

Enhancement of 5D_0 – 7F_J Emissions of Eu^{3+} Ions in the Vicinity of Polymer-Protected Au Nanoparticles in Sol–Gel-Derived B_2O_3 – SiO_2 Glass

Tomokatsu Hayakawa,* Kazunori Furuhashi, and Masayuki Nogami

Ceramic Division, Department of Materials Science and Engineering, Nagoya Institute of Technology, Gokiso, Showa, Nagoya 466-8555, Japan

Received: April 22, 2004

The paper reports the enhanced photoluminescence of trivalent europium (Eu^{3+}) ions in the vicinity of nanometer-sized metal particles in glasses. In this study, gold (Au) colloids synthesized in solution were introduced into a borosilicate (B_2O_3 – SiO_2) glass matrix together with Eu^{3+} ions via a sol–gel route. To control in size and protect Au colloids two types of polymers—poly(vinylpyrrolidone) (PVP) and polyacrylonitrile (PAN)—were used. As a result, PVP provided excellent homogeneity of Au nanoparticles in the hybrid gel. A novel finding on its photoluminescence properties is that a heat treatment of the gel at 400 °C led to a great enhancement of Eu^{3+} fluorescence under a long ultraviolet excitation, six times higher than without Au nanoparticles, accompanied with a broad fluorescence around 400 nm due to carbonyl moieties of PVP. On the contrary, when the carbonyl groups were removed by heat treatment at temperatures higher than 400 °C, a conspicuous decrease in the relative intensity of Eu^{3+} fluorescence was observed, regardless of the presence of Au nanoparticles. Since a group of carbonyl terminations, when coordinated to Eu^{3+} ions ($>\text{C}=\text{O} \rightarrow \text{Eu}^{3+}$), can be an energy donor, these observations showed that Eu^{3+} ions were strongly excited by the energy transfer from $\text{C}=\text{O}$ groups, which were placed in enhanced local fields near Au nanoparticles due to higher orders of surface plasmon resonances.

1. Introduction

Surface plasmon resonance (SPR) of metal nanoparticles, which is of importance as one of the elemental excitations in solid-state quantum physics,^{1,2} provides not only a great enhancement of optical third-order nonlinear susceptibility of metal (“inside-field effect”)^{3–5} but also an increase in the local electromagnetic field strength around a nanometer-sized metal (“outside-field effect”).^{6–8} So far we are focused on the latter effect, and the enhancement of trivalent europium (Eu^{3+}) fluorescence was previously reported.^{9–11} In several of our research studies, silver (Ag^+) and Eu^{3+} ions were co-doped in silica glass by a sol–gel method and the following heat-treatment in a reduction atmosphere gave rise to the precipitation of Ag nanoparticles.

In this work a nanometer size of gold (Au) colloids in solution were incorporated into a borosilicate amorphous network of $10\text{B}_2\text{O}_3$ – 90SiO_2 together with Eu^{3+} ions via a sol–gel route. To control Au nanoparticles in size and obtain their good stability through the whole of synthesis, two different types of polymers (poly(vinylpyrrolidone) (PVP) and polyacrylonitrile (PAN))^{12,13} were utilized. The sizes of the particles in a starting solution and gel–glasses were monitored by optical extinction spectra and analyzed by Mie-Drude theory.^{14–16} One of the aims of this work is to investigate the effects of Au nanoparticles on Eu^{3+} fluorescence under photoexcitation of a long ultraviolet (UV) light, which is higher in energy than the first order (dipole) of SPR of Au nanoparticles. Influences of carbon-related phosphors originating from protective polymers of Au nanoparticles are also studied at various temperatures in heat-treatment for the fabricated gels of $10\text{B}_2\text{O}_3$ – 90SiO_2 doped with

Eu^{3+} ions and polymer-protected Au nanoparticles. It will be shown that Eu^{3+} fluorescence of 5D_0 – 7F_J transitions is enhanced in the presence of such Au nanoparticles and we shall discuss the enhanced Eu^{3+} fluorescence in terms of energy transfer from a carbonyl chromophore ($>\text{C}=\text{O}$) to Eu^{3+} ions in a local field enhancement near nanometer-sized Au particles due to higher orders of SPR excitations.

