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# Is the prediction of $pK_a$ values by constant-pH molecular dynamics being hindered by inherited problems?

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## ABSTRACT

In this study, we investigate two factors that can hinder the performance of constant-pH molecular dynamics methods in predicting protein  $pK_a$  values, using hen egg white lysozyme as a test system. The first factor is related to the molecular definition and  $pK_a$  value of model compounds in the Poisson-Boltzmann framework. We address this by defining the model compound as a molecular fragment with an associated  $pK_a$  value that is calibrated against experimental data, which results in a decrease of 0.12 units in  $pK_a$  errors. The second addressed factor is the possibility that detrimental structural distortions are being introduced in the simulations by the underlying molecular mechanics force field. This issue is investigated by analyzing how the gradual structural rearrangements affect the predicted  $pK_a$  values. The two GROMOS force fields studied here (43A1 and 53A6) yield good  $pK_a$  predictions, although a time-dependent performance is observed: 43A1 performs better after a few nanoseconds of structural reorganization ( $pK_a$  errors of  $\sim 0.45$ ), while 53A6 gives the best prediction right at the first nanosecond ( $pK_a$  errors of 0.42). These results suggest that the good performance of constant-pH molecular dynamics methods could be further improved if these force field limitations were overcome.

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**Key words:** pentapeptides; hen egg white lysozyme; simulation; Poisson-Boltzmann; Monte Carlo; GROMOS; 43A1; 53A6.

## INTRODUCTION

The importance of pH in the structure and function of proteins is well illustrated by the fact that about 25% of its residues contain ionizable side chains.<sup>1</sup> The ionizations that arise from changes in pH can generate strong electrostatic interactions, which will inevitably have a direct influence on molecular structure and binding. It is well accepted that the pH effect on proteins is of complex nature. This is mainly due to the multiplicity of titrable sites, which are not only subjected to different environments but also coupled to one another in complex ways.

As described in the first article of this special issue, the development and improvement of methods to model protonation equilibrium and to predict  $pK_a$  values in proteins has been a continued effort by many research groups over the years but problems still remain. One of the major factors affecting the modeling of protonation processes is certainly the coupling between protonation and conformation, which is explicitly addressed by constant-pH molecular dynamics (MD) methods.<sup>2–25</sup> Still, constant-pH MD methods necessarily inherit the problems of the underlying molecular mechanics (MM)/MD simulations concerning force field accuracy and sampling efficiency. Furthermore, several of these methods rely on Poisson-Boltzmann (PB) or generalized Born methods to compute the protonation free energies, and thus may inherit also some of the parametrization problems of these methods, concerning dielectric constants, charges, radii, model compound  $pK_a$ s, etc.

The stochastic titration method<sup>4–9</sup> developed at our group to perform constant-pH MD simulations tries to derive all parameters required for PB calculations from the underlying MM/MD force field. Thus, atomic partial charges are the MM/MD ones, radii are directly computed from the Lennard-Jones interactions with the water oxygen, and the dielectric constant of the “instantly-frozen” protein is taken as 2 to account only for electronic polarization. Nonetheless, its treatment of model compounds retains some of the theoretical vagueness present in PB methods with respect to their molecular definition and  $pK_a$  value. Thus, the first

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aim of the present study is to provide a clear procedure to derive model compound  $pK_a$  values directly from experimental data, following a rationale similar to the one adopted in the parametrization of MM/MD force fields.

As seen in this special issue, current constant-pH MD methods predict  $pK_a$  values with an accuracy similar to that obtained with other good methods, but recent evidences on some structural distortions induced by MM/MD force fields<sup>26–29</sup> may lead us to ask whether such problems may be actually limiting their performance. Thus, the second aim of the present study is to study the cumulative effect of conformational relaxation during the simulations, analyzing whether it consistently improves predictions or not.

Although we started this study using two of the Staphylococcal nuclease mutants from Garcia-Moreno's Lab, serious issues were found involving the parametrization of neutral Arg and Lys groups, which are currently being investigated in detail. Therefore, we decided to use hen egg white lysozyme (HEWL), which has become, over the years, a standard test system for  $pK_a$  prediction methodologies,<sup>7,14–17,22,24,30–50</sup> because of the availability of accurate experimental data, and because it has a number of sites with  $pK_a$  values that differ markedly from their reference values, mainly in the acidic range.

## THEORY AND METHODS

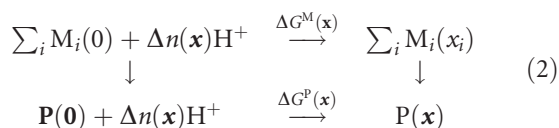
### Model compounds and their $pK_a$ values

The protonation state of a protein with  $N$  protonatable sites is represented here as  $\mathbf{x} = (x_1, x_2, \dots, x_N)$ , where  $x_i$  denotes the protonation state of site  $i$ . For concreteness, we consider a framework using a (unique) ionized reference state and proton tautomerism,<sup>31</sup> but the discussion and conclusions in this section are easily extended to other frameworks (e.g., without tautomerism). Thus,  $x_i = 0$  refers to the ionized state, while the remaining values of  $x_i$  refer to the alternative neutral tautomers with different proton positions. The protein in state  $\mathbf{x}$  is denoted as  $P(\mathbf{x})$  and its probability of occurrence follows a semigrand canonical distribution<sup>4</sup>

$$p(\mathbf{x}) = \exp\{-\ln(10)pH\Delta n(\mathbf{x}) - \Delta G^P(\mathbf{x})/kT\}\Xi^{-1}, \quad (1)$$

where  $\Delta G^P(\mathbf{x})$  is the Gibbs free energy change for the reaction leading from  $P(\mathbf{0})$  to  $P(\mathbf{x})$ ,  $\Delta n(\mathbf{x})$  is the associated change in the number of protons,  $k$  is Boltzmann's constant,  $T$  is the absolute temperature, and  $\Xi$  is the pH-dependent partition function normalizing the probabilities. This relation is exact, meaning that an estimation of the reaction free energies  $\Delta G^P(\mathbf{x})$  allows for a complete characterization of the protonation equilibrium of the protein, either from an exact or approximate (e.g., Monte Carlo) calculation.

