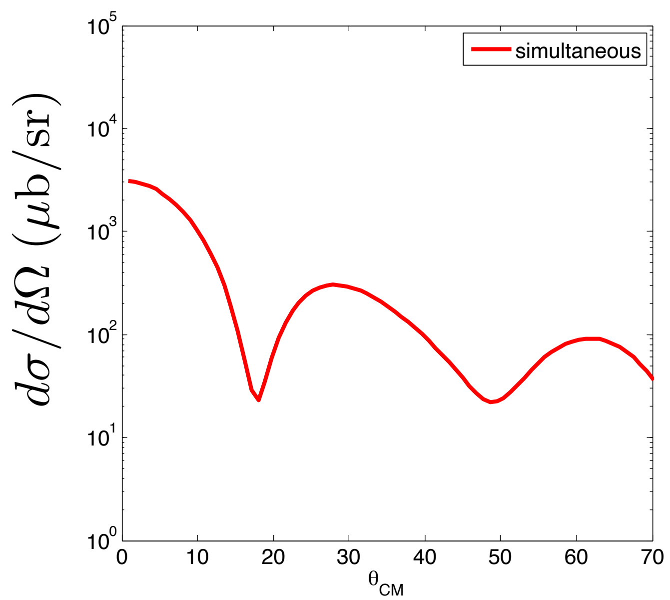
The 2 neutron transfer process is very delocalized



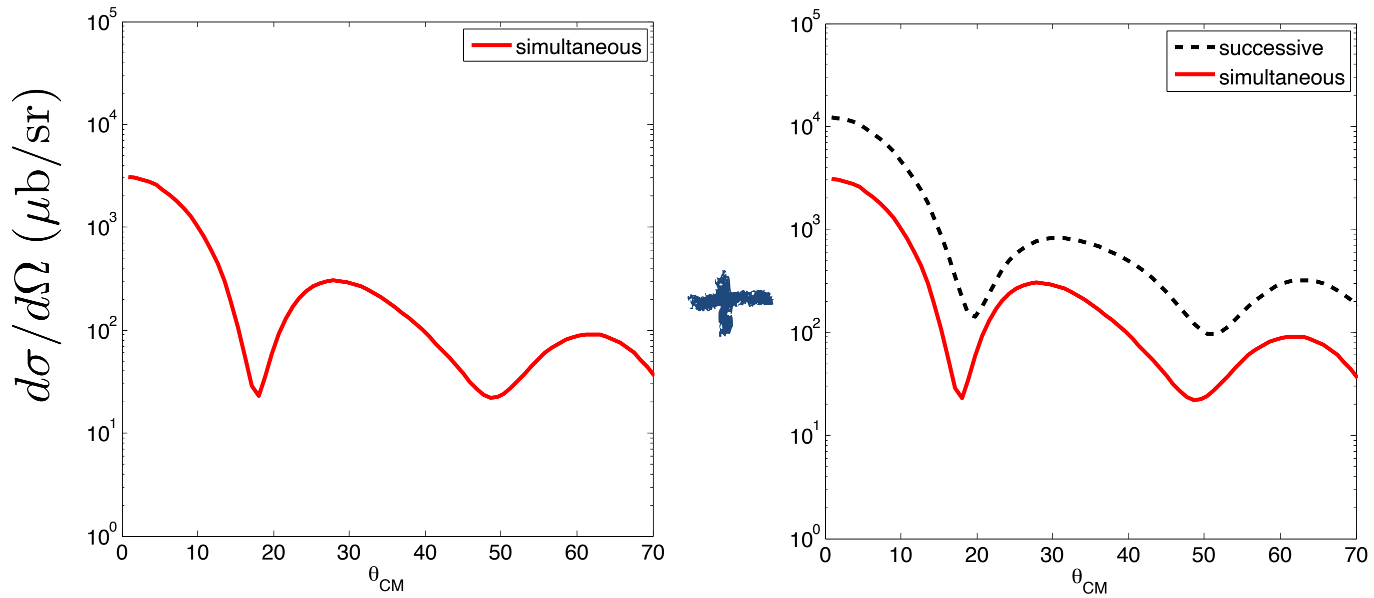
The 2 neutron transfer process is very delocalized



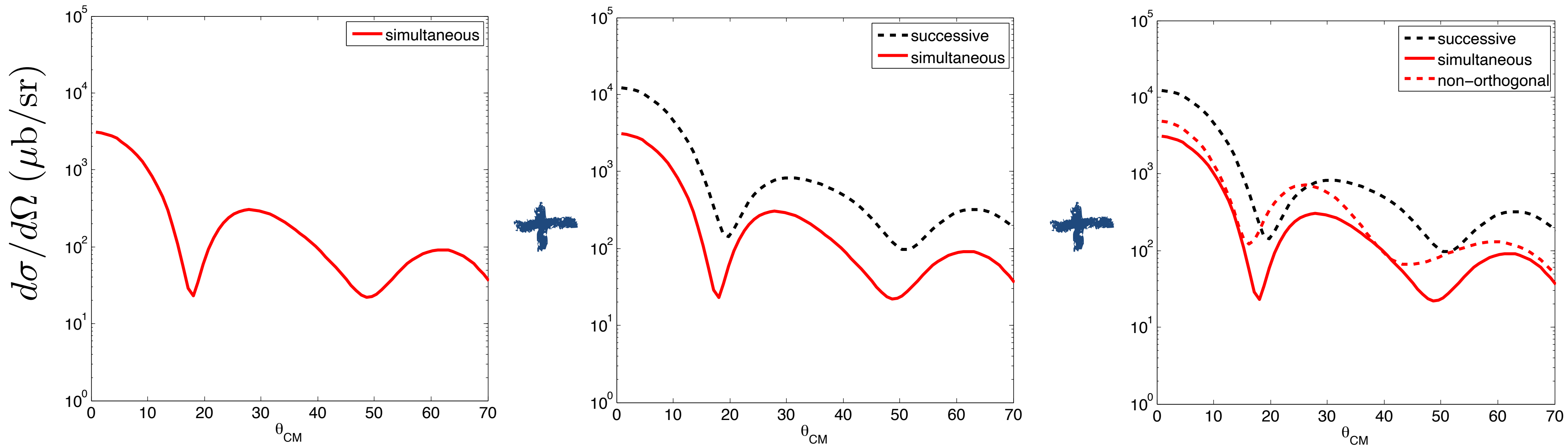
Computing $^{112}Sn(p,t)^{110}Sn$ in second order DWBA



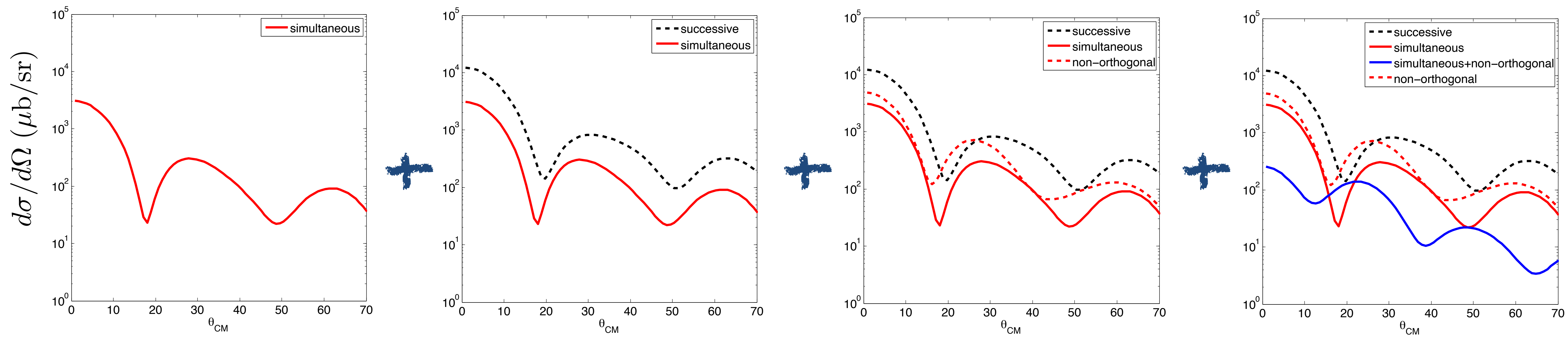
Computing $^{112}Sn(p,t)^{110}Sn$ in second order DWBA



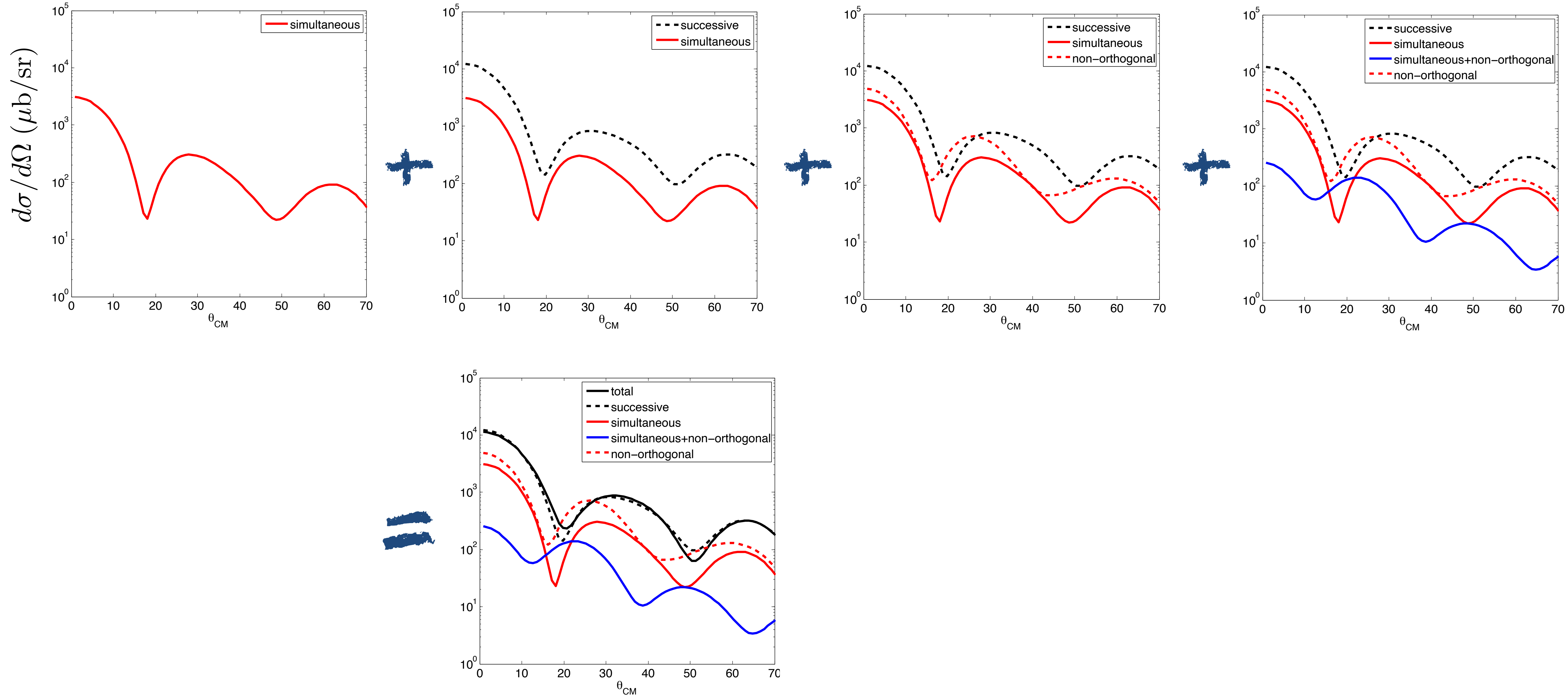
Computing $^{112}Sn(p,t)^{110}Sn$ in second order DWBA



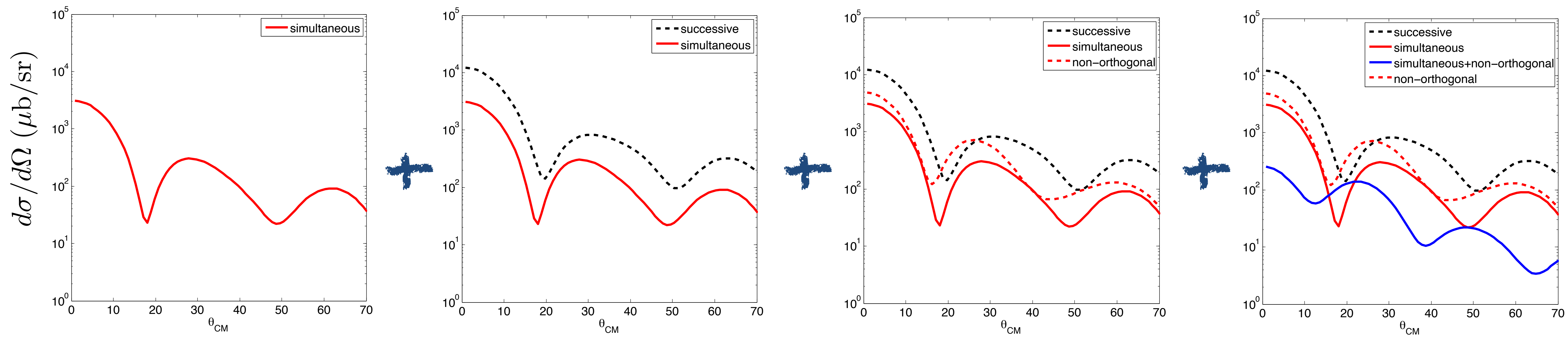
Computing $^{112}Sn(p,t)^{110}Sn$ in second order DWBA

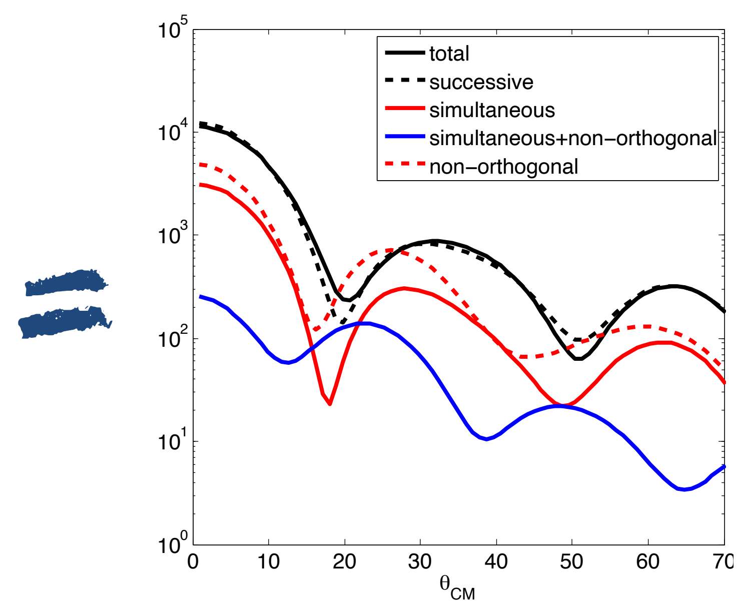


Computing ^{112}Sn(p,t)^{110}Sn in second order DWBA



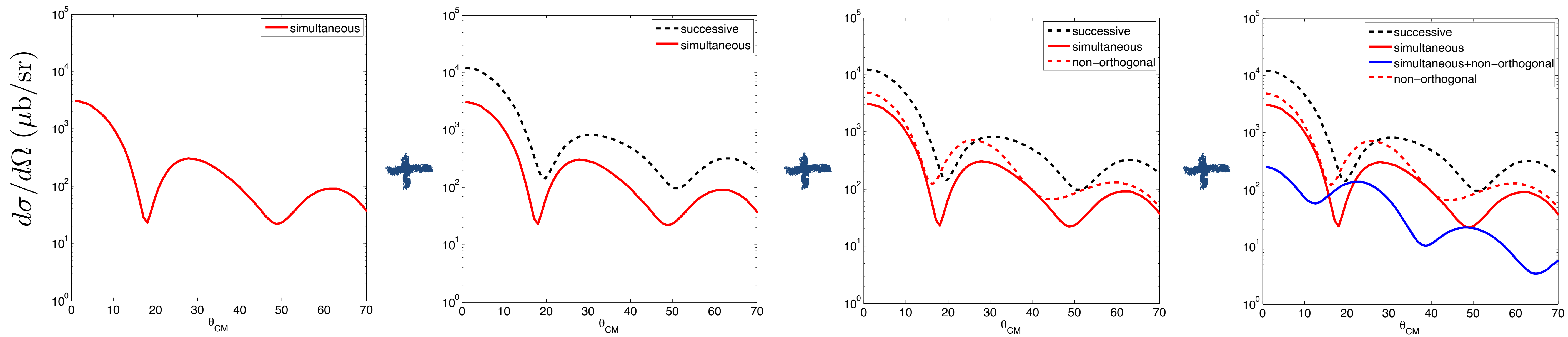
Computing $^{112}Sn(p,t)^{110}Sn$ in second order DWBA

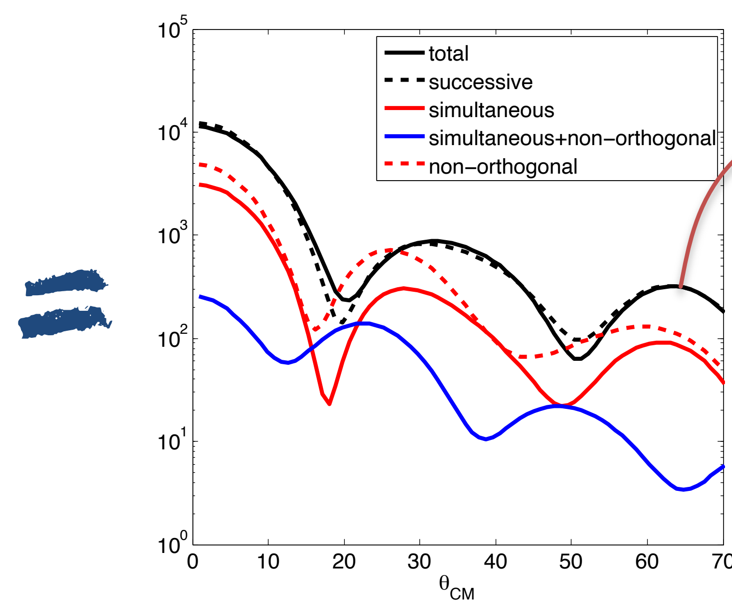




The (t,p) process is essentially a successive transfer of two correlated neutrons

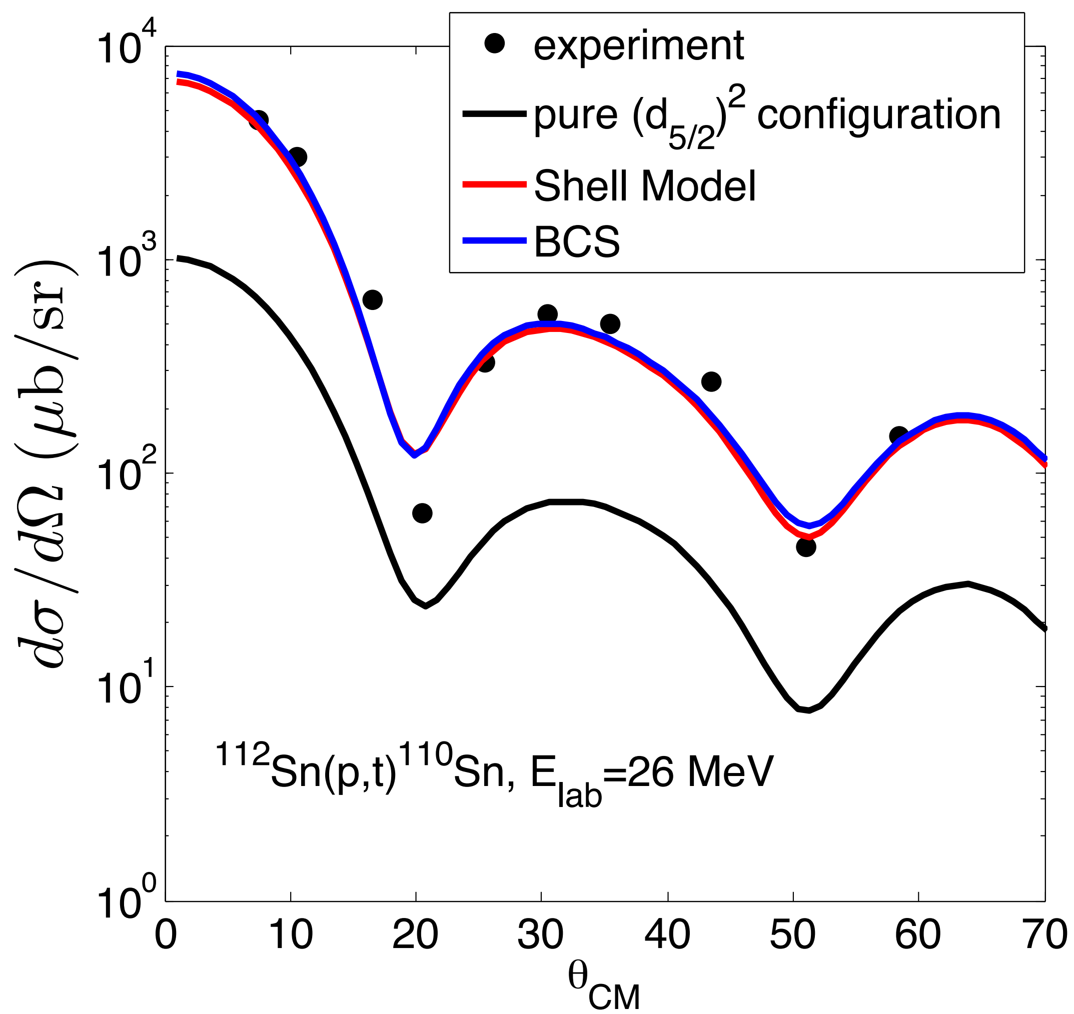
Computing ^{112}Sn(p,t)^{110}Sn in second order DWBA





The (t,p) process is essentially a successive transfer of two correlated neutrons

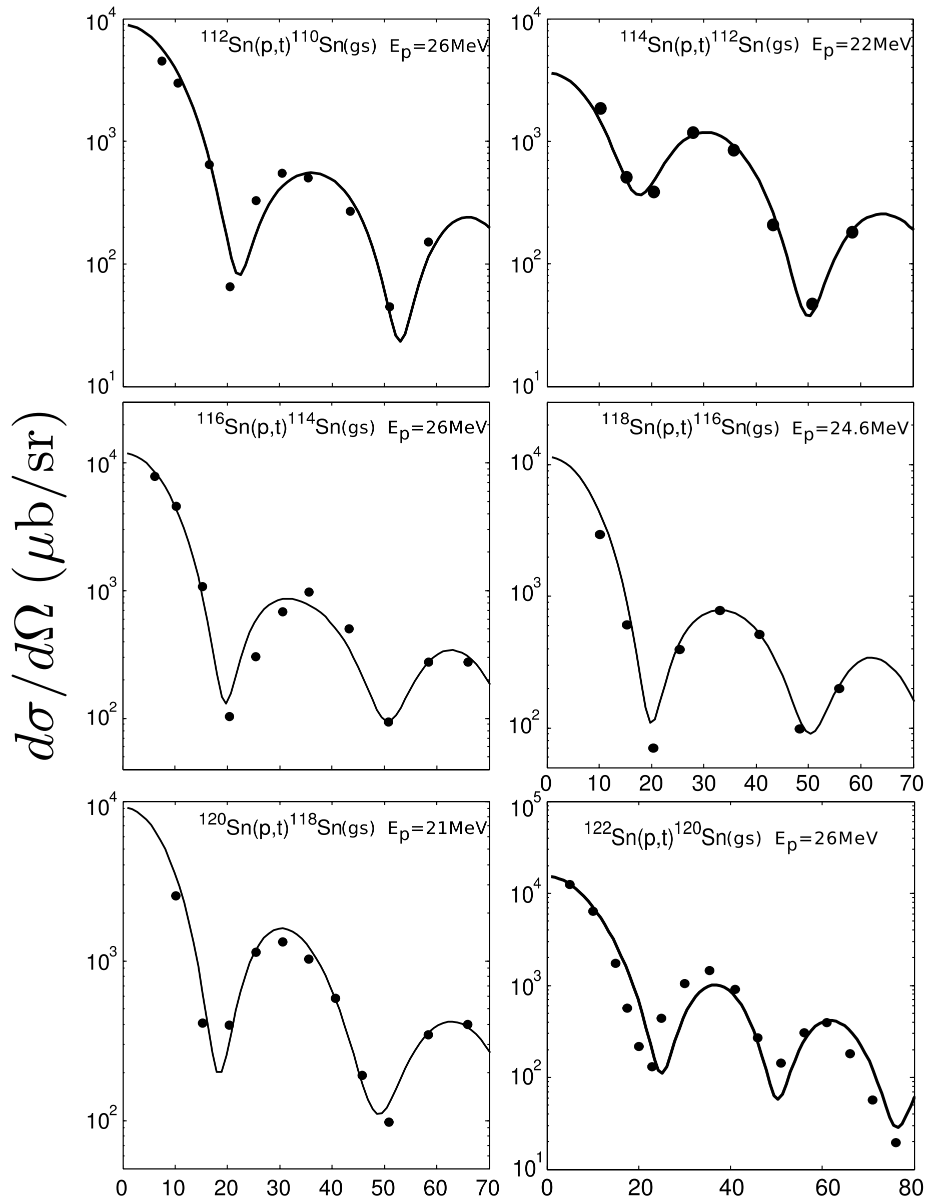
Theory should account for the *absolute value* of the cross section



enhancement factor with respect to the transfer of uncorrelated neutrons:
$\varepsilon = 20.6$

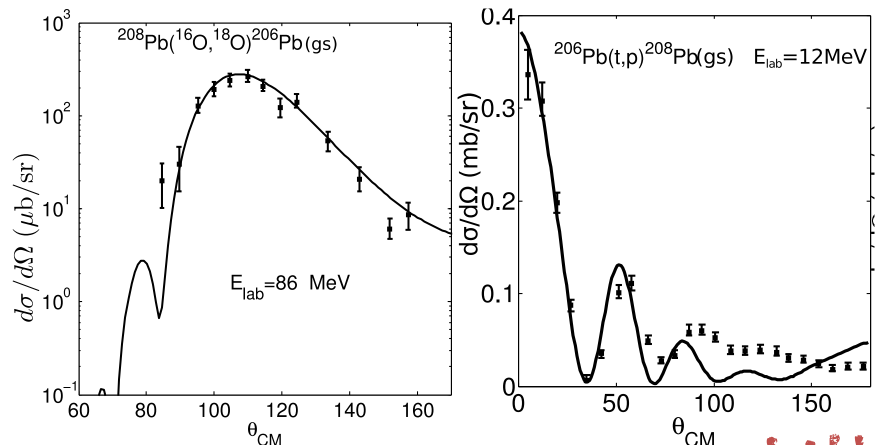
Experimental data and shell model wavefunction from Guazzoni *et al.* PRC **74** 054605 (2006)

Reaction+Structure theory works well across the nuclear chart



Sn isotopes (BCS)

Pb ground state and excited states (QRPA)

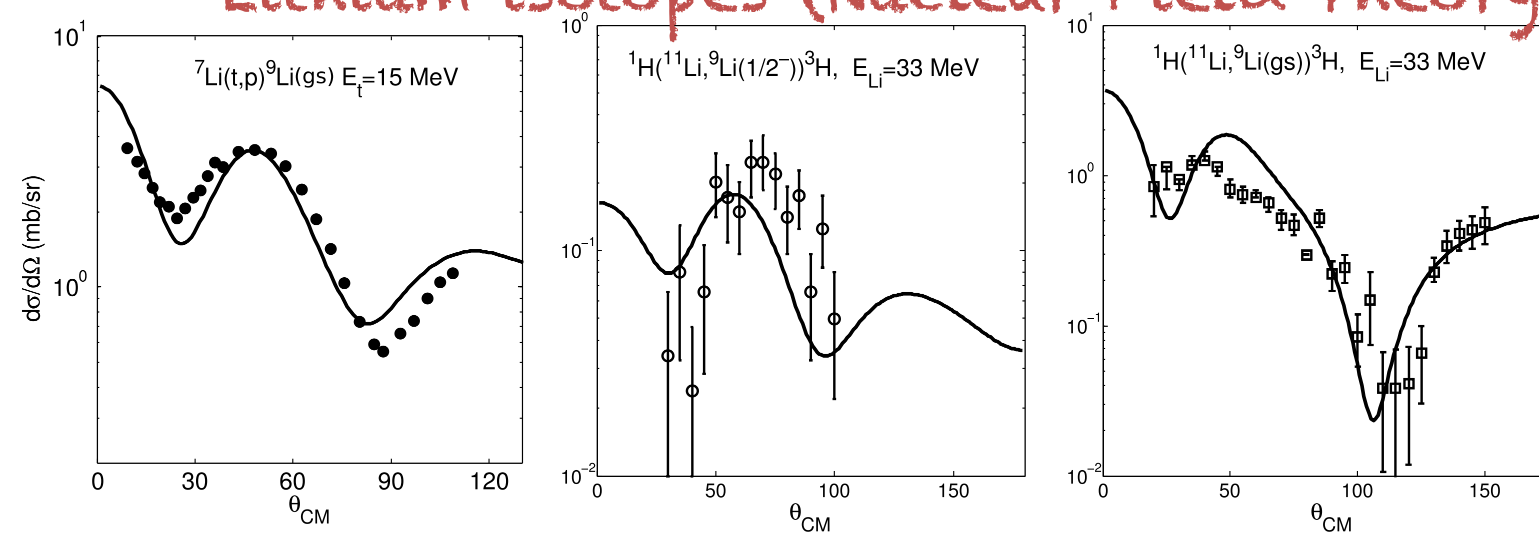


Rep. Prog. Phys. **76** (2013) 106301 (21pp)

Cooper pair transfer in nuclei

G Potel[1], A Idini[2,3,4], F Barranco[5], E Vigezzi[4] and R A Broglia[3,4,6,7]

Lithium isotopes (Nuclear Field Theory)



Looking for something new in the nuclear spectrum: The Giant Pairing Vibration (GPV)

Volume 69B, number 2 PHYSICS LETTERS 1 August 1977

HIGH-LYING PAIRING RESONANCES★

R.A. BROGLIA

The Niels Bohr Institute, University of Copenhagen, DK-2100 Copenhagen Ø, Denmark[1]

State University of New York, Department of Physics, Stony Brook, New York 11794, USA

and

D.R. BES[2]

NORDITA, DK-2100 Copenhagen Ø, Denmark

Pairing vibrations based on the excitation of pairs of particles and holes across major shells are predicted at an excitation energy of about $70/A^{1/3}$ MeV and carrying a cross section which is 20%–100% the ground state cross section.

Collective pairing mode predicted almost 50 years ago, awaiting experimental confirmation?

