Università degli Studi di Firenze



Facoltà d' Ingegneria Dipartimento di Elettronica e Telecomunicazioni

Radar penetranti

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Radar penetranti



SPR (Surface Penetrating Radar)

GPR (Ground Penetrating Radar)

GeoRADAR

Radar penetranti



I radar penetranti impiegano onde e.m. per sondare materiali dielettrici fortemente attenuanti (lossy)

Frequenze impiegate: 1 MHz – 2GHz

Radar penetranti



Accessibilità solo da un lato

TX RX	TX	RX

Configurazione monostatica

Configurazione bistatica

Equazioni di Maxwell



$$\nabla \times \underline{E} = -\frac{\partial \underline{B}}{\partial t}$$

$$\nabla \times \nabla \times \underline{E} = -\frac{\partial}{\partial t} (\nabla \times \underline{B})$$

$$\nabla \times \underline{B} = \mu \underline{J} + \varepsilon_0 \mu \frac{\partial \underline{E}}{\partial t}$$

$$\nabla \cdot \underline{E} = \frac{\rho}{\varepsilon}$$

$$\nabla \cdot \underline{B} = 0$$

$$\nabla \times \nabla \times \underline{E} = -\frac{\partial}{\partial t} (\mu \underline{J} + \varepsilon_0 \mu \frac{\partial \underline{E}}{\partial t})$$

$$\nabla \times \nabla \times \underline{E} + \varepsilon_0 \mu \frac{\partial^2 \underline{E}}{\partial t^2} (\mu \underline{J} + \varepsilon_0 \mu \frac{\partial \underline{J}}{\partial t}) = 0$$
Equazione d'onda

Corrente nel mezzo

Correnti



$$J=J_{\it free}+J_{\it pol}$$

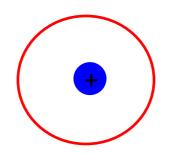
Correnti ohmiche

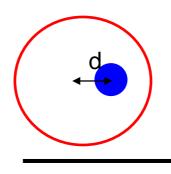
$$J_{free} = \sigma' E$$

Correnti



$$J = J_{\mathit{free}} + J_{\mathit{pol}}$$





$$D = qd$$

$$D = qd \qquad \frac{\partial D}{\partial t} = q \frac{\partial d}{\partial t} = qv = J_{pol}$$

$$D = \chi' E$$

$$J_{pol} = \frac{\partial D}{\partial t} = \chi' \frac{\partial E}{\partial t}$$

Equazioni di Maxwell

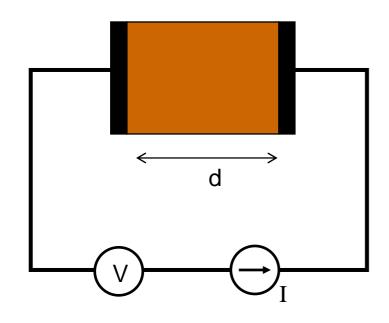


$$J = J_{free} + J_{pol} = \sigma' E + \chi' \frac{\partial E}{\partial t}$$

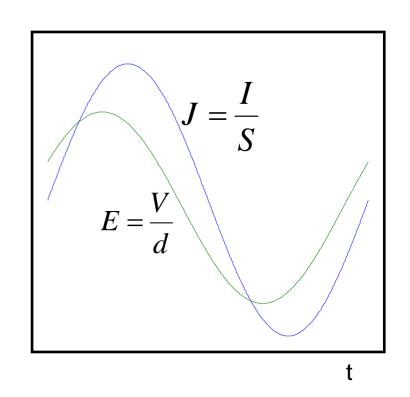
Per una sinusoide:

$$J=\sigma'E+\chi'j\omega E=\sigma E$$
 complesso





$$J = \sigma E$$
 $\sigma = \sigma' - j\sigma''$



Definiamo:

$$\chi = \frac{\sigma}{j\omega\varepsilon_0} \longrightarrow J = \sigma E = j\omega\varepsilon_0 \chi E = \varepsilon_0 \chi \frac{\partial E}{\partial t}$$

Equazioni di Maxwell

$$\nabla \times \nabla \times \underline{E} + \varepsilon_0 \mu \frac{\partial^2 \underline{E}}{\partial t^2} + \mu \frac{\partial \underline{J}}{\partial t} = 0 \qquad J = \varepsilon_0 \chi \frac{\partial E}{\partial t}$$

$$\varepsilon = \varepsilon_0 (1 + \chi)$$

$$\nabla \times \nabla \times \underline{E} + \mu \varepsilon \frac{\partial^2 E}{\partial t^2} = 0$$

$$\varepsilon = \varepsilon_0 (1 + \chi) = \varepsilon_0 \left(1 + \frac{\sigma}{j \omega \varepsilon_0} \right) = \varepsilon_0 + \frac{\sigma' + j \sigma''}{j \omega} = \left(\varepsilon_0 + \frac{\sigma''}{\omega} \right) - j \left(\frac{\sigma'}{\omega} \right)$$

$$\varepsilon = \varepsilon' - j\varepsilon''$$

$$\varepsilon' = \varepsilon_0 + \frac{\sigma''}{\omega}$$

$$\varepsilon'' = \frac{\sigma'}{\omega}$$

Equazioni di Maxwell



$$\nabla \times \nabla \times \underline{E} + \varepsilon \mu \frac{\partial^2 \underline{E}}{\partial t^2} = 0$$

$$\frac{\partial^2 E}{\partial x} + \varepsilon \mu \frac{\partial^2 E}{\partial t^2} = 0$$

$$E(x,t) = E(x)e^{j\omega t}$$

$$\frac{\partial^2 E(x)}{\partial x} e^{j\omega t} + \varepsilon \mu (j\omega)^2 E(x) e^{j\omega t} = 0$$

$$\frac{\partial^2 E(x)}{\partial x} - \varepsilon \mu \omega^2 E(x) = 0 \qquad k = \omega \sqrt{\mu \varepsilon}$$

