

# Reactivity of Lattice Oxygens Present in $V_2O_5(010)$ : A Periodic First-Principles Investigation

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The present paper deals with the periodic first-principles density functional (DFT) calculations on hydrogen adsorption on the  $V_2O_5(010)$  surface. The calculated results reveal that the vanadyl oxygen is the most active site for hydrogen adsorption among the three kinds of structurally different lattice oxygens in terms of their coordination with vanadiums. Our calculated harmonic vibrational frequencies of the three types of OH species are found to correlate with their corresponding bond strengths and support their respective reactivity. It is confirmed that the lattice relaxation contributes to the reactivity of the oxygens significantly in all cases while the local environment of adsorption site affects both the geometry and reactivity of the adsorption system only in the case of tricoordinated oxygen. Hydrogen adsorption at these oxygens reduces the surface in different ways. Comparison of the desorption ability of the OH species from the surface shows that removal of the  $O_1H$  species, formed by H adsorption at vanadyl oxygen, is energetically preferable. The present work demonstrates that the periodic approach gets rid of the artifacts of the cluster method, and thus, first-principles DFT methodology demonstrates its reliability to investigate geometric, electronic, and catalytic properties of transition metal oxide catalysts.

## 1. Introduction

Vanadium pentoxide ( $V_2O_5$ ) is a very important material that has been extensively used as n-type semiconductors, cathode material in lithium batteries, and industrial catalysts.<sup>1–3</sup> It is mainly used as a catalyst in selective oxidation processes of hydrocarbons and sulfur dioxide, selective catalytic reduction (SCR) of nitric oxide by ammonia, and also photocatalytic reactions.<sup>2–4</sup> It is believed that the catalytic properties of a  $V_2O_5$ -based catalyst depend strongly on their ability to provide lattice oxygens for the reactions. Therefore, the investigation of the  $V_2O_5$  lattice oxygen is crucial for understanding the active sites and reaction mechanisms. Three kinds of structurally different lattice oxygens exist in layered  $V_2O_5$ : singly coordinated  $O_1$ , which is a vanadyl oxygen ( $V=O$  species); dicoordinated  $O_2$  and tricoordinated  $O_3$ , which bridge two and three vanadium atoms, respectively.<sup>5</sup>

One of the key steps in both selective oxidation of hydrocarbons<sup>2</sup> and SCR reactions<sup>3</sup> is the abstraction of an H atom from a reactant by one of the surface oxygens to form surface OH species. Hence, a study of adsorption of the H atom at  $V_2O_5$  lattice oxygens is necessary to understand the active sites and reaction mechanisms. Furthermore, it is well-known that the interaction between reactant adsorbates and the surface is determined not only by the local chemical environment but also by the contribution of lattice relaxation. In this study, to rationalize the active oxygen sites as well as to study the influence of local environment and lattice relaxation, we performed the periodic first-principles density functional (DFT) calculations with respect to hydrogen adsorption.

Despite the extensive interest in  $V_2O_5$  as a catalyst, some essential points regarding its active site in catalytic reactions

still remain unresolved. Identification of the exact active site is one of the key issues, and different authors have proposed different active sites.<sup>6–23</sup> It has been generally accepted<sup>6–16</sup> that the vanadyl oxygen plays the most important role in catalysis. Tarama et al. has predicted for the first time the key role of the  $O_1$  atom by studying the adsorption state on the  $V_2O_5$  surface using infrared (IR) spectroscopy,<sup>6</sup> and later it has been widely supported by other experimental<sup>7–12</sup> and theoretical<sup>13–16</sup> workers. One of the present authors has investigated the selective oxidation of hydrocarbons and SCR of NO catalyzed by supported and unsupported  $V_2O_5$  catalysts experimentally<sup>10,11</sup> and theoretically<sup>15</sup> and pointed out that the  $O_1$  oxygen acts as the active site. Using an empirical formula, Andersson<sup>13</sup> calculated the electronic structure of  $V_2O_5$  and suggested that the  $V=O$  species on the surfaces are responsible for the catalytic oxidation of hydrocarbons. Recently, employing the monolayer cluster model, Zhanpeisov et al.<sup>14</sup> have investigated the adsorption of both the H atom and proton on the  $V_2O_5$  surface as well as the formation of oxygen vacancies by means of the SINDO1 approach. They concluded that the active site of  $V_2O_5$  is probably associated with the  $O_1$  site (terminal  $V=O$  group). Gilardoni et al.<sup>16</sup> have also proposed that the  $O_1$  oxygen plays a key role in the SCR reactions from their calculations using a cluster model composed of two vanadiums.