The use of linear PB models to compute  $\Delta G^P(\mathbf{x})$  values is based on the use of model compounds. A model compound is a single-site compound containing the same chemical group as a protonatable site in the protein, and whose  $pK_a$  value is presumably known. To each protein site corresponds a model compound, making possible to write the following thermodynamic cycle<sup>33,51</sup>:



where  $M_i(x_i)$  is the model compound of site  $i$  in state  $x_i$ , and  $\Delta G^M(\mathbf{x})$  is the counterpart of  $\Delta G^P(\mathbf{x})$  for the  $N$  model compounds in solution; a negative  $\Delta n(\mathbf{x})$  corresponds to deprotonation. The fundamental assumption behind the use of model compounds is that the quantum contribution for the (de)protonation of a site in the protein is the same as in its corresponding model compound, so that only classical contributions (e.g., from a PB model) need to be considered when addressing this thermodynamic cycle. Thus, we can write

$$\begin{aligned} \Delta G^P(\mathbf{x}) &= \Delta G^M(\mathbf{x}) + U^P(\mathbf{x}) - U^P(\mathbf{0}) - \sum_i U_{ii}^M(x_i) \\ &\quad + \sum_i U_{ii}^M(0), \end{aligned} \quad (3)$$

where  $U^P(\mathbf{x})$  and  $U_{ii}^M(x_i)$  are the PB energies of, respectively,  $P(\mathbf{x})$  and  $M_i(x_i)$ . The  $N$  model compounds in solution contribute independently to  $\Delta G^M(\mathbf{x})$ , which is thus given by

$$\Delta G^M(\mathbf{x}) = \ln(10)kT \sum_i' \gamma_i pK_i^M(x_i), \quad (4)$$

where the prime indicates that the sum does not include  $i$  values for which  $x_i = 0$ ,  $\gamma_i = \pm 1$  is the charge of the ionized form of site  $i$ , and  $pK_i^M(x_i)$  is the  $pK_a$  value of the dissociation reaction involving the neutral form  $M_i(x_i)$ ,<sup>31,32</sup> related to its global  $pK_a$ ,  $pK_i^M$ , as

$$pK_i^M(x_i) = pK_i^M - \gamma_i \log f_i(x_i), \quad (5)$$

where  $f_i(x_i)$  is the fraction of tautomer  $x_i$  among all neutral tautomers of  $M_i$ . Furthermore, if we conceptually split the protein into  $N+1$  nonoverlapping fragments corresponding to the  $N$  protonatable sites plus the remaining nonprotonatable background (b), the linearity of the PB equation implies that the superposition principle<sup>52</sup> holds for these fragments, giving

$$\begin{aligned} U^P(\mathbf{x}) &= U_{bb}^P + \sum_i U_{ib}^P(x_i) + \sum_i U_{ii}^P(x_i) \\ &\quad + \sum_i \sum_{j < i} U_{ij}^P(x_i, x_j), \end{aligned} \quad (6)$$

where  $U_{uv}^P$  denotes the PB interaction term between fragments  $u$  and  $v$ . The previous equations, together with the relation  $\Delta n(\mathbf{x}) = -\sum_i \gamma_i$  give

$$p(\mathbf{x}) = \exp \left\{ \ln(10) \sum_i [\gamma_i (\text{p}K_i^M - \text{pH}) - \log f_i(x_i)] \right. \\ + \beta \sum_i [U_{ib}^P(x_i) + U_{ii}^P(x_i) - U_{ii}^M(x_i)] \\ - \beta \sum_i [U_{ib}^P(0) + U_{ii}^P(0) - U_{ii}^M(0)] \\ \left. + \beta \sum_i \sum_{j \neq i} [U_{ij}^P(x_i, x_j) - U_{ij}^P(0, 0)] \right\} \Xi^{-1} \quad (7)$$

where  $\beta = 1/kT$ , and whose terms in the second and third sums can be regarded as the effect of the protein environment on the model compound  $i$  in states  $x_i$  and 0, respectively.

The above derivation could have considered a background fragment for each model compound, which would result in additional subtractive terms  $U_{ib}^M(x_i)$  and  $U_{ib}^M(0)$  in the second and third sums, respectively. Instead, we intentionally considered each model compound to be identical to the corresponding fragment in the protein, as it is usual practice and implemented in most software tools. This procedure reduces the computational needs but the model compound becomes an unphysical “molecule” whose  $\text{p}K_a$  value cannot be measured. Moreover, most sets of  $\text{p}K_a$  values of model compounds currently in use seem to derive from the Nozaki and Tanford set,<sup>53</sup> which is a set of “typical” values inferred from various chemical species, rather than a set of experimental values of real compounds. This imprecise treatment of model compounds is further complicated by the fact that their properties,  $\text{p}K_i^M$  and  $f_i(x_i)$ , must be modified if they are considered fixed in a conformation that is differently favored by the different protonation states. More exactly, the  $\text{p}K_a$  of  $M_i$  in conformation  $c$  is<sup>3</sup>

$$\text{p}K_i^M(c) = \text{p}K_i^M + \gamma_i \log \frac{p_i(c|x_i=0)}{p_i(c|\text{any } x_i \neq 0)}, \quad (8)$$

where the numerator and denominator are the conditional probabilities of conformation  $c$  when  $M_i$  is, respectively, charged and neutral. Similarly, the fraction of tautomer  $x_i$  when  $M_i$  is in conformation  $c$  is

$$f_i(x_i|c) = \frac{p_i(x_i|c)}{p_i(\text{any } x_i' \neq 0|c)} = f_i(x_i) \frac{p_i(c|x_i)}{p_i(c|\text{any } x_i' \neq 0)}. \quad (9)$$

These conformational dependences have direct implications to the constant-pH MD method used here, which assumes that the PB calculations are done at fixed protein conformation<sup>4</sup> and thus should use  $\text{p}K_i^M(c)$  and  $f_i(x_i|c)$  values. These dependences could in principle be avoided

by restricting the model compound to a small-enough fragment essentially devoid of conformational flexibility,<sup>54</sup> but this cannot be achieved with charge sets where (de)protonation affects the charge of atoms located at the flexible region of the residue side chain.

