DFT Study of CH₄ Adsorption on the (5,0), (4,4), (5,5) and (6,6) Single-Walled Carbon Nanotubes. Calculated Binding Energies, NMR and NQR Parameters

B.B. Shirvani^a, M.B. Shirvani^b, J. Beheshtian^a and N.L. Hadipour^{a,*}

^aDepartment of Chemistry, Tarbiat Modares University, P.O. Box: 14115-175, Tehran, Iran

^bDepartment of Mathematics, Payame Noor University, Borujen, Iran

(Received 24 January 2010, Accepted 29 July 2010)

Behavior of a single CH₄ molecule adsorbed on external surface of H-capped (4,4) (5,5), (6,6) and (5,0) single-walled carbon nanotubes (SWCNTs) is studied *via* B97D hybrid density functional and 6-31G* basis set. Binding energies clearly exhibit adsorption dependence on tube diameter. ¹³C and ¹H chemical shielding tensors are calculated at the B971 level using GIAO method. The ¹H and ¹³C(CH₄) NMR results reveal that chemical shielding due to CH₄ molecule adsorption are also dependent upon the nanotube electronic structure, and radius. ²H nuclear quadrupole coupling constants, C_Q , and asymmetry parameter, η , reveal the remarkable effect of CH₄ adsorption on electronic structure of the SWCNTs.

Keywords: Single-walled carbon nanotube, Binding energy, Chemical shielding, NMR

INTRODUCTION

A SWNT can be viewed as a rolled graphene sheet and is classified by a vector connecting the two points that meet upon rolling. There are three classes of nanotube: armchair, zigzag, and chiral. These are labeled by (n,m), (n,0), and (n,m), respectively. One of the most striking features of SWNTs is that their electronic structures depend on the chirality and diameter [1]. SWNTs are (1) metals when n = m, (2) narrowgap semimetals when n-m is a multiple of three, and (3) moderate-gap semiconductors otherwise, where (n,m) is the chiral index [1,2]. The adsorption of gases on and inside carbon nanotubes at low temperatures has been extensively studied. Hydrogen is of particular interest for energy applications. Additionally, the physisorption of Xe [3,4] and CF₄ [5] has recently been examined. The adsorption of gases, particularly CH₄, C₂H₂, C₂H₄, CH₃OH [6]; CO and NO [7];

and NH₃ [8] on a similar substrate, C₆₀, has also been studied. For CH₄, C₂H₂, C₂H₄, CO and NO, the adsorption was shown to be caused by dispersion forces (*i.e.*, physisorption). Because of the weakness of the interaction, these gases would adsorb only below room temperature.

A variety of techniques such as scanning tunneling microscopy (STM) [9], optical adsorption [10], fluorescence spectroscopy [11] and Raman scattering [12] have already been applied to the characterization of SWNT samples. Unfortunately, these techniques even in combination do not provide a full characterization [13]. Thus, the development of alternative techniques is desired. ¹³C nuclear magnetic resonance (NMR) spectroscopy is one possible choice. Nuclear experimental techniques such as nuclear magnetic resonance (NMR) [14] and nuclear quadrupolar resonance (NQR) [15] are widely used to study the geometry and electronic structure of molecules and solids. NMR measures the local magnetic fields on a nucleus, generated by response of electrons to an external uniform magnetic field. In order to

^{*}Corresponding author. E-mail: hadipour@modares.ac.ir

form a non-degenerated energy spectrum, both NMR and NQR techniques are applied to produce high external magnetic fields and some kind of internal interaction; *i.e.*, the spin-spin interaction in NMR [16-18] and the interaction of quadrupole moment with electron field gradient (EFG) in NQR [19-21]. Both of the interactions are determined by internal properties of crystal. However, the field has recently started to bloom, and an increasing amount of experimental and theoretical data is becoming available.

