

Marino, “Instantons and Large N ”

Okuto Morikawa

*Department of Physics, Osaka University, Toyonaka, Osaka 560-0043,
Japan*

**E-mail: o-morikawa@het.phys.sci.osaka-u.ac.jp*

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Subject Index

1 Definitions

We consider

$$f(z) = \int \mathcal{D}x e^{-S[x;z]}, \quad (1.1)$$

where

$$S[x; z] = S_0[x] + z S_I[x]. \quad (1.2)$$

Perturbation series:

$$f_{\text{PT}}(z) = \sum_{k=0}^{\infty} a_k z^k, \quad (1.3)$$

$$a_k = \frac{1}{k!} \int \mathcal{D}x e^{-S_0[x]} (-S_I[x])^k. \quad (1.4)$$

$$\text{PT} : \int \mathcal{D}x \sum_{k=0}^{\infty} \rightarrow \underbrace{\sum_{k=0}^{\infty} \int \mathcal{D}x}_{\text{Asymptotic}} \quad (1.5)$$

1.1 Borel resummation

◦ Borel transform:

$$B(u) \equiv \sum_{k=0}^{\infty} \frac{a_k}{k!} u^k. \quad (1.6)$$

◦ Borel resummation:

$$f_{\text{Borel}}(z) \equiv \frac{1}{z} \int_0^{\infty} du B(u) e^{-u/z} \quad (1.7)$$

$$= \int_0^{\infty} du B(zu) e^{-u} \quad (1.8)$$

Noting that

$$k! = \int_0^{\infty} du u^k e^{-u}, \quad (1.9)$$

we have

$$f(z) = \int \mathcal{D}x \sum_{k=0}^{\infty} \int_0^{\infty} du e^{-S_0[x]} \frac{1}{(k!)^2} (-z S_I[x])^k u^k e^{-u}. \quad (1.10)$$

$$\text{Borel} : \int \mathcal{D}x \sum_{k=0}^{\infty} \int_0^{\infty} du \rightarrow \underbrace{\sum_{k=0}^{\infty} \int \mathcal{D}x \int_0^{\infty} du}_{\text{Asymptotic}} \rightarrow \int_0^{\infty} du \underbrace{\sum_{k=0}^{\infty} \int \mathcal{D}x}_{u \sim 0} \quad (1.11)$$

1.2 't Hooft

Suppose that $S[x; z] \geq 0$ for $z > 0$.

$$f(z) = \int \mathcal{D}x \int_0^\infty dt \delta(t - S[x; z]) e^{-S[x; z]} \quad (1.12)$$

$$f_\delta(z) \equiv \int_0^\infty dt F_\delta(t; z) e^{-t}, \quad (1.13)$$

$$F_\delta(t; z) \equiv \int \mathcal{D}x \delta(t - S[x; z]) \quad (1.14)$$

$$= \sum_{x_i} \left| \frac{\delta S[x; z]}{\delta x} \right|^{-1} \bigg|_{x=x_i \text{ s.t. } S[x_i; z] = t} \quad (1.15)$$

When $x \in \mathbb{R}$, this is a change of variables: $x \mapsto t$.

1.2.1 Saddle points

For $x \sim x_0$ such that $S[x_0; z] = 0$ and $\delta S[x; z]/\delta x|_{x=x_0} = 0$, we have

$$a_0 \equiv \frac{\delta^2 S[x; z]}{\delta x^2} \bigg|_{x=x_0} > 0 \quad (1.16)$$

$$S[x_0 + \epsilon; z] \sim \frac{1}{2} a_0 \epsilon^2 = t \quad \epsilon = \pm \sqrt{2t/a_0} \equiv \epsilon_\pm \quad (1.17)$$

and then

$$\sum_{\epsilon=\epsilon_\pm} \left| \frac{\delta S[x; z]}{\delta x} \right|_{x=x_0+\epsilon}^{-1} \quad (1.18)$$

$$= \sum_{\epsilon=\epsilon_\pm} \left| \epsilon \frac{\delta}{\delta x} \frac{\delta S[x; z]}{\delta x} \right|_{x=x_0}^{-1} = \sqrt{\frac{2}{a_0 t}}. \quad (1.19)$$

