

Gradient Structures from Classical to Quantum

Oliver Tse

Eindhoven University of Technology

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Abstract.

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1 Stochastic Schrödinger Equation

As we have seen, a quantum circuit is essentially a noisy controlled Schrödinger equation. In this section, we consider classical noise sources arising from the control mechanism in terms of light-matter interactions. Without going into the details of the physics, we introduce the stochastic Schrödinger equation, which captures the evolution of a quantum system driven by classical noise sources.

Let \mathcal{H} be a Hilbert space and $\mathcal{B}(\mathcal{H})$

$$dU_t = HU_t dt + \sum_j V_j U_t \circ d\alpha_t,$$

where H, V_j are

1.1 Quantum measurements

Let $\varrho \in \mathcal{S}(\mathcal{H})$ be a state on a Hilbert space \mathcal{H} and $t \mapsto U_t \in \mathcal{U}(\mathcal{H})$ be a group of unitary actions on \mathcal{H} . Suppose Q is an \mathbb{R} -valued observable with spectral measure $\pi^Q : \mathcal{B}(\mathbb{R}) \rightarrow \mathcal{P}(\mathcal{H})$, i.e., π^Q is a projection-valued measure (PVM). Then the probability of measuring an event $A \in \mathcal{B}(\mathbb{R})$ at time $t^- > 0$ is given by Born's rule:

$$\mathbb{P}_\varrho(\lambda_{t^-}^Q \in A) = \text{tr}[\pi^Q(A)\varrho_{t^-}] = \text{tr}[U_t^\dagger \pi^Q(A)U_t \varrho].$$

As measurements generally change the state of the system, Lüders' rule postulates that, *conditioned* on the event $A \in \mathcal{B}(\mathbb{R})$, the state “collapses” to

$$\varrho_{t^+} = \mathcal{S}(\pi^Q(A), \varrho_{t^-}), \quad \mathcal{S}(\pi, \varrho) := \frac{\pi \varrho \pi^\dagger}{\text{tr}[\pi \varrho \pi^\dagger]} \in \mathcal{S}(\mathcal{H}).$$

Notice that by definition, $\mathcal{S}(\pi, U\varrho U^\dagger) = \mathcal{S}(\pi U, \varrho)$ and $U\mathcal{S}(\pi, \varrho)U^\dagger = \mathcal{S}(U\pi, \varrho)$.

The probability of measuring the event $B \in \mathcal{B}(\mathbb{R})$ at time $t^+ > 0$ given that the event $A \in \mathcal{B}(\mathbb{R})$ was measured at time $t > 0$ is

$$\mathbb{P}_\varrho(\lambda_{t^+}^Q \in B | \lambda_{t^-}^Q \in A) = \text{tr}[\pi^Q(B)\mathcal{S}(\pi^Q(A), \varrho_{t^-})] = \frac{\text{tr}[\pi^Q(A \cap B)\varrho_{t^-}]}{\text{tr}[\pi^Q(A)\varrho_{t^-}]}.$$

In particular, the conditional probability of measuring event $A \in \mathcal{B}(\mathbb{R})$ at time $t^+ > 0$ given it was measured at time $t^- > 0$ is 1. If R is another \mathbb{R} -valued observable, then

$$\mathbb{P}_\varrho(\lambda_{t^+}^R \in B | \lambda_{t^-}^Q \in A) = \text{tr}[\pi^R(B)\varrho_{t^+}] = \frac{\text{tr}[\pi^R(B)\pi^Q(A)\varrho_{t^-}\pi^Q(A)\pi^R(B)]}{\text{tr}[\pi^Q(A)\varrho_{t^-}]}.$$

When R and Q commute, then $\pi^R(B)\pi^Q(A) = \pi^Q(A)\pi^R(B)$, and hence,

$$\mathbb{P}_\varrho(\lambda_{t^+}^R \in B | \lambda_{t^-}^Q \in A)\mathbb{P}_\varrho(\lambda_{t^-}^Q \in A) = \mathbb{P}_\varrho(\lambda_{t^+}^Q \in A | \lambda_{t^-}^R \in B)\mathbb{P}_\varrho(\lambda_{t^-}^R \in B),$$

i.e., Bayes' theorem holds and the usual notion of joint probability is well defined. Clearly, this does not hold when R and Q do not commute. It is this failure of the Bayes theorem that distinguishes quantum probability from classical probability.

In a similar fashion, one determines the probability of measuring the event A_2 at time $t_2^- > 0$ for R given that the event A_1 was measured at time $t_1 > 0$ for Q from

$$\begin{aligned}\mathbb{P}_\varrho(\lambda_{t_2^-}^R \in A_2 | \lambda_{t_1^-}^Q \in A_1) &= \text{tr}[\pi^R(A_2)U_{t_2-t_1}\mathcal{S}(\pi^Q(A_1), \varrho_{t_1^-})U_{t_2-t_1}^\dagger] \\ &= \text{tr}[\pi^R(A_2)U_{t_2-t_1}\mathcal{S}(\pi^Q(A_1)U_{t_1}, \varrho)U_{t_2-t_1}^\dagger] \\ &= \text{tr}[\pi^R(A_2)\mathcal{S}(U_{t_2-t_1}\pi^Q(A_1)U_{t_1}, \varrho)] = \text{tr}[\pi^R(A_2)\varrho_{t_2^-}].\end{aligned}$$

As before, upon measuring the event A_2 at time $t_2 > 0$, Lüders' rule implies

$$\varrho_{t_2^+} = \mathcal{S}(\pi^R(A_2), \varrho_{t_2^-}) = \mathcal{S}(\pi^R(A_2)U_{t_2-t_1}\pi^Q(A_1)U_{t_1}, \varrho).$$

Restricting ourselves to a single observable Q,

2 Noncommutative probability

2.1 From commutative to noncommutative

Consider the Hilbert space $\mathcal{H} = L^2_{\mathbb{C}}(\Omega, \mu)$ over the complete probability space $(\Omega, \mathcal{F}, \mu)$. Then any function $f \in L^\infty_{\mathbb{C}}(\Omega, \mu)$ gives rise to a multiplication operator $M_f \in \mathcal{B}(\mathcal{H})$:

$$M_f g = fg \in \mathcal{H} \quad \forall g \in \mathcal{H},$$

with $\|M_f\|_\infty = \|f\|_{L^\infty(\mu)}$. The collection of all such multiplication operators

$$\mathcal{A} := \{M_f : f \in L^\infty_{\mathbb{C}}(\Omega, \mu)\} \subset \mathcal{B}(\mathcal{H})$$

forms a *commutative* subalgebra of $\mathcal{B}(\mathcal{H})$.

In fact, this subalgebra is a possibly noncommutative *von Neumann* algebra:

Definition 2.1 (von Neumann algebra) A (unital) *von Neumann algebra* is a $*$ -subalgebra $\mathcal{A} \subset \mathcal{B}(\mathcal{H})$ that contains $I_{\mathcal{H}}$ and is closed in the *weak operator topology* (WOT), i.e.,

$$\text{WOT-lim } a_n = a \iff \langle f, a_n g \rangle_{\mathcal{H}} \rightarrow \langle f, a g \rangle_{\mathcal{H}} \quad \forall f, g \in \mathcal{H}.$$

Equivalently, is it a C^* -algebra with a predual \mathcal{A}_* that is a Banach space.

There are sub-families of \mathcal{A} that play a distinguished role, namely,

$$\begin{aligned}\mathcal{A}_+ &:= \{a \in \mathcal{A} : a \succeq 0\} && \text{(nonnegative operators)} \\ \mathcal{O} &:= \{a \in \mathcal{A} : a^\dagger = a\} && \text{(Hermitian operators)} \\ \mathcal{P} &:= \{a \in \mathcal{O} : a^2 = a\} && \text{(projection operators)}\end{aligned}$$

Here, $a \succeq 0$ if and only if $\langle f, a f \rangle_{\mathcal{H}} \geq 0$ for all $f \in \mathcal{H}$. Additionally, for von Neumann algebras, we also have the following useful density result:

Proposition 2.2. Let \mathcal{A} be a von Neumann algebra. Then

$$\mathcal{A} = \overline{\text{span}(\mathcal{P})}^{\text{WOT}},$$

i.e., \mathcal{A} is the WOT-closure of the linear span of its projections.

Let us return to our previous example with $\mathcal{A} = \{\mathsf{M}_f : f \in L_{\mathbb{C}}^\infty(\Omega, \mu)\}$. It is not difficult to see that \mathcal{A} is a W^* -algebra. Notice that since M_f is self-adjoint for $f \in L_{\mathbb{R}}^\infty(\Omega, \mu)$, the family of self-adjoint operators is given by $\mathcal{O} = \{\mathsf{M}_f : f \in L_{\mathbb{R}}^\infty(\Omega, \mu)\}$ and the family of projection operators is given by $\mathcal{P} = \{\mathsf{M}_f \in \mathcal{O} : f = \mathbf{1}_A, A \in \mathcal{F}\}$.

On a von Neumann algebra \mathcal{A} , we can talk about special types of continuous linear functionals on \mathcal{A} , called states.

Definition 2.3 A *state* on a von Neumann algebra \mathcal{A} is a linear functional $\psi : \mathcal{A} \rightarrow \mathbb{C}$ that is positive and normalized, i.e., $\psi(a^\dagger a) \geq 0$ for all $a \in \mathcal{A}$ and $\psi(I_{\mathcal{B}(\mathcal{H})}) = 1$.

A state ψ is said to be

faithful if $\psi(a^\dagger a) = 0 \Leftrightarrow a = 0$,

tracial if $\psi(ab) = \psi(ba)$ for all $a, b \in \mathcal{A}$, and

normal if $\psi \in \mathcal{A}_*$, i.e., if it is an element of the predual \mathcal{A}_* .

For any probability measure $\nu \ll \mu$, we set

$$\psi_\nu(\mathsf{M}_f) := \int_{\Omega} f \, d\nu, \quad f \in L_{\mathbb{C}}^\infty(\Omega, \mu),$$

we find that ψ is a linear functional that is positive and normalized, i.e., ψ is a state. Moreover, it is tracial. It is normal if $\omega := d\nu/d\mu \in L^1(\Omega, \mu)$ and faithful if $\omega > 0$.

Normal states play an essential role, serving as a counterpart to classical measures, as made explicit by the following proposition.

Proposition 2.4. *Let ψ be a state on a von Neumann algebra \mathcal{A} . The following are equivalent:*

(i) ψ is a normal state.

(ii) (σ -additivity) If $(a_n) \subset \mathcal{P}$ are mutually orthogonal projections, i.e., $a_n(\mathcal{H}) \perp a_m(\mathcal{H})$ for all $n \neq m$, and $a = \vee_n a_n$ being the projection on the smallest closed subspace containing $\cup_n a_n(\mathcal{H})$, then

$$\psi(a) = \sum_n \psi(a_n).$$

(iii) (Continuity from below) For any increasing net $0 \preceq a_n \uparrow a$ in \mathcal{A}_+ , one has the increasing limit $\psi(a_n) \uparrow \psi(a)$.

(iv) There exists a family $\{\xi_n\} \subset \mathcal{H}$ with $\sum_n \|\xi_n\|_{\mathcal{H}}^2 = 1$ such that¹

$$\psi = \sum_n \langle \xi_n, \cdot \xi_n \rangle \quad \text{in the sense of norm convergence.}$$

¹[Theorem 7.1.8, Fundamentals of the Theory of Operator Algebras, V.II, Kadison-Ringrose]

Consider an arbitrary normal state ψ and set

$$\mu(A) := \psi(M_{1_A}) \quad \text{for every } A \in \mathcal{F}.$$

Then, clearly, $\mu(\emptyset) = 0$, $\mu(\Omega) = \psi(I_{\mathcal{B}(\mathcal{H})}) = 1$ and μ is σ -additive due to the equivalent characterization of a normal state given by Proposition 2.4(ii). In particular, one obtains a classical measure on (Ω, \mathcal{F}) . In this sense, a state on a noncommutative von Neumann algebra generalizes that of a classical measure.

2.2 Observables

In classical probability, we are often interested in computing expressions like $\mathbb{P}(X \in A)$, where X is a random variable and $A \subset \mathbb{R}$ is a Borel set.

In the quantum world, a random variable is modelled by a Hermitian operator $a \in \mathcal{O}$ and is called an *observable*. Observables have their spectrum in \mathbb{R} and can therefore be *measured*.

Given a ψ on a von Neumann algebra \mathcal{A} , the expectation of an observable a w.r.t. the state ψ is given by $\psi(a)$. Since every observable $a \in \mathcal{O}$ has a spectral decomposition

$$a = \int_{\mathbb{R}} \lambda E_a(d\lambda), \quad E_a \cong \mathcal{P}\text{-valued measure},$$

we can associate a classical probability measure with the observable a and state ψ :

$$\mathbb{P}_\psi(a \in A) := \psi(E_a(A)) \quad \text{for all Borel set } A \subset \mathbb{R}.$$

For two observables $a, b \in \mathcal{O}$ that do not commute, one would like to be able to write $\mathbb{P}_\psi(a \in A, b \in B)$ for two Borel sets $A, B \subset \mathbb{R}$. The simple argument for this is that if $[a, b] \neq 0$, then $E_a(A)$ and $E_b(B)$ may not commute for all Borel sets A, B . In particular,

$$\psi(E_a(A)E_b(B)) \neq \psi(E_b(B)E_a(A)),$$

so there is no consistent way of writing $\mathbb{P}_\psi(a \in A, b \in B)$. Yet, when they do commute, then $\psi(E_a(A)E_b(B)) = \psi(E_b(B)E_a(A))$ and a joint distribution exists and we are back to the classical scenario.

However, this turns out not to be possible, as the following example portrays.

Example 2.5 (Stern-Gerlach measurements) Take a beam of atoms (each with spin- $\frac{1}{2}$) and perform the following steps:

SG_z Measure spin in z -direction: Send the beam through a Stern-Gerlach magnet oriented along z . The beam splits into spin-up $|\uparrow\rangle$ and spin-down $|\downarrow\rangle$ paths. Keep only the spin-up branch, which gives a pure state $|\uparrow_z\rangle$.

SG_x Measure spin in x -direction: Now send the filtered beam through a second magnet, oriented along x . The beam splits again into $|\uparrow_x\rangle$ and $|\downarrow_x\rangle$ outcomes with 50%–50% probability. So far, this can *still* be explained with classical probability.