(t,p) is an ideal process to populate the elusive Giant Pairing Vibration

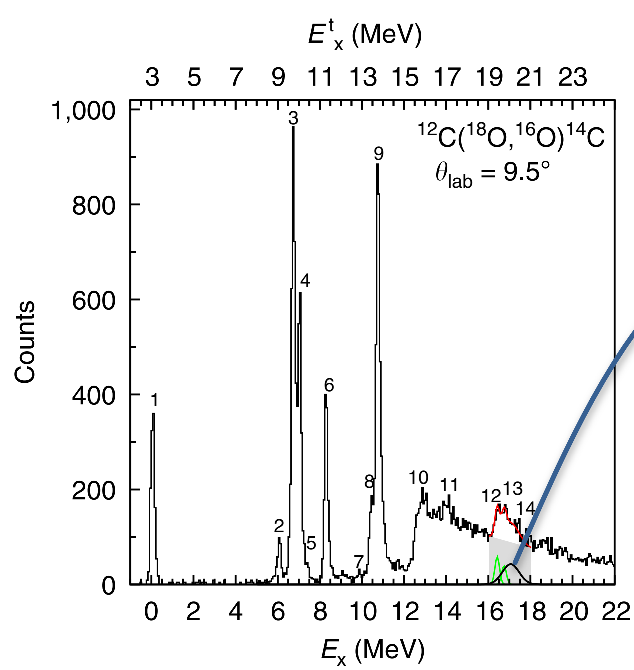
nature COMMUNICATIONS

ARTICLE

Received 28 Dec 2014 | Accepted 24 Feb 2015 | Published 27 Mar 2015 DOI: 10.1038/ncomms7743 OPEN

Signatures of the Giant Pairing Vibration in the ^{14}C and ^{15}C atomic nuclei

F. Cappuzzello[1,2], D. Carbone[2], M. Cavallaro[2], M. Bondi[1,2], C. Agodi[2], F. Azaiez[3], A. Bonaccorso[4], A. Cunsolo[2], L. Fortunato[5,6], A. Foti[1,7], S. Franchoo[3], E. Khan[3], R. Linares[8], J. Lubian[8], J.A. Scarpaci[9] & A. Vitturi[5,6]



bump in the continuum populated by the reaction $^{12}C(^{18}O,^{16}O)^{14}C$

(t,p) is an ideal process to populate the elusive Giant Pairing Vibration

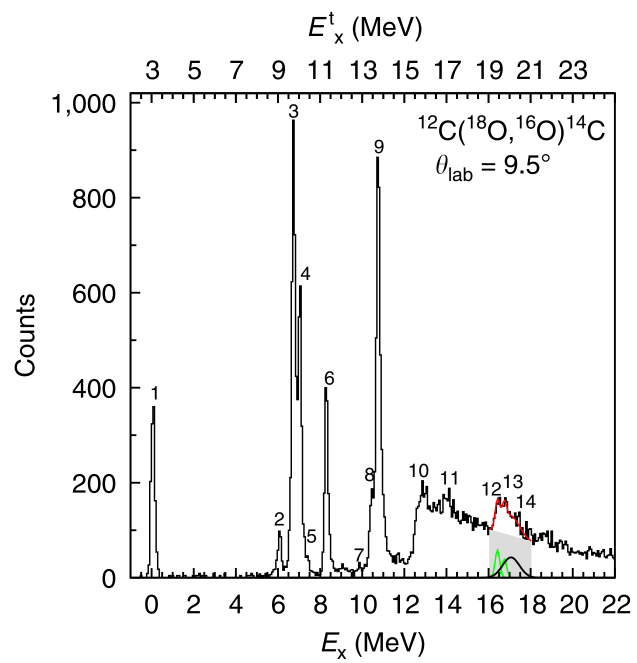
nature COMMUNICATIONS

ARTICLE

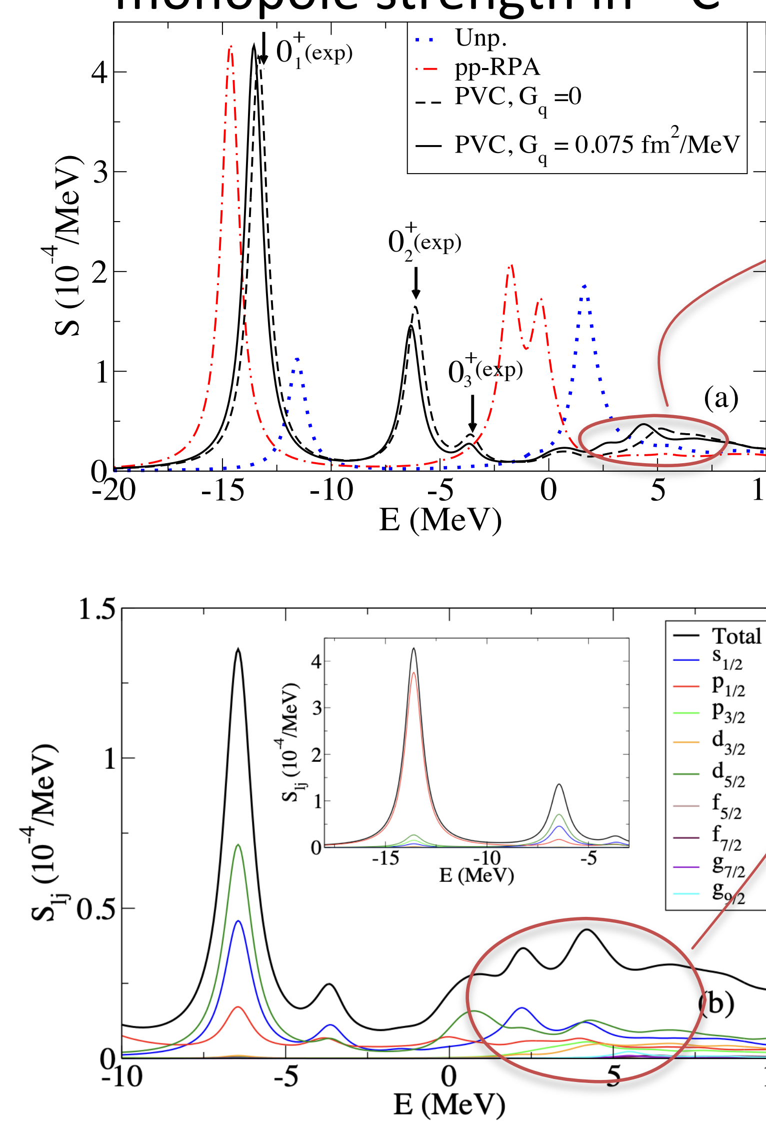
Received 28 Dec 2014 | Accepted 24 Feb 2015 | Published 27 Mar 2015 DOI: 10.1038/ncomms7743 OPEN

Signatures of the Giant Pairing Vibration in the ^{14}C and ^{15}C atomic nuclei

F. Cappuzzello[1,2], D. Carbone[2], M. Cavallaro[2], M. Bondi[1,2], C. Agodi[2], F. Azaiez[3], A. Bonaccorso[4], A. Cunsolo[2], L. Fortunato[5,6], A. Foti[1,7], S. Franchoo[3], E. Khan[3], R. Linares[8], J. Lubian[8], J.A. Scarpaci[9] & A. Vitturi[5,6]



monopole strength in ^{14}C



we predict a rather broad structure in the continuum

PHYSICAL REVIEW LETTERS **134,** 062501 (2025)

Fragmentation of the Giant Pairing Vibration in ^{14}C Induced by Many-Body Processes

F. Barranco,[1] G. Potel,[1,2] and E. Vigezzi[3]

(t,p) is an ideal process to populate the elusive Giant Pairing Vibration

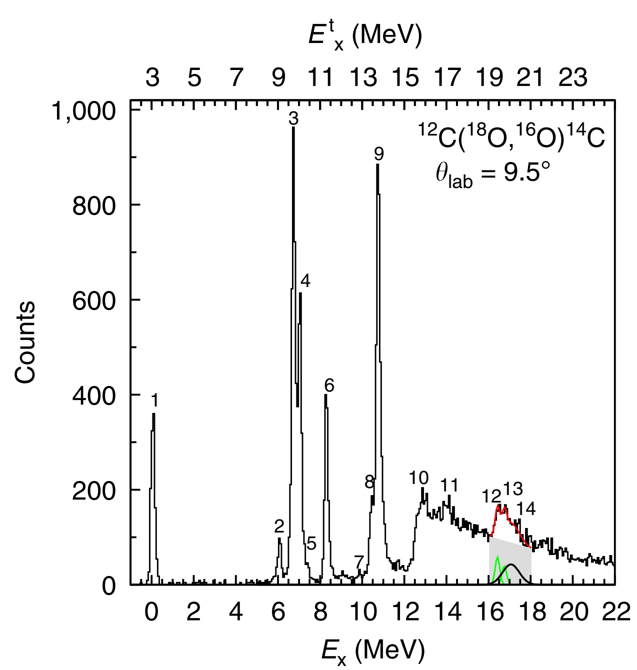
nature COMMUNICATIONS

ARTICLE

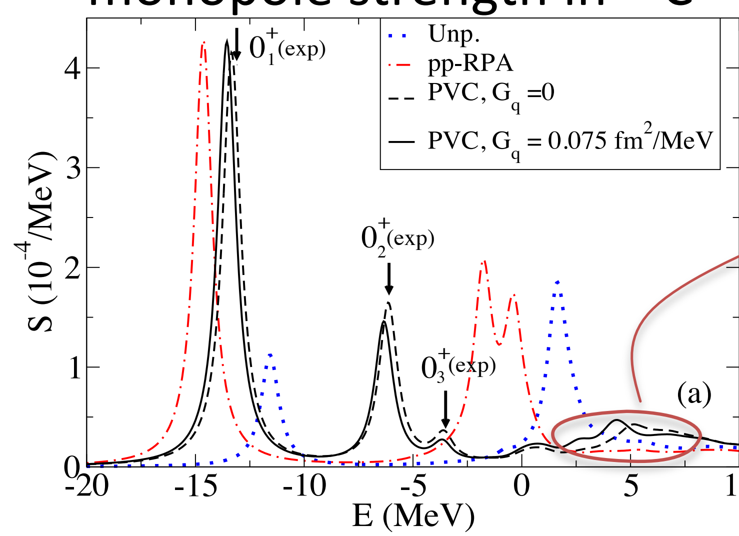
Received 28 Dec 2014 | Accepted 24 Feb 2015 | Published 27 Mar 2015 DOI: 10.1038/ncomms7743 OPEN

Signatures of the Giant Pairing Vibration in the ^{14}C and ^{15}C atomic nuclei

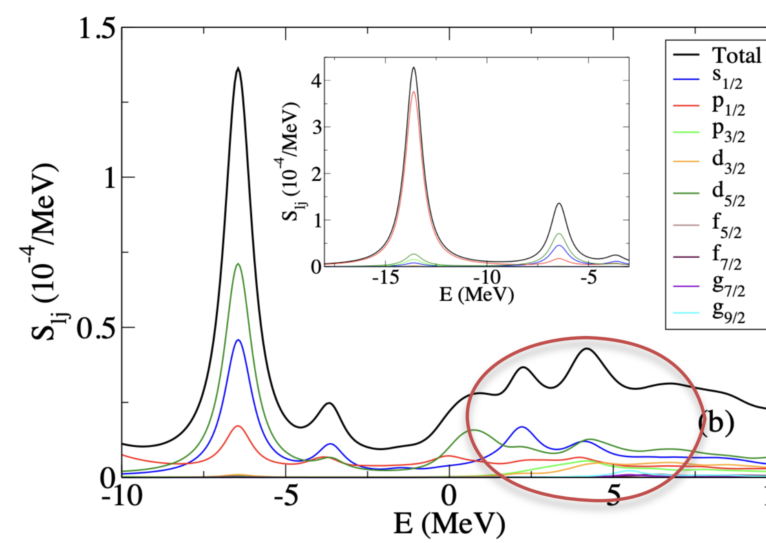
F. Cappuzzello[1,2], D. Carbone[2], M. Cavallaro[2], M. Bondi[1,2], C. Agodi[2], F. Azaiez[3], A. Bonaccorso[4], A. Cunsolo[2], L. Fortunato[5,6], A. Foti[1,7], S. Franchoo[3], E. Khan[3], R. Linares[8], J. Lubian[8], J.A. Scarpaci[9] & A. Vitturi[5,6]



monopole strength in ^{14}C



we predict a rather broad structure in the continuum



working on a theoretical estimate of ^{12}C(t,p)^{14}C(GPV)

PHYSICAL REVIEW LETTERS **134,** 062501 (2025)

Fragmentation of the Giant Pairing Vibration in ^{14}C Induced by Many-Body Processes

F. Barranco,[1] G. Potel,[1,2] and E. Vigezzi[3]

The Pygmy Dipole Resonance (PDR) as a two-quasiparticle mode

- The PDR is rather well described in the harmonic approximation (RPA, QRPA) as a two-quasiparticle mode.
- Therefore, PDR in a nucleus A_0 can be better probed with two-quasiparticle fields, i.e., particle-hole (*ph*), particle-particle (*pp*), and hole-hole (*hh*) fields.

ph

Coulomb, inelastic, and γ-induced excitation on A_0:

- $A_0(d,d')A_0(PDR)$
- $A_0(p,p')$ $A_0(PDR)$
- $A_0(\alpha,\alpha')$ $A_0(PDR)$
- $A_0(\gamma,\gamma')$ $A_0(PDR)$
- $A_0(n,n')$ $A_0(PDR)$
- $A_0(X,X')$ $A_0(PDR)$

one-nucleon transfer on A_0-1:

- A_0-1(d,p) A_0(PDR)

Spieker et al., PRL (2020)

Weinert et al., PRL (2021)

pp

two-nucleon transfer on A_0-2:

- A_0-2(t,p) A_0(PDR)

proposed in this talk

The Pygmy Dipole Resonance (PDR) as a two-quasiparticle mode

- The PDR is rather well described in the harmonic approximation (RPA, QRPA) as a two-quasiparticle mode.
- Therefore, PDR in a nucleus A_0 can be better probed with two-quasiparticle fields, i.e., particle-hole (*ph*), particle-particle (*pp*), and hole-hole (*hh*) fields.

complementary classification of dipole modes

isovector $\Longleftrightarrow$ isoscalar

ph $\Longleftrightarrow$ pp

The Pygmy Dipole Resonance (PDR) as a two-quasiparticle mode

- The PDR is rather well described in the harmonic approximation (RPA, QRPA) as a two-quasiparticle mode.
- Therefore, PDR in a nucleus A_0 can be better probed with two-quasiparticle fields, i.e., particle-hole (*ph*), particle-particle (*pp*), and hole-hole (*hh*) fields.

complementary classification of dipole modes

isovector $\Longleftrightarrow$ isoscalar

ph $\Longleftrightarrow$ pp

GDR $\Longleftrightarrow$ PDR?

Probing the ^{11}Li PDR with 2-neutron transfer

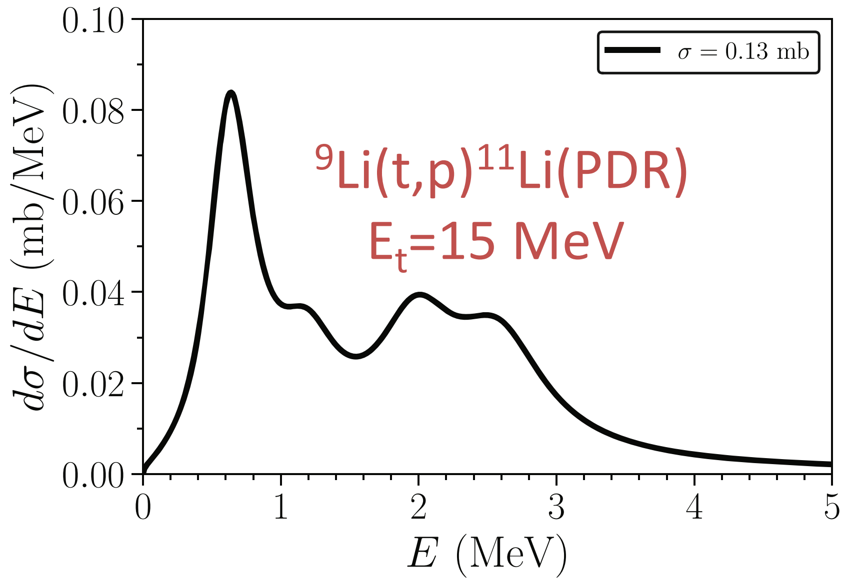
Eur. Phys. J. A (2019) **55**: 243
DOI 10.1140/epja/i2019-12789-y

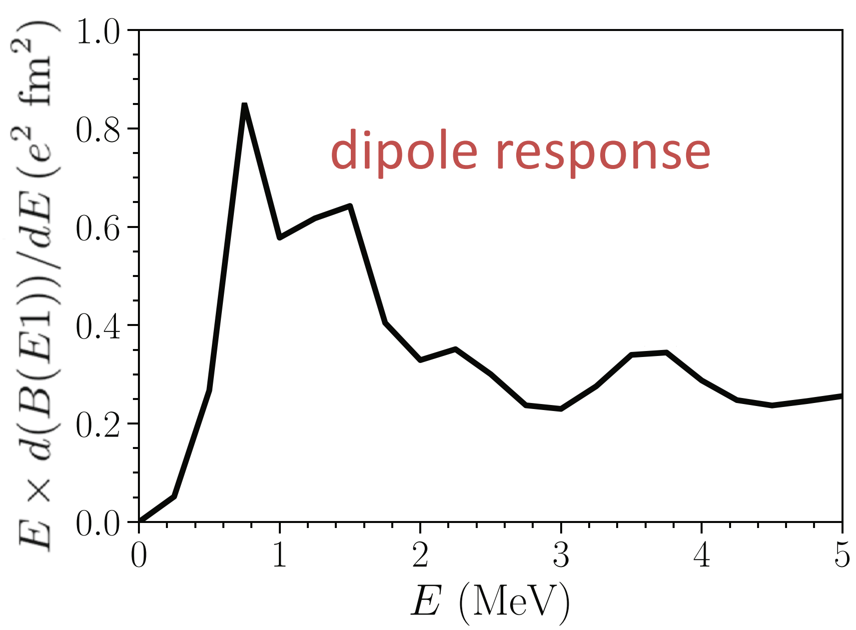
THE EUROPEAN PHYSICAL JOURNAL A

Regular Article – Theoretical Physics

Characterization of vorticity in pygmy resonances and soft-dipole modes with two-nucleon transfer reactions*

R.A. Broglia[1,2], F. Barranco[3], G. Potel[4,a], and E. Vigezzi[5]





EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

Probing the ^{11}Li low-lying dipole strength via ^{9}Li(t,p) with the ISS

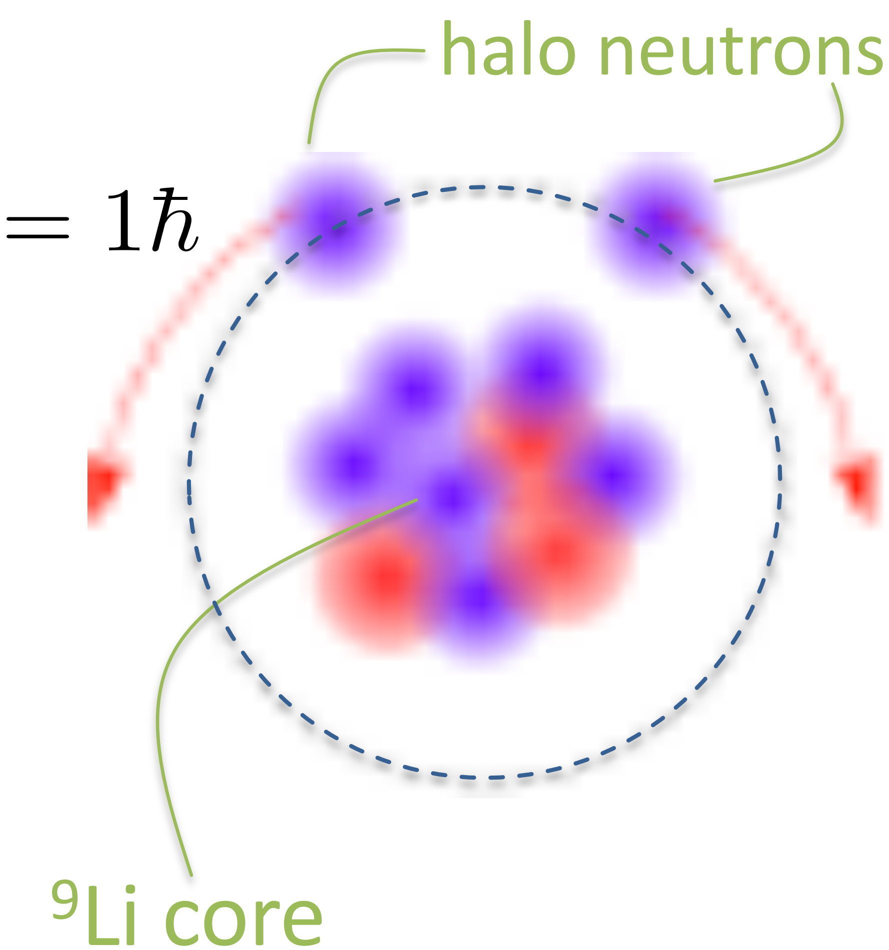
Y. Ayyad[1], E. Vigezzi[2], G. Potel[3], R. Broglia[4,5], B.P. Kay[6], A.O. Macchiavelli[7], H. Alvarez-Pol[8], F. Barranco[9], D. Bazin[1,10], M. Caamaño[8], A. Ceulemans[11], J. Chen[1], H.L. Crawford[7], B. Fernández-Domínguez[8], S.J. Freeman[12], L.P. Gaffney[13], C.R. Hoffman[6], R. Kanungo[14], C. Morse[7], O .Poleshchuk[11], R. Raabe[11], C.A. Santamaria[7], D.K. Sharp[12], T. L. Tang[6], K .Wimmer[15], A.H. Wuosmaa[16]

experiment approved at ISOLDE facility (CERN). Spokepersons: Ayyad, Vigezzi

the ^{11}Li PDR has the structure of an elementary quantum vortex

structure of a multipolar (1$^-$) Cooper pair:
elementary quantum vortex

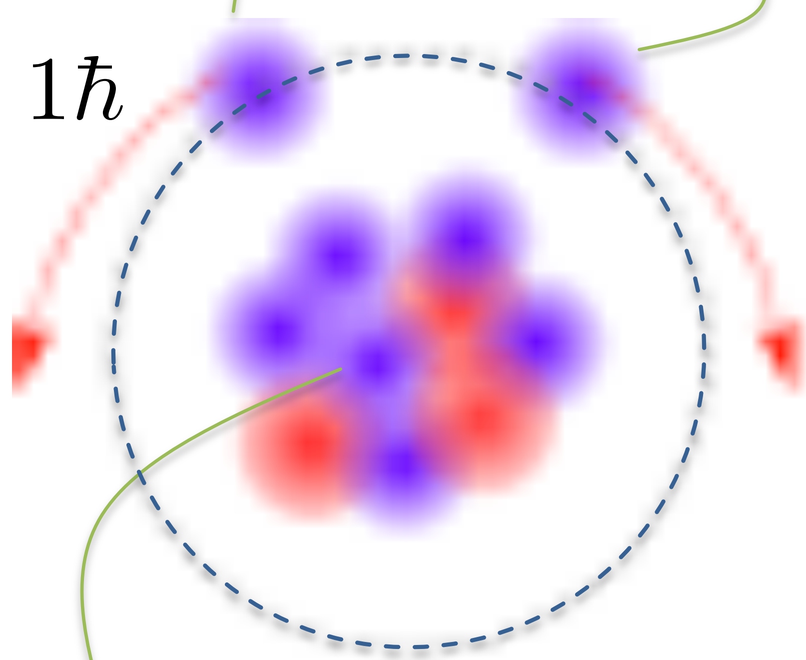
$L = 1\hbar$



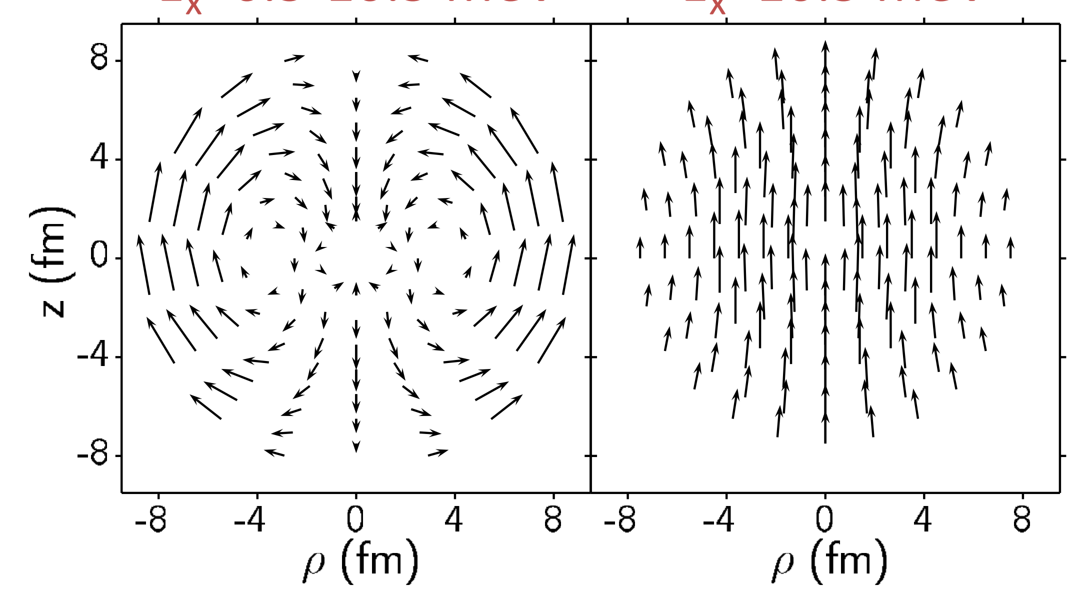
the ^{11}Li PDR has the structure of an elementary quantum vortex

structure of a multipolar (1^-) Cooper pair:
elementary quantum vortex

$$L = 1\hbar$$



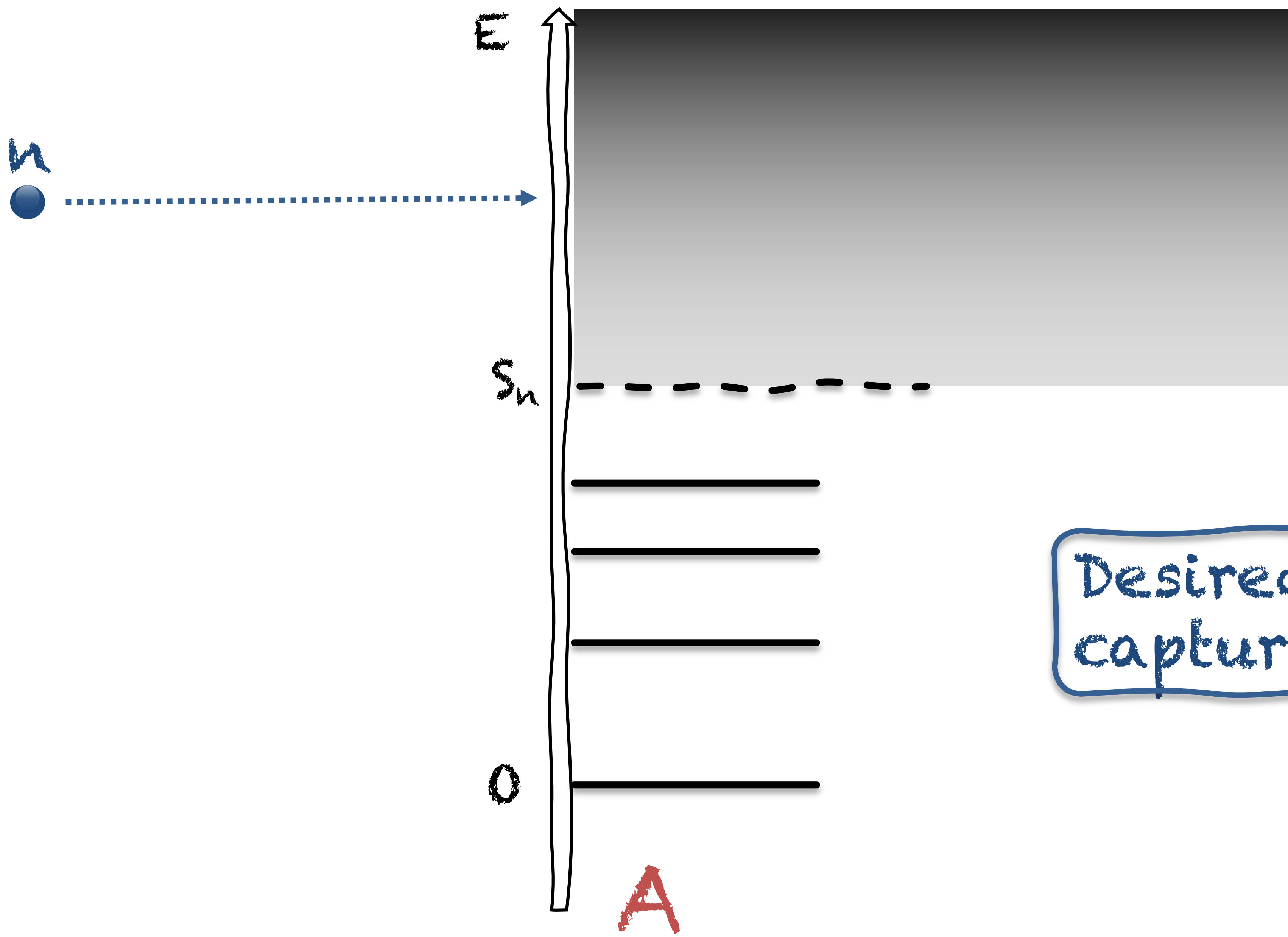
velocity field of ^{208}Pb dipole states



Ryezayeva *et al.* PRL **89** (2002) 272502

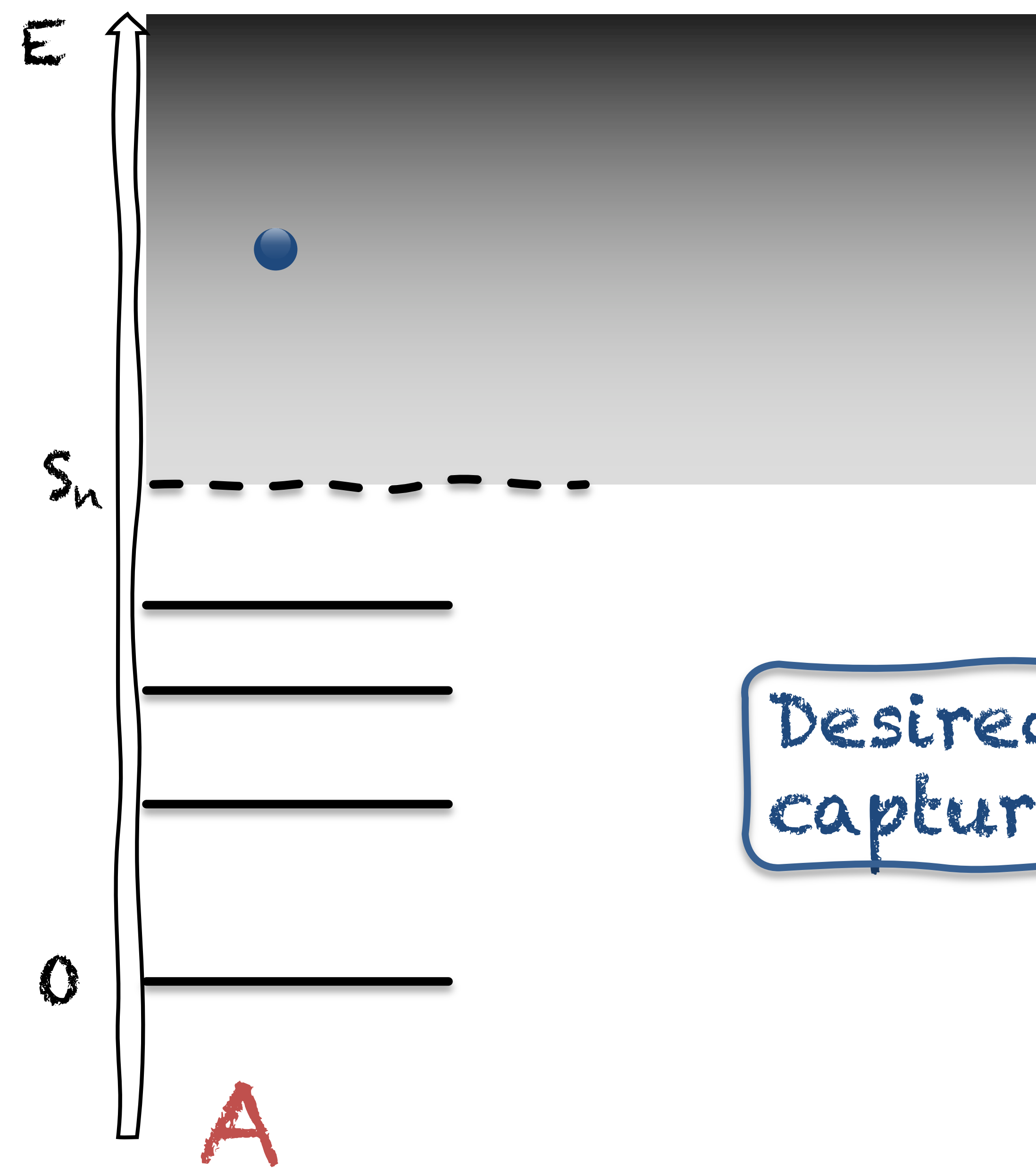
- Is vorticity a signature of PDR?
- Is there an experimental signature for it?

