



## Effects of Peripheral Substituents on the Electronic Structure and Properties of Unligated and Ligated Metal Phthalocyanines, Metal = Fe, Co, Zn

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**Abstract:** The effects of peripheral, multiple  $-F$  as well as  $-C_2F_5$  substituents, on the electronic structure and properties of unligated and ligated metal phthalocyanines,  $PcM$ ,  $PcM(acetone)_2$  ( $M = Fe, Co, Zn$ ),  $PcZn(Cl)$ , and  $PcZn(Cl^-)$ , have been investigated using a DFT method. The calculations provide a clear explanation for the changes in the ground state, molecular orbital (MO) energy levels, ionization potentials (IP), electron affinities (EA), charge distribution on the metal ( $Q_M$ ), axial binding energies, and in electronic spectra. While the strongly electron-withdrawing  $-C_2F_5$  groups on the  $Pc$  ring change the ground state of  $PcFe$ , they do not influence the ground state of  $PcCo$ . The IP is increased by  $\sim 1.3$  eV from  $H_{16}PcM$  to  $F_{16}PcM$  and by another  $\sim 1.1$  eV from  $F_{16}PcM$  to  $F_{48}PcM$ . A similar increase in the EA is also found on going from  $H_{16}PcM$  to  $F_{48}PcM$ . Substitution by the  $-C_2F_5$  groups also considerably increases the binding strength between  $PcM$  and the electron-donating axial ligand(s). Numerous changes in chemical and physical properties observed for the  $F_{64}PcM$  compounds can be accounted for by the calculated results.

### 1. Introduction

Metal phthalocyanines ( $PcMs$ ), with their planar square structure, are interesting molecules that have been studied extensively in the literature. ( $Pc$  is used here to refer to any phthalocyanine, regardless of substituents.) They provide a versatile chemical system: almost every metal in the periodic table can combine with the  $Pc$  ring, and most of these compounds are very stable. The diversity of phthalocyanines, together with their high thermal and chemical stability, has made them suitable for many technological applications such as dyes, pigments, semiconductors,<sup>1</sup> energy conversion (photovoltaic and solar cells),<sup>2</sup> electrophotography,<sup>3</sup> photo-

sensitizers,<sup>4</sup> gas sensors,<sup>5</sup> low-dimensional metals,<sup>6</sup> electrochromism,<sup>7</sup> Langmuir–Blodgett (LB) films,<sup>8</sup> liquid crystals,<sup>9</sup> and nonlinear optics.<sup>10</sup>

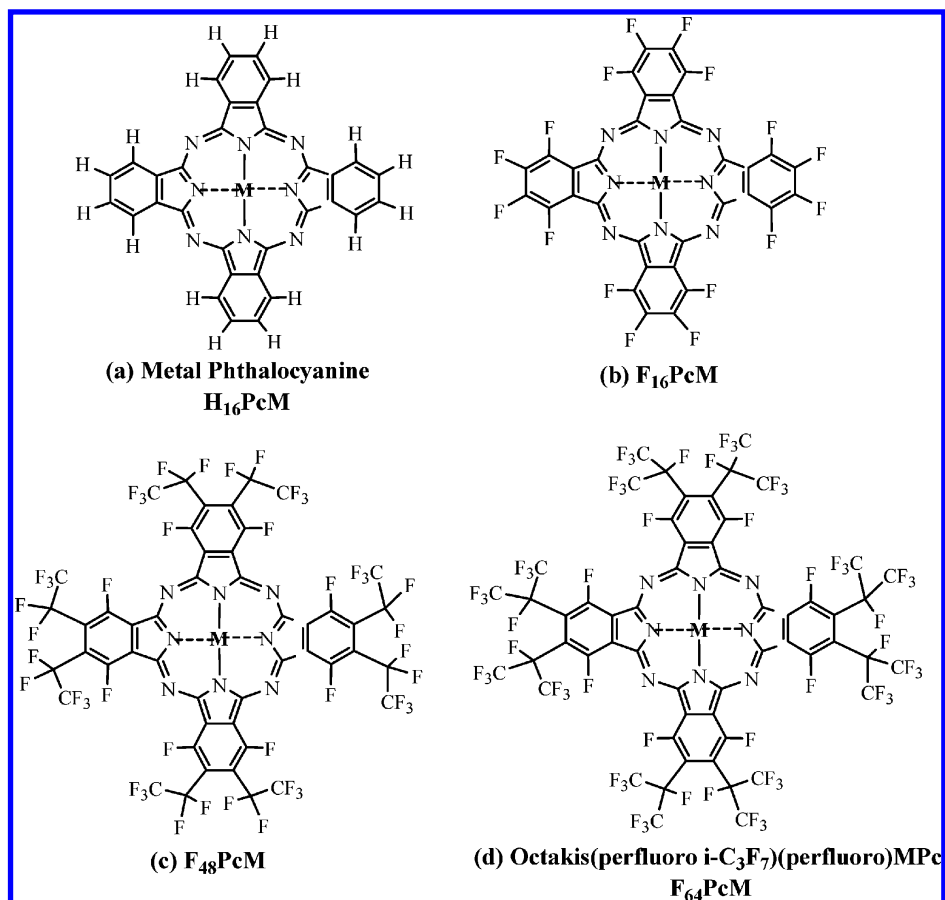
Consequently, a great deal of effort has been expended to develop new phthalocyanine materials that display certain characteristics useful for the particular application. It has been shown that the properties of phthalocyanines can be effectively modulated with different substituents on the periphery of the macrocycle. A useful methodology is the replacement of the hydrogens ( $Hs$ ) of  $Pc$  with halogens ( $Xs$ ), which can greatly increase the catalytic activity and stability of  $PcMs$ .<sup>11,12</sup> However, halogenated  $PcMs$  have the disadvantage of poor solubility, which impedes many new applications of the materials. Recently, several novel octakis-(perfluoro  $i$ - $C_3F_7$ )(perfluoro)- $PcM$  compounds ( $F_{64}PcM$ ) ( $M = Zn, Co, Fe$ ) have been synthesized and characterized by one of us,<sup>13–16</sup> where the aromatic  $X$ -atoms of halogenated

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**Figure 1.** Molecular structures of metal phthalocyanine (H<sub>16</sub>PcM) and its fluorosubstituted derivatives.

PcM (X<sub>16</sub>PcM) are partially replaced by perfluoro isopropyl (*i*-C<sub>3</sub>F<sub>7</sub>) groups (R<sub>f</sub>). These R<sub>f</sub> groups are much more electronegative than the F atoms and were shown to increase the solubility of the compounds and also promote novel catalytic oxidations while resisting self-oxidation.<sup>13</sup>

In a recent paper,<sup>17</sup> density functional theory (DFT) calculations were carried out to investigate the electronic structure and properties of unligated and ligated F<sub>64</sub>PcFe complexes. One important observation was the change of the ground state of Fe<sup>II</sup> caused by the strongly electron-withdrawing peripheral substituents at the Pc ring. To shed more light on the effects of the R<sub>f</sub> substituents, the previous theoretical study is extended in this report to include systems where M, Co, or Zn is four-, five-, or six-coordinated.