$\int_0 dt/\sqrt{t}$ is convergent near $x = x_0$! (We see a branch point at $t = 0$.) Similarly, for $x \sim x_*^\pm$ such that $\delta S[x; z]/\delta x|_{x=x_*^\pm} = 0$, and

$$c_*^\pm \equiv S[x_*^\pm; z] > 0 \quad a_*^\pm \equiv \frac{\delta^2 S[x; z]}{\delta x^2} \bigg|_{x=x_*^\pm} \geq 0, \quad (1.20)$$

$\int_{c_*^+} dt/\sqrt{t - c_*^+}$ and $\int_{c_*^-} dt/\sqrt{c_*^- - t}$ are convergent near $x = x_*$.

Example.

$$S[x] = (x^2 - 1)^2 \quad f \sim 1.97 \quad (1.21)$$

$$F_\delta(t) = \begin{cases} \frac{1}{2\sqrt{t}} \left(\frac{1}{\sqrt{1+\sqrt{t}}} + \frac{1}{\sqrt{1-\sqrt{t}}} \right) & \text{for } 0 < t < 1 \\ \frac{1}{2\sqrt{t}} \frac{1}{\sqrt{1+\sqrt{t}}} & \text{for } t > 1 \end{cases} \quad (1.22)$$

$$f_\delta \sim 1.97 \quad (1.23)$$

$$\int_{-\sqrt{2}}^{\sqrt{2}} e^{-S} \sim 1.88 \quad \int_0^1 dt F_\delta(t) \sim 1.88 \quad (1.24)$$

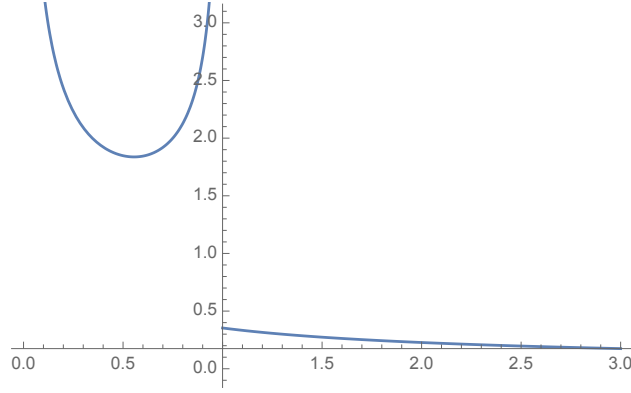


Fig. 1 Eq. (1.22). $F(t) \xrightarrow{x \rightarrow 0} \frac{1}{\sqrt{t}}$, $F(t) \xrightarrow{x \rightarrow 1^-} \frac{1}{\sqrt{2(1-t)}} + \frac{1}{2\sqrt{2}}$, and $F(t) \xrightarrow{x \rightarrow 1^+} \frac{1}{2\sqrt{2}}$

2. 't Hooft \leftrightarrow Borel

2.1. Inverse Laplace transform

“Borel”.

$$f(z) = \int \mathcal{D}x \sum_{k=0}^{\infty} \int_0^{\infty} \frac{du}{z} e^{-S_0[x]} \frac{1}{(k!)^2} (-zS_I[x])^k \left(\frac{u}{z}\right)^k e^{-u/z} \quad (2.1)$$

$$= \frac{1}{z} \int_0^{\infty} du \int \mathcal{D}x e^{-S_0[x]} e^{-u/z} \sum_{k=0}^{\infty} \frac{1}{(k!)^2} (-uS_I[x])^k \quad (2.2)$$

$$= \frac{1}{z} \int_0^{\infty} du e^{-u/z} \underbrace{\int \mathcal{D}x e^{-S_0[x]} J_0(2\sqrt{uS_I[x]})}_{\mathcal{B}(u)} \quad (2.3)$$

Not asymptotic expansion!