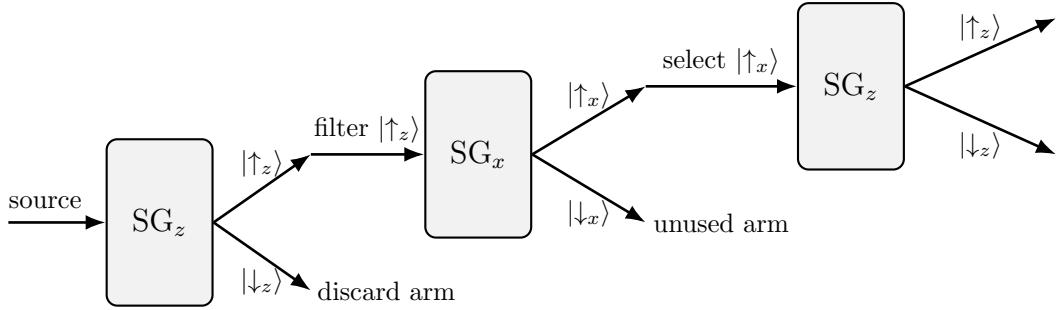


Figure 1: Stern-Gerlach experiment

SG_z Measure spin in z -direction again: Send either of the filtered beams $|\uparrow_x\rangle$ or $|\downarrow_x\rangle$ through a z -magnet again. You do *not* get back the original result, i.e., instead of remaining spin-up, you see another 50%–50% split between $|\uparrow_z\rangle$ and $|\downarrow_z\rangle$

Morally, if spin- x and spin- z were classical random variables, measuring x would *erase* knowledge of z . In classical probability, observing one property never randomizes another unless there is *hidden causal disturbance*.

Quantum mechanically, the two observables (or random variables)

$$\sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \quad \text{spin in } z\text{-direction}, \quad \sigma_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad \text{spin in } x\text{-direction}.$$

do not commute, i.e., $[\sigma_x, \sigma_y] \neq 0$, and measuring σ_x *changes* information about σ_z . In this sense, they cannot be jointly sampled, i.e., there is no classical joint probability $p(\sigma_x, \sigma_z)$ for these observables, which illustrates the need for noncommutative probability.

3 Stochastic Dilation

3.1 Noncommutative random variables

Definition 3.1 A *noncommutative \mathfrak{X} -valued random variable* on a quantum probability space (\mathfrak{A}, μ) is an identity preserving $*$ -homomorphism

$$\mathfrak{z}: \mathfrak{X} \rightarrow \mathfrak{A}.$$

Correspondingly, a *noncommutative \mathfrak{X} -valued stochastic process* on a quantum probability space (\mathfrak{A}, μ) is a family of random variables

$$\mathfrak{z}_t: \mathfrak{X} \rightarrow \mathfrak{A}, \quad t \in \mathbb{T},$$

with \mathbb{T} being a (possibly uncountable) index set.

Example 3.2 (Classical random variable) Let X be a classical E -valued random variable on (Ω, \mathbb{P}) . Setting

$$\mathfrak{X} := B_b(E), \quad \mathfrak{A} := L^\infty(\Omega, \mathbb{P}), \quad \mu(g) := \int_\Omega g \, d\mathbb{P}, \quad g \in \mathfrak{A},$$

we have that

$$\mathfrak{X} \ni f \mapsto \mathfrak{z}(f) := f \circ X \in L^\infty(\Omega, \mathbb{P})$$

is a noncommutative \mathfrak{X} -valued random variable. We see here that the noncommutative notion of a random variable is ‘dual’ to the classical notion of a random variable.

ex:stochastic-classical

Example 3.3 (Classical stochastic process) Consider the space of continuous paths (or trajectories) $\Omega := \mathcal{C}_0(\mathbb{R}_+; E)$ starting at 0 and the Wiener measure R . Let $X = (X_t)_{t \in \mathbb{T}}$ be the canonical stochastic process $X_t(\omega) = \omega(t)$, $t \in \mathbb{R}_+$. As before, we set

$$\mathfrak{X} := B_b(E), \quad \mathfrak{A} := L^\infty(\Omega, \mathsf{R}), \quad \mu(G) := \int_\Omega G \, d\mathsf{R}, \quad G \in \mathfrak{A}.$$

Then,

$$\mathfrak{X} \ni f \mapsto \mathfrak{z}_t(f) := f \circ X_t \in L^\infty(\Omega, \mathsf{R})$$

defines a noncommutative \mathfrak{X} -valued stochastic process on (\mathfrak{A}, μ) .

Definition 3.4 A stochastic process $\mathfrak{z}_t: \mathfrak{X} \rightarrow (\mathfrak{A}, \mu)$, $t \in \mathbb{T}$ admits a *time translation* if there are $*$ -homomorphisms $\alpha_t: \mathfrak{A} \rightarrow \mathfrak{A}$, $t \in \mathbb{T}$ such that

- (1) $\alpha_{t+s} = \alpha_t \circ \alpha_s$ for all $s, t \in \mathbb{T}$, and
- (2) $\mathfrak{z}_t = \alpha_t \circ \mathfrak{z}_0$ for all $t \in \mathbb{T}$.

Example 3.5 Consider the Wiener space (Ω, R) in Example 3.3. Then, R is stationary under the time-shift operator $\mathfrak{s}_t \omega = \omega(\cdot - t)$, i.e., $(\mathfrak{s}_t)_\# \mathsf{R} = \mathsf{R}$ for every $t \in \mathbb{T}$. Hence, the $*$ -homomorphism

$$\alpha_t(G) = G \circ \mathfrak{s}_t, \quad G \in \mathfrak{A}, \quad t \in \mathbb{T}$$

is a time translation for the corresponding process $(\mathfrak{z}_t)_{t \in \mathbb{T}}$ and leaves the state μ invariant. Indeed, property (1) holds trivially, while property (2) follows from

$$\alpha_t \circ \mathfrak{z}_0(f) = f \circ X_0 \circ \mathfrak{s}_t = f \circ X_t = \mathfrak{z}_t(f), \quad f \in \mathfrak{X}.$$

Finally, we conclude with

$$\mu(\alpha_t(G)) = \int_{\Omega} \alpha_t(G) dR = \int_{\Omega} G \circ \mathfrak{s}_t dR = \int_{\Omega} G d(\mathfrak{s}_t)_\# R = \int_{\Omega} G dR = \mu(G),$$

thus proving the invariance of μ .

For $h \in L^2(\mathbb{T})$ we consider the exponential martingale

$$Z_t = \exp \left(- \int_0^t h_r dB_r - \frac{1}{2} \int_0^t |h_r|^2 dr \right), \quad t \in \mathbb{T}.$$

It is easy to see that $(Z_t)_{t \in \mathbb{T}}$ satisfies

$$dZ_t = -h_t Z_t dB_t, \quad Z_0 = 1.$$

Setting $Q := Z_T P$, we claim that

$$W_t = B_t + \int_0^t h_r dr \quad \text{is a } Q\text{-martingale.}$$

To see this, we note that W is a Q -martingale if and only if $Z_T W$ is a P -martingale. Since

$$\begin{aligned} d(Z_t W_t) &= dZ_t W_t + Z_t dW_t + d[Z, W]_t \\ &= -h_t Z_t dB_t + Z_t dB_t = -(h_t - 1) Z_t dB_t, \end{aligned}$$

we conclude that $Z_T W$ is indeed a P -martingale.

$$dU_t = -(a^* a dt + i\sqrt{2} a dB_t) U_t$$

In the classical case, the notion of adaptedness plays an important role for stochastic processes. In this regard, we choose to simply consider the natural filtration induced by a stochastic process.

Definition 3.6 Let $\mathfrak{z}_t: \mathfrak{X} \rightarrow (\mathfrak{A}, \mu)$, $t \in \mathbb{T}$ be a stochastic process. For $\mathbb{I} \subset \mathbb{T}$, we denote by $\mathfrak{A}_{\mathbb{I}}$ the subalgebra of \mathfrak{A} generated by $\{\mathfrak{z}_t(x) : x \in \mathfrak{X}, t \in \mathbb{I}\}$. The filtration induced by $(\mathfrak{z}_t)_{t \in \mathbb{T}}$ is then given by $\mathfrak{F} = (\mathfrak{A}_{[0,t]})_{t \in \mathbb{T}}$.

3.2 From commutative to noncommutative

Consider the Hilbert space $\mathcal{H} = L^2_{\mathbb{C}}(\Omega, \mu)$ over the complete probability space $(\Omega, \mathcal{F}, \mu)$. Then any function $f \in L^\infty_{\mathbb{C}}(\Omega, \mu)$ gives rise to a multiplication operator $M_f \in \mathcal{B}(\mathcal{H})$:

$$M_f g = fg \in \mathcal{H} \quad \forall g \in \mathcal{H},$$

with $\|M_f\|_\infty = \|f\|_{L^\infty(\mu)}$. The collection of all such multiplication operators

$$\mathcal{A} := \{M_f : f \in L^\infty_{\mathbb{C}}(\Omega, \mu)\} \subset \mathcal{B}(\mathcal{H})$$

forms a commutative subalgebra of $\mathcal{B}(\mathcal{H})$.

In fact, this subalgebra is a *von Neumann algebra*:

Definition 3.7 (von Neumann algebra) A (unital) *von Neumann algebra* (or W^* -algebra) is a $*$ -subalgebra $\mathcal{A} \subset \mathcal{B}(\mathcal{H})$ that contains $I_{\mathcal{H}}$ and is closed in the *weak operator topology* (WOT), i.e.,

$$\text{WOT-lim } a_n = a \iff \langle f, a_n g \rangle_{\mathcal{H}} \rightarrow \langle f, ag \rangle_{\mathcal{H}} \quad \forall f, g \in \mathcal{H}.$$

Equivalently, is it a C^* -algebra with a predual \mathcal{A}_* that is a Banach space.

There are sub-families of \mathcal{A} that play a distinguished role, namely,

$$\begin{aligned} \mathcal{A}_+ &:= \{a \in \mathcal{A} : a \succeq 0\} && \text{(nonnegative operators)} \\ \mathcal{O} &:= \{a \in \mathcal{A} : a^* = a\} && \text{(self-adjoint operators)} \\ \mathcal{P} &:= \{a \in \mathcal{O} : a^2 = a\} && \text{(projection operators)} \end{aligned}$$

Here, $a \succeq 0$ if and only if $\langle f, af \rangle_{\mathcal{H}} \geq 0$ for all $f \in \mathcal{H}$.

On a von Neumann algebra \mathcal{A} , we can talk about special types of continuous linear functionals on \mathcal{A} , called states.

Definition 3.8 A *state* on a von Neumann algebra \mathcal{A} is a linear functional $\psi : \mathcal{A} \rightarrow \mathbb{C}$ that is positive and normalized, i.e., $\psi(a^*a) \geq 0$ for all $a \in \mathcal{A}$ and $\psi(I_{\mathcal{B}(\mathcal{H})}) = 1$.

A state ψ is said to be

faithful if $\psi(a^*a) = 0 \Leftrightarrow a = 0$,

tracial if $\psi(ab) = \psi(ba)$ for all $a, b \in \mathcal{A}$, and

normal if $\psi \in \mathcal{A}_*$, i.e., if it is an element of the predual \mathcal{A}_* .

Let us return to our previous example with $\mathcal{A} = \{M_f : f \in L^\infty_{\mathbb{C}}(\Omega, \mu)\}$. It is not difficult to see that \mathcal{A} is a W^* -algebra. Notice that since M_f is self-adjoint for $f \in L^\infty_{\mathbb{R}}(\Omega, \mu)$, the family of self-adjoint operators is given by $\mathcal{O} = \{M_f : f \in L^\infty_{\mathbb{R}}(\Omega, \mu)\}$ and the family of projection operators is given by $\mathcal{P} = \{M_f : f = \mathbf{1}_A, A \in \mathcal{F}\}$. Additionally,

For any probability measure $\nu \ll \mu$, we set

$$\psi_\nu(M_f) := \int_{\Omega} f \, d\nu, \quad f \in L^\infty_{\mathbb{C}}(\Omega, \mu),$$

we find that ψ is a linear functional that is positive and normalized, i.e., ψ is a state. Moreover, it is tracial. It is normal if $\omega := d\nu/d\mu \in L^1(\Omega, \mu)$ and faithful if $\omega > 0$.

Normal states play an essential role, serving as a counterpart to classical measures, as made explicit by the following proposition.

Proposition 3.9. *Let $\psi : \mathcal{A} \rightarrow \mathbb{C}$ be a state on a von Neumann algebra \mathcal{A} . The following are equivalent:*

- (i) ψ is a normal state.
- (ii) (σ -additivity) If $(a_n) \subset \mathcal{P}$ are mutually orthogonal projections, i.e., $a_n(\mathcal{H}) \perp a_m(\mathcal{H})$ for all $n \neq m$, and $a = \vee_n a_n$ being the projection on the smallest closed subspace containing $\cup_n a_n(\mathcal{H})$, then

$$\psi(a) = \sum_n \psi(a_n).$$

- (iii) (Continuity from above) For any increasing net $0 \preceq a_n \uparrow a$ in \mathcal{A}_+ $\Rightarrow \psi(a_n) \uparrow \psi(a)$.

Consider an arbitrary normal state ψ and set

$$\mu(A) := \psi(\mathbf{M}_{1_A}) \quad \text{for every } A \in \mathcal{F}.$$

In particular, we recover the probability \mathbb{P} by

$$\begin{aligned} \mathbb{P}(A) &= \psi(\mathbf{M}_{1_A}) \\ \psi(\mathbf{M}_{1_A} \mathbf{M}_{1_B}) &= \psi(\mathbf{M}_{1_{A \cap B}}) = \mathbb{P}(A \cap B) \\ \mathbb{P}(\cup_i A_i) &= \psi(\mathbf{M}_{1_{\cup_i A_i}}) = \psi(\sum_i \mathbf{M}_{1_{A_i}}) = \sum_i \psi(\mathbf{M}_{1_{A_i}}) = \sum_i \mathbb{P}(A_i) \end{aligned}$$

Let $(X_t)_{t \in \mathbb{T}}$ be a stochastic process on the Wiener space (Ω, \mathcal{R})

Consider a path measure P on $\Omega := \mathcal{D}(\mathbb{R}_+; E)$ and the canonical process $(X_t)_{t \geq 0}$ given by $X_t(\omega) = \omega(t)$, $\omega \in \Omega$, $t \in \mathbb{R}_+$. Suppose that for every $f \in B_b(E)$,

$$f(X_t) - f(X_0) - \int_0^t Lf(X_{r-}) dr \quad \text{is an } (\mathfrak{F}, \mathsf{P})\text{-martingale,} \quad (3.1)$$

where $\mathfrak{F} = (\mathfrak{F}_t)_{t \geq 0}$ is the canonical filtration $\mathfrak{F}_t = \sigma(X_s : s \leq t)$, and $L: B_b(E) \rightarrow B_b(E)$ is a bounded Markov generator such that $L^* \pi = 0$ for some stationary measure $\pi \in \mathcal{P}(E)$.