In contrast to the aforementioned work regarding the most important role of the  $O_1$  atom in catalysis, some authors have reported both from experiments<sup>17,18</sup> and theoretical calculations<sup>19–22</sup> that the  $O_2$  oxygen participates in the catalytic process. Atomic force microscopy observation<sup>17</sup> revealed that the outermost (010) surface adsorbed by water molecules is not composed only of vanadyl oxygens. On the basis of this observation and previous theoretical studies as well as from simple geometrical considerations, they have inferred that the

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O<sub>2</sub> oxygen could act as the active site. Gasior et al.<sup>18</sup> have proposed that the O<sub>2</sub> oxygen plays the crucial role in SCR reactions. On the theoretical side, Witko and co-workers<sup>19–22</sup> have performed the semiempirical and ab initio calculations using the cluster approach. Their model clusters exhibit some artifacts: (1) the V–O bond lengths for all O atoms except for the O<sub>1</sub> atom were set equal to 1.83 Å, which is the respective average value in bulk V<sub>2</sub>O<sub>5</sub>; (2) the substrate unit of V<sub>2</sub>O<sub>5</sub> was kept fixed at the geometry assumed by (1), ignoring the lattice relaxation; and (3) hydrogen was added in a direction perpendicular to the surface at each oxygen site, neglecting the influence of the local environment of the OH group. Owing to the above-mentioned reasons, their conclusion indicating the O<sub>2</sub> atom to be the active site becomes doubtful. So far, very few studies asserting that the O<sub>3</sub> atom acts as an active site have been reported. Using the semiempirical calculations, Ramirez et al.<sup>23</sup> have indicated that the most stable vacancy corresponds to the O<sub>3</sub> site rather than to the O<sub>1</sub> and/or O<sub>2</sub> site. Subsequently, they have found using IR data that the O<sub>3</sub> atoms participate more readily in the oxidation of dimethyl sulfoxide.

Mainly because of the limitations of experiments and also because of the restriction of computational techniques, the role of lattice oxygens present in V<sub>2</sub>O<sub>5</sub> surface has not yet been clarified to formulate an understanding of the reactivity. Theoretically so far, studies concerning this issue have been performed by the cluster approach.<sup>13–16</sup> However, a difficult problem within the cluster approach for surface systems is the selection of size and shape of the model. Clusters provide a testing ground for the evaluation of surface properties, especially in small clusters where almost all atoms considered are surface atoms, although their coordination can be very different from surfaces of the corresponding bulk materials. Therefore, one must be aware of the artifacts that are introduced into the system by the calculation of physical parameters as functions of cluster size and geometry. Since the convergence to the bulk limit for the cluster and the adsorption is predictable by extrapolation with increasing cluster size, it was believed that such a large cluster can give reliable predictions.<sup>24–26</sup> However, in many cases, especially when cluster size is not large enough, the cluster approach may provide only a part of the reliable information, and it fails to supply comprehensive one. To have comprehensive predictions, one must attempt an artifact-free approach. Here, in heterogeneous catalysis, the periodic boundary method was found to be able to rid the problem associated with the localized cluster approach for surface systems.<sup>27–33</sup> Such an attempt of course would require far more computational effort than is needed for the application of the cluster approach.

In recent years, with the rapid development of computational techniques, it has become feasible to deal with the adsorption systems using periodic boundary first-principles calculations. Our previous study<sup>29</sup> shows that the periodic DFT approach is very reliable for investigating V<sub>2</sub>O<sub>5</sub>, since the results are in very good agreement with experimental observations.

The present work is organized as follows. In section 2 we describe the computational details and model, and section 3 presents the results and discussion, which include a comparison and evaluation of the systematic optimization techniques. Finally, section 4 summarizes the highlights of this paper.

## 2. Methodology

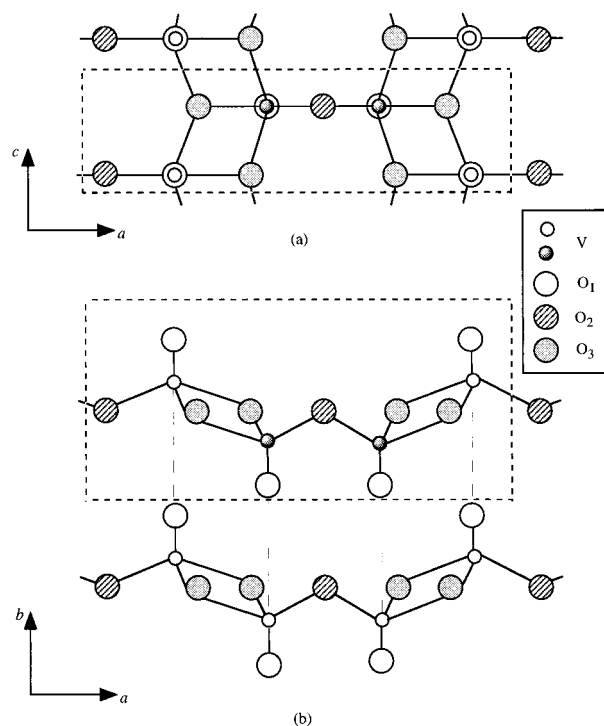
**2.1. Computational Details.** The periodic DFT calculations were performed using the DSolid program<sup>34</sup> within the Kohn–Sham formalism.<sup>35</sup> In this program one-electron Schrödinger equations are solved only at the  $k = 0$  wave vector point of the