In this study, we address the problems described in the previous paragraph by regarding the model compound simply as a conceptual device to account for the contributions to (de)protonation that cannot be captured by PB energies. The adopted rationale is essentially to view the  $\text{p}K_i^M$  and  $f_i(x_i)$  as conformation-independent parameters whose values must be determined from accurate experimental data. Since, as noted above, effectively rigid model fragments are generally not possible to obtain, we require instead that the selected flexible fragments have no functional groups whose positional change may significantly affect (de)protonation. Given that the usual protonatable side chains of amino acids are monofunctional, we take them as the model compounds. For the sites corresponding to the N- and C-terminal sites, we use the fragments  $-\text{COC}^{\text{NH}_3}$  and  $-\text{COOH}$ , respectively. These model compound definitions ensure that a peptide or protein is split into nonoverlapping fragments, as required by Eq. (6) (note that defining the model compound as the whole residue containing the site<sup>34</sup> may lead to overlapping fragments and invalidate the superposition principle). The conformation-independent parameters  $\text{p}K_i^M$  and  $f_i(x_i)$  corresponding to these fragments are thus expected to provide a good modeling of the system over all the conformations explored by the protein during a constant-pH MD simulation. This is similar to the modular approach adopted in MM force fields, where the conformation-independence of the bonded and nonbonded parameters is obviously an approximation. The suitability of this approach can be inferred from the results obtained.

The fractions of the neutral tautomers,  $f_i(x_i)$ , were assigned as previously described.<sup>31,32</sup> In the case of carboxyl sites (Asp, Glu, and C-terminus), we consider four tautomers, corresponding to a proton bonded in either syn or anti geometry to each of the carboxyl oxygen atoms; each syn:anti pair was assigned the ratio 94.5:5.5 inferred from experimental and theoretical studies of acetic acid.<sup>32</sup> For His, we consider two tautomers, corresponding to have the proton in either the  $\text{N}^{\epsilon 2}$  or  $\text{N}^{\delta 1}$  atoms; the  $\text{N}^{\epsilon 2}:\text{N}^{\delta 1}$  pair was assigned the ratio 70:30 measured for His with blocked termini.<sup>55</sup> For all the remaining sites, we considered equiprobable tautomers, corresponding to two in-plane proton positions for Tyr, three  $\text{sp}^3$  proton positions for Cys, and three  $\text{sp}^3$  proton vacancies for Lys and N-terminus.

To determine the  $\text{p}K_i^M$  for each type of model compound, we consider cycle (2) again, but now ignoring tautomerism and regarding the top reaction as the unknown. We then select for P a simple protein-like molecule with a single site with known  $\text{p}K_a$ , perform

constant-pH MD simulations for it at different pH values, and then compute the corresponding  $pK_i^M$  value that gives the best Henderson–Hasselbalch fit to the resulting titration curve. For this purpose, we used the eight blocked pentapeptides studied by Pace and co-workers<sup>56,57</sup> (Ala<sub>5</sub>-NH<sub>2</sub>, Ac-Ala<sub>5</sub> and Ac-Ala<sub>2</sub>-X-Ala<sub>2</sub>-H<sub>2</sub>, where Ac denotes acetyl and X = Asp, Glu, His, Cys, Tyr, and Lys), which should provide a suitable protein-like chemical environment for the fragments selected as model compounds. As seen in Eq. (7), the populations observed in the simulations would depend on  $pK_i^M$  only through the difference  $pK_i^M - \text{pH}$ , meaning that we can perform the simulations of each peptide using an arbitrary  $pK_i^M$  value which is corrected afterward by simply adding the same constant required to bring the pentapeptide titration curve to its experimental position. As noted above, the  $pK_i^M$  values thus obtained should already reflect the overall conformational effects on the  $M_i$  fragments.

### Constant-pH MD methodology

All constant-pH MD simulations were performed using the previously described stochastic titration method.<sup>4–8</sup> The method relies on different sequential blocks. The first block is a PB/Monte Carlo (PB/MC) calculation where the protonation states resulting from the last MC step are assigned to the protein. The second block is the solvent relaxation dynamics, a short MM/MD (MM/MD) simulation of the system with frozen protein, which allows the solvent to adapt to the new protonation states (the duration of this block is hereafter designated as  $\tau_{\text{rlx}}$ ). The last block is a full MM/MD of the unconstrained system (the duration of this block is hereafter designated as  $\tau_{\text{prt}}$ ). A variant of the reduced titration method is used.<sup>5,7,58</sup>

PB and MC calculations were performed with the programs MEAD<sup>59</sup> and PETIT,<sup>32,60</sup> respectively. The atomic charges and radii used in the PB calculations were derived from the GROMOS 43A1 and 53A6 force fields,<sup>61–63</sup> as previously described.<sup>33</sup> All PB calculations consisted of finite-difference linear PB calculations performed with the program MEAD (version 2.2.0)<sup>59</sup> using a temperature of 300 K, a molecular surface defined with a solvent probe radius of 1.4 Å, and a Stern (ion exclusion) layer of 2.0 Å. The dielectric constants were 80 for solvent and 2 for the protein/peptide. The ionic strength used was 0.1 M. A two-step focusing procedure<sup>64</sup> was used, with consecutive grid spacings of 1.0 and 0.25 Å. The MC runs were performed using 10<sup>5</sup> MC cycles, one cycle consisting of sequential state changes over all individual sites and also all pairs of sites with at least one interaction term above 2.0  $pK_a$  units.<sup>60</sup>

The MD simulations were performed with the GROMACS 3.2.1 distribution<sup>65,66</sup> (with some modifications).<sup>5</sup> Explicit water was treated with the single point

charge (SPC) model.<sup>67</sup> Periodic boundary conditions were used with a rhombic dodecahedral unit cell. Long-range electrostatic interactions were treated using the generalized reaction field (GRF)<sup>68</sup> method. A twin-range cutoff was used, with short- and long-range cutoffs of 8 and 14 Å, respectively, and with neighbor lists updated every five steps. All bond lengths were constrained using the LINCS algorithm, and a time step of 2 fs was used. Two Berendsen's temperature couplings to baths at 300 K and with relaxation times of 0.1 ps<sup>69</sup> were used for both the solute and solvent. A Berendsen's pressure coupling<sup>69</sup> at 1 bar was used with a relaxation time of 0.5 ps and a compressibility of  $4.5 \times 10^{-5} \text{ bar}^{-1}$ . The ionic strength in the simulations was set to 0.1 M using a GRF-modified version of the GROMACS 3.2.1 distribution.<sup>5</sup>

### Simulation settings and analysis

The eight pentapeptides (Ala<sub>5</sub>-NH<sub>2</sub>, Ac-Ala<sub>5</sub>, and Ac-Ala<sub>2</sub>-X-Ala<sub>2</sub>-NH<sub>2</sub>, where Ac denotes acetyl and X = Asp, Glu, His, Cys, Tyr, and Lys) were built in an unstructured conformation using PyMOL.<sup>70</sup> The HEWL starting structure was the Brookhaven Protein Data Bank (PDB) entry 4LZT, obtained at pH 4.5.