2. Experiment

Figure 1 shows a flowchart of the sol–gel synthesis for $10\text{B}_2\text{O}_3$ – 90SiO_2 gels containing both polymer-protected Au nanoparticles and Eu^{3+} ions. As a protective organic agent of Au nanoparticles, we selected two different types of organic polymers—PVP (K-90, $M_w \sim 360\,000$; Kishida Chem. Co.) and PAN ($M_w \sim 86\,200$; Aldrich), for which it is known to have affinitive moieties of a lone electron pair of tri-coordinated nitrogen and cyano-group $-\text{CN}$ to Au ions, respectively. The starting materials for the synthesis of nanometer-sized Au colloids in solution were HAuCl_4 (Kishida Chem. Co.), PVP (or PAN), and KBH_4 (Kishida Chem. Co.). For PVP-protected Au sol, 2.5 mL of 2.0 mM HAuCl_4 solution was added to a mixture of ethanol and distilled water with PVP (ethanol: 15 mL, water: 30 mL, PVP: 22.4 mg). For the reduction of AuCl_4^- , 2.5 mL of 20 mM KBH_4 solution was injected into the solution so that the color of the Au sol was turned from yellow to light red (A-PVP; see Figure 2a). ($[\text{vp}]/[\text{Au}] = 40$, where $[\text{vp}]$ = a molar amount of vinylpyrrolidone groups.) For PAN-protected Au sol, 11.2 mg of PAN was added to 45 mL of *N,N*-dimethylformamide (DMF; Nacalai Tesque, Inc.), which was mixed with 2.5 mL of 2.0 mM HAuCl_4 solution. Addition of 2.5 mL of 20 mM KBH_4 solution allowed PAN-protected Au sol to be formed (A-PAN; see Figure 2b). ($[\text{an}]/[\text{Au}] = 40$, where $[\text{an}]$ = a molar amount of $-\text{CH}_2-\text{CH}(\text{CN})-$.) $[\text{vp}]/[\text{Au}]$

* Corresponding author. E-mail: hayatomo@nitech.ac.jp.

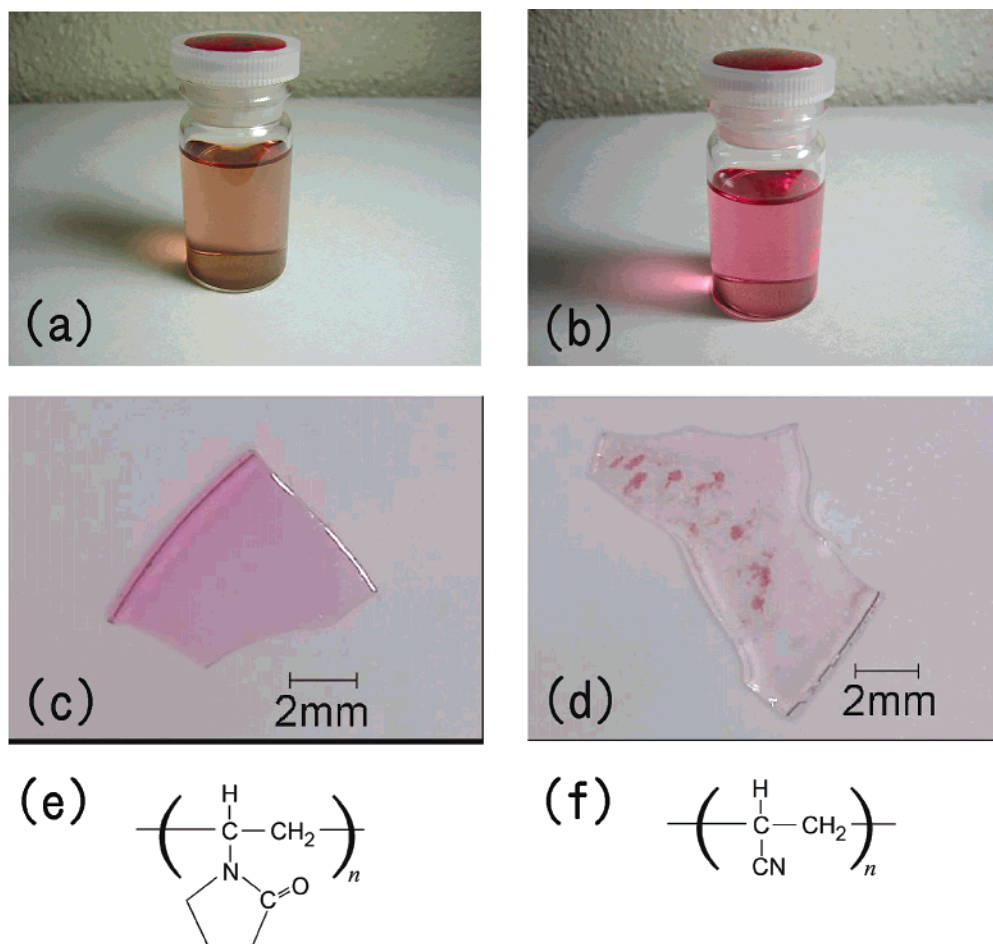


Figure 2. Photographs of (a) A-PVP sol, (b) A-PAN sol, and (c), (d) the fabricated borosilicate glasses incorporated with A-PVP and A-PAN, respectively, which were heated at 700 °C for 2 h in air; (e), (f) chemical structures of PVP and PAN, respectively.

