

Large Area Growth and Electrical Properties of p-Type WSe₂ Atomic Layers

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S Supporting Information

ABSTRACT: Transition metal dichalcogenides represent a unique class of two-dimensional layered materials that can be exfoliated into single or few atomic layers. Tungsten diselenide (WSe₂) is one typical example with p-type semiconductor characteristics. Bulk WSe₂ has an indirect band gap (~1.2 eV), which transits into a direct band gap (~1.65 eV) in monolayers. Monolayer WSe₂, therefore, is of considerable interest as a new electronic material for functional electronics and optoelectronics. However, the controllable synthesis of large-area WSe₂ atomic layers remains a challenge. The studies on WSe₂ are largely limited by relatively small lateral size of exfoliated flakes and poor yield, which has significantly restricted the large-scale applications of the WSe₂ atomic layers. Here, we report a systematic study of chemical vapor deposition approach for large area growth of atomically thin WSe₂ film with the lateral dimensions up to ~1 cm². Microphotoluminescence mapping indicates distinct layer dependent efficiency. The monolayer area exhibits much stronger light emission than bilayer or multilayers, consistent with the expected transition to direct band gap in the monolayer limit. The transmission electron microscopy studies demonstrate excellent crystalline quality of the atomically thin WSe₂. Electrical transport studies further show that the p-type WSe₂ field-effect transistors exhibit excellent electronic characteristics with effective hole carrier mobility up to 100 cm² V⁻¹ s⁻¹ for monolayer and up to 350 cm² V⁻¹ s⁻¹ for few-layer materials at room temperature, comparable or well above that of previously reported mobility values for the synthetic WSe₂ and comparable to the best exfoliated materials.

KEYWORDS: layered materials, semiconductor, tungsten diselenide, field effect transistor, large area growth



Graphene has attracted considerable interest for applications in diverse electronic and optoelectronic devices due to its unique electronic properties and atomically thin geometry.^{1–12} However, the gapless band structure limits the potential of graphene for digital electronic devices.^{2,4} It has been shown the energy band structure of the transition metal dichalcogenides (TMD) materials exhibit a unique indirect-to-direct band gap transition as their layer number is reduced to one.^{13–18} For example, bulk WSe₂ is a p-type semiconductor with an indirect band gap of ~1.2 eV, whereas its monolayer exhibits a direct band gap of ~1.65 eV.^{17,19} The direct band gap of atomically thin TMDs can offer exciting opportunities for potential applications in both digital electronic and optoelectronic devices.^{20–24} For example, it has been recently reported that exfoliated monolayer WSe₂ can be used to create a high performance p-type field-effect transistor (FET).²⁵ However, the size of the monolayer materials obtained from mechanical exfoliation method is limited in a few to a few tens of micrometers. It has also been recently reported that a chemical vapor deposition approach can be used to grow WSe₂ atomic flakes, but only in separated domains with the domain size on the order of 10 μm.^{26–28} Therefore, the preparation of large-

area monolayer WSe₂ film is essential for practical applications yet remains a significant challenge.

Here, we report a systematic study of chemical vapor deposition approach for the preparation of large area atomically thin WSe₂ films directly on SiO₂/Si substrates with areal size of monolayer WSe₂ film up to 1 cm². Microphotoluminescence mapping demonstrates distinct layer-number dependent photoluminescence, with the monolayer area exhibit much stronger emission than bilayer or multilayers. The transmission electron microscopy (TEM) studies reveal excellent crystalline quality of the atomically thin WSe₂ and electrical transport studies further demonstrate that the p-type WSe₂ field-effect transistors exhibit excellent electronic characteristics with hole carrier mobility up to 100 cm² V⁻¹ s⁻¹ for monolayer and up to 350 cm² V⁻¹ s⁻¹ for few-layer materials, comparable or well above that of previously reported mobility values or the synthetic WSe₂ and comparable to the best exfoliated materials.^{26,29–31}

The growth processes are performed in a homemade tube furnace with detailed growth procedures described in the