Using the Surrogate Reaction Method (SRM) to infer $^{A}X(n,\gamma)^{A+1}X$ from $^{A}X(d,p\gamma)^{A+1}X$



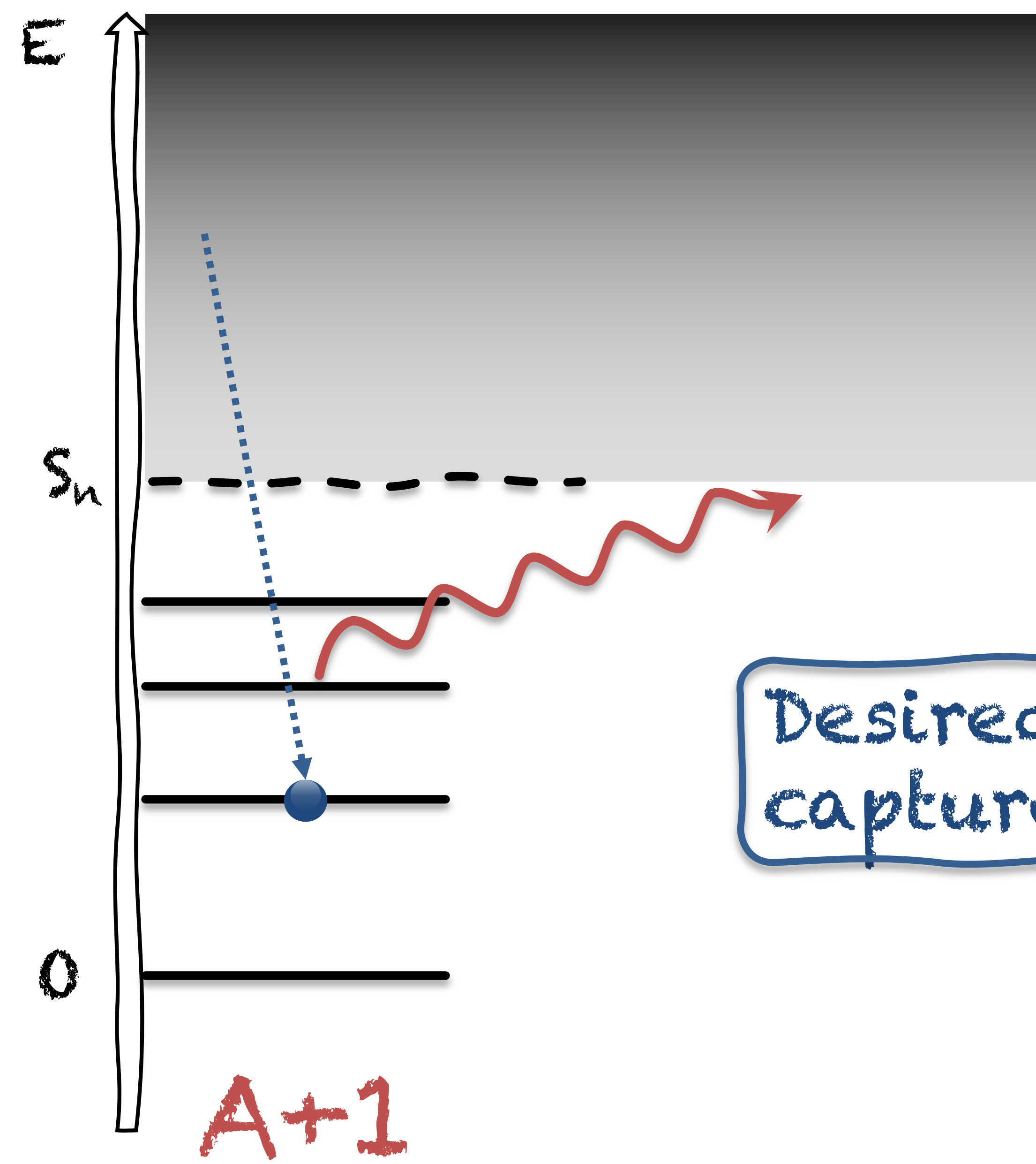
Desired reaction: neutron capture on target A

Using the Surrogate Reaction Method (SRM) to infer $^{A}X(n,\gamma)^{A+1}X$ from $^{A}X(d,p\gamma)^{A+1}X$



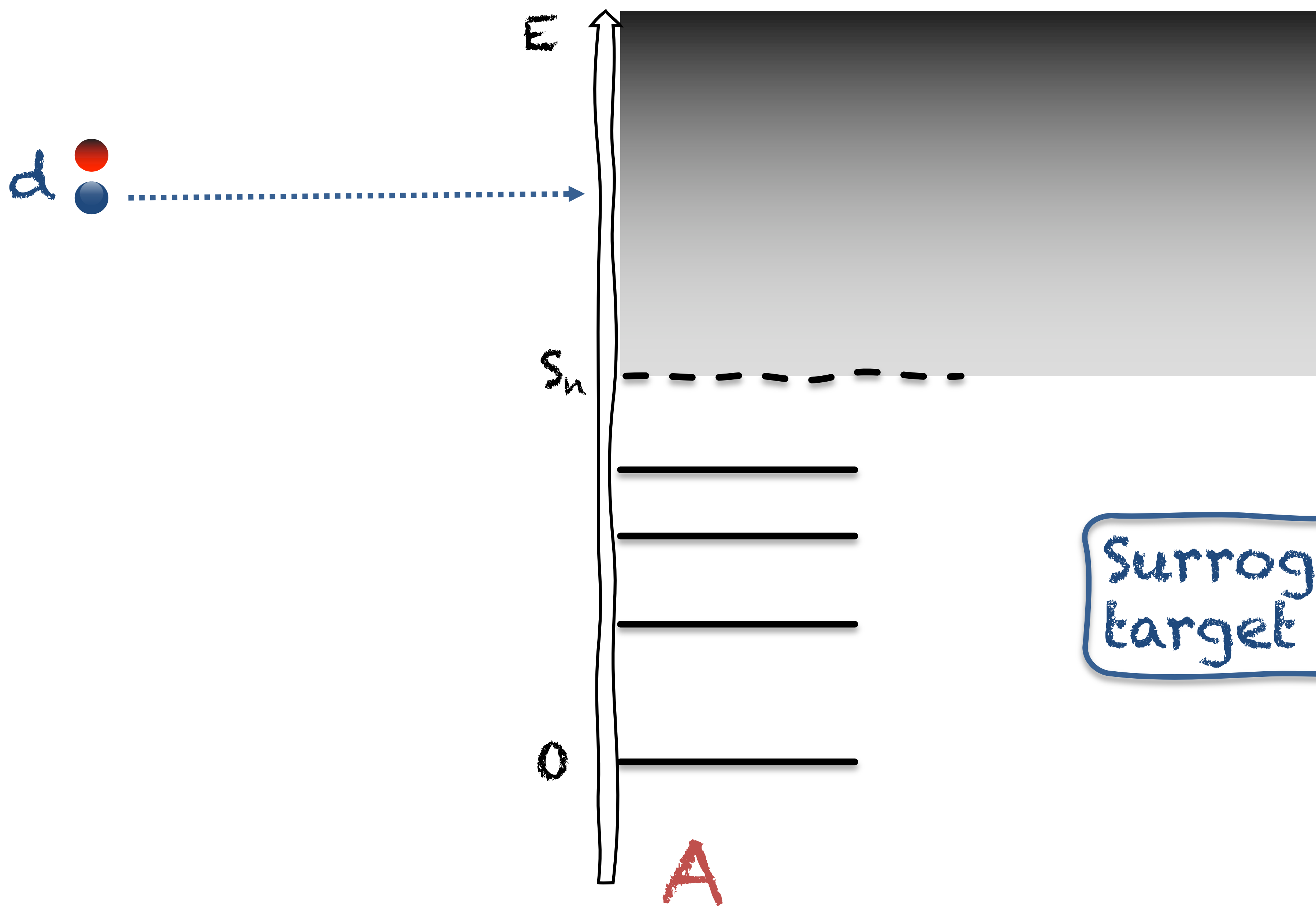
Desired reaction: neutron capture on target A

Using the Surrogate Reaction Method (SRM) to infer $^{A}X(n,\gamma)^{A+1}X$ from $^{A}X(d,p\gamma)^{A+1}X$



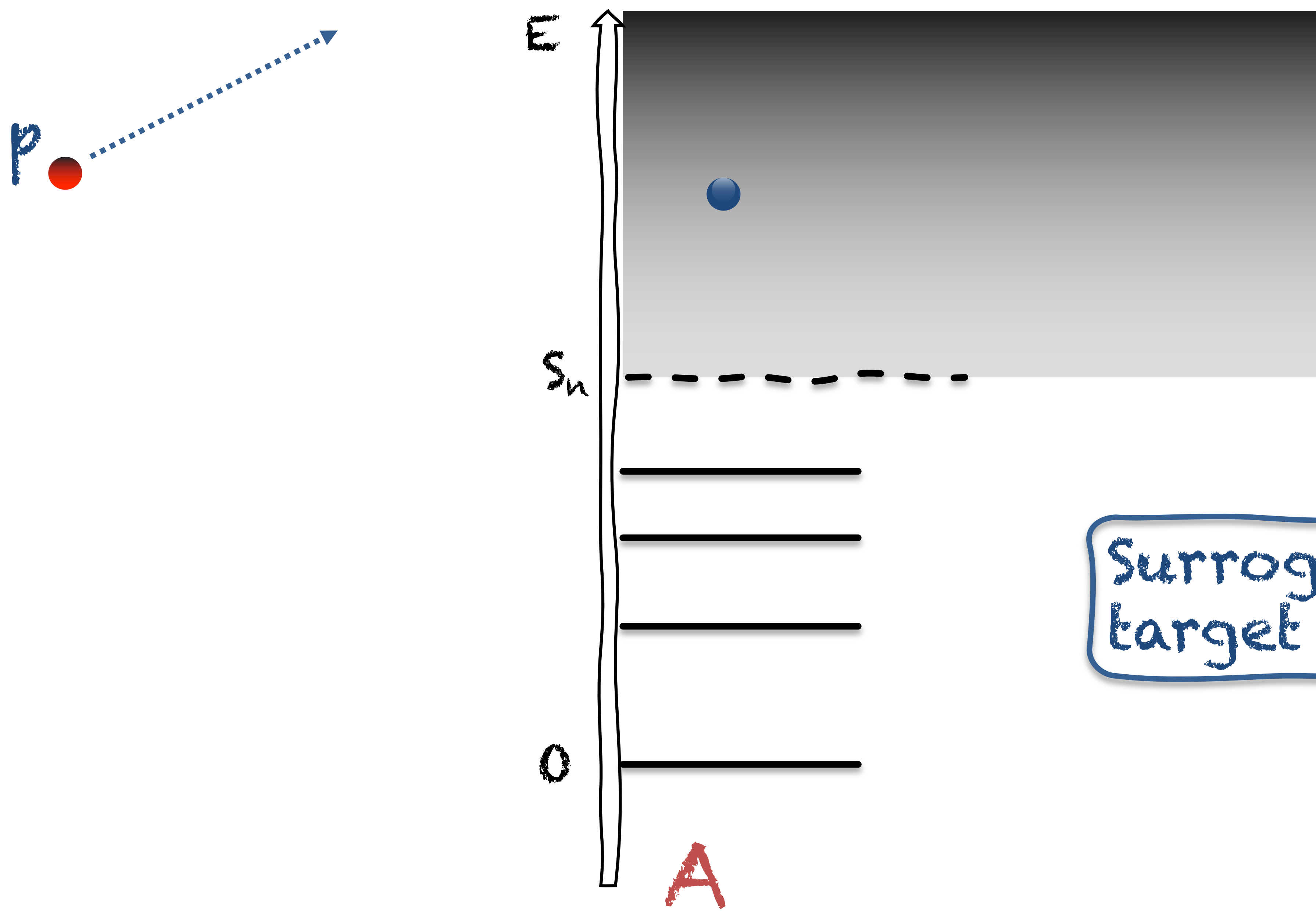
Desired reaction: neutron capture on target A

Using the Surrogate Reaction Method (SRM) to infer $^{A}X(n,\gamma)^{A+1}X$ from $^{A}X(d,p\gamma)^{A+1}X$



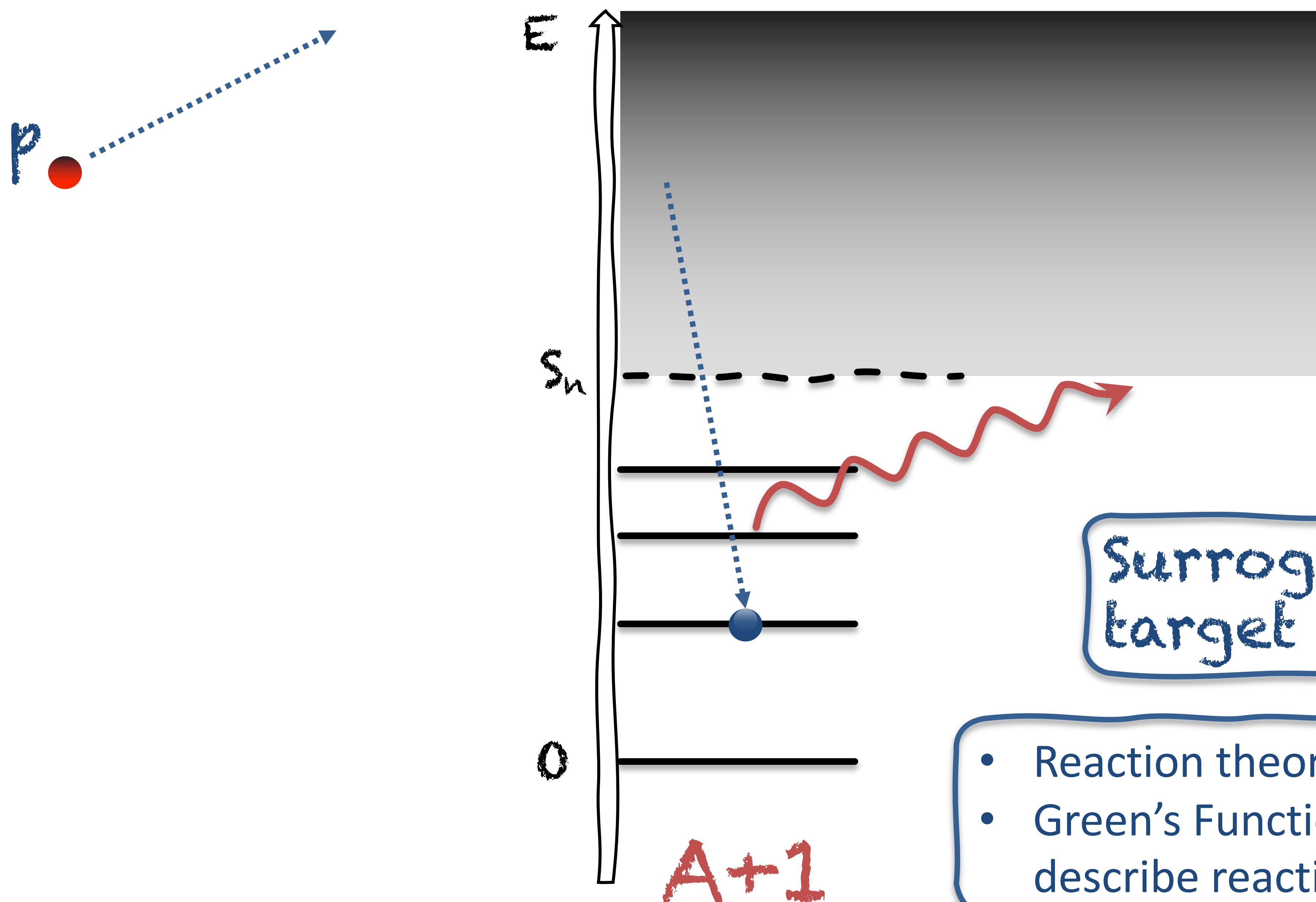
Surrogate: (d,p) reaction on target A

Using the Surrogate Reaction Method (SRM) to infer $^{A}X(n,\gamma)^{A+1}X$ from $^{A}X(d,p\gamma)^{A+1}X$



Surrogate: (d,p) reaction on target A

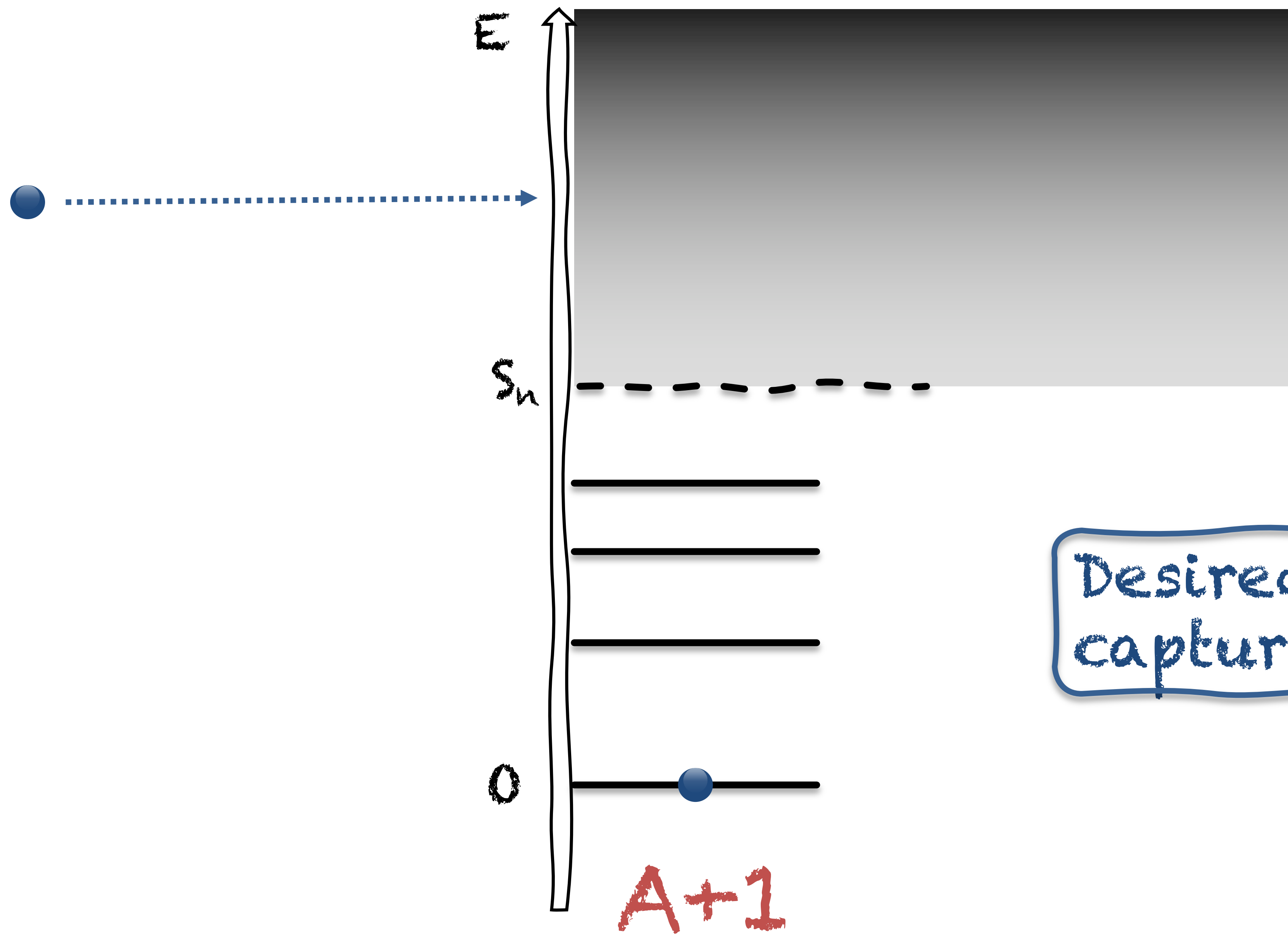
Using the Surrogate Reaction Method (SRM) to infer $^{A}X(n,\gamma)^{A+1}X$ from $^{A}X(d,p\gamma)^{A+1}X$



Surrogate: (d,p) reaction on target A

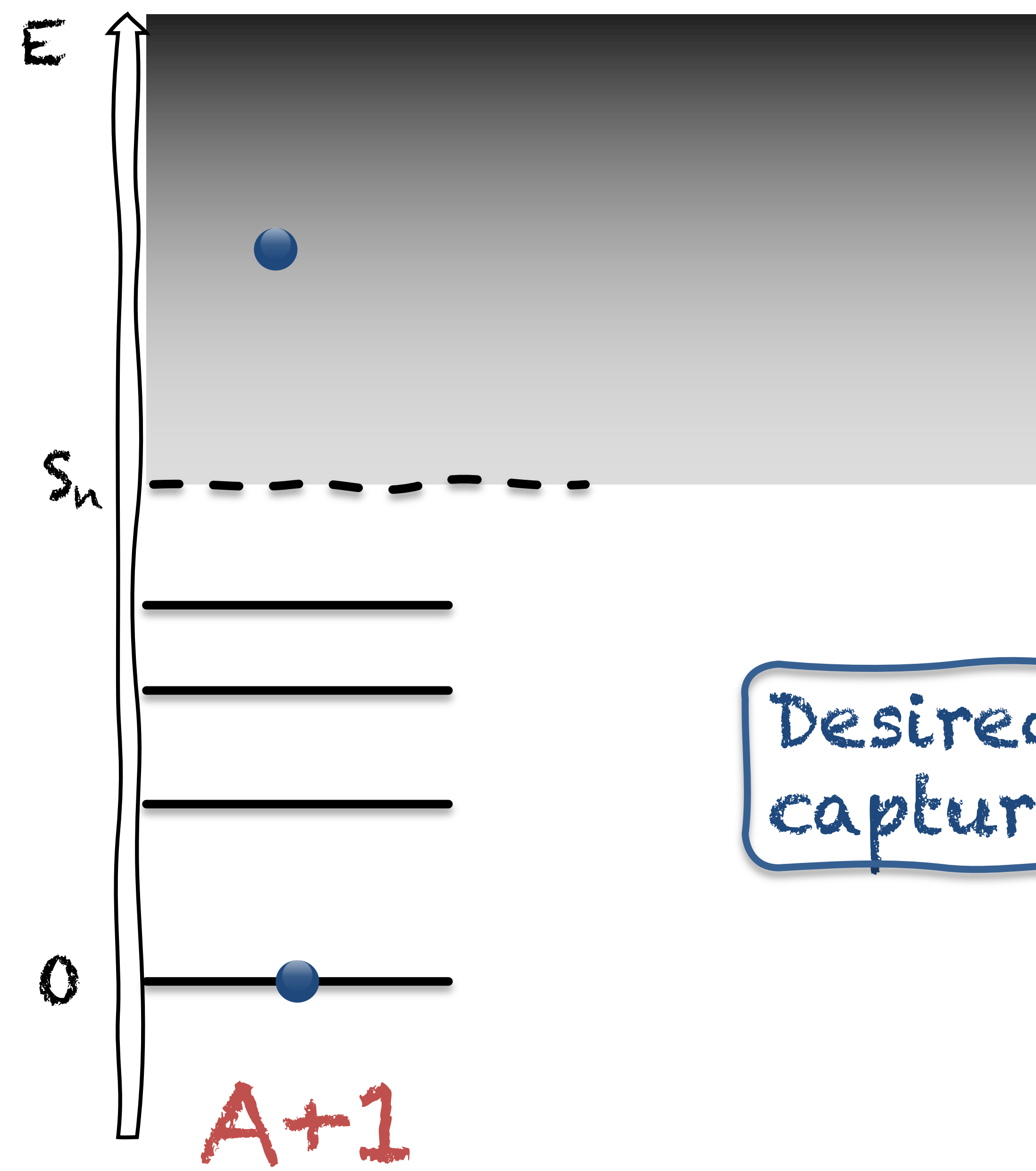
- Reaction theory needed to determine spin distribution
- Green's Function Transfer (GFT) formalism used to describe reaction process

Using the SRM to infer $^{A+1}X(n,\gamma)^{A+2}X$ from $^{A}X(t,p\gamma)^{A+2}X$



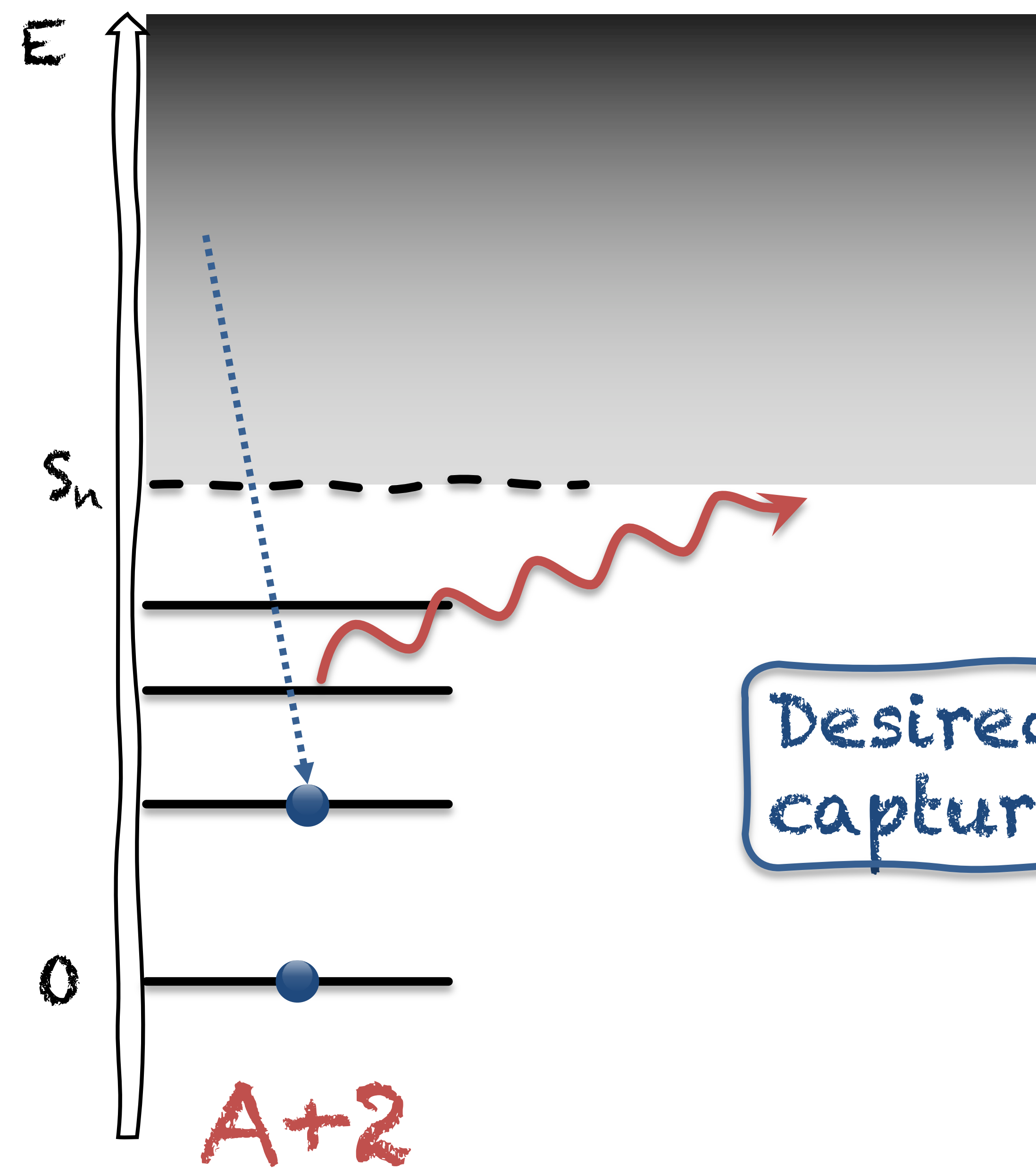
Desired reaction: neutron capture on target A+1

Using the SRM to infer $^{A+1}X(n,\gamma)^{A+2}X$ from $^{A}X(t,p\gamma)^{A+2}X$



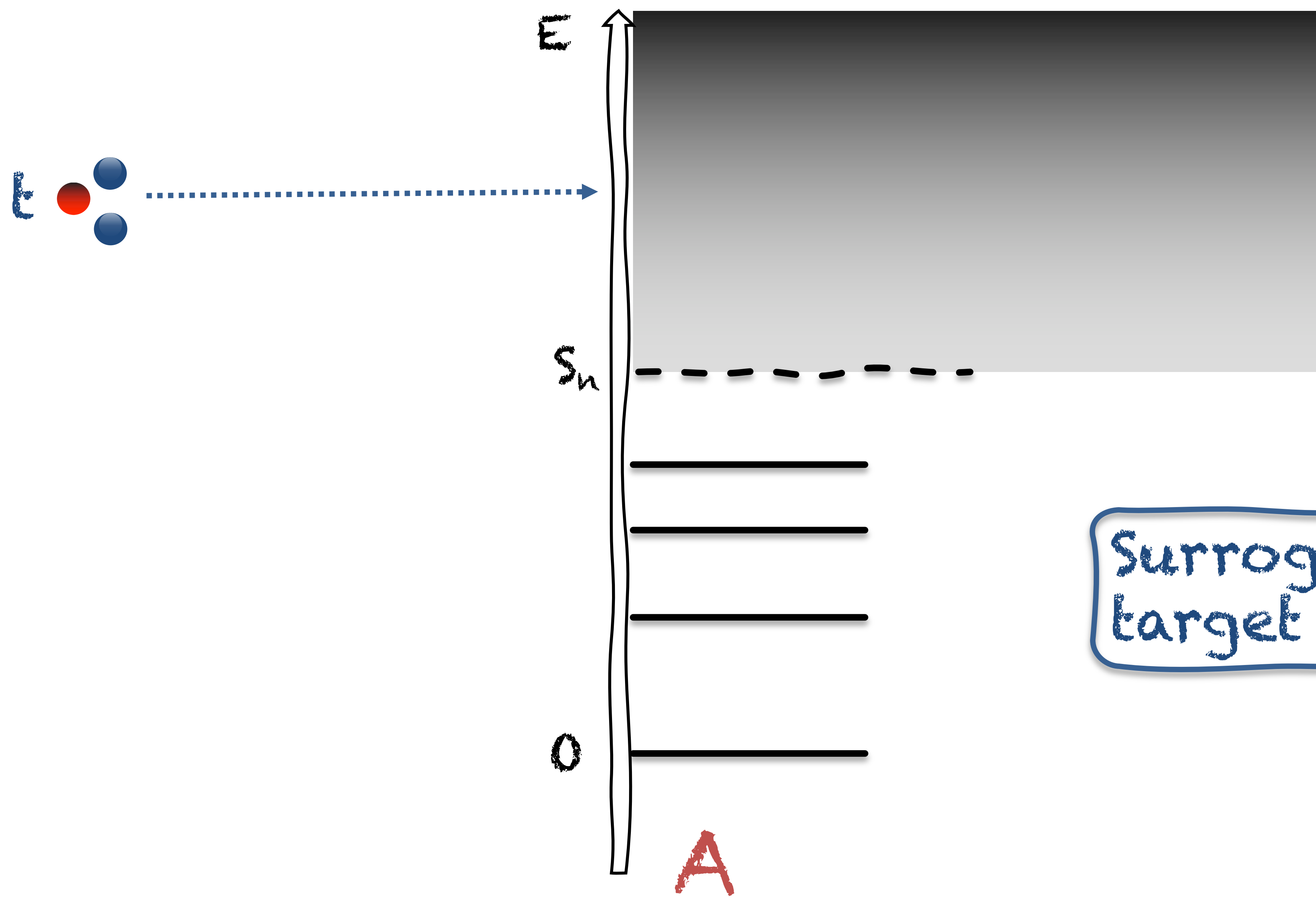
Desired reaction: neutron capture on target A+1

