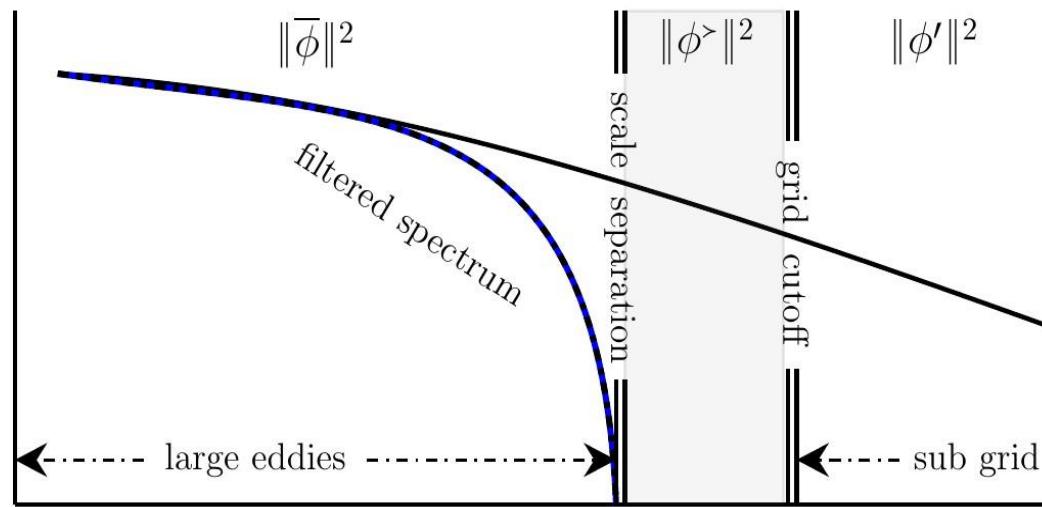


Large Eddy Simulation of Burgulence: a Synthesis of Filtering, Modelling and Discretization

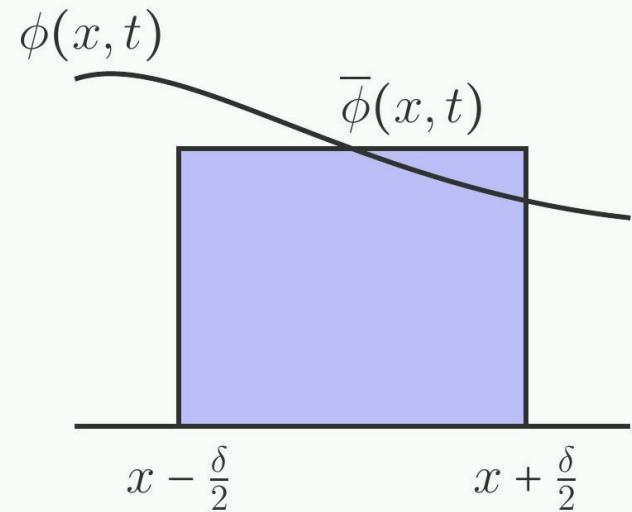


Roel Verstappen

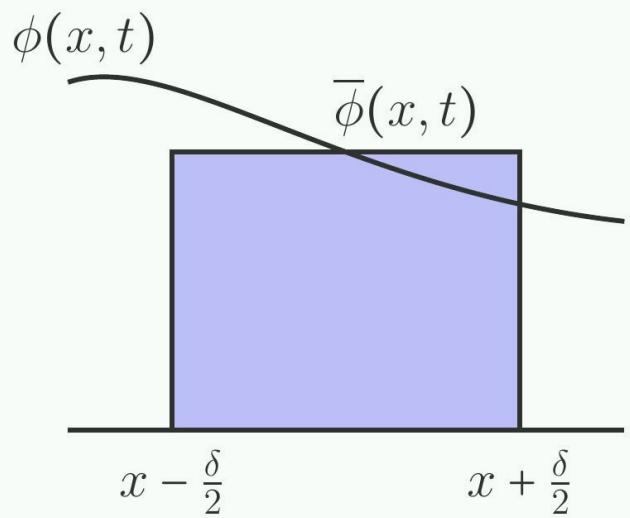


university of
groningen

Filter



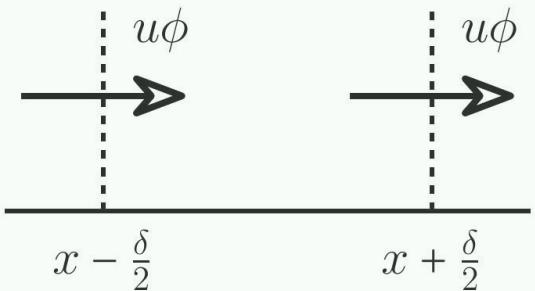
$$\bar{\phi}(x, t) = \frac{1}{\delta} \int_{x-\frac{\delta}{2}}^{x+\frac{\delta}{2}} \phi(x, t) dx$$

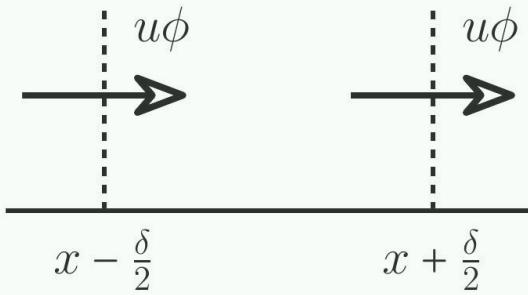


$$\bar{\phi}(x, t) = \frac{1}{\delta} \int_{x-\frac{\delta}{2}}^{x+\frac{\delta}{2}} \phi(x, t) dx$$

Conservation

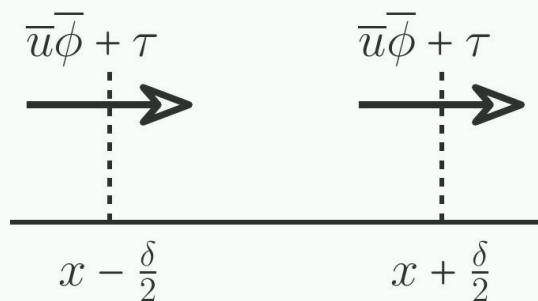
$$\delta \partial_t \bar{\phi}(x, t) = (u\phi)(x - \frac{\delta}{2}, t) - (u\phi)(x + \frac{\delta}{2}, t)$$





$$\delta \partial_t \bar{\phi}(x, t) = (u\phi)(x - \frac{\delta}{2}, t) - (u\phi)(x + \frac{\delta}{2}, t)$$

Closure

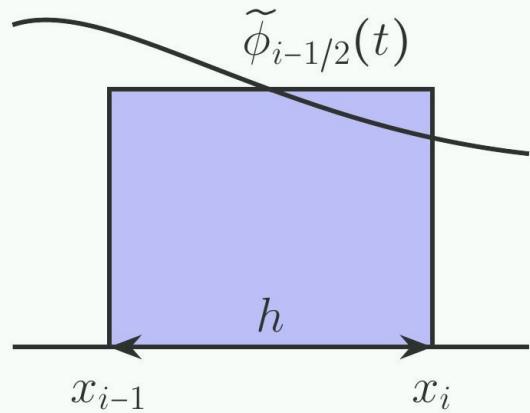


$$\delta \partial_t \bar{\phi}(x, t) = (\bar{u}\phi + \tau)(x - \frac{\delta}{2}, t) - (\bar{u}\phi + \tau)(x + \frac{\delta}{2}, t)$$