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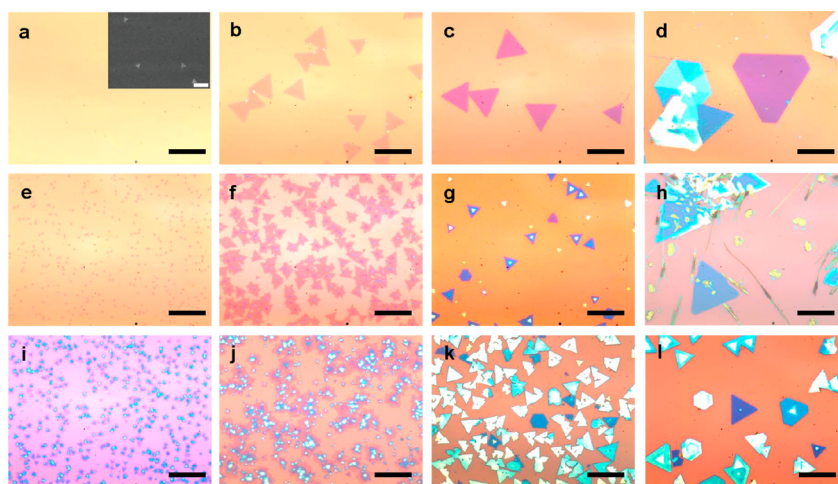


Figure 1. Reactor-conditions-dependent WSe₂ growth with varying temperature and flow rate of argon. Optical microscope images of WSe₂ samples grown at different temperature under designed flow rate for 20 min: (a) 750 °C, 100 sccm (inset: SEM image of same size sample, indicating high density but small nucleations (~300 to 500 nm); scale bar, 5 μm); (b) 765 °C, 100 sccm; (c) 780 °C, 100 sccm; (d) 795 °C, 100 sccm; (e) 750 °C, 150 sccm; (f) 765 °C, 150 sccm; (g) 780 °C, 150 sccm; (h) 795 °C, 150 sccm; (i) 750 °C, 200 sccm; (j) 765 °C, 200 sccm; (k) 780 °C, 200 sccm; (l) 795 °C, 200 sccm. All of the scale bars are 20 μm.

Experimental section in Supporting Information. Briefly, the WSe₂ powders are placed in an alumina boat at the center of a quartz tube inside a one-inch tube furnace, and the clean SiO₂/Si substrates are used as the growth substrate for the deposition of WSe₂ atomic layers at the downstream end with variable substrate temperature. Argon is continuously supplied through the reactor with designed flow rate as the carrier gas. We have first conducted systematic studies to investigate the effect of substrate temperature and flow rate of carrier gas with a constant source temperature at 1060 °C (Figure 1). Most of the WSe₂ domains exhibit a triangular shape as shown by the optical microscope images in Figure 1b, c and Figure 2a–c, and some of them show a hexagonal shape as shown in Figure 1d, g,

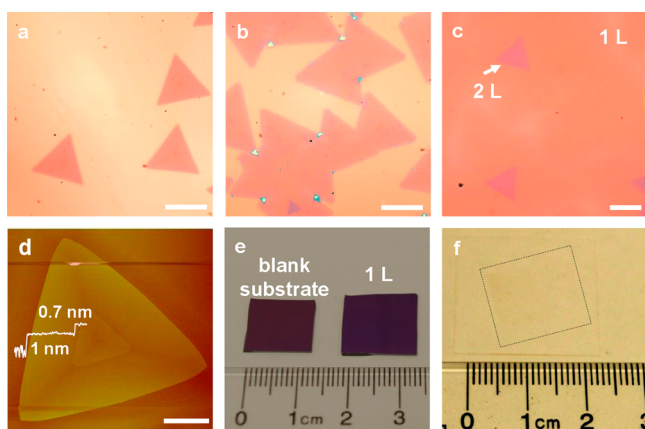


Figure 2. Time dependent growth of large area monolayer WSe₂ single crystal domains. (a) Typical optical microscope images of the monolayer WSe₂ domains taken after 20 min of growth, (b) 30 min of growth, and (c) 40 min of growth. (d) Atomic force microscope image of a WSe₂ monolayer and bilayer domains and their line scan profile. (e) Photograph of the ~1 cm sized monolayer WSe₂ films obtained with 40 min growth time. Left: blank 300 nm SiO₂/Si substrate. Right: fully covered WSe₂ monolayer on blank 300 nm SiO₂/Si substrate. (f) Photograph of the ~1 cm sized monolayer WSe₂ films transferred onto a glass template by chemical etching of SiO₂ layers. Scale bars of panels a–d are all 5 μm.

and l. In general, with a fixed source temperature, the lower substrate temperature typically results in a higher nucleation density, and a higher flow rate of the carrier gas produces a similar effect due to higher supersaturation of the precursors over the substrate surface. On the other hand, the higher substrate temperature promotes the nucleation of the extra atomic layers to produce multilayers.