Most computational solid-state and quantum-chemical studies of SWNTs have so far focused on geometric and electronic structure as well as mechanical properties. It is important to study electric and magnetic response properties as well, since they are the observables of some of the most powerful and frequently applied spectroscopic methods used for characterizing molecules. Among those, NMR is of high practical importance. Previous theoretical work proposed that separation of metallic from semiconducting SWNTs may be possible using high resolution ¹³C NMR, based on a predicted 11 ppm difference in their chemical shifts [22,23]. Solid-state ¹³C NMR has been applied to characterize nanotube carbons in the functionalized SWCNT samples [24,25]. The signals corresponding to sp² carbons of nanotube are generally broad, centered around 120-130 ppm, and similar to those of unfunctionalized SWNTs [26-29].

In this article, DFT calculations were used to study the adsorption of CH₄ on metallic armchair (4,4), (5,5), (6,6) and semi metallic zigzag (5,0), carbon nanotubes. CH₄ effects on ¹³C and ¹H NMR chemical shielding tensors, and difference in nuclear quadrupole resonance parameters, NQR, for ²H nuclei are calculated to obtain useful information about the nature of interactions.

COMPUTATIONAL DETAILS

Full geometry optimizations are performed for the (5,0), (4,4), (5,5) and (6,6) pristine and functionalized SWCNTs using *Gamess* and *Gaussian* suite of programs [30]. Geometry optimizations are carried out using B97D hybrid density functional and 6-31G* basis set [31]. The B97D functional was chosen because, among the presently available functionals, it give a better description of dispersion interactions, and hence, of the physical adsorption. The

starting geometries are generated using *TubeGen* tool [32] based on a hexagonal unit cell. Optimized diameters are 4.10 Å for zigzag and 5.69, 7.03, 8.44 Å for armchair nanotubes, respectively. Length of nanotubes are 13.8 Å and 10.44, 10.46, 10.47 Å, respectively, with an average C-C bond length of 1.43 Å. The adsorption energy is defined by $E_{\rm ads} = E$ (adsorbate and CNT) - E(CNT) - E(adsorbate). Application of localized basis sets largely reduces the amount of computational work required while considering large vacuum regions in the unit cell. The bottom dangling bonds are saturated by hydrogen atoms. Chemical shielding tensors are computed for the pristine and SWCNT-molecule systems using Slater-type gauge-including atomic orbital (GIAO) by B971 hybrid density functional and 6-31G* basis set [33].

The interaction between nuclear electric quadrupole moment and EFG at quadrupole nucleus is described with Hamiltonian:

$$\hat{H} = \frac{e^2 Q V_{zz}}{4I(2I-1)} [(3\hat{I}_z^2 - \hat{I}^2) + \eta_{Q}(\hat{I}_x^2 - I_y^2)]$$

Here eQ is the nuclear electric quadrupole moment, I is the nuclear spin, and V_{zz} is the largest component of EFG tensor. The principal components of the EFG tensor, V_{ii} , are computed in atomic unit (1 $au = 9.717365 \times 10^{21} \ V \ m^{-2}$), with $|V_{zz}| \geq V_{yy} \geq |V_{xx}|$ and $V_{xx} + V_{yy} + V_{zz} = 0$. These diagonal elements relate to each other by the asymmetry parameter; $y = |V_{yy} - V_{xx}|/|V_{zz}|$, $0 \leq y \leq 1$, that measures the deviation of EFG tensor from axial symmetry. The computed V_{zz} component of EFG tensor is used to obtain nuclear quadrupole coupling constant from the equation; $C_Q = e^2 Q V_{zz}/h$ [34], where standard values of $eQ(^2H)$ reported by Pyykkö are employed [35]. EFG calculations at hydrogen nuclei are determined using B971/6-31G* level of theory.

RESULTS AND DISCUSSION

Interaction of Methane with SWCNTs

Interaction of (5,0) zigzag and (4,4), (5,5), (6,6) armchair SWCNTs with a single CH₄ molecule is considered. To find adsorption behavior of CH₄ on SWCNTs, one methane molecule is moved toward the hollow or bridge site of the tube wall, and equilibrium distance of molecule tube are then

found, Fig. 1. Results for binding energies, equilibrium distance and dipole moments of zigzag and armchair tubes with different diameters are presented in Table 1. Calculated adsorption energies tube-CH₄ for (5,0), (4,4), (5,5) and (6,6) are negative. These results are in agreement with previous results [36]. In general, CH₄ molecules are weakly bound to the nanotubes and the tube-molecules interactions can be identified as physisorption. Binding energy is calculated to be -3.22 kcal mol⁻¹ for zigzag (5.0) SWCNT. Effect of the tube diameter on stability of the tube-molecules is interesting. For armchair CNTs, Stability of tube-CH4 decreases as their diameter increases. These results are in agreement with previous results indicating that gaseous molecules binding energy decreases as the diameter of (n,0) and (n,n) CNTs increases [37,38]. Results for armchair tubes reveal that R is magnified while moving from a small diameter (4,4) SWCNT (d = 5.70 Å) to a larger diameter (6,6) tube (d = 8.44 Å).