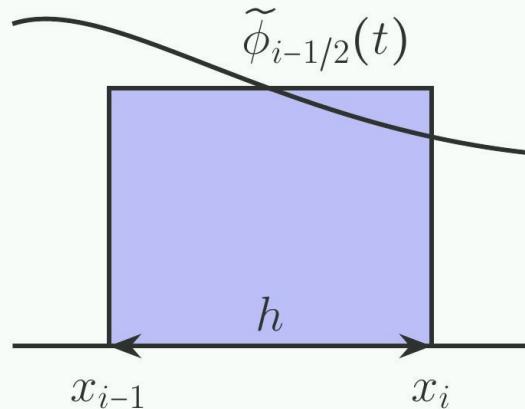
$$\tau(\bar{u}, \bar{\phi}) \simeq u\phi - \bar{u}\phi$$

U. Schumann JCP 1975

Grid filter



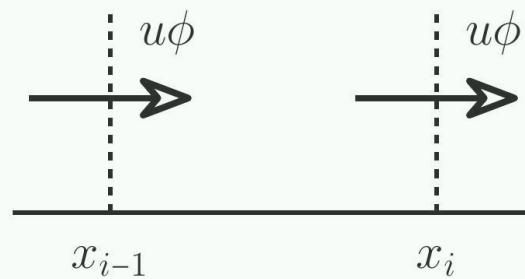
$$\tilde{\phi}_{i-1/2}(t) = \frac{1}{h} \int_{x_{i-1}}^{x_i} \phi(x, t) dx$$

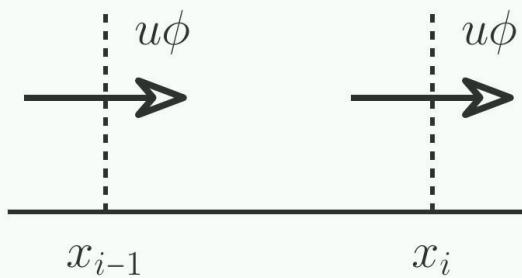


$$\tilde{\phi}_{i-1/2}(t) = \frac{1}{h} \int_{x_{i-1}}^{x_i} \phi(x, t) dx$$

Conservation

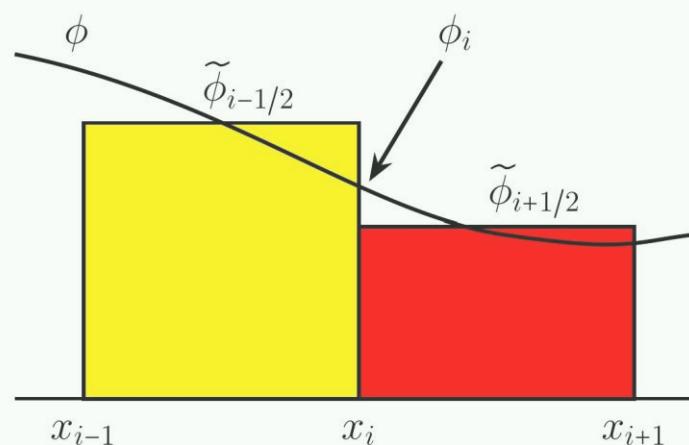
$$h \partial_t \tilde{\phi}_{i-1/2} = (u\phi)_{i-1} - (u\phi)_i$$





$$h \partial_t \tilde{\phi}_{i-1/2} = (u\phi)_{i-1} - (u\phi)_i$$

Interpolation



$$\phi_i \simeq \frac{1}{2}(\tilde{\phi}_{i-1/2} + \tilde{\phi}_{i+1/2})$$

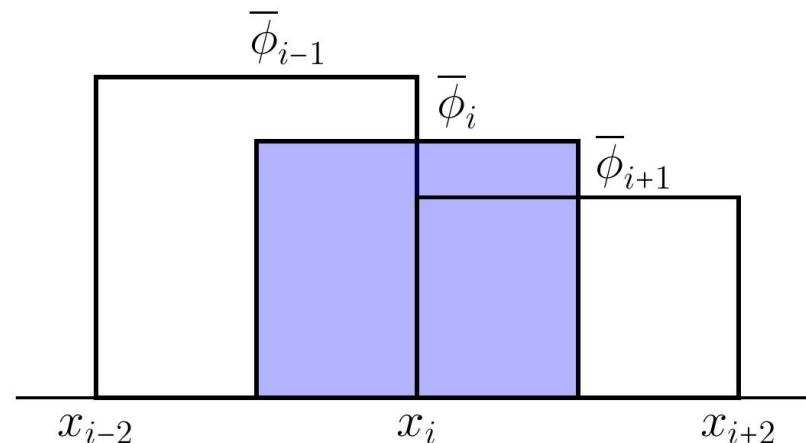
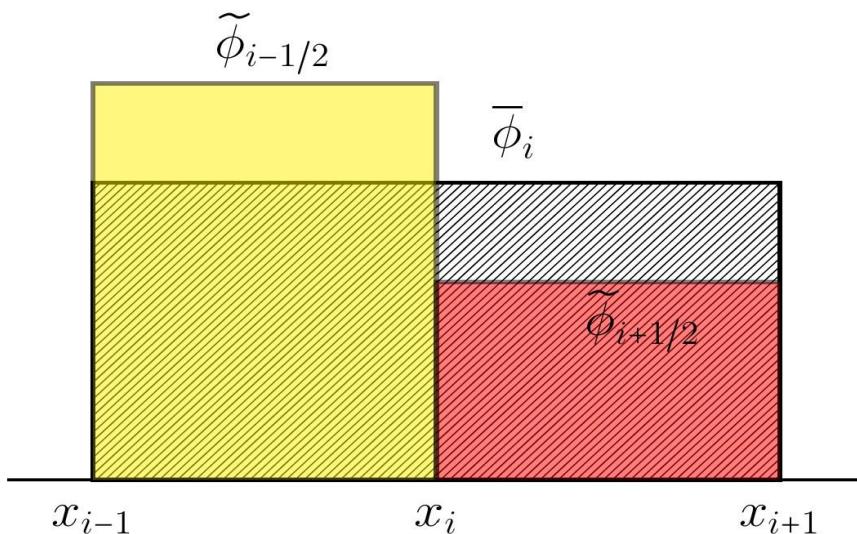
$$= \frac{1}{2h} \int_{x_{i-1}}^{x_{i+1}} \phi(x, t) dx = \bar{\phi}_i$$