We now write (3.1) in the form above. Set $\mathfrak{X} := L^\infty(E, \pi)$ and $\mathfrak{A} := L^\infty(\Omega, \mathsf{P}_\pi)$, which are both commutative von Neumann algebras and where

$$\mathsf{P}_\pi(A) = \int_E \mathsf{P}(x + A) \pi(dx).$$

Now, define the stochastic process

$$\mathfrak{z}_t: \mathfrak{X} \rightarrow (\mathfrak{A}, \mu), \quad \mathfrak{z}_t(f) = f \circ X_t,$$

with the state

$$\mu(G) := \int_\Omega G(\omega) \mathsf{P}_\pi(d\omega), \quad G \in \mathfrak{A},$$

which makes (\mathfrak{A}, μ) a quantum probability space.

$$\mathsf{E}_t[\mathfrak{z}]$$

and the $*$ -homomorphism

$$\begin{aligned} \mathfrak{z}_t(f) - \mathfrak{z}_0(f) - \int_0^t \mathfrak{z}_{r-}(Lf) dr \\ \alpha_t(F) = F \circ \mathfrak{s}_t \\ \alpha_t \circ \mathfrak{z}_0(f) = \alpha_t(f \circ X_0 \otimes \mathbf{1}_{\mathfrak{Y}}) = f \circ X_0 \circ \mathfrak{s}_t(\omega) \\ \alpha_t: \mathcal{A} \times \Omega \rightarrow \mathcal{A}; \quad \alpha_t(M_f, \omega) = M_{f \circ X_t(\omega)}. \end{aligned}$$

The martingale identity (3.1) above can then be expressed as

$$\mathfrak{m}_t[M_f] := \alpha_t(M_f) - \alpha_0(M_f) - \int_0^t \alpha_{r-}(M_{Lf}) dr.$$

Setting $\mathcal{L}M_f := M_{Lf}$, the previous identity allows one to write (3.1) as

$$\mathfrak{m}_t[A] = \alpha_t(A) - \alpha_0(A) - \int_0^t \alpha_{r-}(\mathcal{L}A) dr, \quad A \in \mathcal{A},$$

thereby generalizing the martingale problem to a non-commutative setting.

$$\mathfrak{m}_t[A] - \mathfrak{m}_s[A] = \alpha_t(A) - \alpha_s(A) - \int_s^t \alpha_{r-}(\mathcal{L}A) dr$$

$$\mathcal{L}_a = a^*[\bullet, a] + [a^*, \bullet]a$$

Example 3.10 (Diffusion) Consider the formal example of the (possibly degenerate) diffusion process with the generator

$$Lf = \operatorname{div}(A\nabla f) - A\nabla V \cdot \nabla f, \quad f \in \mathcal{C}_c^\infty(\mathbb{R}^d),$$

i.e., it is the generator of the Itô diffusion

$$dX_t = -\sigma(X_t)\sigma^\top(X_t)\nabla V(X_t)dt + \sigma(X_t) \circ dB_t,$$

with $A = \sigma\sigma^\top \in \mathbb{R}^{d \times d}$, $\sigma \in \mathbb{R}^{d \times m}$

$$a_i = \sum_{j=1}^d \sigma_{ji} \partial_j, \quad a_i^* = -\sum_{j=1}^d \partial_j(\sigma_{ji} \bullet), \quad i = 1, \dots, m.$$

It is not difficult to see that $[a_i^*, M_f] = -[a_i, M_f]$. Moreover, for any $g \in C_c^\infty(\mathbb{R}^d)$,

$$\begin{aligned} a_i^*[M_f, a_i]g &= \sum_{j=1}^d \partial_j(\sigma_{ji}g \sum_{k=1}^d \sigma_{ki} \partial_k f) \\ &= \sum_{j=1}^d g \partial_j(\sigma_{ji} \sum_{k=1}^d \sigma_{ki} \partial_k f) + \sum_{j,k=1}^d \sigma_{ji} \sigma_{ki} \partial_j g \partial_k f \\ &= \sum_{j=1}^d g \partial_j(\sigma_{ji} \sum_{k=1}^d \sigma_{ki} \partial_k f) - [M_f, a_i]a_i g. \end{aligned}$$

Consequently, we find that

$$\sum_{i=1}^m \mathcal{L}_{a_i}(M_f) = \sum_{i=1}^m (a_i^*[M_f, a_i] + [a_i^*, M_f]a_i) = M_{\operatorname{div}(A\nabla f)}$$

Moreover, setting

$$H_i = \sum_{j,k=1}^d \sigma_{ji} \sigma_{ki} \partial_j V \partial_k, \quad i = 1, \dots, m,$$

we have that

$$\begin{aligned} \sum_{i=1}^m [H_i, M_f] &= M_{\sigma\sigma^\top \nabla V \cdot \nabla f} \\ &= [a_i, M_V][a_i, \bullet] \end{aligned}$$

Therefore, we obtain

$$\mathcal{L}M_f = \sum_{i=1}^m ([H_i, M_f] + \mathcal{L}_{a_i} M_f) = M_{Lf} \quad \forall f \in \mathcal{C}_c^\infty(\mathbb{R}^d).$$

However, we know that the diffusion process constitutes a gradient flow with driving energy $\mathcal{F} = \operatorname{Ent}(\bullet|\pi)$, where $\pi = e^{-V} \operatorname{Leb}$ is the invariant measure. Formally, one defines a state on $\mathcal{C}_c^\infty(\mathbb{R}^d)$ by

$$\tau(M_f) := \int_{\mathbb{R}^d} f \, d\pi, \quad f \in \mathcal{C}_c^\infty(\mathbb{R}^d).$$

$$\begin{aligned}\tau_\pi(M_fM_g) &= \tau(M_fM_gM_\pi)=\tau(M_{fg\pi})=\tau_\pi(M_gM_f)\\ \Delta_\pi A &= M_\pi A M_\pi^{-1}\end{aligned}$$

$$\begin{aligned}(\Delta_\pi V_j)(g)&=M_\pi V_jM_\pi^{-1}g=\pi\sum_{j=1}^d\sigma_{ji}\partial_j(\pi^{-1}g)-\sigma_{ji}\partial_jV(\pi^{-1}g)\\&=\sum_{j=1}^d\sigma_{ji}[g\partial_jV+\partial_jg]-\sigma_{ji}\partial_jVg=\\V_j&=\sum_{j=1}^d\sigma_{ji}\partial_j-M_{\sigma_{ji}\partial_jV}\\ \tau_\sigma(M_g\mathscr{L}M_f)&=\hat{\sigma}(M_{gLf})=\int_{\mathbb{R}^d}gLf\,\mathrm{d}\sigma=\end{aligned}$$

$$\begin{aligned}\mathscr{L}M_{e^{-V}}g&=\big(\text{div}(A\nabla e^{-V})+A\nabla V\cdot\nabla e^{-V}\big)g\\&=\big(-\text{div}(A\nabla Ve^{-V})-A\nabla V\cdot\nabla Ve^{-V}\big)g\end{aligned}$$

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$$\begin{aligned}
[a_i, M_f]g &= \sum_{j=1}^d \sigma_{ji} \partial_j(fg) - \sum_{j=1}^d f \sigma_{ji} \partial_j g = \sum_{j=1}^d \sigma_{ji} g \partial_j f = M_{a_i(f)} g \\
[a_i^*, M_f]g &= - \sum_{j=1}^d \partial_j(\sigma_{ji} fg) + \sum_{j=1}^d f \partial_j(\sigma_{ji} g) = - \sum_{j=1}^d \sigma_{ji} \partial_j fg = -M_{a_i(f)} g = -[a_i^*, M_f]g
\end{aligned}$$

Consequently, we find that

$$\begin{aligned}
\mathcal{L}_{a_i} M_f &= \\
[a_i, M_f]a_i g &= \sum_{j,k=1}^d \sigma_{ji} \sigma_{ki} \partial_j f \partial_k g \\
[a_i, M_V] &= \sum_{j=1}^d \sigma_{ji} \partial_j V \\
[a_i, [a_i, M_f]] &= [a_i, M_{a_i(f)}] = M_{a_i(a_i(f))},
\end{aligned}$$

where we used the fact that $[a_i, M_f] = M_{a_i(f)}$ for every $i = 1, \dots, d$.

$$\begin{aligned}
\sum_{i=1}^m a_i \circ a_i(f) &= \sum_{i=1}^m \sum_{j=1}^d \sigma_{ji} \partial_j \left(\sum_{k=1}^d \sigma_{ki} \partial_k f \right) = \sum_{i=1}^m \sum_{j,k=1}^d \sigma_{ji} \partial_j (\sigma_{ki} \partial_k f) \\
&= \sum_{i=1}^m \sum_{j,k=1}^d \sigma_{ji} \partial_j (\sigma_{ki} \partial_k f)
\end{aligned}$$

$$\begin{aligned}
\text{div}(\sigma \sigma^\top \nabla f) &= \sum_{j=1}^d \sum_{i=1}^m \sum_{k=1}^d \partial_j (\sigma_{ji} \sigma_{ki} \partial_k f) \\
&= \sum_{j=1}^d \sum_{i=1}^m \sum_{k=1}^d \partial_j (\sigma_{ji} \sigma_{ki} \partial_k f)
\end{aligned}$$

$$\begin{aligned}
\langle \psi | M_{Lf} | \psi \rangle &= \int (Lf)(x) |\psi(x)|^2 \pi(dx) \\
&= - \int \sigma^2 \partial_x f \partial_x (\psi^2 \pi) dx - \int \sigma^2 \partial_x f \partial_x V \psi^2 d\pi \\
&= -2 \int \sigma^2 \partial_x f \psi \partial_x \psi d\pi \\
&= -2 \int
\end{aligned}$$

$$a\psi = \sigma \partial_x \psi$$

$$\begin{aligned}
\langle \varphi | a\psi \rangle &= \int \varphi \sigma \partial_x \psi \, d\pi \\
&= - \int \partial_x(\sigma \varphi \pi) \psi(x) \, dx \\
&= - \int \partial_x(\sigma \varphi) \psi \, d\pi + \int \varphi \psi \sigma \partial_x V \, d\pi \\
&= \int [-\partial_x(\sigma \varphi) + \varphi \sigma \partial_x V] \psi \, d\pi = \langle a^* \varphi | \psi \rangle
\end{aligned}$$

$$\begin{aligned}
(a^* M_f a \psi)(x) &= -\partial_x(f \sigma^2 \partial_x \psi) + f \sigma^2 \partial_x \psi \partial_x V \\
(a^* a M_f \psi)(x) &= -\partial_x(\sigma^2 \partial_x(f \psi)) + \sigma^2 \partial_x(f \psi) \partial_x V \\
(M_f a^* a \psi)(x) &= -f \partial(\sigma^2 \partial_x \psi) + \sigma^2 f \partial_x \psi \partial_x V
\end{aligned}$$

$$\int \psi(x) (a^* M_f a \psi)(x) \, d\pi = \int f |\sigma \partial_x \psi|^2 \, d\pi$$

3.3 Stochastic processes on matrix algebras

In this section, we consider $\mathfrak{X} = \text{Mat}(\mathbb{C}, n)$. Let $a \in \mathcal{O}(\mathfrak{X})$ be an observable on \mathfrak{X} and set

$$\mathcal{L}x = a^*[x, a] + [a^*, x]a, \quad x \in \mathfrak{X}.$$

We then consider the solution of the

$$d\mathcal{U}_t = (-a^* a \, dt + i\sqrt{2}a \, dB_t) \mathcal{U}_t,$$

which, due to $a = a^*$, can be explicitly expressed as

$$\mathcal{U}_t(x, e(f)) := \exp(i\sqrt{2}a B_t(\omega))x, \quad t \in \mathbb{T},$$

which is a $*$ -automorphism on \mathfrak{X} , which extends to an isometry

$$\mathcal{U}_t: \mathfrak{X} \otimes \mathcal{A}$$

$$B \in \mathcal{A}$$

$$\mathfrak{z}_t(x) - \mathfrak{z}_0(x) - \int_0^t \mathfrak{z}_r(\mathcal{L}x) \, dr$$

3.4 Stochastic processes on the unitary group

3.4.1 The unitary group and its Lie structure

The unitary group is defined as

$$U(n) = \{U \in \mathbb{C}^{n \times n} : U^* U = I\}.$$

It is a compact Lie group with Lie algebra (tangent space at identity)

$$\mathcal{U}(n) = \mathsf{T}_I U(n) = \{\mathfrak{a} \in \mathbb{C}^{n \times n} : \mathfrak{a}^* = -\mathfrak{a}\},$$

i.e., the space of skew-Hermitian matrices. A vector field $\mathfrak{a} : U(n) \rightarrow \mathcal{U}(n)$ is *left-invariant* if $\mathfrak{a}(U) = U\mathfrak{a}(I) \in \mathcal{U}(n)$ for all $U \in U(n)$.

The Lie algebra $\mathcal{U}(n)$ can be equipped with a real inner product given by the Hilbert-Schmidt inner product

$$\langle \mathfrak{a}, \mathfrak{b} \rangle_{\mathcal{U}(n)} = -\text{tr}(\mathfrak{a}^* \mathfrak{b}) = \Re \text{tr}(\mathfrak{a} \mathfrak{b}^*),$$

which is positive-definite on $\mathcal{U}(n)$. The associated norm is then

$$|\mathfrak{a}|_{\mathcal{U}(n)}^2 = \langle \mathfrak{a}, \mathfrak{a} \rangle_{\mathcal{U}(n)} = \text{Tr}(\mathfrak{a} \mathfrak{a}^*).$$

3.4.2 Unitary-valued Brownian motion (rank- r noise)

Let $\mathfrak{a}_1, \dots, \mathfrak{a}_r \in \mathcal{U}(n)$ be orthonormal under the inner product $\langle \cdot, \cdot \rangle_{\mathcal{U}(n)}$, where $1 \leq r \leq n^2$. Further, let $\beta_t^1, \dots, \beta_t^r$ be independent standard real Brownian motions.