Using the SRM to infer $^{A+1}X(n,\gamma)^{A+2}X$ from $^{A}X(t,p\gamma)^{A+2}X$



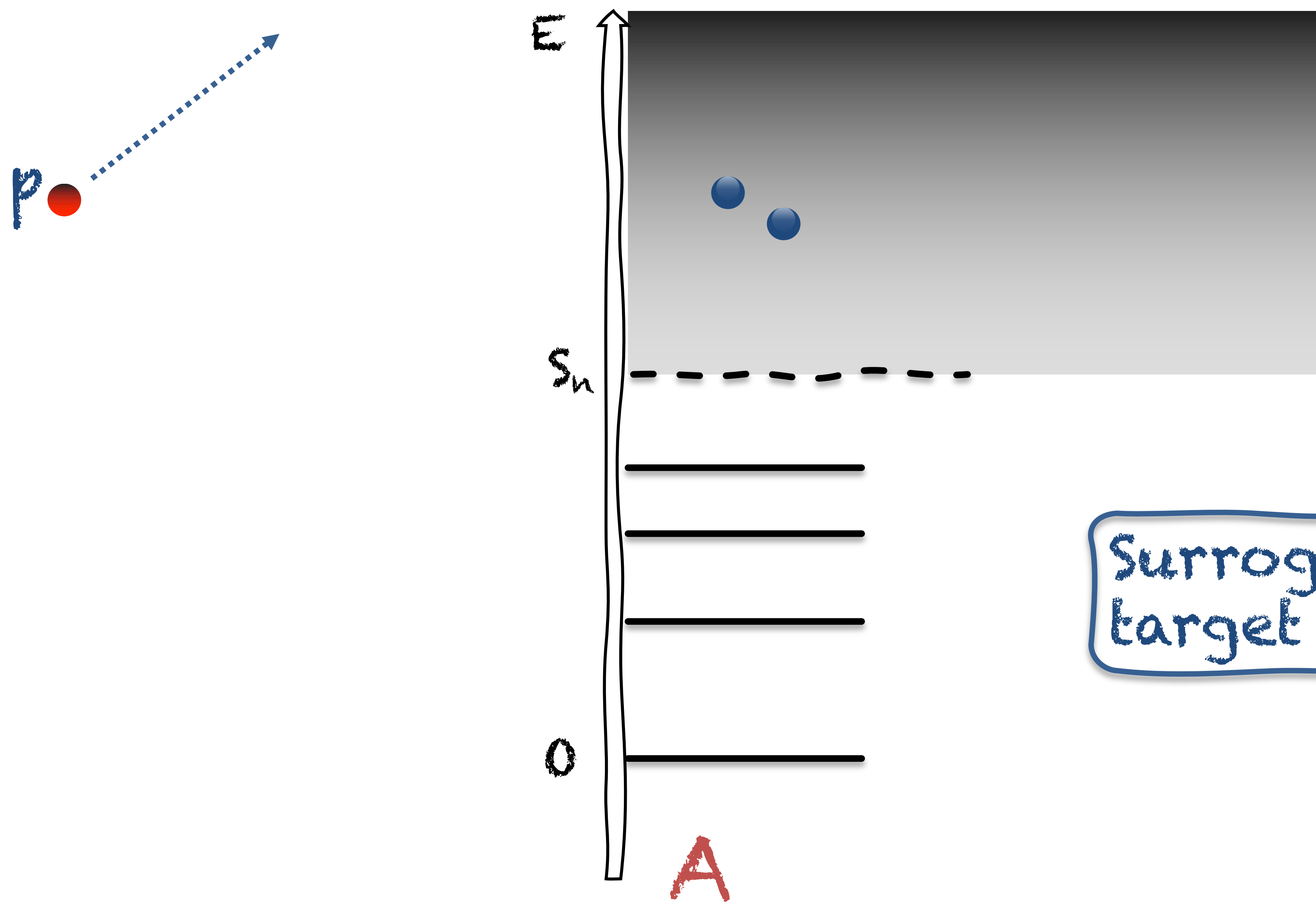
Desired reaction: neutron capture on target A+1

Using the SRM to infer $^{A+1}X(n,\gamma)^{A+2}X$ from $^{A}X(t,p\gamma)^{A+2}X$



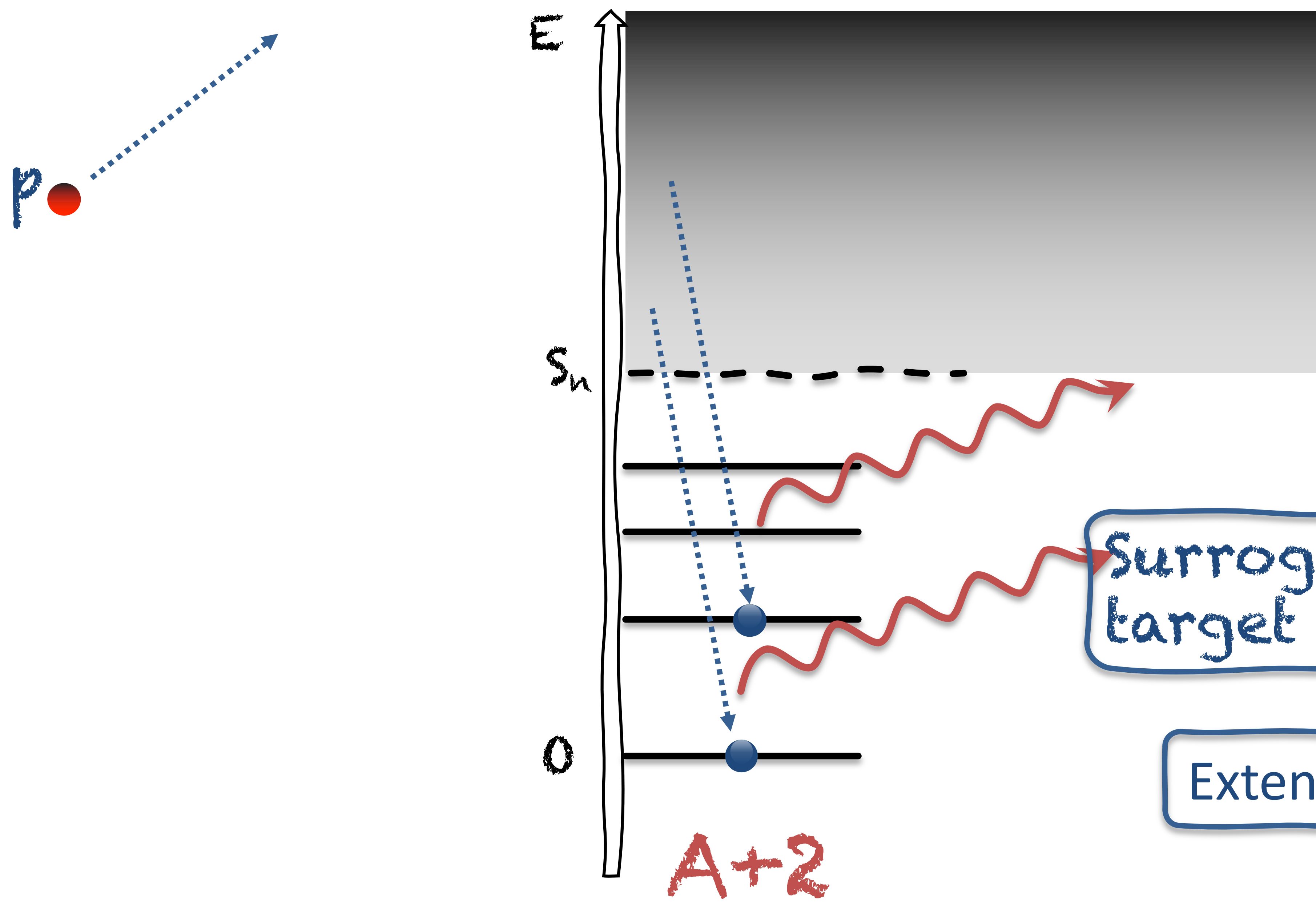
Surrogate: (t,p) reaction on target A+1

Using the SRM to infer $^{A+1}X(n,\gamma)^{A+2}X$ from $^{A}X(t,p\gamma)^{A+2}X$



Surrogate: (t,p) reaction on target A+1

Using the SRM to infer $^{A+1}X(n,\gamma)^{A+2}X$ from $^{A}X(t,p\gamma)^{A+2}X$



Surrogate: (t,p) reaction on target A+1

Extension of GFT to (t,p) reactions

Conclusions

- Pairing correlations leading to particle-number symmetry breaking are ubiquitous in Fermionic systems.
- In nuclei, (t,p) reactions are the specific probe of pairing correlations.
- Some interesting applications:
  - population of the GPV
  - population of the PDR (^{9}Li(t,p)^{11}Li(PDR) @ ISOLDE)
  - surrogate reactions

Collaborators:

- E. Vigezzi (INFN, Milano)
- F. Barranco (University of Seville)
- Y. Ayyad (IGFAE, USC)

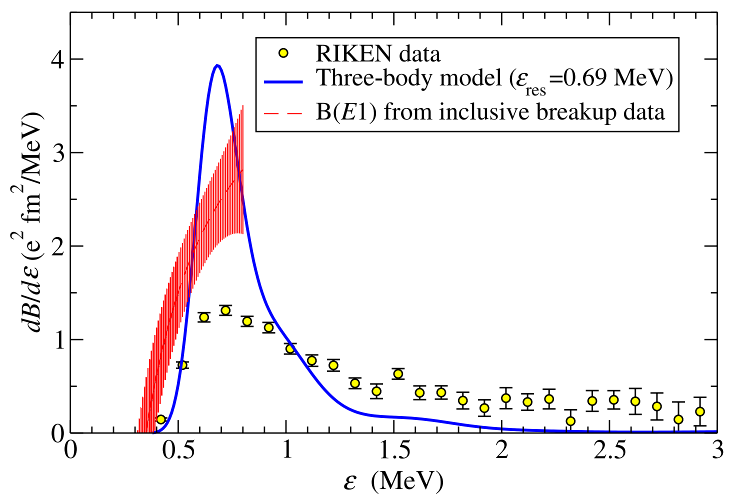
Thank you!

UNIVERSIDAD D SEVILLA 1505

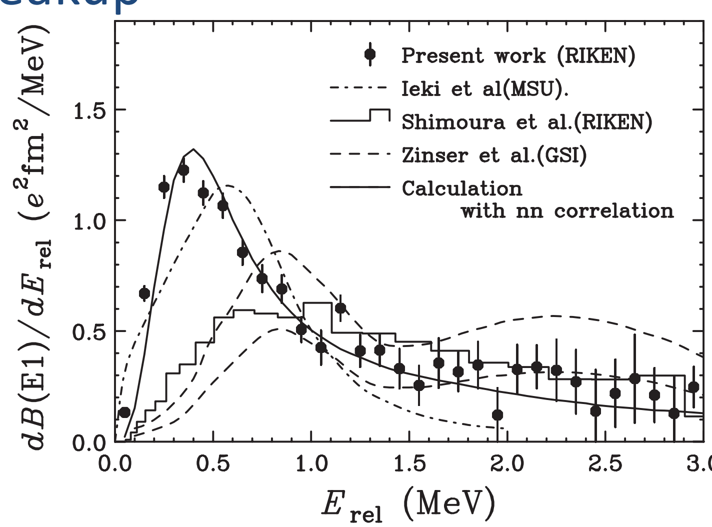
BACKUP SLIDES

Is there a pygmy resonance in ^{11}Li? What's its structure?

Coulomb breakup

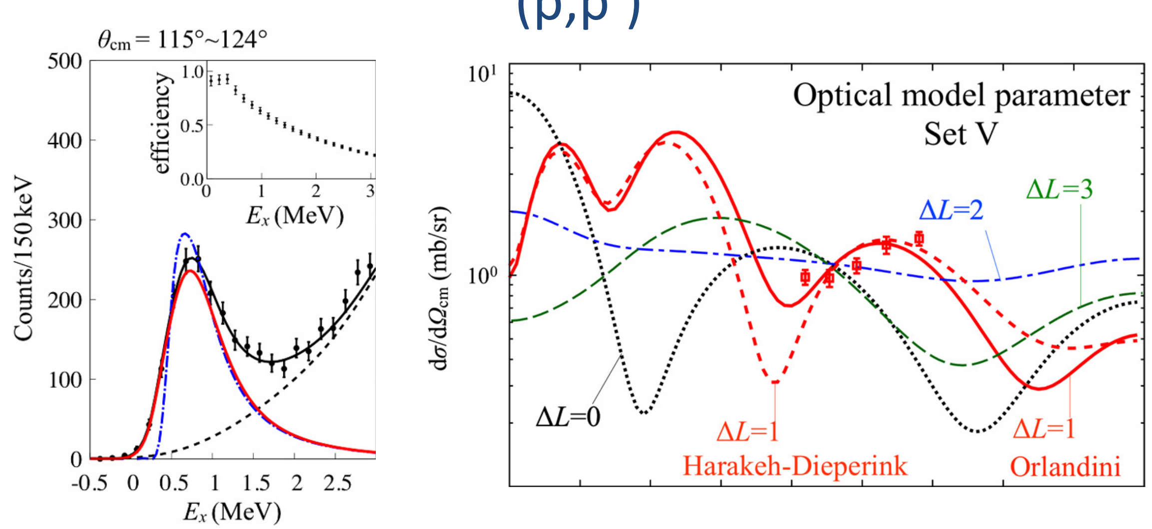


Fernández *et al*. PRL **110**, 142701 (2013)



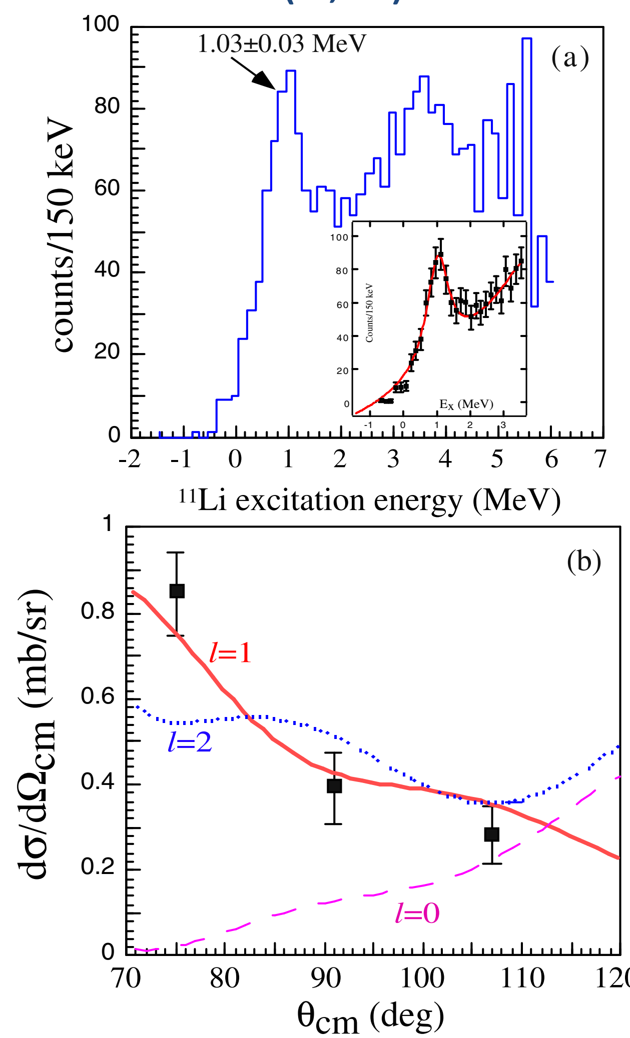
Nakamura *et al*. PRL **96**, 252502 (2006)

(p,p')



J. Tanaka et al. / Physics Letters B 774 (2017) 268–272

(d,d')

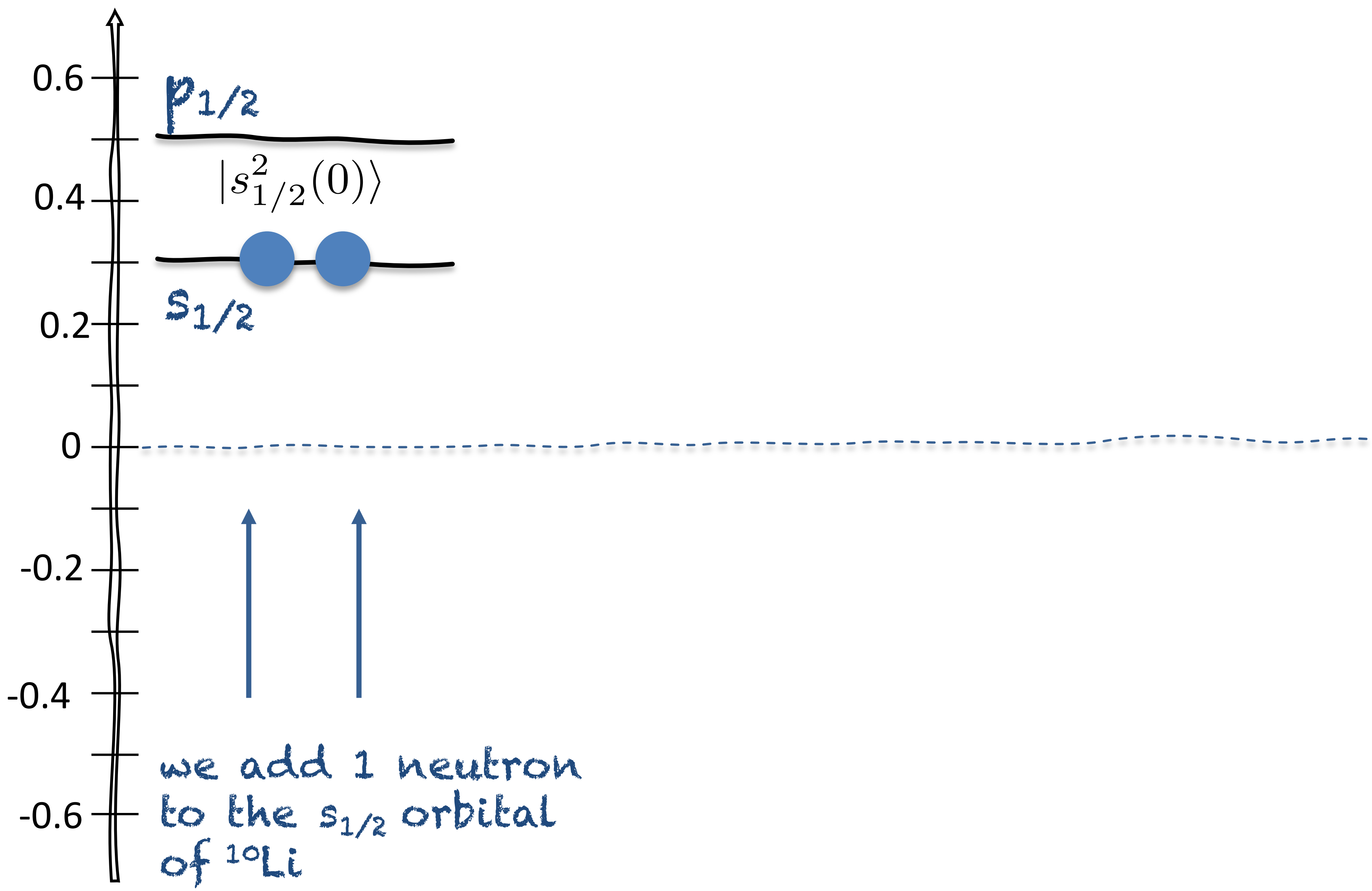


Kanungo *et al.* PRL **114**, 192502 (2015)

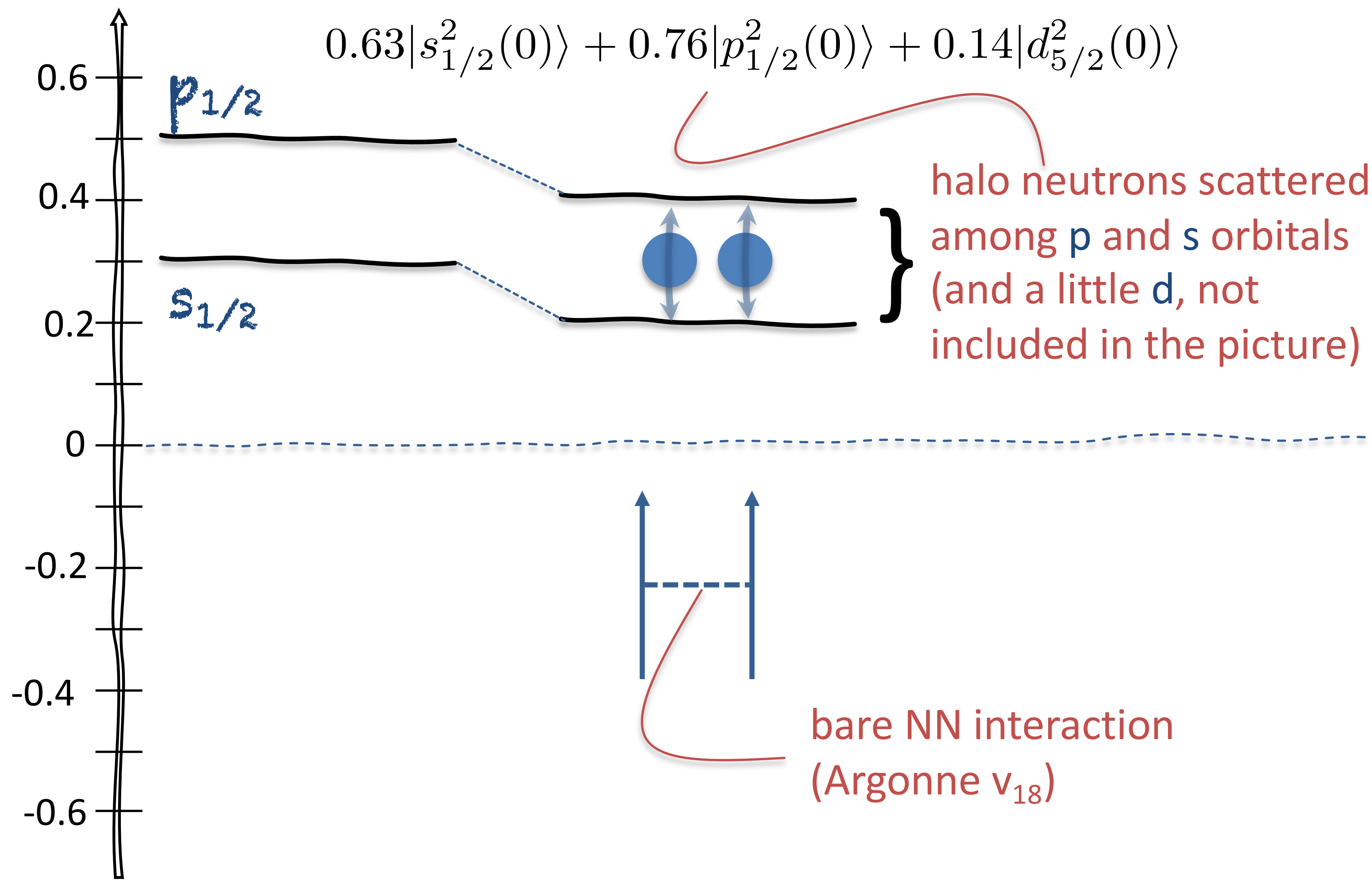
some questions to address:

- How do we characterize the PDR?
- Is it distinct from the GDR?
- How does it compare with theory?

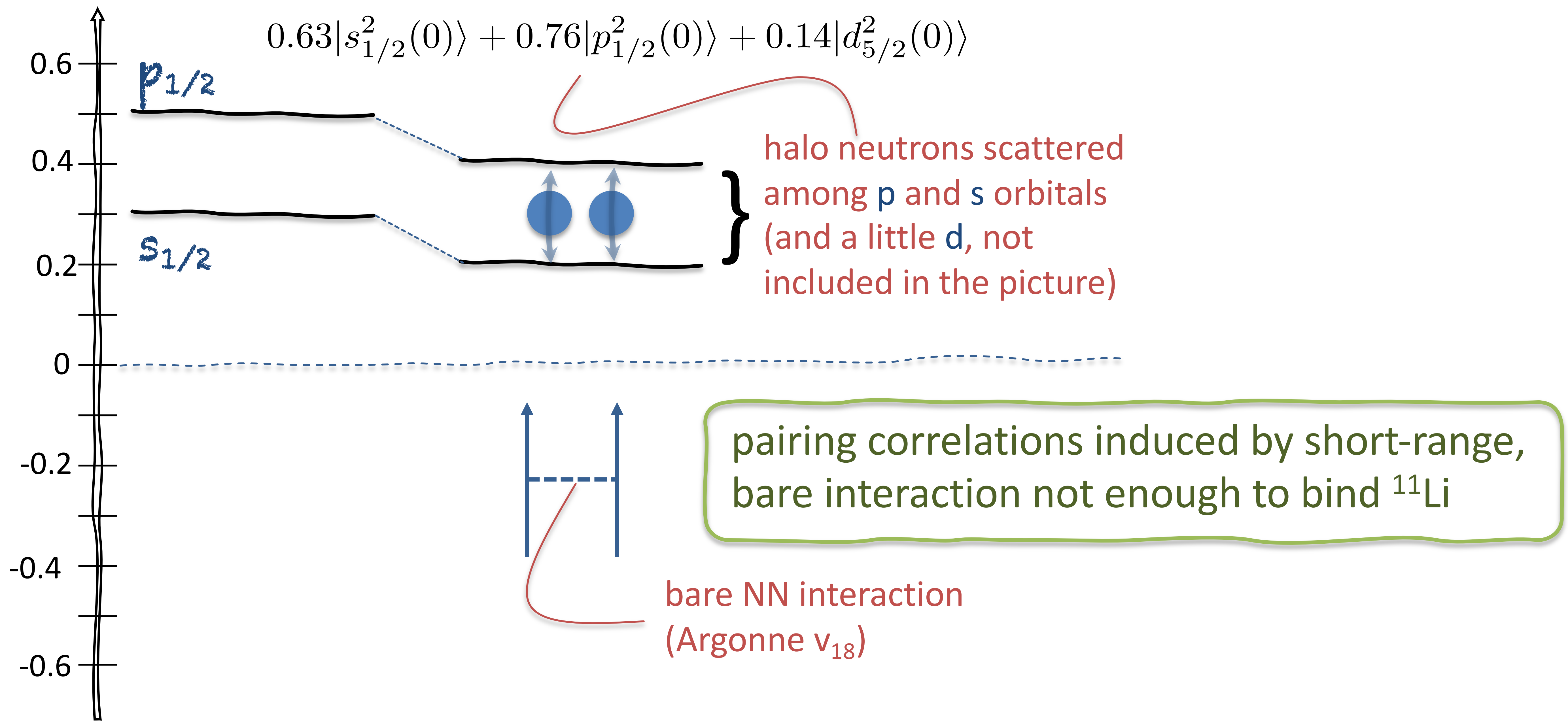
Nuclear field theory (NFT) highlights the role of the PDR in ^{11}Li structure



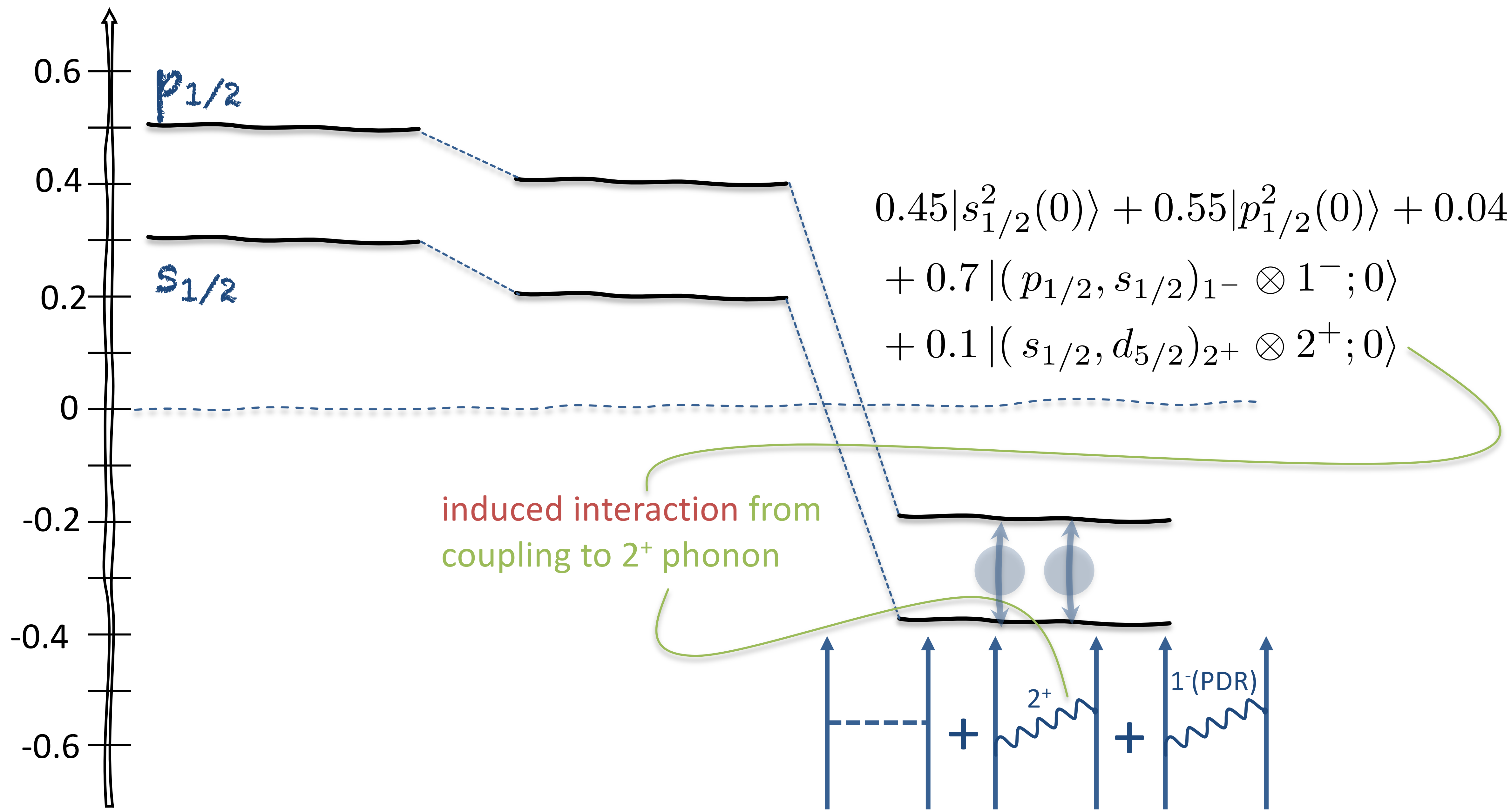
Nuclear field theory (NFT) highlights the role of the PDR in ^{11}Li structure