F<sub>64</sub>PcZn has been studied in detail using single-crystal X-ray diffraction, optical and photoacoustic spectroscopy, and cyclic voltammetry, and it exhibits a number of interesting properties. (1) The ultraviolet–visible (UV–vis) spectra of F<sub>64</sub>PcZn differ significantly from those of the parent, unsubstituted H<sub>16</sub>PcZn species,<sup>13</sup> indicating that the electron-withdrawing property of the peripheral substituents affects the ring  $\pi$  molecular orbitals (MOs) to differing extents. (2) There are large shifts in the ring redox (oxidation and reduction) potentials on going from H<sub>16</sub>PcZn to F<sub>64</sub>PcZn, and the increase in oxidation potentials provides extra stability toward oxidative destruction for the R<sub>f</sub>-substituted phthalocyanine.<sup>13</sup> (3) It is difficult to photochemically oxidize F<sub>64</sub>PcZn, in stark contrast with the behavior of H<sub>16</sub>PcZn, for which the photooxidation reaction results in the formation

of a [H<sub>16</sub>PcZn]<sup>+</sup> species.<sup>13</sup> (4) Axial ligation to F<sub>64</sub>PcZn is favored relative to H<sub>16</sub>PcZn; the X-ray structure of the former complex reveals the formation of F<sub>64</sub>PcZn(Ace)<sub>2</sub> (Ace = acetone), but H<sub>16</sub>PcZn as well as F<sub>16</sub>PcZn do not retain solvent when crystallized.<sup>14</sup> Anionic ligands such as Cl<sup>−</sup>, coordinate even more strongly. Thus, electrospray ionization mass spectrometry data for a DMF solution of F<sub>64</sub>PcZn indicated the existence of F<sub>64</sub>PcZn(Cl) and F<sub>64</sub>PcZn(Cl)<sup>−</sup>.<sup>13</sup>

The main goals of the present report are to examine the effects of F<sup>−</sup> and R<sub>f</sub>-substituents on various properties that include electronic structure, oxidation/reduction properties, ionization potentials, electron affinities, etc. This work extends previous semiempirical ZINDO and MOPAC calculations on F<sub>64</sub>PcZn,<sup>13</sup> which were aimed mainly at interpretation of the absorption spectra of the metal complexes. The axial bonding properties of acetone, Cl, or Cl<sup>−</sup> to the metal complexes are also examined. For comparison among the different metals, the results of iron phthalocyanines are also reported here.

## 2. Computational Details

The molecular structure of the parent metal phthalocyanine H<sub>16</sub>PcM is illustrated in Figure 1a and that of (perfluoro *i*-C<sub>3</sub>F<sub>7</sub>)-(perfluoro)phthalocyanine, F<sub>64</sub>PcM, is in Figure 1d. For computational convenience, the isopropyl –C<sub>3</sub>F<sub>7</sub> substituent was modeled by ethyl –C<sub>2</sub>F<sub>5</sub>. The use of the smaller F<sub>48</sub>PcM (Figure 1c) as an accurate electronic mimic of the larger F<sub>64</sub>PcM was demonstrated in our previous calculations.<sup>17</sup> This simplification is also supported by earlier

semiempirical ZINDO calculations,<sup>13</sup> which show that the spectrum predicted for F<sub>48</sub>PcZn is almost identical to that of F<sub>64</sub>PcZn. For the sake of comparison and completion, the parent perfluorinated metal phthalocyanines, F<sub>16</sub>PcM (Figure 1b), have also been considered here.

All calculations were carried out using the Amsterdam Density Functional (ADF) program package (version 2000.02) developed by Baerends and co-workers.<sup>18–21</sup> A triple- $\zeta$  STO basis was used for the metal 3s–4s shells plus one 4p polarization function, a triple- $\zeta$  basis for C/N/O 2s–2p and Cl 3s–3p shells plus one 3d polarization function, a double- $\zeta$  basis for F 2s–2p shells, and a double- $\zeta$  basis for the H 1s shell. It has been shown that high-quality basis sets (triple- $\zeta$  plus one polarization function) are required for the atoms within the macrocycle ring of the phthalocyanine in order to obtain the correct ground states of H<sub>16</sub>PcFe and its derivatives.<sup>22</sup> The inner orbitals, i.e., 1s–2p for Fe/Cl and 1s for C/N/O/F, were considered as core and kept frozen according to the frozen-core approximation.<sup>18</sup> Among the various exchange-correlation potentials available, the density-parametrization form of Vosko, Wilk, and Nusair (VWN)<sup>23</sup> plus Becke's gradient correction for exchange (B)<sup>24</sup> and Perdew's gradient correction for correlation (P)<sup>25</sup> were employed. The combined VWN-B-P functional has been shown to provide accurate bonding energies for both main-group<sup>26</sup> and transition metal<sup>27</sup> systems. Relativistic corrections of the valence electrons were calculated by the quasi-relativistic (QR) method.<sup>28</sup> For the open-shell states, the unrestricted Kohn–Sham (UKS) spin-density functional approach was adopted.

Electron excitation energies related to the electronic absorption spectra were calculated using the time-dependent density functional response theory (TDDFT)<sup>29</sup> as implemented in the ADF program. TDDFT provides a first-principles method for the calculation of excitation energies and presents an excellent alternative to the conventional highly correlated configuration interactions (CI) method. Applications of TDDFT to excitation energy calculations can be found in recent work.<sup>30,31</sup>

### 3. Results and Discussion

Unsubstituted transition-metal H<sub>16</sub>PcMs have been shown to have square planar *D*<sub>4h</sub> symmetry.<sup>30</sup> X-ray crystal structure data<sup>13–16</sup> indicate that the Pc ring in substituted F<sub>64</sub>PcM is quite planar, and this ring planarity is maintained in the solid state even in the presence of axial acetone ligands. The solution UV–vis spectra<sup>32a–c,33</sup> are also consistent with *D*<sub>4h</sub> symmetry. The invariability of the MN<sub>4</sub> chromophore geometry upon substitution of the peripheral H-atoms by *i*-C<sub>3</sub>F<sub>7</sub> groups strongly suggests that the geometry will not change for steric reasons when the *i*-C<sub>3</sub>F<sub>7</sub> groups are replaced by less bulky groups in F<sub>16</sub>PcM and F<sub>48</sub>PcM. This argument is indeed supported by calculations.<sup>13</sup> Placing the molecule in the *xy* plane, the five metal 3d-orbitals transform as a<sub>1g</sub> (d<sub>z<sup>2</sup></sub>), b<sub>1g</sub> (d<sub>x<sup>2</sup>–y<sup>2</sup></sub>), b<sub>2g</sub> (d<sub>xy</sub>), and e<sub>g</sub> (d<sub>z</sub>, i.e., d<sub>xz</sub> and d<sub>yz</sub>). For Fe<sup>II</sup> and Co<sup>II</sup>, different occupations of electrons in these d-orbitals may yield a number of possible low-lying states. To determine the ground state, relative energies of several selected configurations in H<sub>16</sub>PcM, F<sub>16</sub>PcM, and F<sub>48</sub>PcM (M

