

# Additions and Corrections

Plasmon Resonance Measurements of the Adsorption and Adsorption Kinetics of a Biopolymer onto Gold Nanocolloids

Dirk Eck, Christiane A. Helm,\* Norman J. Wagner, and Abraham Vaynberg *Langmuir* **2001**, *17*, 957–960.

Equation 2 given in Eck et al.<sup>1</sup> is erroneous. Below are given the correct equations as well as some simulated extinction spectra.

The polarizability  $\alpha$  of gold nanocolloids (radius  $R$ , number density  $N$ ) with a shell characterized by a thickness  $d$  and a dielectric constant  $\epsilon_s$  immersed in a medium with dielectric constant  $\epsilon_m$  is given by<sup>2,3</sup>

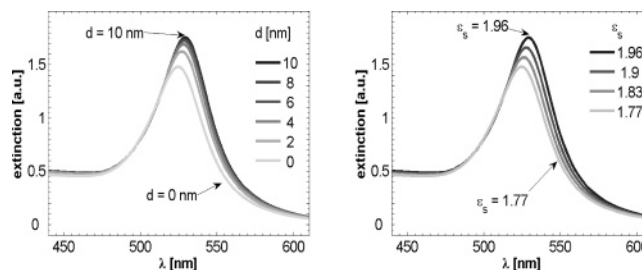
$$\alpha = \frac{\epsilon_0 (\epsilon_s - \epsilon_m)(\epsilon + 2\epsilon_s) + \left(\frac{R}{R+d}\right)^3 (\epsilon - \epsilon_s)(\epsilon_m + 2\epsilon_s)}{(\epsilon_s + 2\epsilon_m)(\epsilon + 2\epsilon_s) + 2\left(\frac{R}{R+d}\right)^3 (\epsilon - \epsilon_s)(2\epsilon_s - 2\epsilon_m)} \times \frac{4}{3}\pi(R+d)^3 \quad (1)$$

$\epsilon$  is the dielectric constant of gold, which is a complex number with a pronounced dependence on the wavelength.<sup>4</sup> Then, one obtains for the effective dielectric constant  $\epsilon_{\text{eff}}$  of the core–shell system in the medium

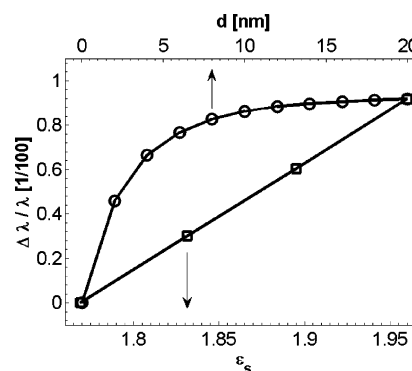
$$\epsilon_{\text{eff}} = \epsilon_m \frac{1 + 2N\alpha/(3\epsilon_0\epsilon_m)}{1 - N\alpha/(3\epsilon_0\epsilon_m)} \quad (2)$$

The experimentally accessible variable is the extinction coefficient  $\gamma = 4\pi\kappa/\lambda$ , with  $\kappa$  as the imaginary part of  $\epsilon_{\text{eff}}$ .<sup>5</sup>

To verify the equations, extinction spectra for gold core–shell particles ( $R = 7.5$  nm) are calculated for two cases (Figure 1) describing the shell: constant dielectric constant  $\epsilon_s$  with increasing thickness  $d$  and constant thickness  $d$  with increasing dielectric constant  $\epsilon_s$ . In both cases, an increase in the maximum extinction as well as a shift of the maximum position to longer wavelengths is observed. Whereas in the former case both effects show saturation for  $d \approx 15$  nm (Figure 2), limiting the sensitivity of the method to shells with 8 nm thickness, in the latter case both the peak maximum and red shift increase linearly with the dielectric constant,  $\epsilon_{\text{eff}}$ , independent of  $d$ . If  $d$  approaches zero, then the equation is that of a naked colloid with radius  $R$ . However,



**Figure 1.** Simulated extinction spectra vs wavelength for gold nanoparticles (radius  $R = 7.5$  nm) with adsorbed polymers in water. Two cases are shown: (left) varying thickness  $d$  with unchanging dielectric constant ( $\epsilon_s = 1.96$ ) and (right) constant thickness ( $d = 10$  nm) with varying dielectric constant  $\epsilon_s$  (right). The simulation parameters are the radius of the nanoparticles,  $R = 7.5$  nm; the dielectric constant of water,  $\epsilon_m = 1.77$ ; the number density  $N = 5.66 \times 10^{19} \text{ m}^{-3} = 56 \mu\text{m}^{-3} = 56 \text{ fL}^{-1}$  (leading to a filling factor of  $10^{-4}$ ).



**Figure 2.** Maximum positions of the extinction spectra shown in Figure 1 as a function of shell thickness  $d$ . Open spheres indicate a fixed-shell dielectric constant of  $\epsilon_s = 1.96$  while the thickness is varied. Open squares symbolize constant thickness  $d = 20$  nm and increasing dielectric constant  $\epsilon_s$ .

if  $d$  exceeds 20 nm, then one obtains the spectrum of a naked colloid immersed in a medium with  $\epsilon_s$ . Finally, we mention that with an increase in the nanoparticle radius a shell thickness exceeding 15 nm can be readily distinguished.

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