Boundary value problems (DRAFT October 8, 2019)

Introduction

Let I = (a, b) be a non-trivial bounded interval. An important class of differential equations (historically, theoretically, didactically and practically) are of the form

$$-(pu')' + qu = f \qquad \text{on} \qquad I \tag{1}$$

where $u: I \to \mathbb{R}$ is the unknown function, while p, q and g are known, sufficiently smooth functions. Typically,

$$p \ge 0$$
 is non-negative. (2)

Example. Consider $u'' + a_1 x' + a_0 x = g$. Set $p(s) := \exp(\int_a^s a_1(\tau) d\tau)$ and $q := -pa_0$.

To streamline the notation we introduce the differential operator L,

$$Lu := -(pu')' + qu. \tag{3}$$

Thus, Eqn. (1) becomes

$$Lu = f. (4)$$

The equation (1) is linear in u and involves two derivatives. From the ODE theory we expect that the homogeneous equation

$$L\varphi = 0 \tag{5}$$

has two (linearly independent) solutions φ_1 and φ_2 , forming a so-called fundamental system. If, moreover, φ_0 is a particular solution to the original non-homogeneous equation (4) then so is any combination, for any $c_1, c_2 \in \mathbb{R}$,

$$L(\varphi_0 + c_1\varphi_1 + c_2\varphi_2) = L\varphi_0 + c_1L\varphi_1 + c_2L\varphi_2 = q.$$
(6)

The solution space is two-dimensional, parameterized by the two constants $c_1, c_2 \in \mathbb{R}$. To fully specify the solution, two additional constraints are needed.

The operator L is self-adjoint (almost)

The particular form of L is significant for the following reason, called "Lagrange identity" in a more general context.

Observation. For any two smooth functions u and v,

$$uLv - vLu = -u(pv')' + v(pu')' = (p(u'v - uv'))'.$$
(7)

The significance of the observation is further seen by defining the bilinear form

$$B(u,v) := \int_{I} (Lu)(x)v(x)dx. \tag{8}$$

Flipping L onto v using (7), we have the identity

$$B(u,v) = B(v,u) + \int_{I} (p(u'v - uv'))'dx,$$
(9)

which says that B is almost symmetric but for the last anti-symmetric term. However, if we restrict u and v to the vector space of functions that vanish on the boundary of I (for example), the bilinear form is symmetric. This unleashes operator-theoretic tools that form the basis of the finite element method for boundary value problems like (4).

Boundary conditions I

In *initial value problems*, we specify two conditions on one end of the interval, such as $u(a) := u_0$ and $u'(a) := u_1$. Such equations typically model oscillatory evolution over time.

In *boundary value problems*, we specify conditions at both ends of the interval, called boundary conditions. For example,

homogeneous Dirichlet condition at
$$a: D_a u := u(a) \stackrel{!}{=} 0$$
 (10)

homogeneous Neumann condition at
$$b: N_b u := u'(b) \stackrel{!}{=} 0.$$
 (11)

More generally, so-called Robin boundary conditions can be imposed:

$$H_a u := \alpha_0 u(a) + \alpha_1 u'(a) \stackrel{!}{=} \eta_a$$
 (12a)

$$H_b u := \beta_0 u(b) + \beta_1 u'(b) \stackrel{!}{=} \eta_b$$
 (12b)

where α_i , β_i and η_x are constants. Note that H_a and H_b are linear operators that take a function and return a real number (if defined). In more general boundary value problems posed on d-dimensional domains, the boundary operator returns a function (called trace of u) on the (d-1)-dimensional boundary.

For the candidate solution in the form $u = \varphi_0 + c_1\varphi_1 + c_2\varphi_2$ we can write the set of boundary conditions in matrix-vector form,

$$\begin{pmatrix} H_a \varphi_1 & H_a \varphi_2 \\ H_b \varphi_1 & H_b \varphi_2 \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \begin{pmatrix} \eta_a \\ \eta_b \end{pmatrix} - \begin{pmatrix} H_a \varphi_0 \\ H_b \varphi_0 \end{pmatrix}, \tag{13}$$

providing conditions on the unknown coefficients c_1 and c_2 . If the determinant of the matrix is non-zero, these coefficients are uniquely determined (but otherwise there may be no solution, one solution or infinitely many solutions).

Example. Take p := 1 and $q := -(\pi k)^2$, where $k \in \mathbb{N}$, and g = 0, on the interval (0, 1) subject to the homogeneous Dirichlet boundary conditions u(0) = 0 and u(1) = 0. There is no need for the particular solution, in other words $\varphi_0 = 0$. Determine c_1 and c_2 .

Boundary conditions II

In practice, that is in algorithms and software packages for boundary value problems, there are other ways to impose the (linear) boundary conditions (12). A relatively easy way is to decompose $u = \tilde{u} + u_0$ where u_0 is any function that satisfies the boundary conditions (12). Then, setting $\tilde{f} := f - Lu_0$, the function \tilde{u} satisfies the non-homogeneous boundary value problem

$$L\tilde{u} = \tilde{f} \tag{14}$$

with homogeneous boundary conditions $H.\tilde{u} = 0$.