The protein and the pentapeptides were placed at the center of different boxes and filled with ~5500 and ~1100 water molecules, respectively. The systems were minimized first with ~50 steps of steepest descent followed by ~10,000 steps using the low-memory Broyden–Fletcher–Goldfarb–Shanno algorithm. The initiations were achieved by harmonically restraining all atoms in a 50 ps MD simulation, followed by another 50 ps simulation with only the C $\alpha$  atoms restrained. Runs of 30 ns were performed for lysozyme and of 50 ns for pentapeptides. The relaxation of the solvent ( $\tau_{\text{rlx}}$ ) was done for 0.2 ps while each full system dynamics segment ( $\tau_{\text{prt}}$ ) was 2.0 ps long.

The constant-pH MD titration of HEWL was simulated only in the acidic region, given that there are a high number of acidic sites with atypical  $pK_a$ s. This also allowed us to increase the number of simulations and, therefore, to improve the sampling. The simulations were performed at pH values from 0.5 to 7.0 with increments of 0.5 units. For each pH value, three replicates were performed. In all simulations, the titrating sites were treated in their tautomeric form.<sup>7</sup> The titration was computed by averaging at each pH value the occupancy states of all titratable sites over the final equilibrated segment. The data was fit to a Hill equation to get the  $pK_a$  and Hill coefficient values for each titratable site. The fits were done with a nonlinear Levenberg–Marquardt least-squares algorithm<sup>71</sup> weighting the average protonations with their correlation-corrected errors.<sup>72</sup> The normal-asymptotic standard errors thus obtained for the  $pK_a$  parameters were used as a measure of their statistical uncertainty.

The pentapeptides were titrated at different pH values. We chose 3–5 pH values close to the titration midpoint



**Table I**pK<sub>a</sub> Values for All Pentapeptides Studied Using Both 43A1 and 53A6 Force Fields

Residue	Tanford pK <sub>a</sub>	Exp. pK <sub>a</sub>	43A1			53A6		
			pK <sub>a</sub>	Shift	pK <sup>M</sup>	pK <sub>a</sub>	Shift	pK <sup>M</sup>
Asp	4.0	3.94	4.37	+0.43	<b>3.57</b>	4.34	+0.40	<b>3.60</b>
Glu	4.4	4.25	4.61	+0.36	<b>4.04</b>	4.53	+0.28	<b>4.12</b>
His	6.45	6.54	5.96	−0.58	<b>7.03</b>	6.21	−0.33	<b>6.78</b>
Cys	9.5	8.55	10.05	+1.50	<b>8.00</b>	9.47	+0.92	<b>8.58</b>
Tyr	9.6	9.84	9.93	+0.09	<b>9.51</b>	9.92	+0.08	<b>9.52</b>
Lys	10.4	10.40	10.19	−0.21	<b>10.61</b>	10.30	−0.10	<b>10.50</b>
CTer	3.8	3.67	4.34	+0.67	<b>3.13</b>	4.49	+0.82	<b>2.98</b>
NTer	7.5	8.00	6.92	−1.08	<b>8.58</b>	7.52	−0.48	<b>7.98</b>

The simulations were performed using the Tanford model compound pK<sub>a</sub>s.<sup>33</sup> Experimental pK<sub>a</sub>s were taken from Refs. 56 and 57. The new calibrated pK<sup>M</sup> values are marked in bold in the table.

of each group. Three replicates were performed for every pH value. The titrations were computed by averaging at each pH value the occupancy states of each titrable site over the whole simulation. The pK<sub>a</sub> values were obtained by fitting the data to the corresponding Henderson–Hasselbalch curve.

All analyses were done using GROMACS,<sup>65,66</sup> Gnuplot,<sup>73</sup> or in-house tools.

## RESULTS AND DISCUSSION

### pK<sup>M</sup> calibration using pentapeptide simulations

The determination of model compound pK<sub>a</sub> values from experimental data, following the procedure described in the Theory and Methods section, used the potentiometric pK<sub>a</sub> measurements done by Pace and co-workers<sup>56,57</sup> on eight blocked pentapeptides (Ala<sub>5</sub>-NH<sub>2</sub>, Ac-Ala<sub>5</sub>, and Ac-Ala<sub>2</sub>-X-Ala<sub>2</sub>-NH<sub>2</sub>, where Ac denotes ace-

tyl and X = Asp, Glu, His, Cys, Tyr, and Lys). We performed constant-pH MD simulations on these pentapeptides with two different force fields, GROMOS 43A1 and 53A6. The initial pK<sup>M</sup> values used in these simulations, shown as “Tanford pK<sub>a</sub>” in Table I, were the ones from Nozaki and Tanford,<sup>53</sup> except for His.<sup>31</sup> From the titration curves produced by these simulations, we obtained for each pentapeptide the pK<sub>a</sub> value and shift to the experimental value (Table I). The new calibrated pK<sup>M</sup> values, shown simply as “pK<sup>M</sup>” in Table I, were obtained by subtracting those shifts from the Tanford values. The new pK<sup>M</sup> values thus obtained already reflect the overall conformational effects on the model compound fragments.

Interestingly, we note that the pK<sup>M</sup> values obtained with the 53A6 force field are, in all cases but CTer, closer to the experimental values for the residue in the pentapeptide template, indicating that the effect of the peptide environment is smaller in that case. When comparing with the previous Tanford pK<sup>M</sup> values, the sites Cys and NTer seem to be the ones with larger shifts, especially with the 43A1 force field, meaning that they would be those more affected by systematic errors when using that set of values.

