

## Effect of H<sub>2</sub>O–CO<sub>2</sub> Organization on Ovalbumin Adsorption at the Supercritical CO<sub>2</sub>–Water Interface

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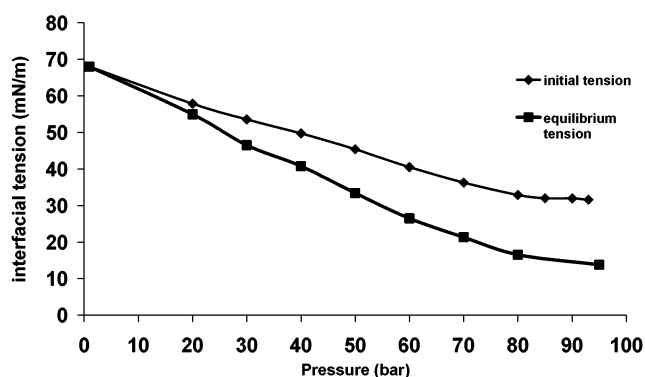
We have studied the formation of water–CO<sub>2</sub> interfaces in the presence of different concentrations of ovalbumin (OVA) by tensiometry and by means of interfacial rheological measurements to obtain some information on the capacity of protein film to stabilize H<sub>2</sub>O in CO<sub>2</sub> emulsion. The formation of pure water–CO<sub>2</sub> interface can be described as a two-step phenomenon.<sup>1</sup> The CO<sub>2</sub> molecules adsorb onto the water surface and then a reorganization of the interface creates a H<sub>2</sub>O–CO<sub>2</sub> cluster network. This organization occurs at a temperature (40 °C) higher than the higher temperature limit (10 °C) allowing the formation of crystalline structure called CO<sub>2</sub> clathrate.<sup>2</sup> Our results show that ovalbumin adsorption from bulk concentrations higher than 0.0229 g/L inhibits the cluster formation for a CO<sub>2</sub> pressure less than 80 bar. However, for lower concentrations, the more the CO<sub>2</sub> pressure is close to 80 bar, the more OVA adsorption is reduced by the H<sub>2</sub>O–CO<sub>2</sub> cluster network. Moreover, from a pressure of 90 bar, the affinity of OVA for the interface increases and mixed films made of protein molecules and clusters are obtained for the OVA concentrations lower than 1 g/L.

### Introduction

Supercritical carbon dioxide (ScCO<sub>2</sub>) is becoming increasingly interesting for many applications such as medium reaction,<sup>3,4</sup> foaming agent, nonmiscible phase for emulsions,<sup>3,5</sup> microemulsions,<sup>6–8</sup> or latexes formation, as well as in cleaning processes.<sup>9</sup> This is due to its peculiar physical properties, which allow the easy tuning of its density and consequently of its solvent power in a small range of temperature and pressure. Moreover, its use answers the need of the substituting toxic and environmentally bad organic solvents in these processes.

In the last 10 years, many studies deal with the formation of water in CO<sub>2</sub> (W/C) and CO<sub>2</sub> in water (C/W) microemulsions<sup>7,8</sup> and emulsions.<sup>5</sup> The first experiments were feasibility studies, with the search of very CO<sub>2</sub>-soluble surfactants (i.e., with fluoroalkyl, fluoroether, or siloxane tails). Almost all studies were based on phase behavior, droplet size determination, conductivity, and interfacial tension ( $\gamma$ ) measurements.<sup>5,10</sup> However, only a few studies have measured  $\gamma$  between H<sub>2</sub>O and CO<sub>2</sub> with surfactant<sup>5,10</sup> while no studies included interfacial rheology.

The dispersed droplets in an emulsion are in constant motion and therefore there are frequent collisions between them in the absence of steric or electrical barrier. These collisions lead to droplet aggregation and, if the interfacial wall is broken, to an irreversible coalescence. In C/W emulsion, this kind of barrier can be created, as this is done in O/W emulsion. Unfortunately, in W/C emulsion, it is very difficult to obtain such a barrier without the use of expensive and unknown biocompatibility surfactants containing fluoroalkyl or fluoroether tails.<sup>11,12</sup> In fact, few polymers able to create a steric barrier are soluble in CO<sub>2</sub>, except fluoropolymers, siloxanes, and polycarbonates. Furthermore, the CO<sub>2</sub> has a low dielectric constant, which renders difficult the creation of an electrostatic barrier.



**Figure 1.** Initial and equilibrium (60 000 s) interfacial tension measured at 40 °C versus pCO<sub>2</sub>, from Tewes et al.<sup>1</sup>

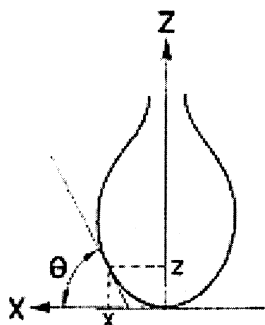
Consequently, the mechanical strength of the interfacial film is one of the prime factors determining the stability of the W/C emulsion.

Proteins strongly adsorb at the oil–water interface, often leading to an interfacial layer with an important elastic contribution.<sup>13–15</sup> Such interfacial films could provide suitable properties to stabilize direct or inverse emulsions composed of water and CO<sub>2</sub>.

We previously investigated the formation of water–CO<sub>2</sub> (W–C) interface by means of a drop tensiometer, in which one can control the temperature, pressure, and drop area.<sup>1</sup>

In this study, we have measured the initial and equilibrium interfacial tension between water and CO<sub>2</sub> (Figure 1) and we have described the formation of this interface as a two-step phenomenon. First, the CO<sub>2</sub> molecules quickly adsorb onto the water surface for equilibrating their chemical potential between the bulk CO<sub>2</sub> and the water surface. This physisorption leads to the interaction of one CO<sub>2</sub> molecule with one H<sub>2</sub>O molecule. From the calculated CO<sub>2</sub> molecular area at saturation and the thermodynamical parameters, we have found that this H<sub>2</sub>O–CO<sub>2</sub> complex would be of H type, as it was found by da Rocha

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**Figure 2.** Profile of drop analyzed with Laplace equation.

et al.<sup>16</sup> by Monte Carlo simulation. The amount of adsorbed CO<sub>2</sub> and the relative orientation of water–CO<sub>2</sub> molecules in the H type complex depend on temperature and pressure.

