



## H<sub>2</sub>, Ne, and N<sub>2</sub> Energies of Encapsulation into C<sub>60</sub> Evaluated with the MPWB1K Functional

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Received December 27, 2005

**Abstract:** The recently suggested MPWB1K functional is tested on H<sub>2</sub>@C<sub>60</sub>, Ne@C<sub>60</sub>, and N<sub>2</sub>@C<sub>60</sub> as a tool for evaluations of stabilization energies upon encapsulation of nonmetallic species into fullerenes. It is found that the MPWB1K (modified Perdew–Wang and Becke functionals) values can be within a few kilocalories per mole from the MP2 or SCS-MP2 (spin-component scaled MP2) values so that further applications of the functional are clearly encouraged. The best estimates of the encapsulation-energy gains found for H<sub>2</sub>@C<sub>60</sub>, Ne@C<sub>60</sub>, and N<sub>2</sub>@C<sub>60</sub> are at least 4 kcal/mol, slightly less than 4 kcal/mol, and about 9 kcal/mol, respectively.

### 1. Introduction

The suitability of density-functional theory (DFT) for evaluations of weak molecular interactions and especially of dispersion energies has been an important computational issue.<sup>1–10</sup> Very recently, Zhao et al.<sup>11–15</sup> performed a series of test DFT calculations with a conclusion<sup>15</sup> that the MPWB1K functional (the modified Perdew and Wang exchange functional MPW<sup>16</sup> and Becke's meta correlation functional<sup>17</sup> optimized against a kinetics database) is the best combination for evaluations of nonbonded interactions with a relative averaged mean unsigned error of only 11%.

Fullerene encapsulations of nonmetal atoms and small molecules represent an interesting case<sup>18–22</sup> for this type of computation because there are also some related observations. A well-known example is represented by complexes of fullerenes with encapsulated rare gas atoms<sup>23–27</sup> prepared by heating under high pressure. Ion bombardment can produce N@C<sub>60</sub><sup>28,29</sup> and P@C<sub>60</sub><sup>30</sup> (and also Li@C<sub>60</sub><sup>31</sup>), though in very low yields. N@C<sub>60</sub> and its derivatives have been studied vigorously,<sup>32–36</sup> inter alia as a candidate for the implementation of qubits in an electron-spin-based quantum computer.<sup>37</sup>

N<sub>2</sub>@C<sub>60</sub> was prepared by pressure heating<sup>38</sup> and was also reported<sup>39,40</sup> in the chromatographic separation after the nitrogen-ion implantation into C<sub>60</sub>. Recently, molecular hydrogen<sup>41–43</sup> (and also water molecule<sup>21</sup>) was placed inside an open-cage fullerene, and the cage with H<sub>2</sub> was subsequently closed.<sup>44</sup>

Previous DFT calculations on nonmetallic fullerene encapsulates used mostly<sup>45–52</sup> the B3LYP functional. The present paper reports calculations of H<sub>2</sub>@C<sub>60</sub>, Ne@C<sub>60</sub>, and N<sub>2</sub>@C<sub>60</sub> with the newly suggested MPWB1K functional in order to judge its performance for such encapsulation energetics.

### 2. Calculations

Three systems known from observations were selected for the present calculations: H<sub>2</sub>@C<sub>60</sub>, Ne@C<sub>60</sub>, and N<sub>2</sub>@C<sub>60</sub>. The computations are carried out at fixed experimental geometries. The C<sub>60</sub> geometry is taken from the gas-phase electron diffraction<sup>53</sup> where the 5/6 (i.e., between a pentagon and a hexagon) and 6/6 bonds (between two hexagons) were determined as 1.458 and 1.401 Å, respectively. The bond lengths of H<sub>2</sub> and N<sub>2</sub> are also from observed data.<sup>54,55</sup> The diatomic molecules in the cage are oriented toward a pair of its parallel pentagons so that the complexes exhibit *D*<sub>5d</sub> symmetry. This arrangement is known from previous calculations<sup>52</sup> as the lowest energy structure, though energy

