# Spectroscopic Probing of Adsorption of Uranyl to Uranyl-Imprinted Silica Sol—Gel Glass via Steady-State and Time-Resolved Fluorescence Measurement

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Steady-state and time-resolved fluorescence spectroscopic methods have been demonstrated to be very informative tools to analyze the microstructure variation induced by the imprinting technique. The structural information obtained provides the rationale for the enhanced uranyl-adsorption capacity and selectivity by the uranyl-imprinted sol—gel glasses.

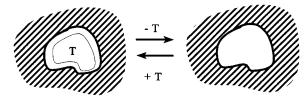
#### Introduction

Fluorescence techniques have been widely used to probe surface characteristics of inorganic solids and organic polymers. Pluorescence signals from probe ions adsorbed or covalently bound to solid matrixes can reveal information about their local environment through changes in their photophysical properties such as fluorescence lifetime and emission spectra. For example, the emission-lifetime distribution of a probe ion or molecule adsorbed on an inorganic sorbent has been directly correlated to the heterogeneity of the material. A wide variety of solid matrixes have been investigated via these techniques, 1,2,7 including silica, 1a,b,c,2 reversed-phase silica, 1h zeolites, titania, clay, and semiconductors.

The uranyl ion exhibits strong fluorescence signals with long emission lifetimes which are sensitive to both the local environment of the ion and its mobility.<sup>3</sup> For these reasons, the uranyl ion is an ideal, in situ photophysics probe for the investigation of the binding of this ion to a surface. Recent studies in our laboratories with sol-gel matrixes imprinted with uranyl ions (UO22+) have shown enhanced selectivity and capacity toward uranyl separations.4 In the present work, fluorescence techniques are used to investigate the variations in the surface microstructure of sol-gel matrixes induced by imprinting techniques. Ultimately we seek to relate models for the surface microstructure of imprinted matrixes to their macroscopic properties such as selective ion binding when competitive ions are present. Understanding these relationships is an essential engineering requirement in the preparation of new materials with tailored properties.

Imprinting methods based on the template approach have been used in cross-linked polymers<sup>5-7</sup> as well as in silica gels<sup>9</sup> to prepare polymeric supports possessing solid-state organized structures. Usually, imprinting processes (Scheme 1) involve three steps: (a) selection of target molecules as templates; (b) incorporation of the templates into rigid solid networks; (c) removal of the templates leaving cavities that later "recognize" or selectively bind the template or target molecule. The imprint cavity is produced by the template which is present during the formation of the solid matrix. After the templating species is removed from the polymer, a predetermined arrangement of ligands and a tailored binding pocket remains. Imprinted materials thus prepared have been shown to combine the binding ability of specifically chosen functional groups or ligands for

**SCHEME 1: Imprinting Process via a Template (T)** 



target substrates, with shape and size characteristics of cavities imprinted into a rigid polymer matrix by a template. Such materials are observed to selectively bind target substrates with high affinity and thus have many applications in the areas of separation science and catalysis. Imprinted organic polymers have been used to resolve racemates<sup>6</sup> and separate mixtures of metal cations.<sup>7</sup>

In imprinting technologies, it is of critical importance to be able to define in detail the properties of the imprinted cavity. These properties include defining the size and shape of the cavity as well as the type and number and functional groups involved in binding target ions. Furthermore, the distribution of templated cavities in the solid matrix must also be known. Previous studies of imprinted materials have relied on indirect measures to characterize the properties of imprinted cavities. No simple spectroscopic techniques have been developed which directly probe the surface microstructures induced by imprinting.

The solid-state silica network used to create imprints of  $\rm UO_2^{2+}$  was prepared by the sol-gel process. The sol-gel process is a technique that can be used to prepare transparent oxide glasses by hydrolysis and condensation of tetraalkyl orthosilicates. Little or no heating is required, and consequently the gel can be doped or templated with molecules whose poor thermal stabilities preclude their incorporation in traditional high-temperature silica glasses. Such molecules become entrapped in the growing covalent silica network rather than being chemically bound to the inorganic matrix.