Second, after this adsorption, CO<sub>2</sub> molecules diffuse into water and then modify the water molecules organization to create an interphase. This reorganization of the interface creates a network of H<sub>2</sub>O–CO<sub>2</sub> clusters and leads to a decrease in the interfacial tension ( $\gamma$ ) until an equilibrium state. These H<sub>2</sub>O–CO<sub>2</sub> clusters are formed at temperatures higher (20–40 °C) than the one usually described as the limit (10 °C) for the formation of crystalline clathrate hydrate.<sup>2,17,18</sup>

Elasticity measurements and macroscopic visualization suggest that the growing of the clusters is driven by the assembly of many small blocks and accelerates with the CO<sub>2</sub> pressure. Their interfacial concentration and their size increase with time until the saturation of the interface.

This organization of the pure H<sub>2</sub>O–CO<sub>2</sub> interface raises an important question. How does the interface react in the presence of usual or polymeric surfactant molecules such as proteins? To answer this question, we have studied the H<sub>2</sub>O–CO<sub>2</sub> interface at 40 °C, in the presence of different concentrations of ovalbumin (OVA), and for various CO<sub>2</sub> pressures, by measuring interfacial tension and film elasticities by means of a pendant drop method.

## Materials and Methods

**Materials.** The vessel was filled with CO<sub>2</sub> (purity 0.99) (Air liquide, France). Ultrapure water was produced by MilliQ plus 188 apparatus (Millipore St Quentin en Yveline, France) and then saturated 12 h with pure CO<sub>2</sub>. Ovalbumin of grade V was purchased from SIGMA, Germany.

**Pendant Drop Tensiometer.** The drop tensiometer (Tracker, IT Concept, Longessaigne, France) allows the determination of the interfacial tension by analyzing the axial symmetric shape (Laplacian profile) of the pendant drop of CO<sub>2</sub> saturated water in pressurized CO<sub>2</sub>. The apparatus has been described in a previous study.<sup>1</sup> Pendant drops were formed at the end of a stainless steel tube, 1-mm inside diameter, connected to a syringe. Drop volume and area were controlled during all experiments by a step-by-step motor. Therefore, it was possible to maintain constant the area of the drop during the time of the experience, so that the surface tension variation was related only to the adsorption of molecules at the interface.

The interfacial tension was determined by digitizing and analyzing the profile of the droplet (Figure 2) using a CCD camera coupled to a video image profile digitizer board connected to a computer. The drop profile was processed according to the fundamental Laplace equation [eq 1] applied to the drop profile:

$$\frac{1}{x} \frac{d}{dx}(x \sin \theta) = \frac{2}{b} - cz \quad (1)$$

Where  $x$  and  $z$  are the Cartesian coordinates at any point of the drop profile.  $b$  is the radius of the curvature of the drop apex.  $\theta$  is the angle of the tangent to the drop profile. In addition,  $c$  is the capillarity constant, equal to “ $(g\rho)/\gamma$ ”, where  $\rho$  is the difference between the densities of the two liquids and  $g$  is the acceleration due to gravity. Five times per second, the computer calculates the characteristic parameters of the drop (area, volume, and interfacial tension).

**Rheological Measurement.** Theoretical aspects have been largely developed and discussed in previous papers.<sup>14,19,20</sup> In summary, the viscoelastic response of an interfacial film submitted to a dilatational mechanical strain was studied using a theoretical approach based on two-dimensional rheology.<sup>21</sup>

This approach consists of realizing two types of continuous and monotonic compression of the equilibrated surface layer of a pendant drop: a slow compression and fast compression. We recorded simultaneously the variation of the interfacial tension (i.e., resulting stress) during the slow compression, and after the fast compression, to measure the relaxation of the interface.

In this case, a convenient theoretical model (generalized Maxwell, Figure 3), corresponding to a solid viscoelastic body,

is applied and allowed us to estimate the three viscoelastic parameters of the interfacial film:  $\tau$ , the characteristic time of relaxation;  $E_e$ , the equilibrium surface dilatational elasticity and  $E_{ne}$ , the nonequilibrium part of the surface dilatational elasticity.

To describe the variation of the total interfacial tension  $\Delta\gamma$ , (equal to the difference between  $\gamma_i$ , the initial interfacial tension before the compression, and  $\gamma(t)$ , the interfacial tension at  $t$  time), during the time of the slow compression at a constant velocity  $U$ , we supposed that at any moment  $\Delta\gamma$  can be written as a sum of equilibrium ( $\Delta\gamma_e$ ) and nonequilibrium ( $\Delta\gamma_{ne}$ ) contributions of stress:

$$\Delta\gamma = \Delta\gamma_e + \Delta\gamma_{ne} \quad (2)$$

The equilibrium part of the resulting stress  $\Delta\gamma_e$  depends on the equilibrium elasticity  $E_e$ . This conservative part of elasticity is represented by the upper branch of the mechanical model in Figure 3, which corresponds to a spring. It is related to a long-range organization of the protein film and consequently to the interaction between the segments anchored at the interface.

The nonequilibrium part of the resulting stress  $\Delta\gamma_{ne}$  depends on the nonequilibrium elasticity  $E_{ne}$  and on the characteristic time of relation  $\tau$ . This viscoelastic behavior is represented by the lower branch of the mechanical model in Figure 3 and corresponds to a spring coupled with a shock absorber. The nonequilibrium elasticity is linked to the dissipation of the accumulated energy through molecular reorganization like expulsion of segments in the adjacent phases occurring during the compression and the relaxation of the interface. The relaxation time  $\tau$  represents the necessary time for the interface to reach a new equilibrium energetic state after a perturbation.

The two branches of the mechanical model are coupled in parallel according to eq 2, corresponding to the additivity of stresses; we obtained the following equation describing the viscoelastic behavior of the monolayer:

$$\frac{\Delta\gamma}{\Delta A} A_i = E_e + E_{ne} \frac{\tau}{t} (1 - e^{-t/\tau}) \quad (3)$$

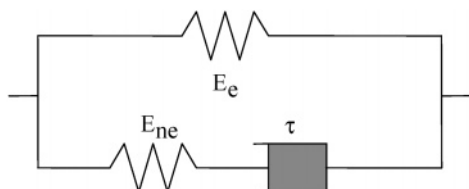


Figure 3. Mechanical model of the monolayer.

where  $\Delta A$  corresponds to the drop area variation and  $A_i$  is the area of the drop before compressing the interfacial film.