**Table 1.** Calculated Relative Energies (*E*, eV) for Selected Configurations in H<sub>16</sub>PcM, F<sub>16</sub>PcM, and F<sub>48</sub>PcM (M = Fe, Co)

configuration <sup>a</sup>				<i>E</i> (R) <sup>b</sup>		
b <sub>2g</sub> / d <sub>xy</sub>	a <sub>1g</sub> / d <sub>z<sup>2</sup></sub>	1e <sub>g</sub> / d <sub>z</sub>	state	H <sub>16</sub> PcM	F <sub>16</sub> PcM	F <sub>48</sub> PcM
M = Fe <sup>c</sup>						
2	2	2	<sup>3</sup> A <sub>2g</sub>	<b>0</b> (1.916)	<b>0</b> (1.918)	0 (1.919)
2	1	3	<sup>3</sup> E <sub>g</sub> (A)	0.05 (1.923)	0.01 (1.922)	–0.04 (1.921)
1	1	4	<sup>3</sup> B <sub>2g</sub>	0.08 (1.921)	0.05 (1.919)	<b>–0.06</b> (1.917)
1	2	3	<sup>3</sup> E <sub>g</sub> (B)	0.52 (1.910)	0.52 (1.911)	0.45 (1.910)
2	0	4	<sup>1</sup> A <sub>1g</sub>	1.43 (1.935)	1.37 (1.934)	1.28 (1.933)
M = Co						
2	1	4	<sup>2</sup> A <sub>1g</sub>	<b>0</b> (1.922)	<b>0</b> (1.924)	<b>0</b> (1.931)
2	2	3	<sup>2</sup> E <sub>g</sub>	0.21 (1.909)	0.24 (1.911)	0.28 (1.916)
1	2	4	<sup>2</sup> B <sub>2g</sub>	0.95 (1.901)	0.94 (1.904)	0.88 (1.904)

<sup>a</sup> Orbital energy levels illustrated in Figure 2. <sup>b</sup> Values in parentheses refer to optimized M–N(Pc) bond length (in Å) for the pertinent state. <sup>c</sup> The calculated *E*'s and *R*'s with the present version ADF program (2000.02) may be slightly different from those obtained with the older version ADF (2.0.1), ref 17, but the trends of the results are not changed.

= Fe, Co) were calculated, wherein geometry optimization was performed separately for each state considered. These energies are reported in Table 1, along with the optimized M–N bond lengths.

Table 2 displays the gross populations of M 3d, 4s, and 4p orbitals (in the ground state), along with the metal's Mulliken atomic charge. Table 3 lists the calculated M–Pc binding energies (*E*<sub>bind</sub>), ionization potentials (IP) for several outer MOs, and the electron affinities (EA). *E*<sub>bind</sub> is defined as the energy required to pull the metal apart from the Pc ring

$$-E_{\text{bind}} = E(\text{PcM}) - \{E(\text{M}) + E(\text{Pc})\}$$

where *E*(PcM), *E*(M), and *E*(Pc) represent the total energies of the indicated species. (The geometries of PcM and Pc are independently optimized.) The IPs and EAs were calculated by the so-called  $\Delta$ SCF method which carries out separate SCF (self-consistent field) calculations for the neutral molecule and its ion, where EA = *E*(X<sup>–</sup>) – *E*(X). The computed relative energies of selected configurations of the ligated iron phthalocyanines are contained in Table 4, and their properties are presented in Table 5.

The electronic structure of H<sub>16</sub>PcFe has been the subject of several experimental studies.<sup>34–36</sup> A <sup>3</sup>B<sub>2g</sub> ground state was originally suggested for H<sub>16</sub>PcFe on the basis of magnetic work,<sup>35</sup> but later magnetic circular dichroism spectra have shown that the ground state is in fact <sup>3</sup>A<sub>2g</sub>.<sup>36</sup> Our previous calculations<sup>37</sup> support the latter assignment. In a recent paper of ours,<sup>17</sup> we further calculated F<sub>16</sub>PcFe, F<sub>48</sub>PcFe, and their complexes with two axial ligands L (L = Ace, H<sub>2</sub>O, pyridine). It is shown that the electronic configuration of each ligated PcFe is determined mainly by the axial ligand-field strength but can also be affected by peripheral substituents. Figure 2 illustrates the changes of the electronic structure from H<sub>16</sub>PcFe to F<sub>16</sub>PcFe to F<sub>48</sub>PcFe. A detailed discussion of the results for the iron–phthalocyanine complexes is reported in ref 17.<sup>38</sup>

**Table 2.** Mulliken Orbital Populations and Atomic Charges on Metal ( $Q_M$ )

	H <sub>16</sub> PcFe	F <sub>16</sub> PcFe	F <sub>48</sub> PcFe	H <sub>16</sub> PcCo	F <sub>16</sub> PcCo	F <sub>48</sub> PcCo	H <sub>16</sub> PcZn	F <sub>16</sub> PcZn	F <sub>48</sub> PcZn
3d	6.58	6.58	6.58	7.60	7.60	7.60	10.00	10.00	10.00
4s	0.42	0.38	0.27	0.32	0.28	0.27	0.55	0.53	0.53
4p	0.25	0.25	0.26	0.42	0.42	0.42	0.82	0.81	0.80
$Q_M$	0.75	0.80	0.88	0.65	0.70	0.70	0.63	0.67	0.67

**Table 3.** Calculated M–N Bond Lengths ( $R_{M-N}$ ), M–Pc Binding Energies ( $E_{bind}$ ), Ionization Potentials (IP), and Electron Affinities (EA) in the Ground State of the Systems

		H <sub>16</sub> PcM	F <sub>16</sub> PcM	F <sub>48</sub> PcM <sup>a</sup>
M = Fe				
$R_{Fe-N}$ , Å		1.916	1.918	1.917
$E_{bind}$ , eV		9.77	9.53	9.95
IP, eV	$a_{1u}$	6.39 (first)	7.67 (first)	8.76 (first)
	$a_{1g}/d_{z^2}$	6.42	7.70	10.63
	$b_{2g}/d_{xy}$	6.58	7.86	9.91
	$1e_g/d_{\pi}$	7.23	8.48	8.87
	$b_{1u}$	8.00	8.88	9.75
EA, eV	$1e_g$	−2.55	−3.90	−5.47 ( $b_{2g}$ )
M = Co				
$R_{Co-N}$ , Å		1.922	1.924	1.931
$E_{bind}$ , eV		11.51	11.29	11.71
IP, eV	$a_{1u}$	6.42 (first)	7.69 (first)	8.72 (first)
	$1e_g/d_{\pi}$	7.19	8.45	9.42
	$b_{2g}/d_{xy}$	7.51	8.73	9.53
EA, eV	$a_{1g}/d_{z^2}$	−2.94	−4.26	−5.35
	$2e_g$	−2.09	−3.44	−4.68
M = Zn				
$R_{Zn-N}$ , Å		2.000	2.001	2.006
$E_{bind}$ , eV		5.64	5.50	5.89
IP, eV	$a_{1u}$	6.44 (first)	7.70 (first)	8.72 (first)
EA, eV	$2e_g$	−2.18	−3.52	−4.74

<sup>a</sup> Note that the ground state of F<sub>48</sub>PcFe is different from those of H<sub>16</sub>PcFe and F<sub>16</sub>PcFe.