## **Experimental Section**

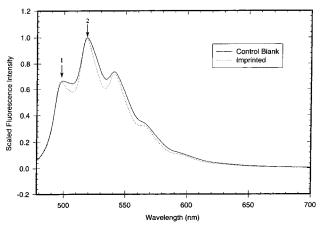
The instrumentation and procedures used for the measurement of both steady-state luminescence spectra and lifetimes have been described previously. 11,12a The 337 nm line of a pulsed nitrogen laser (pulse duration = 5 ns) was used as the excitation source in these studies. Luminescence decay times were analyzed using both the Kohlraush—Williams—Watts (KWW) model and direct numerical inverse Laplace transform. The CurveFit program in SigmaPlot (Jandel Industries Corp.) was employed to fit experimental decay curves to the KWW model,

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**Figure 1.** Comparison of experimental uranyl fluorescence spectra of reloaded control blank (solid line); and reloaded imprinted sample (doted line).

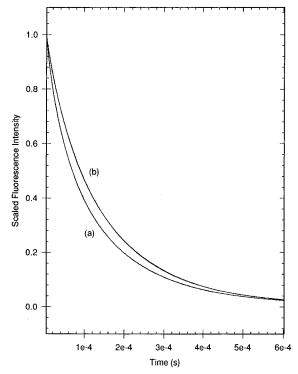
whereas the continuous model analysis program, CONTIN, <sup>12</sup> was used to conduct the numerical inversion of Fredholm and Volterra integral equations of the first kind (inverse Laplace transform). Solutions to the inverse Laplace transform were constrained to positive distribution coefficients. The line widths of emission spectra were obtained via CurveFit program in Grams/386 (Galactic Industries Corp.).

Literature procedures<sup>13</sup> were used to prepare uranyl-doped sol-gel glasses. In a typical run, 1 mL of an aqueous UO<sub>2</sub><sup>2+</sup> solution (0.06 M UO<sub>2</sub>(NO<sub>3</sub>)<sub>2</sub>•6H<sub>2</sub>O dissolved in 2 N HNO<sub>3</sub>) was mixed with 1 mL of tetramethyl orthosilicate (Aldrich Chemical Co., 99%), and 1 mL of methanol (Baker Chemical Co., HPLC grade). The final mixture was cured at the ambient temperature over a period of 4 weeks. Control blank samples were prepared in a similar manner except that no uranyl was added to the nitric acid solution.

It is well-known that the adsorption capability of silica-based sorbents toward UO<sub>2</sub><sup>2+</sup> is very small below the pH of 2.<sup>14</sup> This characteristic in uranyl binding to silicates and highly porous microstructure of sol—gel glasses<sup>9</sup> prompted us to adopt acid washing as a method to efficiently remove the uranyl template ions from both surfaces and bulks of the glasses. Uranyl-free imprinted glasses were obtained by soaking the UO<sub>2</sub><sup>2+</sup>-doped glasses in nitric acid (1 N) overnight. Control blank glasses were also washed with the identical nitric acid over the same period of time to ensure that both types of glasses had identical treatments of their surfaces. No UO<sub>2</sub><sup>2+</sup> fluorescence signals could be observed in any of the washed samples.

### **Results and Discussion**

# Steady-State Fluorescence Studies of Uranyl Binding. Both imprinted and control blank glasses were immersed in identical aqueous solutions containing dissolved uranyl ion ( $[UO_2^{2+}]$ 10<sup>-4</sup> M, pH = 5.2) for reloading of uranyl ions. This reloading uranyl solution was prepared by titrating 10<sup>-4</sup> M uranyl nitrate solution with tetrabutylammonium hydroxide (Fluka, 55-60% in water) to pH 5.2. Following overnight immersion, the samples were subsequently washed with hexane and dried in the air. This adsorption treatment gives rise to the concomitant appearance of fluorescence spectra attributable to uranyl complex species in both glasses (Figure 1a,b). Although the general spectral patterns and peak positions of spectra a and b are very similar, the line widths of the vibronic emission components are clearly different. Assuming the line shape of each vibronic emission line to be Gaussian, we can determine the corresponding line width through nonlinear least-squares fitting (CurveFit program in Grams/386). The line widths of



**Figure 2.** Comparison of experimental uranyl fluorescence decays (solid line) of (a) reloaded control blank and (b) reloaded imprinted sample with those obtained by fitting to KWW model (dotted line) and inverse Laplace transform (dashed line). All experimental and fitting decay curves are superimposed well with each other so that they cannot be easily distinguished from the figure.