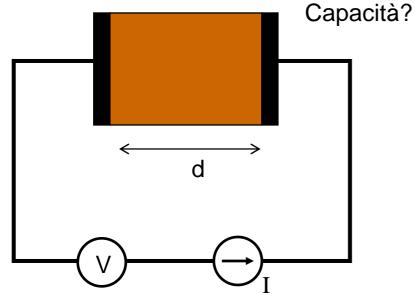
$$k = \omega \sqrt{\mu \varepsilon}$$

$$\frac{\partial^2 E(x)}{\partial x} - k^2 E(x) = 0 \quad \longrightarrow E(x) = E_0 e^{\pm jkx} \longrightarrow E(x) = E_0 e^{-jkx}$$

$$E(x,t) = E_0 e^{j(\omega t - kx)}$$

Mezzi dispersivi





$$J = \frac{I}{S}$$

$$E = \frac{V}{d}$$

$$J = \sigma E$$

Conducibilità complessa $\sigma = \sigma' - j\sigma''$

$$\sigma = \sigma' - j\sigma''$$

Permittività complessa

$$\varepsilon = \varepsilon' - j\varepsilon''$$

$$\varepsilon' = \varepsilon_0 + \frac{\sigma''}{\omega}$$

$$\varepsilon'' = \frac{\sigma'}{\omega}$$

$$\varepsilon'' = \frac{\sigma'}{\omega}$$



$$E(x,t) = E_0 e^{j(\omega t - kx)} \qquad k = \omega \sqrt{\mu (\varepsilon' - j\varepsilon'')}$$

$$jk = \alpha + j\beta$$
 $E(x,t) = E_0 e^{j(\omega t - \beta x)} e^{-\alpha x}$

$$\frac{\varepsilon''}{\varepsilon'} = \gamma$$

$$\beta = \omega \left[\frac{\mu \varepsilon'}{2} \left(\sqrt{1 + \gamma^2} + 1 \right) \right]^{\frac{1}{2}}$$

$$\alpha = \omega \left[\frac{\mu \varepsilon'}{2} \left(\sqrt{1 + \gamma^2} - 1 \right) \right]^{\frac{1}{2}}$$



$$e^{-j\beta x} \qquad \gamma = \frac{\varepsilon''}{\varepsilon'} << 1$$

$$\beta = \omega \left[\frac{\mu \varepsilon'}{2} \left(\sqrt{1 + \gamma'^2} + 1 \right) \right]^{\frac{1}{2}} = \omega \left[\frac{\mu \varepsilon'}{2} 2 \right]^{\frac{1}{2}} = \omega \sqrt{\mu \varepsilon'}$$

$$\varepsilon_r = \frac{\varepsilon'}{\varepsilon_0}$$
 $c = \frac{1}{\sqrt{\mu \varepsilon_0}}$

$$\beta = \omega \frac{\sqrt{\varepsilon_r}}{c} = \frac{\omega}{v}$$

$$v = \frac{C}{\sqrt{\mathcal{E}_r}}$$



$$e^{-\alpha x} \qquad \gamma = \frac{\varepsilon''}{\varepsilon'} << 1$$

$$\alpha = \omega \left[\frac{\mu \varepsilon'}{2} \left(\sqrt{1 + \gamma'^2} - 1 \right) \right]^{\frac{1}{2}} = \omega \frac{\sqrt{\mu \varepsilon'}}{\sqrt{2}} \left(1 + \frac{1}{2} \gamma^2 - 1 \right)^{\frac{1}{2}} = \frac{\omega \sqrt{\mu \varepsilon'}}{2} \gamma$$

$$\alpha = \frac{\omega}{2v}\gamma$$

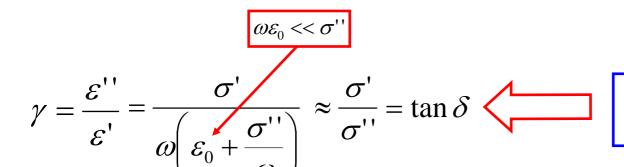


$$E(x,t) = E_0 e^{j\omega\left(t - \frac{1}{v}x\right)} e^{-\frac{\omega\gamma}{2v}x} \qquad v = \frac{c}{\sqrt{\varepsilon_r}} \qquad \gamma = \frac{\varepsilon''}{\varepsilon'}$$

La tangente di perdita

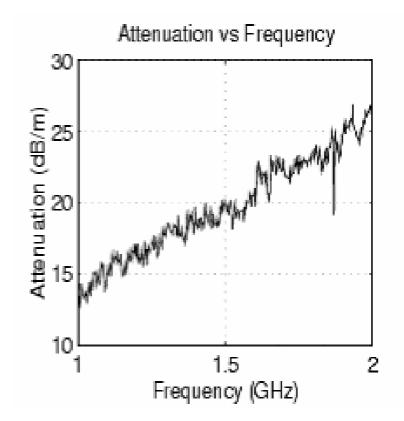


 $\tan \delta = \frac{1}{2}$



Costante [100MHz-2GHz]

$$\alpha = \frac{\omega \gamma}{2v} = \frac{\omega}{2v} \tan \delta$$



$$\alpha = \frac{\pi f}{v} \tan \delta$$

$$\alpha = \frac{\pi f}{Qv}$$

$$\alpha = \frac{\pi f}{f}$$

Velocità di propagazione e attenuazione



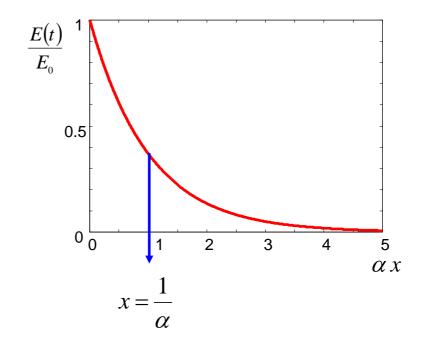
Velocità di propagazione

$$v = \frac{c}{\sqrt{\mathcal{E}_r}} \qquad \mathcal{E}_r = \frac{\mathcal{E}'}{\mathcal{E}_0} \qquad \mathcal{E}_r > 1 \qquad \qquad \qquad \text{Correnti di polarizzazione}$$