**3.1. Cobalt Phthalocyanines.** While the peripheral substituents have substantial influence on the electronic structure of the iron phthalocyanines, the Co analogues are much less sensitive. H<sub>16</sub>PcCo, F<sub>16</sub>PcCo, and F<sub>48</sub>PcCo complexes all have a <sup>2</sup>A<sub>1g</sub> ground state, with a fully occupied  $1e_g/d_{\pi}$  level. The <sup>2</sup>E<sub>g</sub> state, arising from the  $(a_{1g})^2(1e_g)^3$  configuration, lies 0.2–0.3 eV higher in energy. The <sup>2</sup>B<sub>2g</sub> state lies considerably higher, nearly a full eV above the ground state. The <sup>2</sup>E<sub>g</sub> – <sup>2</sup>A<sub>1g</sub> energy gap increases gradually from H<sub>16</sub>PcCo to F<sub>16</sub>PcCo to F<sub>48</sub>PcCo. This comparison between Fe and Co is illustrated more explicitly in Figure 3, where it may be noted first that the metal d-orbitals lie much lower in H<sub>16</sub>PcCo than in H<sub>16</sub>PcFe. Unlike the Fe case, the HOMO in PcCo (Pc  $a_{1u}$ ) is no longer a metal 3d-orbital, and the unoccupied  $b_{1g}$  ( $d_{x^2-y^2}$ ) lies below the  $b_{1u}$  orbital of the ring.

Similar to H<sub>16</sub>PcFe, the first ionization of PcCo also occurs from the Pc  $a_{1u}$  orbital, and the first IP (6.42 eV) is very close to that obtained for H<sub>16</sub>PcFe (6.39 eV). The calculated results are in agreement with experimental gas-phase photoelectron spectra (PES) of a series of metal phthalocyanines H<sub>16</sub>PcM with M = Mg, Fe, Co, Ni, Cu, and Zn,<sup>39</sup> which find a sharp first IP at ~6.4 eV for all the H<sub>16</sub>PcMs and conclude that the orbitals of the first ionization are ringlike

and not metal 3d-like in all cases. While for H<sub>16</sub>PcFe the IPs from  $a_{1u}$  and  $a_{1g}/d_{z^2}$  are very close, for H<sub>16</sub>PcCo the IP from  $a_{1u}$  is 0.77 eV smaller than that from a metal 3d-orbital ( $d_{\pi}$ ), and so the one-electron oxidation of unligated H<sub>16</sub>PcCo clearly occurs from the Pc ring. On the other hand, electrochemical studies of cobalt phthalocyanines in solution<sup>40,41</sup> find that PcCo<sup>II</sup> is oxidized to [PcCo<sup>III</sup>]<sup>+</sup>; i.e., one-electron oxidation of PcCo occurs from the metal. For example, Lever et al.<sup>40</sup> observed that oxidation of Co<sup>II</sup> phthalocyanine by halogen (X) leads to the formation of a complex containing the trivalent metal PcCo<sup>III</sup>(X)<sub>2</sub>. The disagreement between the calculation and these sorts of experiments may be attributed to the effect of solvent. In solution, the cation species is ligated. The axial ligands raise the energy of the metal  $a_{1g}/d_{z^2}$  orbital (see Figure 3) so that the electron in this orbital may be ionized first. Our calculations on H<sub>16</sub>PcCo(Py)<sub>2</sub> (Py = pyridine) show that the first ionization for such a complex indeed occurs from the metal  $a_{1g}/d_{z^2}$  orbital (see Table 5). Our further calculations on H<sub>16</sub>PcCo(Cl)<sub>2</sub> show that the latter complex has a ground state of <sup>2</sup>A<sub>1u</sub> [ $(1e_g/d_{\pi})^4(a_{1u})^1(a_{1g}/d_{z^2})^0$ ], where the oxidation state of Co is III, in agreement with experiment;<sup>40,41</sup> a configuration of  $(1e_g/d_{\pi})^3(a_{1u})^2(a_{1g}/d_{z^2})^0$  is 0.57 eV higher in energy than the ground state.

The similarities between Fe and Co extend (trends are preserved upon fluorination) to the various degrees of fluorosubstitution reported in Table 3. On the other hand, the calculated EAs of PcCo are generally more negative than those of PcFe. For the reduction of PcCo, the added electron goes into the  $a_{1g}/d_{z^2}$  orbital, while it enters the  $1e_g/d_{\pi}$  orbital for H<sub>16</sub>PcFe and F<sub>16</sub>PcFe. Since F<sub>48</sub>PcFe has a different electronic structure (as compared to H<sub>16</sub>PcFe/F<sub>16</sub>PcFe), the calculated EA of this molecule is even slightly larger (0.12 eV) than that of F<sub>48</sub>PcCo. The calculated Co–Pc binding energy is significantly larger (~1.7 eV) than that of Fe–Pc, despite the fact that the Co–N bond lengths are somewhat (0.01 Å) longer than Fe–N. Note that for both Fe and Co, it is the F<sub>16</sub>PcM species that has the lowest M–Pc binding energy.

The perturbations caused by the Ace ligands in the MO energy diagrams of PcM are illustrated in Figure 3 with PcFe-(Ace)<sub>2</sub> on the left and PcCo(Ace)<sub>2</sub> on the right. The coordination of the two axial ligands lowers the symmetry of the system from  $D_{4h}$  to  $D_{2h}$  and splits the  $d_{xz}$ – $d_{yz}$  degeneracy. Acetone is an electron-donating ligand, shifting the MOs upward. The  $a_{1g}/d_{z^2}$  orbital is particularly raised, owing to the strong repulsive interaction between the ligand HOMO and the metal  $d_{z^2}$ . As a result, the IPs of PcM(Ace)<sub>2</sub> are notably decreased as compared to those of PcM, suggesting that the former will be easier to oxidize than unligated PcM. Unlike PcFe(Ace)<sub>2</sub> where the first ionization



**Table 4.** Calculated Relative Energies ( $E$ , eV) for Selected Configurations in  $H_{16}PcFe(Ace)_2$ ,  $F_{16}PcFe(Ace)_2$ , and  $F_{48}PcFe(Ace)_2$ 

configuration					$E_{\text{relative}}$		
$d_{xy}$	$d_{z^2}$	$d_{yz}$	$d_{xz}$	state <sup>a,b</sup>	$H_{16}PcFe(Ace)_2$	$F_{16}PcFe(Ace)_2$	$F_{48}PcFe(Ace)_2$
2	1	2	1	$^3B_{2g} [^3E_g(A)]$	<b>0</b> (1.933/2.373) <sup>c</sup>	<b>0</b> (1.939/2.300)	<b>0</b> (1.937/2.279)
2	1	1	2	$^3B_{3g} [^3E_g(A)]$	0.05 (1.928/2.483)	0.07 (1.932/2.452)	0.09 (1.933/2.405)
1	1	2	2	$^3B_{1g} (^3B_{2g})$	0.15 (1.925/2.503)	0.18 (1.926/2.468)	0.17 (1.924/2.437)
2	0	2	2	$^1A_{1g} (^1A_{1g})$	0.16 (1.946/1.969)	0.12 (1.947/1.973)	0.06 (1.942/1.974)

<sup>a</sup> States in parentheses are the corresponding states in unligated species. <sup>b</sup> No minimum or very long Fe–O(Ace) distance was found for the  $(d_{xy})^2(d_{z^2})^2(d_{yz})^1(d_{xz})^1 - ^3B_{1g} (^3A_{2g})$  and  $(d_{xy})^1(d_{z^2})^2(d_{yz})^2(d_{xz})^1 - ^3B_{3g} [^3E_g(B)]$  states. <sup>c</sup> The values in parentheses represent the optimized Fe–N(Pc) and Fe–O(Ace) bond lengths (in Å), respectively.