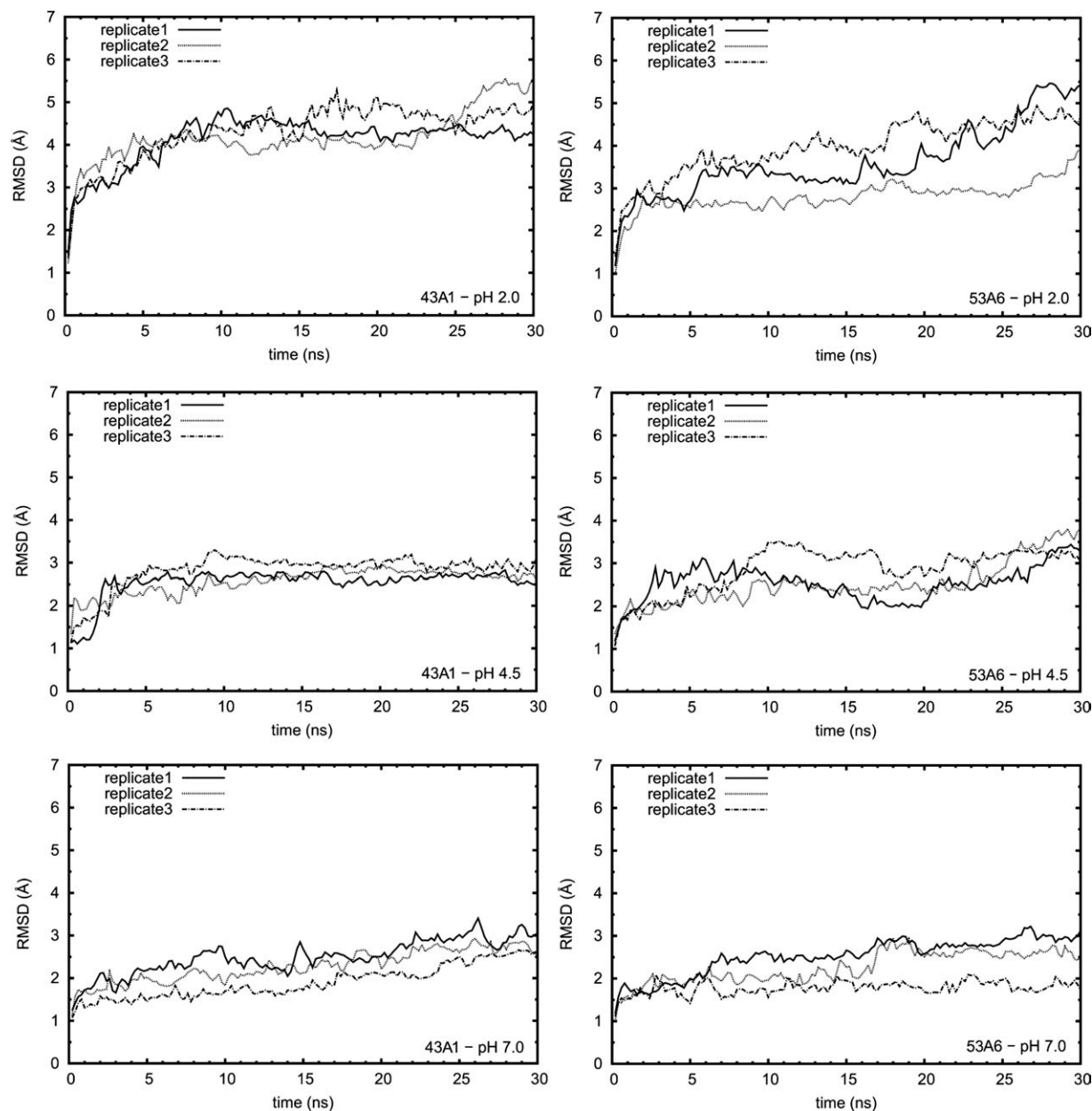
### HEWL titration

The constant-pH MD simulations of HEWL were performed using both the 43A1 and 53A6 GROMOS force fields. In Table II, there is a collection of pK<sub>a</sub> values obtained for the equilibrated segment of the simulations (10–30 ns). A comparison between the 43A1 results and the ones from previous work<sup>7</sup> allows for a measure of the improvement only due to the redefinition of the model compound fragments and to the pK<sup>M</sup> calibration against experiment. With the new approach we obtained an overall pK<sub>a</sub> root mean square error (RMSE) of 0.70, corresponding to an improvement of 0.12 pK units

**Table II**pK<sub>a</sub> Values for All Sites Titrating at Acidic pH Using the Equilibrated Segment of the Simulations (last 20 ns) in Both Force Fields

Residue	Exp. pK <sub>a</sub> (avg.)	Prev. 43A1 <sup>7</sup>		43A1		53A6	
		pK <sub>a</sub>	Error	pK <sub>a</sub>	Error	pK <sub>a</sub>	Error
Glu-7	2.85	4.04	<b>1.19</b>	3.51 ± 0.03	0.66	3.35 ± 0.05	0.50
His-15	5.36	4.22	<b>1.14</b>	4.81 ± 0.16	0.55	5.68 ± 0.15	0.32
Asp-18	2.66	4.11	<b>1.45</b>	3.28 ± 0.06	0.62	3.32 ± 0.07	0.66
Glu-35	6.20	5.50	0.70	4.86 ± 0.10	<b>1.34</b>	5.04 ± 0.11	<b>1.16</b>
Asp-48	1.60	2.56	0.96	2.87 ± 0.10	<b>1.27</b>	2.66 ± 0.10	<b>1.06</b>
Asp-52	3.68	3.96	0.28	3.62 ± 0.07	0.06	2.78 ± 0.06	0.90
Asp-66	0.90	1.50	0.60	1.35 ± 0.11	0.45	2.40 ± 0.04	<b>1.50</b>
Asp-87	2.07	2.68	0.61	2.31 ± 0.05	0.24	2.37 ± 0.05	0.30
Asp-101	4.09	3.77	0.32	3.62 ± 0.06	0.47	3.59 ± 0.05	0.50
Asp-119	3.20	2.97	0.23	2.76 ± 0.06	0.44	3.04 ± 0.08	0.16
CTer-129	2.75	3.23	0.48	3.16 ± 0.04	0.41	3.15 ± 0.06	0.40
RMSE			0.82		0.70		0.79

All experimental pK<sub>a</sub> values were taken from Ref. 74. Results from previous work using the 43A1 force field<sup>7</sup> are presented for a better comparison. All errors over 1 pK unit are marked in bold. RMSE was computed relative to the average value of the experimental range and is presented in pK units. The statistical uncertainties (shown as ± values) of the pK<sub>a</sub> values were computed as described in the Theory and Methods section.