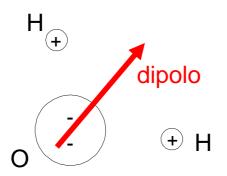
Attenuazione

$$\alpha = \frac{\mu v \sigma'}{2} \qquad \sigma' \quad \longrightarrow$$

Correnti di conduzione







Molecola di acqua

Elevata suscettività

$$\varepsilon_r \cong 80$$

Correnti di polarizzazione

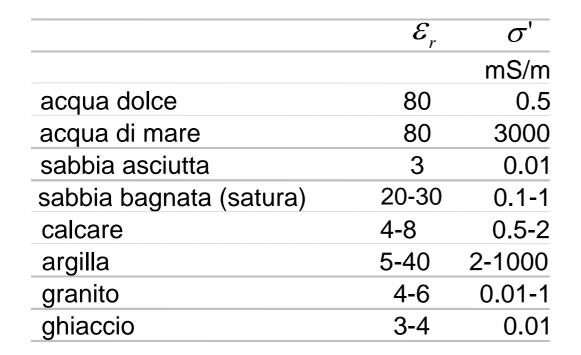
Molecola polare -

Dissolve ioni

 σ

Correnti di conduzione

Materiali



$$\varepsilon_r = \frac{\varepsilon'}{\varepsilon_0}$$

$$\varepsilon' = \varepsilon_0 + \frac{\sigma''}{\omega}$$

$$\varepsilon'' = \frac{\sigma'}{\omega}$$

Dispersione



$$v = \frac{c}{\sqrt{\varepsilon_r(\omega)}}$$

Impulso

 Δt = durata

 \mathcal{O}_0 = frequenza portante



Banda

$$B \approx \frac{1}{\Delta t}$$

 $B \approx \frac{1}{\Delta t}$ $\omega_0 = \text{frequenza portante}$

Componenti spettrali a frequenze diverse si propagano con velocità diversa



L'impulso si allarga

Dispersione



$$\delta t = \frac{L}{v_2} - \frac{L}{v_1} = \frac{L(v_1 - v_2)}{v_1 v_2} \approx -\frac{L\Delta v}{v_1^2}$$

$$v = \frac{c}{\sqrt{\mathcal{E}_r}}$$

$$\Delta v = \frac{\partial v}{\partial \varepsilon_r} \Delta \varepsilon_r = c \left(-\frac{1}{2} \right) \varepsilon_r^{-\frac{3}{2}} \Delta \varepsilon = -\frac{c}{2} \frac{1}{\varepsilon_r \sqrt{\varepsilon_r}} \Delta \varepsilon_r = -\frac{v}{2\varepsilon_r} \Delta \varepsilon_r$$

$$\delta t = \frac{L}{2v\varepsilon_r} \Delta \varepsilon_r$$

$$\delta t = \frac{L}{2v\varepsilon} \frac{\partial \varepsilon}{\partial f} B$$

Interazioni con discontinuità



1) Discontinuità localizzate

Teoria di Mie

$$P_{scatt} = I\sigma_{RCS}$$

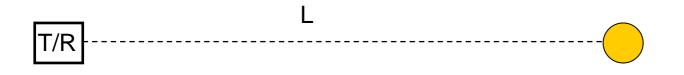
2) Superfici di discontinuità

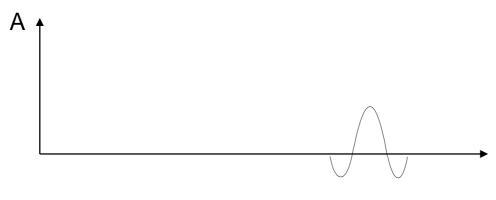
Incidenza perpendicolare

$$R = \frac{\sqrt{\left(\varepsilon_r\right)_2} - \sqrt{\left(\varepsilon_r\right)_1}}{\sqrt{\left(\varepsilon_r\right)_1} + \sqrt{\left(\varepsilon_r\right)_2}}$$