The slow compression, typically lower than  $0.003 \text{ s}^{-1}$ , was performed until reaching a total relative strain ( $\Delta A/A_i$ ) of 1–5% of the initial area, to determine the dilatational elasticity parameters. The fast compression, typically higher than  $0.005 \text{ s}^{-1}$ , was performed to determine  $\tau$ .

In a pure elastic behavior ( $E_{ne} = 0$ ), we estimate  $E_e$ , which is equal to

$$E_e = -A_i \frac{d\gamma}{dA} \quad (4)$$

## Results and Discussion

**1. Interfacial Tension Measurement.** We have measured the  $\gamma$  kinetics at the  $\text{H}_2\text{O}-\text{CO}_2$  interface at  $40^\circ\text{C}$ , in the presence of OVA in the bulk water phase at concentrations ([OVA]) varying from 0.002 g/L to 4.75 g/L, and for various  $\text{CO}_2$  pressures (50–90 bar).

We show in Figure 4 the interfacial tension profile versus  $\text{CO}_2$  pressure obtained for the pure water–carbon dioxide system as we previously described<sup>1</sup>, to have a better comparison with the results measured for the ovalbumin.

Furthermore, we only represent in Figure 5 some significant kinetics curves obtained in the presence of OVA in order not to complicate the text.

For an [OVA] equal to 0.002 g/L and for  $p\text{CO}_2$  comprised between 60 and 80 bar (Figure 5A), the evolution of  $\gamma$  in the presence of OVA is similar to the kinetics we previously obtained for the pure  $\text{CO}_2-\text{H}_2\text{O}$  interface (Figure 4). Actually, the curve obtained for a  $p\text{CO}_2$  of 50 bar in the presence of OVA is slightly below the curve corresponding to the pure system. Furthermore, for a  $p\text{CO}_2$  of 90 bar (Figure 5B), the shape of the kinetics differs from that obtained for the pure interface,

except for the initial rapid decrease of the interfacial tension. In particular, for 90 bar, we can notice the presence of a plateau during about 40 000 seconds, for which  $\gamma$  value slightly increases. This plateau is followed by a novel decrease of  $\gamma$ , which indicates a new organization of the interface.

From the [OVA] of 0.0229 g/L, the kinetics obtained for the  $p\text{CO}_2$  comprised between 50 and 80 bar are different from those obtained at the pure  $\text{CO}_2-\text{H}_2\text{O}$  interface, specially for the  $\gamma$  equilibrium value. For 90 bar, an increase in the [OVA] to 0.0229 g/L, 0.0477 g/L, 0.523 g/L, and above, leads to a decrease of the plateau duration, which was, respectively, 8000 s, 2000 s, and unapparent.

For dilute protein solutions, the “classical” shape of kinetics at fluid interfaces is generally described as the resultant of phenomena modeled with three characteristic times.<sup>22–26</sup> The first time is attributed to the protein adsorption controlled by the diffusion from the bulk. The two next steps are generally related to the protein film organization via intermolecular interactions and molecular unfolding. In our case, from 60 to 80 bar, the superimposition of the kinetics acquired in the presence of OVA at 0.002 g/L with those obtained at pure  $\text{CO}_2-\text{H}_2\text{O}$  interface suggests the possibility of an inhibition of the OVA adsorption by the formation of the  $\text{H}_2\text{O}-\text{CO}_2$  cluster network. On the other hand, increasing the [OVA] over 0.0229 g/L in this range of  $p\text{CO}_2$  leads to kinetics becoming different from those obtained for the pure system, indicating the adsorption of the protein.

For a  $p\text{CO}_2$  of 90 bar, the peculiar shape of the kinetics obtained for the [OVA] in the range 0.002–0.0477 g/L shows the contribution of both the  $\text{CO}_2$  cluster formation and the OVA adsorption on the interfacial tension. The first rapid  $\gamma$  decrease, which looks like that obtained for the pure interface, is related to the  $\text{H}_2\text{O}-\text{CO}_2$  organization. The presence of a plateau in the kinetics, having a duration which decreases when the [OVA] increases, reflects OVA adsorption and indicates some interaction between OVA and the  $\text{H}_2\text{O}-\text{CO}_2$  cluster network. The fact that, at 90 bar, the protein adsorption occurs for an [OVA] of 0.002 g/L, whereas no adsorption is visible for lower  $p\text{CO}_2$ , shows an increase of the protein affinity for the interface. This could be related to the better solvation of the  $\text{CO}_2$ -phile segments of the protein by the  $\text{ScCO}_2$ . This increase in affinity could also be attributed to a change in the protein conformation