The calculated bond lengths for (5,5) and (6,6) SWCNTs

are 3.38 and 3.39 Å, respectively. These values are close to the corresponding those for other armchair SWCNTs [37-39]. When CH₄ molecule is adsorbed on the CNT, dipole moment of the entire system is significantly increased by a factor of about 0.20-0.25 with respect to the vacuum values. This effect can be understood by polarization of the conducting electrons due to physisorption. However, due to different CH₄ configurations on the surface of nanotubes, no regular pattern is detected for magnitude of their dipole moment (Table 1).

¹³C NMR Chemical Shifts

Previously, it has been indicated that for H-capped fragments, the calculated ¹³C chemical shielding values are smaller at the ends of tube, compared to the center, if the carbon is directly bound to hydrogen; otherwise, it is larger [39]. To assess the dependence of NMR results on carbon atom position, ¹³C chemical shielding values of armchair and zigzag tubes are examined. Three different parts of tube axis,

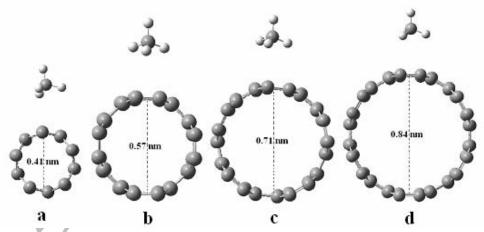


Fig. 1. Equilibrium geometries and pictorial view of the adsorption sites of CH₄-SWCNTs structures: a) (5,0), b) (4,4), c) (5,5) and d) (6,6) CNTs.

Table 1. Calculated Adsorption Energies E_b (kcal mol⁻¹), Equilibrium Tube-Molecule Distance^b R (Å), and Dipole Moment^a of the CH₄ Adsorbed on Zigzag (5,0) and Armchair (n,n), n = 4, 5, 6 Nanotubes

Parameter	$(5,0) + CH_4$	$(4,4) + CH_4$	$(5,5) + CH_4$	$(6,6) + CH_4$
E _b	-3.22	-3.49	-3.40	-3.29
R	3.40	3.38	3.38	3.39
Dipole moment	0.25	0.20	0.21	0.21

^aCalculated dipole moment of single CH₄ is 0.00 Debye. ^bCC equilibrium distance.

S112 www.SID.ir

(two ends and middle parts) are considered (Fig. 2).

Isotropic chemical shielding. Table 2 exhibits the calculated ¹³C isotropic chemical shielding tensors for various CNTs. CH₄ molecule adsorption on external surface of CNTs has a remarkable influence on ¹³C NMR tensors, which is in complete accordance with the facts mentioned above. Interesting trends are evidenced: For (5,0) and (4,4) CNTs, the isotropy of ¹³C shielding tensor are larger at the ends compared to the center (Fig. 2). It can be also seen from Fig. 2 that the isotropy of ¹³C chemical shielding decreased at the ends compared to the center for armchair tubes with increase of the tube diameter. The shifts are clearly sensitive to the nanotube radius. DFT study of ¹³C chemical shielding tensors on small-to-medium-diameter infinite SWCNTs reveal a roughly inversely proportional relationship between chemical shielding and tube diameter [40]. It may be noted that ¹³C chemical shielding tensor at the carbon sites depends remarkably on the tube size and nature of frontier orbitals [41].