Risposta radar



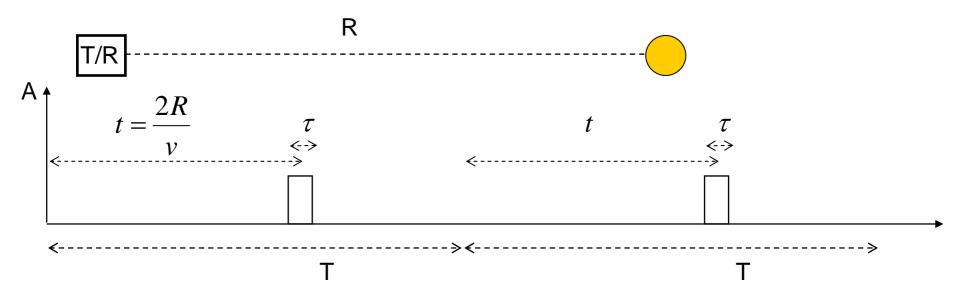




$$t = \frac{2L}{v}$$

Radar ad impulsi





Distanza target:
$$R = \frac{vt}{2}$$

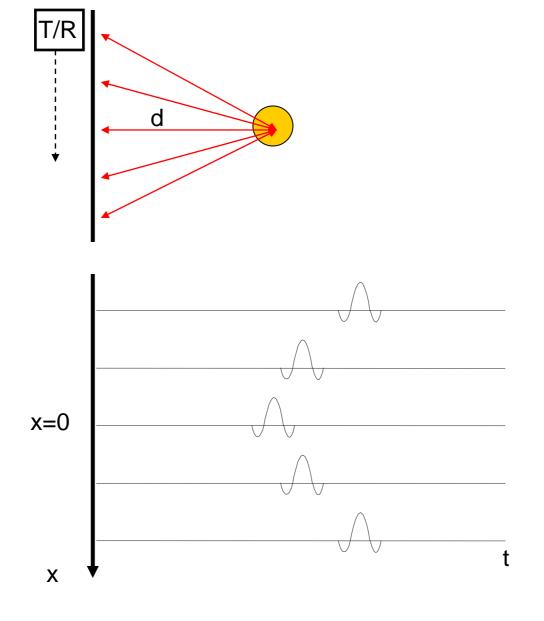
Risoluzione in range:
$$\Delta R = \frac{v\tau}{2}$$

Range non ambiguo:
$$R_U = \frac{vT}{2}$$

$$B = \frac{1}{\tau}$$

Risposta radar



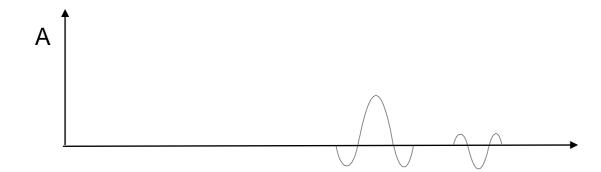


$$t = \frac{2\sqrt{x^2 + d^2}}{v}$$

Polarità





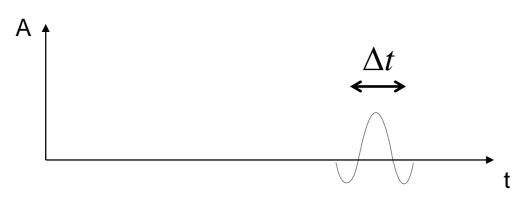


Incidenza perpendicolare

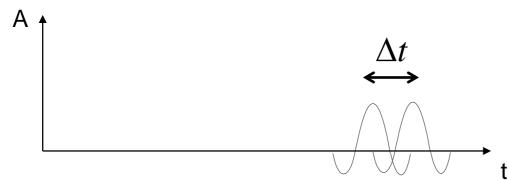
$$R = \frac{\sqrt{(\varepsilon_r)_2} - \sqrt{(\varepsilon_r)_1}}{\sqrt{(\varepsilon_r)_1} + \sqrt{(\varepsilon_r)_2}}$$

Risoluzione in range





$$\Delta t = \frac{1}{B}$$



$$\Delta R = \frac{v}{2} \Delta t$$

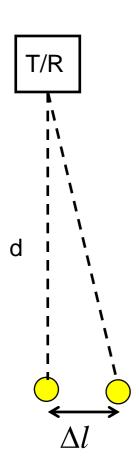
Bassa velocità di propagazione



$$\Delta R = \frac{v}{2B}$$

Risoluzione laterale (cross-range)





$$t_1 = \frac{2d}{v}$$

$$t_2 = \frac{2\sqrt{d^2 + (\Delta l)^2}}{v}$$

$$\Delta t = t_2 - t_1 = \frac{2}{v} \left(\sqrt{d^2 + (\Delta l)^2} - d \right)$$

$$\approx \frac{2}{v} \left[d \left(1 + \frac{1}{2} \left(\frac{\Delta l}{d} \right)^2 \right) - d \right] = \frac{(\Delta l)^2}{vd}$$