Brillouin zone. This program can be applied to systems that have small unit cells simply by ignoring this and treating them as having a larger cell, as described in the DSolid User Guide.<sup>34</sup> We used the Vosko–Wilk–Nusair local type functional (LDA)<sup>36</sup> for the geometry optimizations, whereas the interaction energies were computed using the Becke–Lee–Yang–Parr nonlocal type functional (NLDA)<sup>37,38</sup> that contains the necessary gradient corrections to eliminate the inaccuracies of the LDA for the estimations of the energy differences. Molecular orbitals were expanded into a set of numerical atomic orbitals of doubly numerical + polarization (DNP) quality provided by the software package. The DNP is comparable in quality to Pople's split-valence 6-31G\*\* basis set and usually yields the most reliable results.<sup>39–41</sup> To eliminate convergence problems in some cases, we allowed partial occupancies for the high-lying orbitals using smear values of less than 0.02 eV. For the numerical integration we used the MEDIUM quality mesh size of the program. To reduce the necessary CPU time, we froze the inner-core orbitals of the atoms except for hydrogen, whose pure potential was used; i.e., 1s was frozen for oxygen and 1s2s2p for vanadium. For geometry optimization, the BFGS routine in which the gradients were computed numerically was employed. The tolerances of energy and gradient convergence are used as their default values:  $1.0 \times 10^{-6}$  and  $2.0 \times 10^{-3}$ . We have also performed the vibrational frequency calculations at the LDA level using the LDA-optimized structures. Two-point differencing and a step size of 0.01 were used.

**2.2. Model.** Owing to the importance in catalysis mainly, the real structure of V<sub>2</sub>O<sub>5</sub>(010) has been discussed well in the literature.<sup>12,13,20–22,42–50</sup> Experimental studies show that V<sub>2</sub>O<sub>5</sub>-(010) plays a key role in catalysis<sup>51,52</sup> and that the supported (010) monolayer has excellent catalytic activity.<sup>53,54</sup> Experimental<sup>47</sup> and theoretical<sup>29,31</sup> studies demonstrate that the (010) surface has very similar physical properties and stability compared with its bulk crystal. Therefore, we considered V<sub>2</sub>O<sub>5</sub>-(010) for our present study. The surface has been modeled by a periodic slab composed of single unit cell (SUC) monolayer that has been successfully employed for band calculations,<sup>29–33</sup> as shown in Figure 1, where the optimized lattice parameters of  $a$  (11.535 Å) and  $c$  (3.603 Å) obtained by the CASTEP program were used.<sup>29</sup> The successive slabs along the  $b$  axis are separated by a vacuum region. The shortest distance between atoms belonging to successive slabs is greater than 8.3 Å, where the interlayer interaction is not significant.

The large unit cell (LUC) method,<sup>55–60</sup> which is mainly investigated by Evarestov and co-workers,<sup>55–58</sup> is believed to be better than the SUC one. Accordingly, it needs much more computational effort for the LUC calculations. Since a V<sub>2</sub>O<sub>5</sub> unit cell consists of 14 atoms, even the smallest LUC model contains 28 atoms, including 8 vanadiums and 20 oxygens and excluding adsorbate atom(s). Such calculations consume enormous time. To understand the difference between the LUC and SUC methods, we made the LUC model, which is twice the V<sub>2</sub>O<sub>5</sub>(010) unit cell monolayer, by doubling the lattice parameter  $c$  (3.603 Å) and performed the calculations. As predicted, the obtained bond lengths and atomic charges (Tables 1 and 2) show that some differences exist between the two models with different sizes; however, the differences are not significant. In particular, the ordering of the bond lengths and atomic charges were found to be the same in the case of SUC. This indicates that the SUC model used in the present study is found to be a good compromise between cost and accuracy.

To investigate the influence of the V<sub>2</sub>O<sub>5</sub>(010) sublayer on the geometric and electronic structures quantitatively, we have



**Figure 1.** Periodic boundary model of V<sub>2</sub>O<sub>5</sub>(010) used in the present study is indicated by broken lines, projection onto (a) the *ac* plane and (b) the *ab* plane. The closed small circles represent the exposed V atoms on the surface, and the opened small ones designate the V atoms beneath the singly coordinated oxygens.

**TABLE 1: O–V Bond Lengths (Å) in V<sub>2</sub>O<sub>5</sub>(010) Monolayer Composed of the Single Unit Cell (SUC) and Large Unit Cell (LUC)**

bond	O <sub>1</sub> –V	O <sub>2</sub> –V	O <sub>3</sub> –V	O <sub>3</sub> –V
SUC (Å)	1.57	1.74	1.97	2.04
LUC (Å)	1.57	1.78	1.91	2.03

**TABLE 2: Mulliken Atomic Charges (au) of V<sub>2</sub>O<sub>5</sub>(010) Monolayer Made Up of the SUC and LUC**

atoms	O <sub>1</sub>	O <sub>2</sub>	O <sub>3</sub>	V
SUC	–0.324	–0.602	–0.646	+1.272
LUC	–0.301	–0.610	–0.709	+1.315

also performed calculations using a bilayer slab composed of one unit cell. The following points were justified through the obtained results (Tables 3 and 4). (1) For the clean bilayer surface, no significant differences of both the V–O bond lengths and atomic charges are observed between the monolayer and bilayer slabs, and the ordering of the bilayer values is in good agreement with that of the monolayer slab. (2) Similarly for the adsorption systems, the major geometrical parameters and atomic charges are also very close to those of the monolayer. In particular, the ordering of the bilayer results almost remains the same as that of the monolayer model. All these good agreements show that our monolayer slab is sufficient for studying the interaction with hydrogen, although little difference was found when the thickness of the slab was increased.