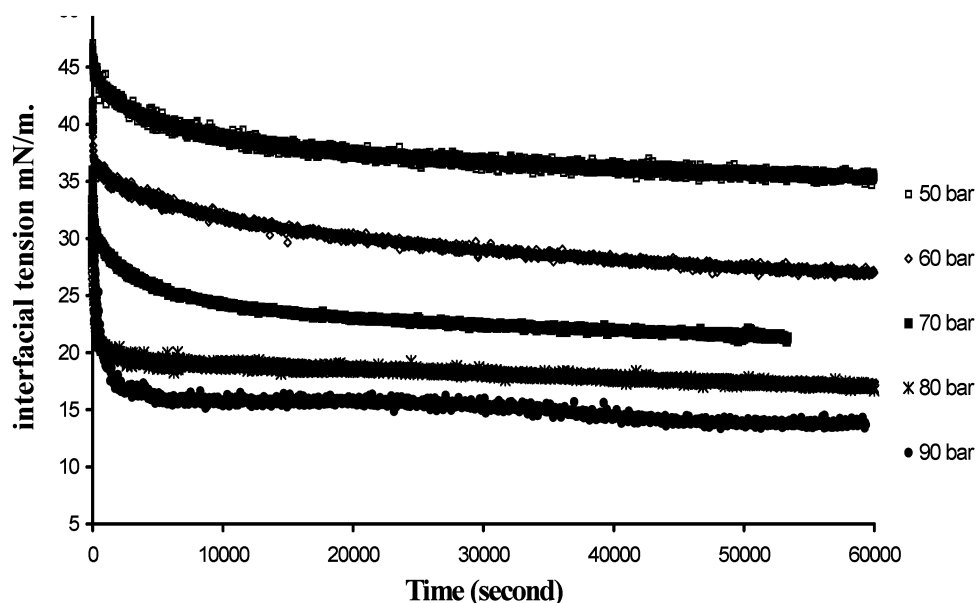
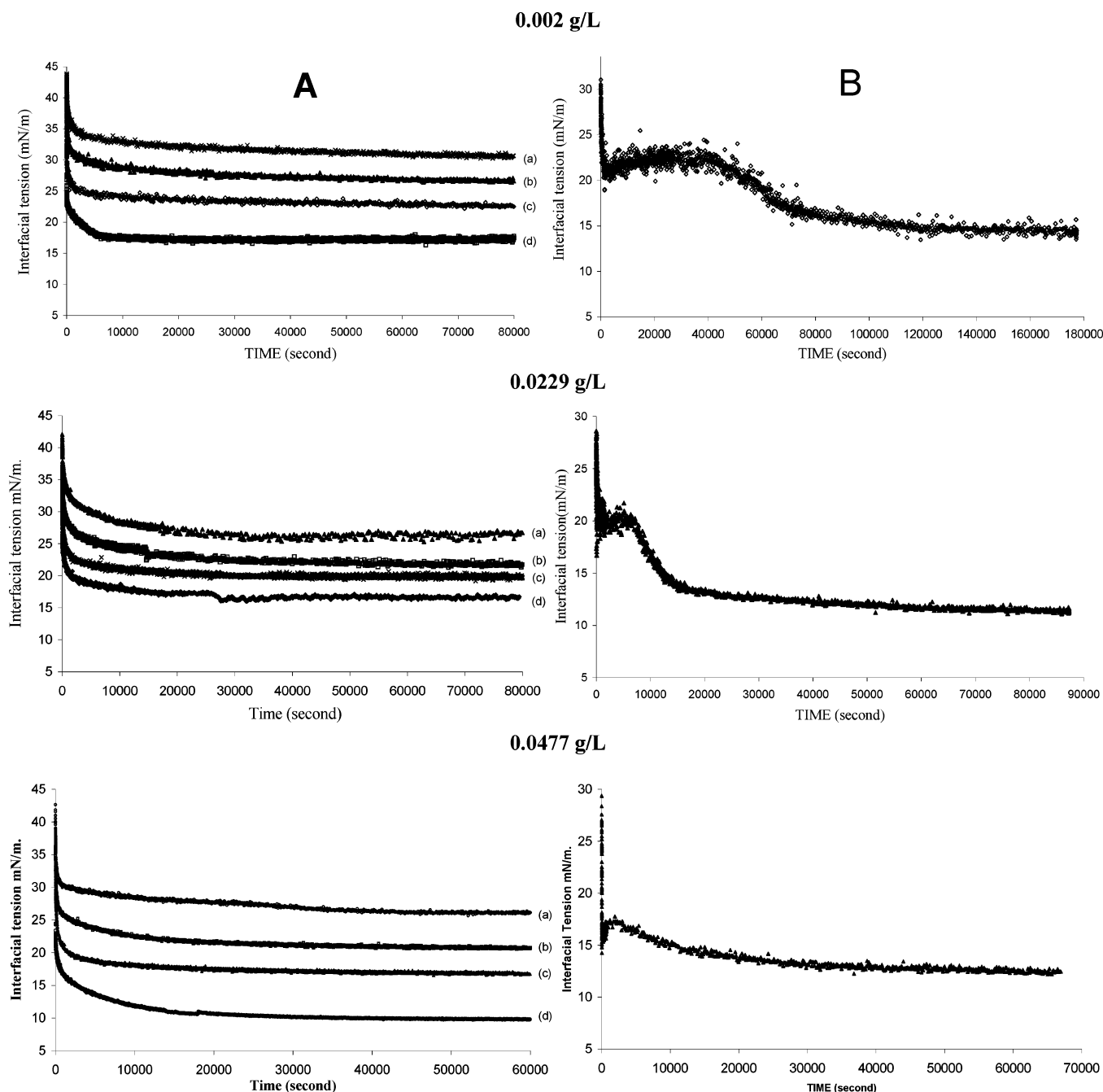


Figure 4. Interfacial tension kinetics measured at  $40^\circ\text{C}$  between the pure water– $\text{CO}_2$  system as a function of  $\text{CO}_2$  pressure, from Tewes et al.<sup>1</sup>



**Figure 5.**  $\gamma$  kinetics at the CO<sub>2</sub>–H<sub>2</sub>O interface measured at 40 °C in the presence of OVA at a concentration of 0.002 g/L, 0.0229 g/L, and 0.0477 g/L for which (A) pCO<sub>2</sub> was 50 bar (a), 60 bar (b), 70 bar (c), and 80 bar (d) and (B) pCO<sub>2</sub> was 90 bar.

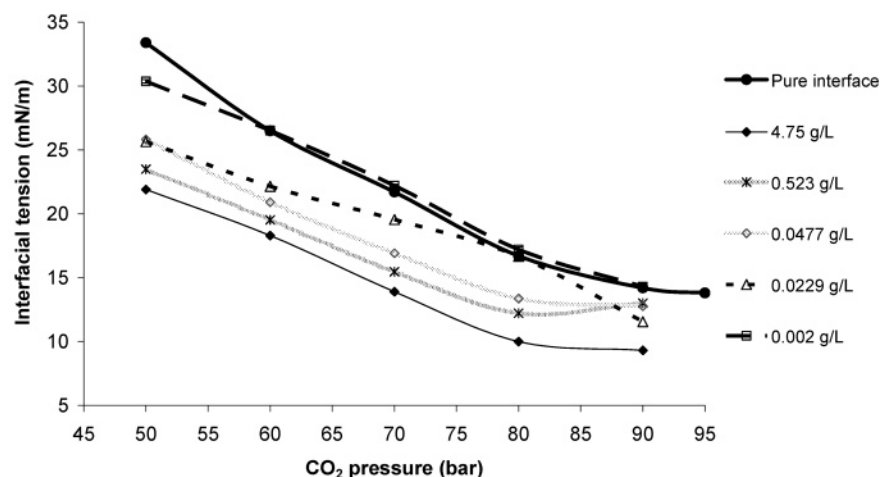
for pCO<sub>2</sub> higher than 80 bar, as it has been shown for bovine serum albumin (BSA) and lysozyme films by FT-IR under ScCO<sub>2</sub> conditions.<sup>27</sup> Furthermore, the treatment of protein solutions (OVA, BSA, and myoglobin) by microbubbling of CO<sub>2</sub> (300 bar and 35 °C) induces an irreversible unfolding of the protein with the loss of  $\alpha$ -helix structure.<sup>28</sup> This result was not only correlated with the protein denaturation due to the pH decrease, induced by the production of carbonic acid by the reaction of CO<sub>2</sub> with H<sub>2</sub>O, but also with the CO<sub>2</sub> sorption onto the amino and guanidium protein groups.

