Solvothermal Preparation and Spectroscopic Characterization of Copper Indium Diselenide Nanorods

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Single-crystalline chalcopyrite CuInSe₂ nanorods (CuInSe₂NRs) of 50–100 nm in diameter and up to a few micrometers in length have been synthesized solvothermally. High-resolution transmission electron microscopic images of the CuInSe₂NRs reveal the *d*-spacing of 0.335 nm for the (112) crystalline planes and a growth direction along [331]. The near-infrared absorption spectrum of the chalcopyrite CuInSe₂NRs shows a peak maximum at 1162 nm and an onset at 1262 nm, indicating no apparent blue-shift compared with those of Cu-rich CuInSe₂ thin films. An intense peak at 175.1 cm⁻¹ in the room-temperature Raman scattering spectrum of CuInSe₂NRs corresponds to the A₁ phonon mode of tetragonal CuInSe₂ chalcopyrite. The narrower full width at half-maximum (fwhm) of 9.5 cm⁻¹ for CuInSe₂NRs, in comparison with fwhm \sim 12 cm⁻¹ for CuInSe₂ films, indicates a uniform size distribution and single crystallization in the nanorods. Analysis of the photoluminescence from the single-crystalline CuInSe₂NRs measured at 10 K has categorized the emission into seven groups of transitions as characterized by free excitons, bound excitons, conduction band to acceptor levels, and bound excitons at different defects.

1. Introduction

Because of the exhausting of fossil fuels, such as petroleum and coal, the development of new energy resources and their related environmental concerns has become an extremely important issue in recent years. Among many other potential energy sources, the renewable ones, such as photovoltaic electric generators, have drawn much attention. In particular, the use of the sun's radiant energy seems to be a sensible long-term goal. To make more efficient use of solar energy, a variety of solar cells have been developed for energy preservation.

The first solar cell was developed with the semiconductor material of single-crystalline silicon at Bell Laboratories by Chapin et al. more than a half century ago. The cell was applied as a power supply in space vehicles with a rough estimate for the energy conversion efficiency of 6%. Because the cost of silicon was a significant proportion for the solar cell, a great deal of effort has been made to search for stable, highly conversion-efficient, and cost-effective substitutes. Ever since, laboratory demonstration cells of reduced minority-carrier recombination rate and trapping light in active layers have made dramatic increases in efficiency. For instance, fundamental GaAs solar cells and the double-junction laboratory cells of single-crystalline AlGaAs/GaAs³ and GaInP₂/GaAs⁴ have reached 24, 27.6, and 29.5% conversion efficiencies, respectively.

As photovoltaic research extended beyond the simple elemental Si of Group IV to Group III-V compounds, such as GaAs, and then to Group II-VI compounds (e.g., CdTe), it has moved with one more step of complexity to I-III-VI₂ compounds, such as copper indium diselenide (CuInSe₂). This CuInSe₂ could be simplistically thought of as the result of substituting both Group I (Cu) and Group III (In) elements to give an average similar to the Group II element in II-VI

crystals. Such a ternary compound represents the first step to more complicated solid solutions of more than three different atoms, with which both desired band gap and proper heterojunction lattice match could, in principle, be obtained to provide maximum efficiency. Research on quaternary and pentenary systems has also been under way; however, to date, these much more complex materials have not met the stringent demands of reproducibility and operating lifetime typically required for a solar-cell material.⁵

The ternary I-III-VI2 semiconductor of CuInSe2 has advantageous applications in thin-film solar cells because of its large absorption coefficient, low band gap (\sim 1.05 eV), and good radiation stability.6-8 Films of CuInSe₂ have been fabricated with various methods, such as molecular beam epitaxy, 9 liquidphase epitaxy, 10 and halogen vapor phase epitaxy. 11 Polycrystalline CuInSe2 films have also been grown by metalorganic chemical vapor deposition using separate Cu and In precursors. 12 To the best of our knowledge about the development of CuInSe₂ nanostructures, there are only very few studies reported so far that include CuInSe₂ nanoparticles¹³ and CuInSe₂ nanowhiskers.¹⁴ In this study, we report high-purity CuInSe₂ nanorods (CuInSe₂NRs) synthesized by a solvothermal method. X-ray diffraction (XRD) and electron microscopic examinations of CuInSe2NRs have shown a single-crystalline structure for the chalcopyrite crystals. Spectroscopic characterizations for the CuInSe₂NRs by near-infrared (NIR) absorption, micro-Raman scattering, and micro-photoluminescence (PL) spectroscopies will also be presented.