According to GIAO calculations performed after adsorption of CH₄ molecule on (5,0) CNT, the isotropy value of the ¹³C NMR shielding tensor decreases, approximately by 0.63, 0.25, 1.04, ppm, at the C1,C2,C3, (carbon atoms located near CH₄ molecule) sites (Table 2). These shifts are much smaller than those calculated for N-H functionalized finite (11,0), (13,0) and (14,0) CNTs (ca., 44 ppm) [42]. As for the armchair (6,6) CNT, the ¹³C chemical shielding isotropy values are slightly deshielded for all carbon atoms after CH₄ adsorption. In case of (4,4) CNT, the isotropy value of the ¹³C NMR shielding tensors of interacted carbon atoms are slightly increases the C1,C2, sites as a result of CH₄ adsorption. Comparing pristine (5,5) and tube-molecule, it becomes clear that the C1 and C2 atoms included in CH₄ adsorption are slightly deshielded. The results are consistent with the weak interaction between the tube and CH₄ molecule. The discrepancy between the 13C chemical shielding tensor variation for the zigzag and armchair nanotubes can be attributed to the different nature of the frontier orbitals.

However, it is obvious that there is a large effect of the HOMO on the energies of the HOMO-1 *etc*. which will have an influence on the 13 C chemical shieldings (the expression for the paramagnetic shielding tensor contains $(\Delta E_{VO})^{-1}$, the inverse of the energy differences between virtual and occupied

orbitals).

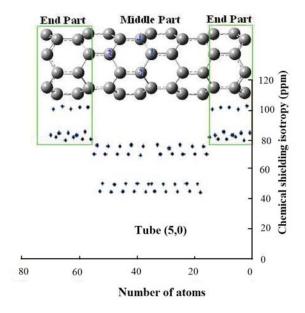
Anisotropic chemical shielding. Interesting trends are evidenced: For (5,0), (4,4), (5,5) and (6,6) CNTs, the anisotropy of ¹³C shielding tensor are smaller at the ends compared to the center (Fig. 2). Table 3 is shown the calculated ¹³C anisotropic chemical shielding tensors for various CNTs. According to GIAO calculations performed after adsorption of CH₄ molecule on (5,0) CNT, the anisotropy value of the ¹³C NMR shielding tensor increases for C1 (2.51 ppm), C3 (3.73 ppm), and decreases 0.42 ppm for C2, sites. For the armchair (4,4), (5,5) and (6,6) CNTs, the ¹³C NMR chemical shielding anisotropy values of interacted carbon atoms are slightly increased after CH₄ adsorption (Table 3). In case of (6,6) CNT, ¹³C NMR tensors of interacted carbon atoms are also modified as a result of CH₄ adsorption.

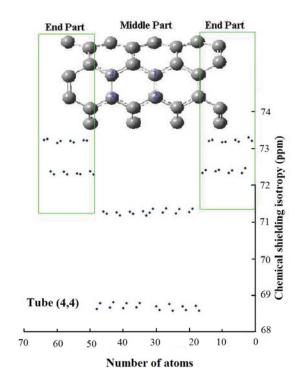
Chemical shielding ¹H and ¹³C methane. The ¹³C and ¹H chemical shielding tensors (isotropy and anisotropy) of adsorbed CH₄ molecule on external surface of zigzag (5,0) and armchair (4,4), (5,5), (6,6) CNTs are shown in Table 4. Results indicate that the ¹³C and ¹H NMR chemical shielding are also dependent upon the nanotube electronic structure, and radius. More especially, the ¹³C NMR isotropy and anisotropy values increase from CH₄ to CH₄-tube. The ¹H NMR isotropy value except for (5,0) increased for armchair tubes from CH₄ to CH₄-tubes. However, small differences are detected in proton NMR calculations of the nanotubes. This is mainly due to the limited interaction of hydrogen atoms with nanotubes.

