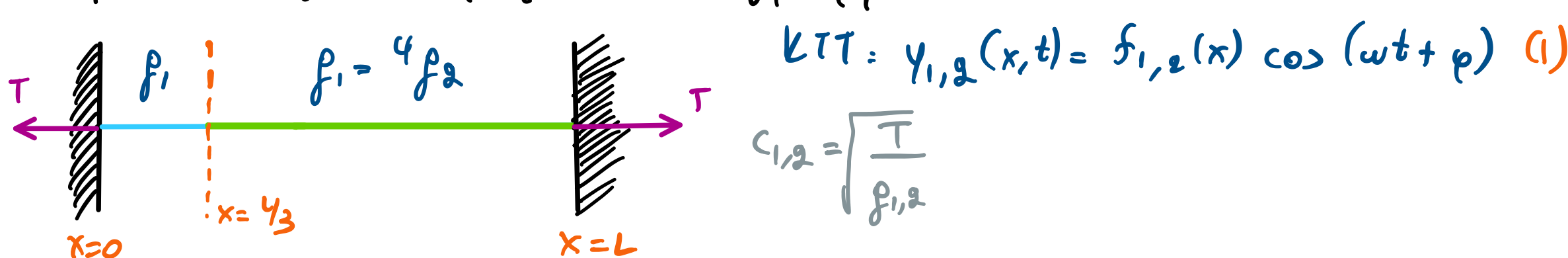


Παράδειγμα

Υπολογισμός των ΚΤΤ (σεάσιμα κύματα) σε συνδυασμό

2 χορδών με διαφορετικές γραμμικές πυκνότητες.



$$\frac{1}{c_{1,2}} \frac{\partial^2 y_{1,2}}{\partial t^2} = \frac{\partial^2 y_{1,2}}{\partial x^2} \quad (2)$$

$$f_1'' + \left(\frac{\omega}{c_1}\right)^2 f_1 = 0 \rightarrow f_1(x) = A \sin(k_1 x + \theta_1)$$

$$f_2'' + \left(\frac{\omega}{c_2}\right)^2 f_2 = 0 \rightarrow f_2(x) = B \sin(k_2 x + \theta_2)$$

Συνοριακές συνθήκες

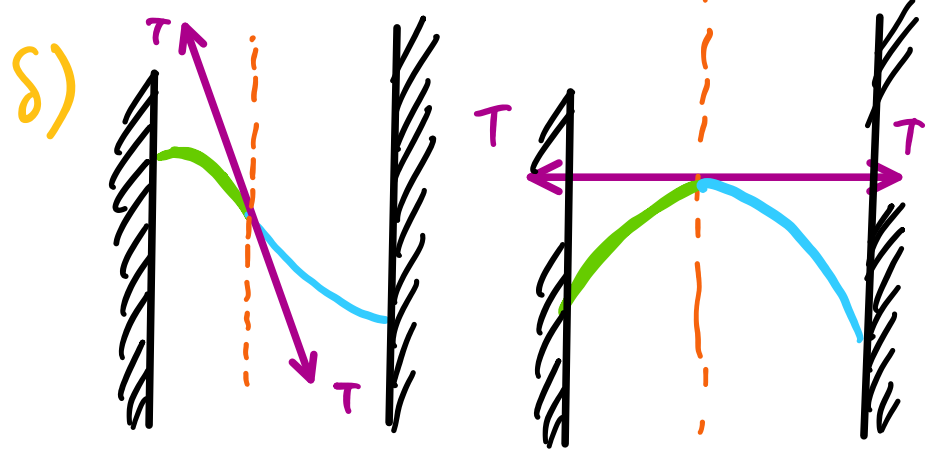
α)  $f_1(x=0) = 0$  (ακλόνητο άκρο)  $\Rightarrow \sin \theta_1 = 0 \Rightarrow \theta_1 = 0$

β)  $f_2(x=L) = 0$  (ακλόνητο άκρο)  $\Rightarrow \sin(k_2 L + \theta_2) = 0 \Rightarrow \theta_2 = -k_2 L$