2. Experimental Section

The synthesis of CuInSe₂NRs was conducted in a solvothermal reaction similar to the method by Li et al. ¹⁴ In a typical experimental procedure, the product of CuInSe₂ could be prepared from a stoichiometric mixture of Se powder (Strem,

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99.99%), In₂Se₃ powder (Aldrich, 99%), and CuCl₂ anhydrous powder (Acros). The powder precursors were dissolved in ethylenediamine (Acros), which also acted as a reactant in this experiment, as will be discussed later. In the beginning of the solvothermal reaction, Se powder (0.07896 g) was added to 50 mL anhydrous ethylenediamine. The temperature was then increased to ~120 °C by a heating plate. The solvothermal reaction was followed by the addition of both In₂Se₃ (0.46652 g) and anhydrous CuCl₂ (0.26890 g) powders. The reaction temperature was kept at \sim 120 °C throughout the reflux experiment. After the reaction, precipitates were filtered and washed several times with distilled water and ethanol to remove byproducts. The reaction product was dried at 60 °C for 6 h and used afterward for powder XRD analysis, electron microscopic measurements, and NIR, Raman, and PL spectroscopic characterizations.

The experimental details for the electron microscopic and spectroscopic characterizations of the as-synthesized CuInSe₂NRs are similar to those of several previous studies in this laboratory, including Si nanowires, 15 CdSe nanobelts and nanosheets, 16 Zn_xCd_{1-x}Se alloy nanowires, ¹⁷ and SiO₂ nanoparticles. ¹⁸ Briefly, scanning electron microscopic (SEM) morphologies of the CuInSe₂NRs were taken in a LEO 1530 field emission gun SEM and energy dispersive spectrometer (FEG-SEM+EDS). Transmission electron microscopic (TEM) images were obtained in a JEOL JEM-1200EX II TEM at 80 kV and a Hitachi H-7100 TEM at 75 kV. High-resolution transmission electron microscopic (HRTEM) measurements were conducted in a JEOL JEM-2010 analytical TEM at 200 kV and a JEOL JEM-4000EX HRTEM at 400 kV.

Spectroscopic examination of NIR absorption was performed in a Jasco V-570 UV/Vis/NIR spectrophotometer with a scanning range from 1100 to 1400 nm. In the NIR spectroscopic investigation, the sample of CuInSe2NRs was suspended in ethylenediamine where ethylenediamine served as a background for reference. Micro-Raman spectra of CuInSe2NRs were recorded with the signal-accumulation time of 180 s in a confocal microscope (Jobin Yvon, MFO) combined with a system of a monochromator (Jobin Yvon, Triax 550) and a liquid-nitrogen-cooled charge-coupled device camera (Jobin Yvon, 1024×256 pixels). The 488-nm line of an Ar⁺-Kr⁺ laser (Coherent, Innova 70C) was used for the Raman excitation with a power level of 50 mW. Crystal structures of the CuInSe2NRs were analyzed using an X-ray diffractometer (Philips, PNAalytical X'Pert PRO) of Cu Ka radiation with a scanning range of 2θ from 20 to 70°. PL from the CuInSe₂NRs, excited at 760 nm by an AlGaAs solid-state laser of 25 mW, was collected by a double-grating monochromator (Jobin Yvon, Spex 1403) together with a liquid-nitrogen-cooled germanium detector (North Coast, EO807). The PL measurement was also conducted in the temperature range of 8-300 K with a closedcycle helium-cooled refrigerator (Edinburgh Instruments, STL20).

3. Results and Discussion

XRD of CuInSe₂NRs. To determine the crystal structures of the as-synthesized CuInSe₂NRs, the XRD pattern was examined as shown in Figure 1a. Seven prominent XRD peaks were identified as being due to the CuInSe₂ crystal planes of (112), (211), (220), (204), (301), (312), and (400), which give a chalcopyrite tetragonal structure with lattice constants of a =0.5949 nm and c = 1.1443 nm, consistent with the standard values of the reported CuInSe₂ data (JCPDS No. 23-0209). Six XRD signals of much weaker intensities, which could not be assigned to chalcopyrite CuInSe2, were attributed to the Se-

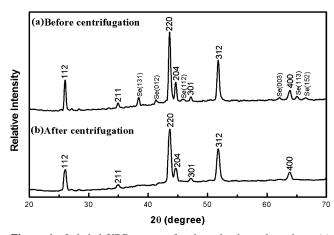


Figure 1. Labeled XRD spectra for the solvothermal products (a) before and (b) after high-speed centrifugation. Before centrifugation, both CuInSe₂ chalcopyrite and Se-related nanocrystals are responsible for the observed XRD signals. After centrifugation, the Se nanocrystals have been removed effectively.