$$\Delta t = \frac{1}{B} = \frac{\left(\Delta l\right)^2}{vd}$$

$$\Delta l = \sqrt{\frac{vd}{B}}$$

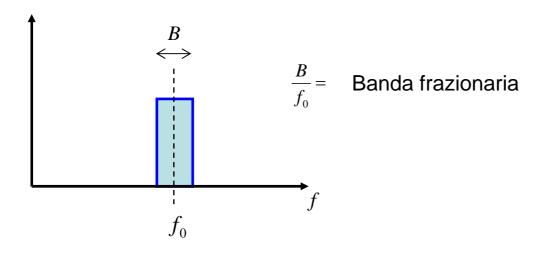
Risoluzione e penetrazione

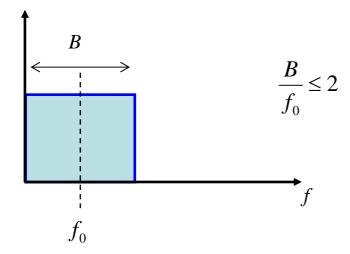


Risoluzione Larga banda Penetrazione **Bassa frequenza**

Banda frazionaria

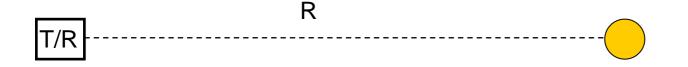


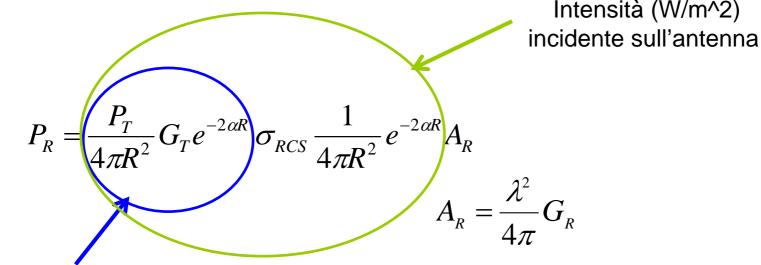




Equazione radar







Intensità (W/m^2) incidente sul bersaglio

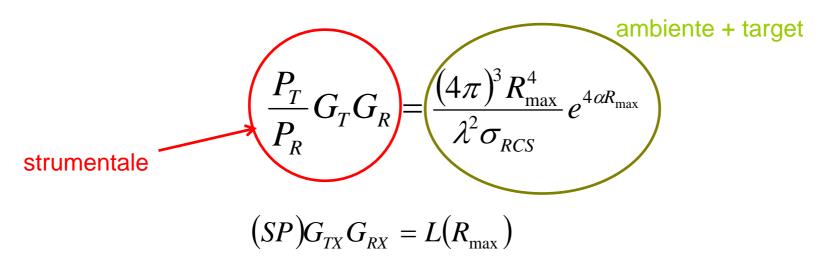
Equazione radar

$$P_R = P_T G_T G_R \frac{\lambda^2 \sigma_{RCS} e^{-4\alpha R}}{(4\pi)^3 R^4}$$

Profondità di penetrazione



$$P_{R \min} = P_{T \max} G_T G_R \frac{\lambda^2 \sigma_{RCS} e^{-4\alpha R_{\max}}}{(4\pi)^3 R_{\max}^4}$$



All'aumentare della frequenza, aumenta linearmente α quindi deve diminuire R_{max} mantenere L costante

Profondità di penetrazione



$$(SP)G_{TX}G_{RX} = L(R_{\max})$$

$$\alpha = \frac{\pi f}{Ov}$$

$$L(R_{\text{max}}) = \frac{(4\pi)^3 R_{\text{max}}^4}{\lambda^2 \sigma_{RCS}} e^{4\alpha R} = \frac{(4\pi)^3 \left(\frac{R_{\text{max}}}{\lambda}\right)^4}{\left(\frac{\sigma_{RCS}}{\lambda^2}\right)} e^{4\frac{\pi}{Q}\left(\frac{R_{\text{max}}}{\lambda}\right)}$$

$$SPG_{TX}G_{RX} = L \text{ [dB]} \begin{array}{c} 180 \\ 160 \\ 140 \\ 100 \\ 80 \\ 60 \\ 40 \\ 20 \\ 0 \\ -20 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \quad Q = 1 \\ 0 \\ 100 \\ 100 \\ -20 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \quad Q = 1 \\ 0 \\ 100 \\ 100 \\ -20 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ 100 \\ -20 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ 100 \\ -20 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ 100 \\ -20 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ 100 \\ -20 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ 100 \\ -20 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array} \begin{array}{c} \frac{\sigma_{RCS}}{\lambda^2} = 1 \\ 0 \\ -3 \end{array}$$

Risoluzione in range



$$\Delta R = \frac{v}{2B}$$

$$B = \frac{\int_{0}^{\infty} A(f) df}{A_{\text{max}}}$$

$$A(f) = e^{-\frac{\pi f}{Qv}R} = e^{-\eta f}$$

$$\int_{c}^{f_{c} + \frac{B_{0}}{2}} \int_{e^{-\eta f} df}^{f_{c} + \frac{B_{0}}{2}} = e^{\eta f_{\min}} \frac{-1}{\eta} \left[e^{-\eta \left(f_{c} + \frac{B_{0}}{2} \right)} - e^{-\eta \left(f_{c} - \frac{B_{0}}{2} \right)} \right] = e^{\eta f_{\min}} \frac{1}{\delta} e^{-\eta f_{c}} \left[e^{\eta \frac{B_{0}}{2}} - e^{-\eta \frac{B_{0}}{2}} \right]$$

$$B = e^{\eta f_{\min}} \frac{1}{\eta} e^{-\eta f_c} 2 \sinh\left(\frac{\eta B_0}{2}\right) = \frac{2}{\eta} e^{-\eta (f_c - f_{\min})} \sinh\left(\frac{\eta B_0}{2}\right) = \frac{2}{\eta} e^{-\frac{\eta B_0}{2}} \sinh\left(\frac{\eta B_0}{2}\right)$$

$$B = \frac{2}{\eta} e^{-\frac{\eta B_0}{2}} \sinh\left(\frac{\eta B_0}{2}\right) \approx \frac{2}{\eta} \left(1 - \frac{\eta B_0}{2}\right) \left(\frac{\eta B_0}{2}\right) = B_0 - \frac{\eta^2 B_0^2}{2}$$

Risoluzione in range



$$\Delta R = \frac{v}{2B}$$

$$B = B_0 - \frac{\eta^2 B_0^2}{2} \qquad \eta = \frac{\pi R}{Qv}$$

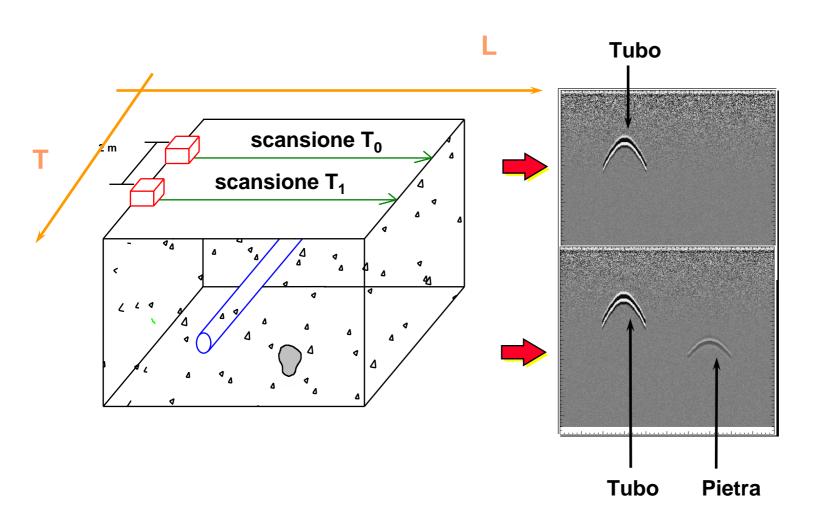
$$\Delta R = \frac{v}{2B_0} = \frac{v}{2\left(B_0 - \frac{\eta B_0^2}{2}\right)} = \frac{v}{2B_0} \frac{1}{1 - \frac{\eta B}{2}} = \frac{v}{2B_0} \left(1 + \frac{\eta B_0}{2}\right)$$