TABLE 1: Line Widths of Two Main Vibronic Emission Peaks Obtained via Nonlinear Least-Squares Fitting

	sample	line width (peak 1) (cm <sup>-1</sup> )	line width (peak 2) (cm <sup>-1</sup> )	
	imprinted	539	467	
	control	566	475	

two main peaks obtained via nonlinear least-squares fitting (marked 1, 2 in Figure 1) are listed in Table 1. As seen qualitatively from Figure 1 and quantitatively from Table 1, the line widths of the vibronic emission lines corresponding to the imprinted sample are systematically smaller than those of the control blank. Wirth and Hahn<sup>16</sup> have pointed out before that the spectroscopic line width of an optical probe can be correlated with the shape selectivity of separation sorbents. This assertion was originally based on an intuitive argument that solutes having shapes incompatible with those of the solvation site would result in inhomogeneous broadening of the spectral line width through various interactions. Later, Wirth<sup>15b</sup> developed a physical model for this phenomenon by showing that structural order is a necessary condition for shape selectivity. Following the argument of Wirth, 15 we can attribute the narrower line widths of the emission lines for the reloaded imprinted sample to the shape-selective binding of uranyl ion in cavities produced by imprinting. In contrast, the adsorption of uranyl ion on the control blank would entail no shape selectivity, accordingly having broader line widths. The adsorption site of the imprinted material is more ordered and homogeneous than that of the control blank.

Fluorescence Lifetime Studies of Uranyl Binding. Fluorescence decay profiles of imprinted and control glass samples are shown in Figure 2. The two decay curves were measured at the maximum emission wavelength (518 nm). To avoid problems caused by noise associated with photon-counting statistics, <sup>1a,11</sup> 20 000 decay traces were acquired for each sample

TABLE 2: Results from Nonlinear Least-Squares Fitting of **Experimental Decays with KWW Model** 

	control b	control blank		imprinted sample	
	parameter value	standard deviation	parameter value	standard deviation	
$egin{array}{c}  au_0 \ eta \ \langle  au  angle \end{array}$	9.700 <i>E</i> -5 (s) 0.740 1.17 <i>E</i> -4 (s)	6.8 <i>E</i> -8 3.5 <i>E</i> -4	1.282 <i>E</i> -8 0.852 1.39 <i>E</i> -4 (s)	1.2 <i>E</i> -7 6.7 <i>E</i> -4	

and summed to increase signal-to-noise ratios (SN). Both decay profiles are nonexponential and similar to those measured for other amorphous systems such as pyrene bound to amorphous silicon surfaces.<sup>1,2</sup>

For many amorphous materials,  $^{1e,1f,1g,2b}$  decay profiles ( $I_t(t)$ ) can be approximated well by the two-parameter Kohlraush-Williams-Watts (KWW) model (eq 1). The parameter  $\beta$  (0 <

$$I_{f}(t)/I_{0} = \exp[-(t/\tau_{0})^{\beta}]$$
 (1)

 $\beta$  < 1) has been related to a distribution of exponential decay constants that are serially linked and is sometimes referred to as a "width parameter". Small  $\beta$  values indicate a broad distribution of decay rates. The parameter  $\tau_0$  is the lifetime at the maximum amplitude of the distribution. The average lifetime can be calculated from the first moment of eq 1 as given in eq 2, where  $\Gamma$  represents the gamma function. The computer

$$\langle \tau \rangle = (\tau_0/\beta)\Gamma(1/\beta)$$
 (2)

program used to evaluate the gamma function has been described in the literature.16 Comparison of the experimental decays and those obtained from the nonlinear least-squares fitting according to the KWW model is shown in Figure 2. Table 2 summarizes the fit parameters obtained from the KWW model along with their standard deviations.