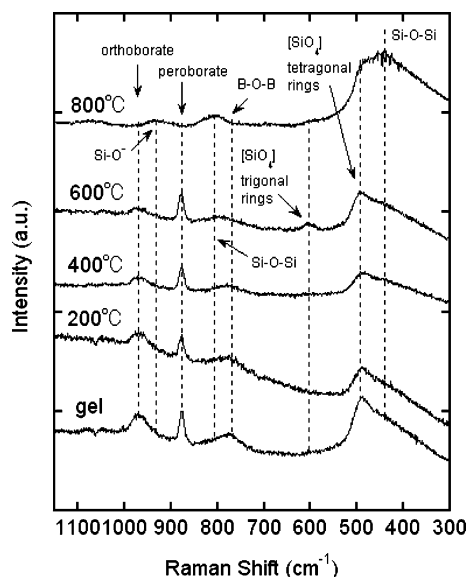


Figure 3. Raman spectra of sol-gel-derived $10\text{B}_2\text{O}_3$ - 90SiO_2 , containing PVP-protected Au nanoparticles, heated at various temperatures for 2 h in air.

(fwhm) $\Delta\omega$ for SPR is correlated with the relaxation time of conduction electrons τ through $\tau^{-1} \sim \Delta\omega$. As Doyle¹⁶ and Kreibig et al.²¹ proposed, for a spherical particle much smaller than the mean free path of the electrons in bulk material, l_∞ (~ 42 nm at 273 K for bulk Au),²² the collisions of the conduction electrons with the particle surfaces become important as an additional relaxation process, and the effective mean free

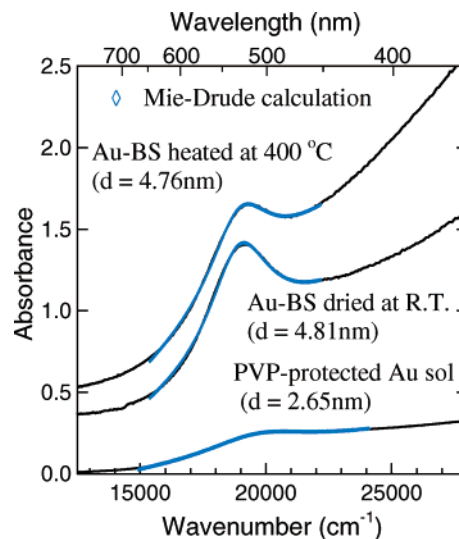


Figure 4. Optical absorption spectra of PVP-protected Au nanoparticles in sol and $10\text{B}_2\text{O}_3$ - 90SiO_2 gel-glass hybridized with PVP-protected Au nanoparticles (Au-BS).

path, l , is restricted by the representative length of the particle, $\sim R$. Thus, the relaxation time τ can be substituted for the collision time,

$$\tau_c = R/v_f \quad (2)$$

where v_f ($= 1.39 \times 10^8$ cm/s) is the Fermi velocity for Au. Consequently, the particle size of metal nanoparticle can be

estimated from the fwhm of Mie's surface plasmon resonance in optical absorption spectra through

$$R = v_f / \Delta\omega \quad (3)$$

Additionally, the Mie resonant frequency is given by $\omega_r = \omega_p / (1 + 2n^2)^{1/2}$, in the first approximation of free electron behavior, where ω_p is a bulk plasmon frequency. However, bulk plasmon frequency for noble metals such as Ag, Au, and Cu does not coincide with free electron plasma frequency ($N_e / em_{\text{opt}})^{1/2}$ (N_e , e , and m_{opt} represent the number, charge, and optical effective mass of conduction electrons, respectively), because of a tremendous contribution of the interband transition from the filled d band to the sp orbital. (The interband transition of Au is above 2.38 eV.^{22,23}) The frequency of the surface plasmon resonance of a noble metal nanoparticle is determined from the original Mie equation, $(\epsilon_1 + \epsilon_M)^2 + \epsilon_2^2 = \text{minimum}$.^{14–16} It is noted that the contribution of interband transition to the peak position is given not only the imaginary part ϵ_2 (energy dissipation) but also the real part ϵ_1 (polarization) of metal. As a consequence, the surface plasmon of Au is located around 2.36 eV (~ 525 nm), moreover slightly dependent on the refractive index of surrounding media through $\epsilon_M = n^2$. It can be seen that, in this case, the peak position of SPR is shifted from 526.8 nm in sol to 534.3 nm in gel and glass.