**Figure 1**

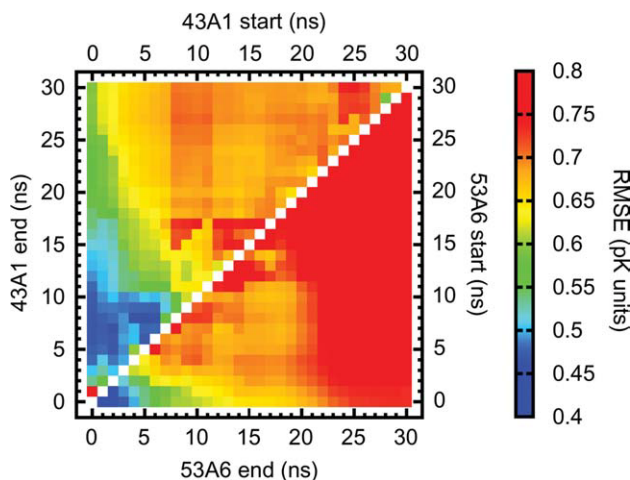
Temporal evolution of the RMSD of the simulations at three different pH values (2.0, 4.5, and 7.0) for both force fields used in this work. There are three replicates for each pH value. The fit and RMSD values were computed using the C $\alpha$  atoms and taken relative to the X-ray structure.

relative to that previous study. From the individual  $pK_a$  values, we got significant improvements in some residues but the predictions for Glu 35 and Asp 48 became significantly worse. The case of Glu 35 is particularly important because this residue is the proton donor in the catalytic site of HEWL and has a highly shifted  $pK_a$ .

The  $pK_a$ s obtained with the 53A6 force field were somewhat disappointing, not due to the RMSE obtained (0.79 is a good overall result), but because the method did a poor job dealing with the internalized residues. In

fact, such residues seemed to lose some of their local environment during the constant-pH MD simulations, behaving more like solvated residues. Interestingly, this loss of environment effects using 53A6 was also observed in the pentapeptide simulations (see above). This suggests that the results for the pentapeptides and HEWL may be both reflecting some tendency of 53A6 to exaggerate solvent exposure.

These results raise several questions regarding the influence of these two force fields on the stability of

**Figure 2**

$pK_a$  RMSE color map obtained with GROMOS force fields 43A1 (upper-left triangle) and 53A6 (lower-right triangle). The RMSE values were computed for all time segments multiple of 1 ns, with each colored square in the plot representing a segment of time comprised by the start and end points indicated in the axes. To present both force fields with a symmetrical orientation, the  $x$  and  $y$  axes were exchanged for 53A6.

HEWL and its impact on the prediction of correct protonation states. From the temporal evolution in the  $C^\alpha$  RMSD plots (illustrated in Fig. 1) and from visual inspection of the structures, we do not see any major conformational transition that could seriously compromise the  $pK_a$  calculations. Nevertheless, the 53A6 simulations seem to exhibit more variations in the observed RMSDs (Fig. 1) and a slight decrease in helical content (data not shown). A key factor may be the significantly high  $C^\alpha$  RMSD values observed after equilibration in the simulations with both force fields. Although the particu-

larly high RMSD values at pH 2.0 certainly reflect the structural destabilization that HEWL starts to experience at very low pH,<sup>75–77</sup> even the values of 2–3 Å observed at less extreme pH values seem excessive considering that the 4LZT X-ray crystallographic structure was determined at pH 4.5–4.6 (close to the optimal pH of  $\sim 5$  for the enzyme activity). Inspection of the RMSD plots reveals that most of the structural deviation takes place early in the simulations (as seen in the examples in Fig. 1), often during the first nanoseconds, suggesting that the conformations used in the  $pK_a$  calculations are all systematically deviated from the 4LZT structure. In other words, this raises the question of whether the structural rearrangements allowed by the force fields have created any significant distortions in addition to the intended structural reorganization. In particular, the aforementioned slight loss of helical content observed with 53A6 is in agreement with a previous study reporting a tendency of this force field to disrupt helical peptides.<sup>28</sup>

To investigate how the gradual structural rearrangements occurring along the simulations have affected the  $pK_a$  predictions, we split our 30 ns of simulations into all possible fragments with sizes multiple of 1 ns. For all these fragments ( $((30^2 - 30)/2 = 435)$ ), we did independent  $pK_a$  calculations of HEWL followed by calculation of the corresponding RMSE values. The results obtained for both force fields were combined into a single color map plot (Fig. 2) for easier comparison.

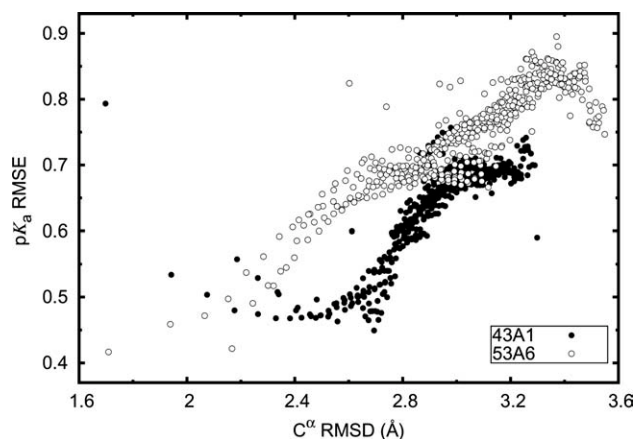
From Figure 2, it is evident that the two force fields have distinct time-dependent RMSE patterns. The GROMOS 43A1 simulations present somewhat poor  $pK_a$  predictions at the very initial time segments, especially the first one (0–1 ns) when the protein conformations sampled are still close to the crystal structure. Nevertheless, after some equilibration time, the RMSE values of the next short equilibrated segments decrease signifi-