According to the KWW model, there are two important differences in the distribution of binding sites between the imprinted and control glasses. First, both  $\langle \tau \rangle$  and  $\tau_0$  for the imprinted sample are greater than those of the control blank. Second, the width of the lifetime distribution for the control blank glass is larger than that of the imprinted glasses. This observation can be explained within the framework of a variation in the microstructure of the glass induced by the imprinting technique. As pointed out previously, lifetime has a direct correlation with the local environment of the probe ion. Two kinds of local environments for adsorbed uranyl ions can be anticipated for reloaded uranyl ions in the sol-gel glasses. The first one is considered to be a hydrated uranyl species associated weakly with sol-gel silica surfaces, whereas the second is uranyl ion strongly bound to a cavity in the silica matrix. The former local environment is closer to uranyl ions in aqueous solutions, while the latter is like uranyl ions solvated by silica (e.g., uranyl ion doped in silica glasses). The fluorescence decay of uranyl in silica and boron-silica glasses has been extensively studied, and the experimental lifetime (10<sup>-4</sup> s) is very long. <sup>17</sup> However, the lifetime of uranyl ion in aqueous solutions is relatively short, in the neighborhood of  $10^{-6}$  s.<sup>18</sup> Accordingly, the uranyl lifetime of the reloaded imprinted system is longer than that of the reloaded control blank, since imprinted glass samples have a greater number of cavity-bound uranyl ions present in comparison to the control samples. The width of the lifetime distribution can be correlated to the distribution of adsorption sites. Most of the adsorption sites in the imprinted glass are considered to be created by the template. Therefore, the distribution is expected to be narrower than that of the control blank where no microorganization exists. In addition

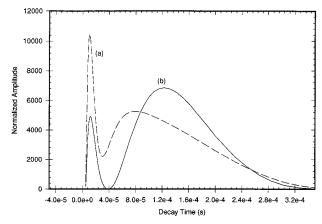


Figure 3. Lifetime distributions of (a) control blank and (b) imprinted sample obtained via inverse Laplace transform.

to the fluorescence decays monitored at the maximum-emission wavelength (518 nm), decay curves at several other wavelengths have been measured. The average lifetimes were found to vary slightly with monitoring wavelengths, being consistent with the heterogeneity of the sorbent materials. At all monitoring wavelengths, the average lifetimes of uranyl fluorescence in the blank sample were shorter than those of the imprinted sample. Similar decay kinetics were obtained for four other sets of imprinted and blank samples prepared using the same experimental conditions.

Further support for the above interpretation comes from the CONTIN analysis of experimental data. Unlike the KWW model discussed above, no a priori assumptions concerning the lifetime distribution are made in the CONTIN inverse Laplace transfer of experimental decay curves. Figure 3 shows the bimodal distribution recovered using the CONTIN program. The comparison of the decays reproduced using the Figure 3 lifetime distribution with experimental decays is given in Figure 2. The sharp peak with short lifetime is considered to be due to the hydrated uranyl species weakly associated with silica surfaces, whereas the broad long-lifetime peak can be attributed to the cavity-bound uranyl ions. As seen in Figure 3, the short-lifetime component contributes more in the control blank than in the imprinted glass. This is consistent with the fact that more reloaded uranyl ion is weakly bonded in the former sample. Furthermore, the width of the longer lifetime component is broader for the control blank than for the imprinted sample, consistent with a narrower distribution of cavities in the imprinted sample.

In conclusion, the fluorescence spectroscopy has been used to probe imprinted materials. This study provides microscopic information concerning the imprinted materials. Such information corroborates well with the fact that imprinted materials possess enhanced selectivity and capacity toward target ions. Experiments are currently underway to investigate low-temperature spectra of these materials.

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