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**Table 1.** Encapsulation Energies Computed in the Observed Geometries (kcal/mol)

level	H <sub>2</sub> @C <sub>60</sub>	Ne@C <sub>60</sub>	N <sub>2</sub> @C <sub>60</sub>
B3LYP/6-31G**	1.56	-2.52	7.70
B3LYP/6-311G(2d, 2p)	1.34	-2.95	7.94
MPWB1K/6-31G**	-7.93 <sup>a</sup>	-3.08	-14.93
MPWB1K/6-31G** & BSSE	-7.29	-0.22	-11.34
MP2/6-31G**	-3.99	-5.05	-17.05
MP2/6-31G** & BSSE	-2.36	-1.85	-8.90
MPWB1K/6-311G(2d, 2p)	-8.09	-7.20	-13.18
MPWB1K/6-311G(2d, 2p) & BSSE	-7.17	-3.83	-9.54
MP2/6-311G(2d, 2p)	-8.63	-5.39	-23.44
MP2/6-311G(2d, 2p) & BSSE	-6.07		-15.54
SCS-MP2/6-311G(2d, 2p)	-6.71	-4.35	-16.35
SCS-MP2/6-311G(2d, 2p) & BSSE	-4.20		-8.57
MPWB1K/(d, p)-6-311G**	-9.44		
MPWB1K/(d, p)-6-311G** & BSSE	-7.81		
MP2/(d, p)-6-311G**	-10.62		
MP2/(d, p)-6-311G** & BSSE	-6.94		

<sup>a</sup> The value after the full geometry optimization is -8.37 kcal/mol.

differences relative to other orientations are less than 0.1 kcal/mol. In the case of Ne@C<sub>60</sub>, the Ne atom is placed in the cage center so that the system still exhibits icosahedral symmetry.

Single-point energy calculations were carried out with the MPWB1K functional in three well-known basis sets: 6-31G\* (6-31G\*\* for H atoms), 6-311G(2d,2p), and the augmented (d,p)-6-311G\*\* basis set.<sup>56</sup> For comparison, some computations were also performed with Becke's<sup>57</sup> three-parameter exchange functional combined with the nonlocal Lee-Yang-Parr<sup>58</sup> correlation functional (B3LYP). The ultrafine grid in numerical integrations of the DFT functional and the tight self-consistent field convergency criterion were used. However, the comparisons are primarily related to the second-order Møller-Plesset (MP2) perturbation treatment<sup>59,60</sup> with the frozen core option carried out in the above three basis sets. Moreover, the conventional MP2 method is also complemented with the spin-component-scaled (SCS) MP2 calculations.<sup>61</sup> Finally, the basis set superposition error (BSSE) was estimated by the Boys-Bernardi counterpoise method.<sup>62</sup> The selected basis sets represent a practical computational limit for calculations of systems comparable to C<sub>60</sub> at DFT and especially MP2 levels.

The Gaussian<sup>63</sup> and PQS<sup>64</sup> program packages were used for all calculations.

### 3. Results and Discussion

Table 1 reports the zero-point-exclusive encapsulation energies (also called the energies of interaction or stabilization) corresponding to the change in potential energy for the process:



where X is H<sub>2</sub>, Ne, or N<sub>2</sub>. To compare this to an experimental term, one would have to add the changes in zero-point energy and heat content function. Let us note that the quantity computed here is a thermodynamic one, although the kinetics of reaction 1 are also of interest.<sup>65-67</sup>

As already known from previous computations,<sup>21,52</sup> the B3LYP functional does not produce reliable values for nonbonding fullerene encapsulations. In the cases of H<sub>2</sub>@C<sub>60</sub> and N<sub>2</sub>@C<sub>60</sub>, their stabilization energies become positive at the B3LYP/6-31G\*\* and B3LYP/6-311G(2d,2p) levels, which does not agree well with the very fact that the endohedrals can be observed. Only for Ne@C<sub>60</sub> is the B3LYP/6-31G\* encapsulation energy negative (binding), though the stabilization energy is too low compared to other calculations in Table 1 and also to previous evaluations.<sup>20,68</sup>

All of the MPWB1K and MP2 calculations in Table 1 show for H<sub>2</sub>@C<sub>60</sub> a substantial stabilization energy. Obviously, the BSSE-corrected energies should be preferred, though it is not necessarily assured<sup>69</sup> that the true energy is closer to the BSSE-corrected term rather than to the uncorrected term. The MPWB1K and MP2 treatments with the 6-31G\*\* basis set differ by some 4 kcal/mol before the BSSE correction. With respect to the remaining computations of H<sub>2</sub>@C<sub>60</sub> in Table 1, it seems however likely that the MP2/6-31G\*\* term underestimates the stabilization for the species. Upon moving to the 6-311G(2d,2p) and (d,p)-6-311G\*\* basis sets, the MPWB1K and MP2 values differ only by about 1 kcal/mol. Application of the SCS approach tested for the 6-311G(2d,2p) basis set suggests, however, a reduction of the stabilization by about 2 kcal/mol compared to the conventional MP2 method. Overall, it can be stated that the encapsulation of H<sub>2</sub> into C<sub>60</sub> brings a potential energy gain of at least 4 kcal/mol.