Two filters, lengths h and $\delta = 2h$

Two filters - large eddies are defined by interpolation filter

$$\tilde{\phi}_{i-1/2}(t) = \frac{1}{h} \int_{x_{i-1}}^{x_i} \phi(x, t) dx$$

$$\bar{\phi}_i(t) = \frac{1}{\delta} \int_{x_{i-1}}^{x_{i+1}} \phi(x, t) dx$$



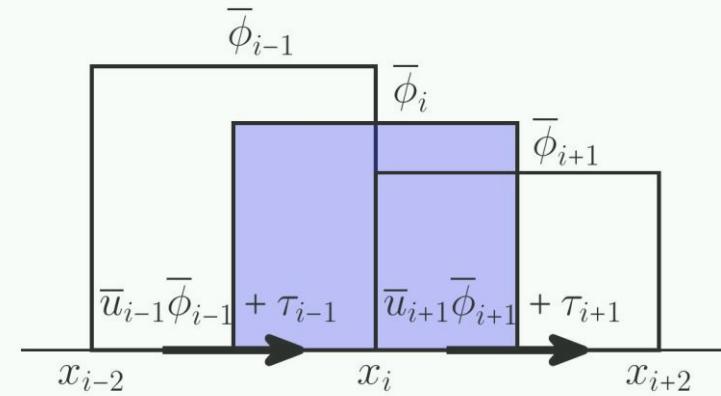
$$\bar{\phi}_i = \frac{1}{2}(\tilde{\phi}_{i-1/2} + \tilde{\phi}_{i+1/2})$$

Note: overlapping boxes

FVM

Adding neighbours

Truncation error

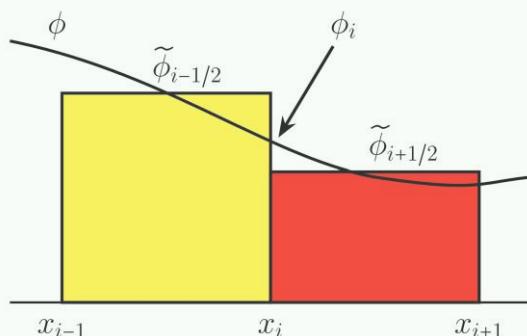


$$\begin{aligned}
 h \partial_t \tilde{\phi}_{i-1/2} + \bar{u}_i \bar{\phi}_i - \bar{u}_{i-1} \bar{\phi}_{i-1} &= -\tau_i + \tau_{i-1} \\
 h \partial_t \tilde{\phi}_{i+1/2} + \bar{u}_{i+1} \bar{\phi}_{i+1} - \bar{u}_i \bar{\phi}_i &= -\tau_{i+1} + \tau_i \\
 \hline
 \delta \partial_t \bar{\phi}_i + \bar{u}_{i+1} \bar{\phi}_{i+1} - \bar{u}_{i-1} \bar{\phi}_{i-1} &= -\tau_{i+1} + \tau_{i-1}
 \end{aligned}$$

The first two equations are crossed out with orange lines, while the third equation is shown in blue.

Recall: interpolation

$$\bar{\phi}_i = \frac{1}{2}(\tilde{\phi}_{i-1/2} + \tilde{\phi}_{i+1/2})$$



FVM

$$\delta \partial_t \bar{\phi}_i + \bar{u}_{i+1} \bar{\phi}_{i+1} - \bar{u}_{i-1} \bar{\phi}_{i-1} = -\tau_{i+1} + \tau_{i-1}$$

Truncation

$$\delta \partial_t \bar{\phi}(x) + (\bar{u}\bar{\phi})(x+\frac{\delta}{2}) - (\bar{u}\bar{\phi})(x-\frac{\delta}{2}) = -\tau(x+\frac{\delta}{2}) + \tau(x-\frac{\delta}{2})$$

Closure

$$\delta \partial_t \bar{\phi}_i + \bar{u}_{i+1} \bar{\phi}_{i+1} - \bar{u}_{i-1} \bar{\phi}_{i-1} = -\tau_{i+1} + \tau_{i-1}$$

Truncation

$$\delta \partial_t \bar{\phi}(x) + (\bar{u}\bar{\phi})(x+\frac{\delta}{2}) - (\bar{u}\bar{\phi})(x-\frac{\delta}{2}) = -\tau(x+\frac{\delta}{2}) + \tau(x-\frac{\delta}{2})$$

Closure

No difference between truncation error and closure model if

$$x = x_i \quad \delta = 2h$$

Box filter

$$\tilde{\phi}(x, t) = \frac{1}{h} \int_{x-\frac{h}{2}}^{x+\frac{h}{2}} \phi(\xi, t) d\xi$$

commutes with differentiation

$$\partial_x \tilde{\phi} = \frac{\phi(x + \frac{h}{2}, t) - \phi(x - \frac{h}{2}, t)}{h} = \widetilde{\partial_x \phi}$$

$$\partial_x \tilde{\phi} = \frac{\phi(x + \frac{h}{2}, t) - \phi(x - \frac{h}{2}, t)}{h} = \widetilde{\partial_x \phi}$$

$$\partial_x \tilde{\phi} \approx \frac{\tilde{\phi}(x + \frac{h}{2}, t) - \tilde{\phi}(x - \frac{h}{2}, t)}{h}$$

Viscous flux

$$\nu \partial_x \phi = \nu \partial_x \tilde{\phi} + \nu \partial_x (\phi - \tilde{\phi})$$

$$= \nu \frac{\tilde{\phi}(x + \frac{h}{2}, t) - \tilde{\phi}(x - \frac{h}{2}, t)}{h} + \nu \partial_x (\phi - \tilde{\phi})$$

unclosed term $\rightarrow \tau$

subgrid modeling and discretization are intrinsically linked 

Filters

$$\tilde{\phi}_{i-1/2}(t) = \frac{1}{h} \int_{x_{i-1}}^{x_i} \phi(x, t) dx \quad \bar{\phi}_i(t) = \frac{1}{\delta} \int_{x_{i-1}}^{x_{i+1}} \phi(x, t) dx$$

Flux

$$\bar{u}_i \bar{\phi}_i = \nu \frac{\tilde{\phi}_{i+1/2} - \tilde{\phi}_{i-1/2}}{h} + \tau_i$$

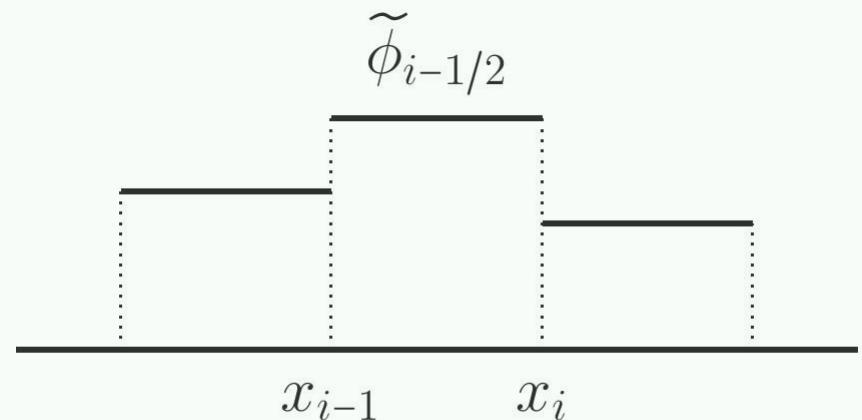
The diagram illustrates the decomposition of the flux term. Two blue arrows point from the terms $\tilde{\phi}_{i+1/2}$ and $\tilde{\phi}_{i-1/2}$ towards the center of the equation, indicating they are inputs to the LES model. Another blue arrow points from the term τ_i towards the center, indicating it is the error from the FVM.