Define the $\mathcal{U}(n)$ -valued (possibly degenerate) Brownian driver

$$\mathfrak{w}_t = \sum_j \mathfrak{a}_j \beta_t^j.$$

If $r = n^2$, the driver is elliptic (non-degenerate). If $r < n^2$, the covariance has rank r and the process is *degenerate* (hypoelliptic), exploring only the connected subgroup

$$H = \exp(\mathcal{H}) \subset U(n), \quad \mathcal{H} = \text{Lie}\{\mathfrak{a}_1, \dots, \mathfrak{a}_r\} \subset \mathcal{U}(n).$$

The intrinsic $U(n)$ -valued diffusion process solves the Stratonovich SDE

$$dU_t = U_t \circ d\mathfrak{w}_t, \quad U_0 = I,$$

or in components

$$dU_t = \sum_j U_t \mathfrak{a}_j \circ d\beta_t^j.$$

The matrix-valued quadratic variation of \mathfrak{w} is given by

$$d[\mathfrak{w}] = \sum_j \mathfrak{a}_j \otimes \mathfrak{a}_j dt.$$

In Itô form, the equivalent SDE reads

$$dU_t = U_t (d\mathfrak{w}_t + \mathfrak{l} dt), \quad \mathfrak{l} := \frac{1}{2} \sum_j \mathfrak{a}_j^2,$$

where \mathfrak{l} is the Laplace-Beltrami generator associated with the left-invariant connection.

In the following, we denote its law on path space $C([0, T]; U(n))$ by P .

3.4.3 Girsanov transform on $U(n)$ (Karandikar type)

Let $\mathbf{u} \in L^2((0, T); \mathcal{U}(n))$ be adapted and square-integrable.

Define the drifted group motion by

$$dU_t^{\mathbf{u}} = U_t^{\mathbf{u}} \circ (d\mathbf{w}_t + \mathbf{u}_t dt), \quad U_0^{\mathbf{u}} = I.$$

Assuming a Novikov condition in the active noise directions,

$$\mathbb{E} \left[\exp \left(\frac{1}{2} \int_0^T |\mathbf{u}_s|^2 ds \right) \right] < +\infty,$$

the process defined below is a true martingale.

Decompose the drift against the active algebra frame:

$$\mathbf{u}_t = \sum_j \mathbf{a}_j v_t^j, \quad v = (v^1, \dots, v^r) \in L^2((0, T); \mathbb{R}^r).$$

Define the exponential Cameron-Martin density process

$$\begin{aligned} M_T &= \exp \left(\int_0^T \langle \mathbf{u}_t, d\mathbf{w}_t \rangle - \frac{1}{2} \sum_j \int_0^T \langle \mathbf{u}_t, \mathbf{a}_j \otimes \mathbf{a}_j \mathbf{u}_t \rangle dt \right) \\ &= \exp \left(\sum_j \int_0^T v_t^j d\beta_t^j - \frac{1}{2} \sum_j \int_0^T |v_t^j|^2 dt \right). \end{aligned}$$

This expression depends only on the directions with nonzero quadratic variation. It is the standard Cameron-Martin shift density on \mathbb{R}^r lifted to the group via left multiplication.

We then define a new path measure on the same probability space by

$$\mathbf{Q} = M_T \mathbf{P}.$$

Under \mathbf{Q} :

(i) The shifted drivers

$$\beta_t^{j, \mathbf{Q}} := \beta_t^j - \int_0^t v_s^j ds$$

are Brownian motions again on \mathbb{R} .

(ii) The algebra noise

$$d\mathbf{w}_t^{\mathbf{Q}} := d\mathbf{w}_t - \mathbf{u}_s dt = \sum_j \mathbf{a}_j d\beta_t^{j, \mathbf{Q}}$$

is a $\mathcal{U}(n)$ -valued Brownian motion again with same covariance

$$d[\mathbf{w}^{\mathbf{Q}}] = d[\mathbf{w}] = \sum_j \mathbf{a}_j \otimes \mathbf{a}_j dt.$$

(iii) The group paths satisfy the same *driftless* intrinsic SDE

$$dU_t^{\mathbf{u}} = U_t^{\mathbf{u}} \circ dW_t^{\mathbf{Q}} = \sum_j U_t^{\mathbf{u}} \mathbf{a}_j \circ d\beta_t^{j, \mathbf{Q}},$$

so $U_t^{\mathbf{u}}$ is an intrinsic Brownian motion on the same symmetry group H .

Hence, M_T is the Radon-Nikodym derivative between the laws of the drifted and driftless group diffusions on path space.

3.4.4 Induced Girsanov on pure states

Rank-1 projections evolve by the conjugation map

$$P_t = U_t P_0 U_t^*, \quad P_0 = \psi_0 \psi_0^*.$$

This gives a valid reference probability on paths of $C([0, T]; \mathcal{M})$, where

$$\mathcal{M} = \{V P_0 V^* : V \in H\} \subset \mathcal{P}_1(n) \simeq \mathbb{CP}^{n-1}$$

is the homogeneous orbit that the degenerate noise can reach.

Linearising conjugation and applying the Stratonovich chain rule yields the SDE

$$dP_t^u = [\circ dW_t^Q, P_t^u] = \sum_j [\mathfrak{a}_j, P_t^u] \circ d\beta_t^{j,Q},$$

so the pure-state path law remains Brownian on the same orbit manifold \mathcal{M} .

If a drift u has components outside $\text{span}\{\mathfrak{a}_j\}$, only the projection of the drift onto that span can be removed. Directions with zero quadratic variation cannot absorb drifts.

$$\int F(U_t P_0 U_t^*) \mathsf{P}(dU) =$$

3.5 Dilation of semigroups (Sz.Nagy)

3.6 Stochastic dilations: How to make a heat bath

3.7 Interpretation of LDPs on unitary evolutions

The task of this subsection is to reinterpret the large deviations of empirical measures in a functional analytic framework.

Consider a family of iid E -valued random variables $(X_i)_{i \in \mathbb{N}}$ with $\sigma = \text{Law } X_1$. Setting $\Omega = E^{\mathbb{N}}$ and $\mathsf{R} = \sigma^{\otimes \mathbb{N}}$, we see that $\text{Law}(X_i)_{i \in \mathbb{N}} = \mathsf{R}$. Moreover, we can disregard the initial probability space for X_i and consider instead the canonical random variable on Ω .

For every $f \in \mathfrak{X} := \mathcal{C}_b(E)$ and $n \in \mathbb{N}$, we define the noncommutative random variable

$$\mathfrak{X} \ni f \mapsto \mathfrak{z}^n(f) = \frac{1}{n} \sum_{i=1}^n f(X_i) \in \mathfrak{A} := L^\infty(\Omega, \mathsf{R}),$$

where \mathfrak{A} is equipped with the state $\mu(G) = \mathsf{E}_{\mathsf{R}}[G]$, $G \in \mathfrak{A}$.

$$\mu(\mathfrak{z}^n(f)) = \frac{1}{n} \sum_{i=1}^n \mathsf{E}_{\mathsf{R}}[f(X_i)] = \frac{1}{n} \sum_{i=1}^n \int_E f \, d\sigma = \int_E f \, d\sigma = \mu(\mathfrak{z}_1(f)).$$

$$\begin{aligned} \mathsf{R}(\sup_f \{\mu^n(f) - \nu(f)\} > \varepsilon) &= \mathsf{R}(n\mu^n(f) > n(\nu(f) + \varepsilon)) \\ &= \mathsf{R}(e^{n\mu^n(f)} > e^{n(\nu(f) + \varepsilon)}) \leq e^{-n(\nu(f) + \varepsilon)} \int_{\Omega} e^{n\mu^n(f)} \, d\mathsf{R} \end{aligned}$$

$$\begin{aligned} \mathfrak{z} &\colon \mathfrak{X} \rightarrow (\mathfrak{X} \otimes \mathfrak{C}, \text{tr} \otimes \mu) \\ \mathfrak{A} &= \mathfrak{X} \otimes \mathfrak{C}^{\otimes \mathbb{Z}}, \quad \varphi = \text{tr} \otimes \mu^{\otimes \mathbb{Z}} \end{aligned}$$

Let X_i

$$\mathfrak{z}_i(x) = x \otimes c_i, \quad c_i = \cdots 1 \otimes c$$

$$\mathfrak{z}^n = \frac{1}{n} \sum_{i=1}^n \mathfrak{z}_i, \quad \mathfrak{z}_i = \cdots \otimes \mathbb{1} \otimes \mathfrak{z} \otimes \mathbb{1} \otimes \cdots$$

$$\varphi(\mathfrak{z}^n(x)) = \frac{1}{n} \sum_{i=1}^n \varphi(\mathfrak{z}_i(x)) = \varphi(\mathfrak{z}_1(x)) = \mu(\mathfrak{z}_1(x))$$

$$\mathfrak{z}_i(x) = \mathbb{1} \otimes \cdots \otimes \mathbb{1} \otimes \mathfrak{z}(x) \otimes \mathbb{1} \otimes \cdots \otimes \mathbb{1}$$

$$\mathfrak{z}^n(x) \in \mathfrak{A} \cong \mathfrak{X} \otimes \mathfrak{Y}$$

$$\mu = \text{tr} \otimes \mu_{\mathfrak{C}}^{\otimes \mathbb{N}}$$

$$\mu(\mathbf{1}_A(\mathfrak{z}^n(x) - \mathfrak{y}_1(x))) = \mu_{\mathfrak{C}}^{\otimes \mathbb{N}}(\text{tr}[\mathbf{1}_A(\mathfrak{z}^n(x) - \mathfrak{y}(x) \otimes \mathbb{1}_{\mathfrak{C}^{\otimes \mathbb{N}}})])$$

$$\sup_f \{\mu^n(f) - \nu(f)\} > \varepsilon \Rightarrow \mu^n(f^\delta) - \nu(f^\delta) > \varepsilon - \delta$$

$$\begin{aligned} \mathsf{R}(\mu^n(f) - \nu(f) > \varepsilon) &= \mathsf{R}(e^{n\mu^n(f^\delta)} > e^{n(\nu(f^\delta)+\varepsilon-\delta)}) \\ &\leq e^{-n(\nu(f^\delta)+\varepsilon-\delta)} \mathsf{E}_{\mathsf{R}}[e^{n\mu^n(f^\delta)}] \end{aligned}$$

$$\begin{aligned} -\frac{1}{n} \log \mathsf{R}(\mu^n(f) - \mu(f) > \varepsilon) &\geq \mu(f) + \varepsilon - \frac{1}{n} \log \int_{\Omega} e^{n\mu^n(f)} d\mathsf{R} \\ &= \mu(f) + \varepsilon - \log \int_E e^f d\mu \\ -\frac{1}{n} \log \mathsf{R}(\mu^n(f) - \mu(f) > \varepsilon) &\geq \end{aligned}$$

Taking the sup

$$\begin{aligned} \int_{\Omega} e^{n\mu^n(f)} d\mathsf{R} &= \int_{\Omega} \prod_{i=1}^n e^{f(X_i)} d\mathsf{R} \\ &= \int_{E^N} \prod_{i=1}^n e^{f(x_i)} \mu(dx_1) \cdots \mu(dx_n) = \left(\int_E e^f(x) \mu(dx) \right)^n \end{aligned}$$

3.8 LDP via Gärtner-Ellis revisited

Let $\Omega = \mathcal{C}(I; \mathfrak{h})$ be the space of \mathfrak{h} -valued continuous path and consider the canonical process $(X_t)_{t \in I}$ defined by $X_t(\omega) := \omega(t)$ for all $t \in I$. Further, let $\mathcal{F} = \cup_{t \geq 0} \mathcal{F}_t$, where the filtration $\mathfrak{F} := \{\mathcal{F}_t := \sigma(X_s : 0 \leq s \leq t)\}_{t \in I}$ is generated by $(X_t)_{t \in I}$. Throughout, we consider a reference path measure $\mathsf{R} \in \mathcal{P}(\Omega)$ such that X admits the R -semimartingale decomposition

$$X = X_0 + B^{\mathsf{R}} + M^{\mathsf{R}} \quad \mathsf{R}\text{-almost surely},$$

where B^{R} is an adapted process with absolutely continuous sample paths R -almost surely, and M^{R} is an R -martingale. For simplicity of presentation, we assume that the quadratic variation of M^{R} is an $\mathcal{S}_+(\mathfrak{X})$ -valued random measure on I with Lebesgue dentisy, i.e.,

$$\frac{d[M^{\mathsf{R}}, M^{\mathsf{R}}]}{d\lambda}(t) = \mathbf{a}_t(X_t) \in \mathcal{S}_+(\mathfrak{h}) \quad \mathsf{R}\text{-almost surely},$$

and also that $t \mapsto \mathbf{a}_t(X_t)$ is an \mathfrak{F} -adapted process.

and has the quadratic variation $\langle M^{\mathsf{R}}[\varphi] \rangle$ for every smooth function φ .

Let X^1, X^2, \dots be independent copies of

$$(X^1, \dots, X^n)$$

$$\begin{aligned} \frac{d\mathsf{Q}^n}{d\mathsf{P}^{\otimes n}} &= \exp \left(\sum_{i=1}^n \int_I \langle \beta_t^i(X_t), dM_t^{\mathsf{R}} \rangle - \frac{1}{2} \sum_{i=1}^n \int_I \langle \beta_t^i(X_t), \mathbf{a}_t(X_t) \beta_t^i(X_t) \rangle dt \right) =: Z^{(n)} \\ Y_t^{(n)} &:= \sum_{i=1}^n \int_0^t \langle \beta_s^i(X_s), dM_s^{\mathsf{R}} \rangle - \frac{1}{2} \sum_{i=1}^n \int_0^t \langle \beta_s^i(X_s), \mathbf{a}_s(X_s) \beta_s^i(X_s) \rangle ds \end{aligned}$$

$$dZ_t^{(n)} = de^{Y_t^{(n)}} = Z_t^{(n)} \left(dY_t^{(n)} + \frac{1}{2} d[Y^{(n)}, Y^{(n)}]_t \right) = \sum_{i=1}^n Z_t^{(n)} \langle \beta_t^i(X_t), dM_t^R \rangle$$

Define $\rho^n := \frac{1}{n} \sum_{i=1}^n X^i$. Then,

$$d\rho^n = \frac{1}{n} \sum_{i=1}^n dX_t^i = \frac{1}{n} \sum_{i=1}^n b_t^R(X_t^i) dt + \frac{1}{n} \sum_{i=1}^n dM_t^{R,i}$$

$$\begin{aligned} \mathbb{P}^{\otimes n} \left(\frac{1}{n} \sum_{i=1}^n X^i \in A \right) &= \int_A e^{-\int_I \langle \beta_t^i(X_t), dM_t^R \rangle + \frac{1}{2} \int_I \langle \beta_t^i(X_t), \mathbf{a}_t(X_t) \beta_t^i(X_t) \rangle dt} d\mathbb{Q}^{(n)} \\ &= \int_A e^{-\int_I \langle \beta_t^i(X_t), dM_t^R \rangle + \frac{1}{2} \int_I \langle \beta_t^i(X_t), \mathbf{a}_t(X_t) \beta_t^i(X_t) \rangle dt} d\mathbb{Q}^{(n)} \end{aligned}$$

$$A := \{\rho : \partial_t \rho + \overline{\operatorname{div}} j = L\rho\}$$

$$D_{\beta^i} \exp \left(\int_I \langle \beta_t^i(X_t), dM_t^R \rangle + \frac{1}{2} \int_I \langle \beta_t^i(X_t), \mathbf{a}_t(X_t) \beta_t^i(X_t) \rangle dt \right) [h^i]$$

$$\int_I \langle \beta_t, j_t \rangle dt - \log \int_A Z^{(n)} d\mathbb{P}^{\otimes n}$$