At lower substrate temperature of 750 °C with low flow rate of 100 sccm, there are no visible WSe₂ domains observed on the SiO₂/Si substrates under optical microscope. However, after analyzing the same sample by using scanning electron microscopy (SEM), very small nucleations (~300 nm of average edge size) were observed with nucleation density of ~4200/mm². By increasing the substrate temperature to 765 °C, the monolayer WSe₂ domains (~10 μm of average edge size) start to appear with a nucleation density of ~1060/mm² (Figure 1b). When the growth temperature reaches 780 °C, the bilayer WSe₂ domains (~20 μm of average edge size) start to appear with an overall lower domain density reduced to ~350/mm² (Figure 1c). As the growth temperature is further increased to 795 °C, even thicker and larger flakes (~50 μm of average edge size) appear with even lower domain density of ~280/mm² (Figure 1d). Overall, higher growth temperature yields a lower density of thicker WSe₂ domains with a larger domain size. By increasing the flow rate of the carrier gas, both the layer number and the domain density of the WSe₂ are increased. For instance, at growth temperature of 765 °C, the nucleation density is ~1060/mm² under 100 sccm flow rate of carrier gas, which increases to ~11 000/mm² under 150 sccm, and ~23 000/mm² under 200 sccm. All of our observations are consistent with the nucleation model of the vapor phase deposition developed by W. K. Burton and N. Cabrera, where they predict that the nucleation probability is proportional to the supersaturation and inversely proportional to the substrate temperature.³²

Based on the above studies, we have identified an optimized condition for the large area growth of WSe₂ monolayers at a substrate temperature of 765 °C and 100 sccm of carrier gas. Figure 2a shows the OM images of typical triangular WSe₂ monolayers after 20 min growth with an average edge length of

approximately 5 μm . With the continued growth, the monolayer domains starts to merge together at 30 min, with bilayer and few-layer domains occasionally appeared on their first layer (Figure 2b). After 40 min growth, the WSe₂ monolayer domains are completely merged together to form a continuous monolayer, with a few second layer triangular domains seen with different optical contrast (Figure 2c). In this case, the SiO₂/Si substrates are completely covered by the monolayer WSe₂ domains with less than 5% areal coverage of the bilayer. The number of WSe₂ atomic layers is also determined by AFM measurements (Figure 2d). The AFM step height of WSe₂ monolayer on the SiO₂ substrate is typically measured between 0.7–1.0 nm, and the step heights of the second layer on monolayer is around ~ 0.7 nm (line scan in Figure 2d), which is consistent with the published reports of exfoliated WSe₂.^{14,15} The larger step height observed in the first layer is commonly seen in graphene or other layered materials such as MoSe₂.¹⁶ The lateral size of the resulted continuous monolayer WSe₂ film is as large as 1 cm² and is only limited by the size of the furnace (Figure 2e). The large area monolayer WSe₂ film can be readily transferred onto another substrate such as the glass (Figure 2f). We have further investigated the layer number dependent morphological and optical properties of the resulting atomically thin WSe₂ films. To this end, we have employed the μ -Raman to study the monolayer/bilayer WSe₂ domains. The bilayer WSe₂ domain shows a triangular shape, which prefers to crystallize from the center of the first layer. The layer number of the WSe₂ domains is further confirmed using μ -Raman studies. The Raman spectrum of monolayer region shows a single peak at 252 cm⁻¹ (green line in Figure 3b), corresponding to the A_{1g} resonance mode of WSe₂; and that of bilayer area shows an additional small peak at 307 cm⁻¹ (red line in Figure 3b), corresponding to the B_{2g} resonance mode of WSe₂. In general, the B_{2g} signature mode is only active on the bilayer or few-layer region, which could reflect the presence of the additional interlayer interaction.^{14,15}