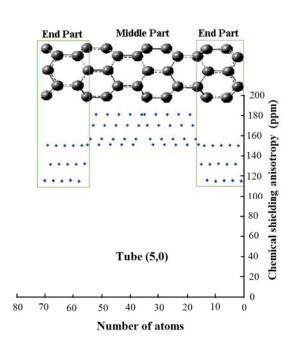
²H NQR Parameters

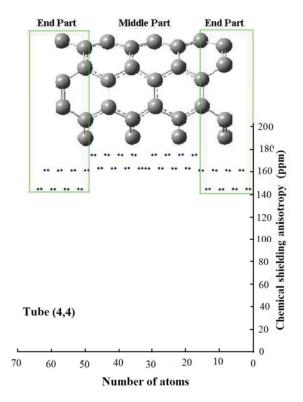
Due to the physisorption, NQR parameters of hydrogen atoms are also altered. Table 5 indicates the calculated NQR parameters (C_Q , and η) at the hydrogen positions for various SWCNTs. A quick look at the results reveals that the C_Q , and η values for 2 H sites are modified by CH₄ adsorption. More especially, the C_Q (2 H) parameter decreased at the H₁, H₄ and increased at the H₂,H₃ sites from single CH₄ molecule to CH₄-tube. The asymmetry parameter η (2 H) parameter increased from single CH₄ molecule to CH₄-tube. The asymmetry parameter increases (by 0.0003-0.007 units) as result of CH₄-tube interaction. In assessing the significance of C_Q (2 H) and η (2 H) values, one must keep in mind that electronic contribution to the EFG at hydrogen sites arises almost entirely from mean carbon atoms other than the hydrogen, since near spherical

DFT Study of CH₄ Adsorption on Single-Walled Carbon Nanotubes









S114 www.SID.ir

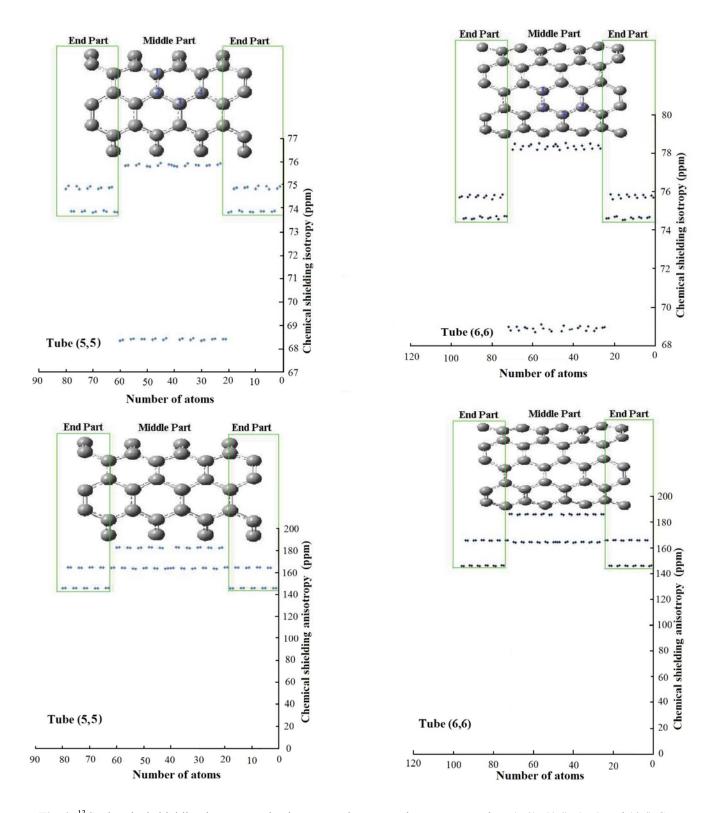


Fig. 2. ¹³C chemical shielding isotropy and anisotropy values *vs.* carbon atoms number. (5,0), (4,4), (5,5) and (6,6) CNTs.

Table 2. Calculated ¹³C Chemical Shielding Isotropy Values (in ppm) for SWCNTs

Species	C_1	C_2	C ₃	C ₄
Tube (5,0)	45.28	45.31	70.42	45.28
Tube $(5,0) + CH_4$	44.65	45.06	69.38	45.35
Tube (4,4)	68.61	68.79	71.18	71.26
Tube $(4,4) + CH_4$	68.89	68.88	70.73	70.71
Tube (5,5)	75.87	75.85	68.37	75.81
Tube $(5,5) + CH_4$	75.79	74.71	68.46	75.87
Tube (6,6)	78.27	78.52	69.12	78.21
Tube $(6,6) + CH_4$	77.41	77.50	68.85	78.21