**Table III**  
Best  $pK_a$  Value Predictions

Residue	Experimental $pK_a$		43A1 (5–8 ns)		53A6 (0–1 ns)	
	Range	Avg	$pK_a$	Error	$pK_a$	Error
Glu-7	2.60–3.10	2.85	$3.48 \pm 0.06$	0.63 (0.38)	$3.34 \pm 0.08$	0.49 (0.24)
His-15	5.29–5.43	5.36	$4.98 \pm 0.11$	0.38 (0.31)	$5.43 \pm 0.06$	0.07 (0.00)
Asp-18	2.58–2.74	2.66	$3.11 \pm 0.08$	0.45 (0.37)	$3.43 \pm 0.05$	0.77 (0.69)
Glu-35	6.1–6.3	6.20	$5.17 \pm 0.09$	<b>1.03</b> (0.93)	$5.57 \pm 0.12$	0.63 (0.53)
Asp-48	1.2–2.0	1.60	$1.73 \pm 0.11$	0.13 (0.00)	$1.46 \pm 0.10$	0.14 (0.00)
Asp-52	3.60–3.76	3.68	$3.65 \pm 0.05$	0.03 (0.00)	$3.37 \pm 0.10$	0.31 (0.23)
Asp-66	0.4–1.4	0.90	$1.02 \pm 0.30$	0.12 (0.00)	$1.30 \pm 0.34$	0.40 (0.00)
Asp-87	1.92–2.22	2.07	$2.12 \pm 0.09$	0.05 (0.00)	$2.38 \pm 0.09$	0.31 (0.16)
Asp-101	4.02–4.16	4.09	$3.82 \pm 0.07$	0.27 (0.20)	$3.80 \pm 0.04$	0.29 (0.22)
Asp-119	3.11–3.29	3.20	$2.90 \pm 0.04$	0.30 (0.21)	$2.87 \pm 0.04$	0.33 (0.24)
CTer-129	2.63–2.97	2.75	$3.21 \pm 0.09$	0.46 (0.24)	$3.06 \pm 0.04$	0.31 (0.09)
RMSE				0.45 (0.36)		0.42 (0.30)

The segment chosen for each force field was the one with the lowest RMSE. All experimental  $pK_a$  ranges were taken from Ref. 74. All errors over 1 pH unit are marked in bold. The errors in parenthesis are calculated relative to the experimental range. RMSE was computed relative to the midpoint of the experimental range (or the range itself when presented in parenthesis). The statistical uncertainties (shown as  $\pm$  values) of the  $pK_a$  values were computed as described in the Theory and Methods section.



**Figure 3**

Scatter plot of  $C^\alpha$  RMSD versus  $pK_a$  RMSE calculated from simulations with GROMOS force fields 43A1 and 53A6. Both RMSD and RMSE were obtained over different time segments following the procedure used in Figure 2. The RMSE values are presented in pK units. The RMSD values were computed using the  $C^\alpha$  atoms and taken relative to the X-ray structure.

cantly. The best  $pK_a$  predictions were obtained in one of these segments (5–8 ns) with a remarkable RMSE value of 0.45 pK units. For longer segments or segments starting after  $\sim 10$  ns, the RMSE values increased consistently, reaching values close to the one reported above for the whole equilibrated segment (0.70).

The GROMOS 53A6 simulations present a very uniform gradient of RMSE values, with the best  $pK_a$  predictions at the very initial time segments. In fact, the best predictions occur at the first segment (0–1 ns) with the even more remarkable RMSE value of 0.42 pK units, which is probably the lowest RMSE value ever reported for HEWL. Unlike the 43A1 force field, the 53A6 seems to perform outstandingly well when the protein conformations are still very close to the crystal structure. This becomes even more obvious from the fact that the RMSE values also increase dramatically with increasing structural RMSDs (see Fig. 1). Segments after the first  $\sim 10$  ns result in poorer  $pK_a$  predictions.

Table III presents the  $pK_a$  values calculated for the best simulation segments of each GROMOS force field. With the 43A1 force field and using the segment 5–8 ns, the method is able to correctly predict four Asp residues inside the experimental ranges, while still showing a relatively high error ( $>1$  pK unit) for the Glu 35 residue. When we used the first nanosecond of the 53A6 simulations, we were able to predict all residues under 1 pK unit error and, in three cases, we obtained  $pK_a$  values inside the experimental range. All the RMSE values mentioned till now were obtained using the usual procedure of taking the midpoint of the measured range as “the” experimental value. As this implies that even computed values lying within the measured experimental range

have an associated “error,” it can be argued that a more meaningful approach is to compute the error to the nearest experimental bound<sup>31</sup> (which obviously gives lower errors). Therefore, Table III also presents the results obtained with this second approach (0.36 and 0.30 for GROMOS 43A1 and 53A6, respectively), which is probably more directly comparable to the accuracy of experimental measurements. It must be stressed that, by showing these best-segment results, we are not arguing that these are the actual predictions of the simulations, but rather pointing how significantly those predictions seem to be affected by the apparently excessive structural rearrangements imparted by the force fields. It is also important to note that the trends seen in Figure 2 clearly indicate a systematic effect and not a mere statistical dependence on the selected segment.

In Figure 2, it is clear that, especially for the simulations using the 53A6 force field, the  $pK_a$  RMSE values are increasing with simulation time and might thus be correlated with the  $C^\alpha$  RMSD. To investigate this, we did scatter plots of these two properties (Fig. 3). Although the 43A1 simulations present a negative correlation at low RMSD values, they exhibit a high positive correlation for RMSD values above  $\sim 2.7$  Å. The simulations using 53A6 force field exhibit a strong correlation between the two properties. These results support the idea that, although structural reorganization is obviously crucial in a constant-pH MD simulation, the significant deviations from the crystal structure that eventually take place in the long run may be detrimental to its ability to predict  $pK_a$  values.