Therefore, the SUC monolayer slab used in the current study can be considered approximately as a realistic model of V<sub>2</sub>O<sub>5</sub>(010) for the investigation of hydrogen adsorption.

### 3. Results and Discussion

To investigate the effect of both the local environment of an adsorption site and lattice relaxation, systematic optimization

**TABLE 3: Selected Bond Lengths (Å) Obtained by Using Monolayer and Bilayer SUC of V<sub>2</sub>O<sub>5</sub>(010)<sup>a</sup>**

adsorption sites	bonds	I	II	III	clean surface
O <sub>1</sub> site	O <sub>1</sub> –H	0.97 (0.97)	0.97 (0.97)	0.97 (0.97)	
	O <sub>1</sub> –V	1.68 (1.70)	1.69 (1.70)	1.70 (1.70)	1.57 (1.56)
O <sub>2</sub> site	O <sub>2</sub> –H	0.98 (0.98)	0.98 (0.98)	0.98 (0.98)	
	O <sub>2</sub> –V	1.85 (1.83)	1.85 (1.83)	1.90 (1.89)	1.74 (1.75)
O <sub>3</sub> site	O <sub>2</sub> –V	1.85 (1.83)	1.85 (1.83)	1.90 (1.89)	1.74 (1.75)
	O <sub>3</sub> –H	0.99 (0.99)	0.99 (0.99)	0.99 (0.99)	
	O <sub>3</sub> –V	1.99 (2.02)	2.16 (2.31)	2.25 (2.30)	2.04 (2.04)
	O <sub>3</sub> –V	2.01 (2.00)	1.99 (1.99)	2.00 (2.00)	1.97 (1.97)
	O <sub>3</sub> –V	2.01 (2.00)	1.99 (1.99)	2.00 (2.00)	1.97 (1.97)
	H–O <sub>2</sub> <sup>b</sup>	3.00 (3.00)	2.32 (2.24)	2.29 (2.23)	

<sup>a</sup> Bilayer results are shown in parentheses brackets. Three types of optimization methods were employed: (I) *z* coordinates of OH group, (II) all three coordinates of OH group, and (III) full geometry including all atoms. <sup>b</sup> This atom is the closest neighbor of the O<sub>3</sub> site among dicoordinated oxygens.

methods were adopted: (I) optimization of the *z*-coordinates of the H atom and its corresponding adsorption site, i.e. the O–H group, constraining these two atoms in an orientation perpendicular to the surface plane, and in this case, other atoms of the adsorbate–substrate system are fixed and the geometry remains identical to that of the clean (010) surface; (II) optimization of all three coordinates of the O–H group using former equilibrium geometries in order to examine the influence of the local chemical environment of the adsorption site; (III) full optimization of all the constituent atoms of the adsorbate–substrate system, based on the geometries obtained by optimization II, to examine the contribution of lattice relaxation.

The adsorption energy (*E*<sub>ads</sub>) has been calculated according to the expression

$$E_{\text{ads}} = E_{(\text{adsorbate}+\text{substrate})} - (E_{\text{adsorbate}} + E_{\text{substrate}})$$

where *E*<sub>(adsorbate+substrate)</sub>, *E*<sub>adsorbate</sub>, and *E*<sub>substrate</sub> are the total energies of the adsorbate–substrate system, isolated adsorbate, and isolated substrate, respectively. A negative *E*<sub>ads</sub> value corresponds to a stable adsorbate–substrate system.

**3.1. Reactivity of Lattice Oxygens.** The results obtained from the three different optimization conditions show that all three lattice oxygens have high reactivity for H adsorption and that the O<sub>1</sub> site is found to be the most active (Table 5), which is in line with the experimental identification.<sup>6–12</sup>

Although there are rather significant differences in reactivity of these oxygen sites after optimization I, these differences cannot reflect the realistic situation precisely, since the contribution of neither local environment nor lattice relaxation was taken into account. After the optimization II, the reactivity of O<sub>1</sub> and O<sub>3</sub> atoms is found to be very close. In particular, no significant difference is found in the reactivity of the three types of oxygens after full geometry optimization III (see Table 5). It was confirmed that the O<sub>1</sub> site acts as the active site in an unsupported V<sub>2</sub>O<sub>5</sub> catalyst. Regarding the supported catalyst, however, the situation is different more or less. Since the support and/or promoter of catalyst may change the geometric and electronic structures of V<sub>2</sub>O<sub>5</sub>, which greatly affect the adsorption state, and considering that no significant adsorption energy difference at the three O sites of unsupported V<sub>2</sub>O<sub>5</sub> is found, one can say that the reactivity ordering of these oxygens may be changed for supported V<sub>2</sub>O<sub>5</sub> catalyst in some cases.