Grid filter

$$\tilde{\phi}(x, t) = \tilde{\phi}_{i-1/2}(t) \quad \text{for} \quad x_{i-1} \leq x < x_i$$

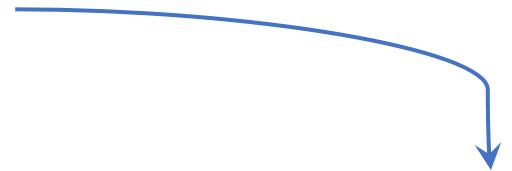
Idempotent

$$\tilde{\tilde{\phi}} = \tilde{\phi} \implies \tilde{\phi}' = \overline{\phi - \tilde{\phi}} = 0$$



Grid filter

$$\implies \tilde{\phi}' = 0$$

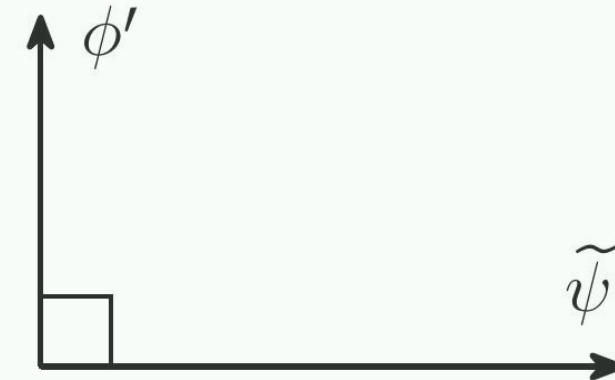


Symmetric

$$(\tilde{\phi}, \phi') = (\phi, \tilde{\phi}') = 0$$

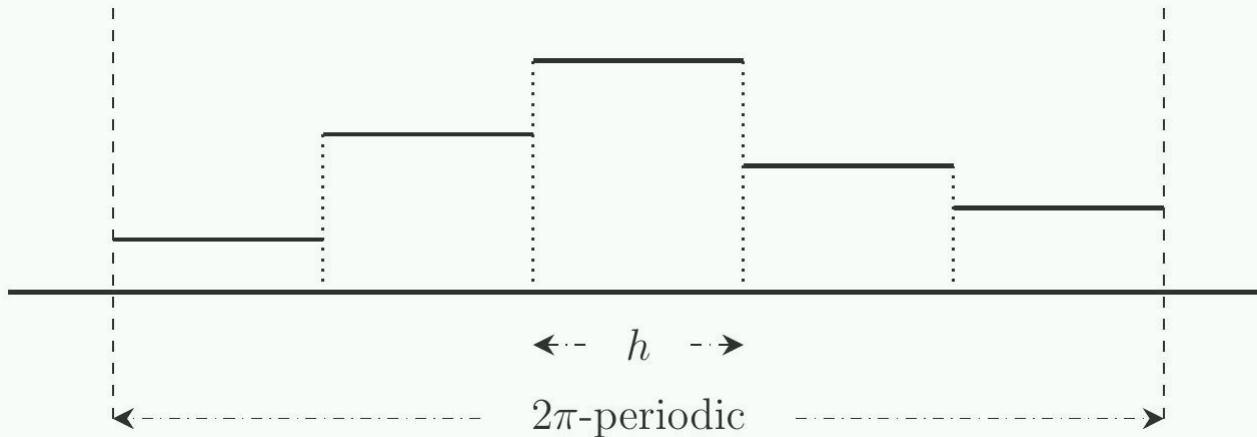
$$(\tilde{\psi}, \phi') = (\psi, \tilde{\phi}') = 0$$

Orthogonal decomposition of function space



Finite volume space

$$\tilde{\mathcal{F}}_h(0, 2\pi)$$



$$\tilde{\phi}(x, t) \in \tilde{\mathcal{F}}_h(0, 2\pi) \implies \tilde{\phi}(x \pm h, t) \in \tilde{\mathcal{F}}_h(0, 2\pi)$$

$$\tilde{u} \in \tilde{\mathcal{F}}_h(0, 2\pi) \text{ and } \tilde{\phi} \in \tilde{\mathcal{F}}_h(0, 2\pi) \implies \tilde{u}\tilde{\phi} \in \tilde{\mathcal{F}}_h(0, 2\pi)$$

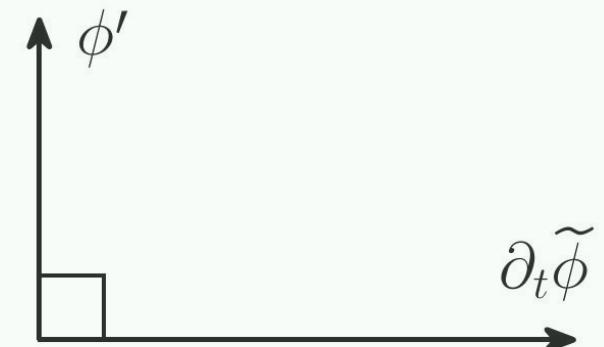
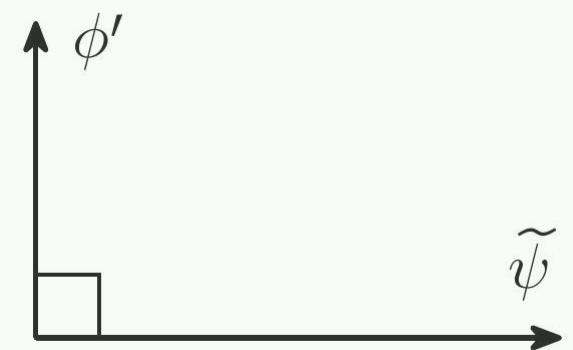
$$\tilde{u} \in \tilde{\mathcal{F}}_h(0, 2\pi) \text{ and } \tilde{\phi} \in \tilde{\mathcal{F}}_h(0, 2\pi) \implies \partial_t \tilde{\phi} \in \tilde{\mathcal{F}}_h(0, 2\pi)$$

No subgrid scales are produced

$$\partial_t \tilde{\phi} \in \tilde{\mathcal{F}}_h(0, 2\pi) \implies (\partial_t \tilde{\phi}, \phi') = 0$$

$$\implies d_t \|\phi'\|^2 = 0$$

$$d_t \|\phi\|^2 = d_t \|\tilde{\phi}\|^2 + d_t \|\phi'\|^2 = d_t \|\tilde{\phi}\|^2$$



Large and small supergrid scales

$$\tilde{\phi}_{i-1/2} = \underbrace{\frac{1}{2}(\tilde{\phi}_{i+1/2} + \tilde{\phi}_{i-1/2})}_{\bar{\phi}_i} + \underbrace{\frac{1}{2}(\tilde{\phi}_{i-1/2} - \tilde{\phi}_{i+1/2})}_{\phi_i^>}$$