**Table 5.** Calculated Properties<sup>a</sup> of  $H_{16}PcM$ ,  $F_{16}PcM$ , and  $F_{48}PcM$  with Two Axial Ligands (L) at the Ground State (M = Fe, Co, Zn; L = Ace, Py)

	$H_{16}PcM(L)_2$ L = Ace	$F_{16}PcM(L)_2$ L = Ace	$F_{48}PcM(L)_2$ L = Ace	$H_{16}PcCo(L)_2$ L = Py
M = Fe				
$R_{M-N(Pc)}$ (Å)	1.933	1.939	1.937	
$R_{M-O(Ace)}$ (Å)	2.373	2.300	2.279	
$E_{\text{bind}}[PcM-(L)_2]$ (eV)	0.25	0.63	0.94	
$Q_M$ (e)	0.79	0.81	0.80	
$Q_{\text{Ace}}$ (e)	0.15	0.19	0.22	
IP (eV)				
	$a_{1g}/d_{z^2}$	6.97	8.01	8.86
	$1b_{3g}/d_{yz}$	<b>5.74</b> (first)	<b>6.99</b> (first)	<b>7.94</b> (first)
	$b_{1g}/d_{xy}$	6.01	7.26	8.14
	$a_{1u}$	6.00	7.21	8.27
EA (eV)				
	$1b_{2g}/d_{xz}$	−2.29	−3.55	−4.72
	$2b_{3g}$	−1.66	−2.96	−4.22
M = Co				
$R_{M-N(Pc)}$ (Å)	1.926	1.929	1.932 (1.924) <sup>c</sup>	1.934
$R_{M-O(Ace)}$ (Å)	2.401	2.366	2.331 (2.314)	2.318
$E_{\text{bind}}[PcM-(L)_2]$ (eV)	0.16	0.49	1.06	0.77
$Q_M$ (e)	0.67	0.70	0.69	0.90
$Q_{\text{Ace}}$ (e)	0.13	0.14	0.16	0.18
IP (eV)				
	$a_{1g}/d_{z^2}$	7.56	8.69	9.45
	$a_{1u}$	<b>6.03</b> (first)	<b>7.25</b> (first)	<b>8.31</b> (first)
EA (eV)				
	$2b_{3g}$	−1.67	−2.98	−4.25
	$a_{1g}/d_{z^2}$	−1.63	−2.82	−3.87
$E_{\text{dis}}^d$ (eV)		−1.11	−0.79	−0.33
M = Zn				
$R_{M-N(Pc)}$ (Å)	2.010	2.014	2.012 (2.006) <sup>c</sup>	
$R_{M-O(Ace)}$ (Å)	2.505	2.459	2.414 (2.411)	
$E_{\text{bind}}[PcM-(L)_2]$ (eV)	0.12	0.44	0.73	
$Q_M$ (e)	0.64	0.67	0.65	
$Q_{\text{Ace}}$ (e)	0.10	0.12	0.13	
IP (eV)				
	$a_{1u}$	<b>6.05</b> (first)	<b>7.26</b> (first)	<b>8.30</b> (first)
EA (eV)				
	$2b_{3g}$	−1.78	−3.08	−4.32

<sup>a</sup>  $R$ : distance,  $E_{\text{bind}}[PcM-(Ace)_2]$ : binding energy between PcM and two Ace ligands,  $Q$ : charge distribution, IP: ionization potential, EA: electron affinity. <sup>b</sup> The first IP is indicated in bold. <sup>c</sup> The values in parentheses are experimental distances for  $F_{64}PcM(Ace)_2$  (refs 14 and 15).

<sup>d</sup> Dissociation energy for  $[PcCo(L)_2]^- \rightarrow [PcCo]^- + 2L$ .

occurs from the central metal ( $1b_{3g}/d_{yz}$ ), Table 5 reveals that of  $PcCo(Ace)_2$  takes place still from the Pc  $a_{1u}$  orbital, similar to unligated PcM. As mentioned above, in the presence of relatively strong field ligands L (e.g. L = Py), the first ionization of  $PcCo(L)_2$  takes place from the central metal. For  $PcCo(Ace)_2$ , the IP from  $a_{1g}/d_{z^2}$  is more than 1 eV larger than that from  $a_{1u}$ . There is about a 0.4 eV decrease of IP from PcCo to  $PcCo(Ace)_2$ , in contrast to  $\sim 0.8$  eV decrease of IP from PcFe to  $PcFe(Ace)_2$ . Note also that for any degree

of substitution, the IP of the  $PcCo(Ace)_2$  is roughly 0.3 eV higher than the corresponding quantity for the Fe analogue.

The axial ligands reduce the electron affinity of PcM considerably (by 1.1–1.3 eV) when M is Co. In contrast to PcCo, the added electron in  $PcCo(Ace)_2$  now occupies a high-lying antibonding Pc  $2b_{3g}$  orbital. Since the added electron in PcCo goes into a low-lying metal orbital,  $[PcCo]^-$  has a relatively low energy. Unless the ligand field of L is very strong, the  $[PcM(L)_2]^-$  complex is expected to be unstable



yielding a  $^2A_{1u}$  ground state. This IP is strikingly similar in magnitude to the Co cases. In the case of reduction, the electron is accommodated in the LUMO  $e_g$  ( $Pc \pi^*$ ). Both IP and EA are increased considerably by the  $-C_2F_5$  substituents, in agreement with the trend in the electrochemical oxidation and reduction potentials.<sup>13</sup> The EA values are slightly more negative for the Zn molecules than for Co.

With two electrons in the  $b_{1g}$  ( $d_{x^2-y^2}$ ) orbital in  $PcZn$ , the repulsive interaction between these electrons and those on the pyrrole nitrogens lengthens the M–N bond length by  $\sim 0.08$  Å relative to Co and considerably decreases the M–Pc binding energy, as is evident by the  $E_{bind}$  entries in Table 3 for  $M = Zn$ .

On the other hand, the presence of two electrons in the  $a_{1g}/d_{z^2}$  orbital in  $PcZn(Ace)_2$  gives rise to longer axial Zn–O(Ace) bond lengths as compared to Co–O(Ace). Therefore the binding energy of  $PcZn-(Ace)_2$  is smaller than the  $PcCo-(Ace)_2$  value. Also, there is a large increase of the  $PcZn-(Ace)_2$  binding energy when more F atoms are added to the system. It was argued that the perfluoro peripheral substituents withdraw electron density from the inner ring and the central metal so that the axial ligation of  $F_{64}PcZn$  is enhanced relative to the parent  $H_{16}PcZn$  complex.<sup>13</sup> According to the calculation, the presence of the F-atoms in  $F_{16}PcZn$  makes the atomic charge on Zn more positive by 0.04 (see Table 2), but no change in  $Q_M$  is found from  $F_{16}PcZn$  to  $F_{48}PcZn$ . The same is true for the cobalt systems. There is a more pronounced shortening of the axial M–O(Ace) bond distance from  $H_{16}PcZn$  to  $F_{48}PcZn$  than from  $H_{16}PcCo$  to  $F_{48}PcCo$ . Again the calculated bond lengths  $R_{Zn-N(Pc)} = 2.012$  Å and  $R_{Zn-O(Ace)} = 2.414$  Å in  $F_{48}PcZn$  agree very well with the 2.006 Å and 2.411 Å measured in the  $F_{64}PcZn$  crystal.<sup>14</sup> An addition of  $(Ace)_2$  to  $PcZn$  produces little geometric change, similar to the  $PcCo$  case.