Now the results obtained from all the three optimization conditions are compared. Between optimizations I and II, there is no significant difference for H adsorption at the O<sub>1</sub> site, as



**TABLE 4: Selected Mulliken Atomic Charges (au) Calculated by Using Monolayer and Bilayer SUC of  $V_2O_5(010)^a$** 

adsorption sites	atoms	I	II	III	clean surface
O <sub>1</sub> site	H	+0.303 (+0.296)	+0.303 (+0.291)	+0.300 (+0.296)	
	O <sub>1</sub>	-0.490 (-0.491)	-0.490 (-0.487)	-0.511 (-0.510)	-0.324 (-0.305)
	V	+1.279 (+1.266)	+1.272 (+1.268)	+1.275 (+1.240)	+1.272 (+1.236)
O <sub>2</sub> site	H	+0.394 (+0.394)	+0.391 (+0.392)	+0.388 (+0.386)	
	O <sub>2</sub>	-0.607 (-0.652)	-0.604 (-0.652)	-0.567 (-0.577)	-0.602 (-0.596)
	V	+1.195 (+1.259)	+1.195 (+1.258)	+1.180 (+1.252)	+1.272 (+1.335)
O <sub>3</sub> site	V	+1.195 (+1.259)	+1.195 (+1.258)	+1.180 (+1.252)	+1.272 (+1.335)
	H	+0.440 (+0.439)	+0.431 (+0.414)	+0.421 (+0.413)	
	O <sub>3</sub>	-0.670 (-0.652)	-0.643 (-0.604)	-0.621 (-0.590)	-0.646 (-0.619)
	V	+1.137 (+1.159)	+1.200 (+1.228)	+1.280 (+1.251)	+1.272 (+1.236)
	V	+1.180 (+1.289)	+1.128 (+1.208)	+1.153 (+1.237)	+1.272 (+1.335)
	V	+1.180 (+1.289)	+1.128 (+1.208)	+1.153 (+1.237)	+1.272 (+1.335)
	O <sub>2</sub> <sup>b</sup>	-0.576 (-0.630)	-0.635 (-0.691)	-0.673 (-0.671)	-0.602 (-0.596)

<sup>a</sup> Bilayer values are listed in parenthesesbrackets. The optimization conditions used were described in Table 3. <sup>b</sup> See Table 3.

**TABLE 5: Adsorption Energies of H Atom at O<sub>1</sub>, O<sub>2</sub>, and O<sub>3</sub> Oxygen Sites Present in Monolayer SUC of  $V_2O_5(010)$  Calculated by the Aforementioned Optimization Methods**

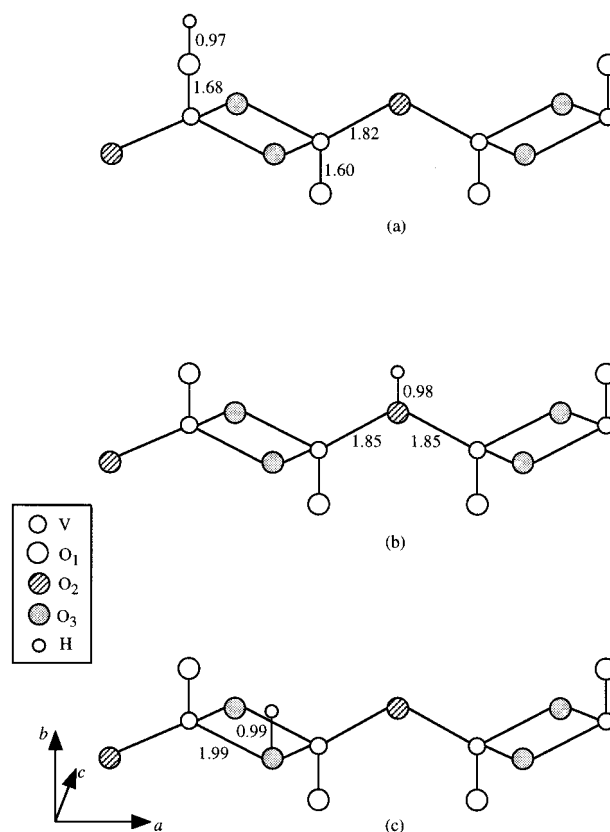
optimization methods	I	II	III
$E_{\text{ads}}, \text{O}_1$ (kcal/mol)	-47.1	-46.4	-62.2
$E_{\text{ads}}, \text{O}_2$ (kcal/mol)	-39.9	-39.0	-60.3
$E_{\text{ads}}, \text{O}_3$ (kcal/mol)	-32.0	-45.8	-59.3

well as at the O<sub>2</sub> site. This indicates that there is almost no effect of local environment at both O<sub>1</sub> and O<sub>2</sub> sites on their reactivity. However, the O<sub>3</sub> site exhibits completely different behavior. For optimization condition II, H atom binding to the O<sub>3</sub> site also interacts with its nearest O<sub>2</sub> site with a hydrogen bond. As a consequence, the adsorption energy is negatively increased by 13.8 kcal/mol. It shows that the local chemical environment of the O<sub>3</sub> site affects both geometry and energy.