The decrease of the duration of the plateau with the increase in [OVA] indicates that the lag time due to the protein adsorption and the development of the first intermolecular interactions decreases. At the air–water interface, a lag time of 10 000 seconds is observable on the kinetics performed with a bulk [OVA] of 0.001 g/L, and none for concentrations higher than

0.01 g/L.<sup>29,30</sup> In our case, the duration of the plateau is 40 000 seconds for a bulk [OVA] equal to 0.002 g/L and it is present up to a concentration of 0.0477 g/L. Therefore, for a pCO<sub>2</sub> of 90 bar, a peculiar phenomenon increases the necessary time for the protein to develop some interactions at the H<sub>2</sub>O–CO<sub>2</sub> interface. The presence of the H<sub>2</sub>O–CO<sub>2</sub> cluster network, which slows down the development of OVA–OVA interactions, combined with the required long time for the protein to modify its secondary structure under ScCO<sub>2</sub><sup>27</sup> could explain this result.

The slight increase of the interfacial tension during this plateau, followed by a second decrease, implies several concomitant phenomena: First, the substitution or the breaking of the cluster network by the protein, whose affinity for the interface has been increased from this pCO<sub>2</sub> and, second, the formation of an interfacial structure (mixed OVA–cluster or





**Figure 6.** Variation of the equilibrium interfacial tension versus pCO<sub>2</sub> in the presence of various [OVA].

OVA alone), which decreases more efficiently the interfacial tension.

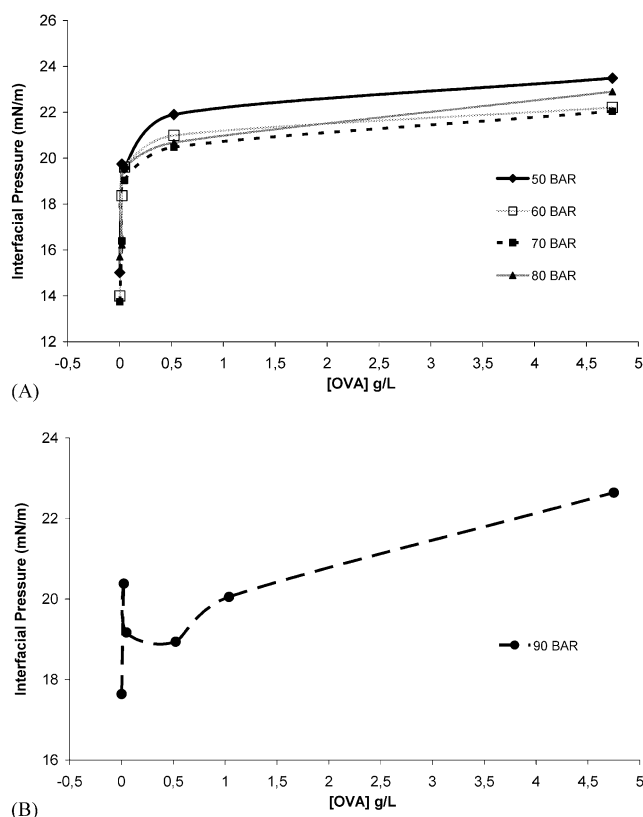
To analyze the interfacial layers that are formed in the different equilibrium conditions, the equilibrium  $\gamma$  values obtained for the several [OVA] were plotted in function of the pCO<sub>2</sub> (Figure 6).

Figure 6 shows three types of behavior accompanied with different organization of the interface: (i) The first behavior is observed between 60 and 80 bar and corresponds to the curves obtained for the lowest [OVA] (<0.0229 g/L). The values in these curves are very close to the  $\gamma$  values that we have obtained for the pure CO<sub>2</sub>–H<sub>2</sub>O interface. (ii) A second type of behavior is limited between 50 and 80 bar and composed of the curves obtained for the high [OVA] (0.0477–4.75 g/L), which are parallel and have very different values than those acquired for the pure interface. The curve obtained at 0.0229 g/L shows an intermediate state between behavior i and ii probably corresponding to a critical [OVA] in our system, which induces a reorganization of the protein film in the subphase. (iii) A third type of behavior includes the values acquired for a pCO<sub>2</sub> of 90 bar. Except for this last behavior, one can remark that the equilibrium  $\gamma$  values decrease when the [OVA] increases.

For an [OVA] equal to 0.002 g/L, the equilibrium  $\gamma$  values are equal to those obtained for the pure interface, except for the pCO<sub>2</sub> equal to 50 bar which gave lower values. We have shown in a previous study<sup>1</sup> that at 50 bar, the H<sub>2</sub>O–CO<sub>2</sub> interface is slightly organized in cluster network. Then, this organization increases when the pCO<sub>2</sub> increases. So, for this [OVA] and a pCO<sub>2</sub> of 50 bar, the OVA adsorption occurs, without being obviously influenced by the organization of the H<sub>2</sub>O and CO<sub>2</sub>. However, above 50 bar, the equilibrium  $\gamma$  values suggest an inhibition of the OVA adsorption by the H<sub>2</sub>O–CO<sub>2</sub> cluster network.