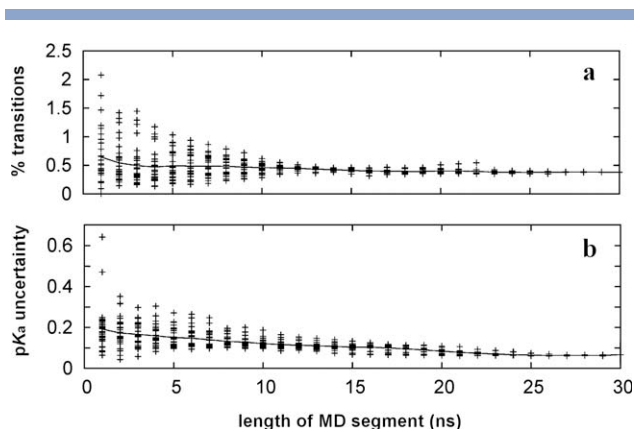
It is also instructive to examine how the length of the 435 MD segments considered in Figures 2 and 3 affects the frequency of protonation/deprotonation transitions and the precision of the subsequent  $pK_a$  calculations. As the number of transitions is expected to be very low near complete protonation or deprotonation, we consider for each site and MD segment the pH value that brings that site closer to its titration midpoint, computing the percentage of transitions relative to the maximum allowed

**Table IV**

Average of the Percentage of Protonation/Deprotonation Transitions Observed for Each Site

Residue	43A1	53A6
Glu-7	26.2	24.6
His-15	15.1	23.7
Asp-18	14.7	8.3
Glu-35	9.0	3.7
Asp-48	11.2	8.8
Asp-52	3.6	1.5
Asp-66	0.5	5.7
Asp-87	11.0	10.5
Asp-101	32.8	32
Asp-119	33.7	33.9
CTer-129	30.9	25.8

Each average was computed over the 435 MD segments considered in Figure 2, taking for each segment the pH value for which the site being considered is closer to midpoint titration.



**Figure 4**

Percentage of protonation/deprotonation transitions (a) and  $pK_a$  statistical uncertainty (b) versus the length of each of the 435 MD segments, for site Asp-66 using the 43A1 force field. Each percentage of transitions corresponds to the pH value for which the site being considered is closer to midpoint titration for that particular MD segment. Lines indicate the average values.

by the deprotonation fraction obtained in that segment; this results in 435 values of the percentage of transitions for each site, which should be representative of the actively titrating site. The averages of those percentages, shown in Table IV, are not strongly dependent on the force field, but they vary widely from “fast” sites like Asp-119 to very “slow” sites like Asp-66. Interestingly, there is not a marked dependence of the percentage of transitions with the length of the MD segment, even for extremely slow sites, as illustrated in Figure 4(a) for Asp-66 using the 43A1 force field—although there is a larger spread of values for short segments, the same average percentage is kept. As seen in Figure 4(b), a similar trend is observed for the  $pK_a$  statistical uncertainties, which already reflect the fast or slow nature of the transitions through the correlation time of the protonation states (see Theory and Methods section). Similar features are observed for the other sites. This indicates that the precision of the  $pK_a$  calculations is remarkably robust to variations in the length of the MD segment being considered. Furthermore, we do not find any general relation between the percentage of transitions and the accuracy of the calculated  $pK_a$  values, which may be positively or negatively correlated, depending on the site. For example, more frequent transitions of Asp-48 are usually associated with an excessive solvent exposure that results in a worse prediction (the transition percentages and the errors of the  $pK_a$  prediction have a correlation around 0.6 using either force field). These results seem to indicate that, as in other Monte Carlo simulations,<sup>72,78</sup> there is no reason to expect that a high frequency of transitions will result in a better sampling of states. Each site probably has a typical transition frequency that simply reflects its local environment, and trying to increase that

frequency may be actually inadvisable—forcing a site to be constantly titrating may end up favoring conformations that are equally favored by both the protonated and deprotonated states, eventually missing conformations that are more characteristic of either state and crucial in determining the titration behavior of that site.

## CONCLUSIONS

Two limitations to the performance of our constant-pH MD methodology were investigated in this work. The first is the ill-defined nature of model compounds in PB calculations, both in terms of molecular definition and  $pK_a$  value. The second limitation is the eventual distortion that MM force fields may impart into the structure in addition to the intended reorganization to (de)protonation events.

The model compound issue was addressed here by inverting the typical thermodynamic cycle (equation 2) in such a way that, using experimental  $pK_a$  values from a set of pentapeptides, we were able to obtain calibrated  $pK^M$  values that properly model our unphysical model compound fragments. Using this calibration, we calculated new  $pK_a$  values for the acidic range of HEWL and obtained a significant improvement on the RMSE value (0.70 instead of the previous 0.82 pK units) using the GROMOS 43A1 force field. Another calculation was done using the newer GROMOS 53A6 force field, which yielded a RMSE value of 0.79 pK units.

The force field influence on the constant-pH MD method's ability to predict  $pK_a$  values was studied here in detail by looking individually at all simulation segments with a length multiple of 1 ns, for the 43A1 and 53A6 GROMOS force fields. We observed that the method performed very distinctly using either force field, especially at the initial nonequilibrated segments of the simulations. With the 43A1 force field, the method requires an initial equilibration period after which it performs with very high precision (RMSE of 0.45 pK units) during short nanosecond segments. In contrast, with the 53A6 force field the performance decreases continually throughout the MD simulation, although the RMSE obtained using only nonequilibrated 1 ns simulations is remarkably low (0.42 pK units). These results seem to indicate that, despite allowing for structural reorganization to (de)protonation, MM force fields may also be inducing some structural distortions, as indicated by recent studies.<sup>26–29</sup> In particular, one of those studies<sup>28</sup> reported that the 53A6 force field led to a destabilization of helical peptides, while no such effect was observed for 43A1; this could explain the slight loss of helical content of HEWL observed here using 53A6 and was presumably solved by a recent reparameterization that we plan to investigate.<sup>79</sup> Thus, although current constant-pH MD methods are among the best approaches to predict  $pK_a$

values, the problems of the underlying MM force fields may be hindering their progression toward an unprecedented level of accuracy. On the other hand, these results suggest that very short constant-pH MD simulations may be a suitable approach to predict  $pK_a$  values for proteins whose structure is not expected to be markedly dependent on pH.

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