In Figure 4, the Mie-Drude fitting curves and estimated particle sizes $d = 2R$ from $R = v_f / \Delta\omega$ are also given. Initially, Au nanoparticles protected by PVP in solution have a small mean size of $d = 2.65$ nm; when incorporated with the B_2O_3 – SiO_2 gel matrix, the particle grows to a size of $d = 4.81$ nm, possibly because of the coalescence of initial nanoparticles. The heat treatment at 400 °C in air leads to the stiffer glass network and combustion of PVP phase (an exothermic peak due to combustion of PVP was observed at 400 °C in DTA-TG measurement). The particle size is estimated to be 4.76 nm. These data show that the protection of Au nanoparticles by the polymer works well in the gel-to-glass transformation and each of the Au nanoparticles also remains to be isolated, even in the heat treatment at 400 °C, though the initial particle size in sol is not kept in the incorporation into the B_2O_3 – SiO_2 gel matrix. Figure 5a displays a light-field image of TEM observation for the gel. The nanometer particles have a mean diameter of 6.4 nm and distribution to ± 2.3 nm (see Figure 5b). The estimated particle size from optical absorption spectra is smaller than that obtained directly from TEM pictures. Such a discrepancy has been reported in several studies, which is thought to be due to chemical interface effect (CIE).²⁴

3.3. Enhanced Eu^{3+} PL in the Vicinity of Au Nanoparticles under a Long UV Excitation. Figure 6 shows PL spectra of AuEu–BS and Eu–BS heated at 200–800 °C in air under a long UV excitation of 337.1 nm ($= 3.68$ eV). In the region between 570 and 700 nm the PL bands due to $^5\text{D}_0$ – $^7\text{F}_j$ ($j = 0, 1, 2, 3, \dots$) transitions of Eu^{3+} ions can be observed (the assignment is given in the figure). It is found that AuEu–BS shows a stronger Eu^{3+} fluorescence than Eu–BS except that heated at 600 °C in air. It should also be stressed that both AuEu–BS and Eu–BS contain the polymer of PVP and exhibit an additional broad luminescence around 400 nm. The peak position of the 400 nm-band is dependent on the heat treatment temperature, and when baked at temperatures higher than 800 °C the luminescence disappears (see Figure 7). The lifetime of the broad band is very fast and is found to be ~ 25 ns. The carbonyl group ($>\text{C}=\text{O}$) of PVP, which is one of the most luminescent organic parts due to the conjugation of π electrons, causes the observed broad luminescence.²⁵

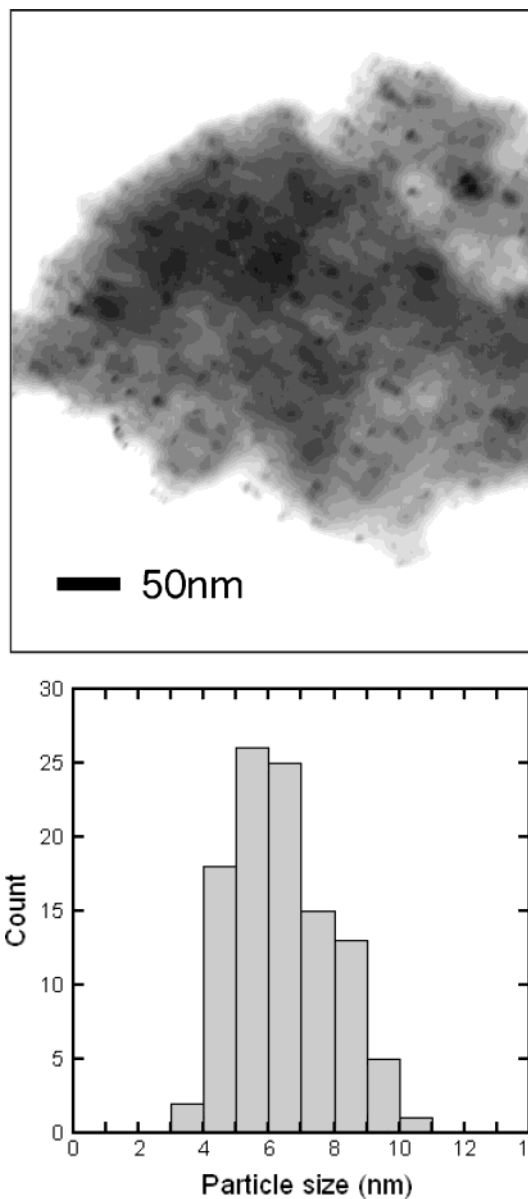


Figure 5. (a) Light-field image in TEM observation for the borosilicate gel containing PVP-protected Au nanoparticles (Au–BS), and (b) size-distribution of Au nanoparticles.