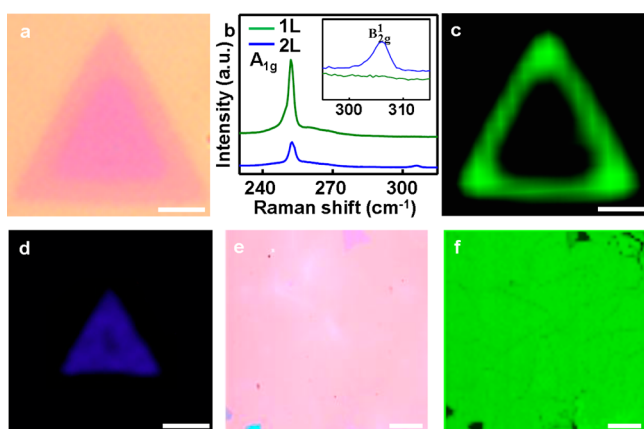


Figure 3. Micro-Raman investigation of the monodomain WSe₂. (a) Optical microscope image of a typical monolayer WSe₂ with bilayer domain on the center. (b) The Raman spectra of the monolayer and bilayer WSe₂. (c) and (d) Typical Raman map of the single domain monolayer (c, center wavenumber: ~ 252 cm⁻¹) with bilayer (d, center wavenumber: ~ 307 cm⁻¹) WSe₂ domain. (e) Optical microscope image of large scale monolayer WSe₂ with several bilayer or few-layer domain. (f) Raman map of the large scale monolayer (center wavenumber: ~ 250 cm⁻¹) WSe₂ with several bilayer and few-layer domain. Scale bars of panels a, c, and d are all 5 μm ; Scale bars of panels e and f are 10 μm .

In contrast, the Raman A_{1g} mode of WSe₂ is less sensitive to layer thickness, only with the intensity increasing with reducing atomic layers. The corresponding Raman map recorded at A_{1g} mode can be further used to determine the layer number of the WSe₂ (Figure 3c–f). The darker triangular region on the Raman map of A_{1g} mode corresponds to the bilayer or few-layer WSe₂ domain. The Raman mapping shows nearly the same color contrast throughout the monolayer and bilayer or few-layer area, indicating a uniform crystal quality.

The optical properties of the monolayer WSe₂ domains and the continuous WSe₂ sheets were further investigated using microphotoluminescence (μ -PL). The WSe₂ monodomain with single layer shows the PL peak located at approximately 767 nm, with full-width-half-maximum (FWHM) values of 25 nm (Figure 4b), which is comparable to the published data for

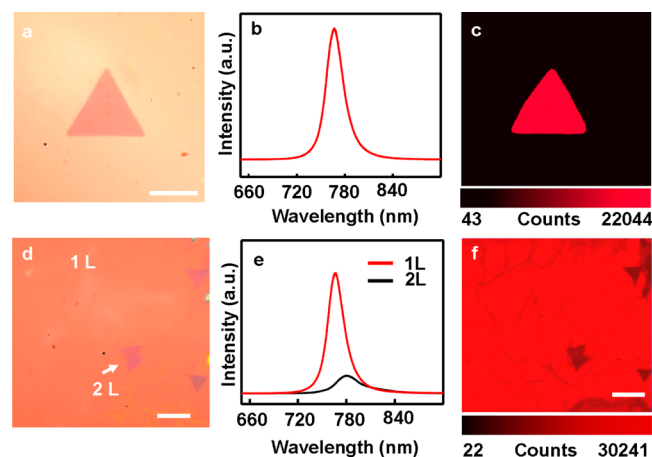


Figure 4. Optical properties of the single domain and fully covered WSe₂ monolayers. (a) Optical microscope image of a typical monolayer WSe₂. Scale bar is 5 μm . (b) PL spectra of the monolayer WSe₂ and PL map (c) of WSe₂ single domain showed in panel a. (d) Optical microscope image of the full covered monolayer WSe₂ with some bilayer domains. Scale bar is 10 μm . (e) PL spectra of the monolayer and bilayer WSe₂. (f) Corresponding PL map of the same region taken at panel d. Scale bar is 10 μm .

exfoliated WSe₂.^{14,15} The corresponding μ -PL map of the monolayer WSe₂ single crystal domain shows a very uniform contrast (Figure 4c), indicating the high crystalline quality and uniformity of the as-grown WSe₂ atomic layers. We have also conducted μ -PL studies on the continuous WSe₂ film that are merged together (Figure 4d). The μ -PL spectra taken of the monolayer region show the characteristic peak around 766 nm (red curve in Figure 4e). The strong light emission from monolayer indicates the high quality of the continuous WSe₂ films. The μ -PL spectra taken at bilayer WSe₂ region (the darker triangular area on the OM image in Figure 4d) shows a wider peak at ~ 790 nm with significantly lower intensity (black curve in Figure 4e). We have further conducted μ -PL mapping studies over a relatively large area ($\sim 60 \times 60$ μm) to evaluate the overall quality of the material (Figure 4f). Importantly, the PL mapping shows rather uniform emission across the entire film, indicating that the film is largely consisted of monolayer crystals. There are a few darker triangles in the PL mapping image, corresponding to the bilayer region. Additionally, it is also noted that there are some slightly darker lines, which can be attributed to the grain boundaries between merged monolayer domains.