γ)  $f_1(x=L/3) = f_2(x=L/3) \Rightarrow A \sin(k_1 L/3) = B \sin(-k_2 L/3)$  (3)

$$\frac{k_2}{k_1} = \frac{c_1}{c_2} \Rightarrow k_2 = \frac{k_1}{2} \quad (4)$$

(3)  $\xrightarrow{(4)} A \sin(k_1 L/3) = -B \sin(k_1 L/3) \Rightarrow \begin{cases} A = -B & (1a) \\ \sin(k_1 L/3) = 0 \Rightarrow k_{1,n} = n \frac{3\pi}{L} & (1\beta) \end{cases}$



$m=0 \Rightarrow f_{01}, y=0 \Rightarrow$

$$\left(\frac{\partial f_1}{\partial x}\right)_{x=L/3} = \left(\frac{\partial f_2}{\partial x}\right)_{x=L/3} \Rightarrow \begin{cases} k_1 A \cos(k_1 L/3) = k_2 B \cos(-k_2 L/3) \\ k_1 A \cos(k_1 L/3) = \frac{k_1}{2} B \cos(-k_1 L/3) \end{cases}$$

(1a), (2)  $\Rightarrow k_{1,n} = (2n-1) \frac{3\pi}{2L}$

(1\beta), (2)  $\Rightarrow A = \pm B/2$

Κύματα σε γραμμές μεταφοράς

$$\begin{cases} n \rightarrow x \\ n+L \rightarrow x+dx \end{cases} \Rightarrow V(x) - (L_0 dx) \frac{\partial I}{\partial t} = V(x) + \frac{\partial V}{\partial x} dx \quad (1)$$

$$I(x-dx) = I(x) + \frac{d}{dt} V \cdot C_0 \cdot dx \xrightarrow{\text{Taylor}}$$

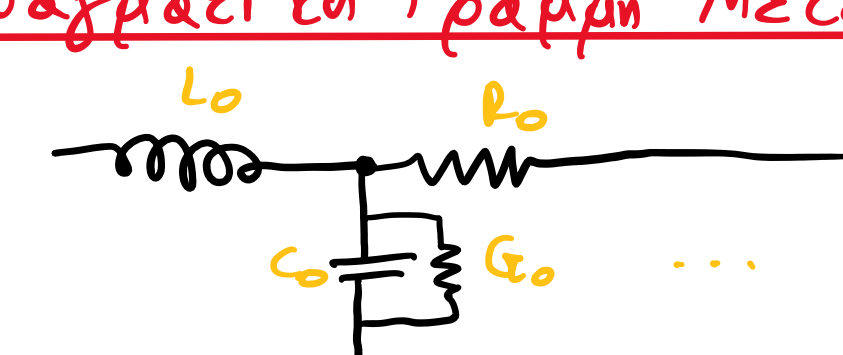
$$I(x) - \frac{\partial I}{\partial x} dx = I(x) + C_0 \frac{\partial V}{\partial t} dx \Rightarrow$$

$$-\frac{\partial I}{\partial x} dx = C_0 \frac{\partial V}{\partial t} dx \quad (2)$$

$$\frac{\partial}{\partial t} (1) \Rightarrow -L_0 \frac{\partial^2 I}{\partial t^2} = \frac{\partial}{\partial t} \frac{\partial V}{\partial x} = \frac{\partial}{\partial x} \frac{\partial V}{\partial t} \Rightarrow$$

$$-L_0 \frac{\partial^2 I}{\partial t^2} = \frac{\partial}{\partial x} \left( -\frac{1}{C_0} \frac{\partial I}{\partial x} \right) \Rightarrow L_0 C_0 \frac{\partial^2 I}{\partial t^2} = \frac{\partial^2 I}{\partial x^2}, \quad c = \frac{1}{\sqrt{L_0 C_0}}$$

$$\frac{\partial}{\partial t} (2) \Rightarrow L_0 C_0 \frac{\partial^2 V}{\partial t^2} = \frac{\partial^2 V}{\partial x^2}, \quad c = \frac{1}{\sqrt{L_0 C_0}}$$

Πραγματική Γραμμή Μεταφοράς

$$\begin{aligned} V(x) - (L_0 dx) \frac{\partial I}{\partial t} - (R_0 dx) I &= V(x+dx) \\ I(x) &= I(x+dx) + \frac{\partial}{\partial t} (C_0 dx V) + (G_0 dx V) \end{aligned} \Rightarrow$$

$$\Rightarrow \frac{\partial^2 V}{\partial x^2} = L_0 C_0 \frac{\partial^2 V}{\partial t^2} + (R_0 C_0 + L_0 G_0) \frac{\partial V}{\partial t} + R_0 G_0 V$$

Φαινόμενα διασποράς

Διασπορά: ταχύτητα φάσης ενός κύματος είναι διαφορετική για κάθε συχνότητα.