The excitation spectra monitored at 614 nm ($^5\text{D}_0$ – $^7\text{F}_2$) for AuEu–BS heated at 400 and 800 °C are also measured, which are shown in Figure 8, parts a and b, respectively. Spectrum (a) in Figure 8 exhibits a broad excitation band due to the excitation of carbonyl chromophores, but spectrum (b) does not. The broad band includes a direct process of excitation–emission of carbonyl chromophores. As for AuEu–BS heated at 800 °C, f – f transitions from $^7\text{F}_0$ to $^5\text{D}_{1-4}$, $^5\text{L}_6$, $^5\text{G}_{2,4}$, and $^5\text{H}_4$ levels of Eu^{3+} ions are clearly observed. In absorption spectra (not shown here), there is not observed a substantial difference between AuEu–BS and Eu–BS except the resonance peak due to the surface plasmon, if they are heated at the same temperature.

To evaluate the enhancement of $^5\text{D}_0$ – $^7\text{F}_j$ luminescence of Eu^{3+} ions, we define an enhancement factor f_E as

$$f_E = \frac{\sum_j I_j^{\text{AuEu-BS}}}{\sum_j I_j^{\text{Eu-BS}}} \quad (4)$$

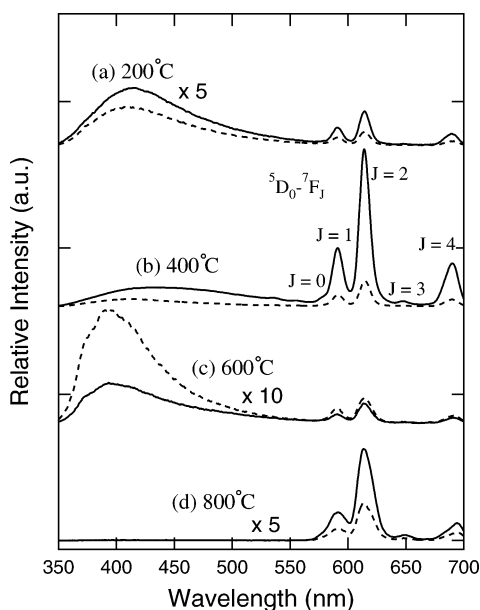


Figure 6. PL spectra of AuEu-BS (solid line) and Eu-BS (dashed line) samples heated at various temperatures for 2 h in air, which were recorded with a longer ICCD gate time of 5 ms.

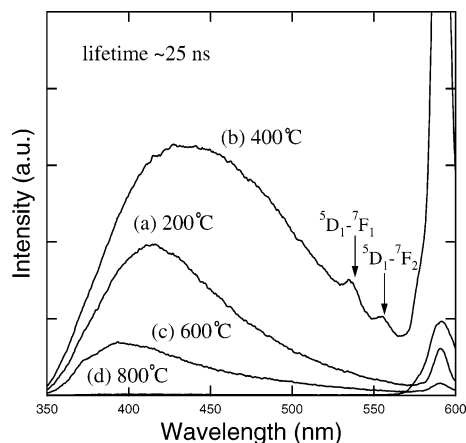


Figure 7. PL spectra of the broad band around 400 nm for AuEu-BS samples heated at various temperatures, which were recorded with a shorter ICCD gate time of 1 μ s.

where I_J denotes the integrated intensity of the 5D_0 - 7F_J band and the superscripts the sample names. Table 1 compiles f_E and the total emission intensities, $R_{AuEu-BS} = \sum J I_J^{AuEu-BS}$ and $R_{Eu-BS} = \sum J I_J^{Eu-BS}$, for AuEu-BS and Eu-BS, respectively, in comparison with data when heated at 200 °C for Eu-BS. When baked at 400 °C the highest enhancement of ~ 6 is observed. However, the samples of AuEu-BS and Eu-BS heated at 600 °C show smaller f_E of 0.76, and the relative intensities of $R_{AuEu-BS}$ and R_{Eu-BS} are decreased. According to Levy et al. (see Table 1 in ref 26), when the SiO₂ gel doped with Eu³⁺ ions was heated at elevated temperatures, the total emission intensities of 5D_0 - 7F_J transitions of Eu³⁺ ions was increased because of the removal of hydroxyl and organic groups from the gel. On the contrary, since in our study the PVP is incorporated for both AuEu-BS and Eu-BS, the relative PL intensity of 5D_0 - 7F_J is drastically decreased at 600 °C, as shown in Table 1. This can be explained in the way that organics coming from PVP are at most decomposed at 600 °C but still remain on pore walls of AuEu-BS and Eu-BS gels, which therefore de-excite Eu³⁺ ions in 5D_0 levels as killer sites, resulting in reducing PL intensity at 600 °C. However, if the gels are heated at temperatures higher than 600 °C, the residual