When we increase the [OVA] to 0.0229 g/L, the values of  $\gamma$  measured for the low pCO<sub>2</sub> (i.e., 50–70 bar) are quite different from those obtained for the pure interface, indicating a protein adsorption. However, at 80 bar  $\gamma$  reaches the value of the pure H<sub>2</sub>O–CO<sub>2</sub> system. This can be attributed to the formation of a stronger H<sub>2</sub>O–CO<sub>2</sub> cluster network that inhibits the OVA adsorption.

For the second behavior observed in Figure 6, the equilibrium  $\gamma$  values decrease linearly with the pCO<sub>2</sub> until 80 bar and the curves are parallel between them. These curves are also parallel with the curve which represents the  $\gamma$  values measured at time zero of the kinetics obtained for the pure H<sub>2</sub>O–CO<sub>2</sub> system versus pCO<sub>2</sub> (Figure 1). This suggests that, from 0.0477 g/L,



**Figure 7.** Interfacial pressure ( $\pi$ ) versus OVA concentration.

the values of  $\gamma$  are no more influenced by the formation of a H<sub>2</sub>O–CO<sub>2</sub> cluster network, but by the protein adsorption. Consequently, for high [OVA] and between 50 and 80 bar, OVA is the main component of the interfacial film.

To better describe the contribution of protein upon the adsorption phenomenon, we represented the interfacial pressure ( $\pi = \gamma_0 - \gamma(t)$ ), where  $\gamma_0$  is the initial interfacial tension and  $\gamma(t)$  the interfacial tension measured at the equilibrium adsorption, in function of the [OVA] (Figure 7).

For pCO<sub>2</sub> comprised between 50 and 80 bar, the curves  $\pi$  versus [OVA] have a hyperbolic shape which is not influenced by the pCO<sub>2</sub> for high [OVA]. This is the shape usually found for protein adsorption at classical fluid–fluid interface,<sup>31</sup> or air–fluid interface, except that in our case, the pure H<sub>2</sub>O–CO<sub>2</sub> interface is influenced by the H<sub>2</sub>O–CO<sub>2</sub> organization leading to  $\pi$  values different from zero.

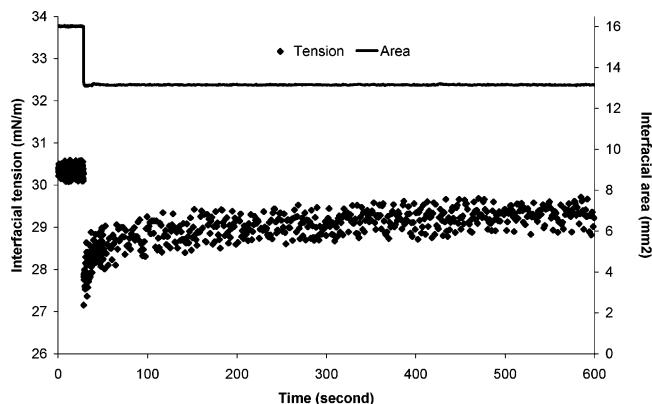
**TABLE 1: Rheological Parameters versus pCO<sub>2</sub> Obtained for the Pure Interface and [OVA] Equal to 0.002 g/L**

pCO <sub>2</sub> (bar)	Ee (mN/m)		Ene (mN/m)		$\tau$ (seconds)	
	water–CO <sub>2</sub>	OVA 0.002 g/L	water–CO <sub>2</sub>	OVA 0.002 g/L	water–CO <sub>2</sub>	OVA 0.002 g/L
50	2 ± 1	5 ± 1	34 ± 5	19 ± 2	23 ± 8	11 ± 2
60	6 ± 1	8 ± 3	29 ± 6	8 ± 3	11 ± 4	9 ± 4
70	8 ± 3	10 ± 3	19 ± 2	18 ± 10	15 ± 8	11 ± 4
80	56 ± 19	65 ± 17	0	0	0	0
90	109 ± 12	49 ± 21	0	0	0	0

At 90 bar, the  $\gamma$  equilibrium values measured in the presence of a [OVA] higher than 0.002 g/L are lower than those measured for the pure interface, indicating a protein adsorption (Figure 6). However, the interfacial pressure values do not increase monotonically with the [OVA], contrary to the other pCO<sub>2</sub> (Figure 7B). We can observe a peculiar accident in the curve for [OVA] between 0.0229 and 0.523 g/L. Since the interfacial pressure reflects the degree of interaction between the two phases, this surprising decrease of  $\pi$  is probably linked to a new organization of the interfacial film, leading to a decrease of the affinity for both phases in contact (i.e., diminution of the water–CO<sub>2</sub> interactions). This can also be explained by the conformational state of the adsorbed protein which is controlled in part by the bulk concentration. To better precise this interfacial organization, we performed rheological measurements.

**2. Rheology Measurements.** We have measured the dilatational rheological comportment of the interfacial layer in the presence of various bulk [OVA] by compressing the drop area (strain) and then measuring the resulting change of  $\gamma$  (stress). For the pure H<sub>2</sub>O–CO<sub>2</sub> interface (Table 1), the equilibrium elasticity increases with the pCO<sub>2</sub>, leading to high elasticity values for the pCO<sub>2</sub> equal to 90 bar, which reveals strong lateral interactions between the H<sub>2</sub>O–CO<sub>2</sub> clusters that can be compared to a solidlike behavior. In the range 50–70 bar, a viscoelastic behavior of the pure interface is observed with a significant Ene value. This can be explained by the structure of the interface, with the presence of a few number of cluster blocks surrounded by a phase made of less organized structures. The decrease of the Ene part of the elasticity with the pCO<sub>2</sub> indicates an organization of the interface and a lowering of the dynamic effects linked to the interaction between the cluster blocks.

At an OVA concentration of 0.002 g/L, the rheological comportment of the interface is also influenced by the pCO<sub>2</sub>. Until a pCO<sub>2</sub> of 70 bar, the interfacial layer has a viscoelastic comportment, evidenced by a relaxation of the interfacial tension after the deformation is maintained (Figure 8). This behavior looks like the one we obtained at the pure interface, but with



**Figure 8.** Fast drop compression after reaching the equilibrium state of the interfacial layer in the presence of OVA at 0.002 g/L and a pCO<sub>2</sub> of 50 bar.

lower Ene values for 50 and 60 bar (Table 1). This last observation shows the presence of OVA intercalated between the cluster blocks. This organization could reduce the lateral motion of the blocks during the compression. At 70 and 80 bar, the rheological comportment in the presence of OVA is the same as that of H<sub>2</sub>O–CO<sub>2</sub> pure system and becomes purely elastic for a pCO<sub>2</sub> of 80 bar. Therefore, protein adsorption is inhibited at these pressures as it was already suggested by the adsorption kinetics. These situation is depicted schematically in Figure 9.