To further evaluate the atomic structure of the WSe₂ atomic layers, we have performed high resolution transmission electron microscopy (HRTEM) and electron diffraction (ED) studies. To this end, the as-grown WSe₂ crystals were transferred onto a carbon-coated TEM grid. Figure 5a is the optical microscopy

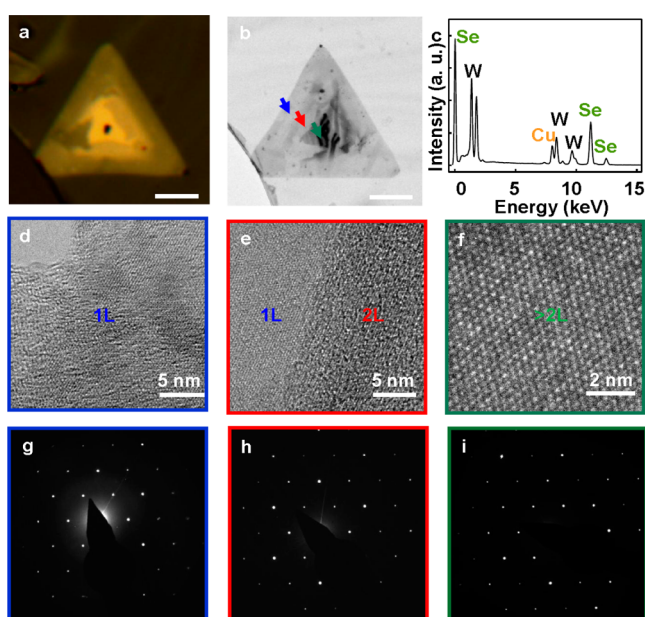


Figure 5. Crystalline structure characterization. (a) Optical microscope image of a typical monolayer WSe₂ with domain of few layers on the center. Scale bar is 10 μm . (b) Low magnification bright field TEM image of a typical transferred WSe₂ triangular crystal. (c) Energy dispersive X-ray spectra showing the ratio between W and Se obtained from the integrated peak area is approximately 1:2. (d–f) HRTEM images of the atomic structure of the WSe₂ monolayer (d), bilayer (e), and multilayers. These images were taken on the regions indicated on panel b by the arrows with blue, red, and dark cyan colors, respectively. Scale bars are 5 nm (d,e) and 2 nm (f). (g–i) SAED pattern of the monolayer (g), bilayer (h) and few-layer (i) WSe₂ with the zone axis of [0001], the 6-fold symmetry in the position of the diffraction spots demonstrates that the triangular monolayer is predominantly single crystal with hexagonal structures.

image. Figure 5b shows a low magnification TEM image of a typical transferred WSe₂ triangular crystal, which shows increasing contrast from the edge to the center, with three distinct regions corresponding to the monolayer, bilayer, and few-layer area. The energy dispersive X-ray studies demonstrate that the atomic ratio between W and Se is approximately 1:2 (Figure 5c), consistent with the expected stoichiometry. We have further conducted HRTEM on different regions of the same WSe₂ flake. Figure 5d–f show the HRTEM images of the atomic structure of the WSe₂ monolayer, bilayer, and multilayers, respectively. These images were taken on the regions indicated on Figure 5b by the arrows with blue, red, and dark cyan colors, respectively. The hexagonal lattice is clearly visible from each atomic resolution image, confirming the excellent crystalline quality of the atomic layered material. The selected-area electron diffraction (SAED) is used to characterize the crystal structure of the film. Figure 5g–i show the SAED patterns of the monolayer, bilayer, and few-layer WSe₂ with the zone axis of [0001]. The single set of diffraction spots with 6-fold symmetry demonstrates that the triangular monolayer is single crystals with hexagonal structures.^{33,34}

We have further evaluated the electronic properties of the WSe₂ atomic layers. To this end, we have fabricated back-gated field-effect transistors (FETs) from the synthetic WSe₂ on the 300 nm SiO₂/Si substrates (Figure 6a). The source and drain