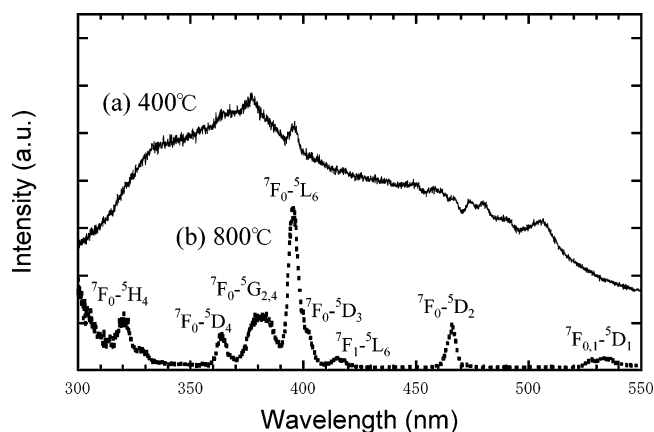


Figure 8. Excitation spectra monitored at 614 nm for AuEu-BS heated at (a) 400 and (b) 800 °C.

TABLE 1: Enhancement Factor f_E ($= R_{AuEu-BS}/R_{Eu-BS}$) and Total Emission Intensities, $R_{AuEu-BS} = \sum J I_J^{AuEu-BS}$ and $R_{Eu-BS} = \sum J I_J^{Eu-BS}$, for AuEu-BS and Eu-BS, respectively, for Various Heating Temperatures^a

temperature (°C)	f_E	$R_{AuEu-BS}$	R_{Eu-BS}	$R_{C=O}$
200	2.56	2.56	1.0	1.0
400	6.08	58.9	9.68	2.35
600	0.76	0.77	1.01	0.33
800	2.45	6.91	2.82	

^a These data are normalized by that of Eu-BS heated at 200 °C. The relative intensity $R_{C=O}$ of the broad band emission around 400 nm due to carbonyl groups, in comparison with that when heated at 200 °C, for various heat-treatment temperatures.

organics are completely removed and the PL intensities are increased again to the values about third times higher than at 200 °C. Additionally, f_E is also increased to ~ 2.45 . Because of the complete removal of PVP at the temperature, the factor f_E of 2.45 shows a pure enhancement due to Au nanoparticles located in the vicinity of Eu³⁺ ions.

4. Discussion

More recently, an energy transfer process from carbonyl groups to Eu³⁺ ions in an organo-silicate gel is reported,²⁷ where the luminescence efficiency of Eu³⁺ ions is greatly affected by forward and backward energy transfers between them. From our PL intensity data (Figure 6) and excitation spectra data (Figure 8) for various heating temperatures, a probable explanation of the enhancement of Eu³⁺ fluorescence when sintered at low temperatures is given by the combination of enhanced local fields in the vicinity of Au nanoparticles and the energy transfer process from excited C=O groups to Eu³⁺ ions in the ground state. The former causes the increasing population of the C=O groups in the excited state as well as the excited Eu³⁺ ions. Figure 9 schematically illustrates the excitation and emission process of Eu³⁺ ions. In the case (a) that no PVP and Au nanoparticles are introduced, Eu³⁺ ions are excited through a charge transfer (CT) state by the wavelength of $\lambda = 337.1$ nm. After relaxation to 5D_0 level, the emissions due to 5D_0 - 7F_J transitions are obtained. If PVP is incorporated and the resultant gel is heated at temperatures around 400 °C (b), the coordinated C=O groups acts as energy donors to Eu³⁺. The absorbed energy by singlet-singlet ($S_0 \rightarrow S_1$) transition of π electrons of C=O groups is transferred to the 5L_6 levels of Eu³⁺ ions through a triplet state of C=O. Interestingly, the 5D_1 - $^7F_{1,2}$ emissions were obtained for the AuEu-BS sample heated at 400 °C, i.e., in case of the incorporation of the polymer-protected Au nanoparticles (see Figure 7). We found at the same time