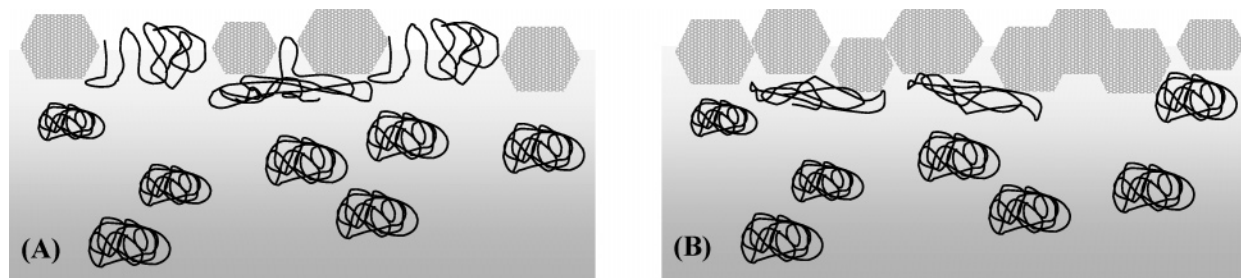
At 90 bar, an important decrease in the equilibrium elasticity is recorded in comparison with the pure interface (Table 1, Figure 10). This indicates the presence of proteins in the interfacial film, resulting probably from the penetration of protein segments into the cluster network. This could be due to a modification of the protein conformation in the water bulk from this peculiar CO<sub>2</sub> pressure as already observed by Striolo et al.<sup>27</sup> Such conformational changes could enhance the adsorption capacity of OVA.

Another reason could be the increasing solubility of CO<sub>2</sub>-phile protein segments into the CO<sub>2</sub> phase due to the increase of the density of the supercritical phase.

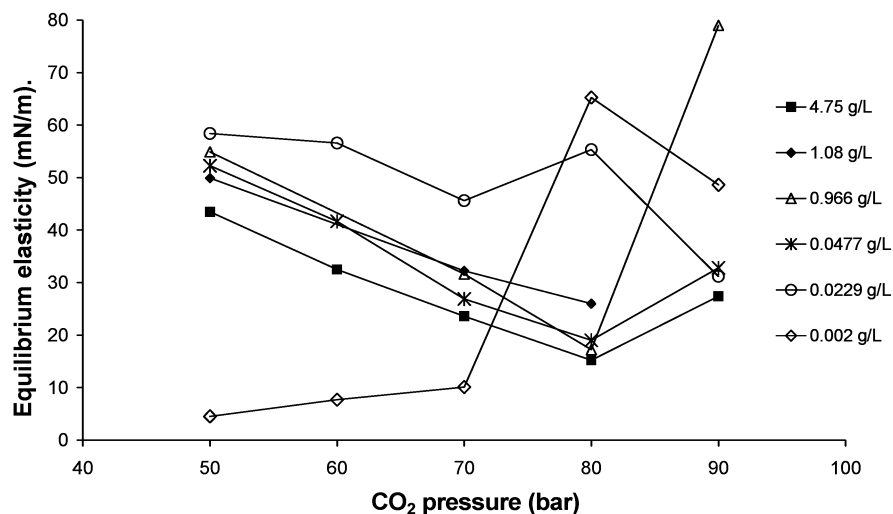
In presence of bulk [OVA] higher than 0.002 g/L, the interfacial layers exhibit a strict elastic comportment for all studied pCO<sub>2</sub>. To compare the results, we have plotted the equilibrium elasticity (Ee) versus pCO<sub>2</sub> (Figure 10).

[OVA] higher than 0.002 g/L lead to high Ee values for the lowest pCO<sub>2</sub> compared to the values obtained for the pure interface. Then, in the range 0.0477–4.75 g/L of [OVA], contrary to the pure system, Ee decreases when the pCO<sub>2</sub> increases until 80 bar. In this range of [OVA], the more pCO<sub>2</sub> increases to 80 bar, the less cohesive is the interfacial layer. A decrease in Ee can be related to a decrease in the number or the energy of the interactions between the OVA segments in the plane of the interface. This can be explained by a competitive increase of the degree of interactions between the CO<sub>2</sub>-phile segments of the OVA and the CO<sub>2</sub> molecules in the subphase because of the increase in the CO<sub>2</sub> density. By comparing adsorption and rheological properties of BSA at air/water interface and dichloromethane (DCM)/water interface, Boury et al.<sup>14</sup> also observed a modification of the interfacial properties, interpreted by a better partitioning of the protein segments at the liquid–liquid (DCM/water) interface. On the other hand, protein adsorption is in part entropically driven as the dehydration of the hydrophobic segments of the protein is favored.<sup>15</sup> When the density of the supercritical CO<sub>2</sub> increases, the water–protein interactions are lowered because of the increase in the CO<sub>2</sub> solubility in the aqueous phase, leading to protein dehydration and then to a change in its conformation.<sup>28</sup> This protein unfolding leads to the exposition of more hydrophobic segments into the water phase, which increases the affinity of the protein for the interface. Consequently, a high CO<sub>2</sub> density probably leads to a better partitioning of the protein between the two adjacent phases.