Table 3. Calculated ¹³C Chemical Shielding Anisotropy Values (in ppm) for SWCNTs

Species	C_1	C_2	C ₃	C ₄
Tube (5,0)	170.26	170.22	151.45	170.38
Tube $(5,0) + CH_4$	172.77	170.22	155.18	169.96
Tube (4,4)	175.52	175.35	163.31	163.12
Tube $(4,4) + CH_4$	175.83	175.82	166.62	166.66
Tube (5,5)	164.26	164.38	183.07	164.33
Tube $(5,5) + CH_4$	166.06	168.55	183.38	164.99
Tube (6,6)	164.76	164.37	186.05	164.63
Tube $(6,6) + CH_4$	168.78	168.52	187.24	165.23

Table 4. Calculated ¹³C and ¹H Chemical Shielding Isotropy (Anisotropy) Values (in ppm) a Single CH₄ Molecule and for CH₄-SWCNTs

Species	¹³ C	¹ H	¹ H	¹ H	¹ H
Tube $(5,0) + CH_4$	191.82(5.83)	31.78(11.76)	31.52(4.04)	31.36(8.23)	31.10(10.65)
Tube $(4,4) + CH_4$	192.71(12.44)	33.10(8.21)	33.04(8.07)	32.91(13.09)	32.76(15.03)
Tube $(5,5) + CH_4$	192.29(13.85)	33.41(10.34)	33.25(9.03)	33.09(13.04)	32.93(15.94)
Tube $(6,6) + CH_4$	193.15(13.58)	33.87(10.55)	33.56(10.47)	33.51(14.10)	33.20(16.76)
Single CH ₄	193.16(0.00)	31.70(8.57)	31.70(8.57)	31.70(8.57)	31.70(8.57)

Table 5. Calculated C_O and η Parameters for SWCNTs and Single CH₄ Molecule^{a,b}

Species	$^{2}\mathrm{H}_{1}$	$^{2}\mathrm{H}_{2}$	$^{2}\mathrm{H}_{3}$	$^{2}\mathrm{H}_{4}$
$(5,0) + CH_4$	236.13(0.007)	238.44(0.003)	238.28(0.002)	237.74(0.002)
$(4,4) + CH_4$	237.36(0.006)	238.14(0.002)	238.12(0.002)	237.36(0.0003)
$(5,5) + CH_4$	237.61(0.004)	238.18(0.002)	238.13(0.002)	237.64(0.0003)
$(6,6) + CH_4$	237.73(0.005)	238.49(0.003)	238.09(0.003)	237.64(0.0006)
Single CH ₄	237.78(0.000)	237.78(0.000)	237.78(0.000)	237.78(0.000)

^aCalculated C_Q values in kHz for ²H nuclei. ^bThe first value in each cell represents C_Q value, while value within parenthesis indicates η value.

electronic cloud around these nuclei has no contribution. As the CH₄-tube interaction sets in, the electron distribution around the hydrogen nucleus becomes less polarized in the direction of carbon atom and makes a significant reduction in magnitude of all the three individual EFG components.

CONCLUSIONS

Theoretical calculations are performed to characterize the behavior of CH₄ molecule adsorption on external surface of zigzag (5,0), and armchair (4,4), (5,5), (6,6) SWCNTs. Results indicate that as the diameter of armchair tubes increases, the binding energy of CH4 molecule decreases. According to GIAO calculations, performed after CH₄ adsorption on (5,0) CNT, isotropy value of ¹³C NMR shielding tensors decreases, approximately by 0.63, 0.25, 1.04, ppm, at the C1,C2,C3, sites whereas the anisotropy value of the ¹³C NMR shielding tensor increases for C1(2.51 ppm), C3(3.73 ppm), and decreases 0.42 ppm for C2, sites, respectively. In addition, for (4,4), (5,5), (6,6) tubes are slightly increased anisotropy value of the ¹³C NMR shielding tensor after CH₄ adsorption. The ¹H and ¹³C(CH₄) NMR results reveal that chemical shielding due to CH₄ molecule adsorption are also dependent upon the nanotube electronic structure, and radius. Due to the physisorption, NQR parameters of hydrogen atoms are also altered. The asymmetry parameter increases (by 0.0003-0.007 units).