On the other hand, a remarkable difference is observed after full geometry optimization at both O<sub>1</sub> and O<sub>2</sub> sites. A negative increment of adsorption energy of 15.8 and 21.3 kcal/mol was calculated, respectively. These comparatively higher energy differences show the significant contribution of lattice relaxation on adsorption systems with respect to both the O<sub>1</sub> and O<sub>2</sub> sites. Similarly, full geometry optimization for H adsorption at the O<sub>3</sub> site also leads to a negative increment of adsorption energy by 13.5 kcal/mol where the surface relaxation is the main contributor.

**3.2. Geometric Structures.** Equilibrium geometries of the adsorption systems obtained by the respective optimization methods are shown in Figures 2–4. The selected bond lengths between adsorption sites and their closest neighbors are shown in Table 3. It is observed that the surface relaxation has significant effect. At all the three optimizations, equilibrium bond lengths of the O–H group increase in the order O<sub>1</sub>–H < O<sub>2</sub>–H < O<sub>3</sub>–H, indicating that the bonding of O<sub>1</sub>–H is the strongest, which further corresponds to the highest reactivity of the vanadyl oxygen. On the other hand, no significant change in the O<sub>1</sub>–H bond length is observed after all the three types of optimizations. The same is true for both O<sub>2</sub>–H and O<sub>3</sub>–H bonds. This justifies again that lattice relaxation contributes significantly to the adsorption energy and not to the newly forming O–H bonds.

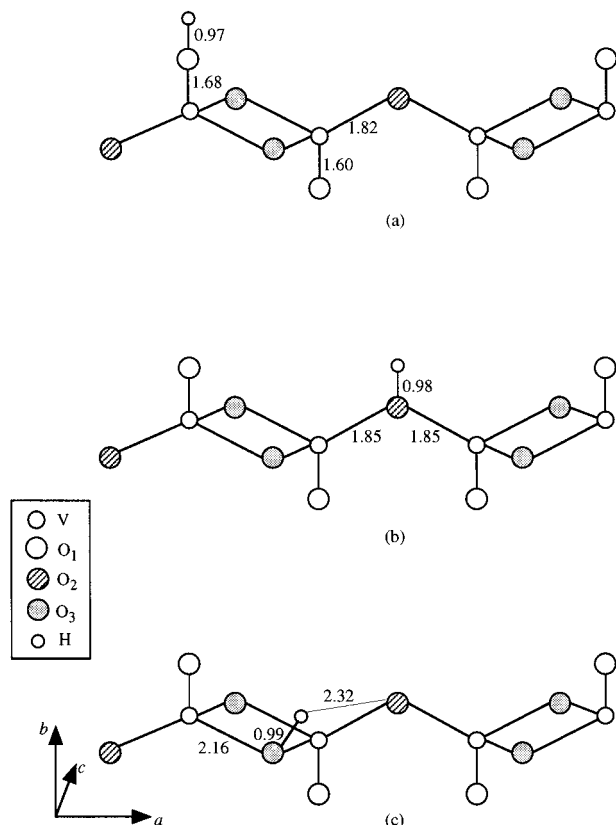
As described above, geometry optimization with step II for the O<sub>3</sub> site shows that a hydrogen bond is formed between the H and its closest neighbor O<sub>2</sub> atom with a distance of 2.32 Å (Figure 3c). This H-bond almost remains unaltered after the full geometry optimization (step III), which indicates that the lattice relaxation has no pronounced influence on the H-bond (Figure 4c) either. To form this hydrogen bond, one V–O<sub>3</sub> bond has been elongated significantly while the other two identical V–O<sub>3</sub>

**Figure 2.** Geometries of hydrogen adsorption at (a) O<sub>1</sub>, (b) O<sub>2</sub>, and (c) O<sub>3</sub> oxygen sites on  $V_2O_5(010)$  obtained by optimizing the *z* coordinates of OH species. The values shown here denote the bond lengths (Å).

bonds almost have not shown any significant changes (Table 3). In contrast, the local environment of either the O<sub>1</sub> or the O<sub>2</sub> site hardly affects the geometry.

**3.3. Electronic Structures.** Major atomic charges of the respective adsorption systems are shown in Table 4. For all the optimization conditions, the charge on the H atom increases with the order of the corresponding adsorption site: O<sub>1</sub> < O<sub>2</sub> < O<sub>3</sub>, which shows the trend of their depleting electron from H atom.

In the step I optimization, adsorption sites accumulate the electrons; the most significant accumulation is found at the O<sub>1</sub> site. The V atoms bonded to the O<sub>2</sub> and O<sub>3</sub> sites are observed to be reduced, whereas the V atom bound to the O<sub>1</sub> site has a charge almost identical to that of the clean surface, and the other vanadiums are reduced rather significantly instead.