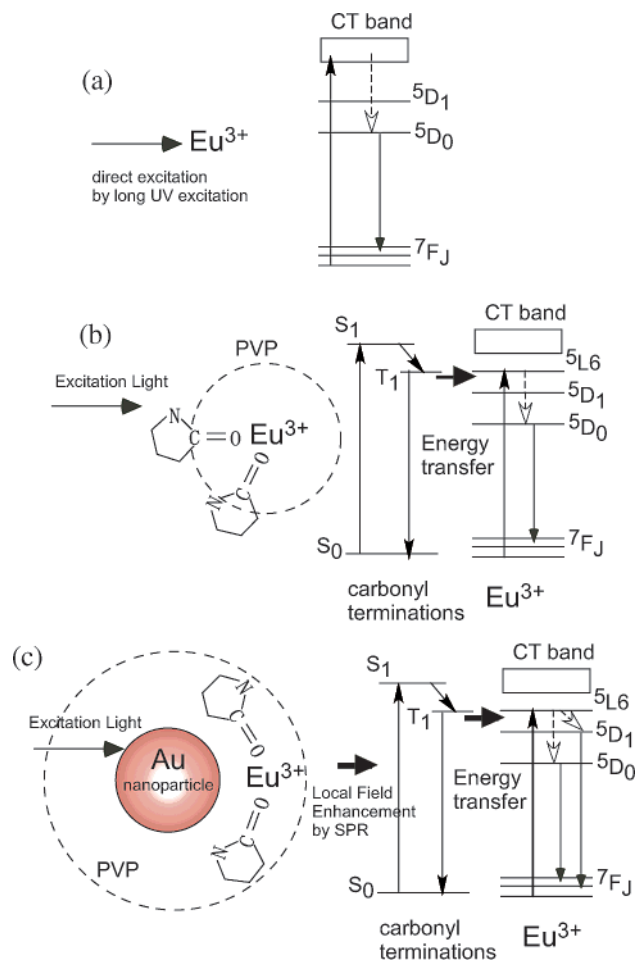


Figure 9. Schematic illustrations for enhanced Eu^{3+} PL in the vicinity of PVP-protected Au nanoparticles. (a) Isolated Eu^{3+} ions, (b) Eu-BS glass, and (c) AuEu-BS glass.

that Eu^{3+} emissions were greatly enhanced in the presence of Au nanoparticles, and thus an influence of Au nanoparticles to Eu^{3+} emissions is sufficient to be considered. Our data is therefore interpreted in such a way as the additional supplement of excitation energy for Eu^{3+} ions was employed from $\text{C}=\text{O}$ groups strongly excited in the local field enhanced by SPR of Au nanoparticles (see Figure 9c). This is also supported by the fact that the broad luminescence of $\text{C}=\text{O}$ chromophores was most intensified in the heat-treatment at 400 °C, as it can be seen as $R_{\text{C}=\text{O}}$ in Table 1. Indeed, the results in Table 1 also include an influence of hydroxyl (OH) vibration on the Eu^{3+} luminescence efficiency. However, the observed enhancement of Eu^{3+} luminescence cannot be explained only by this de-excitation.

It is important to point out that the incident light for excitation used in this work does not match with any resonance mode of surface plasmon, whose frequency is expressed by $\omega_L = \omega_p / (1 + (L + 1)/L \cdot \epsilon_M)^{1/2}$; the excitation energy is 3.68 eV, while a plasma oscillation of $L \rightarrow \infty$ occurs at the energy of 3.2 eV. Nevertheless, the strong excitation produce hot electrons through the $d \rightarrow sp$ interband transition, resulting in an increasing number of density of conduction electrons within a very short time of electron-phonon coupling time in picosecond order.²⁸ This simultaneously means a probable blue-shift of plasmon frequency ω_p . Thus, we confidently propose that it is possible for the long UV light to derive a higher order of surface plasmon resonance for nanometer-sized Au particles, in the vicinity of which luminescent centers such as $\text{C}=\text{O}$

chromophore or rare-earth ions can be strongly excited through the induced local field.

5. Conclusion

Nanometer-sized gold (Au) colloids in solution were incorporated into a borosilicate amorphous network of $10\text{B}_2\text{O}_3$ – 90SiO_2 together with Eu^{3+} ions via a sol–gel route. To control the size and obtain good stability of Au nanoparticles a protective polymer of poly(vinylpyrrolidone) (PVP) was utilized throughout the synthesis. The sizes of the particles in a starting solution and gel-glasses were monitored by optical extinction spectra and analyzed by Mie-Drude theory. It was shown that Eu^{3+} fluorescence of $5D_0$ – $7F_J$ transitions in the vicinity of Au nanoparticles was enhanced by a local field enhancement due to higher orders of SPR excitations. The enhancement factor f_E was ~ 2.5 . Influences of carbon-related phosphors originating from the protective polymer for Au nanoparticles were also studied at various temperatures in heat-treatment for the fabricated gels of $10\text{B}_2\text{O}_3$ – 90SiO_2 doped with Eu^{3+} ions. It was revealed that carbonyl groups around Au nanoparticles played a role of energy donor to Eu^{3+} ions and further enhancement of Eu^{3+} fluorescence, $f_E \sim 6$, was conclusively achieved. The enhancement will further be improved by modifying space-dispersion and size-distribution of Au nanoparticles in glasses.

Acknowledgment. A part of this work was carried out by the Nanotechnology Glass Project as part of Nanotechnology Materials Program supported by New Energy and Industrial Technology Development Organization (NEDO). T.H. gratefully acknowledges Nippon Sheet Glass Foundation for Materials Science and Engineering for financial support.