The MO levels are shifted up from  $F_{48}PcZn$  to  $F_{48}PcZn-(Ace)_2$ . The increase in the energy of  $a_{1u}$  leads to a relatively small IP from this orbital. Therefore, solvent coordination to the central metal is able to decrease the oxidation potential of  $PcM$ , as observed experimentally.<sup>13</sup> While attempts to oxidize  $F_{64}PcZn$  were unsuccessful, there is evidence that in the presence of strong electron-donating ligands such as imidazole,  $F_{64}PcZn$  can be oxidized electrochemically.<sup>13</sup> The experimental observation can be accounted for by the calculated results. Here the IP and EA of  $PcZn(Ace)_2$  are about 0.4 eV smaller (lower) than those of  $PcZn$ .

**3.2.2. Electronic Spectra.** Electronic spectra of both  $H_{16}PcZn$  and  $F_{48}PcZn$  have been calculated using the semiempirical ZINDO method,<sup>13</sup> with the aim of understanding the changes in energy of the spectral features when the electron-withdrawing peripheral substituents ( $-C_2F_5$ ) are introduced. However, the semiempirical method yields some excitation energies which are considerably too large (more than 1 eV) as compared to the experimental spectra. In this study, we describe the results of TDDFT calculations for the series of  $PcZn$  complexes from  $H_{16}PcZn$  to  $F_{16}PcZn$  to  $F_{48}PcZn$  to  $F_{48}PcZn(Ace)_2$ . Our calculated excitation energies and oscillator strengths for several lowest, spin-allowed singlet  $^1E_u$  excited states are displayed in Table 6, together with available experimental data.<sup>13,33</sup>

The experimental spectrum of  $PcZn$  is characterized by an intense absorption band in the visible (Q-band) and two other strong absorption bands in the near-UV ( $B_1$  and  $B_2$  bands).<sup>33</sup> The multiple bands in the B region is supported by the calculations. The Q-band is assigned to the  $1^1E_u$  state, which is nearly pure (90%) HOMO  $\rightarrow$  LUMO transition. For  $H_{16}PcZn$ ,  $1^1E_u$  is calculated to be 1.96 eV, in good agreement with the experimental value of 1.85 eV.<sup>33</sup> The close-lying  $6^1E_u$  and  $7^1E_u$  states are responsible for the  $B_2$  and  $B_1$  bands, respectively; their oscillator strengths are both quite large. Different from  $1^1E_u$ , the latter two states involve significant mixture of several transitions. Again, the calculated  $E^{exc}$  values (3.40 and 3.74 eV) for the B bands agree very well with experiment (3.16 and 3.71 eV).<sup>33</sup> The same is true for  $F_{48}PcZn$ .

The  $1^1E_u$  state (Q-band) in  $F_{16}PcZn$  is predicted at 1.84 eV, a red shift of 0.12 eV relative to that of  $H_{16}PcZn$ . For the B bands, the red shift is somewhat larger, about 0.2 eV. These results reveal that the introduction of the electron-withdrawing F-atoms into the Pc ring shifts the Q and B bands to the red.

Turning to  $F_{48}PcZn$ , its  $1^1E_u$  has nearly the same  $E^{exc}$  as that of  $F_{16}PcZn$ . So the partial replacement of eight F-atoms by eight  $-C_2F_5$  groups does not shift the Q-band. Here the main effect of the  $-C_2F_5$  substituents is a notable red shift of the  $B_1$  band.

From  $F_{48}PcZn$  to  $F_{48}PcZn(Ace)_2$ , the  $E^{exc}$  values for the various  $^1E_u$  states remain nearly unchanged, indicating that the presence of axial acetone ligands does not change the positions of the absorption bands. These results are also in accord with experimental observations.<sup>13</sup> Note that in Figure 4, the LUMO–HOMO energy gap is not shown to be changed from one system to another, but the peripheral substituents may still change the calculated  $E^{exc}$  for the HOMO  $\rightarrow$  LUMO transition. This is because the  $1^1E_u$  state is not a pure HOMO  $\rightarrow$  LUMO transition; it contains some contributions ( $\sim 10\%$ ) from other transitions owing to configuration interaction. As a matter of fact, the change in  $E^{exc}$  is insignificant from one system to another.

Similar to the calculations on  $H_{16}PcNi$  by Rosa et al.,<sup>31</sup> our TDDFT calculations on  $PcZn$  also predict several allowed  $\pi \rightarrow \pi^*$  transitions located between the Q and B bands. These  $\pi \rightarrow \pi^*$  transitions may be responsible for the intensity and broadness of the B bands. The ZINDO calculations predicted a HOMO  $\rightarrow$  LUMO+1 transition (called second  $\pi \rightarrow \pi^*$  transition) located to the red of the B band for both  $H_{16}PcZn$  and  $F_{48}PcZn$ . Our results show that the HOMO  $\rightarrow$  LUMO+1 transition is located to the blue of the B band for  $F_{48}PcZn$ . No analogous transition is found for  $F_{16}PcZn$ .

**3.2.3.  $PcZn(Cl)$  and  $PcZn(Cl^-)$  Complexes.** Finally, calculations were also performed on the various  $PcZn(Cl)$  and  $PcZn(Cl^-)$  complexes. When only a single axial ligand is attached to the system, significant out-of-plane displacement of the metal is expected and in fact observed. For each molecule, the geometry was optimized assuming  $C_{4v}$  symmetry. The calculated  $PcZn-L$  binding energies ( $L = Cl, Cl^-$ ) and three critical coordination parameters in  $PcZn(L)$  are displayed in Table 7.  $R_{Cl \cdots N}$  (distance between the center of the ring and pyrrole nitrogen atom) is a measure of the ring