$$\Delta R = \Delta R_0 \left(1 + \left(\frac{\pi R}{Q v} \right) \frac{B}{2} \right) = \Delta R_0 + \frac{v}{2B_0} \frac{\pi R}{Q v} \frac{B_0}{2}$$

$$\Delta R = \Delta R_0 + \frac{\pi}{4} \frac{1}{Q} R$$

Modalità operativa







Elettronica

Elaborazione

- on line
- post-processing

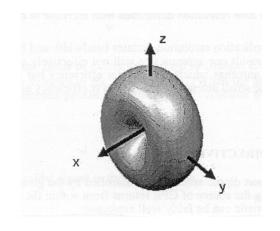




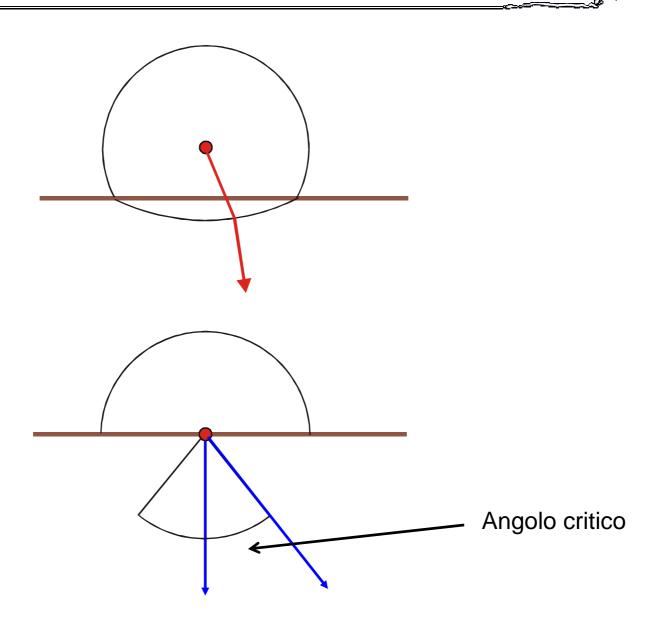
Larga Banda Frequenza bassa

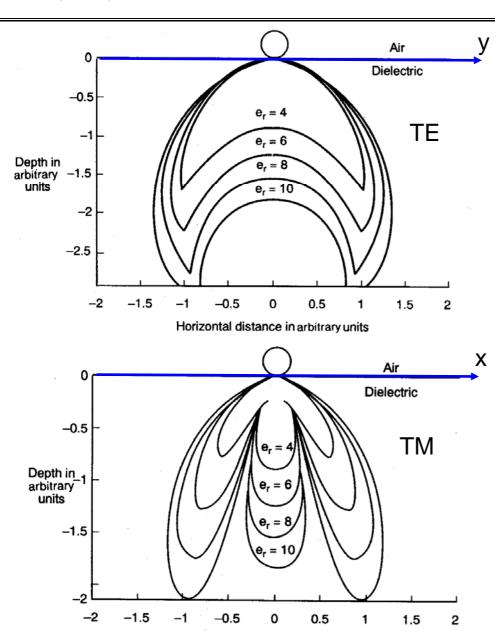
Bow-tie





Schermatura



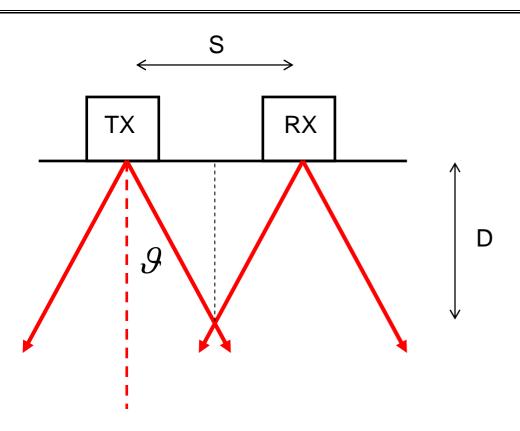


Horizontal distance in arbitrary units

(x asse del dipolo)

Antenne bow-tie

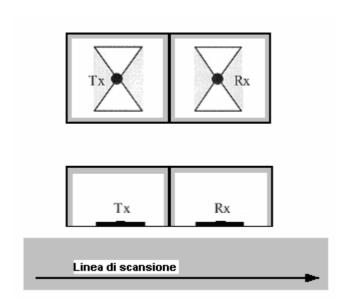


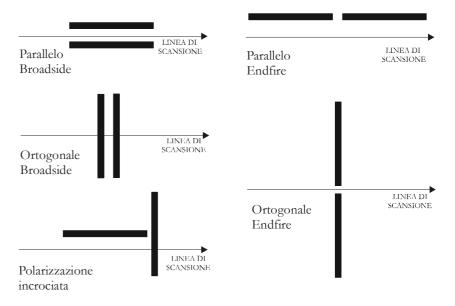


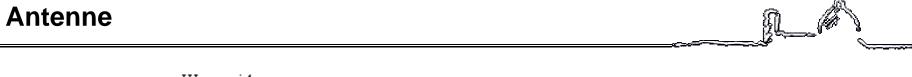
$$\sin(\mathcal{G}) = \frac{1}{\sqrt{\mathcal{E}_r}}$$

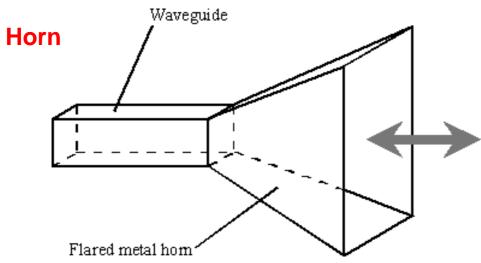
$$D = \frac{S}{2\sqrt{\varepsilon_r - 1}}$$