't Hooft. Suppose that there exists a mapping,

$$x \mapsto y = y(x), \quad (2.4)$$

such that

$$f(z) = \int \mathcal{D}x e^{-S[x;z]} = \sigma(z) \int \mathcal{D}y e^{-\frac{1}{z}S[y]}. \quad (2.5)$$

$$f(z) = \sigma(z) \int_0^\infty dt e^{-t/z} \underbrace{\int \mathcal{D}y \delta(t - S[y])}_{\mathcal{F}_\delta(t)} \quad (2.6)$$

We define the inverse Laplace transform of the overall factor

$$z\sigma(z) = \int_0^\infty dt e^{-t/z} \mathcal{G}(t). \quad (2.7)$$

Then,

$$f(z) = \frac{1}{z} \int_0^\infty dt e^{-t/z} (\mathcal{F}_\delta * \mathcal{G})(t) \quad (2.8)$$

Here “*” denotes the convolution,

$$(f * g)(t) = \int_0^t ds f(s)g(t-s) \quad (2.9)$$

Example.

$$S[x; z] = \frac{x^2}{2} + z \frac{x^4}{4} \quad (2.10)$$

$$\frac{dS}{dx} = x + zx^3 \quad (2.11)$$

$$t = S[x; z] = \frac{z}{4} \left(\frac{1}{z} + x^2 \right)^2 - \frac{1}{4z} \quad (2.12)$$

$$x^2 + \frac{1}{z} = \pm \frac{1}{z} \sqrt{(4tz + 1)} \quad (2.13)$$

$$x = \pm \sqrt{-\frac{1}{z} \pm \frac{1}{z} \sqrt{4tz + 1}} \quad (2.14)$$

$$\mathcal{F}_\delta(t) = \int dx \delta(t - S[y]) = \frac{2}{x_*(1 + x_*^2)} = 2 \sqrt{\frac{1}{-1 + \sqrt{4t + 1}}} \frac{1}{\sqrt{4t + 1}} \quad (2.15)$$

for $t > 0$, and

$$\mathcal{G}_\delta(t) = \frac{1}{\sqrt{\pi t}} \quad (2.16)$$

$$\mathcal{B}(u) = \int dx e^{-x^2/2} J_0(2\sqrt{ux^4/4}) = 2\sqrt{\frac{2}{\pi}}(1+4u)^{-1/4} K(\frac{1}{2} - \frac{1}{2\sqrt{1+4u}}) \quad (2.17)$$

where $K(z)$ is the complete elliptic integral of the first kind.

In the figure below, we can see $\mathcal{F} * \mathcal{G} = \mathcal{B}$ for $t > 0$. We have $f = f_{\mathcal{B}} = f_\delta \sim 1.94$ ($z = 1$).

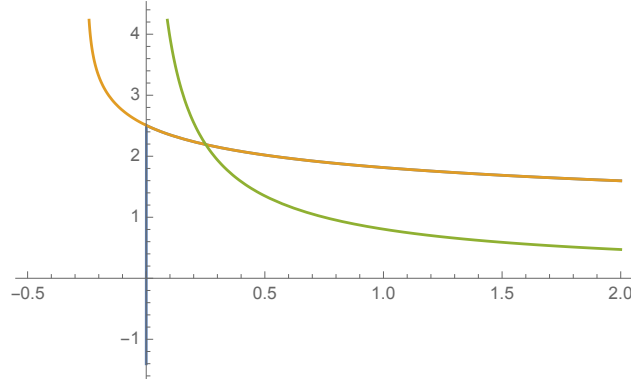


Fig. 2 $(\mathcal{F} * \mathcal{G})(t)$ (blue) and $\mathcal{B}(t)$ (orange) [and $\mathcal{F}(t)$ (green)].

2.2. 0D models

$$S[x; g^2] = \frac{x^2}{2} + \frac{1}{g^2} V[gx], \quad (2.18)$$

$$V[x] = V_3 x^3 + V_4 x^4 + \dots \quad (2.19)$$

and $S[x] \equiv S[x; 1]$.

$$f(z) = \int dx e^{-S[x; g^2]} \quad (2.20)$$

$$= \frac{1}{\sqrt{g^2}} \int dx e^{-\frac{1}{g^2} S[x]} \quad (2.21)$$

$$= \frac{1}{g^2} \sqrt{g^2} \int_0^\infty dt e^{-t/g^2} \underbrace{\int dx \delta(t - S[x])}_{F(t)} \quad (2.22)$$