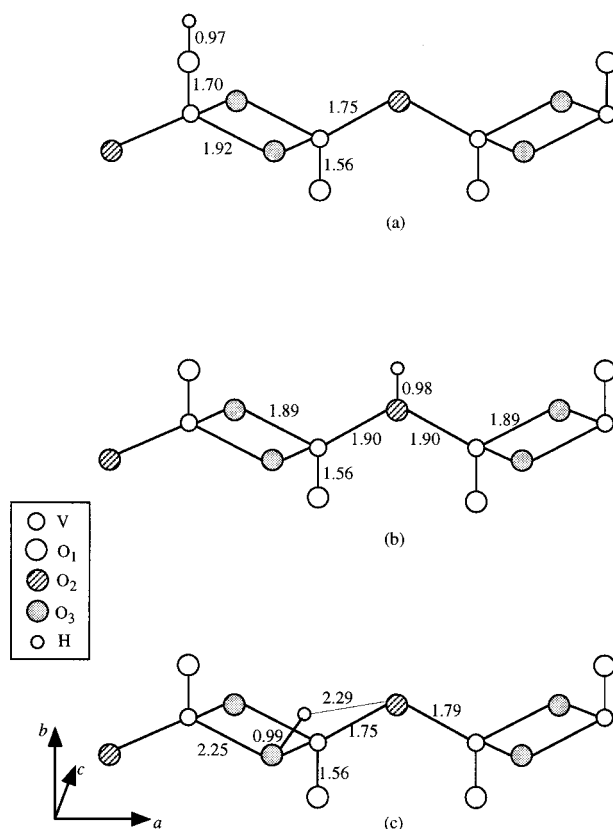


**Figure 3.** Geometries of hydrogen adsorption at (a) O<sub>1</sub>, (b) O<sub>2</sub>, and (c) O<sub>3</sub> oxygen sites on  $\text{V}_2\text{O}_5(010)$  obtained by optimizing all three coordinates of OH species. The values shown here denote the bond lengths (Å).

Electronic structures obtained by optimization II are compared with that of optimization I. Atomic charges on both the O<sub>1</sub> and O<sub>2</sub> sites remain almost the same under both conditions. In the case of O<sub>3</sub>, however, a significant change is observed. At this site, charges on the H and O<sub>3</sub> atoms are decreased, charges on one V atom is increased by 0.063 e, while those on the other two V atoms that are located in the same local environment are decreased by 0.052 e. Since the H-bond was formed between the H bound to O<sub>3</sub> site and its closest O<sub>2</sub> atom, the charge on this O<sub>2</sub> increased correspondingly.

At the full optimization (step III), charge redistribution occurs when compared to that of the former two optimization methods. Charges on both the O<sub>2</sub> and O<sub>3</sub> sites are depleted remarkably; however, the O<sub>1</sub> site has accumulated more charge. Our previous investigation of the  $\text{V}_2\text{O}_5(010)$  surface<sup>29</sup> has revealed that the vanadyl oxygen is the least ionic among these three oxygens, which shows that the O<sub>1</sub> site still has potential to gain extra electrons, and hence, the O<sub>1</sub> site exhibits a different behavior. Simultaneously, the charge on the V atom beneath the adsorption site O<sub>1</sub> is almost identical to that of the clean surface, while the other three V atoms get reduced instead. On the other hand, the two V atoms bonded to the O<sub>2</sub> site and the two V atoms closer to the O<sub>3</sub> site are observed to get reduced very significantly in the cases of adsorption at O<sub>2</sub> and O<sub>3</sub> sites, respectively.

**3.4. Vibrational Frequencies of the Surface Hydroxyl Species.** The spectra of hydroxyl groups raise a question regarding the location of the groups. What is the nature of the local surface environment responsible for their differences in frequency and in chemical behavior?<sup>61</sup> As catalytically important Brønsted acid sites, hydroxyl groups play a key role in the catalysis process. Bond et al.<sup>62</sup> revealed that the vibrational



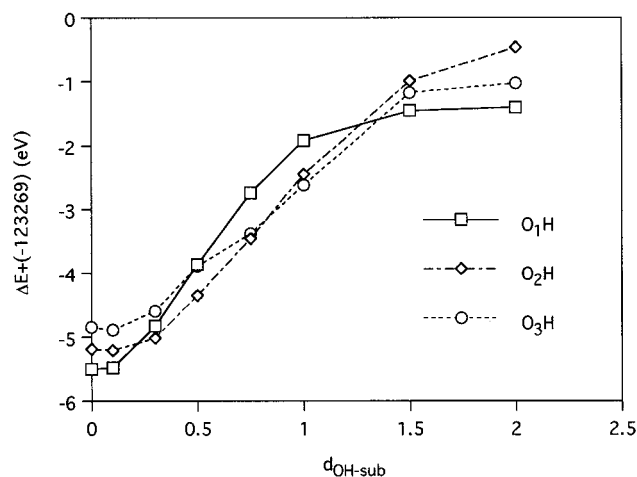
**Figure 4.** Fully optimized geometries of hydrogen adsorption at (a) O<sub>1</sub>, (b) O<sub>2</sub>, and (c) O<sub>3</sub> oxygen sites on  $\text{V}_2\text{O}_5(010)$ . The values shown here denote the bond lengths (Å).

frequency of surface OH groups on  $\text{V}_2\text{O}_5$  ranges from 3300 to 3700  $\text{cm}^{-1}$  by means of IR study. Although some authors<sup>9,63</sup> pointed out that the surface hydroxyl species connected to the O<sub>1</sub> atom shows a frequency of 3648 or 3645  $\text{cm}^{-1}$ , there are very few reports with respect to the frequencies of O<sub>2</sub>H and O<sub>3</sub>H species in the literature. To elucidate their difference, harmonic vibrational frequencies are evaluated on the fully optimized geometries. Since the anharmonic vibrational frequency is not included, overestimation of the frequency is reasonable.