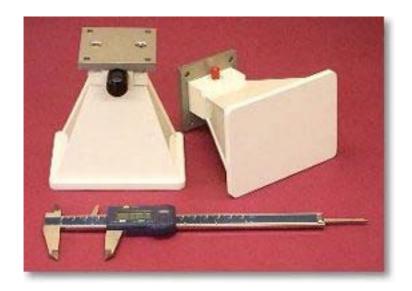














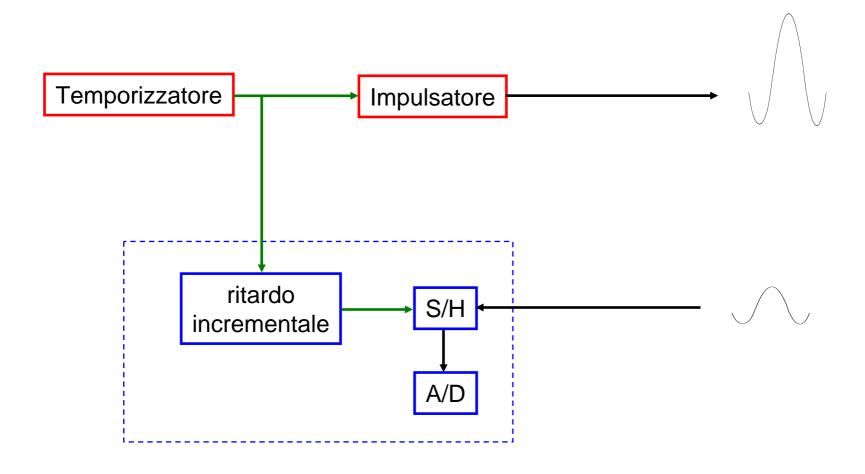
Radar penetranti



Elettronica

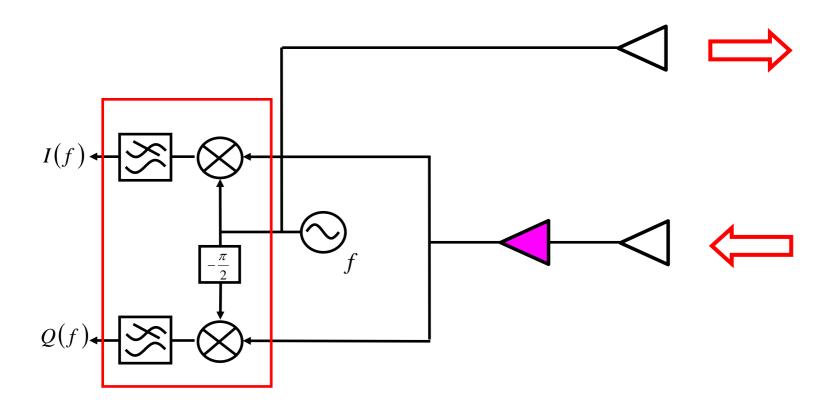
Radar impulsati





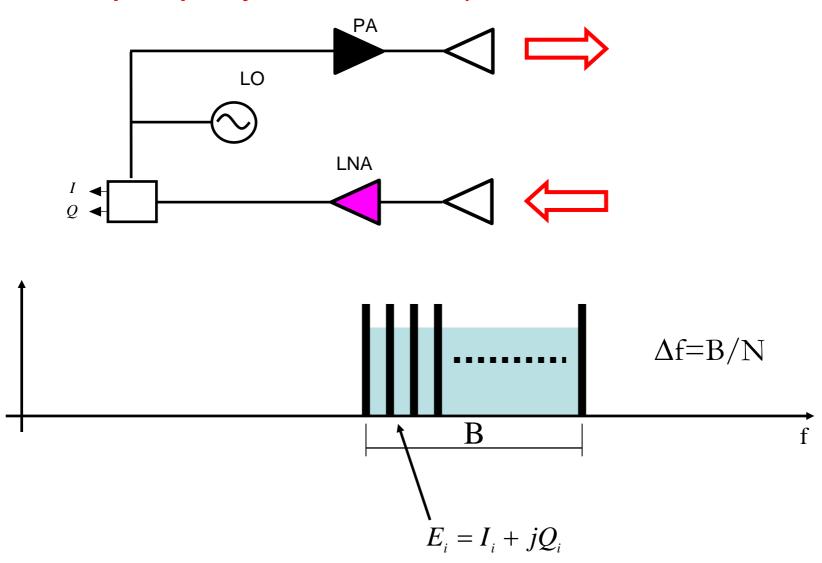
Radar ad onda continua





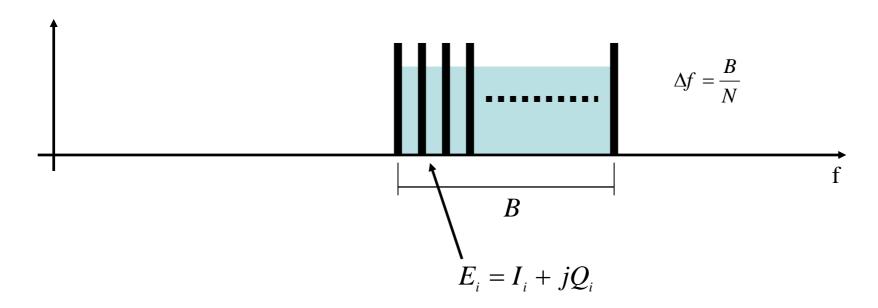


(SF-CW: Step Frequency Continuous Wave)