Ταχύτητα φάσης:  $v_m = \frac{\omega}{k} = \frac{\lambda}{T_{\text{φ}}} = \sqrt{\frac{T_{\text{scw}}}{\rho}}$

Για ιδανική χορδή:  $v_m = \frac{\omega}{k} = c = \sqrt{\frac{T}{\rho}} \neq f(\omega)$

π.χ.  $\frac{1}{c^2} \frac{\partial^2 y}{\partial t^2} = \frac{\partial^2 y}{\partial x^2} + a \frac{\partial^4 y}{\partial x^4}$  (1)

Διάδοση αρμονικού κύματος  $y = A e^{i(kx - \omega t)}$

(2)  $\Rightarrow (1) \Rightarrow -\frac{\omega^2}{c^2} e^{i(kx - \omega t)} = -k^2 A e^{i(kx - \omega t)} + a k^4 A e^{i(kx - \omega t)} \Rightarrow$

$$\Rightarrow \omega^2 = c^2 k^2 - c^2 a k^4 \Rightarrow \omega = c k \sqrt{1 - a k^2} = \omega(k) \quad \text{σχέση διασποράς}$$

$k = k(\omega)$ : σχέση διασποράς

$$v_{ph} = \frac{\omega}{k} = \frac{\omega}{k(\omega)} = v_{ph}(\omega)$$

$$k = \frac{2\pi}{\lambda} = k(\omega) \Rightarrow \lambda = \frac{2\pi}{k(\omega)}$$

Παράδειγμα διασποράς:

Ιδανική χορδή ( $\rho, T$ ) μέσα σε ένα "ελαστικό περιβάλλον" με σταθερά σκληρότητας ανά μονάδα μήκους  $\frac{ds}{dx} = \sigma$

$$dm = \rho dx \Rightarrow$$

$$dm \frac{\partial^2 y}{\partial t^2} = T \frac{\partial^2 y}{\partial x^2} dx - (\sigma dx) y \Rightarrow$$

$$\frac{1}{c^2} \frac{\partial^2 y}{\partial t^2} = \frac{\partial^2 y}{\partial x^2} - \frac{\sigma}{T} y \Rightarrow \begin{cases} -\frac{\omega^2}{c^2} \cancel{y} = -k^2 \cancel{y} - \frac{\sigma}{T} \cancel{y} \Rightarrow \\ \omega^2 = c^2 k^2 + c^2 \frac{\sigma}{T} \end{cases}$$

$$k^2 = \frac{\omega^2}{c^2} - \frac{\sigma}{T} \Rightarrow \frac{2\pi}{\lambda} = \sqrt{\frac{\omega^2}{c^2} - \frac{\sigma}{T}}$$

$$\lambda = \frac{2\pi}{\sqrt{\frac{\omega^2}{c^2} - \frac{\sigma}{T}}} \quad v_{ph} = \frac{\omega}{k} = \frac{\omega}{\sqrt{\frac{\omega^2}{c^2} - \frac{\sigma}{T}}} = v_{ph}(\omega)$$

$$v_{ph}: \text{πραγματική} \Rightarrow \frac{\omega^2}{(c)^2} > \frac{\sigma}{T} > \omega > \sigma \sqrt{\frac{\sigma}{T}} = \omega_c$$

$\omega > \omega_c$ : διάδοση κύματος

$$k = \sqrt{\frac{\omega^2}{c^2} - \frac{\sigma}{T}}: \text{πραγματικό } \omega > \omega_c$$

$$\omega < \omega_c \Rightarrow k = ia \Rightarrow y = A e^{i(ia - \omega t)} = (A e^{-ax}) e^{-i\omega t}$$