Orthogonal decomposition

$$\bar{\phi} \cdot \phi^> = 0$$

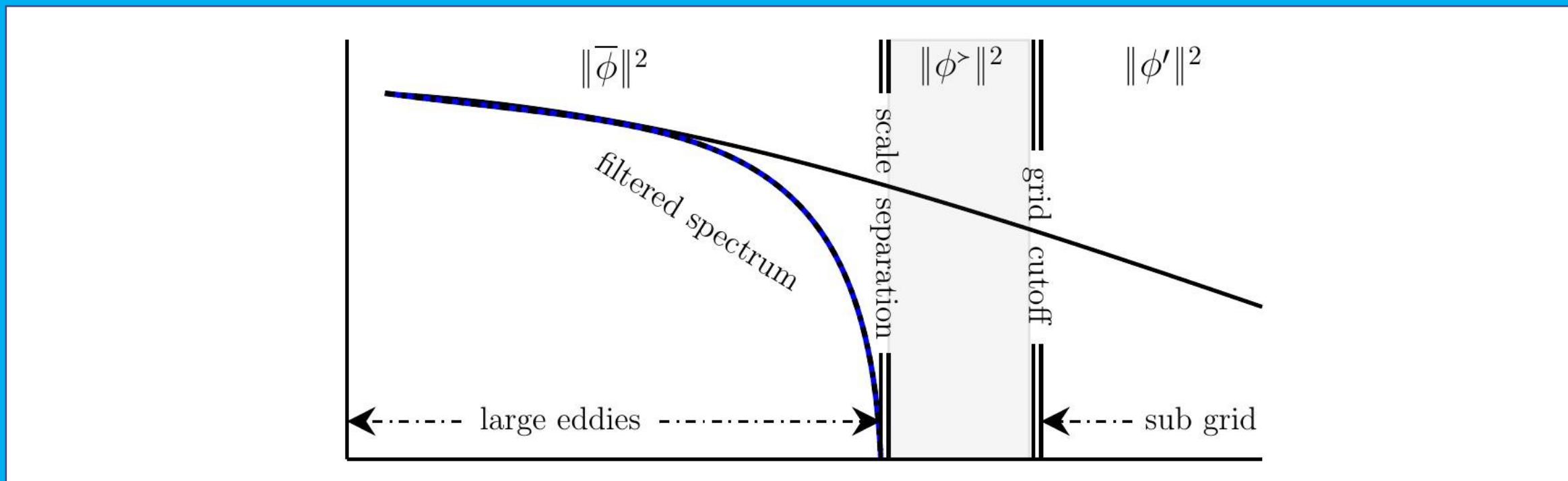
Dissipation

$$d_t \|\phi\|^2 = d_t \|\tilde{\phi}\|^2 = d_t \|\bar{\phi}\|^2 + d_t \|\phi^>\|^2$$