Nuclear field theory (NFT) highlights the role of the PDR in ^{11}Li structure

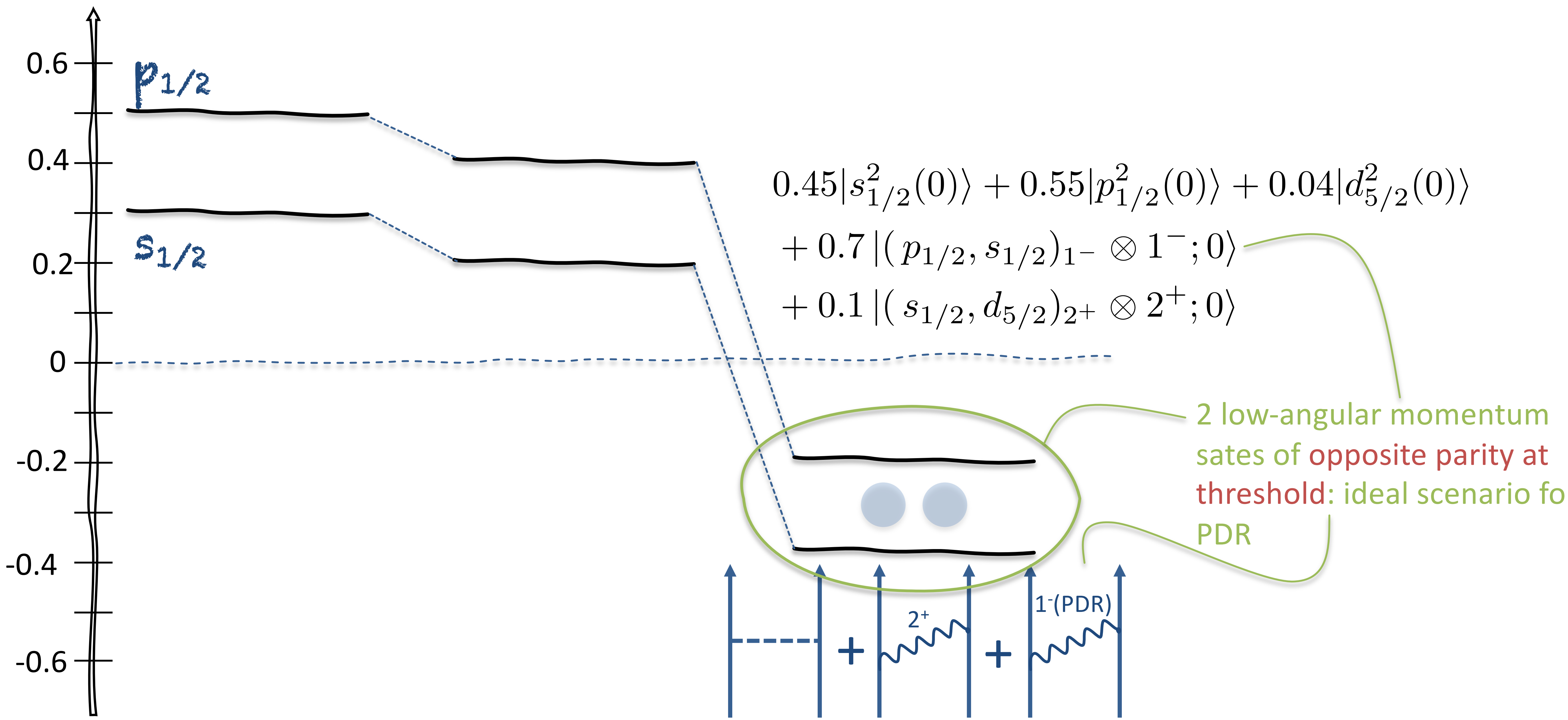


Nuclear field theory (NFT) highlights the role of the PDR in ^{11}Li structure



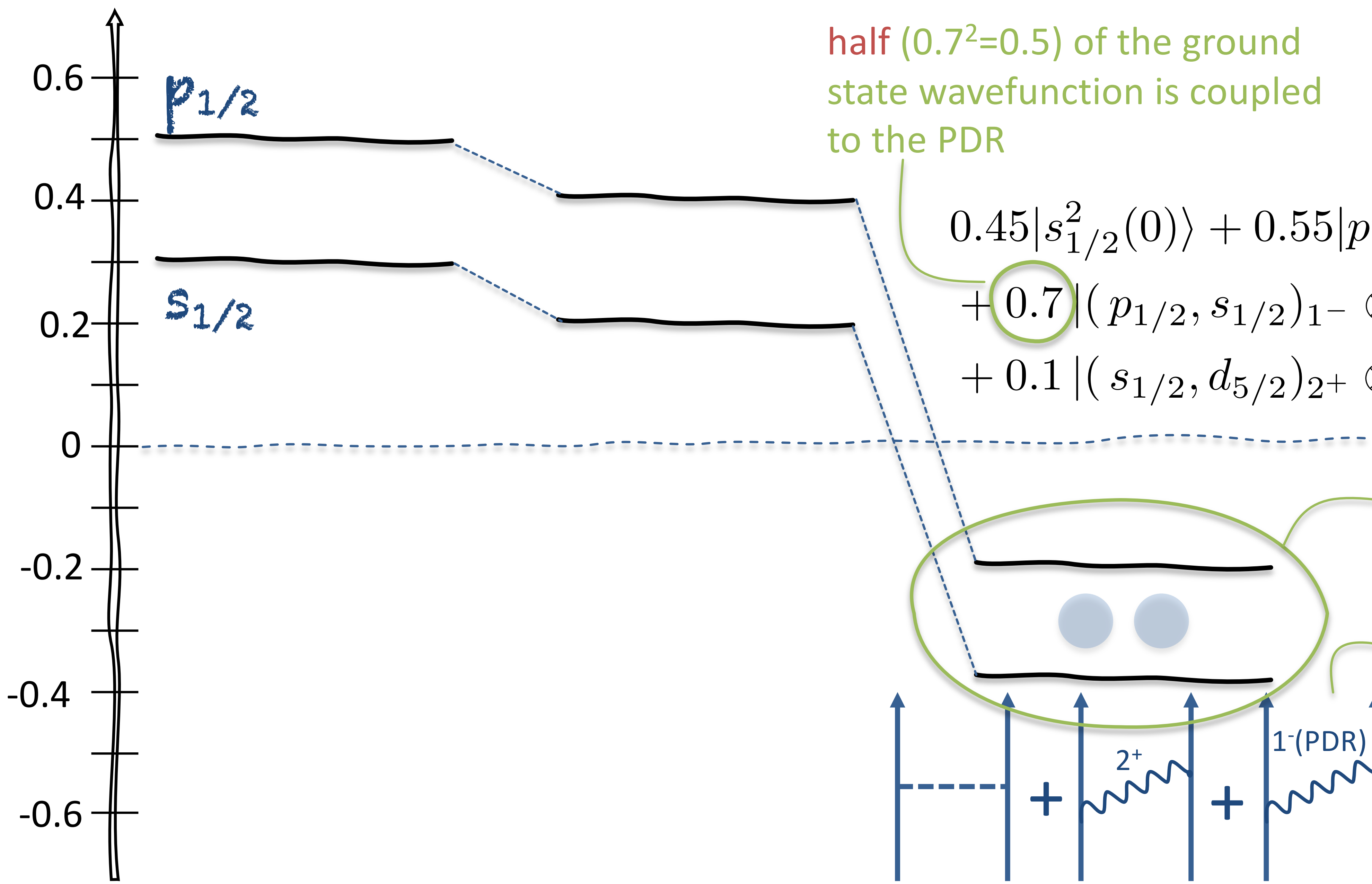
$$0.45|s^2_{1/2}(0)\rangle + 0.55|p^2_{1/2}(0)\rangle + 0.04|d^2_{5/2}(0)\rangle$$
$$+\,0.7\,|(p_{1/2}, s_{1/2})_{1^-} \otimes 1^-; 0\rangle$$
$$+\,0.1\,|(s_{1/2}, d_{5/2})_{2^+} \otimes 2^+; 0\rangle$$

Nuclear field theory (NFT) highlights the role of the PDR in ^{11}Li structure



$$0.45|s^2_{1/2}(0)\rangle + 0.55|p^2_{1/2}(0)\rangle + 0.04|d^2_{5/2}(0)\rangle$$
$$+\,0.7\,|(\,p_{1/2}, s_{1/2})_{1^-} \otimes 1^-; 0\rangle$$
$$+\,0.1\,|(\,s_{1/2}, d_{5/2})_{2^+} \otimes 2^+; 0\rangle$$

Nuclear field theory (NFT) highlights the role of the PDR in ^{11}Li structure

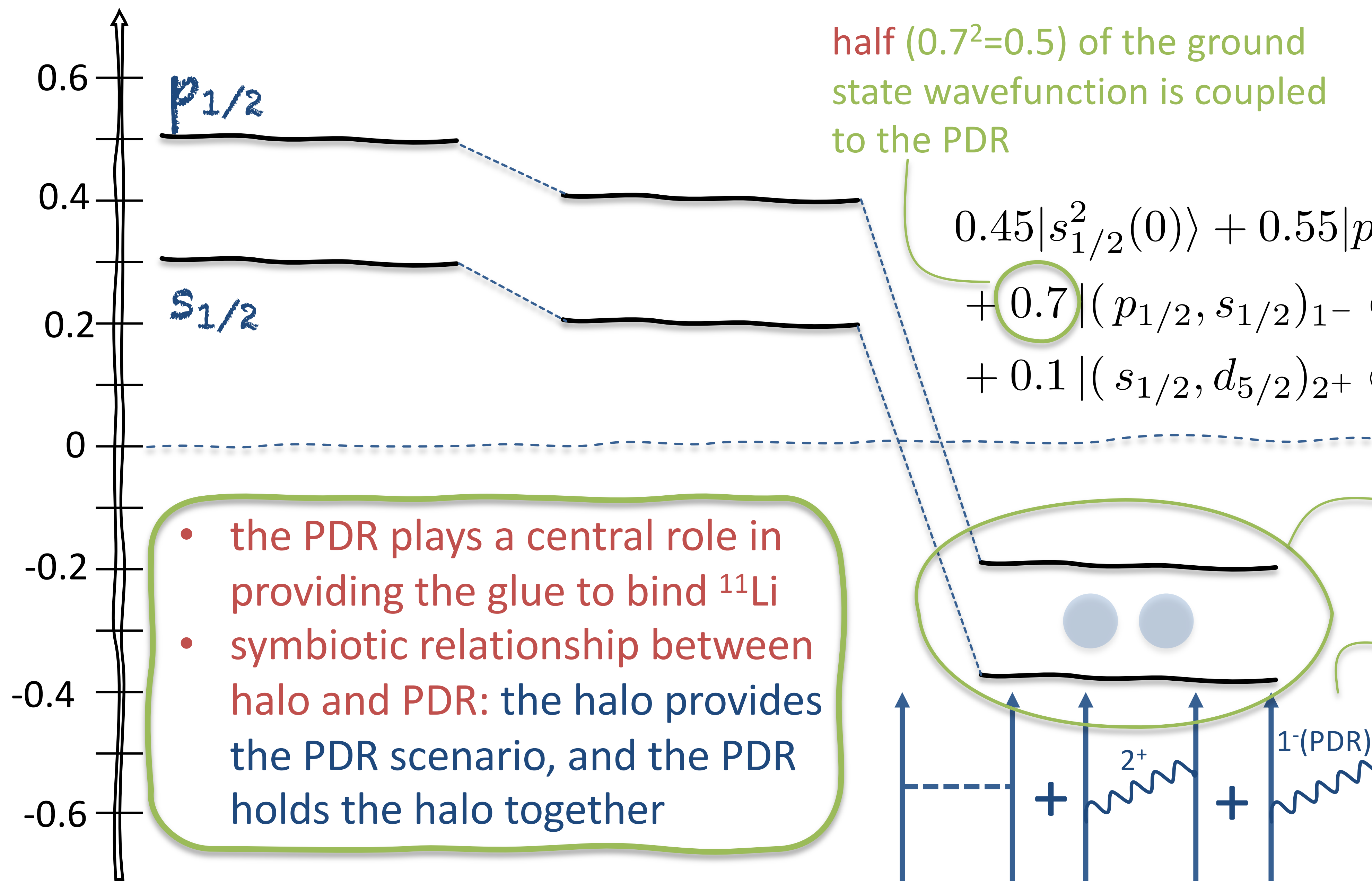


half (0.7^2=0.5) of the ground state wavefunction is coupled to the PDR

$$0.45|s^2_{1/2}(0)\rangle + 0.55|p^2_{1/2}(0)\rangle + 0.04|d^2_{5/2}(0)\rangle$$
$$+ 0.7\,|(p_{1/2}, s_{1/2})_{1^-} \otimes 1^-; 0\rangle$$
$$+ 0.1\,|(s_{1/2}, d_{5/2})_{2^+} \otimes 2^+; 0\rangle$$

2 low-angular momentum sates of opposite parity at threshold: ideal scenario for PDR

Nuclear field theory (NFT) highlights the role of the PDR in ^{11}Li structure



half (0.7^2=0.5) of the ground state wavefunction is coupled to the PDR

$$
\begin{aligned}
&0.45|s^2_{1/2}(0)\rangle + 0.55|p^2_{1/2}(0)\rangle + 0.04|d^2_{5/2}(0)\rangle \\
&+ 0.7\,|(p_{1/2}, s_{1/2})_{1^-} \otimes 1^-; 0\rangle \\
&+ 0.1\,|(s_{1/2}, d_{5/2})_{2^+} \otimes 2^+; 0\rangle
\end{aligned}
$$

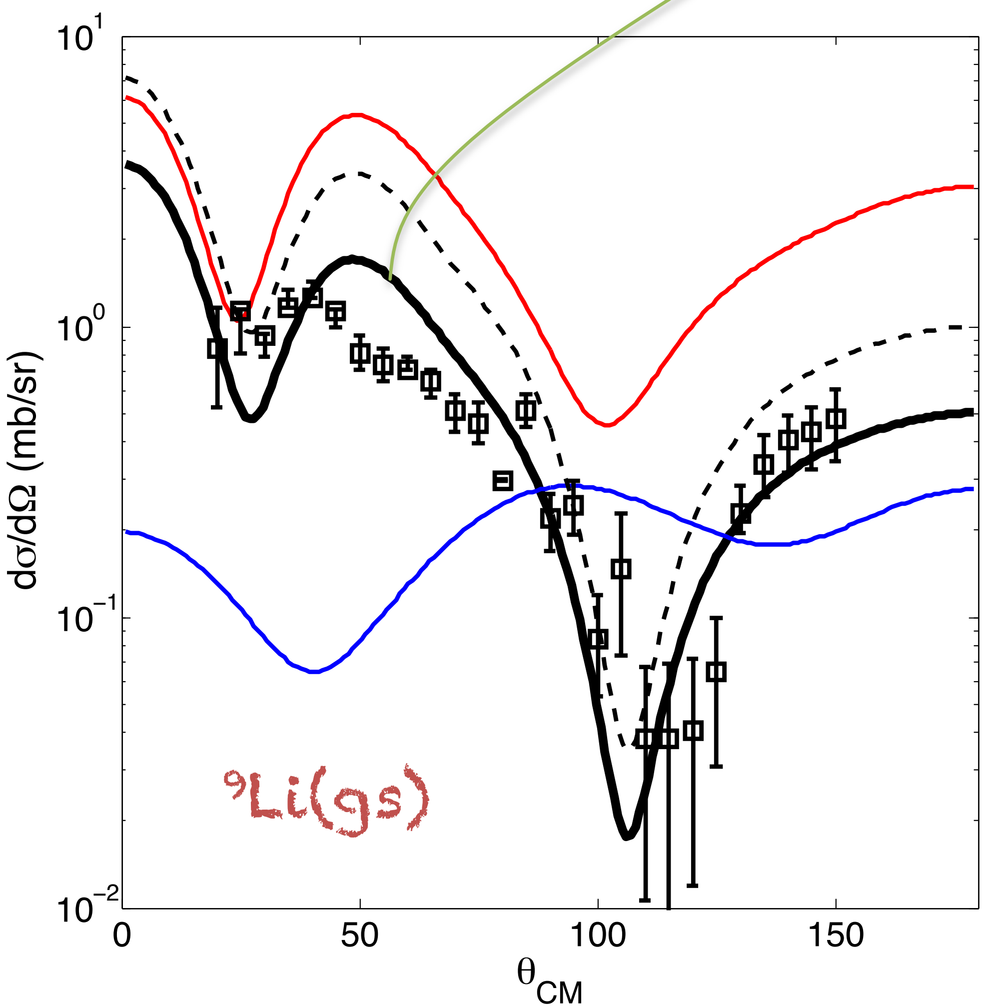
2 low-angular momentum sates of opposite parity at threshold: ideal scenario for PDR

- the PDR plays a central role in providing the glue to bind ^{11}Li
- symbiotic relationship between halo and PDR: the halo provides the PDR scenario, and the PDR holds the halo together

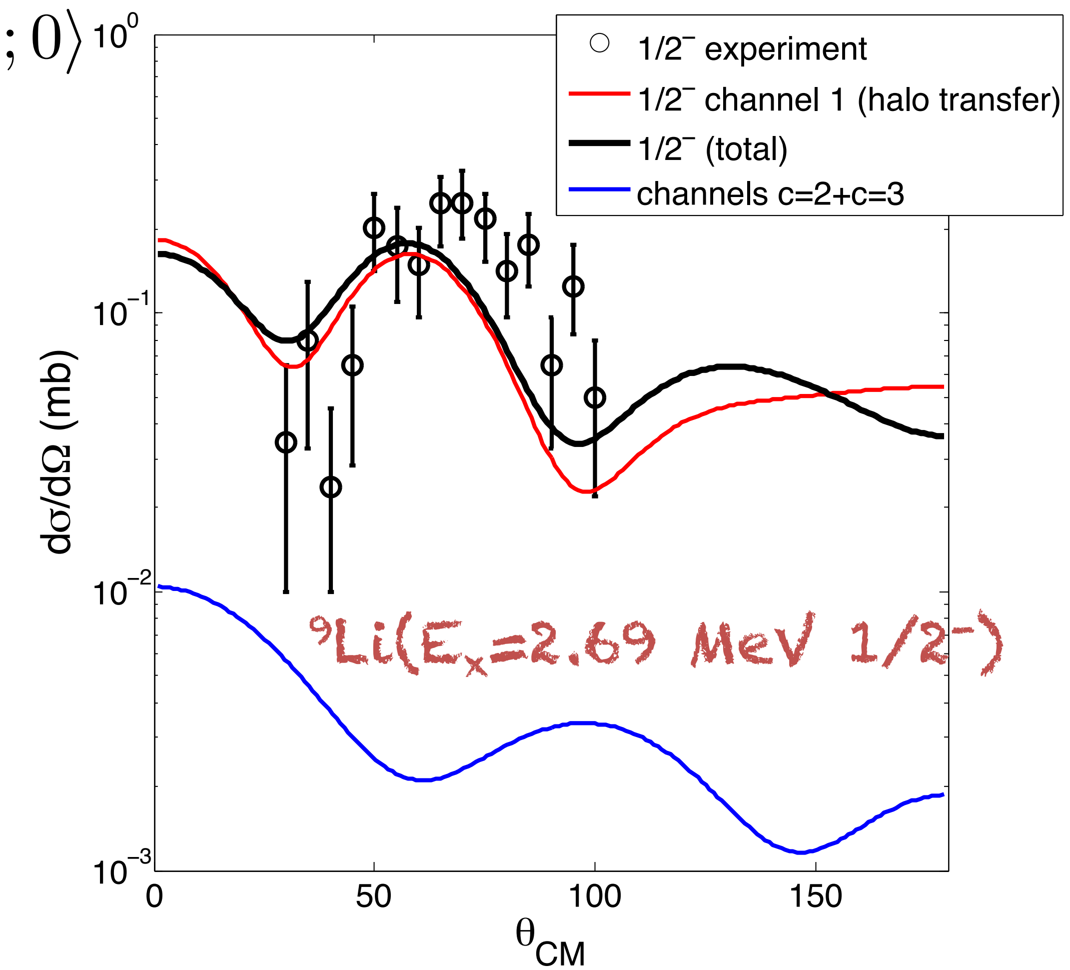
theory confirmed by $^{11}Li(p,t)^9Li(gs;E_x=2.69\ \text{MeV}\ 1/2^-)$

GP, Barranco, Vigezzi, Broglia
PRL **105** 172502 (2010)

$$0.45|s_{1/2}^2(0)\rangle + 0.55|p_{1/2}^2(0)\rangle + 0.04|d_{5/2}^2(0)\rangle$$
$$+\,0.7\,|(p_{1/2}, s_{1/2})_{1^-} \otimes 1^-; 0\rangle$$
$$+\,0.1\,|(s_{1/2}, d_{5/2})_{2^+} \otimes 2^+; 0\rangle$$



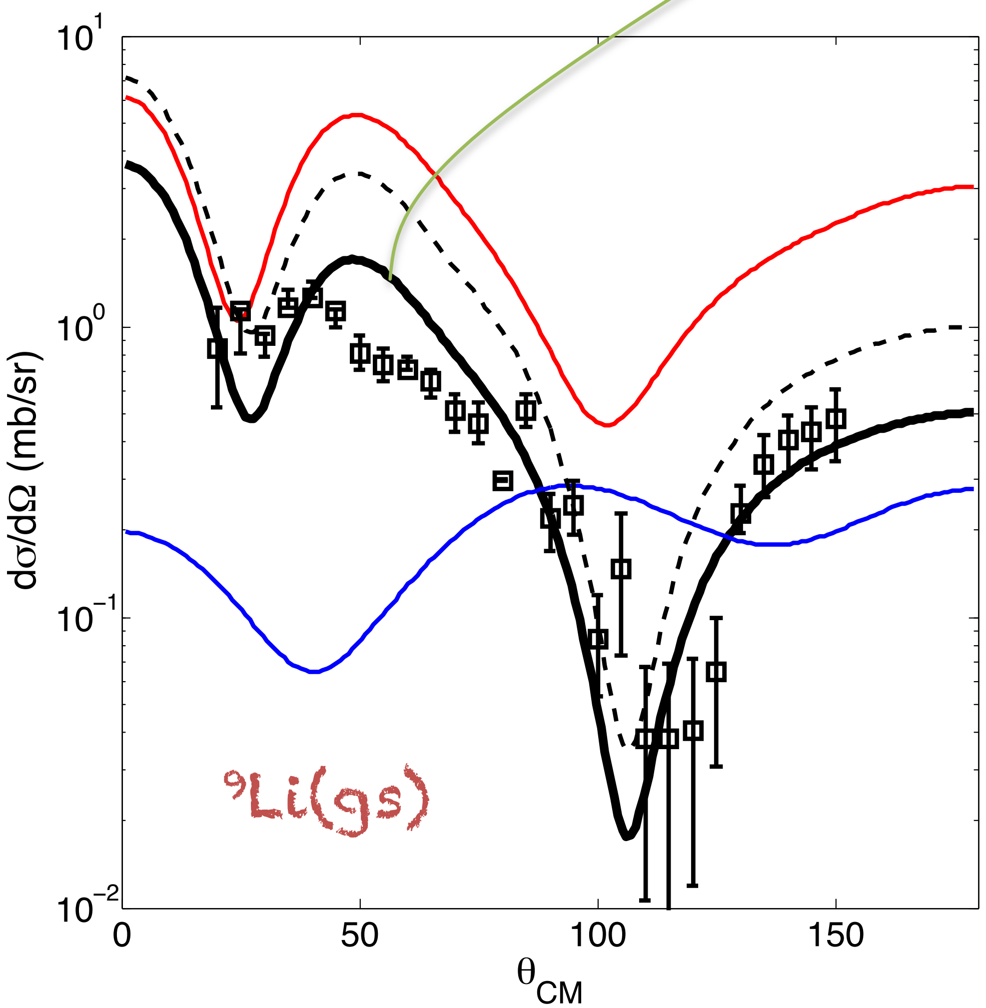
reaction calculation in 2-order DWBA, dominated by successive transfer of the 2 neutrons (E. Vigezzi talk yesterday)

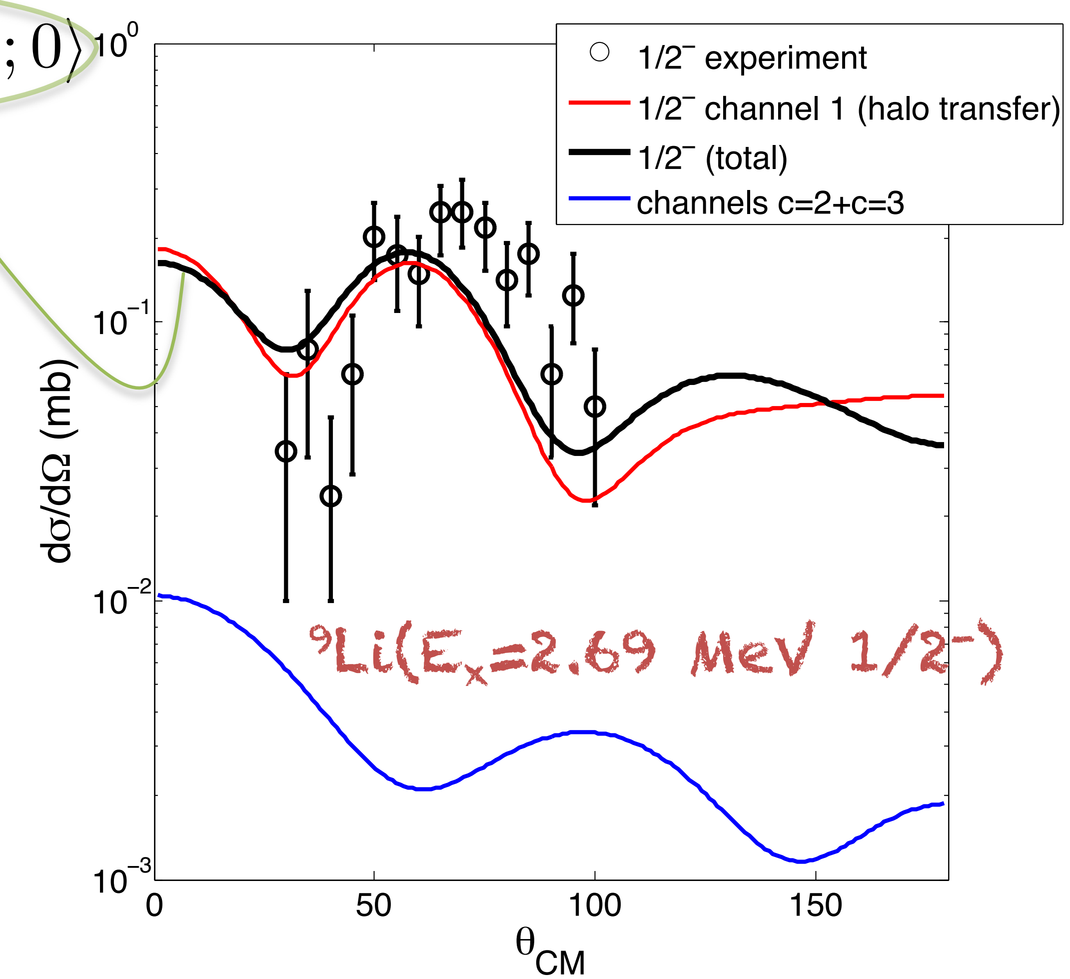


theory confirmed by ^{11}Li(p,t)^{9}Li(gs;E_x=2.69 MeV 1/2$^-$)

GP, Barranco, Vigezzi, Broglia
PRL **105** 172502 (2010)

$$0.45|s^2_{1/2}(0)\rangle + 0.55|p^2_{1/2}(0)\rangle + 0.04|d^2_{5/2}(0)\rangle$$
$$+\,0.7\,|(p_{1/2}, s_{1/2})_{1^-} \otimes 1^-; 0\rangle$$
$$+\,0.1\,|(s_{1/2}, d_{5/2})_{2^+} \otimes 2^+; 0\rangle$$





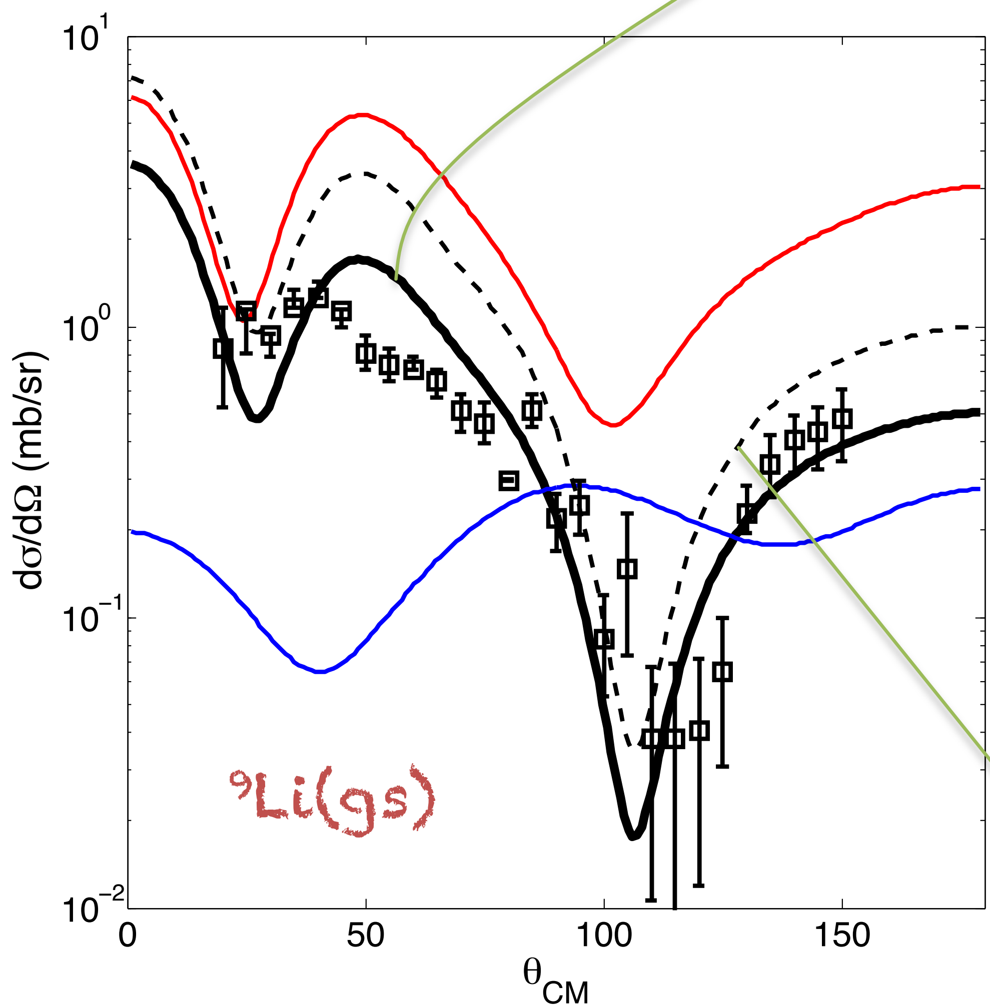
theory confirmed by ^{11}Li(p,t)^{9}Li(gs;E_x=2.69 MeV 1/2$^-$)