**Table 6.** Calculated Excitation Energies ( $E^{\text{exc}}$ , in eV) and Oscillator Strengths ( $f$ )

system	state	contribution <sup>a</sup>	$E^{\text{exc}}$ , eV		$f$	assignment
			calcd	exptl <sup>c</sup>		
H <sub>16</sub> PcZn	1 <sup>1</sup> E <sub>u</sub>	91% (2a <sub>1u</sub> → 7e <sub>g</sub> )	1.96	1.85	0.6482	Q
	2 <sup>1</sup> E <sub>u</sub>	97% (3b <sub>2u</sub> → 7e <sub>g</sub> )	2.80		0.0280	$\pi \rightarrow \pi^*$
	3 <sup>1</sup> E <sub>u</sub>	39% (2b <sub>1u</sub> → 7e <sub>g</sub> ); 33% (5a <sub>2u</sub> → 7e <sub>g</sub> ); 25% (5a <sub>2u</sub> → 7e <sub>g</sub> )	3.02		0.0014	$\pi \rightarrow \pi^*$
	4 <sup>1</sup> E <sub>u</sub>	54% (2b <sub>1u</sub> → 7e <sub>g</sub> ); 23% (4a <sub>2u</sub> → 7e <sub>g</sub> ); 18% (5a <sub>2u</sub> → 7e <sub>g</sub> )	3.07		0.3390	$\pi \rightarrow \pi^*$
	5 <sup>1</sup> E <sub>u</sub>	76% (2a <sub>1u</sub> → 8e <sub>g</sub> )	3.35		0.0242	$\pi \rightarrow \pi^*$
	6 <sup>1</sup> E <sub>u</sub>	36% (1a <sub>1u</sub> → 7e <sub>g</sub> ); 23% (4a <sub>2u</sub> → 7e <sub>g</sub> ); 18% (2a <sub>1u</sub> → 8e <sub>g</sub> ); 17% (5a <sub>2u</sub> → 7e <sub>g</sub> )	3.40	3.16	0.6840	B <sub>2</sub>
	7 <sup>1</sup> E <sub>u</sub>	53% (1a <sub>1u</sub> → 7e <sub>g</sub> ); 17% (4a <sub>2u</sub> → 7e <sub>g</sub> ); 16% (5a <sub>2u</sub> → 7e <sub>g</sub> )	3.74	3.71	1.0402	B <sub>1</sub>
	8 <sup>1</sup> E <sub>u</sub>	95% (6e <sub>g</sub> → 3b <sub>1u</sub> )	4.22		0.0332	
	1 <sup>1</sup> E <sub>u</sub>	91% (4a <sub>1u</sub> → 11e <sub>g</sub> )	1.84		0.6332	Q
	2 <sup>1</sup> E <sub>u</sub>	84% (4b <sub>1u</sub> → 11e <sub>g</sub> ); 11% (5b <sub>2u</sub> → 11e <sub>g</sub> )	2.60		0.1442	$\pi \rightarrow \pi^*$
F <sub>16</sub> PcZn	3 <sup>1</sup> E <sub>u</sub>	87% (5b <sub>2u</sub> → 11e <sub>g</sub> ); 10% (4b <sub>1u</sub> → 11e <sub>g</sub> )	2.71		0.1360	$\pi \rightarrow \pi^*$
	4 <sup>1</sup> E <sub>u</sub>	80% (7a <sub>2u</sub> → 11e <sub>g</sub> ); 15% (6a <sub>2u</sub> → 11e <sub>g</sub> )	2.88		0.2928	$\pi \rightarrow \pi^*$
	5 <sup>1</sup> E <sub>u</sub>	67% (3a <sub>1u</sub> → 11e <sub>g</sub> ); 29% (6a <sub>2u</sub> → 11e <sub>g</sub> )	3.16		0.1038	B <sub>2</sub>
	6 <sup>1</sup> E <sub>u</sub>	43% (6a <sub>2u</sub> → 11e <sub>g</sub> ); 23% (3a <sub>1u</sub> → 11e <sub>g</sub> )	3.58		1.4316	B <sub>1</sub>
	7 <sup>1</sup> E <sub>u</sub>	83% (10e <sub>g</sub> → 5b <sub>1u</sub> )	4.00		0.0220	
	8 <sup>1</sup> E <sub>u</sub>	89% (9e <sub>g</sub> → 5b <sub>1u</sub> )	4.15		0.0786	
	1 <sup>1</sup> E <sub>u</sub>	91% (12a <sub>1u</sub> → 27e <sub>g</sub> )	1.83	1.80	0.7984	Q
	2 <sup>1</sup> E <sub>u</sub>	95% (12b <sub>1u</sub> → 27e <sub>g</sub> )	2.51		0.3138	$\pi \rightarrow \pi^*$
F <sub>48</sub> PcZn	3 <sup>1</sup> E <sub>u</sub>	75% (13b <sub>2u</sub> → 27e <sub>g</sub> ); 13% (11a <sub>1u</sub> → 27e <sub>g</sub> )	2.89		0.1374	$\pi \rightarrow \pi^*$
	4 <sup>1</sup> E <sub>u</sub>	59% (11a <sub>1u</sub> → 27e <sub>g</sub> ); 22% (13b <sub>2u</sub> → 27e <sub>g</sub> ); 17% (15a <sub>2u</sub> → 27e <sub>g</sub> )	2.95		0.0086	$\pi \rightarrow \pi^*$
	5 <sup>1</sup> E <sub>u</sub>	91% (14a <sub>2u</sub> → 27e <sub>g</sub> )	3.13	2.89	0.0124	B <sub>2</sub>
	6 <sup>1</sup> E <sub>u</sub>	48% (15a <sub>2u</sub> → 27e <sub>g</sub> ); 20% (12a <sub>1u</sub> → 28e <sub>g</sub> ); 12% (11a <sub>1u</sub> → 27e <sub>g</sub> )	3.33	3.10	1.7998	B <sub>1</sub>
	7 <sup>1</sup> E <sub>u</sub>	62% (12a <sub>1u</sub> → 28e <sub>g</sub> ); 27% (26e <sub>g</sub> → 13b <sub>1u</sub> )	3.41		0.0356	$\pi \rightarrow \pi^*$
	8 <sup>1</sup> E <sub>u</sub>	60% (26e <sub>g</sub> → 13b <sub>1u</sub> ); 12% (12a <sub>1u</sub> → 28e <sub>g</sub> )	3.62		0.0058	
	9 <sup>1</sup> E <sub>u</sub>	70% (25e <sub>g</sub> → 13b <sub>1u</sub> ); 22% (12a <sub>1u</sub> → 29e <sub>g</sub> )	4.02		0.0700	
	1 <sup>1</sup> E <sub>u</sub>	90% (25a <sub>1u</sub> → 30e <sub>g</sub> )	1.87		0.8075	Q
	2 <sup>1</sup> E <sub>u</sub>	91% (24a <sub>1u</sub> → 30e <sub>g</sub> )	2.65		0.2024	$\pi \rightarrow \pi^*$
	3 <sup>1</sup> E <sub>u</sub>	62% (33b <sub>1u</sub> → 30e <sub>g</sub> ); 22% (23a <sub>1u</sub> → 30e <sub>g</sub> )	2.86		0.5055	$\pi \rightarrow \pi^*$
F <sub>48</sub> PcZn(Ace) <sub>2</sub> <sup>b</sup>	4 <sup>1</sup> E <sub>u</sub>	89% (32b <sub>1u</sub> → 30e <sub>g</sub> )	2.99		0.0021	$\pi \rightarrow \pi^*$
	5 <sup>1</sup> E <sub>u</sub>	56% (23a <sub>1u</sub> → 30e <sub>g</sub> ); 27% (31b <sub>1u</sub> → 30e <sub>g</sub> )	3.13		0.1025	B <sub>2</sub>
	6 <sup>1</sup> E <sub>u</sub>	51% (25a <sub>1u</sub> → 31e <sub>g</sub> ); 38% (31b <sub>1u</sub> → 30e <sub>g</sub> )	3.30		0.7449	B <sub>1</sub>
	7 <sup>1</sup> E <sub>u</sub>	41% (25a <sub>1u</sub> → 31e <sub>g</sub> ); 28% (31b <sub>1u</sub> → 30e <sub>g</sub> )	3.37		0.3465	$\pi \rightarrow \pi^*$
	8 <sup>1</sup> E <sub>u</sub>	75% (28e <sub>g</sub> → 26a <sub>1u</sub> )	3.59		0.1101	

<sup>a</sup> Contribution of less than 10% is not listed. <sup>b</sup> For convenience, we also use the label e<sub>g</sub> (of D<sub>4h</sub> symmetry) for the lower symmetry system. <sup>c</sup> References 13 and 33.