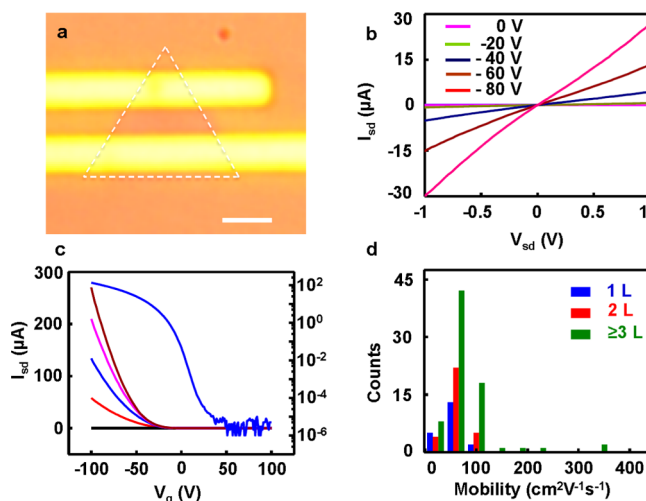


Figure 6. Electronic properties of WSe₂ atomic layers. (a) Optical microscope image of a monolayer WSe₂ transistor; scale bar is 2 μm . (b) $I_{\text{sd}}-V_{\text{sd}}$ output characteristics of the WSe₂ transistor shown in panel a. (c) $I_{\text{sd}}-V_{\text{g}}$ transfer characteristics of the device shown in (a) at $V_{\text{sd}} = 0, 1, 2, 3,$ and 4 V. (d) Summary of the mobility values obtained in WSe₂ field effect transistors with different number of atomic layers, demonstrating a mobility range of $10\text{--}350\text{ cm}^2\text{ V}^{-1}\text{ s}^{-1}$ and maximum on/off ratio over 10^8 .

electrodes (100 nm Au) defined by the electron beam lithography (EBL) and deposited using thermal evaporation. More than 100 transistors were fabricated on WSe₂ monolayer, bilayer, and few-layer monodomains. The standard transistor measurements were conducted under ambient conditions to obtain the on-currents, carrier mobilities, and on/off ratios for all devices. Figure 6b shows the $I_{\text{ds}}-V_{\text{ds}}$ output characteristic of a monolayer WSe₂ FETs at various gate voltages. The linear and symmetric curves suggested that ohmic contacts were formed at the source and drain electrodes. The $I_{\text{ds}}-V_{\text{bg}}$ curve of the same device was measured at different drain bias from 0 to 4 V, with the back-gate voltage sweeping from -100 to 100 V (Figure 6c). Additionally, the $I_{\text{ds}}-V_{\text{bg}}$ curve of the typical WSe₂ monolayer, bilayer and few-layer FETs are shown in Supporting Information Figure S2. All these IV curves show typical p-type semiconductor characteristic, consistent with mechanically exfoliated WSe₂ materials. The highest mobility of the monolayer WSe₂ FETs is $100\text{ cm}^2\text{ V}^{-1}\text{ s}^{-1}$. The measured on/off ratio of this device reaches a maximum value above 1×10^8 for gate voltages swept in the range from -100 to 100 V with a source-drain bias of 2 V, which is comparable to the best reported values for exfoliated samples, and greatly higher than synthetic samples reported previously. After analysis of all the fabricated WSe₂ FETs, we have plotted a histogram to show the mobility distribution for all devices (Figure 6d). Importantly, the highest mobility of the few-layer WSe₂ FETs can reach up to $350\text{ cm}^2\text{ V}^{-1}\text{ s}^{-1}$, which represents the highest value observed in these atomically thin TMDs at room temperature,^{26,35–40} further highlighting the high crystalline quality of these as-grown materials.

In summary, we have reported a systematic investigation on the chemical vapor deposition growth of large area p-type WSe₂

atomic layers. By systematically tuning the synthetic parameters, we show that high quality monolayer WSe₂ thin film can be prepared over large area with the lateral dimensions up to ~1 cm². Microphotoluminescence studies demonstrate uniform PL emission from the monolayer materials and the transmission electron microscopy studies confirm the excellent crystalline quality. Electrical transport studies further demonstrate the p-type WSe₂ back-gated field-effect transistors exhibit excellent electronic characteristics with hole mobility over 350 cm² V⁻¹ s⁻¹ and on/off ratio above 10⁸. The availability of high quality large area p-type atomically thin materials offers an important building block for the design of future van de Waals heterostructures for electronic and photonic devices.

■ ASSOCIATED CONTENT

Supporting Information

Experimental details; Figures S1–S2. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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Notes

The authors declare no competing financial interest.