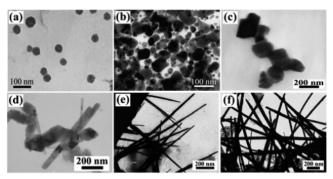


Figure 2. TEM morphologies for the CuInSe₂ nanocrystals in a series of solvothermal reactions with reflux times of (a) 12, (b) 24, (c) 28, (d) 32, (e) 36, and (f) 48 h. The morphologies of the as-synthesized CuInSe₂ products evolving from nanoparticles to nanorods during the experimental course are manifested in panels a-f.

related structures of (131), (012), (112), (003), (113), and (152). This XRD result indicates a mixture of CuInSe₂ and Se crystals in the product. To extract the CuInSe2NRs out of the mixture, we carried out a physical separation in a high-speed centrifuge (typically $\sim 10~000$ rpm). Figure 1b shows the XRD pattern of tetragonal CuInSe2 crystallites after the high-speed centrifugation, indicating an effective separation.

Electron Microscopic Characterizations of CuInSe2NRs. Figure 2a-f presents the TEM images of six samples produced in the solvothermal reaction after a series of reflux times of 12, 24, 28, 32, 36, and 48 h, respectively. With a relatively shorter reaction time (12 h, Figure 2a), the product is mainly spherical nanoparticles. In a longer experimental course, these nanoparticles were recrystallized to develop new morphologies, as shown in Figure 2b. These irregularly shaped nanoparticles with a diameter range of 20-80 nm played an important role as initiators (or "seeds") for the growth of one-dimensional CuInSe₂NRs. Figure 2c shows the start of elongation of the nanoparticles to develop to nanorods at the reflux time of 28 h. In the subsequent reaction period of 32 h of reflux (Figure 2d), some of the nanorods already possessed quite large aspect ratios and had grown to hundreds of nanometers in length. When the reaction time reached at least 36 h, CuInSe₂NRs were found in the ethylenediamine solution with a high production yield. Shown in Figure 2e are the CuInSe₂NRs 50-100 nm in diameter and up to a few micrometers in length. A TEM image of the as-synthesized CuInSe2NRs after 48 h of reaction is presented

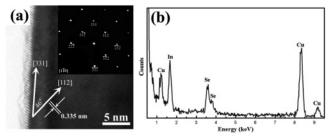


Figure 3. (a) HRTEM image of the CuInSe₂NRs. The *d*-spacing of 0.335 nm in the (112) planes and a growth direction along [331] have been identified on the basis of the SAED pattern (inset) with the [1 $\overline{10}$] zone axis. (b) EDS analysis for the CuInSe₂NRs indicates the molar ratio of Cu/In/Se = 0.98:0.87:1.95.

in Figure 2f, in which the morphology of CuInSe₂NRs is almost the same as that in Figure 2e, except more nanorods having been produced.

An HRTEM image of the CuInSe₂NRs (Figure 3a) reveals their single-crystalline structure of 0.335 nm d-spacing due to the (112) crystal planes. The growth direction of CuInSe₂NRs along [331] has been analyzed from the selected-area electron diffraction (SAED) pattern, as shown in the inset of Figure 3a, with a zone axis of [110]. The angle between the rod-axial growth direction of [331] and the <112> lattice growth direction is determined to be \sim 46°. An EDS analysis (Figure 3b) for the chemical compositions of CuInSe₂NRs further confirms the stoichiometric ratio of Cu/In/Se \approx 1:1:2 (strictly, 0.98:0.87:1.95).