(SF-CW: Step Frequency Continuous Wave)



$$f(t) = \sum_{i} E_{i} e^{j\omega_{i}t}$$

Formula per la sintesi dell'impulso



$$R_0$$

$$E(f) = e^{-j\frac{2\pi}{\lambda}(2R_0)}$$

$$\widetilde{E}^{-1}(t) = \frac{1}{2\pi(f_2 - f)_1} \int_{f_1}^{f_2} E(f) e^{j2\pi f t} df = \frac{1}{2\pi(f_2 - f_1)} \int_{f_1}^{f_2} e^{-j\frac{2\pi}{c}f(2R_0)} e^{j2\pi f t} df =$$

$$=\frac{1}{2\pi(f_{2}-f_{1})}\int_{f_{1}}^{f_{2}}e^{-j2\pi f\left(\frac{2R_{0}}{c}\right)}e^{j2\pi ft}df=...=e^{j2\pi\frac{f_{1}+f_{2}}{2}\left(\frac{2R_{0}}{c}\right)}\frac{\sin\left[2\pi\left(\frac{f_{2}-f_{1}}{2}\right)\left(t-\frac{2R_{0}}{c}\right)\right]}{2\pi\left(\frac{f_{2}-f_{1}}{2}\right)\left(t-\frac{2R_{0}}{c}\right)}$$

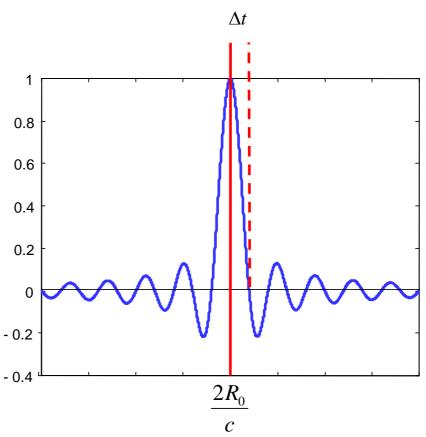


$$\widetilde{E}^{-1}(t) = e^{j2\pi f_c\left(\frac{2R_0}{c}\right)} \frac{\sin\left[2\pi\left(\frac{B}{2}\right)\left(t - \frac{2R_0}{c}\right)\right]}{2\pi\left(\frac{B}{2}\right)\left(t - \frac{2R_0}{c}\right)}$$

$$2\pi \left(\frac{B}{2}\right) \Delta t = \pi$$

$$\Delta t = \frac{1}{B}$$

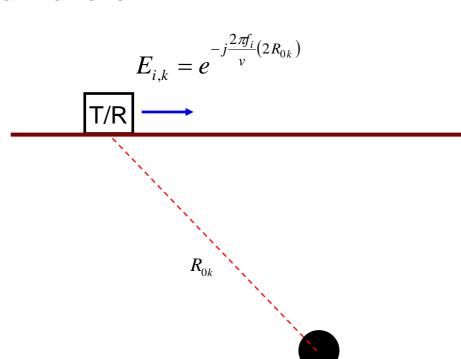
$$\Delta R = \frac{c}{2B}$$



Impulso sintetico



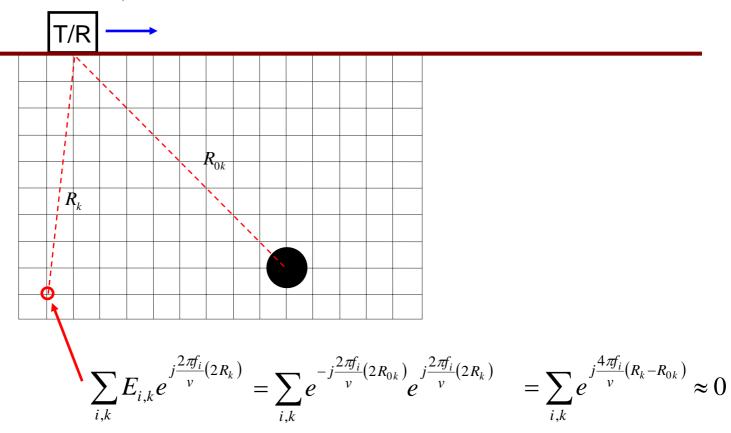
Focalizzazione





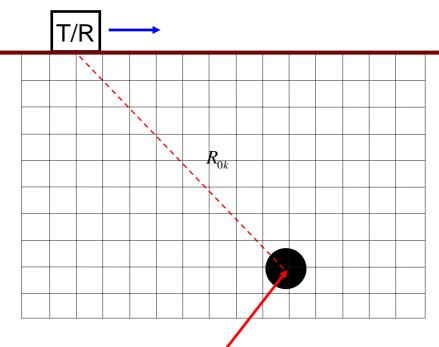
Focalizzazione

$$E_{i,k} = e^{-j\frac{2\pi f_i}{v}(2R_{0k})}$$





$$E_{i,k} = e^{-j\frac{2\pi f_i}{v}(2R_{0k})}$$

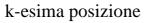


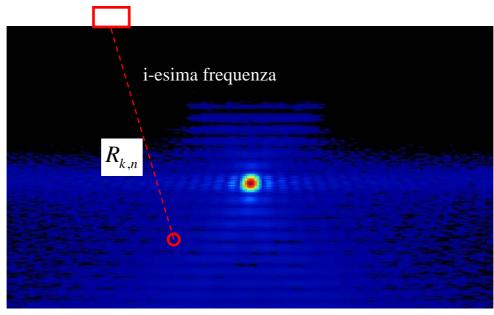
$$R_k = R_{0k}$$

$$\sum_{i,k} E_{i,k} e^{j\frac{2\pi f_i}{v}(2R_k)} = \sum_{i,k} e^{-j\frac{2\pi f_i}{v}(2R_{0k})} e^{j\frac{2\pi f_i}{v}(2R_k)} = \sum_{i,k} 1 = N_f N_p$$

Impulso sintetico e apertura sintetica







Impulso sintetico e apertura sintetica



Diversità in frequenza

Risoluzione in range

Diversità spaziale — Risoluzione angolare (cross-range)