Noting that

$$\sqrt{g^2} = \int_0^\infty dt e^{-t/g^2} \underbrace{\frac{1}{\sqrt{\pi t}}}_{G(t)}, \quad (2.23)$$

we have

$$f(z) = \frac{1}{g^2} \int_0^\infty dt e^{-t/g^2} (F * G)(t) \quad (2.24)$$

$$(F * G)(t) = \int_0^t ds F(s) G(t-s) \quad (2.25)$$

$$= \int_0^t ds \int dx \delta(s - S[x]) \frac{1}{\sqrt{\pi(t-s)}} \quad (2.26)$$

$$= \sum_{x_i} \int_0^t ds \left| \frac{dS[x]}{dx} \right|^{-1} \Big|_{x=x_i \text{ s.t. } S[x]=s} \frac{1}{\sqrt{\pi(t-s)}} \quad (2.27)$$

◦ For $x \sim x_0$ ($t \gtrsim 0$),

$$\int_0^t ds \frac{1}{\sqrt{s(t-s)}} = \pi \quad (\text{const. and finite}) \quad (2.28)$$

◦ For $x \sim x_*^+$ ($t \gtrsim c_*^+$)

$$\int_{c_*^+}^t ds \frac{1}{\sqrt{(s-c_*^+)(t-s)}} = \pi \quad (\text{const. and finite}) \quad (2.29)$$

◦ For $x \sim x_*^-$ ($t \lesssim c_*^-$)

$$\int_{(0)}^t ds \frac{1}{\sqrt{(c_*^- - s)(t-s)}} = 2 \operatorname{arccosh} \left(\sqrt{\frac{c_*^-}{c_*^- - t}} \right) \quad (2.30)$$

$$= \ln \left(\frac{4c_*^-}{c_*^- - t} \right) + O(c_*^- - t) \quad (2.31)$$

Example.

$$S[x] = (x^2 - 1)^2 \quad (2.32)$$

$$F(t) = \begin{cases} \frac{1}{2\sqrt{t}} \left(\frac{1}{\sqrt{1+\sqrt{t}}} + \frac{1}{\sqrt{1-\sqrt{t}}} \right) & \text{for } 0 < t < 1 \\ \frac{1}{2\sqrt{t}} \frac{1}{\sqrt{1+\sqrt{t}}} & \text{for } t > 1 \end{cases} \quad (2.33)$$

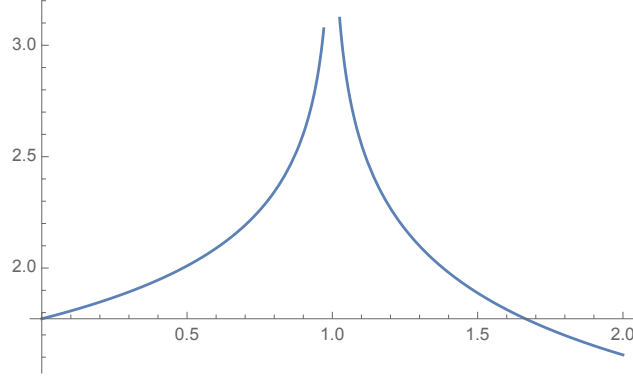


Fig. 3 $F * G$

Example.

$$S[x] = \frac{x^2}{2} - 2\gamma \frac{x^3}{3} + \frac{x^4}{4} \quad (2.34)$$

$$\frac{dS[x]}{dx} = x[1 - \gamma^2 + (x - \gamma)^2] \quad (2.35)$$

$$x_0 = 0, \quad x_*^\pm = \gamma \pm \sqrt{\gamma^2 - 1} \quad (2.36)$$

$$S[x_0] = 0 \quad (2.37)$$

$$S[x_*^\pm] = -\frac{1}{4} + \gamma^2 - \frac{2}{3}\gamma^4 \mp \frac{2}{3}\gamma(\gamma^2 - 1)^{3/2} \quad (2.38)$$

$$t = S[x] = x^2 \left(\frac{1}{2} - 2\gamma \frac{x}{3} + \frac{x^2}{4} \right) \quad (2.39)$$