Example. On I = (-1, 1), take

$$Lu := -u'' + u = 0 \quad \text{with} \quad u(\pm 1) = \pm 1.$$
 (15)

Then $u_0(x) := x$ satisfies the boundary conditions. Therefore, if \tilde{u} solves

$$(L\tilde{u})(x) \stackrel{!}{=} -(Lu_0)(x) = x \quad \forall x \in (-1,1) \quad \text{with} \quad \tilde{u}(\pm 1) = 0,$$
 (16)

then $u := \tilde{u} + u_0$ solves the original problem.

In this way, we do not need the fundamental system $\{\varphi_1, \varphi_2\}$, which is in general hard to find anyway. In the following we therefore focus on homogeneous boundary conditions.

Green's function

We consider boundary value problems Lu = f (4) on an interval I with homogeneous boundary conditions. The Green's function is largely a theoretical device for representing solutions that is based on the following idea (that also works in higher dimensions).

Let δ denote the Dirac functional at 0 that formally verifies

$$\int_{I} f(y)\delta(y-x)dy = f(x) \tag{17}$$

for all smooth functions f.

Example. Let H be the Heaviside function with H(x) = 0 for x < 0 and H(x) = 1 for x > 0. Show that

$$\int_{\mathbb{R}} f(y)H'(y-x)dy = f(x) \quad \forall f \in C^1(\mathbb{R}).$$
 (18)

The Green's function to the operator L is a scalar-valued function $G: I \times I \to \mathbb{R}$ of two variables such that

$$L[G(\cdot, y)](x) = \delta(y - x) \quad \forall x \in I, \tag{19}$$

where we agree that L acts on the first variable of G. Now set

$$u(x) := \int_{I} G(x, y) f(y) dy. \tag{20}$$

Then, bravely exchanging differentiation in x and integration in y,

$$(Lu)(x) \stackrel{(20)}{=} \int_{I} L[G(\cdot, y)](x) f(y) dy \stackrel{(19)}{=} \int_{I} \delta(y - x) f(y) dy \stackrel{(17)}{=} f(x) \quad \forall x \in I.$$
 (21)

Remark. If $x \mapsto G(x,y)$ satisfies the homogeneous boundary conditions then so does (20).

Green's function for Lu = -(pu')' + qu

For the one-dimensional problem (4) with homogeneous boundary conditions a Green's function can be constructed from the fundamental system $\{\varphi_1, \varphi_2\}$. Recall that this means in particular that $L\varphi_i = 0$.

First, note that the Lagrange identity (7) implies that

$$c := (\varphi_1 \varphi_2' - \varphi_1' \varphi_2) p \quad \text{is constant.}$$
 (22)

We assume that this constant is non-zero.

Now set

$$G(x,y) := \begin{cases} \frac{1}{c}\varphi_1(x)\varphi_2(y) & \text{if } x \ge y\\ \frac{1}{c}\varphi_2(x)\varphi_1(y) & \text{if } x \le y. \end{cases}$$
 (23)

Remark. The function u defined by (20) satisfies the homogeneous boundary conditions because $G(\cdot, y)$ does (for any $y \in I$).

Remark. The function $x \mapsto G(x, y)$ is continuous but likely not differentiable at x = y. We write G' to mean the derivative with respect to the first variable.

Now we verify that G is indeed a Green's function for L by checking (19). Specifically, we check that

$$\int_{y-\varepsilon}^{y+\varepsilon} L[G(\cdot,y)](x)dx \to 1 \quad \text{as} \quad \varepsilon \searrow 0 \qquad \forall y \in I.$$
 (24)

Since the integrand is zero whenever $x \neq y$, this implies that it equals $\delta(x - y)$, which is equivalent to (19).

We assume that φ_i , φ_i' and p are continuous. Fix $y \in I$. Now, for $\varepsilon > 0$ small,

$$LHS(24) = \int_{y-\varepsilon}^{y+\varepsilon} \left\{ -(pG'(\cdot, y))'(x) + q(x)G(x, y) \right\} dx$$
 (25)

$$= \int_{y+\varepsilon}^{y-\varepsilon} (pG'(\cdot,y))'(x)dx + \mathcal{O}(\varepsilon)$$
 (26)

$$= pG'(x,y)|_{x=y+\varepsilon}^{x=y-\varepsilon} + \mathcal{O}(\varepsilon)$$
(27)

$$= \frac{1}{c} (p(y-\varepsilon)\varphi_2'(y-\varepsilon)\varphi_1(y) - p(y+\varepsilon)\varphi_1'(y+\varepsilon)\varphi_2(y)) + \mathcal{O}(\varepsilon)$$
 (28)

$$\rightarrow \frac{1}{c}p(y)(\varphi_2'\varphi_1 - \varphi_1'\varphi_2)(y) \quad \text{as} \quad \varepsilon \searrow 0$$
 (29)

$$= 1$$
 by (22) . (30)

Example. TODO: hat/tent