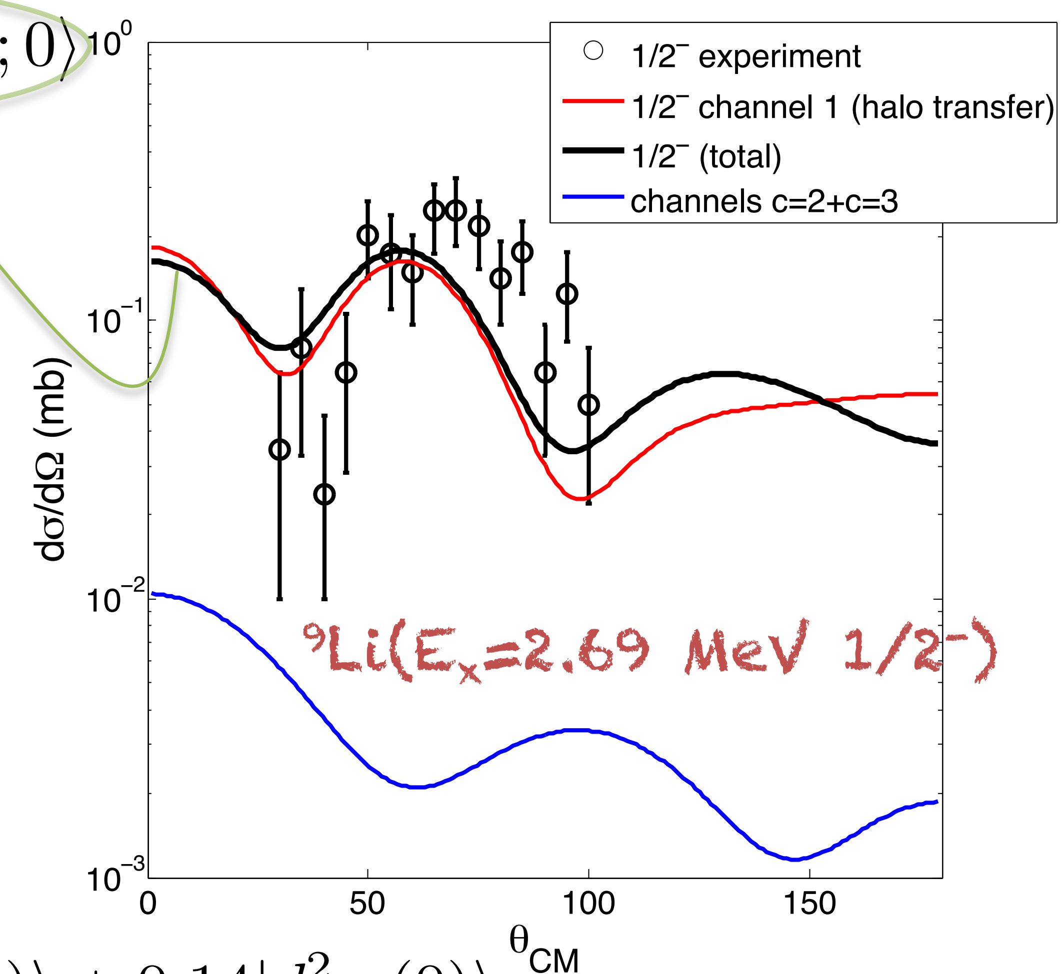
GP, Barranco, Vigezzi, Broglia
PRL **105** 172502 (2010)

$$0.45|s^2_{1/2}(0)\rangle + 0.55|p^2_{1/2}(0)\rangle + 0.04|d^2_{5/2}(0)\rangle$$

$$+\,0.7\,|(p_{1/2}, s_{1/2})_{1^-} \otimes 1^-; 0\rangle$$

$$+\,0.1\,|(s_{1/2}, d_{5/2})_{2^+} \otimes 2^+; 0\rangle$$



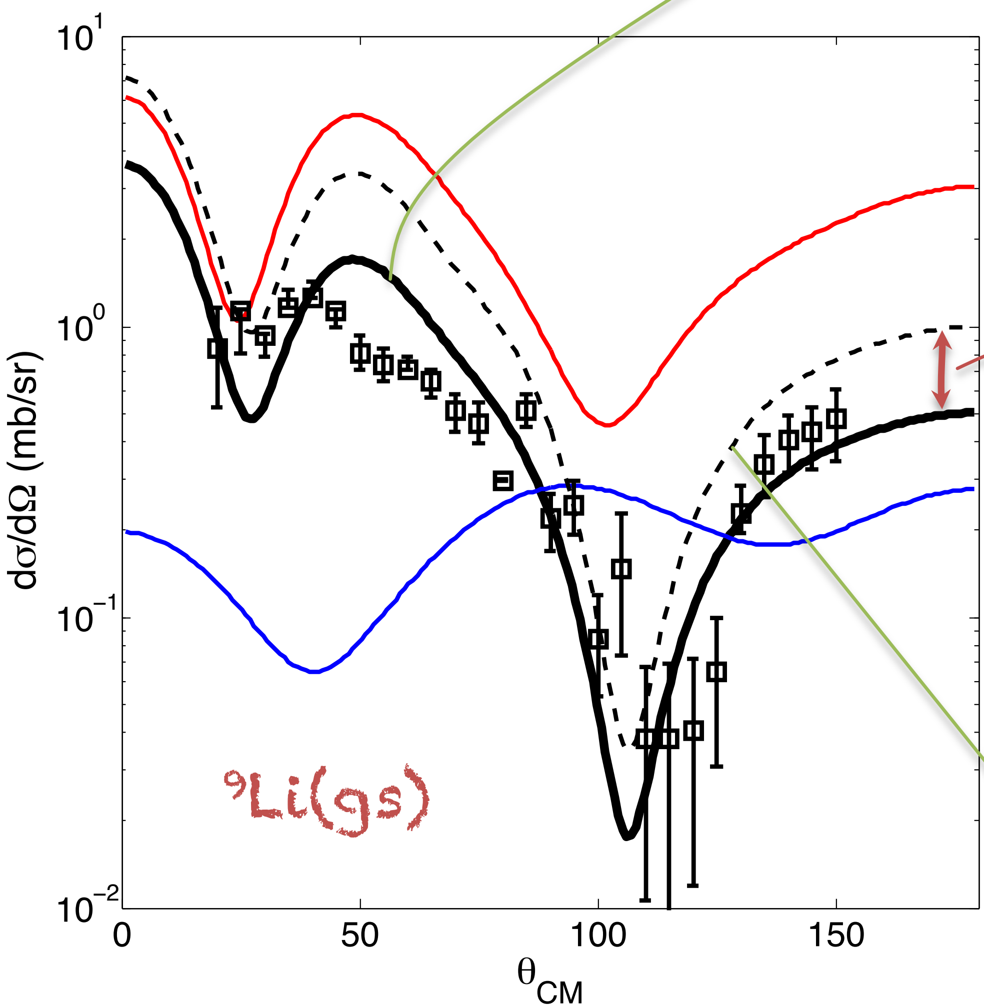


$$0.63|s^2_{1/2}(0)\rangle + 0.76|p^2_{1/2}(0)\rangle + 0.14|d^2_{1/2}(0)\rangle$$

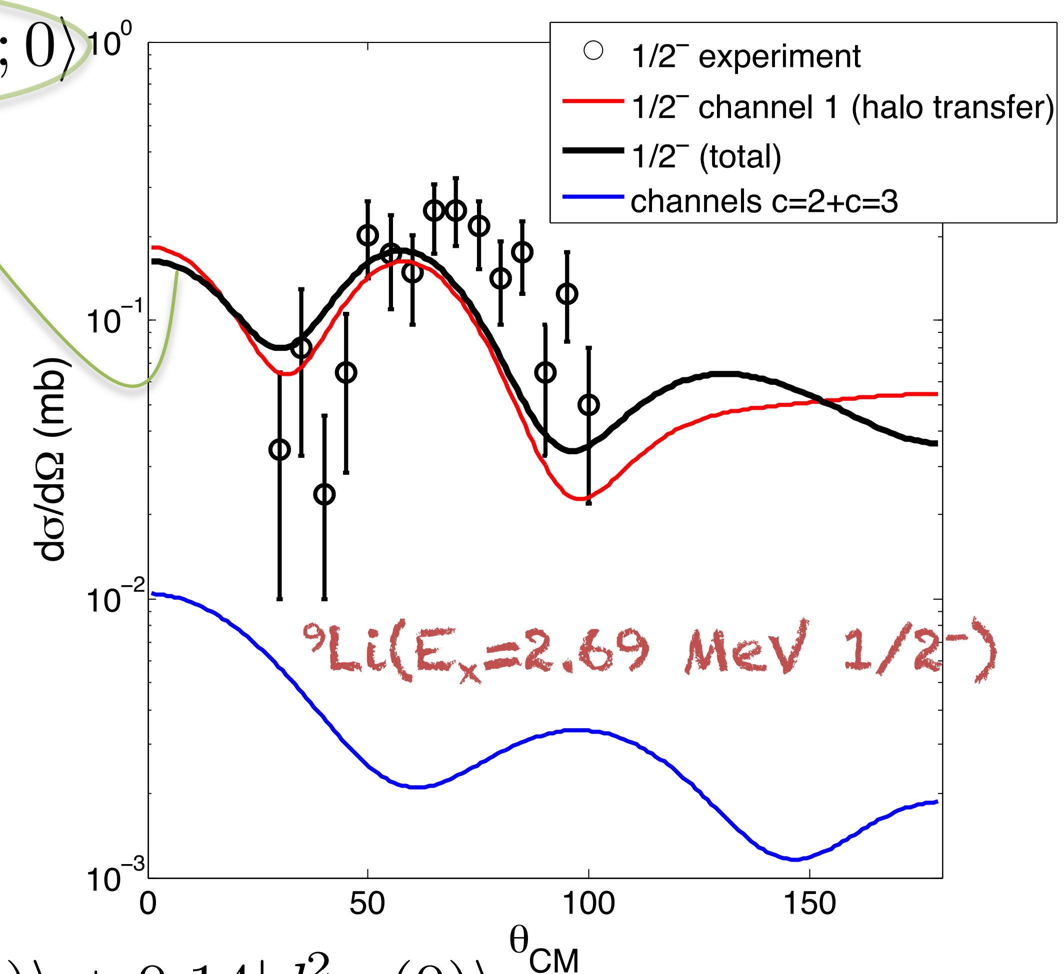
theory confirmed by ^{11}Li(p,t)^{9}Li(gs;E_x=2.69 MeV 1/2$^-$)

GP, Barranco, Vigezzi, Broglia
PRL **105** 172502 (2010)

$$0.45|s^2_{1/2}(0)\rangle + 0.55|p^2_{1/2}(0)\rangle + 0.04|d^2_{5/2}(0)\rangle$$
$$+\,0.7\,|(p_{1/2}, s_{1/2})_{1^-} \otimes 1^-; 0\rangle$$
$$+\,0.1\,|(s_{1/2}, d_{5/2})_{2^+} \otimes 2^+; 0\rangle$$



depletion of pp strength due to coupling with PDR

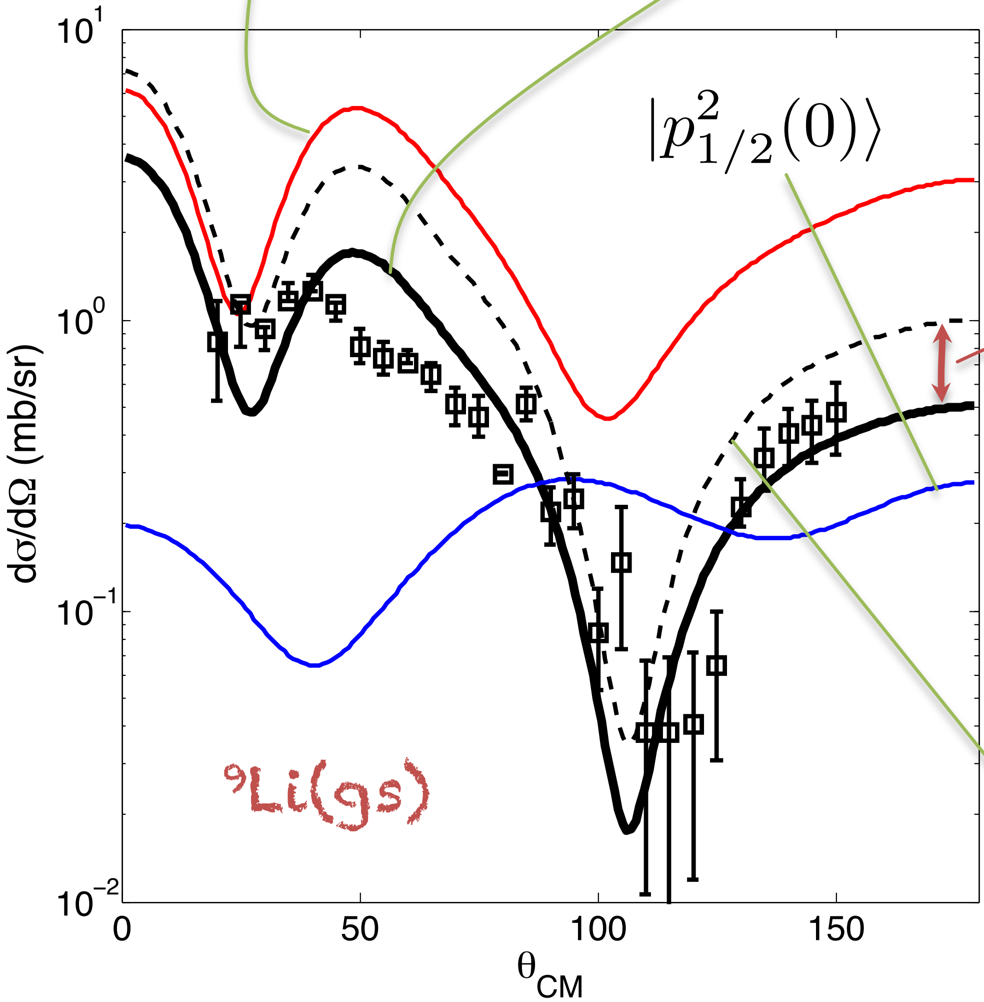


$$0.63|s^2_{1/2}(0)\rangle + 0.76|p^2_{1/2}(0)\rangle + 0.14|d^2_{1/2}(0)\rangle$$

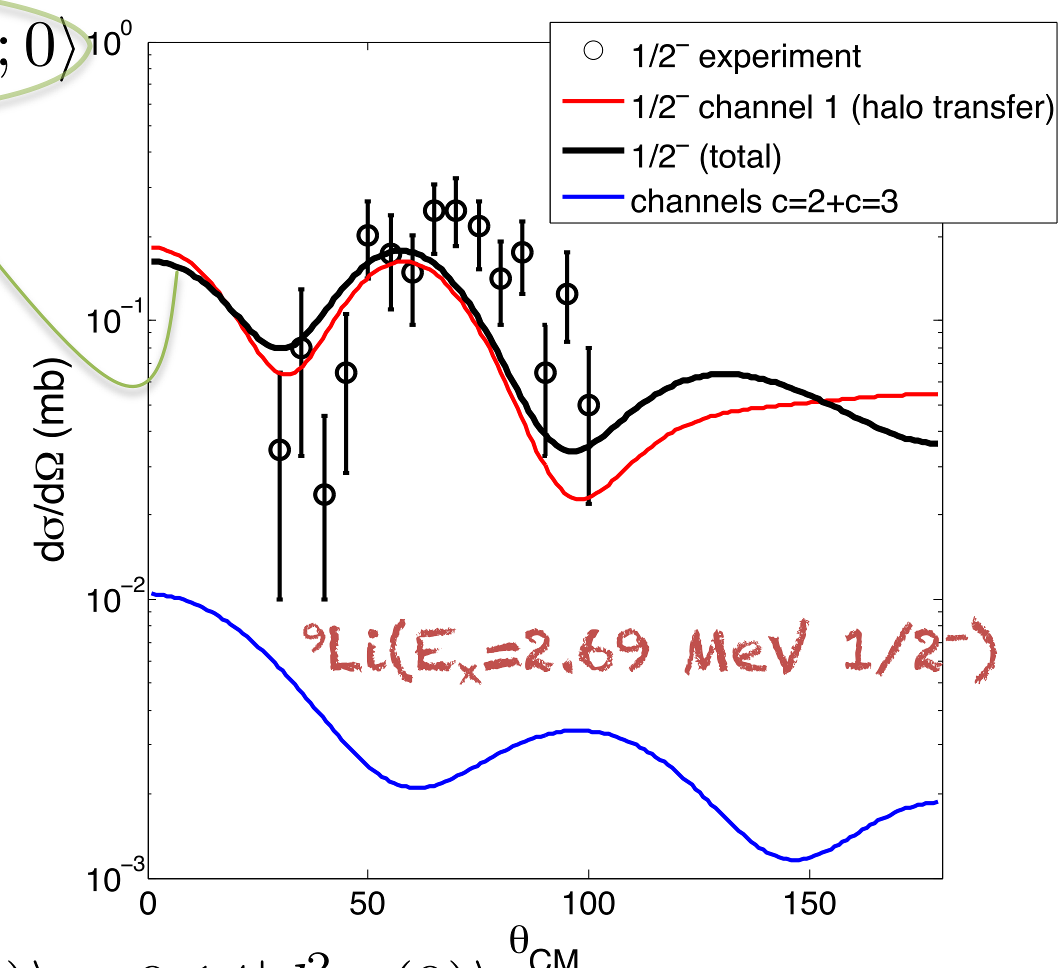
theory confirmed by ^{11}Li(p,t)^{9}Li(gs;E$_x$=2.69 MeV 1/2$^-$)

$$0.45|s^2_{1/2}(0)\rangle + 0.55|p^2_{1/2}(0)\rangle + 0.04|d^2_{5/2}(0)\rangle$$
$$+\,0.7\,|(p_{1/2}, s_{1/2})_{1^-} \otimes 1^-; 0\rangle$$
$$+\,0.1\,|(s_{1/2}, d_{5/2})_{2^+} \otimes 2^+; 0\rangle$$

$|s^2_{1/2}(0)\rangle$



depletion of pp strength due to coupling with PDR



$$0.63|s^2_{1/2}(0)\rangle + 0.76|p^2_{1/2}(0)\rangle + 0.14|d^2_{1/2}(0)\rangle$$

we compute the ^{11}Li PDR structure in RPA

3 representative low-lying dipole RPA peaks

E=0.65 MeV

	i	j	X_{ij}	Y_{ij}
ν	$2s_{1/2}$	$1p_{1/2}$	−0.780	0.078
ν	$3s_{1/2}$	$1p_{1/2}$	0.479	0.108
ν	$4s_{1/2}$	$1p_{1/2}$	0.220	0.106
ν	$5s_{1/2}$	$1p_{1/2}$	0.144	0.093
ν	$6s_{1/2}$	$1p_{1/2}$	0.106	0.080
ν	$1p_{3/2}$	$4d_{5/2}$	0.166	0.139
ν	$1p_{3/2}$	$5d_{5/2}$	0.241	0.208
ν	$1p_{3/2}$	$6d_{5/2}$	0.250	0.221
ν	$1p_{3/2}$	$7d_{5/2}$	0.199	0.180
ν	$1p_{3/2}$	$8d_{5/2}$	0.148	0.135
ν	$1p_{3/2}$	$9d_{5/2}$	0.110	0.102
ν	$1p_{1/2}$	$4d_{3/2}$	0.103	0.075
ν	$1p_{1/2}$	$5d_{3/2}$	0.119	0.095
ν	$1p_{1/2}$	$6d_{3/2}$	0.128	0.108
ν	$1p_{1/2}$	$7d_{3/2}$	0.128	0.112
ν	$1p_{1/2}$	$8d_{3/2}$	0.117	0.106
π	$2s_{1/2}$	$1p_{3/2}$	−0.136	−0.131
π	$1p_{3/2}$	$1d_{5/2}$	0.337	0.322

E=1.21 MeV

	i	j	X_{ij}	Y_{ij}
ν	$2s_{1/2}$	$1p_{1/2}$	−0.119	0.048
ν	$3s_{1/2}$	$1p_{1/2}$	−0.748	0.074
ν	$4s_{1/2}$	$1p_{1/2}$	0.410	0.080
ν	$5s_{1/2}$	$1p_{1/2}$	0.181	0.075
ν	$6s_{1/2}$	$1p_{1/2}$	0.117	0.067
ν	$1p_{3/2}$	$4d_{5/2}$	0.170	0.121
ν	$1p_{3/2}$	$5d_{5/2}$	0.243	0.183
ν	$1p_{3/2}$	$6d_{5/2}$	0.249	0.196
ν	$1p_{3/2}$	$7d_{5/2}$	0.196	0.161
ν	$1p_{3/2}$	$8d_{5/2}$	0.144	0.122
ν	$1p_{3/2}$	$9d_{5/2}$	0.107	0.093
ν	$1p_{1/2}$	$2d_{3/2}$	0.168	0.024
ν	$1p_{1/2}$	$3d_{3/2}$	0.114	0.043
ν	$1p_{1/2}$	$4d_{3/2}$	0.117	0.063
ν	$1p_{1/2}$	$5d_{3/2}$	0.126	0.081
ν	$1p_{1/2}$	$6d_{3/2}$	0.131	0.094
ν	$1p_{1/2}$	$7d_{3/2}$	0.128	0.099
ν	$1p_{1/2}$	$8d_{3/2}$	0.116	0.094
π	$2s_{1/2}$	$1p_{3/2}$	−0.130	−0.12
π	$1p_{3/2}$	$1d_{5/2}$	0.322	0.294

E=2 MeV

	i	j	X_{ij}	Y_{ij}
ν	$3s_{1/2}$	$1p_{1/2}$	−0.118	0.040
ν	$4s_{1/2}$	$1p_{1/2}$	−0.821	0.046
ν	$5s_{1/2}$	$1p_{1/2}$	0.250	0.046
ν	$6s_{1/2}$	$1p_{1/2}$	0.116	0.043
ν	$1p_{3/2}$	$4d_{5/2}$	0.144	0.081
ν	$1p_{3/2}$	$5d_{5/2}$	0.201	0.125
ν	$1p_{3/2}$	$6d_{5/2}$	0.201	0.135
ν	$1p_{3/2}$	$7d_{5/2}$	0.156	0.112
ν	$1p_{3/2}$	$8d_{5/2}$	0.113	0.085
ν	$1p_{1/2}$	$9d_{3/2}$	−0.126	0.014
ν	$1p_{1/2}$	$10d_{3/2}$	0.187	0.026
ν	$1p_{1/2}$	$11d_{3/2}$	0.121	0.040
ν	$1p_{1/2}$	$12d_{3/2}$	0.113	0.053
ν	$1p_{1/2}$	$13d_{3/2}$	0.111	0.064
ν	$1p_{1/2}$	$14d_{3/2}$	0.104	0.068
π	$1p_{3/2}$	$1d_{5/2}$	0.245	0.210

IOP Publishing Physica Scripta
Phys. Scr. **94** (2019) 114002 (18pp) https://doi.org/10.1088/1402-4896/ab2431

Pygmy resonances: what's in a name?

R A Broglia[1,2,7], F Barranco[3], A Idini[4], G Potel[5] and E Vigezzi[6]

we compute the ^{11}Li PDR structure in RPA

3 representative low-lying dipole RPA peaks

E=0.65 MeV

	i	j	X_{ij}	Y_{ij}
ν	$2s_{1/2}$	$1p_{1/2}$	−0.780	0.078
ν	$3s_{1/2}$	$1p_{1/2}$	0.479	0.108
ν	$4s_{1/2}$	$1p_{1/2}$	0.220	0.106
ν	$5s_{1/2}$	$1p_{1/2}$	0.144	0.093
ν	$6s_{1/2}$	$1p_{1/2}$	0.106	0.080
ν	$1p_{3/2}$	$4d_{5/2}$	0.166	0.139
ν	$1p_{3/2}$	$5d_{5/2}$	0.241	0.208
ν	$1p_{3/2}$	$6d_{5/2}$	0.250	0.221
ν	$1p_{3/2}$	$7d_{5/2}$	0.199	0.180
ν	$1p_{3/2}$	$8d_{5/2}$	0.148	0.135
ν	$1p_{3/2}$	$9d_{5/2}$	0.110	0.102
ν	$1p_{1/2}$	$4d_{3/2}$	0.103	0.075
ν	$1p_{1/2}$	$5d_{3/2}$	0.119	0.095
ν	$1p_{1/2}$	$6d_{3/2}$	0.128	0.108
ν	$1p_{1/2}$	$7d_{3/2}$	0.128	0.112
ν	$1p_{1/2}$	$8d_{3/2}$	0.117	0.106
π	$2s_{1/2}$	$1p_{3/2}$	−0.136	−0.131
π	$1p_{3/2}$	$1d_{5/2}$	0.337	0.322

E=1.21 MeV

	i	j	X_{ij}	Y_{ij}
ν	$2s_{1/2}$	$1p_{1/2}$	−0.119	0.048
ν	$3s_{1/2}$	$1p_{1/2}$	−0.748	0.074
ν	$4s_{1/2}$	$1p_{1/2}$	0.410	0.080
ν	$5s_{1/2}$	$1p_{1/2}$	0.181	0.075
ν	$6s_{1/2}$	$1p_{1/2}$	0.117	0.067
ν	$1p_{3/2}$	$4d_{5/2}$	0.170	0.121
ν	$1p_{3/2}$	$5d_{5/2}$	0.243	0.183
ν	$1p_{3/2}$	$6d_{5/2}$	0.249	0.196
ν	$1p_{3/2}$	$7d_{5/2}$	0.196	0.161
ν	$1p_{3/2}$	$8d_{5/2}$	0.144	0.122
ν	$1p_{3/2}$	$9d_{5/2}$	0.107	0.093
ν	$1p_{1/2}$	$2d_{3/2}$	0.168	0.024
ν	$1p_{1/2}$	$3d_{3/2}$	0.114	0.043
ν	$1p_{1/2}$	$4d_{3/2}$	0.117	0.063
ν	$1p_{1/2}$	$5d_{3/2}$	0.126	0.081
ν	$1p_{1/2}$	$6d_{3/2}$	0.131	0.094
ν	$1p_{1/2}$	$7d_{3/2}$	0.128	0.099
ν	$1p_{1/2}$	$8d_{3/2}$	0.116	0.094
π	$2s_{1/2}$	$1p_{3/2}$	−0.130	−0.12
π	$1p_{3/2}$	$1d_{5/2}$	0.322	0.294

E=2 MeV

	i	j	X_{ij}	Y_{ij}
ν	$3s_{1/2}$	$1p_{1/2}$	−0.118	0.040
ν	$4s_{1/2}$	$1p_{1/2}$	−0.821	0.046
ν	$5s_{1/2}$	$1p_{1/2}$	0.250	0.046
ν	$6s_{1/2}$	$1p_{1/2}$	0.116	0.043
ν	$1p_{3/2}$	$4d_{5/2}$	0.144	0.081
ν	$1p_{3/2}$	$5d_{5/2}$	0.201	0.125
ν	$1p_{3/2}$	$6d_{5/2}$	0.201	0.135
ν	$1p_{3/2}$	$7d_{5/2}$	0.156	0.112
ν	$1p_{3/2}$	$8d_{5/2}$	0.113	0.085
ν	$1p_{1/2}$	$9d_{3/2}$	−0.126	0.014
ν	$1p_{1/2}$	$10d_{3/2}$	0.187	0.026
ν	$1p_{1/2}$	$11d_{3/2}$	0.121	0.040
ν	$1p_{1/2}$	$12d_{3/2}$	0.113	0.053
ν	$1p_{1/2}$	$13d_{3/2}$	0.111	0.064
ν	$1p_{1/2}$	$14d_{3/2}$	0.104	0.068
π	$1p_{3/2}$	$1d_{5/2}$	0.245	0.210

IOP Publishing Physica Scripta
Phys. Scr. 94 (2019) 114002 (18pp) https://doi.org/10.1088/1402-4896/ab2431

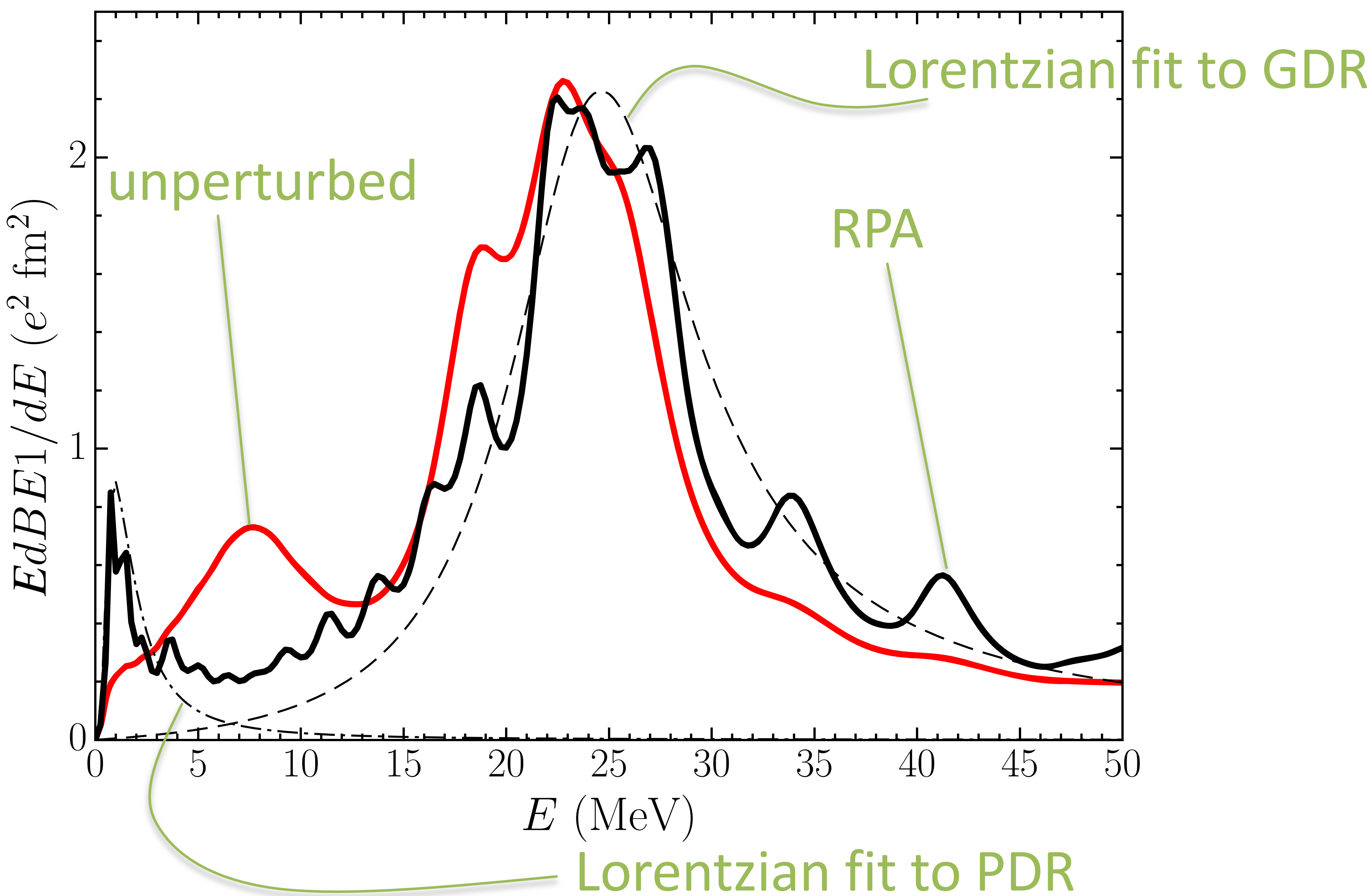
Pygmy resonances: what's in a name?