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■ REFERENCES

- (1) Novoselov, K. S.; Fal'ko, V. I.; Colombo, L.; Gellert, P. R.; Schwab, M. G.; Kim, K. *Nature* **2012**, 490 (7419), 192–200.
- (2) Schwierz, F. *Nat. Nanotechnol.* **2010**, 5 (7), 487–496.
- (3) Bonaccorso, F.; Sun, Z.; Hasan, T.; Ferrari, A. C. *Nat. Photonics* **2010**, 4 (9), 611–622.
- (4) Weiss, N. O.; Zhou, H. L.; Liao, L.; Liu, Y.; Jiang, S.; Huang, Y.; Duan, X. F. *Adv. Mater.* **2012**, 24 (43), 5782–5825.
- (5) Liao, L.; Lin, Y. C.; Bao, M. Q.; Cheng, R.; Bai, J. W.; Liu, Y. A.; Qu, Y. Q.; Wang, K. L.; Huang, Y.; Duan, X. F. *Nature* **2010**, 467 (7313), 305–308.
- (6) Bao, Q. L.; Loh, K. P. *ACS Nano* **2012**, 6 (5), 3677–3694.
- (7) Lin, Y. M.; Valdes-Garcia, A.; Han, S. J.; Farmer, D. B.; Meric, I.; Sun, Y. N.; Wu, Y. Q.; Dimitrakopoulos, C.; Grill, A.; Avouris, P.; Jenkins, K. A. *Science* **2011**, 332 (6035), 1294–1297.
- (8) Britnell, L.; Gorbachev, R. V.; Jalil, R.; Belle, B. D.; Schedin, F.; Mishchenko, A.; Georgiou, T.; Katsnelson, M. I.; Eaves, L.; Morozov, S. V.; Peres, N. M. R.; Leist, J.; Geim, A. K.; Novoselov, K. S.; Ponomarenko, L. A. *Science* **2012**, 335 (6071), 947–950.
- (9) Yang, H.; Heo, J.; Park, S.; Song, H. J.; Seo, D. H.; Byun, K. E.; Kim, P.; Yoo, I.; Chung, H. J.; Kim, K. *Science* **2012**, 336 (6085), 1140–1143.
- (10) Xia, F. N.; Mueller, T.; Lin, Y. M.; Valdes-Garcia, A.; Avouris, P. *Nat. Nanotechnol.* **2009**, 4 (12), 839–843.
- (11) Liu, Y.; Cheng, R.; Liao, L.; Zhou, H. L.; Bai, J. W.; Liu, G.; Liu, L. X.; Huang, Y.; Duan, X. F. *Nat. Commun.* **2011**, 2, 579.
- (12) Mueller, T.; Xia, F. N. A.; Avouris, P. *Nat. Photonics* **2010**, 4 (5), 297–301.
- (13) Splendiani, A.; Sun, L.; Zhang, Y. B.; Li, T. S.; Kim, J.; Chim, C. Y.; Galli, G.; Wang, F. *Nano Lett.* **2010**, 10 (4), 1271–1275.
- (14) Li, H.; Lu, G.; Wang, Y. L.; Yin, Z. Y.; Cong, C. X.; He, Q. Y.; Wang, L.; Ding, F.; Yu, T.; Zhang, H. *Small* **2013**, 9 (11), 1974–1981.
- (15) Tonndorf, P.; Schmidt, R.; Bottger, P.; Zhang, X.; Borner, J.; Liebig, A.; Albrecht, M.; Kloc, C.; Gordan, O.; Zahn, D. R. T.; de Vasconcellos, S. M.; Bratschitsch, R. *Opt. Express* **2013**, 21 (4), 4908–4916.
- (16) Xu, Y. X.; Huang, X. Q.; Lin, Z. Y.; Zhong, X.; Huang, Y.; Duan, X. F. *Nano Res.* **2013**, 6 (1), 65–76.
- (17) Terrones, H.; Lopez-Urias, F.; Terrones, M. *Sci. Rep.* **2013**, 3, 1549.
- (18) Lee, Y. H.; Yu, L. L.; Wang, H.; Fang, W. J.; Ling, X.; Shi, Y. M.; Lin, C. T.; Huang, J. K.; Chang, M. T.; Chang, C. S.; Dresselhaus, M.; Palacios, T.; Li, L. J.; Kong, J. *Nano Lett.* **2013**, 13 (4), 1852–1857.
- (19) Kumar, A.; Ahluwalia, P. K. *Eur. Phys. J. B* **2012**, 85, (6).
- (20) Najmaei, S.; Liu, Z.; Zhou, W.; Zou, X. L.; Shi, G.; Lei, S. D.; Yakobson, B. L.; Idrobo, J. C.; Ajayan, P. M.; Lou, J. *Nat. Mater.* **2013**, 12 (8), 754–759.