Dissipation condition

$$d_t \|\phi\|^2 = d_t \|\bar{\phi}\|^2 + d_t \|\phi^>\|^2$$

~~$d_t \|\phi^>\|^2$~~



Small supergrid scales

$$\phi_i^> = \frac{1}{2}(\tilde{\phi}_{i-1/2} - \tilde{\phi}_{i+1/2})$$

$$\phi^> = S\tilde{\phi}$$

$$\partial_t \tilde{\phi} = \nabla^\top (\underbrace{\Phi + \tau}_\text{flux})$$

$$\partial_t \phi^> = \underbrace{S \nabla^\top (\Phi + \tau)}_{\nabla_s^\top}$$

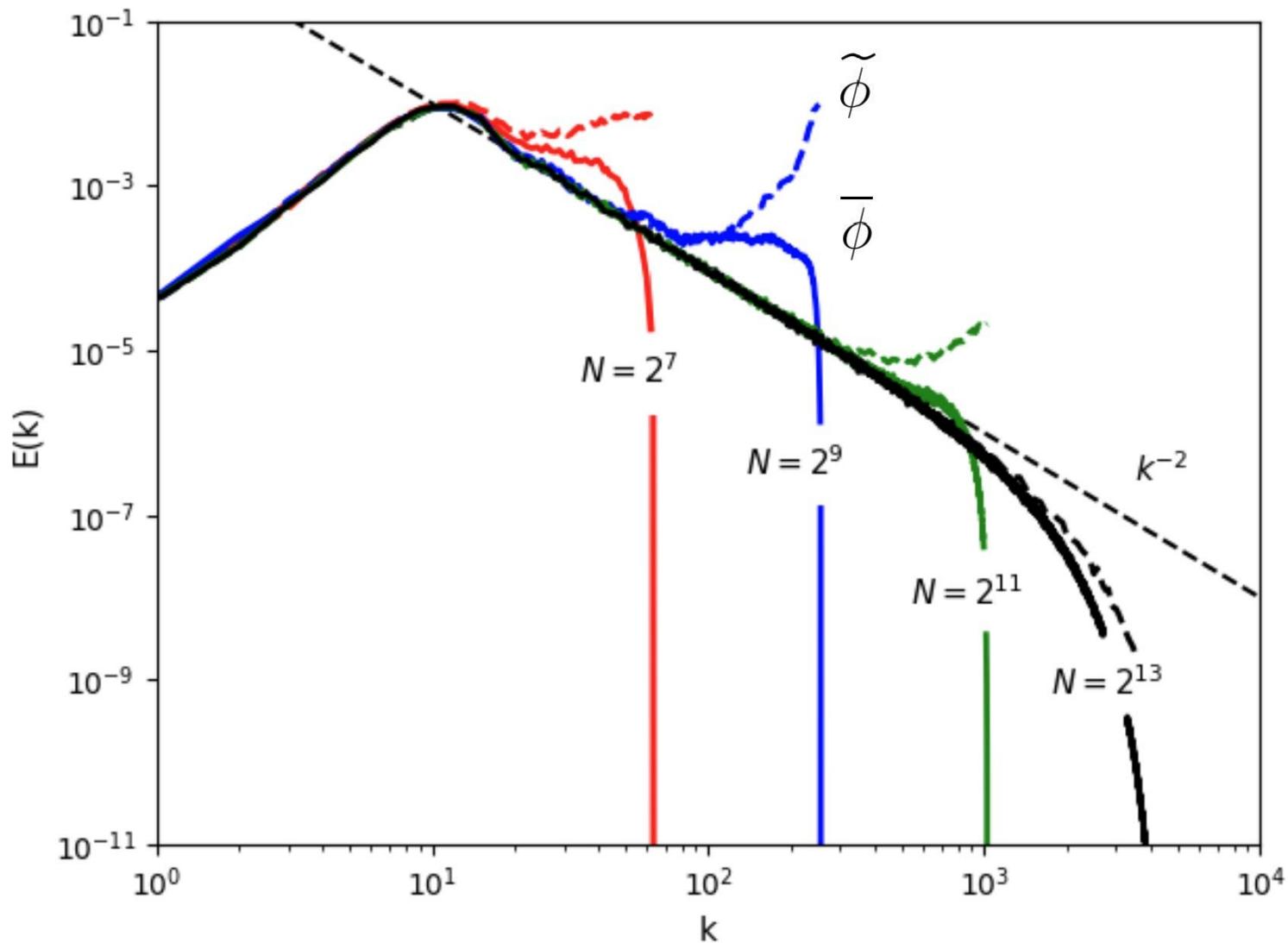
Model

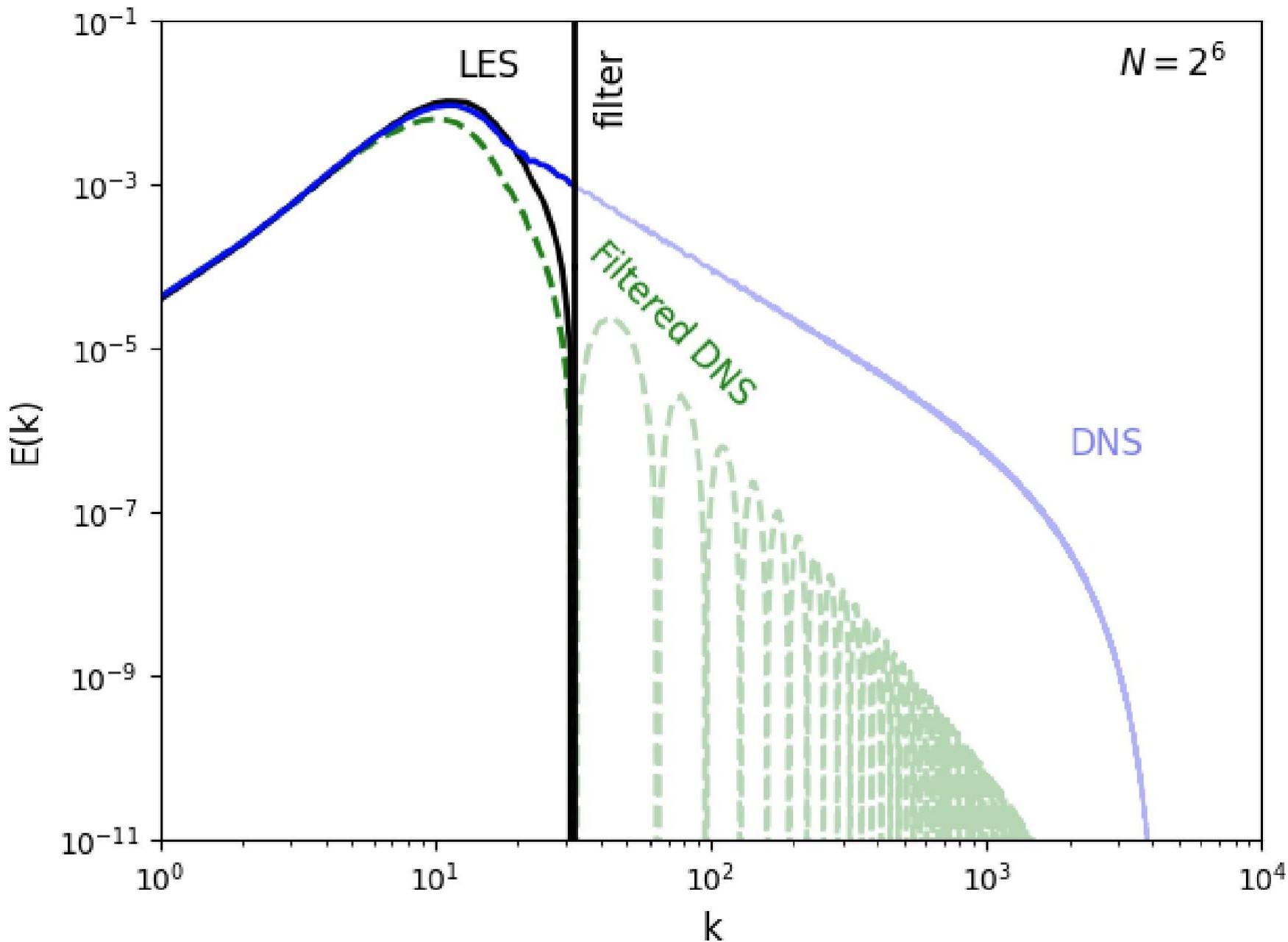
$$\partial_t \phi^> = \nabla_s^\top (\Phi + \tau)$$

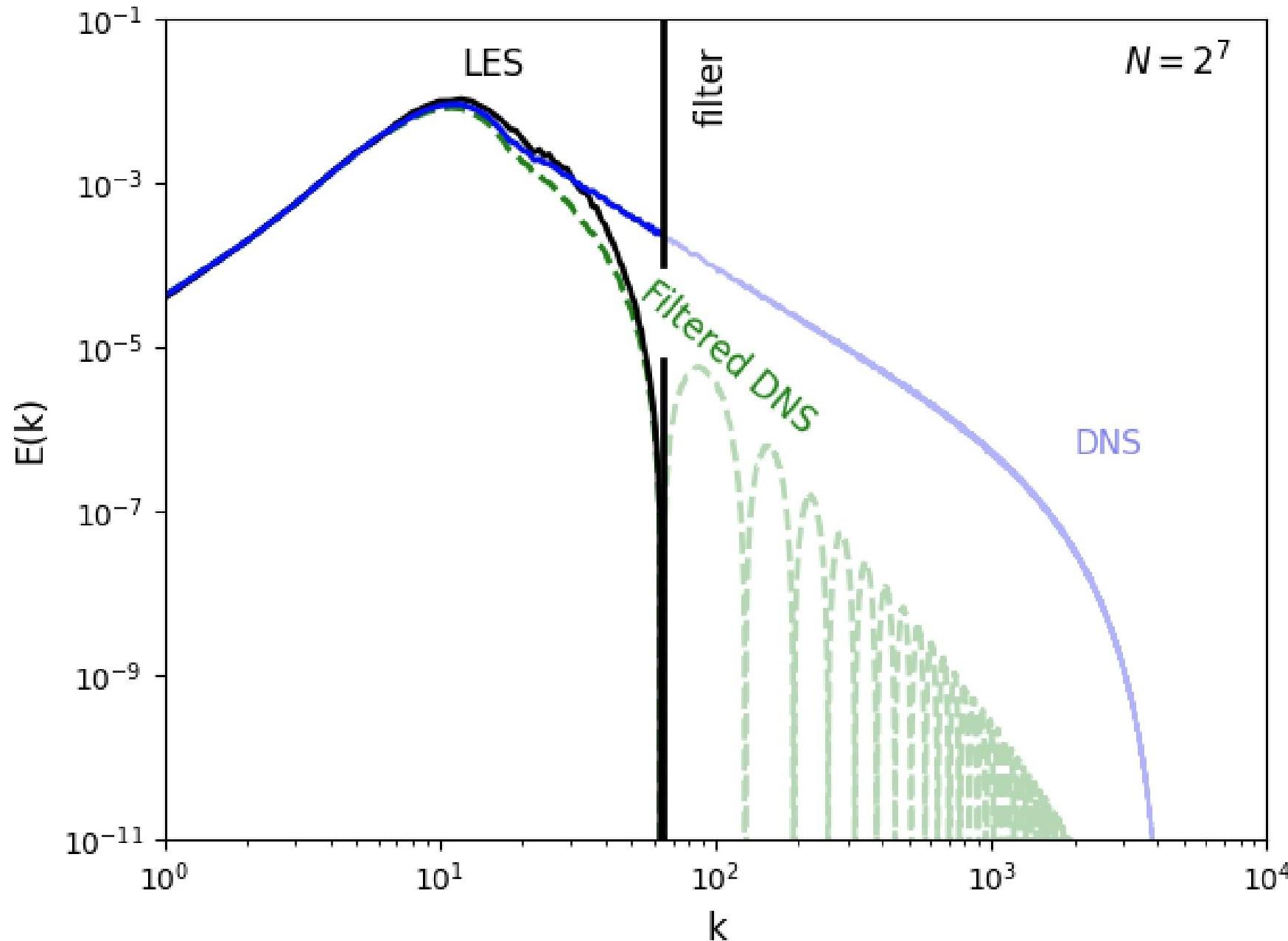
$$d_t \phi^> \cdot \phi^> = (\Phi + \tau) \cdot \nabla_s \phi^> = 0$$