3.8.1 Example

We consider the Hilbert space $\mathfrak{h} = \mathbb{M}$ equipped with the scalar product $\langle A, B \rangle = \operatorname{tr}[A^* B]$. Consider the stochastic process

$$P_t = P_0 + \int_0^t LP_s ds + \int_0^t \overline{\nabla} P_s dW_s,$$

where $\overline{\nabla} P = [i\sigma_X, P]$ and $LP = (1/2)[i\sigma_X, [i\sigma_X, P]]$. In other words,

$$B_t^R = \int_0^t LP_s ds, \quad M_t^R = \int_0^t \overline{\nabla} P_s dW_s,$$

where M^R has the quadratic variation

$$\mathbf{a}(P) = \langle (\overline{\nabla} P)^*, \bullet \rangle \overline{\nabla} P \quad \text{for every } P \in \mathfrak{h}.$$

Let P^1, P^2, \dots be independent copies and define $\rho^n := \frac{1}{n} \sum_{i=1}^n P^i$. Then,

$$d\rho_t^n = L\rho_t^n dt + \frac{1}{n} \sum_{i=1}^n \overline{\nabla} P_t^i dW_t^i = L\rho_t^n dt + \frac{1}{n} dM_t^{R,n}.$$

$$\begin{aligned} \frac{1}{n} \sum_{i=1}^n \bar{\nabla} P_t^i dW_t^i &= \frac{1}{n} \sum_{i=1}^n \left(\bar{\nabla} P_t^i dW_t^i - \mathbf{a}(P_t^i) H_t dt \right) + \frac{1}{n} \sum_{i=1}^n \mathbf{a}(P_t^i) H_t dt \\ d\rho^n &= L\rho^n dt + \frac{1}{n} \sum_{i=1}^n \mathbf{a}(P_t^i) H_t dt + \frac{1}{n} dM_t^{\mathbf{Q},n} \end{aligned}$$

For each $n \in \mathbb{N}$, we set

$$\frac{d\mathbf{Q}^n}{d\mathbf{R}^{\otimes n}} = \exp \left(\sum_{i=1}^n \int_I \langle \bar{\nabla} P_t^i, H_t \rangle dW_t^i - \frac{1}{2} \sum_{i=1}^n \int_I \langle H_t, \mathbf{a}(P_t^i) H_t \rangle dt \right) =: Z^{(n)}.$$

Then, for curves η satisfying $\partial_t \eta + \overline{\operatorname{div}} J^\eta = L\eta$ for some J^η , we find that

$$J^\eta = \sum \langle \sigma_j, J^\eta \rangle \sigma_j$$

Suppose

$$-\overline{\operatorname{div}} = \sum_{j \in J} \alpha_j [i\sigma_j, \bullet],$$

then

$$-\overline{\operatorname{div}} J^\eta = \sum_{j,k} \alpha_j \langle \sigma_k, J^\eta \rangle [i\sigma_j, \sigma_k] = \sum_{j,k \neq j} \alpha_j \langle \sigma_k, J^\eta \rangle [i\sigma_j, \sigma_k] \in \operatorname{span} \{ [i\sigma_j, \sigma_k] : j \in J, k \neq j \}$$

$$\begin{aligned} \mathbf{R}^{\otimes n}(\rho^n \in A) &= \int_A \exp \left(- \sum_{i=1}^n \int_I \langle \bar{\nabla} P_t^i, H_t \rangle dW_t^i + \frac{1}{2} \sum_{i=1}^n \int_I \langle H_t, \mathbf{a}(P_t^i) H_t \rangle dt \right) d\mathbf{Q}^{(n)} \\ &= \int_A \exp \left(- \int_I \langle dM_t^{\mathbf{R},n}, H_t \rangle + \frac{1}{2} \sum_{i=1}^n \int_I \langle H_t, \mathbf{a}(P_t^i) H_t \rangle dt \right) d\mathbf{Q}^{(n)} \\ &= \int_A \exp \left(- \int_I \langle dM_t^{\mathbf{Q},n}, H_t \rangle - \frac{1}{2} \sum_{i=1}^n \int_I \langle H_t, \mathbf{a}(P_t^i) H_t \rangle dt \right) d\mathbf{Q}^{(n)} \end{aligned}$$

$$H^\eta \in \operatorname{argmin}_H \left\{ \mathbf{E}_{\mathbf{Q}} \left[\frac{1}{2} \int_I \langle H_t, \mathbf{a}(P_t) H_t \rangle dt \right] - \int_I \langle H_t, \overline{\operatorname{div}} J_t^\eta \rangle dt \right\}$$

$$\mathbf{E}_{\mathbf{Q}} \left[\int_I (\operatorname{tr}[P_t \bar{\nabla} H_t])^2 dt \right] = \int_I \mathbf{E}_{\mathbf{Q}} \left[(\operatorname{tr}[P_t \bar{\nabla} H_t])^2 \right] dt$$

$$\operatorname{tr}[\rho_t^{\mathbf{Q}} \bar{\nabla} H_t] = \sum \lambda_i(t) \langle v_i(t), \bar{\nabla} H_t v_i(t) \rangle \geq \gamma \sum \lambda_i(t)$$

$$\mathbf{E}_{\mathbf{R}} \left[\int_I \langle G_t, \mathbf{a}(P_t) H_t \rangle dt \right] = \int_I \langle G_t, \overline{\operatorname{div}} J_t \rangle dt$$

$$\int_I \langle H_t, \overline{\operatorname{div}} J_t \rangle dt - \log \mathbf{E}_{\mathbf{R}} \left[\exp \left(\int_I \langle H_t, dM_t^{\mathbf{R},n} \rangle - \frac{1}{2} \sum_{i=1}^n \int_I \langle H_t, \mathbf{a}(P_t^i) H_t \rangle dt \right) \right]$$

$$\int_I \langle G_t, \overline{\operatorname{div}} J_t \rangle dt = -\mathsf{E}_{\mathsf{R}^{\otimes n}} \left[\left(\int_I \langle G_t, dM_t^{\mathsf{R},n} \rangle - \sum_{i=1}^n \int_I \langle G_t, \mathsf{a}(P_t^i) H_t \rangle dt \right) \exp \left(\int_I \langle H_t, dM_t^{\mathsf{R},n} \rangle - \frac{1}{2} \sum_{i=1}^n \int_I \langle H_t, \mathsf{a}(P_t^i) H_t \rangle dt \right) \right]$$

Let $Y_t^{(n)} = \sum_{i=1}^n \int_0^t \langle \overline{\nabla} H_s, \overline{\nabla} P_s^i \rangle dW_s^i$ and $Z_t^{(n)} := \exp(Y_t^{(n)})$. Then,

$$\begin{aligned} dZ_t^{(n)} &= Z_t^{(n)} \left(\sum_{i=1}^n \langle \overline{\nabla} H_t, \overline{\nabla} P_t^i \rangle dW_t^i + \frac{1}{2} \sum_{i=1}^n \langle \overline{\nabla} H_t, \mathsf{a}(P_t) \overline{\nabla} H_t \rangle dt \right) \\ Z_t^{(n)} &= 1 + \sum_{i=1}^n \int_0^t Z_s^{(n)} \langle \overline{\nabla} H_s, \overline{\nabla} P_s^i \rangle dW_s^i + \frac{1}{2} \sum_{i=1}^n \int_0^t Z_s^{(n)} \langle \overline{\nabla} H_s, \mathsf{a}(P_s) \overline{\nabla} H_s \rangle ds \end{aligned}$$

Let $\varrho \in \mathcal{S}(\mathcal{H})$ be state and $(\varrho^n)_{n \in \mathbb{N}} \subset \mathcal{S}(\mathcal{H})$ be a sequence of states. We say that $\varrho^n \rightarrow \varrho$ setwise if

$$\varrho^n(E) \rightarrow \varrho(E) \quad \text{for all projections } E \in \mathcal{P}(\mathcal{H}).$$

Further, let $(\mathfrak{z}_t)_{t \in \mathbb{T}}$ be a stochastic process in the sense of Definition 3.1 with

$$\mathfrak{z}_t: \mathfrak{X} \rightarrow (\mathfrak{X} \otimes \mathfrak{C}, \varphi), \quad \varphi = \varrho \otimes \mu.$$

$$\mathfrak{z}_t(x) = \alpha_t^* x \alpha_t$$

and $\mathfrak{z}_0(x) = x \otimes 1$, $x \in \mathfrak{X}$. By construction, we have that

$$\varphi(\mathfrak{z}_0(x)) = \varphi(x \otimes 1) = \varrho(x)\mu(1) = \varrho(x) \quad \forall x \in \mathfrak{X}.$$

We suppose that

$$\begin{aligned} \mathcal{S}(\mathfrak{X} \otimes \mathfrak{C}^{\otimes n}) &\rightarrow \mathcal{S}(\mathfrak{X} \otimes \mathfrak{C}) \\ \mu_\varphi &= \frac{1}{n} \sum_{i=1}^n \varphi_i = \frac{1}{n} \sum_{i=1}^n \rho \otimes \varphi \\ \mathbf{x}_i : \mathfrak{C}^{\otimes n} &\rightarrow \mathfrak{C}; \quad \mathbf{x}_i(\omega_1, \dots, \omega_n) = \omega_i \\ \delta_{\mathbf{x}_i}(E) \end{aligned}$$

3.9 Groups to algebras

Let \mathbb{T} be the 1-dimensional torus. Further, let \mathcal{H} be a Hilbert space and $\mathsf{h} \in \mathcal{O}(\mathcal{H})$ be an observable on \mathcal{H} . Consider the unitary $\mathfrak{u}: \mathbb{T} \rightarrow \mathcal{U}(\mathcal{H})$ such that $\mathfrak{u}(\exp(tX)) = e^{ith}$.

$$\frac{d}{dt} f(\exp(tX))|_{t=0} = X(f)$$

$$dX_t =$$

$$\mathcal{B}(\mathcal{H}) \rightarrow L^\infty(\mathbb{T}; \mathcal{B}(\mathcal{H})); \quad a \mapsto \mathfrak{u}^*(z)a\mathfrak{u}(z)$$

Let $\rho_0 \in \mathcal{S}(\mathcal{H})$ be a given fixed state

$$\begin{aligned} \varrho_\mu(x) &= \int_{\mathbb{T}} \text{tr}[\mathfrak{u}^*(z)x\mathfrak{u}(z)\varrho_0] \mu(dz) = \int_{\mathbb{T}} f_{\rho_0,x}(z) \mu(dz) \\ &= \text{tr} \left[x \int_{\mathbb{T}} \mathfrak{u}(z)\varrho_0\mathfrak{u}^*(z) \mu(dz) \right] = \text{tr}[x\varrho_\mu] \end{aligned}$$

Consider iid random variables $(Z^i)_{i \in \mathbb{N}}$, $Z^i \sim \mu \in \mathcal{P}(\mathbb{T})$ and the corresponding sequence of empirical measures

$$\mu^n = \frac{1}{n} \sum_{i=1}^n \delta_{Z^i}, \quad \varrho_{\mu^n}(x) = \frac{1}{n} \sum_{i=1}^n \mathfrak{u}(Z^i)\varrho_0\mathfrak{u}^*(Z^i)$$

$$\begin{aligned}\frac{1}{n} \log \mathbb{E}[e^{n\langle f, \mu^n \rangle}] &= \frac{1}{n} \log \mathbb{E} \left[\exp \left(\sum_{i=1}^n f(Z^i) \right) \right] \\ &= \log \left(\int_{\mathbb{T}} e^{f(z)} \mu(dz) \right)\end{aligned}$$

$$\sup_{f=f_{\varrho_0,x}} \left\{ \langle f, \nu \rangle - \log \int_{\mathbb{T}} e^{f(z)} \mu(dz) \right\} = \sup_{x \in \mathcal{O}(\mathcal{H})} \left\{ \text{tr}[x \varrho_\nu] - \log \int_{\mathbb{T}} e^{\text{tr}[\mathfrak{u}^*(z) x \mathfrak{u}(z) \varrho_0]} \mu(dz) \right\}$$

$$\begin{aligned}& \sum_{j,k} x_{jk} \text{tr}[E_{jk} \varrho_\nu] - \log \int_{\mathbb{T}} e^{\sum_{j,k} x_{jk} \text{tr}[\mathfrak{u}^*(z) E_{jk} \mathfrak{u}(z) \varrho_0]} \mu(dz) \\&= \sum_{j,k} x_{jk} \text{tr}[E_{jk} \varrho_\nu] - \log \int_{\mathbb{T}} \prod_{j,k} e^{x_{jk} \text{tr}[\mathfrak{u}^*(z) E_{jk} \mathfrak{u}(z) \varrho_0]} \mu(dz) \\& \text{tr}[E_{jk} \varrho_\nu] = \frac{1}{c} \int_{\mathbb{T}} \text{tr}[\mathfrak{u}^*(z) E_{jk} \mathfrak{u}(z) \varrho_0] e^{\sum_{j,k} x_{jk} \text{tr}[\mathfrak{u}^*(z) E_{jk} \mathfrak{u}(z) \varrho_0]} \mu(dz) \\& b = \frac{1}{c} \int_{\mathbb{T}} v(z) e^{\langle \mathbf{x}, v(z) \rangle} \mu(dz)\end{aligned}$$

Consider the map

$$\mathbf{x} \mapsto \Lambda(\mathbf{x}) = \int_{\mathbb{T}} v(z) e^{\langle \mathbf{x}, v(z) \rangle} \mu(dz)$$

$$\begin{aligned}\langle \Lambda(\mathbf{x}) - \Lambda(\mathbf{y}), \mathbf{x} - \mathbf{y} \rangle &= \int_0^1 \langle D\Lambda((1-\tau)\mathbf{x} + \tau\mathbf{y})[\mathbf{x} - \mathbf{y}], \mathbf{x} - \mathbf{y} \rangle d\tau \geq 0 \\ \frac{d}{d\lambda} \int_{\mathbb{T}} v(z) e^{\langle \lambda, v(z) \rangle} \mu(dz) &= \int_{\mathbb{T}} v(z) \otimes v(z) e^{\langle \lambda, v(z) \rangle} \mu(dz)\end{aligned}$$

$$\text{tr}[E_{jk} \mathfrak{u}(z) \varrho_0 \mathfrak{u}^*(z)] = \langle \psi_k | \mathfrak{u}(z) \varrho_0 \mathfrak{u}^*(z) | \psi_j \rangle = \sum_i \lambda_i \langle \psi_k | \mathfrak{u}(z) | \psi_i \rangle \langle \psi_i | \mathfrak{u}^*(z) | \psi_j \rangle$$