Growth Mechanism for CuInSe2NRs. At room temperature, elemental Se is soluble in ethylenediamine. This dissolution is accompanied not only by a color change of the solution from opaque to brown, but also the instant disappearance of the Se powder. The increasing reactivity of Se is enhanced by its dispersion into ethylenediamine. In the solvothermal reaction, dissolved elemental Se was reduced to negatively charged Se²⁻ by amines in an organic nucleophilic attack, similar to how elemental S can be activated by amines or hydroxides.¹⁹ When temperature increased, the In₂Se₃ reactant, which has low solubility in ethylenediamine, reacted with the Se²⁻ ions to form InSe₂⁻ (eq 1 in the following summary reactions). According to previous studies on the fabrication of nanomaterials, ethylenediamine can facilitate the growth of one-dimensional nanostructures, 20,21 in which the strong chelation ability of alkylamine is essential for the formation of chelated compounds as a template to build up one-dimensional structures. In this work, Cu²⁺ from CuCl₂ was first reduced to Cu⁺ as discussed previously by Li et al.,14 which was then chelated by bidentate ethylenediamine (denoted by "en" in the following eq 2) to form the complex of [Cu(en)₂]⁺. According to the proposed structure of [Cu(en)₂]⁺ by Li et al., ¹⁴ the ethylenediamine binds strongly with Cu⁺ to form a two five-membered-ring chelated structure in which the Cu⁺ bridges the amino groups of the two ethylenediamine moieties. In a recent theoretical study by Alcamí et al.,²² the five-membered-ring structure, composed of Cu⁺ and ethylenediamine, has further been calculated to be most stable in energy among possible isomeric structures. As mentioned by Li et al., 14 InSe₂⁻ in ethylenediamine was self-organized with the [Cu(en)₂]⁺ complex to start nucleation and crystalline growth to generate the one-dimensional CuInSe₂ nanostructures. The [Cu(en)₂]⁺ complex was also regarded as serving as a molecular template to guide the growth of CuInSe2NRs (eq 3), similar to that reported by Morris et al.²³ on the study of a template growth mechanism. The proposed mechanism involved in this solvothermal reaction for the growth of CuInSe2NRs can be sum-

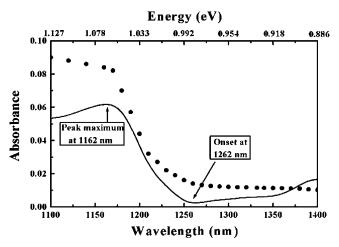


Figure 4. NIR absorption spectrum of the as-synthesized CuInSe₂NRs with a peak maximum at \sim 1162 nm and an onset at \sim 1262 nm. The dotted curve represents the absorption coefficient of Cu-rich CuInSe₂ proceeded films taken from refs 25 and 26.

marized in the following eqs 1-3:

$$In2Se3 + Se2- \rightarrow 2InSe2-$$
 (1)

$$Cu^{+} + 2en \rightarrow [Cu(en)_{2}]^{+}$$
 (2)

$$InSe_2^- + [Cu(en)_2]^+ \rightarrow CuInSe_2 + 2en$$
 (3)

NIR Absorption of CuInSe2NRs. An NIR absorption spectrum of the CuInSe2NRs as illustrated in Figure 4 was recorded from 1100 to 1400 nm, with CuInSe₂NRs being dispersed in ethylenediamine solution. Shown in Figure 4 are a peak maximum at \sim 1162 nm and an onset at \sim 1262 nm, whose separation can be regarded as a guideline to evaluate the size distribution of the as-prepared semiconductor nanocrystallites.²⁴ In comparison with the absorption spectrum, the dotted curve in Figure 4 of Cu-rich CuInSe₂ proceeded films (sample C in ref 25 synthesized by Mooney et al.) observed by Kavcar, 26 there is no apparent blue-shift (i.e., quantum confinement effect) in the absorption spectrum of CuInSe₂NRs. It is noted that the reported Bohr radius of 1.2 nm²⁷ or 2.02 nm¹³ for the CuInSe₂ compound is much smaller than the radii (typically 25-50 nm) of the as-synthesized chalcopyrite CuInSe2NRs, and a quantum confinement effect is unlikely to be observed in the present experiment.

Raman Scattering of CuInSe₂NRs. Figure 5 shows the room-temperature Raman scattering spectrum of CuInSe₂NRs. The very strong signal at 175.1 cm⁻¹ corresponds to the A_1 optical phonon mode in tetragonal CuInSe₂ crystal, consistent with the vibrational observation of CuInSe₂ films where the A_1 mode is much stronger than any other phonon mode. This A_1 mode relates to the vibration of two pairs of anions, one in the direction of the a-axis and the other in the direction of the b-axis.²⁸ An analysis of the line widths further indicates that the full width at half-maximum (fwhm) of \sim 9.3 cm⁻¹ for CuInSe₂NRs is narrower than the typical fwhm of \sim 12–13 cm⁻¹ for CuInSe₂ thin films,²⁹ which could be attributed to more uniformity in size-distribution and better crystallization in the single-crystalline CuInSe₂NRs.