R A Broglia[1,2,7], F Barranco[3], A Idini[4], G Potel[5] and E Vigezzi[6]

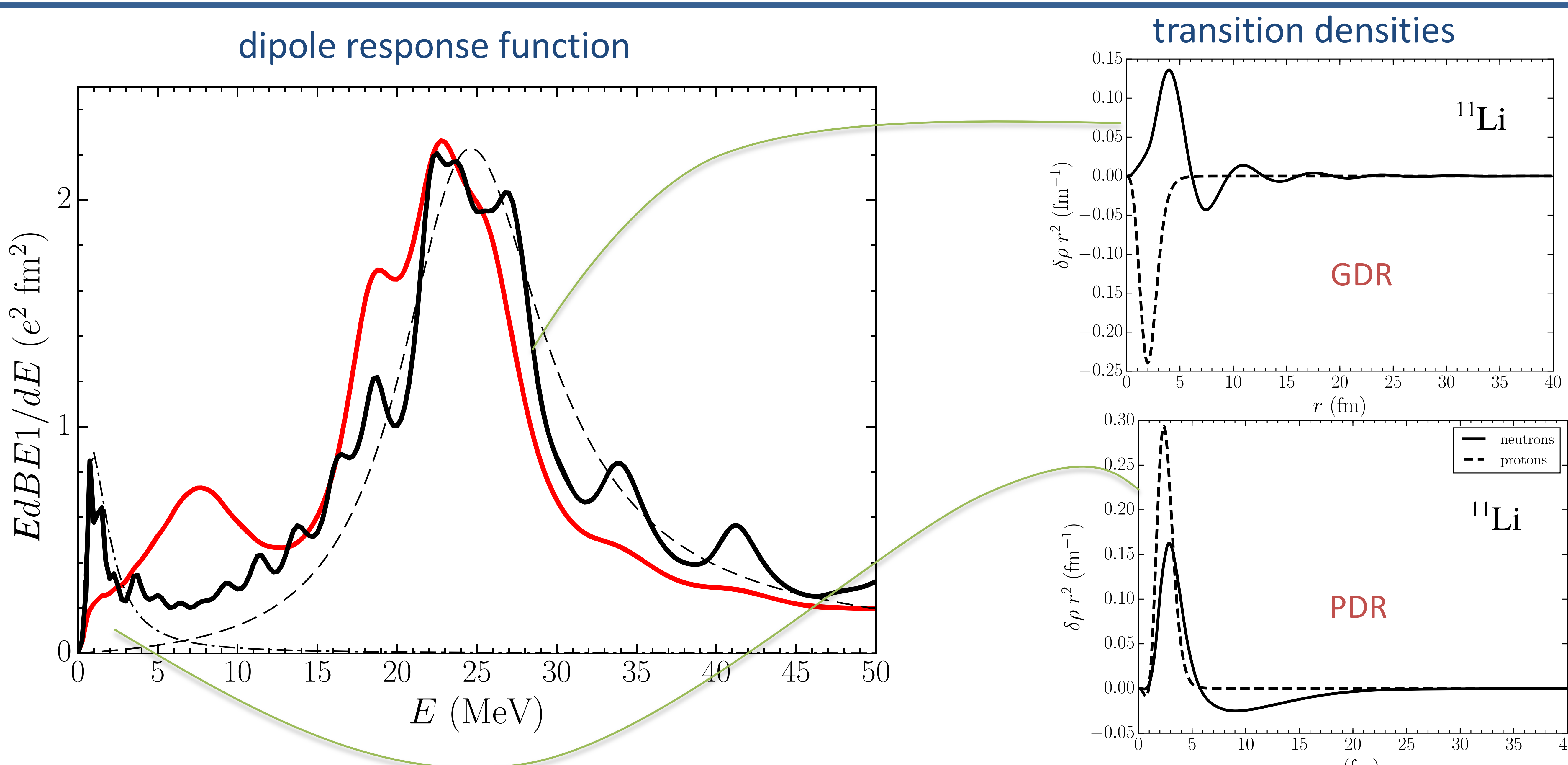
largest components are 2-quasiparticle neutron halo $(s_{1/2}\, p_{1/2})_{1^-}$ states

the PDR exhausts about 8% of the EWSR

dipole response function

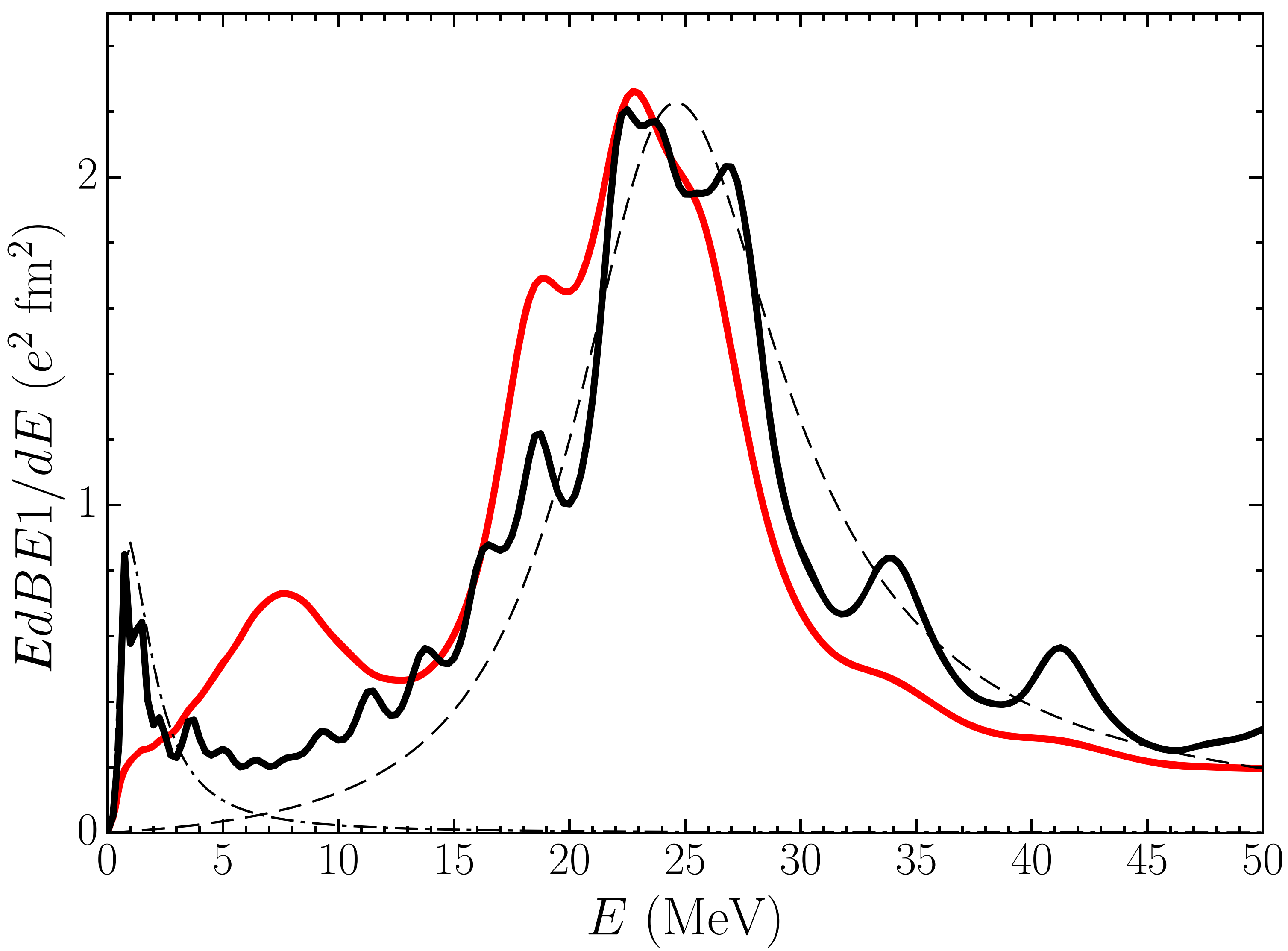


the PDR exhausts about 8% of the EWSR



the PDR exhausts about 8% of the EWSR

dipole response function



experiment approved at FRIB to probe the whole dipole response with (*p*,*p*').
Spokepersons: Ayyad, Zamora

Conclusions

- the PDR plays an important role in the structure of the exotic two-neutron halo nucleus ^{11}Li: halo-PDR symbiotic nature
- our calculations point to a strong ***pp*** component of the PDR, as opposed to the more ***ph*** nature of the GDR

complementary classification of dipole modes

isovector	⟺	isoscalar
ph	⟺	pp
GDR?	⟺	PDR?

talks by Vandebrouck, Spieker, Weinert, Khumalo

- PDR of ^{11}Li as a vortical excitation of the halo. Extrapolable to neutron skins?
- Approved experiments: ^{11}Li(p,p')^{11}Li* @ FRIB, and ^{9}Li(t,p)^{11}Li(PDR) @ ISOLDE

- along with (d,p) and (n,n'), (t,p) to join the ranks of novel probes to the PDR
- personal wish: (t,p) measurements on nuclei with neutron skin. Maybe with new FSU triton beam?

we compute the ^{11}Li PDR structure and the ^{9}Li $(t,p)^{11}$Li(PDR) cross section

3 representative low-lying dipole RPA peaks

E=0.65 MeV

	i	j	X_{ij}	Y_{ij}
ν	$2s_{1/2}$	$1p_{1/2}$	-0.780	0.078
ν	$3s_{1/2}$	$1p_{1/2}$	0.479	0.108
ν	$4s_{1/2}$	$1p_{1/2}$	0.220	0.106
ν	$5s_{1/2}$	$1p_{1/2}$	0.144	0.093
ν	$6s_{1/2}$	$1p_{1/2}$	0.106	0.080
ν	$1p_{3/2}$	$4d_{5/2}$	0.166	0.139
ν	$1p_{3/2}$	$5d_{5/2}$	0.241	0.208
ν	$1p_{3/2}$	$6d_{5/2}$	0.250	0.221
ν	$1p_{3/2}$	$7d_{5/2}$	0.199	0.180
ν	$1p_{3/2}$	$8d_{5/2}$	0.148	0.135
ν	$1p_{3/2}$	$9d_{5/2}$	0.110	0.102
ν	$1p_{1/2}$	$4d_{3/2}$	0.103	0.075
ν	$1p_{1/2}$	$5d_{3/2}$	0.119	0.095
ν	$1p_{1/2}$	$6d_{3/2}$	0.128	0.108
ν	$1p_{1/2}$	$7d_{3/2}$	0.128	0.112
ν	$1p_{1/2}$	$8d_{3/2}$	0.117	0.106
π	$2s_{1/2}$	$1p_{3/2}$	-0.136	-0.131
π	$1p_{3/2}$	$1d_{5/2}$	0.337	0.322

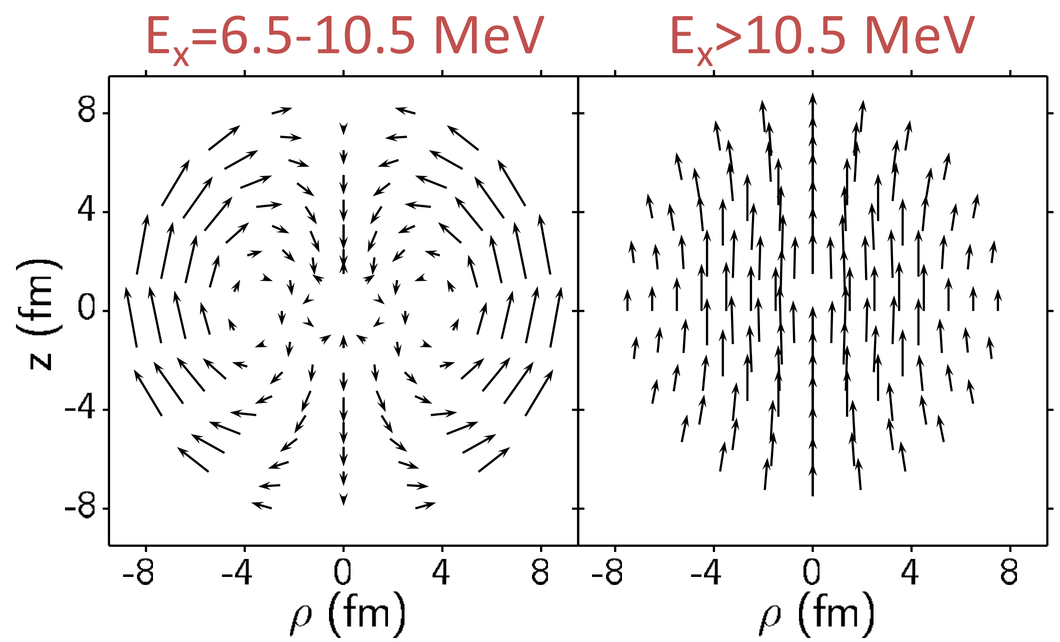
E=1.21 MeV

	i	j	X_{ij}	Y_{ij}
ν	$2s_{1/2}$	$1p_{1/2}$	-0.119	0.048
ν	$3s_{1/2}$	$1p_{1/2}$	-0.748	0.074
ν	$4s_{1/2}$	$1p_{1/2}$	0.410	0.080
ν	$5s_{1/2}$	$1p_{1/2}$	0.181	0.075
ν	$6s_{1/2}$	$1p_{1/2}$	0.117	0.067
ν	$1p_{3/2}$	$4d_{5/2}$	0.170	0.121
ν	$1p_{3/2}$	$5d_{5/2}$	0.243	0.183
ν	$1p_{3/2}$	$6d_{5/2}$	0.249	0.196
ν	$1p_{3/2}$	$7d_{5/2}$	0.196	0.161
ν	$1p_{3/2}$	$8d_{5/2}$	0.144	0.122
ν	$1p_{3/2}$	$9d_{5/2}$	0.107	0.093
ν	$1p_{1/2}$	$2d_{3/2}$	0.168	0.024
ν	$1p_{1/2}$	$3d_{3/2}$	0.114	0.043
ν	$1p_{1/2}$	$4d_{3/2}$	0.117	0.063
ν	$1p_{1/2}$	$5d_{3/2}$	0.126	0.081
ν	$1p_{1/2}$	$6d_{3/2}$	0.131	0.094
ν	$1p_{1/2}$	$7d_{3/2}$	0.128	0.099
ν	$1p_{1/2}$	$8d_{3/2}$	0.116	0.094
π	$2s_{1/2}$	$1p_{3/2}$	-0.130	-0.12
π	$1p_{3/2}$	$1d_{5/2}$	0.322	0.294

E=2 MeV

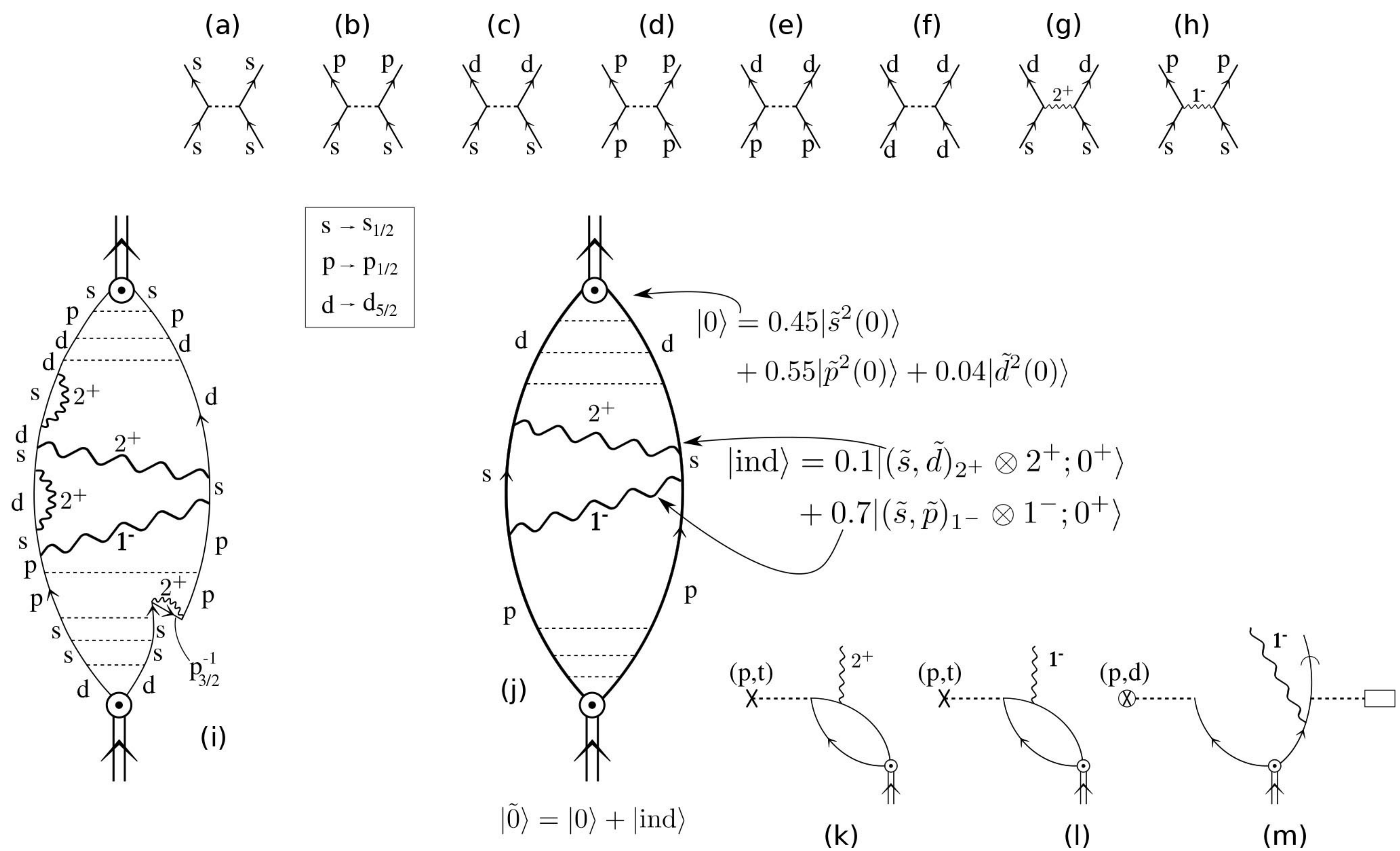
	i	j	X_{ij}	Y_{ij}
ν	$3s_{1/2}$	$1p_{1/2}$	-0.118	0.040
ν	$4s_{1/2}$	$1p_{1/2}$	-0.821	0.046
ν	$5s_{1/2}$	$1p_{1/2}$	0.250	0.046
ν	$6s_{1/2}$	$1p_{1/2}$	0.116	0.043
ν	$1p_{3/2}$	$4d_{5/2}$	0.144	0.081
ν	$1p_{3/2}$	$5d_{5/2}$	0.201	0.125
ν	$1p_{3/2}$	$6d_{5/2}$	0.201	0.135
ν	$1p_{3/2}$	$7d_{5/2}$	0.156	0.112
ν	$1p_{3/2}$	$8d_{5/2}$	0.113	0.085
ν	$1p_{1/2}$	$9d_{3/2}$	-0.126	0.014
ν	$1p_{1/2}$	$10d_{3/2}$	0.187	0.026
ν	$1p_{1/2}$	$11d_{3/2}$	0.121	0.040
ν	$1p_{1/2}$	$12d_{3/2}$	0.113	0.053
ν	$1p_{1/2}$	$13d_{3/2}$	0.111	0.064
ν	$1p_{1/2}$	$14d_{3/2}$	0.104	0.068
π	$1p_{3/2}$	$1d_{5/2}$	0.245	0.210

velocity field of ^{208}Pb dipole states



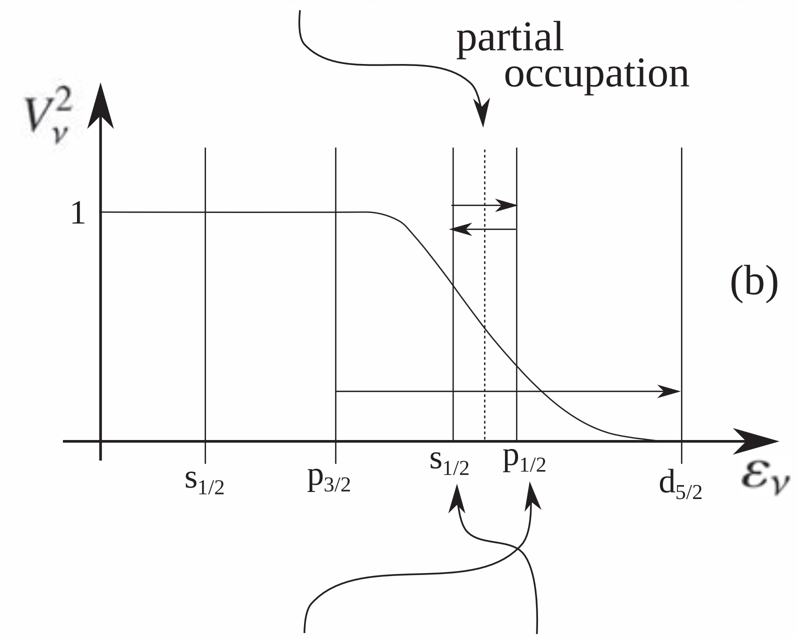
Ryezayeva *et al.* PRL **89** (2002) 272502

Ground state of ^{11}Li



Low-lying dipole strength

$$|0\rangle_\nu = |0\rangle + 0.7|(p_{1/2}, s_{1/2})_{1^-} \otimes 1^-; 0\rangle + 0.1|(s_{1/2}, d_{5/2})_{2^+} \otimes 2^+; 0\rangle$$



$$|0\rangle = 0.55|p^2_{1/2}\rangle + 0.45|s^2_{1/2}\rangle + 0.04|d^2_{5/2}\rangle$$

$$|1^-, \mathrm{pygmy}\rangle = \alpha\, \Gamma^\dagger_{\mathrm{pygmy}}|\mathrm{halo}\rangle + \beta\, \Gamma^\dagger_{\mathrm{GDR}}|\mathrm{core}\rangle$$

$\alpha^2 \gg \beta^2$

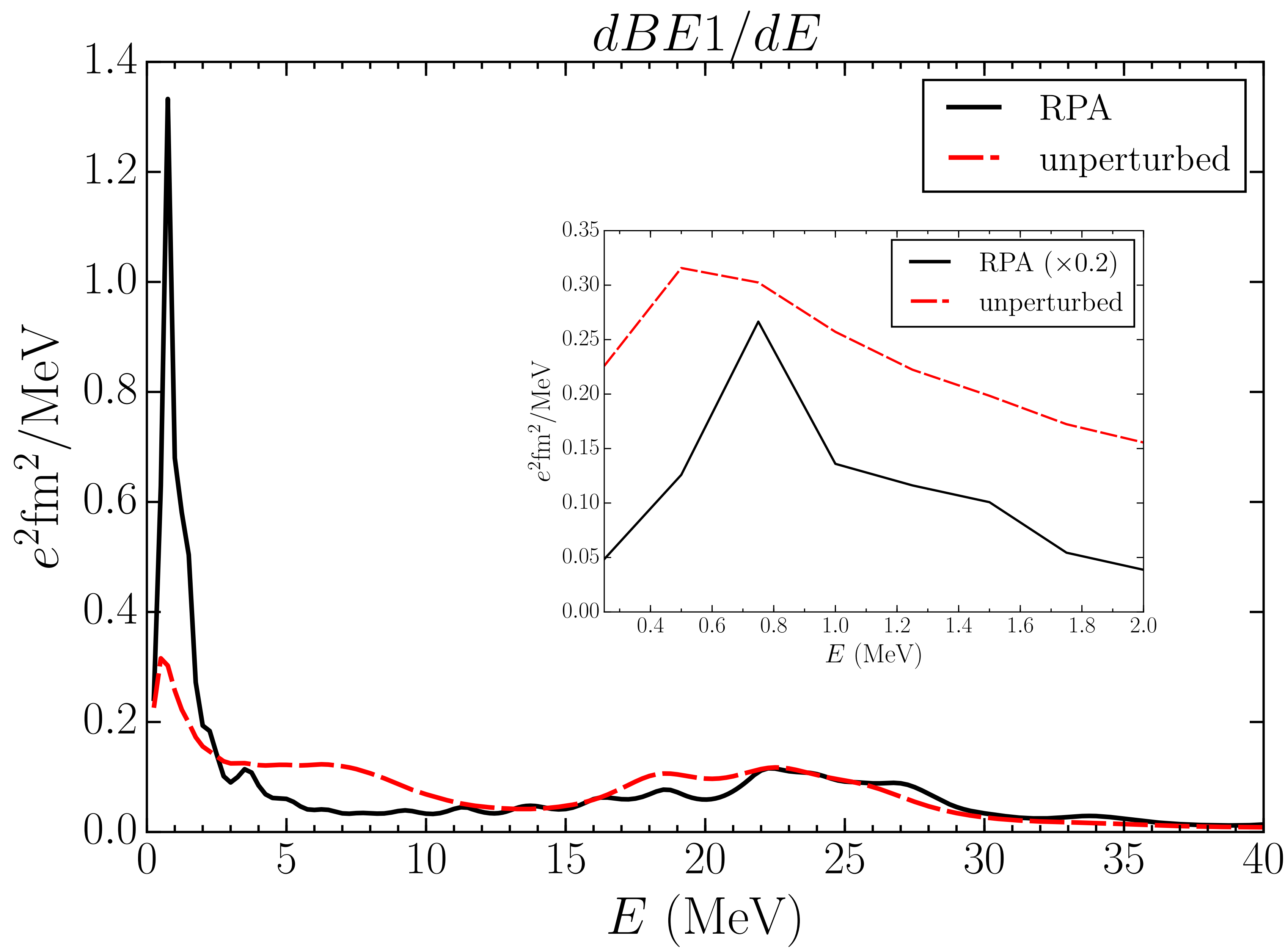
(a)

	$1p^{-1}_{1/2}2s_{1/2}$	$1p^{-1}_{1/2}3s_{1/2}$	$1p^{-1}_{1/2}4s_{1/2}$	$1p^{-1}_{1/2}1d_{3/2}$	$1p^{-1}_{3/2}4d_{5/2}$	$1p^{-1}_{3/2}5d_{5/2}$	$1p^{-1}_{3/2}6d_{5/2}$
X	-0.780	0.479	0.220	0.103	0.166	0.241	0.250
Y	0.078	0.108	0.106	0.075	0.139	0.208	0.221