**Table 7.** Calculated Properties of H<sub>16</sub>PcZn, F<sub>16</sub>PcZn, and F<sub>48</sub>PcZn with Axial Ligands of Cl and of Cl<sup>−</sup>

	H <sub>16</sub> PcZn(Cl)	F <sub>16</sub> PcZn(Cl)	F <sub>48</sub> PcZn(Cl)	H <sub>16</sub> PcZn(Cl <sup>−</sup> )	F <sub>16</sub> PcZn(Cl <sup>−</sup> )	F <sub>48</sub> PcZn(Cl <sup>−</sup> )
$R_{\text{Ct}\cdots\text{N(Pc)}}^a$ (Å)	1.976	1.971	1.978	1.987	1.986	1.986
$R_{\text{Ct}\cdots\text{Zn}}^b$ (Å)	0.761	0.774	0.736	0.710	0.695	0.695
$R_{\text{Zn}-\text{Cl}}$ (Å)	2.232	2.214	2.200	2.269	2.250	2.249
$Q_{\text{Zn}}$ (e)	0.54	0.54	0.52	0.57	0.56	0.55
$Q_{\text{Cl}}$ (e)	−0.45	−0.41	−0.44	−0.51	−0.45	−0.44
$E_{\text{bind}}(\text{PcZn}-\text{Cl})^c$ (eV)	1.77	1.62	1.32	1.81	2.96	3.79

<sup>a</sup> Ct denotes the center of the Pc ring and  $R_{\text{Ct}\cdots\text{N(Pc)}}$  denotes the distance between Ct and N(Pc). <sup>b</sup>  $R_{\text{Ct}\cdots\text{Zn}}$  denotes displacement of the Zn atom out of the Pc plane. <sup>c</sup> Binding energy between PcZn and Cl or Cl<sup>−</sup>.

core size,  $R_{\text{Ct}\cdots\text{Zn}}$  represents the displacement of the Zn atom out of the N<sub>4</sub>-plane toward the L ligand, and  $R_{\text{Zn}-\text{Cl}}$  refers to the Zn–Cl bond length.

The calculated H<sub>16</sub>PcZn–Cl binding energy  $E_{\text{bind}}$  is large, 1.77 eV, indicating a high affinity of H<sub>16</sub>ZnPc for Cl. As an electron-withdrawing ligand, the binding of neutral Cl to F<sub>48</sub>–

PcZn becomes significantly weaker. A large Zn out-of-plane displacement is found when Cl is coordinated to the system. In this case, the Ct⋯N distance becomes somewhat shorter than  $R_{\text{Zn}-\text{N}}$  of unligated PcZn.

For the chloride ion, there is a clear trend of increasing binding energy in the order H<sub>16</sub>PcZn–Cl<sup>−</sup> < F<sub>16</sub>PcZn–Cl<sup>−</sup>



$< F_{48}PcZn-Cl^-$ , where the  $F_{48}PcZn-Cl^-$  binding energy is much larger than that of  $H_{16}PcZn-Cl^-$ . The values of  $E_{bind}$  correlate well with the relative stabilities of these  $PcZn(Cl^-)$  complexes.<sup>13</sup>

#### 4. Conclusions

The following main conclusions may be drawn from the calculated results.

(1) While the strongly electron-withdrawing  $-C_2F_5$  groups on the Pc ring are able to change the ground state of  $PcFe$ , they do not influence the ground state of  $PcCo$ .

(2) Corresponding to the downshifts of the valence MOs caused by the peripheral substituents, the first ionization potential (IP) is increased by  $\sim 1.3$  eV from  $H_{16}PcM$  to  $F_{16}PcM$  and by another  $\sim 1.1$  eV from  $F_{16}PcM$  to  $F_{48}PcM$ . A similar increase in the EA is also found on going from  $H_{16}PcM$  to  $F_{48}PcM$ . These results account for the fact that the  $F_{64}PcM$  compounds are difficult to oxidize but easy to reduce.<sup>13</sup> The change from  $H_{16}PcZn$  to  $F_{48}PcZn$  does not produce an obvious change in the LUMO–HOMO energy gap, but large changes in the relative MO energies are found between the  $b_{1g}$  (HOMO-1) and  $b_{2g}$  (HOMO-2) orbitals.

(3) The axial acetone (Ace) ligands in  $PcM(Ace)_2$  shift the MO energy levels upward. There is a decrease of  $\sim 0.8$  eV in IP from  $PcFe$  to  $PcFe(Ace)_2$ , and a decrease of  $\sim 0.4$  eV when the metal M is Co and Zn. Therefore, axial coordination by (strong)  $\sigma$ -donors renders easier the oxidation of the compound, as observed experimentally.<sup>13</sup>

(4) The calculated  $H_{16}PcM-(Ace)_2$  binding energy ( $E_{bind}$ ) is small ( $< 0.2$  eV), particularly for  $M = Zn$ , indicating a low affinity of the parent metal phthalocyanine for acetone. The  $E_{bind}$  is increased to nearly 1 eV when the  $-C_2F_5$  groups are introduced at the periphery of the Pc ring. This rise accounts for the experimental result<sup>13</sup> that there is formation of a  $F_{64}PcZn(Ace)_2$  compound, but no  $H_{16}PcZn(Ace)_2$  was detected. On the other hand, the calculated  $PcZn-Cl^-$  binding energies can also account for the relative stabilities of those  $PcZn(Cl^-)$  compounds.

(5) The TDDFT calculated excitation energies for the  $PcZn$  complexes are in quantitative agreement with available experimental data. The substitution of 16 H-atoms in  $H_{16}PcZn$  by F-atoms (i.e. from  $H_{16}PcZn$  to  $F_{16}PcZn$ ) shifts the Q and B bands to the red. A further red shift in the B<sub>1</sub> band occurs with partial replacement of eight F-atoms in  $F_{16}PcZn$  by  $-C_2F_5$  groups, but there is no shift in the Q-band from  $F_{16}PcZn$  to  $F_{48}PcZn$ . On the other hand, the calculated electronic spectrum of  $F_{48}PcZn$  is nearly the same as that of  $F_{48}PcZn(Ace)_2$ , supporting an assumption that the results (spectra) for the fluorinated phthalocyanines are virtually independent of the presence or absence of axial ligands.<sup>13</sup>

(6) The calculated redox properties of gas-phase cobalt phthalocyanines may be different from those measured in solution. These differences are likely due to association of specific solvent molecules with the complex.

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