If $\varrho_0 = \mathbb{1}$, then

$$\text{tr}[E_{jk} \mathfrak{u}(z) \varrho_0 \mathfrak{u}^*(z)] = \delta_{kj}.$$

Hence,

$$\text{tr}[E_{jk} \varrho_\nu] = \frac{1}{c} \delta_{kj} e^{\sum_j x_j}$$

4 Boson Fock spaces and quantum noise

4.1 Boson Fock spaces

Consider a (complex) Hilbert space \mathcal{H} with scalar product $\langle \cdot, \cdot \rangle_{\mathcal{H}}$. The symmetric Fock space associated with \mathcal{H} is

$$\mathfrak{F} = \mathfrak{F}_{\text{sym}}(\mathcal{H}) := \bigoplus_{n \in \mathbb{N}_0} \mathcal{H}^{\odot n}, \quad \mathcal{H}^{\odot 0} = \mathbb{C},$$

where \odot denotes the symmetric tensor product such that

$$\mathcal{H}^{\odot n} = \left\{ f \in \mathcal{H}^{\otimes n} : f(x_{\sigma_1}, \dots, x_{\sigma_n}) = f(x_1, \dots, x_n) \text{ for every permutation } \sigma \right\}.$$

The Fock space \mathfrak{F} inherits the scalar product from \mathcal{H} defined by

$$\langle \bigoplus f^{(n)}, \bigoplus g^{(n)} \rangle_{\mathfrak{F}} = \sum_{n \in \mathbb{N}_0} \langle f^{(n)}, g^{(n)} \rangle_{\mathcal{H}^{\otimes n}}.$$

We define the *vacuum vector* $\Omega = 1 \oplus 0 \oplus 0^{\otimes 2} \oplus \dots \in \mathfrak{F}$, and the *exponential vectors*

$$\mathbf{e}(f) = \bigoplus_{n \in \mathbb{N}_0} \frac{1}{\sqrt{n!}} f^{\otimes n}, \quad f \in \mathcal{H}.$$

It turns out that the family of exponential vectors \mathbf{E} is *total* in \mathfrak{F} , i.e., the linear span of \mathbf{E} is dense in \mathfrak{F} . This fact will be helpful for us in the future. Since,

$$\langle \mathbf{e}(f), \mathbf{e}(g) \rangle_{\mathfrak{F}} = \sum_{n \in \mathbb{N}_0} \frac{1}{n!} \langle f, g \rangle_{\mathcal{H}}^n = e^{\langle f, g \rangle_{\mathcal{H}}}, \quad f, g \in \mathcal{H},$$

the exponential vectors are normalizable. These normalized exponential vectors

$$\Psi(f) = e^{-\frac{1}{2}\|f\|_{\mathcal{H}}^2} \mathbf{e}(f), \quad f \in \mathcal{H},$$

are commonly known as *coherent vectors*.

4.1.1 Weyl and field operators

For any $f \in \mathcal{H}$, we define the *Weyl operator* on exponential vectors by

$$W(f)\mathbf{e}(g) := \exp \left(-\langle f, g \rangle_{\mathcal{H}} - \frac{1}{2} \|f\|_{\mathcal{H}}^2 \right) \mathbf{e}(f+g), \quad g \in \mathcal{H}.$$

Weyl operators play an essential role in the setup of Fock spaces. For one, they generate coherent states by acting on the vacuum state, i.e.,

$$W(f)\Omega = e^{-\frac{1}{2}\|f\|_{\mathcal{H}}^2} \mathbf{e}(f) = \Psi(f), \quad f \in \mathcal{H}.$$

Moreover, they give the means to map any element $f \in \mathcal{H}$ to unitary operators on \mathfrak{F} that satisfy the *canonical commutation relation* (CCR):

Proposition 4.1. *The Weyl operator $\mathsf{W}(f)$ is a unitary operator and satisfies*

- (i) $\mathsf{W}^\dagger(f) = \mathsf{W}(-f)$.
- (ii) $\mathsf{W}(f)\mathsf{W}(g) = e^{-i\text{Im}(\langle f, g \rangle_{\mathcal{H}})}\mathsf{W}(f+g) = e^{-2i\text{Im}(\langle f, g \rangle_{\mathcal{H}})}\mathsf{W}(g)\mathsf{W}(f)$.

Property (ii) is the Weyl form of the canonical commutation relation (CCR).

Proof. For any $g \in \mathcal{H}$,

$$\begin{aligned}\langle \mathsf{W}(f)\mathbf{e}(g), \mathsf{W}(f)\mathbf{e}(g) \rangle_{\mathfrak{F}} &= \exp(-2\langle f, g \rangle_{\mathcal{H}} - \|f\|_{\mathcal{H}}^2) \langle \mathbf{e}(f+g), \mathbf{e}(f+g) \rangle_{\mathfrak{F}} \\ &= \exp(-2\langle f, g \rangle_{\mathcal{H}} - \|f\|_{\mathcal{H}}^2 + \|f+g\|_{\mathcal{H}}^2) \\ &= e^{\|g\|_{\mathcal{H}}^2} = \langle \mathbf{e}(g), \mathbf{e}(g) \rangle_{\mathfrak{F}}.\end{aligned}$$

Hence, $\mathsf{W}(f)$ preserves inner products on \mathfrak{E} . Since \mathfrak{E} is dense in \mathfrak{F} , $\mathsf{W}(f)$ extends uniquely to an isometry on \mathfrak{F} .

In a similar fashion, we compute

$$\mathsf{W}(-f)\mathsf{W}(f)\mathbf{e}(g) = \mathbf{e}(g) = \mathsf{W}(f)\mathsf{W}(-f)\mathbf{e}(g) \quad \forall g \in \mathcal{H},$$

i.e., $\mathsf{W}(-f)\mathsf{W}(f)$ is the identity on the dense set \mathfrak{E} . In particular, $\mathsf{W}(f)$ is surjective and an isometry, i.e., $\mathsf{W}(f)$ is unitary with $\mathsf{W}^\dagger(f) = \mathsf{W}(-f)$.

As for the last property, we observe that

$$\begin{aligned}\mathsf{W}(f)\mathsf{W}(g)\mathsf{W}(-(f+g))\mathbf{e}(h) &= e^{\langle f+g, h \rangle_{\mathcal{H}} - \frac{1}{2}\|f+g\|_{\mathcal{H}}^2}\mathsf{W}(f)\mathsf{W}(g)\mathbf{e}(-(f+g) + h) \\ &= e^{\langle f+g, h \rangle_{\mathcal{H}} - \frac{1}{2}\|f\|_{\mathcal{H}}^2 - \langle g, -f+h \rangle_{\mathcal{H}} - \text{Re}(\langle g, f \rangle_{\mathcal{H}})}\mathsf{W}(f)\mathbf{e}(-f+h) \\ &= e^{-i\text{Im}(\langle f, g \rangle_{\mathcal{H}})}\mathbf{e}(h),\end{aligned}$$

and hence, $\mathsf{W}(f)\mathsf{W}(g)\mathsf{W}(-(f+g)) = e^{-i\text{Im}(\langle f, g \rangle_{\mathcal{H}})}I_{\mathfrak{F}}$. We then conclude by using property (i) of Weyl operators. \square

Since $\mathsf{W}(f)$ is unitary for every $f \in \mathcal{H}$, the family $\{\mathsf{W}(tf)\}_{t \in \mathbb{R}}$ forms a one-parameter (strongly continuous) group of unitaries. In particular, due to Stone's theorem, it has a corresponding Hermitian operator $\mathsf{P}(f)$ such that

$$\mathsf{W}(tf) = \exp(it\mathsf{P}(f)).$$

We further define the following operators

$$\mathsf{Q}(f) := -\mathsf{P}(if), \quad \mathsf{A}^-(f) := \frac{\mathsf{Q}(f) + i\mathsf{P}(f)}{2}, \quad \mathsf{A}^+(f) := \frac{\mathsf{Q}(f) - i\mathsf{P}(f)}{2}.$$

The operators A^\pm are called the *field operators* and will play an essential role as they form the creation/annihilation operators on Fock spaces.

Proposition 4.2. *The following are true: For any $f, g \in \mathcal{H}$,*

- (i) \mathfrak{E} is a core for $\mathsf{P}(f)$ and $[\mathsf{P}(f), \mathsf{P}(g)] = 2i\text{Im}(\langle f, g \rangle_{\mathcal{H}})\mathsf{I}_{\mathfrak{F}}$.

- (ii) $\mathbf{A}^-(f)\mathbf{e}(g) = \langle f, g \rangle_{\mathcal{H}} \mathbf{e}(g)$, $\mathbf{A}^+(f)\mathbf{e}(g) = \frac{d}{dt}\mathbf{e}(g + tf)|_{t=0}$.
 - (iii) $\mathbf{W}^\dagger(f)\mathbf{A}^-(g)\mathbf{W}(f) = \mathbf{A}^-(g) + \langle g, f \rangle_{\mathcal{H}} \mathbf{I}_{\mathfrak{F}}$, $\mathbf{W}^\dagger(f)\mathbf{A}^+(g)\mathbf{W}(f) = \mathbf{A}^+(g) + \langle f, g \rangle_{\mathcal{H}} \mathbf{I}_{\mathfrak{F}}$.
 - (iv) $[\mathbf{A}^-(f), \mathbf{A}^-(g)] = [\mathbf{A}^+(f), \mathbf{A}^+(g)] = 0$, $[\mathbf{A}^-(f), \mathbf{A}^+(g)] = \langle f, g \rangle_{\mathcal{H}} \mathbf{I}_{\mathfrak{F}}$,
- i.e., the field operators \mathbf{A}^\pm satisfy the canonical commutation relation.

Remark 4.3 On the finite particle vectors, the field operators act as

$$\mathbf{A}^-(f)g^{\otimes n} = \sqrt{n}\langle f, g \rangle_{\mathcal{H}} g^{\otimes(n-1)}, \quad \mathbf{A}^+(f)g^{\otimes(n-1)} = \frac{1}{\sqrt{n}} \sum_{k=0}^{n-1} g^{\otimes k} \otimes f \otimes g^{\otimes(n-1-k)}.$$

Example 4.4 Let $\mathcal{H} = L^2_{\mathbb{C}}([-\pi, \pi])$ be the space of square-integrable functions. Then the countable set $\{\psi_\ell(x) = e^{i\ell x} : \ell \in \mathbb{Z}\}$ forms an orthonormal basis for \mathcal{H} .

$$\psi(x) = \sum_{\ell} e^{i\ell x} \mathbf{A}_\ell^-$$

4.1.2 Second quantization

The term *second quantization* is associated with the action of lifting operators on an k -particle space $\mathcal{H}^{\otimes k}$ to an associated operator on the Fock space.

We begin our discussion with the *1-particle* case. For any bounded operator $A \in \mathcal{B}(\mathcal{H})$ one defines the map

$$\Gamma(A) := I \oplus \bigoplus_{n \in \mathbb{N}} A^{\otimes n},$$

which acts on $\mathcal{H}^{\otimes n}$ by

$$\Gamma(A)g_1 \otimes \cdots \otimes g_n = Ag_1 \otimes \cdots \otimes Ag_n.$$

Clearly, if A is unitary, then so is $\Gamma(A)$. Indeed, in this case, one has

$$\langle \Gamma(A)e(g), \Gamma(A)e(g) \rangle_{\mathfrak{F}(\mathcal{H})} = \sum_{n \in \mathbb{N}_0} \frac{1}{n!} \langle Ag, Ag \rangle_{\mathcal{H}}^n = \sum_{n \in \mathbb{N}_0} \frac{1}{n!} \|g\|_{\mathcal{H}}^n = \langle e(g), e(g) \rangle_{\mathfrak{F}(\mathcal{H})}.$$

Now let H be a self-adjoint operator on \mathcal{H} and $U_t := \exp(itH)$ be its unitary evolution. Then, $\Gamma(U_t)$ is a one-parameter group of unitary operators on $\mathfrak{F}(\mathcal{H})$. Stone's theorem then provides the existence of a densely defined Hermitian operator $d\Gamma(H)$ such that

$$\Gamma(U_t) = \exp(it d\Gamma(H)).$$

The generator $d\Gamma(H)$ is called the *second quantization* of H , and takes the explicit form

$$d\Gamma(H)g^{(n)} = \sum_{j=1}^n g_1 \otimes \cdots \otimes g_{j-1} \otimes Hg_j \otimes g_{j+1} \otimes \cdots \otimes g_n = \sum_{j=1}^n H_j g^{(n)},$$

for any $g^{(n)} = g_1 \otimes \cdots \otimes g_n \in \mathcal{H}^{\otimes n}$ with $g_j \in D(H)$ and

$$H_j = I_{\mathcal{H}} \otimes \cdots \otimes I_{\mathcal{H}} \otimes H \otimes I_{\mathcal{H}} \otimes \cdots \otimes I_{\mathcal{H}},$$

where H acts on the j -th tensor product. The special case $H = I_{\mathcal{H}}$ yields

$$\mathsf{N}g^{(n)} := \mathsf{d}\Gamma(I_{\mathcal{H}})g^{(n)} = ng^{(n)}, \quad n \in \mathbb{N},$$

and is called the *number operator* due to its diagonal nature, with eigenvalues representing the number of particles in each configuration. Its domain is given by

$$D(\mathsf{N}) = \left\{ \{f^{(n)}\}_{n \in \mathbb{N}_0} : \sum_{n \in \mathbb{N}_0} n^2 \|f^{(n)}\|_{\mathcal{H}^{\otimes n}}^2 < +\infty \right\}.$$

$$\begin{aligned} \mathsf{A}_j^- \mathsf{A}_j^+ \psi_j^{\otimes n} &= \sqrt{n} \mathsf{A}_j^- \psi_j^{\otimes(n-1)} = \sum_{k=0}^{n-1} \psi_j^{\otimes k} \otimes \psi_j \otimes \psi_j^{\otimes(n-1-k)} = n \psi_j^{\otimes n} \\ \mathsf{A}_j^+ \mathsf{A}_j^- \mathbf{e}(\psi_k) &= \delta_{jk} \mathsf{N}_j \mathbf{e}(\psi_j) \end{aligned}$$

This construction can be performed similarly for the general k -particle case.

$$\mathsf{d}\Gamma(H^{(k)})f^{(n)} = \sum_{j_1 \neq \dots \neq j_k} H_{j_1 \dots j_k} f^{(n)},$$

where H_{j_1, \dots, j_k} denotes the operator where $H^{(k)}$ acts on the (j_1, \dots, j_k) -th tensor product.