The effect of geometry optimizations was also checked with H<sub>2</sub>@C<sub>60</sub>. It turns out that, if the full geometry optimization is carried out at the MPWB1K/6-31G\*\* level, the encapsulation energy is changed only by some 0.44 kcal/mol. The perturbation of the cage upon encapsulation changes the C-C bonds only on the fourth digit. The 5/6 C-C bonds in the cage optimized at the MPWB1K/6-31G\*\* level vary between 1.4377 and 1.4388 Å, while for the 6/6 bonds, it is from 1.3792 to 1.3798 Å. It should also be noted that rotating the H<sub>2</sub> molecule toward two hexagons (instead of toward two pentagons) changes the energy by less than 0.1 kcal/mol so that the molecule can indeed nearly freely rotate inside the cage.<sup>70</sup>

Let us move now to the Ne@C<sub>60</sub> system. With the 6-31G\*\* basis set, the MPWB1K and MP2 terms differ by some 2 kcal/mol both before and after the BSSE correction, Ne@C<sub>60</sub> showing up as more stable in the MP2/6-31G\* treatment. However, with the 6-311G(2d,2p) basis set, the MPWB1K functional produces greater stabilization compared to the conventional MP2 value, which is further reduced by about 1 kcal/mol upon application of the SCS approach. Although the BSSE computations could not be finished in this case owing to computational difficulties, Ne@C<sub>60</sub> is likely stabilized by slightly less than 4 kcal/mol.

The differences between the conventional MP2 and SCS-MP2 treatments are more significant in the N<sub>2</sub>@C<sub>60</sub> system, amounting to some 7 kcal/mol. However, if the MPWB1K/6-311G(2d,2p) and SCS-MP2/6-311G(2d,2p) values are compared, they are still within 3 kcal/mol before the BSSE correction and within 1 kcal/mol after correcting for BSSE. The best estimate gives about a 9 kcal/mol gain in the

potential energy upon  $\text{N}_2@C_{60}$  formation. Interestingly enough, the MP2/6-31G\* encapsulation energy without the BSSE correction for  $\text{N}@C_{60}$  was computed<sup>50</sup> to be  $-7.3$  kcal/mol, that is, roughly one-half of the value found here at the same level for  $\text{N}_2@C_{60}$ .

While the performance of the MPWB1K exchange-correlation potential is reasonably good for the inclusion compounds considered here, it does not perform well for  $\pi$  stacking. For instance, the binding energy predicted for the prototype benzene dimer, in the parallel displaced configuration, is only 0.11 kcal/mol after counterpoise correction, using the 6-311+G(2df,2pd) basis set. The best current estimate is<sup>71,72</sup> 2.78 kcal/mol. For the sandwich conformer, the counterpoise corrected MPWB1K potential is repulsive, although the best estimate of this quantity is 1.81 kcal/mol. Zhao and Truhlar have successfully used<sup>13a</sup> an empirical combination of density-functional and wave-function-based correlation methods to calculate this quantity accurately at a reasonable cost, and they also suggested<sup>13b</sup> an improved version of MPWB1K called PWB6K, with a good performance for stacking interactions. Although the benzene dimer issue is not directly related to the fullerene encapsulation energies, it indicates the fundamental limitations of local or semi-local DFT to describe the dispersion interaction.

It can be concluded that the values for nonmetallic encapsulations into  $C_{60}$  evaluated with the MPWB1K functional closely follow the results from the MP2 and SCS-MP2 treatments with the same basis set. This finding suggests the MPWB1K functional as a practical tool for the evaluation of encapsulations of small molecules into fullerenes and model nanotubes, though larger basis sets<sup>73,74</sup> and also other functionals<sup>74,75</sup> are to be still tested, especially for specific cases<sup>76,77</sup> like interactions with metals.

**Acknowledgment.** This research has been supported by a grant-in-aid for the NAREGI Nanoscience Project and for Scientific Research on Priority Area (A) from the Ministry of Education, Culture, Sports, Science and Technology of Japan. P.P. would like to thank the Institute for Molecular Science for a visiting professorship. The constructive comments from the reviewers are highly appreciated, too.

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