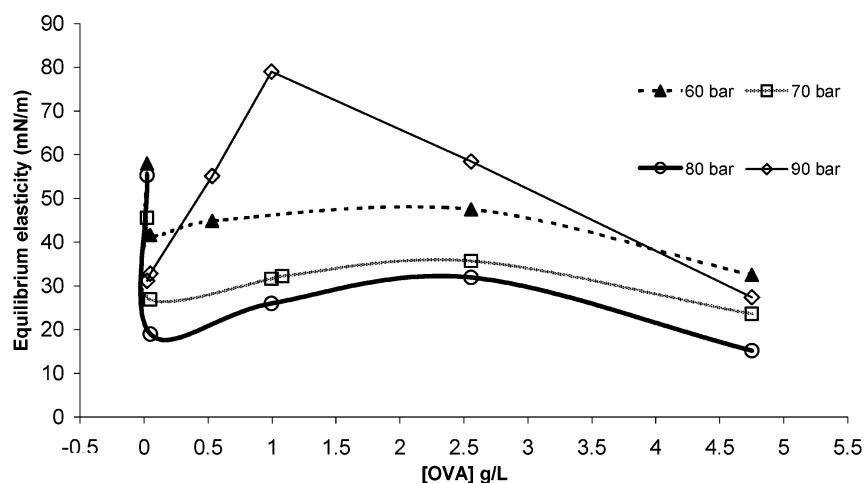
After 80 bar, the variation of Ee with pCO<sub>2</sub> is inverted, showing an increase in the cohesiveness of the interface. A



**Figure 9.** Schematic representation of the interfacial layer obtained at [OVA] equal to 0.002 g/L and under 50–60 bar (A) and 70–80 bar (B).



**Figure 10.** Equilibrium elasticity versus pCO<sub>2</sub> in the presence of various concentrations of OVA.



**Figure 11.** Equilibrium elasticity versus [OVA] concentration.

change in the OVA conformation alone cannot totally explain this increase in Ee because unfolded proteins usually give less elastic interfaces. For instance, slightly organized  $\beta$ -caseine gives a less elastic layer than stable lysozyme at hexadecane/water interface.<sup>32</sup> Therefore, additional effects must cause the formation of a peculiar layer. In fact, another reason can be the formation of a mixed layer, composed of H<sub>2</sub>O–CO<sub>2</sub> clusters and OVA, in which the interactions between protein segments and clusters increase the interfacial elasticity. These clusters–protein interactions could be compared to the clathrate formed with various kind of organic compounds adsorbed onto different proteins.<sup>33,34</sup> Water clathrates in contact with proteins were observed earlier by the X-ray method in the Q-chymotrypsin crystals around bound *n*-hexane molecules.<sup>35</sup> The formation of these “protein–water–organic solvent” clathrate can induce an increase in the water uptake.<sup>36,37</sup> Furthermore, adsorption of CO<sub>2</sub>

onto proteins under supercritical conditions (74 bar and 313 °K) was reported to be close to 10 wt %.<sup>38</sup> Thus, in our case, formation of protein–H<sub>2</sub>O–CO<sub>2</sub> cluster probably occurs and reduces the protein dehydration, leading to a lower partitioning between the two phases compared to that observed at 80 bar, explaining the increase in Ee.

To interpret the effect of [OVA] on the elasticity behavior of the interface, we represented Ee versus the [OVA] for different CO<sub>2</sub> pressures. We have plotted the values measured for [OVA] higher than 0.002 g/L, since at this pressure the protein adsorption did not occur. At pressures lower than 90 bar, Ee is maximum for the lowest [OVA]. Then, Ee abruptly decreases and slightly evolves with the concentration. The effect of the concentration on the interfacial equilibrium elasticity could be explained as it was previously done at the dichloromethane/water interface in the presence of BSA.<sup>14</sup> At low

protein concentrations, the interfacial molecular area is high and the protein is greatly unfolded. This allows the anchored segments to maximize their lateral interactions, explaining the strong elastic behavior. At high protein concentrations, proteins at the interface are less unfolded, promoting intramolecular interactions rather than interactions between the anchored segments.

For 90 bar, the comportment of Ee versus protein concentration is different to that obtained for the other pressures. Ee increases strongly and linearly with the [OVA] until 1 g/L, and then it decreases. This increase in Ee shows an increase of the interface cohesiveness and it is concomitant with the accident observed on the  $\pi$ –[OVA] curve (Figure 7B). This increase in Ee can be due to a poor partitioning of the protein through the two adjacent phases because of the interactions with the H<sub>2</sub>O–CO<sub>2</sub> clusters reducing the protein dehydration. Finally, the interactions of the clusters with the protein modulate the degree of the water–CO<sub>2</sub> interaction and lead to a decrease in the interfacial pressure. At low protein concentrations, the time for the proteins to reach the interface is high and the clusters have enough time to form, allowing their further interaction with the protein. At high protein concentrations (above 1 g/L), proteins can saturate the interface before the initiation of the cluster formation.

## Conclusion

We demonstrated that protein interfacial layers formed at CO<sub>2</sub>–water interface can be analyzed by means of a pendant drop technique. This technique allowed us to measure the adsorption kinetics and dilatational properties of the films. Results show that OVA adsorption is inhibited by H<sub>2</sub>O–CO<sub>2</sub> cluster formation for a bulk concentration of 0.002 g/L and for pCO<sub>2</sub> comprised between 50 and 80 bar. For higher concentration, and in these ranges of pCO<sub>2</sub>, OVA is able to adsorb with an organization pCO<sub>2</sub> dependent. Moreover, for a pCO<sub>2</sub> of 90 bar, and for all OVA concentrations, we obtained peculiar effects probably related to the formation of mixed films made of protein and H<sub>2</sub>O–CO<sub>2</sub> clusters.

This complex adsorption process involves four main phenomena which are dependent on pressure and protein concentration: (i) the formation of a cluster network at the H<sub>2</sub>O–CO<sub>2</sub> interface, (ii) the adsorption of proteins involving classical mechanisms such as unfolding and multilayer formation, (iii) the solvation properties of the CO<sub>2</sub> toward protein segments, and (iv) the conformational modifications of the protein molecules in the bulk water.

Our results show that, for H<sub>2</sub>O–CO<sub>2</sub> interfacial studies performed in the presence of macromolecular surfactant, the formation of water–CO<sub>2</sub> network should be taken into account.

Proteins could be good candidates to create a cohesive film at the H<sub>2</sub>O–CO<sub>2</sub> interface, which may be useful to inhibit the coalescence of water droplets in a CO<sub>2</sub> surrounding phase (inverse emulsions). Studies are in course to establish the ability of protein to stabilize such emulsions. This involves feasibility diagrams and analysis of the structural conformation of the proteins in emulsion, in particular by Raman spectroscopy.

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