Example 4.5 Consider the Hilbert space $\mathcal{H} = L^2_{\mathbb{C}}(\mathbb{T})$ and the Hamiltonians,

$$H^{(1)} = -\Delta \in \mathcal{O}(\mathcal{H}), \quad H^{(2)} = \mathsf{M}_W \in \mathcal{O}(\mathcal{H}^{\otimes 2}),$$

where $W : \mathbb{T} \times \mathbb{T} \rightarrow \mathbb{R}$ is an interaction potential, and M_f denotes the multiplication operator corresponding to f .

Then, their second quantization is given by

$$\mathsf{d}\Gamma(H^{(1)}) = \sum_{j=1}^n H_j^{(1)}, \quad \mathsf{d}\Gamma(H^{(2)}) = \sum_{j \neq \ell} H_{j\ell}^{(2)} \quad \text{on } \mathcal{H}^{\otimes n}.$$

$$\mathsf{d}\Gamma(H^{(1)})\psi_j^{\otimes n} = nj^2 \psi_j^{\otimes n} = j^2 \mathsf{A}_j^+ \mathsf{A}_j^- \psi_j^{\otimes n} = j^2 \mathsf{A}_j^+ \mathsf{A}_j^- \psi_j^{\otimes n}$$

$$\psi^{(n)} = \psi_{k_1} \otimes \dots \otimes \psi_{k_n}$$

$$\mathsf{d}\Gamma(H^{(1)})\psi^{(n)} = \sum_{j=1}^n k_j^2 \psi^{(n)}$$

4.1.3 Free field operators

For each $f \in \mathcal{H}$, we consider the pair $\{a^-(f), a^+(f)\}$ of operators on $\mathfrak{F}_{\text{sym}}(\mathcal{H})$ defined by

$$a^-(f)\Psi =$$

satisfying the commutation relations: For any $f, g \in \mathcal{H}$,

$$[a^\pm(f), a^\pm(g)] = 0, \quad [a^-(f), a^+(g)] = \langle f, g \rangle,$$

where $f \mapsto a^-(f)$ is conjugate linear and $f \mapsto a^+(f)$ is linear. Moreover, if $\Omega \in \mathfrak{F}_{\text{sym}}(\mathcal{H})$ is the *vacuum vector*, then $a^-(f)\Omega = 0$ for every $f \in \mathcal{H}$. The *field operators* a^- and a^+ are called *creation* and *annihilation* operators, respectively. On appropriate domains, the field operators are adjoints of one another, i.e., $(a^-(f))^\dagger = a^+(f)$ for every $f \in \mathcal{H}$.

It is common in the physics literature to consider operator-valued distributions $\{a_x^-, a_x^+\}$ instead, where if $\mathcal{H} = L^2_{\mathbb{C}}(\mathbb{R}^d)$, then $x \in \mathbb{R}^d$ and

$$a^-(f) = \int_{\mathbb{R}^d} \overline{f(x)} a_x \, dx, \quad a^+(f) = \int_{\mathbb{R}^d} f(x) a_x \, dx.$$

The commutation relations then simply read

$$[a_x^\pm, a_y^\pm] = 0, \quad [a_x^-, a_y^+] = \delta(x - y).$$

The *number operator* is formally defined by $N_x = a_x^+ a_x$, $x \in \mathbb{R}^d$.

On the rigorous note, if \mathcal{H} is separable with orthonormal basis $\{\psi_i\}$, then one obtains a family of field operators $\{a_i^-, a_i^+\}$ with $a_i^\pm : a^\pm(\psi_i)$.

4.1.4 Gaussian states

Definition 4.6 (Gaussian states) A state Ψ on $\mathfrak{F}_{\text{sym}}(\mathcal{H})$ is said to be a mean-zero *Gaussian* (or *quasi-free*) state if it can be uniquely determined from the field operators $\{a, a^\dagger\}$ by its covariance

$$\Sigma_\Psi(f, g) := \begin{pmatrix} \Psi(a^+(f) a^-(g)) & \Psi(a^-(f) a^-(g)) \\ \Psi(a^+(f) a^+(g)) & \Psi(a^-(f) a^+(g)) \end{pmatrix}, \quad f, g \in \mathcal{H}.$$

If the off-diagonal elements of the covariance are zero, the state Ψ is called *gauge-invariant* since it is invariant under the so-called gauge transformations of the first kind, i.e.,

$$a^\pm(f) = e^{\pm i\alpha} a^\pm(f) \quad \text{for any } \alpha \in \mathbb{R}.$$

If the off-diagonal elements of the covariance are nonzero, the state Ψ is called *squeezed*.

4.1.5 States invariant under free evolutions

Consider a Hamiltonian H on $\mathfrak{F}_{\text{sym}}(\mathcal{H})$ with $\mathcal{H} = L^2_{\mathbb{C}}(\mathbb{R}^d)$ and its associated 1-parameter automorphism group

$$\mathfrak{u}_t(a) = e^{itH} a e^{-itH}, \quad t \in \mathbb{R}.$$

Definition 4.7 The Hamiltonian H on $\mathfrak{F}_{\text{sym}}(\mathcal{H})$ is called *free* if there exists a real-valued function $\omega: \mathbb{R}^d \rightarrow \mathbb{R}$ such that

$$\mathfrak{u}_t(a_x^-) = e^{-it\omega(x)}a_x^-, \quad x \in \mathbb{R}^d.$$

In this case, the function ω is called the *free 1-particle Hamiltonian*, and H is said to be the *second quantization* of ω . Accordingly, \mathfrak{u}_t is called a *free evolution*.

Example 4.8 Consider a Hamiltonian on $\mathcal{H} = L^2_{\mathbb{C}}(\mathbb{R})$ and its eigensystem $\{(\lambda_i, \psi_i)\}_i$ such that $\{\psi_i\}_i$ forms an orthonormal basis for \mathcal{H} . Setting $a_i^\pm = a^\pm(\psi_i)$, we see that

$$[a_i^\pm, a_j^\pm] = 0, \quad [a_i^-, a_j^+] = \delta_{ij},$$

i.e., a_i^\pm define field operators on the symmetric Fock space $\mathfrak{F}_{\text{sym}}(\mathcal{K})$ with $\mathcal{K} = \ell^2_{\mathbb{C}}$

Defining the field operators

$$a^- = \frac{1}{\sqrt{2}}(Q + iP), \quad a^+ = \frac{1}{\sqrt{2}}(Q - iP), \quad N := a^+a^-$$

such that $[a^-, a^+] = I$, we find that

$$H = a^+a^- + \frac{1}{2} = N + \frac{1}{2}.$$

Clearly, the eigenvectors of N and H coincide

Definition 4.9 A Gaussian state Ψ on $\mathfrak{F}_{\text{sym}}(\mathcal{H})$ with $\mathcal{H} = L^2(\mathbb{R}^d)$ is said to be invariant under a free evolution \mathfrak{u}_t if

$$\Psi(\mathfrak{u}_t(a_x^-) \mathfrak{u}_s(a_y^+)) = \Psi(\mathfrak{u}_{t-s}(a_x^-) a_y^+) = \Psi(a_x^- \mathfrak{u}_{s-t}(a_y^+)) \quad \text{for all } t \in \mathbb{R}, x \in \mathbb{R}^d.$$

A *Gaussian free state* is a Gaussian state that is invariant under *all* free evolutions.

Theorem 4.10. *A Gaussian state Ψ on $\mathfrak{F}_{\text{sym}}(\mathcal{H})$ is a Gaussian free state if and only if it is gauge-invariant and its diagonal correlations are supported on the diagonal, i.e.,*

$$\Psi(a_x^- a_y^+) = m(x)\delta(x - y), \quad \Psi(a_x^+ a_y^-) = n(x)\delta(x - y).$$

Theorem 4.11. *The field operators $\{a_x^+, a_x^-\}$ satisfying the commutation relations*

$$[a_x^\pm, a_y^\pm] = 0, \quad [a_x^-, a_y^+] = m(x)\delta(x - y).$$

are mean-zero Gaussian random variables w.r.t. the Fock vacuum state $\Psi_\Omega = \langle \Omega, \cdot \Omega \rangle$, where $\Omega \in \mathfrak{F}_{\text{sym}}(\mathcal{H})$ is the Fock vacuum vector, with covariance

$$\Sigma_\Psi(x, y) = \begin{pmatrix} 0 & 0 \\ 0 & m(x) \end{pmatrix} \delta(x - y), \quad x, y \in \mathbb{R}^d.$$

Conversely, if $\{a_x^+, a_x^-\}$ are random variables with these properties, then they satisfy the commutation relations above.

4.1.6 Boson Fock white noise

Definition 4.12 A boson Fock white noise on \mathbb{R}^d is a boson Fock field $\{b_{t,x}^+, b_{t,x}^-\}$ on \mathbb{R}^{d+1} with vacuum vector Ω satisfying the commutation relations

$$[b_{t,x}^\pm, b_{s,y}^\pm] = 0, \quad [b_{t,x}^-, b_{s,y}^+] = \delta(t-s)m(x)\delta(x-y), \quad b_{t,x}^- \Omega = 0.$$

Theorem 4.13. Let $\{a_x^+, a_x^-\}$ be Gaussian free fields w.r.t. the Fock vacuum state Ψ_Ω with

$$\mathfrak{u}_t(a_x^-) = e^{-it\omega(x)}a_x^-.$$

Then the rescaled field operators

$$b_{t,x}^{\lambda,\pm} := \frac{1}{\lambda}\mathfrak{u}_{t/\lambda^2}(a_x^\pm)$$

converges in the sense of correlator distributions to a boson Fock white noise, i.e.,

$$\lim_{\lambda \rightarrow 0} \Psi(b_{t,x}^{\lambda,\varepsilon_1} b_{t,y}^{\lambda,\varepsilon_2}) = \Psi(b_{t,x}^{\varepsilon_1} b_{t,y}^{\varepsilon_2}) \quad \varepsilon_1, \varepsilon_2 \in \{+, -\},$$

where $b_{t,x}^\pm$ as defined in Definition 4.12 with $m = 2\pi\delta(\omega)$.

Proof. Using the invariance of Ψ_Ω under free evolutions, we find

$$\begin{aligned} \Psi_\Omega(b_{t,x}^{\lambda,-} b_{s,y}^{\lambda,+}) &= \frac{1}{\lambda^2} \Psi_\Omega(\mathfrak{u}_{(t-s)/\lambda^2}(a_x^-) a_y^+) \\ &= \frac{1}{\lambda^2} e^{-i\omega(x)(t-s)/\lambda^2} \Psi_\Omega(a_x^- a_y^+) = \frac{1}{\lambda^2} e^{-i\omega(x)(t-s)/\lambda^2} \delta(x-y). \end{aligned}$$

Passing to the limit $\lambda \rightarrow 0$ recovers the desired limit. All the other terms vanish. \square

Remark 4.14 Let $\mathcal{K} \subset L^2(\mathbb{R}^d)$ be a set of functions for which

$$\int_{\mathbb{R}} |\langle f, e^{it\omega} g \rangle| dt < +\infty \quad \forall f, g \in \mathcal{K}.$$

Since $t \mapsto \langle f, e^{it\omega} f \rangle$ is positive definite for each $f \in \mathcal{K}$, Bochner's theorem implies that the sesquilinear form

$$\langle f, 2\pi\delta(\omega)g \rangle := \int_{\mathbb{R}} \langle f, e^{it\omega} g \rangle dt,$$

is a pre-scalar product. With $(\cdot|\cdot)$, the set \mathcal{K} becomes a pre-Hilbert space, which can be completed to obtain a Hilbert space, still denoted by \mathcal{K} . The function $m = 2\pi\delta(\omega)$ has to be understood in this sense, and only makes sense for functions in \mathcal{K} .

The operators

$$B_t^-(f) := \int_0^t \int_{\mathbb{R}^d} \overline{f(x)} b_{s,x}^- dx ds, \quad B_t^+(f) := \int_0^t \int_{\mathbb{R}^d} f(x) b_{s,x}^+ dx ds,$$

define *quantum Brownian motions*. The self-adjoint (*momentum*) operators

$$P_t(f) := \frac{1}{i} [B_t^-(f) - B_t^+(f)],$$

form a commuting family of classical random variables whose statistics in the Fock vacuum state $\Psi_\Omega = \langle \Omega, \cdot | \Omega \rangle$ is completely determined by the relation

$$\Psi(e^{iP_t(f)}) = \exp\left(-\frac{t}{2} \|f\|_{L^2(\mathbb{R}^d)}^2\right). \quad \text{check!}$$

4.1.7 Gaussian equilibrium states: The KMS condition

For any states $a, b \in \mathcal{A}$, the map $t \mapsto \Psi(a \mathfrak{u}_t(b))$ can be analytically continued and satisfies the so-called *KMS condition* at inverse temperature $\beta > 0$:

$$\Psi(a \mathfrak{u}_{t+i\beta}(b)) = \Psi(\mathfrak{u}_t(a) b) \quad \forall a, b \in \mathcal{A}. \quad \text{check!}$$

4.2 Composite systems

Definition 4.15 A composite system of two given quantum dynamical systems $S = \{\mathcal{H}_S, H_S\}$, $R = \{\mathcal{H}_R, H_R\}$ is a quantum dynamical system of the form

$$\{\mathcal{H}_S \otimes \mathcal{H}_R, H_{SR}\}, \quad H_{SR} = H_S \otimes 1_R + 1_S \otimes H_R + H_I,$$

where H_I is called the *interaction Hamiltonian* and contains all the new physics of the composite system, while $H_0 := H_S \otimes 1_R + 1_S \otimes H_R$ is called the free Hamiltonian.

We will consider *scaled* total Hamiltonians

$$H^\lambda := H_0 + \lambda H_I,$$

and the following unitary evolutions:

$$\begin{aligned} \text{free evolution} \quad V_t^0 &= e^{-itH_0}, & \text{total evolution} \quad V_t^\lambda &= e^{-itH^\lambda}, \\ \text{interacting representation evolution} \quad U_t^\lambda &= (V_t^0)^\dagger V_t^\lambda, \end{aligned}$$

where U_t^λ satisfies the Schrödinger equation in the interaction picture:

$$\partial_t U_t^\lambda = -i\lambda H_I(t) U_t^\lambda, \quad U_0^\lambda = I,$$

with the time dependent Hamiltonian $H_I(t) = (V_t^0)^\dagger H_I V_t^0$.