Our calculated vibrational frequencies of the O<sub>1</sub>H, O<sub>2</sub>H, and O<sub>3</sub>H groups are 3792.17, 3658.26, and 3561.13  $\text{cm}^{-1}$ , respectively, which indicates that the least coordinated hydroxyl of O<sub>1</sub>H shows the highest frequency value and that the most coordinated hydroxyl of O<sub>3</sub>H exhibits the lowest one. This trend is very consistent with that of the hydroxyl species present in the alumina surface.<sup>64</sup> Also, these values are almost within the range of the observation of Bond and co-workers.<sup>62</sup> In particular, the calculated frequency of the O<sub>1</sub>H group is very close to the experimental values,<sup>9,63</sup> which again justified our methodology. Furthermore, it is found that these frequencies correlate with the reactivity of the surface oxygens as well as with the bond strengths of the hydroxyl species.

### 3.5. Desorption Ability of the Surface Hydroxyl Species.

The reactivity of surface oxygens depends not only on their adsorption ability but also on their desorption ability of the hydroxyl group from the surface to form the OH<sup>-</sup> species and O vacancy. A series of calculations were performed to evaluate the desorption ability. The obtained potential curves describing the desorption of surface OH species obtained by optimization condition I are shown in Figure 5. The OH desorption process was investigated by keeping the surface and OH group fixed at



**Figure 5.** Potential curves describing the desorption of surface OH species obtained by optimizing the  $z$  coordinates of OH group.  $d_{\text{OH-sub}}$  represents the distance between the hydroxyl O atom and its equilibrium position of 0.0 Å. The value of  $[\Delta E + (-123\,269)]$  (eV) corresponds to the total energy of adsorption systems.

their equilibrium coordinates, respectively, and varying only the position of the OH species particularly on the surface. The potential was described by the distance  $d_{\text{OH-sub}}$  between the hydroxyl O atom and its equilibrium position. The results (Figure 5) show that the  $\text{O}_1\text{H}$  and  $\text{O}_3\text{H}$  species essentially reach their free positions where they are out of the interaction with the (010) surface when the distance reached is 2.0 Å, whereas the  $\text{O}_2\text{H}$  species still has significant interaction with the surface. From their equilibrium of 0.0 Å to the position of 2.0 Å, energy differences are found to be 4.100, 4.727, and 3.819 eV corresponding to  $\text{O}_1\text{H}$ ,  $\text{O}_2\text{H}$ , and  $\text{O}_3\text{H}$ , respectively. These correspond to the minimum energies that should be exerted to remove the respective OH species from the equilibrium adsorption system. These values demonstrate that removal of the  $\text{O}_3\text{H}$  species is the most favorable energetically in the case of optimization condition I, but practically, it is necessary to overcome the energy of ca. 0.6 eV contributed by the H-bonding, as discussed above. Thus, removal of the  $\text{O}_1\text{H}$  species from the (010) surface is found to be the cheapest energetically if one considers the effect contributed by the local environment of adsorption site.

It was found that H adsorption, OH desorption, and O vacancy formation are correlated processes. This investigation indicates that the vanadyl oxygen not only has the strongest binding ability on H atom but also the greatest desorption ability of hydroxyl group to form the  $\text{OH}^-$  species and oxygen vacancy.

#### 4. Conclusions

We performed the first-principles DFT calculations under the periodic boundary condition to rationalize the active oxygen sites of  $\text{V}_2\text{O}_5$ , as well as to study the influence of local environment and lattice relaxation on reactivity. The major conclusions obtained from the calculations are as follows. Vanadyl oxygen (singly coordinated  $\text{O}_1$ ) acts as the most active site toward hydrogen adsorption among the three kinds of lattice oxygens present in unsupported  $\text{V}_2\text{O}_5(010)$ , in good agreement with many experimental observations. Lattice relaxation contributes to the reactivity of the surface oxygens significantly in all cases, while the local environment of the adsorption site affects both geometry and reactivity of the adsorption system only for tricoordinated oxygen. Hydrogen adsorption reduces the surface, whereas the adsorption at  $\text{O}_1$  shows different

behavior; the  $\text{O}_1$  site accumulates much charge, and the V atom beneath the  $\text{O}_1$  is almost identical to that of clean surface, while the other three V atoms get reduced. The calculated harmonic vibrational frequencies reveal the difference among the three types of hydroxyl species and correlate with bond strengths as well as support their reactivity. When the influence of the local environment of the adsorption site is considered, desorption of the  $\text{O}_1\text{H}$  species from  $\text{V}_2\text{O}_5(010)$  is the cheapest energetically. It is found that the processes of H adsorption, OH desorption, and O vacancy formation are correlated. The present study indicates that the periodic approach rids the artifacts of the cluster method and that the periodic first-principles DFT methodology is found to be reliable for investigating not only geometric and electronic structures but also the reactivity of transition metal oxide catalysts.

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