Photoluminescence of CuInSe₂NRs. The PL spectrum (Figure 6) of the as-synthesized single-crystal CuInSe₂NRs, freely dispersed on a copper substrate, was obtained at 10 K with an excitation at 760 nm by a 25 mW solid-state laser. The well-resolved spectral features in Figure 6 are categorized into

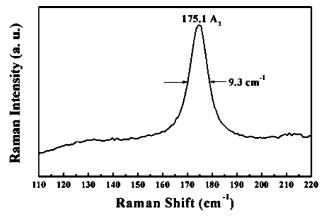


Figure 5. Raman-scattering spectrum of the CuInSe₂NRs. An intense A_1 phonon mode at 175.1 cm⁻¹ (fwhm = 9.3 cm⁻¹) is observed with laser excitation at 488 nm.

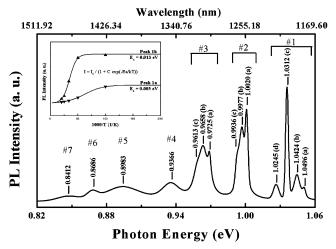


Figure 6. PL of the CuInSe₂NRs at 10 K with laser excitation at 760 nm. Inset shows the fitted activation/binding energies (E_a) for peaks #1a and #1b.

seven groups denoted #1-#7. The fine features in #1-#3 are further labeled by a, b, c, and d.

A CuInSe₂ semiconductor has a direct band gap of ~1.05 eV at low temperature. It has been found that the tetragonal degeneracy in the valence-band maximum is lifted under both crystal-field and spin-orbital interactions. The splitting results in two closely spaced levels (with a separation of only a few millielectronvolts) and another one lying 233 meV below.^{30,31} The #1 group, composed of one prominent and three weak features in Figure 6, is related to the transitions of bound excitons and free excitons. Yakushev et al.32,33 studied the PL spectra of single-crystal CuInSe2 at different temperatures and found that the intensities of both #1a (1.0496 eV) and #1b (1.0424 eV) decrease slowly with increasing temperature; in contrast, those of #1c (1.0312 eV) and #1d (1.0245 eV) decrease rapidly. On the basis of these results, while #1a and #1b are assigned to free-exciton transitions, #1c and #1d are due to bound-exciton transitions. To determine the energy band gap for the #1a and #1b transitions, the PL intensities of these two transitions were examined from 10 to 50 K as illustrated in the inset of Figure 6. Fittings in the plots of PL intensity vs T⁻¹ yield the activation/binding energies (E_a) of 5 and 13 meV for #1a and #1b, respectively. The binding energies of #1a and #1b are further confirmed by their peak locations at 5.4 and 12.6 meV, respectively, below the band gap energy of 1.055 eV.

Features in groups #2-#7 are similar to those reported in previous PL observations of CuInSe₂ crystals.^{34–36} Both #2 and

#3 have one sharp peak at the leading edge, followed by a broad overlapping. The separations of #1c, #2a, and #3a are 29.2-29.5 meV, close to the reported LO phonon mode of 28.9 meV (233 cm⁻¹) determined by Raman scattering spectroscopy.³⁷ In view of the relative intensities of #1c, #2a, and #3a with a gradual decrease in intensity and a comparable phonon-energy separation, these three peaks should belong to a series of vibrational progression. Therefore, #2a and #3a are assigned to the 1LO and 2LO replicas of the strong #1c signal of boundexciton emission.

Other transitions of groups #2 and #3 in Figure 6 are ascribed to the emissions related to different defect centers. The #2b (0.9977 eV) and #3b (0.9658 eV) peaks correspond to free-tobound transitions with the bound states lying at 57 and 89 meV, respectively, below the energy band gap of 1.055 eV. Theoretical prediction for the energies of defect formations can be used as a guideline to identify probable defect centers responsible for the two defect levels.³⁸ According to the theoretical estimation of Wasim,³⁸ the formation energies of the two lowest intrinsic defects are 1.4 eV for In_{Cu} (indium on copper antisite) and 1.5 eV for Cu_{In} (copper on indium antisite). However, an ab initio self-consistent calculation of the electronic structures of CuInSe2 chalcogenide by Zhang et al.39 showed that it is easier to form a Cu vacancy (V_{Cu}) than the antisites of In_{Cu} and Cu_{In} . The calculated formation energies of V_{Cu} , In_{Cu} , and Cu_{In} are 0.60, 1.85, and 1.54 eV, respectively. On the basis of the study of activation energies for the native point defects of CuInSe₂ by Rincón et al., 27 the observed acceptor level of #2b with E_a 57 meV is likely due to the acceptor defects of V_{Cu} , Se_{In} (selenium on indium antisite), or Cu_{Se} (copper on selenium antisite). By taking into account both the favorably lower formation energy^{38,39} and the $E_a = 45 \pm 12$ meV calculated for acceptor levels, ²⁷ #2b is assigned to the vacancy defect of V_{Cu}.