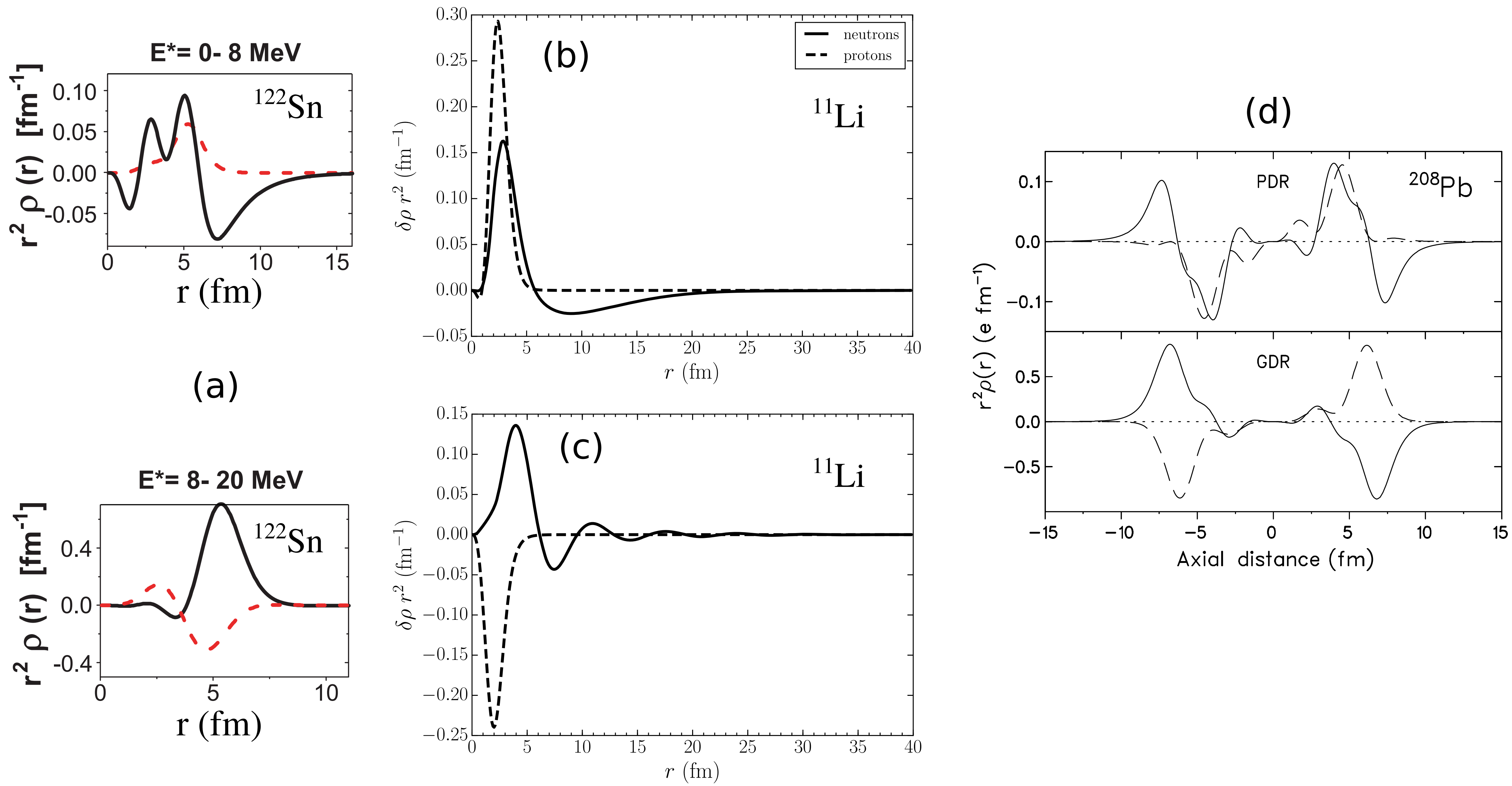
8% EWSR

$E_{1^-} \approx 0.7$ MeV

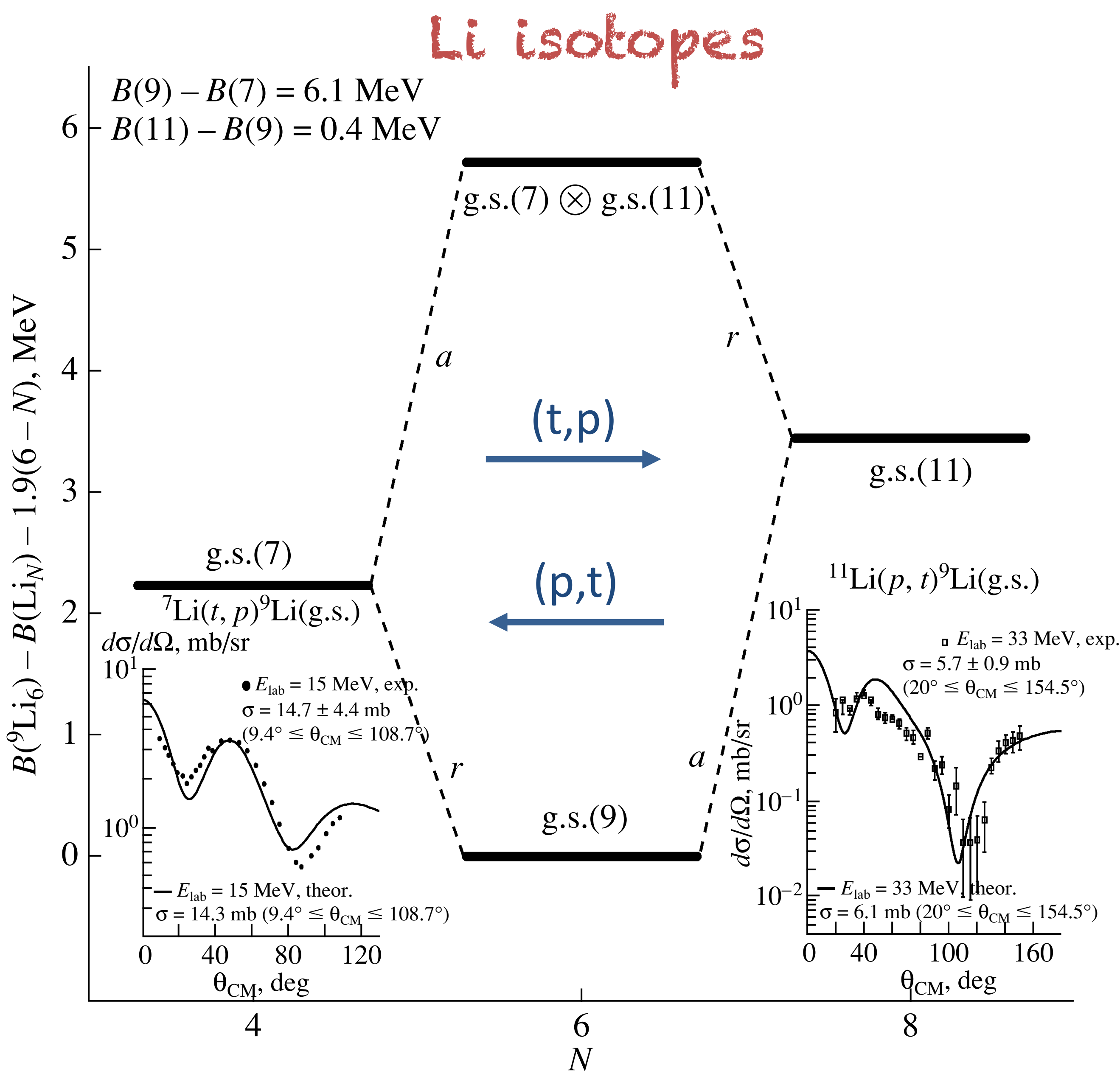
Full dipole strength

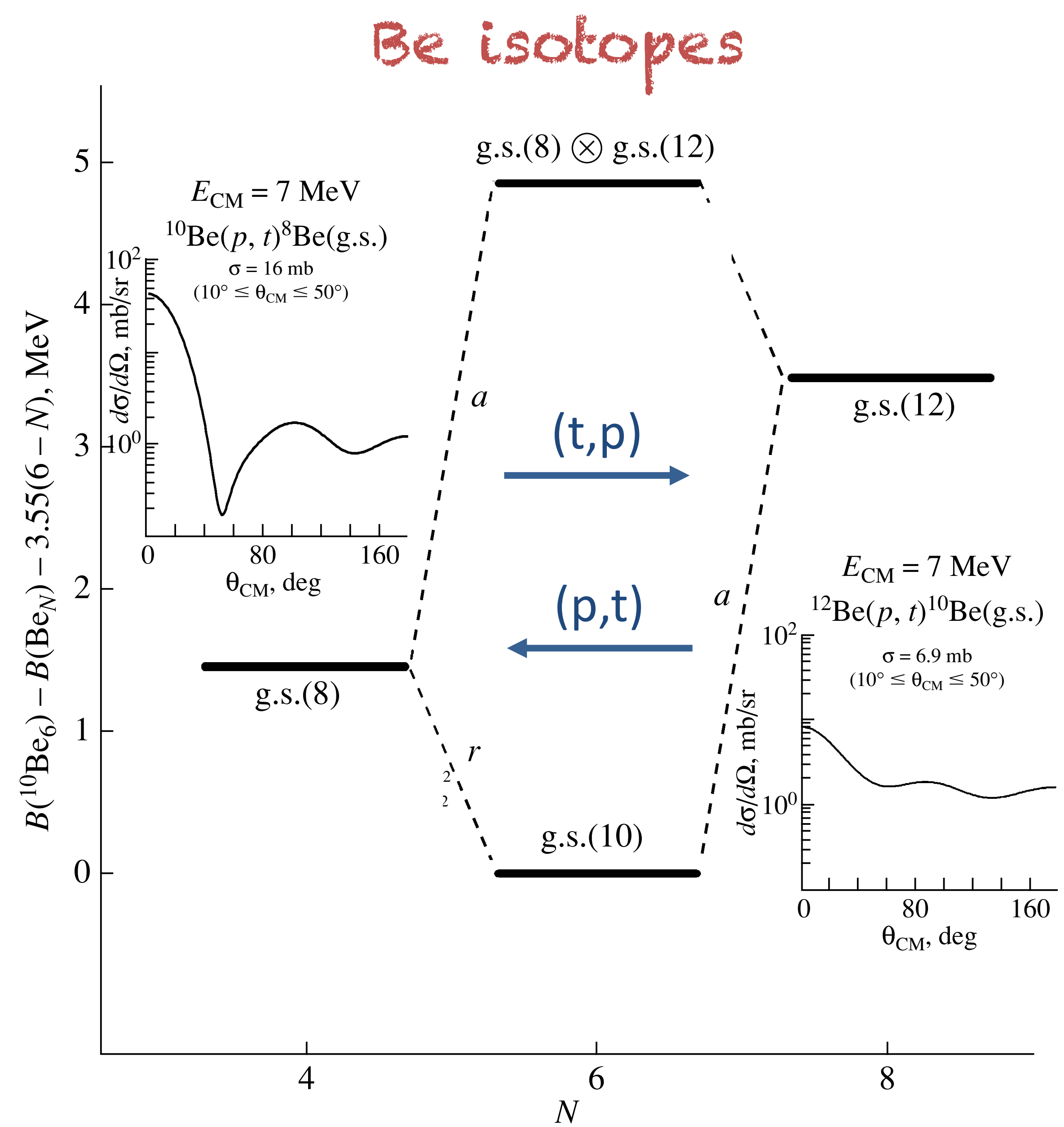


Transition densities

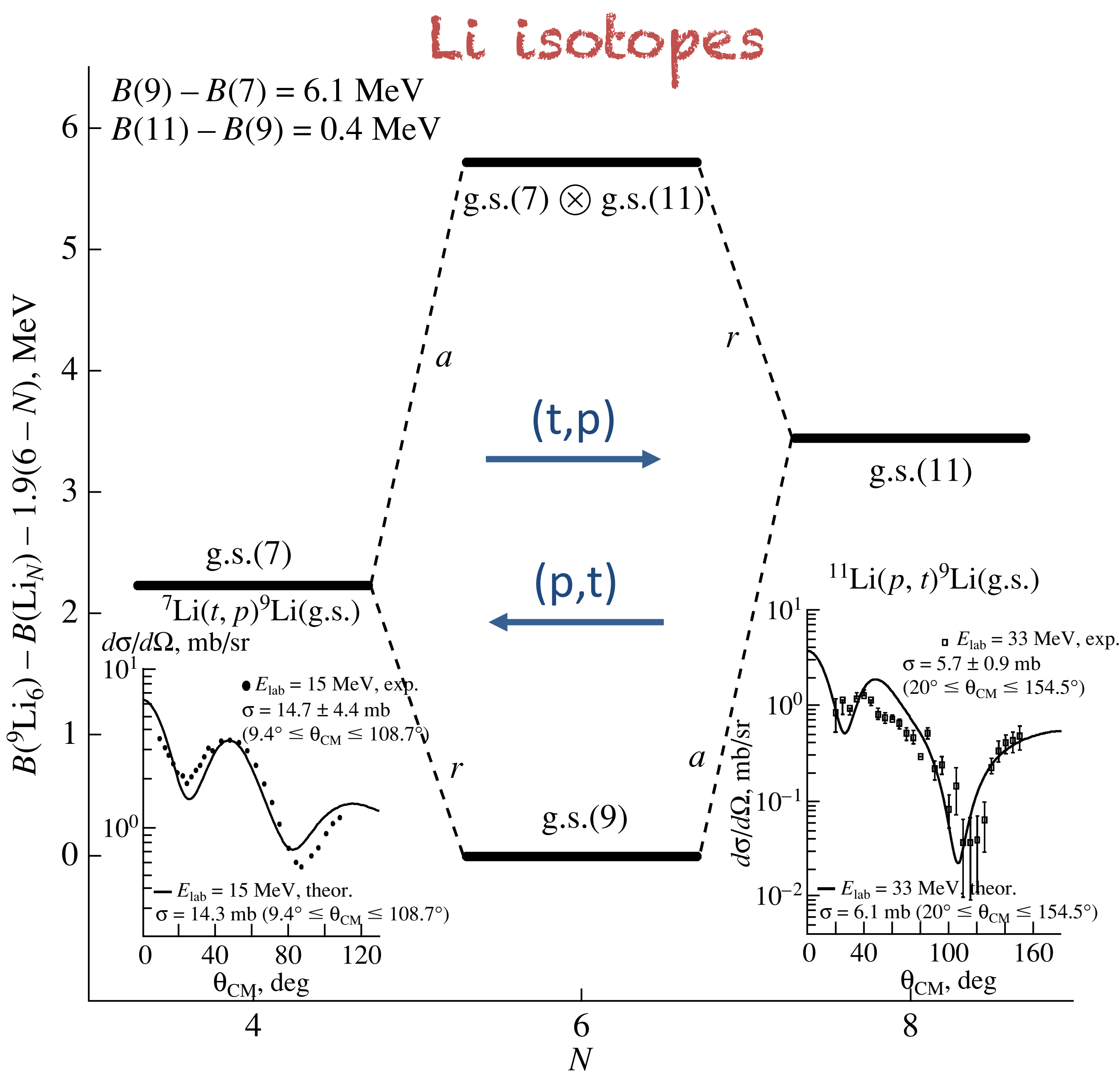


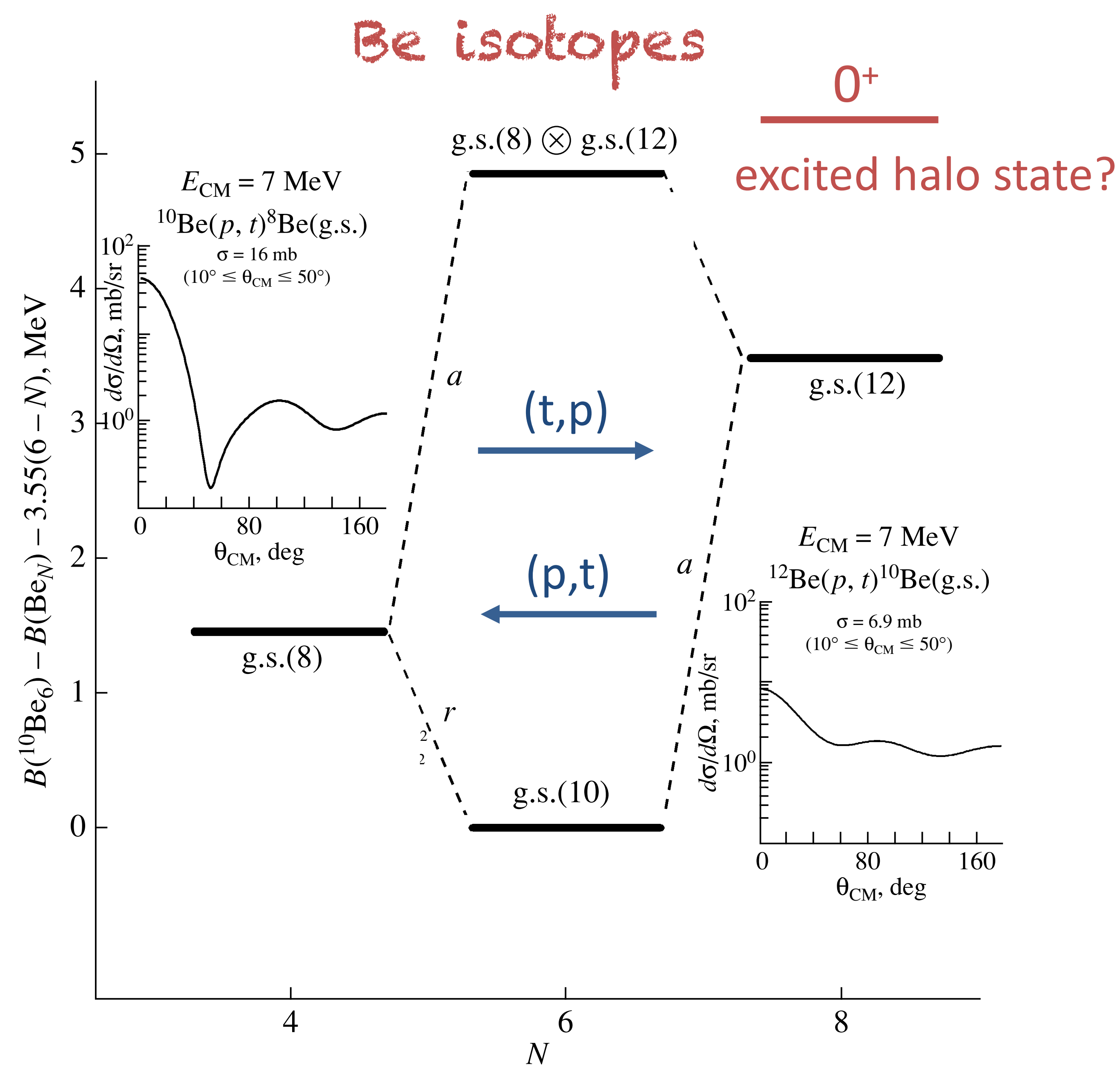
Excited halo state in ^{12}Be (0^+_2)



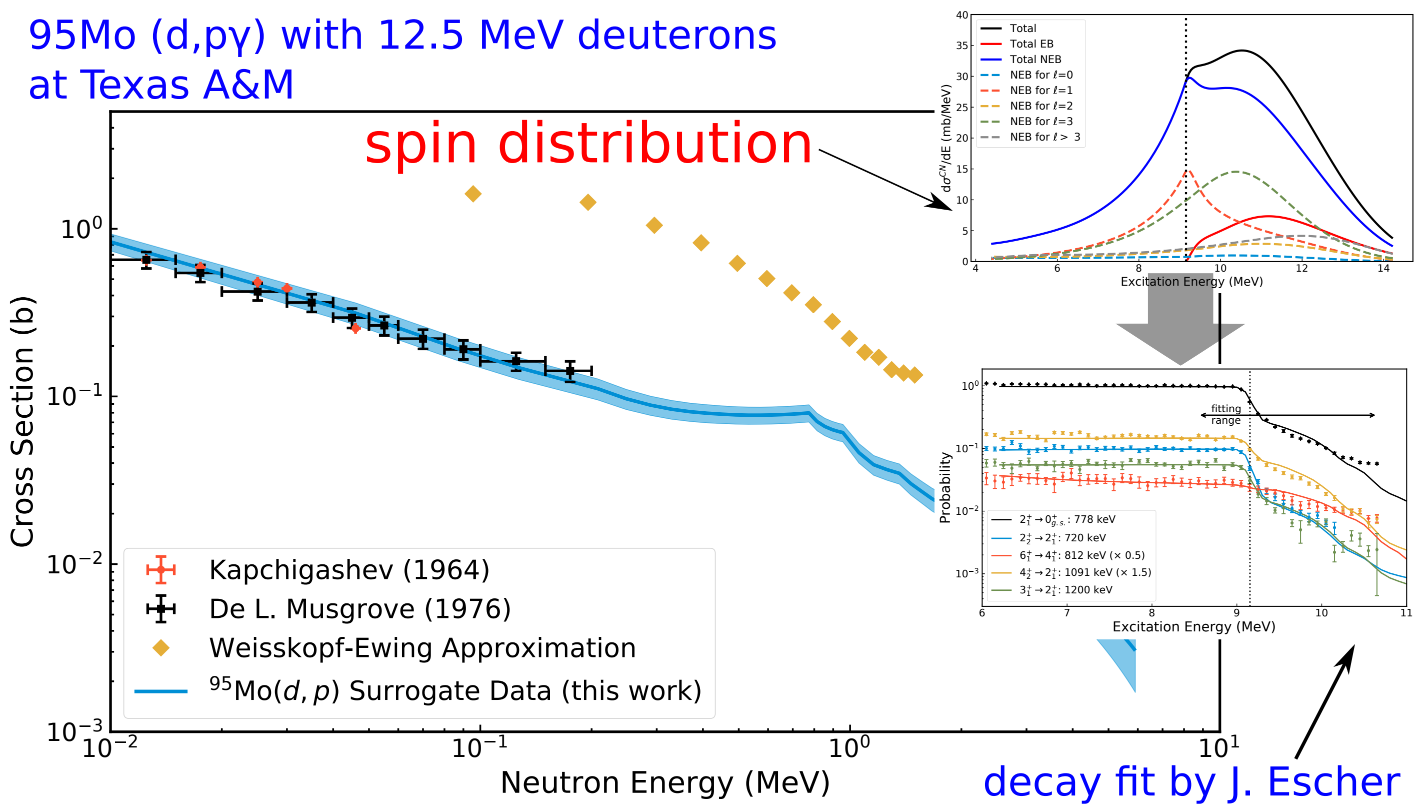


Excited halo state in ^{12}Be (0^+_2)





An opportunity to thoroughly benchmark the SRM with ^{95}Mo(n,γ)

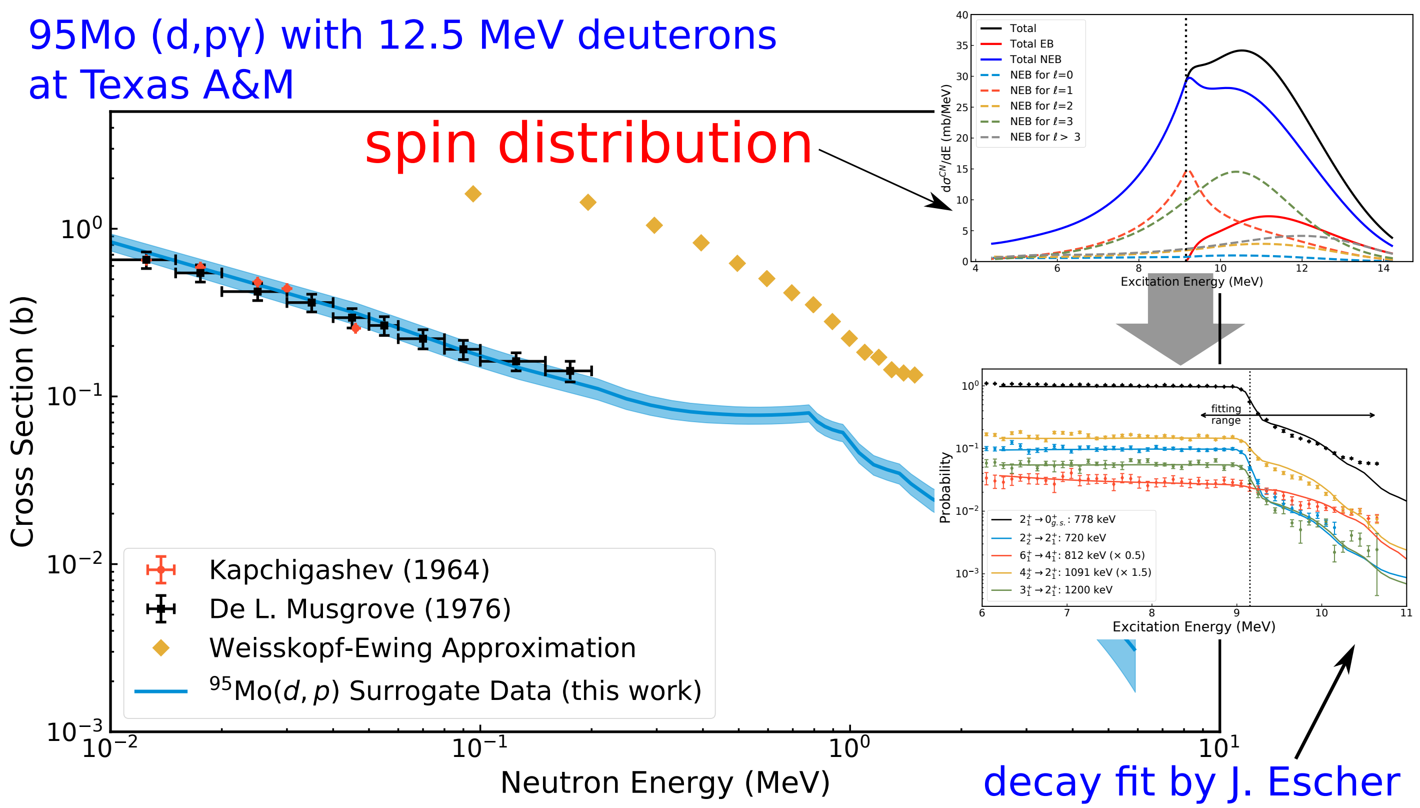


Ratkiewicz, Cizewski, Escher, GP, et al. Phys. Rev. Lett. 122052502 (2019)

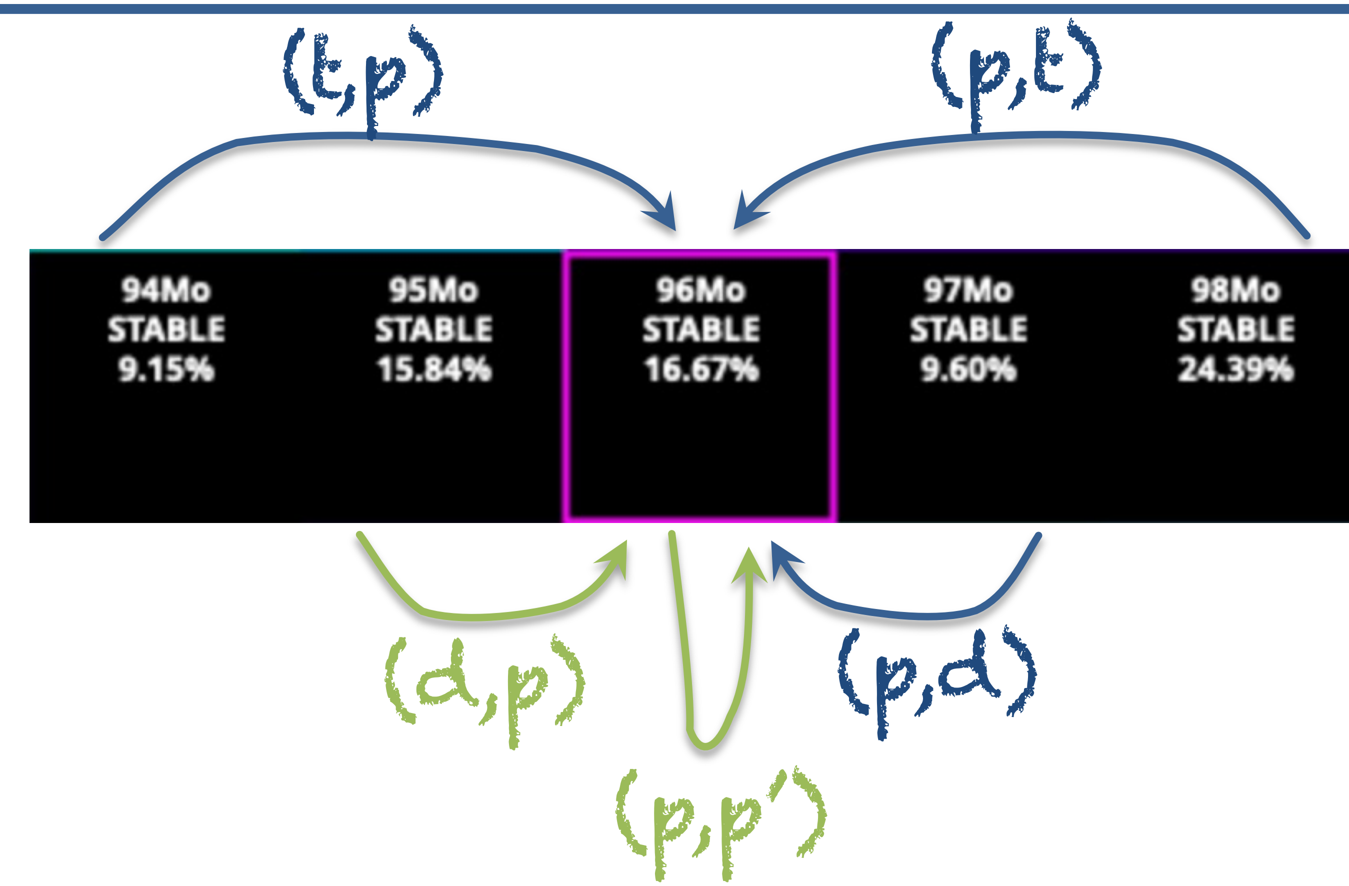
- Excellent agreement with (n,γ) data.
- The fitted Hauser-Feshbach decay is used to infer (n,γ) rates.

- No previous knowledge of D_0, and/or $\langle\Gamma_\gamma\rangle$ is needed.
- No need for separate determination of NLD and γSF.

An opportunity to thoroughly benchmark the SRM with ^{95}Mo(n,γ)



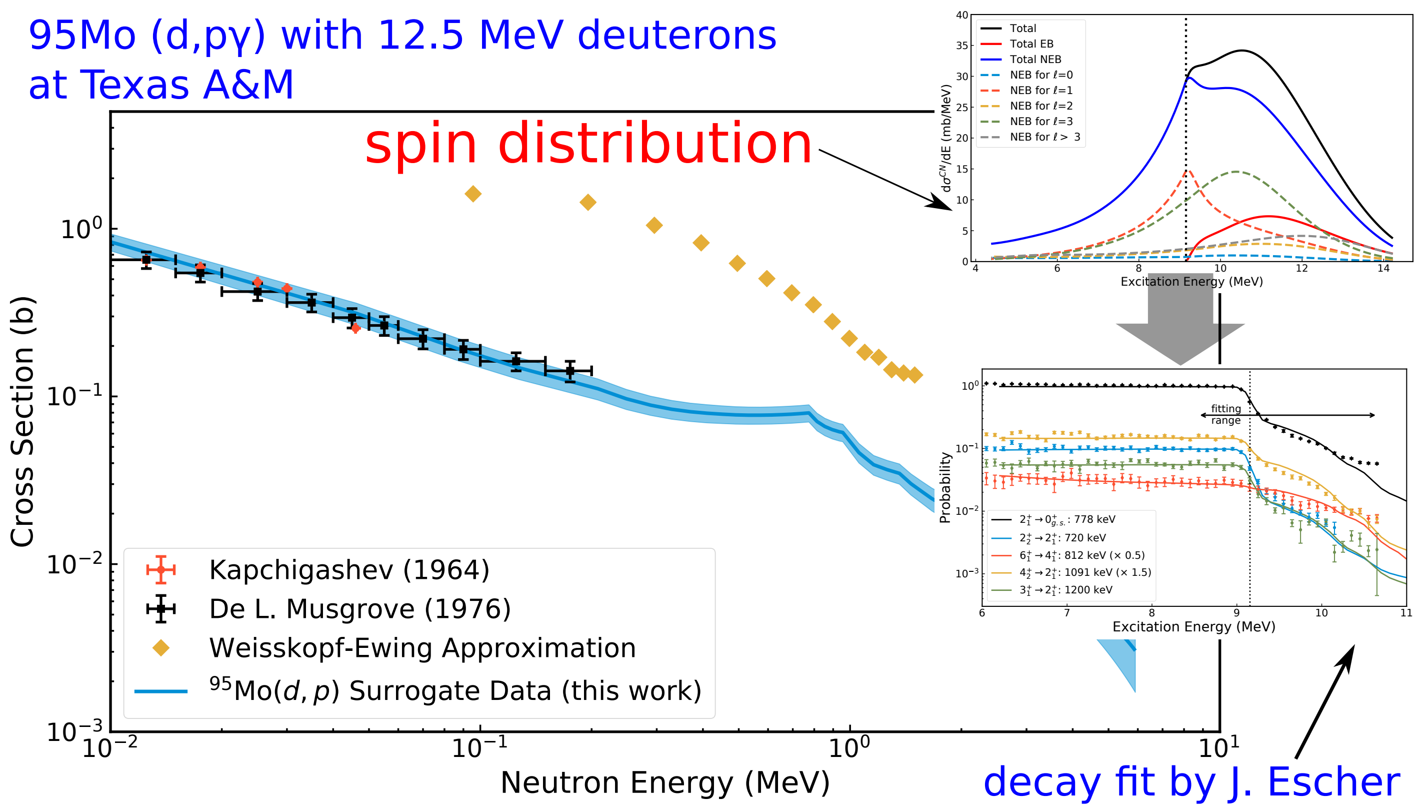
Ratkiewicz, Cizewski, Escher, GP, et al. Phys. Rev. Lett. 122052502 (2019)



- Excellent agreement with (n, γ) data.
- The fitted Hauser-Feshbach decay is used to infer (n, γ) rates.

- No previous knowledge of D_0, and/or $\langle \Gamma_\gamma \rangle$ is needed.
- No need for separate determination of NLD and γSF.

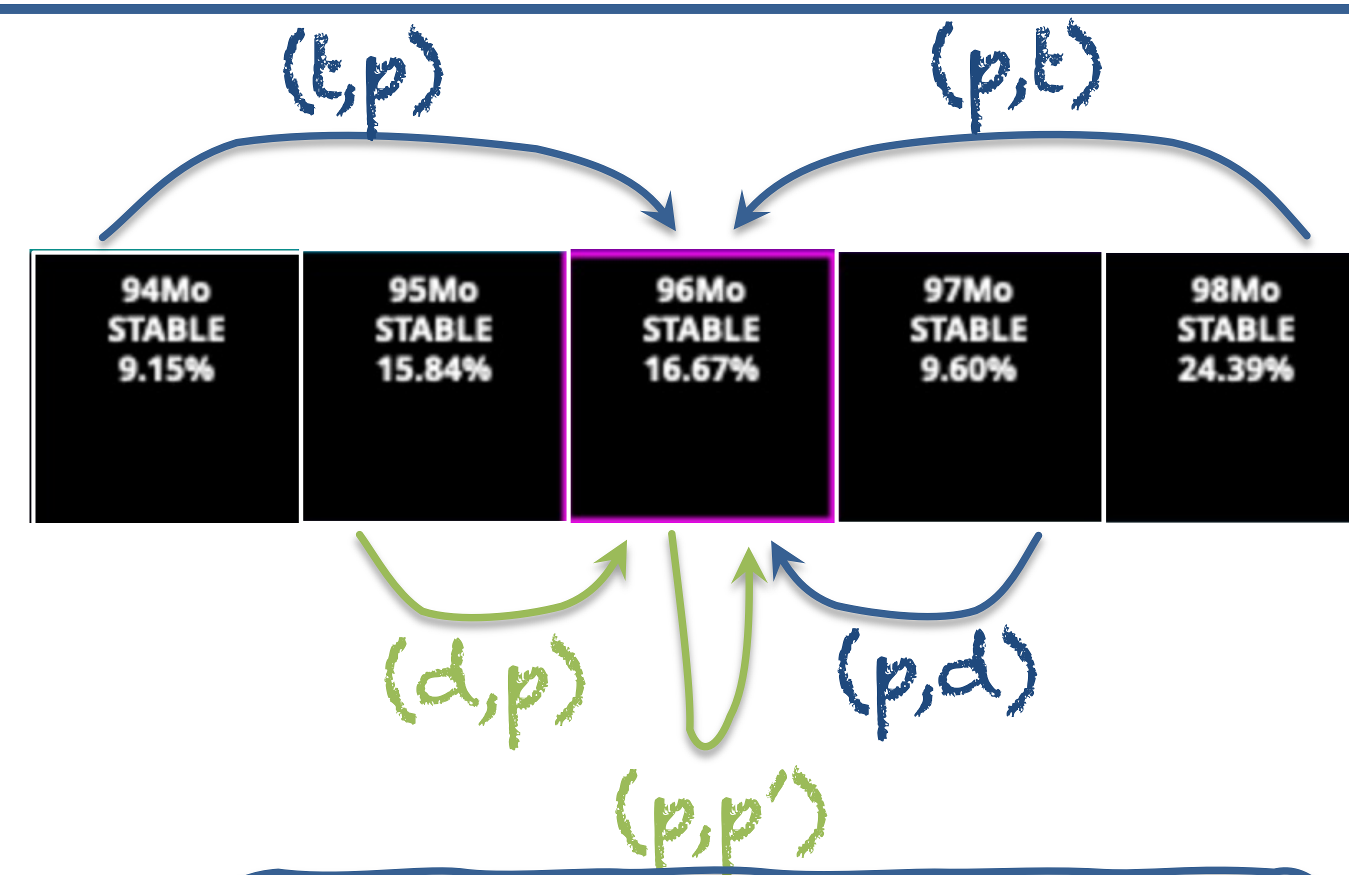
An opportunity to thoroughly benchmark the SRM with ^{95}Mo(n,γ)



Ratkiewicz, Cizewski, Escher, GP, et al. Phys. Rev. Lett. 122052502 (2019)

- Excellent agreement with (n, γ) data.
- The fitted Hauser-Feshbach decay is used to infer (n, γ) rates.

- No previous knowledge of D_0, and/or $\langle \Gamma_\gamma \rangle$ is needed.
- No need for separate determination of NLD and γSF.



- $^{94\text{-}98}$Mo are all stable
- ^{95}Mo(n,γ) is known
- ^{95}Mo(d,pγ) and ^{95}Mo(p,p'γ) have been measured
- Opportunity to benchmark many SRM techniques (A. McIntosh talk)