For simplicity, we will make the following assumptions on system $S = \{\mathcal{H}_S, H_S\}$, the reservoir $R = \{\mathcal{H}_R, H_R\}$, and the interaction Hamiltonian H_I .

4.2.1 The reservoir

The reservoir $R = \{\mathcal{H}_R, H_R\}$ is given by the Hilbert space $\mathfrak{F}_{\text{sym}}(\mathcal{H})$ with $\mathcal{H} = L^2_{\mathbb{C}}(\mathbb{R}^d)$, a free Hamiltonian H_R with continuous spectrum \mathbb{R} , and a mean-zero Gaussian free state Ψ_R such that

$$\int_{\mathbb{R}} |\Psi_R(a_x^{\varepsilon_1} \mathfrak{u}_t(a_y^{\varepsilon_2}))| dt < +\infty, \quad \varepsilon_1, \varepsilon_2 \in \{+, -\}.$$

In particular, Ψ_R is characterized by the covariances

$$\Psi_R(a_x^- a_y^+) = m(x) \delta(x - y), \quad \Psi_R(a_x^+ a_y^-) = n(x) \delta(x - y).$$

Since H_R is free, there exists a function ω , for which

$$\mathfrak{u}_t(a_x^-) = e^{itH_R} a_x^- e^{-itH_R} = e^{-it\omega(x)} a_x^-,$$

where ω describes the 1-particle evolution.

An example of a free reservoir Hamiltonian is given by

$$H_R = \int_{\mathbb{R}^d} \omega(x) a_x^+ a_x^- dx,$$

where ω is a smooth cutoff function.

4.2.2 The system Hamiltonian

For simplicity, we shall assume that the system Hamiltonian H_S has a discrete spectrum such that

$$H_S = \sum_j \lambda_j P_j,$$

where λ_j are the eigenvalues and P_j are their corresponding spectral projections.

4.2.3 The interaction Hamiltonian

We consider dipole-type interaction Hamiltonians of the form

$$H_I = \int_{\mathbb{R}^d} [D(x) \otimes a_x^+ + D^\dagger(x) \otimes a_x^-] dx,$$

where $\{D(x) : x \in \mathbb{R}^d\}$ is a family of system operators called the *response terms*.

With the spectral projections of H_S , we may express H_I as

$$H_I = \sum_{j,k} \int_{\mathbb{R}^s} [P_j D(x) P_k \otimes a_x^+ + P_k D^\dagger(x) P_j \otimes a_x^-] dx,$$

and hence, the time-dependent Hamiltonian reads

$$\begin{aligned} H_I(t) &= \sum_{j,k} \int_{\mathbb{R}^d} [P_j D(x) P_k \otimes e^{it(\omega(x)+\lambda_j-\lambda_k)} a_x^+ + P_k D^\dagger(x) P_j \otimes e^{-it(\omega(x)+\lambda_j-\lambda_k)} a_x^-] dx \\ &= \sum_q \sum_{\lambda_k - \lambda_j = \eta_q} \int_{\mathbb{R}^d} [P_j D(x) P_k \otimes e^{it(\omega(x)-\eta_q)} a_x^+ + P_k D^\dagger(x) P_j \otimes e^{-it(\omega(x)-\eta_q)} a_x^-] dx \\ &= \sum_q \int_{\mathbb{R}^d} [D_q(x) \otimes e^{it(\omega(x)-\eta_q)} a_x^+ + D_q^\dagger(x) \otimes e^{-it(\omega(x)-\eta_q)} a_x^-] dx, \end{aligned}$$

where the system operators

$$D_q(x) := \sum_{\lambda_k - \lambda_j = \eta_q} P_j D(x) P_k \quad \text{satisfy} \quad e^{itH_S} D_q(x) e^{-itH_S} = e^{-it\eta_q} D_q(x).$$

To simplify things drastically, we assume that $q = 1$ and $D_1(x) = \chi(x)D$ for some smooth cutoff function χ and a fixed system operator D . In this case, we obtain

$$H_I(t) = \int_{\mathbb{R}^d} [D \otimes \chi(x) e^{it(\omega(x)-\eta)} a_x^+ + D^\dagger \otimes \overline{\chi(x)} e^{-it(\omega(x)-\eta)} a_x^-] dx.$$

4.3 The weak interaction stochastic limit

Altogether, we arrive at the rescaled Schrödinger equation in the interaction picture

$$U_{t/\lambda^2}^\lambda = I - i \int_0^t H_I^\lambda(s) U_{s/\lambda^2}^\lambda ds, \tag{4.1} \quad \text{eq:rescaled-schrodinger-interaction}$$

with $H_I^\lambda(t) = D \otimes b_t^{\lambda,+} + D^\dagger \otimes b_t^{\lambda,-}$, where, due to Theorem 4.13,

$$b_t^{\lambda,\pm} = \frac{1}{\lambda} a^\pm(\chi e^{i(t/\lambda^2)(\omega-\eta)}) \longrightarrow b_t^\pm \quad \text{in the sense of correlators,}$$

and therefore,

$$H_I^\lambda(t) \longrightarrow H_t = D \otimes b_t^+ + D^\dagger \otimes b_t^-, \quad U_{t/\lambda^2}^\lambda \longrightarrow U_t,$$

where U_t satisfies the formal stochastic differential equation

$$dU_t = -i(D \otimes dB_t^+ + D^\dagger \otimes dB_t^-)U_t, \quad B_t^\pm = \int_0^t b_s^\pm ds. \quad (4.2)$$

Under certain assumptions on ω , this SDE has a unique solution.

The SDE (4.2) in natural-time order (or Itô form) is the SDE given by

$$dU_t = -i(D dB_t^+ U_t + D^\dagger U_t dB_t^-) - \gamma_- D^\dagger D U_t dt, \quad (4.3)$$

obtained by commuting $dB_t^- U_t$ and using the fact that the solution U_t to (4.2) satisfies

$$\begin{aligned} [b_t^-, U_t] &= -i\gamma_- D U_t, \\ [U_t^\dagger, b_t^+] &= i\bar{\gamma}_- U_t^\dagger D^\dagger, \quad \text{where } \gamma_- := \int_{-\infty}^0 \langle \chi, e^{-it(\omega-\eta)} \chi \rangle dt. \\ [b_t^-, U_t^\dagger] &= i\gamma_- U_t^\dagger D, \end{aligned}$$

The additional term is known as the Itô correction term in quantum stochastic calculus. Morally, the natural-time order is the order induced by the filtration for which one expects U_t to be adapted to. **Expand on the concept of normal order!**

With the unitary evolution U_t one can derive an evolution for any system observable $X_t = U_t^\dagger (X \otimes 1_{\mathcal{H}_R}) U_t$ given by

$$dX_t = -i dB_t^+ [D, X_t] + i[X_t, D^\dagger] dB_t^- + L X_t dt$$

where

$$LX = 2\Re(\gamma_-) D^\dagger X D - \gamma_- D^\dagger D X - \bar{\gamma}_- X D^\dagger D,$$

is the corresponding Lindblad operator.

Remark 4.16 The rescaled equation (4.1) may be formally expressed as

$$U_{t/\lambda^2}^\lambda = I + \sum_{n=1}^{\infty} (-i)^n \lambda^n \int_0^{t/\lambda^2} \cdots \int_0^{t_{n-1}} H_I(t_1) \cdots H_I(t_n) dt_1 \cdots dt_n.$$

In the case when $H_I(t)$ commutes for every $t \geq 0$, the integrand would be a symmetric function of t_1, \dots, t_n , and gives

$$U_{t/\lambda^2}^\lambda = I + \sum_{n=1}^{\infty} (-i)^n \frac{\lambda^n}{n!} \left(\int_0^{t/\lambda^2} H_I(s) ds \right)^n = \exp \left(-i\lambda \int_0^{t/\lambda^2} H_I(s) ds \right),$$

i.e., the expectation of U_{t/λ^2}^λ is the characteristic function of the process

$$W_t^\lambda := \lambda \int_0^{t/\lambda^2} H_I(s) ds.$$

4.4 Damped harmonic oscillator

We consider a simple setup in which a single atom interacts with an electromagnetic field. The total system is given by the composite of the atom and the reservoir systems

$$S = \{\mathfrak{F}_{\text{sym}}(\mathbb{C}^2), H_S\}, \quad R = \{\mathfrak{F}_{\text{sym}}(L^2(\mathbb{R}^d)), H_R\},$$

with the free Hamiltonians

$$H_S = \omega_0 c^+ c^-, \quad H_R = \int_{\mathbb{R}^d} \omega(x) a_x^+ a_x^- dx,$$

where $\{c^+, c^-\}$ is the field operator for $\mathcal{H}_S = \mathfrak{F}_{\text{sym}}(\mathbb{C}^2)$, $\{a_x^+, a_x^-\}$ are the field operators for $\mathcal{H}_R = \mathfrak{F}_{\text{sym}}(L^2(\mathbb{R}^d))$, and ω is a suitable cutoff function.

For the interaction Hamiltonian H_I , we consider a dipole approximation of the form

$$H_I = \int_{\mathbb{R}^d} \chi(x) [c^- \otimes a_x^+ + c^+ \otimes a_x] dx = c^- \otimes A^+ + c^+ \otimes A^-,$$

where χ are suitable cutoff function and

$$A^\pm = \int_{\mathbb{R}^d} \chi(x) a_x^\pm dx.$$

such that the rescaled total Hamiltonian is given by

$$H^\lambda = H_S \otimes 1_{\mathcal{H}_R} + 1_{\mathcal{H}_S} \otimes H_R + \lambda H_I.$$

Define the evolution

$$\mathfrak{u}_t^\lambda(a) = e^{itH^\lambda} a e^{-itH^\lambda}.$$

Then, the Heisenberg equation for $c(t) = \mathfrak{u}_t^\lambda(c \otimes 1_{\mathcal{H}_R})$ and $a_x^-(t) = \mathfrak{u}_t^\lambda(1_{\mathcal{H}_S} \otimes a_x)$ reads

$$\begin{aligned} \frac{d}{dt} c^-(t) &= -i\omega_0 c^-(t) - i\lambda \int_{\mathbb{R}^d} \chi(x) a_x^-(t) dx \\ \frac{d}{dt} a_x^-(t) &= -i\omega(x) a_x^-(t) - \lambda \chi(x) c(t). \end{aligned}$$

Solving for $a_x^-(t)$, one obtains

$$a_x^-(t) = a_x^- e^{-it\omega(x)} - i\lambda \chi(x) \int_0^t e^{-i(t-s)\omega(x)} c^-(s) ds.$$

Inserting $a_x^-(t)$ into the equation for $c(t)$, we find

$$\frac{d}{dt} c^-(t) = -i\omega_0 c^-(t) - \int_0^t \gamma(t-s) c^-(s) ds - i\xi_t^-.$$

where we defined the quantities

$$\gamma(r) := \lambda^2 \int_{\mathbb{R}^d} e^{-ir\omega(x)} \chi^2(x) dx, \quad \xi_t^- := \lambda \int_{\mathbb{R}^d} e^{-it\omega(x)} \chi(x) a_x^- dx.$$

Observe that ξ_t^- depends only on the reservoir field operator a_x^- , and therefore, acts as an *external force* to the atomic system.

Now let us consider the statistics of ξ_t^- for the Fock vacuum state free state $\Psi_\Omega = \langle \Omega, \cdot \Omega \rangle$ on \mathcal{H}_R . Since Ψ_Ω is a Gaussian free state, the statistics of ξ_t^- are uniquely determined by the 2-points correlation functions. Clearly, Ψ_Ω has zero mean, and therefore,

$$\Psi_\Omega(\xi_t^-) = \lambda \int_{\mathbb{R}^d} e^{-it\omega(x)} \chi(x) \Psi_\Omega(a_x^-) dx = 0, \quad \Psi_\Omega((\xi_t^-)^\dagger) = 0.$$

Moreover, we find

$$\Psi_\Omega((\xi_t^-)^\dagger \xi_s^-) = \Psi_\Omega(\xi_t^- \xi_s^-) = \Psi_\Omega((\xi_t^-)^\dagger (\xi_s^-)^\dagger) = 0,$$

and

$$\begin{aligned} \Psi_\Omega(\xi_t^- (\xi_s^-)^\dagger) &= \lambda^2 \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} e^{-it(\omega(x)-\omega(y))} \chi(x) \chi(y) \Psi_\Omega(a_x^- a_y^+) dx dy \\ &= \lambda^2 \int_{\mathbb{R}^d} e^{-i(t-s)\omega(x)} \chi^2(x) dx = \gamma(t-s). \end{aligned}$$

Hence, ξ_t^- is a mean-zero γ -correlated *random process* under Ψ .

To obtain the stochastic limit, we start by rescaling time $t \mapsto t/\lambda^2$ and consider the rescaled quantity $c^{\lambda,\pm}(t) := c^\pm(t/\lambda^2)$, which yields

$$\dot{c}^{\lambda,-}(t) = -i\omega_0 c^{\lambda,-}(t) - \int_0^t \gamma^\lambda(t-s) c^{\lambda,-}(s) ds - i \xi_t^{\lambda,-}, \quad \text{eq:quantum-pre-langevin} \quad (4.4)$$

with

$$\gamma^\lambda(r) := \frac{1}{\lambda^2} \langle \chi, e^{-i(r/\lambda^2)\omega} \chi \rangle, \quad \xi_t^{\lambda,-} := \lambda^{-2} \xi_{t/\lambda^2}^-.$$

From Theorem 4.13, we then establish that

$$\gamma^\lambda(r) \longrightarrow \gamma \delta(r), \quad \xi_t^{\lambda,-} \longrightarrow b_t^- \quad \text{in the sense of correlators,}$$

with $\gamma = \langle \chi, \delta(\omega) \chi \rangle$. Consequently, we obtain

$$dc_t^- = -(i\omega_0 + \gamma) c_t^- dt - i dB_t^-, \quad B_t^- = \int_0^t b_s^- ds.$$

$$dU_t = -i(c^- dB_t^+ U_t + c^- U_t dB_t^-) - \gamma_- c^+ c^- U_t dt$$

$$\begin{aligned} d(U_t^\dagger c^- U_t) &= i(c^- dB_t^+ U_t + c^+ U_t dB_t^-)^\dagger c^- U_t - \bar{\gamma}_- U_t^\dagger c^- c^+ c^- U_t dt \\ &\quad - i U_t^\dagger c^-(c^- dB_t^+ U_t + c^+ U_t dB_t^-) - \gamma_- U_t^\dagger c^- c^+ c^- U_t dt \\ &= i U_t^\dagger dB_t^- c^+ c^- U_t + i dB_t^+ U_t^\dagger c^- c^- U_t - \bar{\gamma}_- U_t^\dagger c^- c^+ c^- U_t dt \end{aligned}$$