Peak #3b, with $E_a = 89$ meV, is located in the energy range of acceptor levels according to previous studies.^{27,35} The $E_a =$ 83 meV for Cu_{In}, calculated from the optical transitions observed in Cu-rich CuInSe₂ thin films, 33 is close to the observed $E_a =$ 89 meV for #3b. However, the defects of interstitial selenium (Se_i) and Cu_{In} in CuInSe₂ with their $E_a = 78 \pm 20$ meV calculated by Rincón et al.²⁷ could also be responsible for peak #3b. While others assigned the acceptor level of this energy range to Cu_{In}, 40 Kazmerski et al. 41 identified an acceptor level of $E_{\rm a}\sim 60$ meV to Se_i. Taking the formation energy into consideration,³⁸ the defect of Cu_{In} is easier to create than Se_i; therefore, #3b is assigned to Cu_{In} in this work. The Cu_{In} defect has also been reported to exist in p-type CuInSe₂ films.^{34–36,40}

Weak and broad peak #4 centered at 0.9366 eV is assigned to the LO phonon replica of #3b based on the relative intensities of these two peaks and their separation of 29.2 meV, comparable to the LO phonon energy. It has been found that the PL resulting from donor-acceptor pair (DAP) transitions³⁴ appears at ~0.9 eV. Peak #5 at 0.8983 eV, close to the DAP transitions, is identified as another free-to-bound transition involving a deep level of $E_a = 156.7$ meV. In this work, the cation/anion ratio for the CuInSe₂NRs is determined to be 0.95 from the EDS analysis. With this cation/anion ratio, the Se_i defect with $E_a =$ $184 \pm 45 \text{ meV}$ calculated by Rincón et al.²⁷ can account for the observed acceptor level of peak #5 with $E_a = 156.7$ meV. Peak #6 at 0.8686 eV is 29.7 meV from peak #5. Of similar separation, peak #7 at 0.8412 eV is 27.4 meV from peak #6. As discussed earlier about the phonon replicas, #6 and #7 are assigned to the 1LO and 2LO of #5. The LO modes of 27.4-29.7 meV (221-240 cm⁻¹) phonon energy in the present experiment is in good agreement with the phonon frequency of

TABLE 1: Observed PL Transitions of CuInSe₂NRs with Their Peak Energies and Spectroscopic Assignments

peak	peak energy (eV)	spectroscopic assignment
#1a	1.0496	free exciton emission
#1b	1.0424	free exciton emission
#1c	1.0312	bound exciton emission
#1d	1.0245	bound exciton emission
#2a	1.0020	phonon replica (1LO) of #1c
#2b	0.9977	copper vacancy (V _{Cu})
#3a	0.9725	phonon replica (2LO) of #1c
#3b	0.9658	copper on indium antisites (Cu _{In})
#4	0.9366	phonon replica (1LO) of #3b
#5	0.8983	interstitial selenium (Se _i)
#6	0.8686	phonon replica (1LO) of #5
#7	0.8412	phonon replica (2LO) of #5

233 cm⁻¹ observed by Raman spectroscopy.³⁷ The observed PL transitions with their peak energies and spectroscopic assignments are tabulated in Table 1.

4. Conclusions

Single-crystalline CuInSe $_2$ NRs have been successfully synthesized in a soft solvothermal reaction with ethylenediamine as the reaction solvent. The growth of CuInSe $_2$ NRs stems from the chelated-complex formation of [Cu(H $_2$ NCH $_2$ CH $_2$ NH $_2$) $_2$] $^+$, which serves as a molecular template to combine with InSe $_2$ $^-$ for the growth of one-dimensional nanostructures. XRD and electron microscopic measurements of CuInSe $_2$ NRs have shown a single-crystalline structure for the chalcopyrite crystals. The NIR absorption spectrum, room-temperature Raman scattering, and cryogenic PL from the CuInSe $_2$ NRs have further been examined to optically characterize the single-crystalline one-dimensional chalcopyrite nanorods.

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