References and Notes

- (1) Born, M.; Wolf, E. *Principles of Optics*, 6th ed.; Pergamon: Oxford, 1980.
- (2) Kittel, C. *Introduction of Solid State Physics*, 7th ed.; John Wiley & Sons: New York, 1996.
- (3) Hache, F.; Ricard, D.; Flytzanis, C. *J. Opt. Soc. Am. B* **1986**, *3*, 1647.
- (4) Ricard, D.; Roussignol, Ph.; Flytzanis, C. *Opt. Lett.* **1985**, *10*, 511.
- (5) Uchida, K.; Kaneko, S.; Omi, S.; Hata, C.; Tanji, H.; Asahara, Y.; Ikushima, A. J.; Tokizaki, T.; Nakamura, A. *J. Opt. Soc. Am. B* **1994**, *11*, 1236.
- (6) Das, P.; Metiu, H. *J. Phys. Chem.* **1985**, *89*, 4680.
- (7) Chew, H. *J. Chem. Phys.* **1987**, *87*, 1355.
- (8) (a) Malta, O. L.; Santos, M. A. Couto dos. *Chem. Phys. Lett.* **1990**, *174*, 13. (b) Malta, O. L.; Santa-Cruz, P. A.; De Sá, G. F.; Auzel, F. *J. Lumin.* **1985**, *33*, 261.
- (9) Hayakawa, T.; Selvan, Tamil, S.; Nogami, M. *Appl. Phys. Lett.* **1999**, *74*, 1513.
- (10) Hayakawa, T.; Selvan, Tamil, S.; Nogami, M. *J. Non-Cryst. Solids* **1999**, *259*, 16.
- (11) Selvan, Tamil, S.; Hayakawa, T.; Nogami, M. *J. Phys. Chem. B* **1999**, *103*, 7064.
- (12) Teranishi, T.; Kiyokawa, I.; Miyake, M. *Adv. Mater.* **1998**, *10*, 596.
- (13) (a) Mayer, A.; Antonietti, M. *Colloid Polym. Sci.* **1998**, *276*, 769. (b) Toshima, N.; Harada, M.; Yamazaki, Y.; Asakura, K. *J. Phys. Chem.* **1992**, *96*, 9927. (c) Teranishi, T.; Hori, H.; Miyake, M. *J. Phys. Chem. B* **1997**, *101*, 5774.
- (14) Mie, G. *Ann. Phys.* **1908**, *25*, 377.
- (15) Doremus, R. H. *J. Chem. Phys.* **1964**, *40*, 2389.
- (16) Doyle, W. T. *Phys. Rev.* **1958**, *111*, 1067.
- (17) Epifani, M.; Carlino, E.; Blasi, C.; Giannini, C.; Tapfer, L.; Vasanelli, L. *Chem. Mater.* **2001**, *13*, 1533.
- (18) Toki, M.; Chow, T. Y.; Ohnaka, T.; Samura, H.; Saegusa, T. *Polym. Bull.* **1992**, *29*, 653.
- (19) Morikawa, A.; Iyoku, Y.; Kakimoto, M.; Imai, Y. *Polym. J.* **1992**, *24*, 687.
- (20) Selvan, Tamil, S.; Hayakawa, T.; Nogami, M.; Möller, M. *J. Phys. Chem. B* **1999**, *103*, 7441.

- (21) Kreibig, U.; Frangstien, C. v. *Z. Physik* **1969**, 224, 307.
- (22) Kreibig, U.; Vollmer, M. *Optical Properties of Metal Clusters*; Springer: Berlin, 1995.
- (23) Cooper, B. R.; Ehrenreich, H.; Philipp, H. R. *Phys. Rev.* **1965**, 2, A494.
- (24) (a) Hovel, J.; Fritz, S.; Hilger, A.; Kreibig, U.; Vollmer, M. *Phys. Rev.* **1993**, 48, 18278. (b) Alemany, P.; Boorse, R. S.; Burlitchmad, J. M.; Hoffman, R. *J. Phys. Chem.* **1993**, 97, 8464. (c) Hayakawa, T.; Ono, Y.; Nogami, M. *Proc. SPIE* **2000**, 3943, 102.
- (25) Canham, L. T.; Loni, A.; Calcott, P. D. J.; Simons, A. J.; Reeves, C.; Houlton, M. R.; Newey, J. P.; Nash, K. J.; Cox, T. I. *Thin Solid Films* **1996**, 276, 112.
- (26) Levy, D.; Reisfeld, R.; Avnir, D. *Chem. Phys. Lett.* **1984**, 109, 593.
- (27) Molina, C.; Dahmouche, K.; Messaddeq, Y.; Ribeiro, S. J. L.; Silva, M. A. P.; de Zea Bermudez, V.; Carlos, L. D. *J. Lumin.* **2003**, 104, 83.
- (28) (a) Hamanaka, Y.; Hayashi, N.; Nakamura, A.; Omi, S. *J. Lumin.* **1998**, 76&77, 221. (b) Inouye, H.; Tanaka, K.; Takahashi, I.; Hirao, K. *Phys. Rev. B* **1998**, 57, 1134.