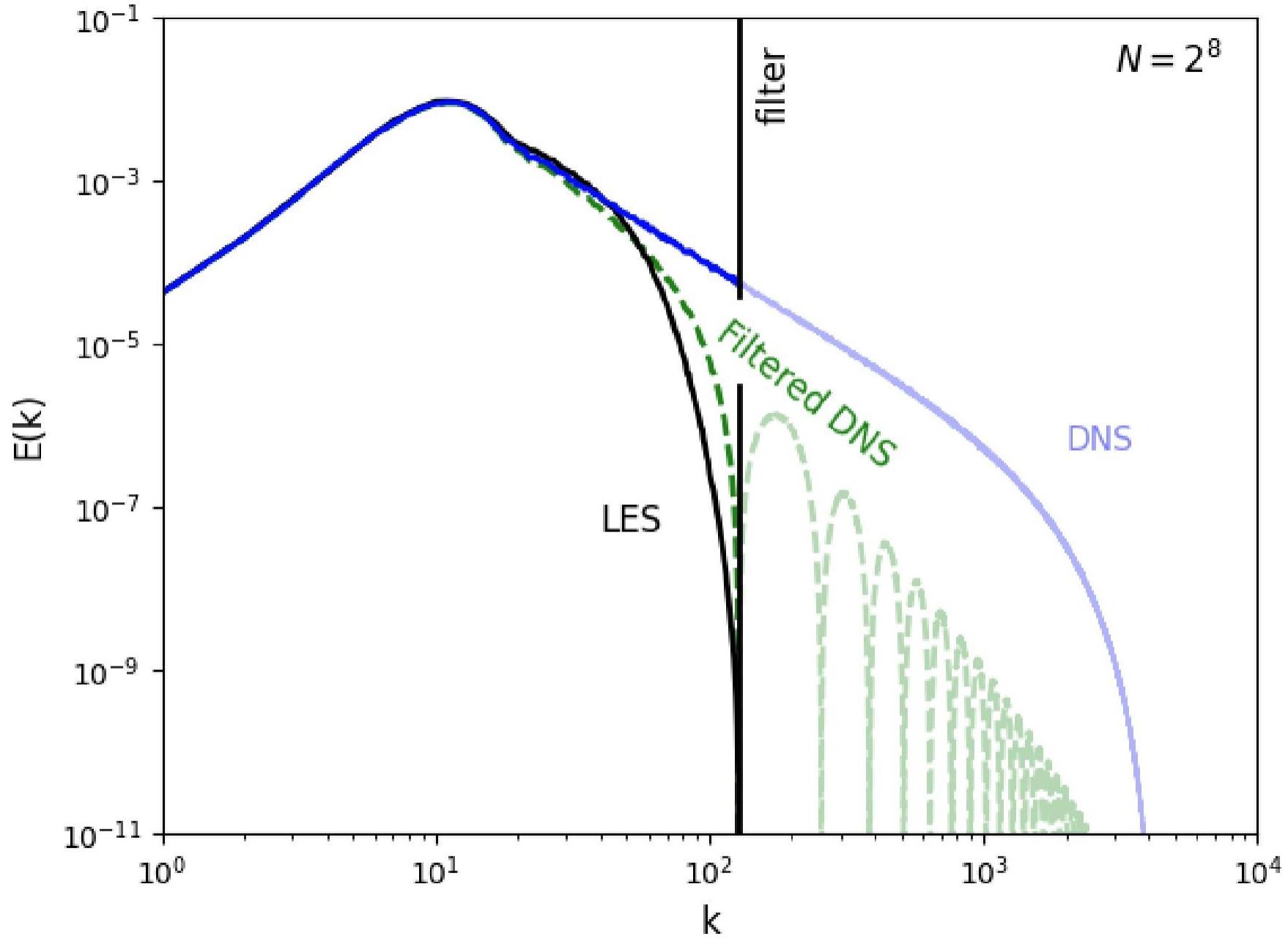
$$\tau = -\alpha \nabla_s \phi^> \quad \alpha = \frac{(\nabla_s \phi^> \cdot \Phi)}{\|\nabla_s \phi\|^2}$$

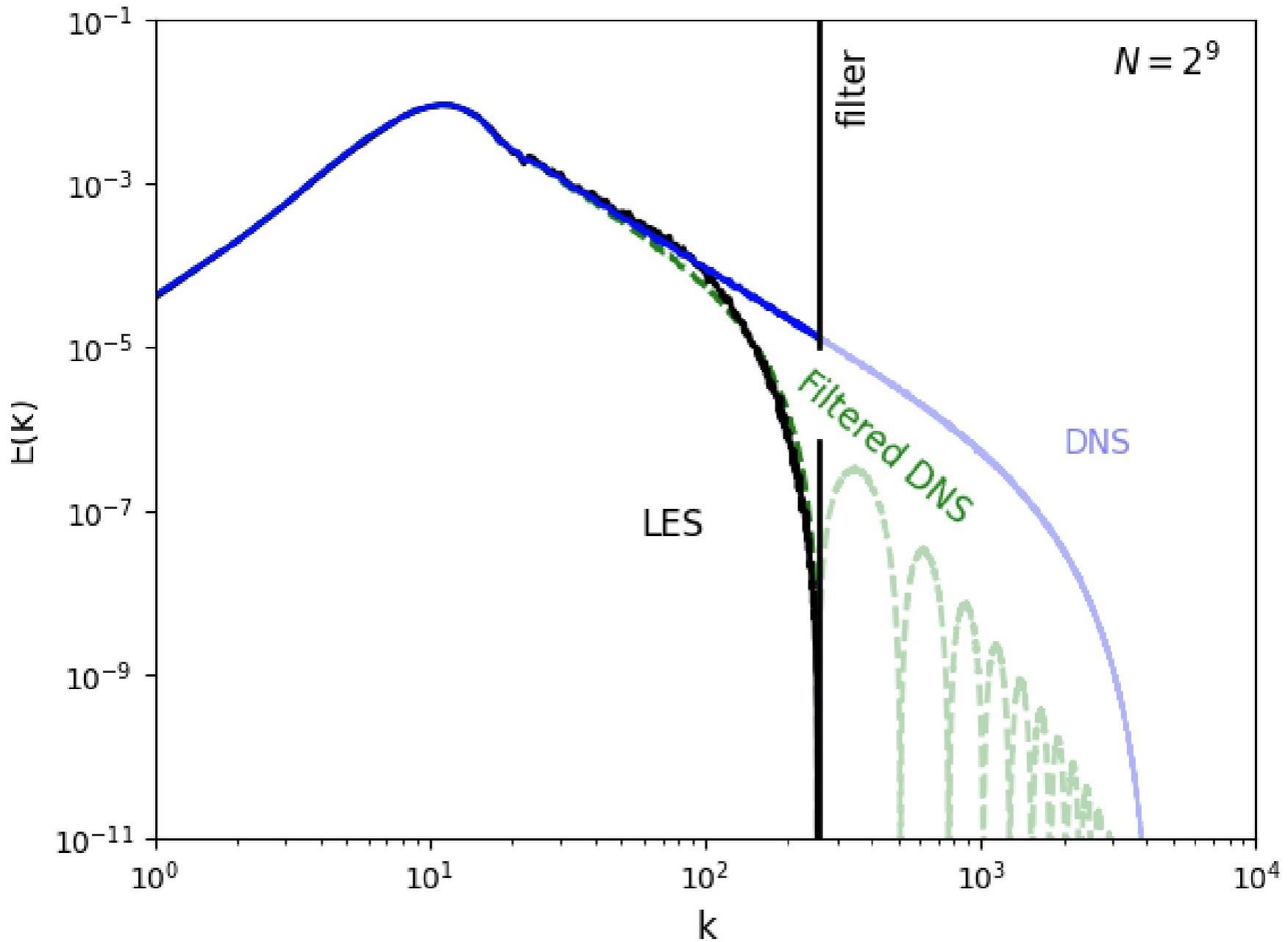
Decaying Burgers' turbulence without model