$$= x^2 \left[\left(\frac{x}{2} - \gamma \frac{2}{3} \right)^2 - \gamma^2 \frac{4}{9} + \frac{1}{2} \right] \quad (2.40)$$

$$\sum_{k=0}^{\infty} \int dx e^{-x^2/2} \frac{1}{k!} \left(2\gamma g \frac{x^3}{3} - g^2 \frac{x^4}{4} \right)^k \quad (2.41)$$

$$= \sum_{k=0}^{\infty} \int dx e^{-x^2/2} \frac{1}{k!} \sum_{m=0}^k {}_k C_m \left(2\gamma g \frac{x^3}{3} \right)^{k-m} \left(g^2 \frac{x^4}{4} \right)^m \quad (2.42)$$

$$= \sum_{k=0}^{\infty} \int dx e^{-x^2/2} \frac{1}{k!} \sum_{m=0}^k {}_k C_m g^{k+m} x^{3k+m} \left(\frac{\gamma}{3} \right)^{k-m} 2^{k-3m} \quad (2.43)$$

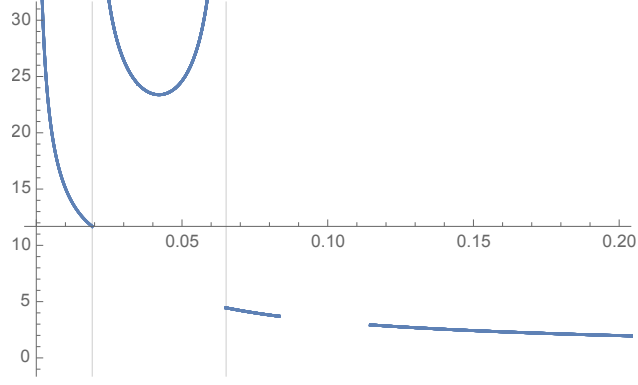


Fig. 4 F with $\gamma = 1.05$

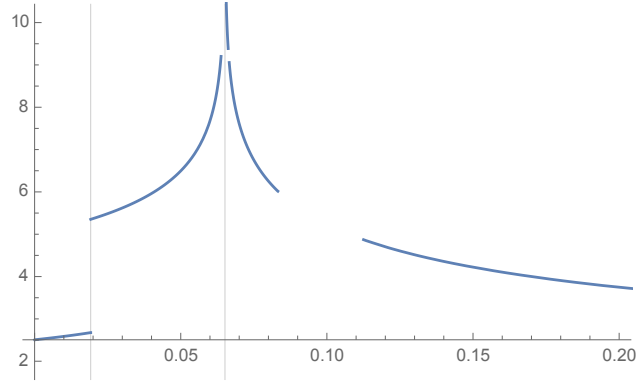


Fig. 5 $F * G$ with $\gamma = 1.05$

$$k' = 2\ell = k + m$$

$$= \sum_{k'=0}^{\infty} \left(\frac{2\gamma}{3} g \right)^{k'} \sum_{m=0}^{\lfloor k'/2 \rfloor} \frac{1}{(k'-m)!} {}^{k'-m}C_m \left(\frac{4\gamma}{3} \right)^{-2m} \int dx e^{-x^2/2} x^{3k'-2m} \quad (2.44)$$

$$= \sum_{\ell=0}^{\infty} \left(\frac{2\gamma}{3} g \right)^{2\ell} \sum_{m=0}^{\ell} \frac{1}{(2\ell-m)!} {}^{2\ell-m}C_m \left(\frac{4\gamma}{3} \right)^{-2m} 2^{3\ell-m+1/2} \Gamma(3\ell-m+\frac{1}{2}) \quad (2.45)$$

$$= \sum_{\ell=0}^{\infty} \left(\frac{2\gamma}{3} g \right)^{2\ell} 2^{3\ell+1/2} \frac{\Gamma(3\ell+\frac{1}{2})}{\Gamma(2\ell+1)} F(\frac{1}{2}-\ell, -\ell; \frac{1}{2}-3\ell; -\frac{9}{8\gamma^2}) \quad (2.46)$$