- (21) Wang, Q. H.; Kalantar-Zadeh, K.; Kis, A.; Coleman, J. N.; Strano, M. S. *Nat. Nanotechnol.* **2012**, 7 (11), 699–712.
- (22) Yu, W. J.; Liu, Y.; Zhou, H. L.; Yin, A. X.; Li, Z.; Huang, Y.; Duan, X. F. *Nat. Nanotechnol.* **2013**, 8 (12), 952–958.
- (23) Yu, W. J.; Li, Z.; Zhou, H. L.; Chen, Y.; Wang, Y.; Huang, Y.; Duan, X. F. *Nat. Mater.* **2013**, 12 (3), 246–252.
- (24) Zhu, C. F.; Zeng, Z. Y.; Li, H.; Li, F.; Fan, C. H.; Zhang, H. J. *Am. Chem. Soc.* **2013**, 135 (16), 5998–6001.
- (25) Fang, H.; Chuang, S.; Chang, T. C.; Takei, K.; Takahashi, T.; Javey, A. *Nano Lett.* **2012**, 12 (7), 3788–3792.
- (26) Huang, J. K.; Pu, J.; Hsu, C. L.; Chiu, M. H.; Juang, Z. Y.; Chang, Y. H.; Chang, W. H.; Iwasa, Y.; Takenobu, T.; Li, L. J. *ACS Nano* **2014**, 8 (1), 923–930.
- (27) Cheng, R.; Li, D.; Zhou, H.; Wang, C.; Yin, A.; Jiang, S.; Liu, Y.; Chen, Y.; Huang, Y.; Duan, X. *Nano Lett.* **2014**, 14 (10), 5590–5597.
- (28) Duan, X.; Wang, C.; Shaw, J. C.; Cheng, R.; Chen, Y.; Li, H.; Wu, X.; Tang, Y.; Zhang, Q.; Pan, A.; Jiang, J.; Yu, R.; Huang, Y. *Nat. Nanotechnol.* **2014**, No. 10.1038/nnano.2014.222.
- (29) Fang, H.; Tosun, M.; Seol, G.; Chang, T. C.; Takei, K.; Guo, J.; Javey, A. *Nano Lett.* **2013**, 13 (5), 1991–1995.
- (30) Chuang, H. J.; Tan, X. B.; Ghimire, N. J.; Perera, M. M.; Chamlagain, B.; Cheng, M. M. C.; Yan, J. Q.; Mandrus, D.; Tomanek, D.; Zhou, Z. X. *Nano Lett.* **2014**, 14 (6), 3594–3601.
- (31) Liu, W.; Kang, J. H.; Sarkar, D.; Khatami, Y.; Jena, D.; Banerjee, K. *Nano Lett.* **2013**, 13 (5), 1983–1990.
- (32) Burton, W. K.; Cabrera, N. *Discuss. Faraday Soc.* **1949**, No. 5, 33–39.
- (33) Liu, L. X.; Zhou, H. L.; Cheng, R.; Yu, W. J.; Liu, Y.; Chen, Y.; Shaw, J.; Zhong, X.; Huang, Y.; Duan, X. F. *ACS Nano* **2012**, 6 (9), 8241–8249.
- (34) Zhou, H. L.; Yu, W. J.; Liu, L. X.; Cheng, R.; Chen, Y.; Huang, X. Q.; Liu, Y.; Wang, Y.; Huang, Y.; Duan, X. F. *Nat. Commun.* **2013**, 4, No. 10.1038/ncomms3096.
- (35) Liu, L. T.; Kumar, S. B.; Ouyang, Y.; Guo, J. *IEEE Trans. Electron Devices* **2011**, 58 (9), 3042–3047.
- (36) Lee, Y. H.; Zhang, X. Q.; Zhang, W. J.; Chang, M. T.; Lin, C. T.; Chang, K. D.; Yu, Y. C.; Wang, J. T. W.; Chang, C. S.; Li, L. J.; Lin, T. W. *Adv. Mater.* **2012**, 24 (17), 2320–2325.
- (37) Huang, X.; Zeng, Z. Y.; Zhang, H. *Chem. Soc. Rev.* **2013**, 42 (5), 1934–1946.
- (38) Zhan, Y. J.; Liu, Z.; Najmaei, S.; Ajayan, P. M.; Lou, J. *Small* **2012**, 8 (7), 966–971.
- (39) Liu, K. K.; Zhang, W. J.; Lee, Y. H.; Lin, Y. C.; Chang, M. T.; Su, C.; Chang, C. S.; Li, H.; Shi, Y. M.; Zhang, H.; Lai, C. S.; Li, L. J. *Nano Lett.* **2012**, 12 (3), 1538–1544.
- (40) Larentis, S.; Fallahazad, B.; Tutuc, E. *Appl. Phys. Lett.* **2012**, 101 (22), 223104.