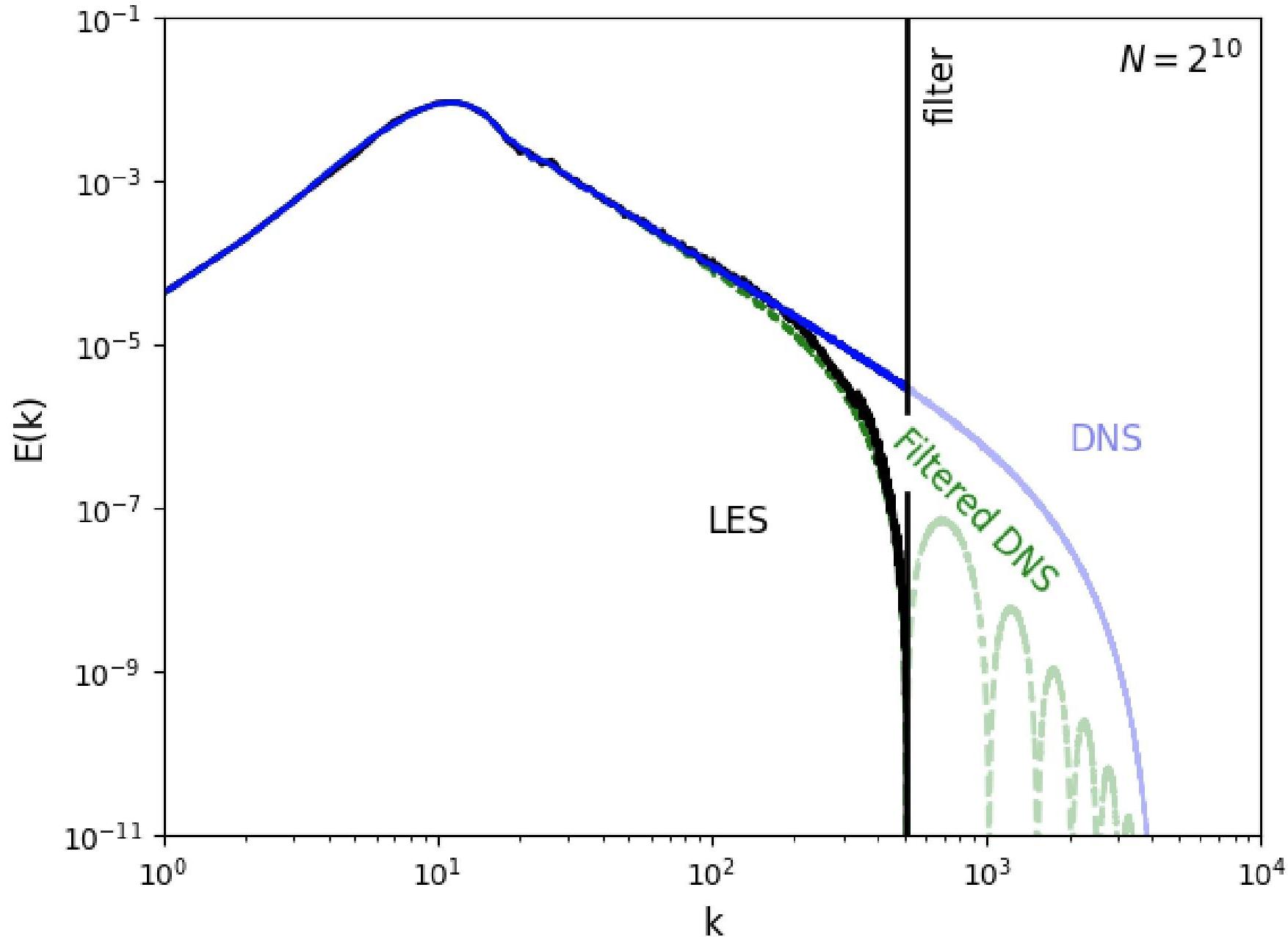








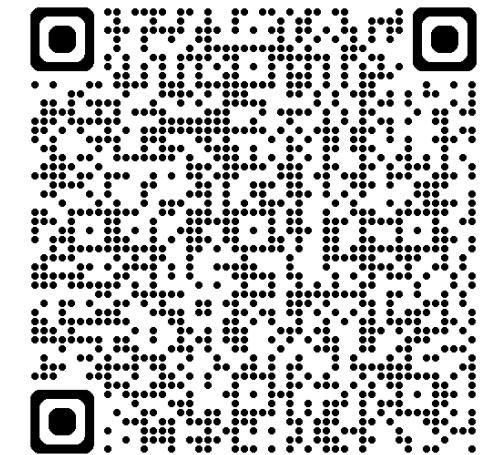
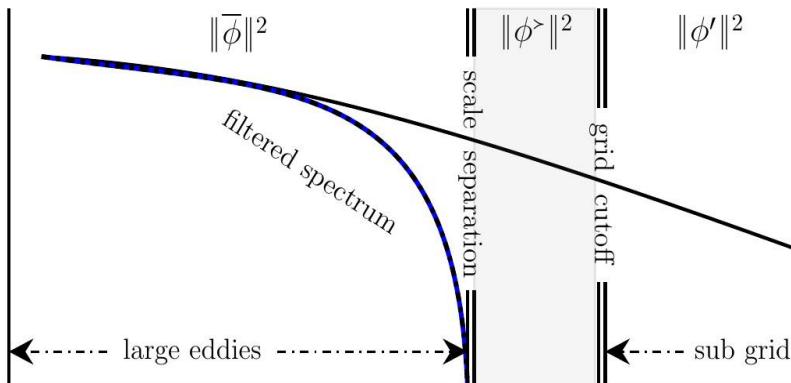




Concluding remarks

FVM-based LES uses two filters that divide the energy into three pieces

Actually no difference between physical and numerical model -> merge



Merging Filtering, Modeling and Discretization to Simulate Large Eddies in Burgers' Turbulence

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Model successfully applied
to decaying Burgers' turbulence

To do: 1D -> 3D, nonuniform grids, etc.