REFERENCES

- [1] N. Hamada, S.-I. Sawada, A. Oshiyama, Phys. Rev. Lett. 68 (1992) 1579.
- [2] L. Lai, J. Lu, W. Song, M. Ni, L. Wang, G. Luo, J. Zhou, W.N. Mei, Z. Gao, D. Yu, J. Phys. Chem. C 112 (2008) 16417.
- [3] A. Kuznetsova, J.T. Yates, J. Liu, R.E. Smalley, J. Chem. Phys. 112 (2000) 9590.
- [4] A. Kuznetsova, J.T. Yates, V.V. Simonyan, J.K. Johnson, C.B. Huffman, R.E. Smalley, J. Chem. Phys. 115 (2001) 6691.
- [5] O. Byl, P. Kondratyuk, S.T. Forth, S.A. FitzGerald, J.K. Chen, J.T. Johnson, J. Yates, J. Am. Chem. Soc. 125 (2003) 5889.
- [6] A. Lubezky, L. Chechelnitsky, M. Folman, J. Chem. Soc., Faraday Trans. 92 (1996) 2269.
- [7] M. Fastow, Y. Kozirovski, M. Folman, J. Heidberg, J. Phys. Chem. 96 (1992) 6126.
- [8] A. Lubezky, L. Chechelnitsky, M. Folman, Surf. Sci. 454 (2000) 147.
- [9] J.W.G. Wildoer, L.C. Venema, A.G. Rinzler, R.E. Smalley, C. Dekker, Nature 391 (1998) 59.
- [10] A.G. Marinopoulos, L. Reining, A. Rubio, N. Vast, Appl. Phys. A: Mater. Sci. Process. 78 (2004) 1157.
- [11] R.B. Weisman, S.M. Bachilo, D. Tsyboulski, Appl. Phys. A: Mater. Sci. Process 78 (2004) 1111.

- [12] M.S. Dresselhaus, G. Dresselhaus, R. Saito, A. Jorio, Phys. Rep. 409 (2005) 47.
- [13] N.A. Besley, J.J. Titman, M.D. Wright, J. Am. Chem. Soc. 127 (2005) 17948.
- [14] D. Canet, Nuclear Magnetic Resonance: Concepts and Methods, Wiley, Chichester, 1996.
- [15] G.K. Semin, T.A. Babushkina, G.G. Yakobson, Nuclear Quadrupole Resonance in Chemistry, Wiley, New York, 1975.
- [16] N.A. Gershenfeld, I.L. Chuang, Science 275 (1997) 350
- [17] D.G. Gory, A.F. Fahmy, T.F. Havel, Proc. Natl. Acad. Sci. USA 94 (1997) 1634.
- [18] J.A. Jones, Fort. Der Physik 48 (2000) 909.
- [19] A.R. Kessel, V.L. Ermakov, JETP Lett. 70(1999) 61.
- [20] A.K. Khitrin, B.M. Fung, J. Chem. Phys. 112 (2000) 6963.
- [21] A.K. Khitrin, H. Song, B.M. Fung, Phys. Rev. A 63 (2001) 1.
- [22] X.-P. Tang, A. Kleinhammes, H. Shimoda, L. Fleming, K.Y. Bennoune, S. Shina, C. Bower, O. Zhou, Y. Wu, Science 288 (2000) 492.
- [23] S. Hayashi, F. Hoshi, T. Ishikura, M. Yumura, S. Ohshima, Carbon 41 (2003) 3047.
- [24] H. Peng, L.B. Alemany, J.L. Margrave, V.N. Khabashesku, J. Am. Chem. Soc. 125 (2003) 15174.
- [25] L.S. Cahill, Z. Yao, A. Adronov, J. Penner, K.R. Moonoosawmy, P. Kruse, G.R. Goward, J. Phys. Chem. B 108 (2004) 11412.
- [26] X.-P. Tang, A. Kleinhammes, H. Shimoda, L. Fleming, K.Y. Bennoune, S. Sinha, C. Bower, O. Zhou, Y. Wu, Science 288 (2000) 492.
- [27] C. Goze-Bac, S. Latil, P. Lauginie, V. Jourdain, J. Conard, L. Duclaux, A. Rubio, P. Bernier, Carbon 40 (2002) 1825.
- [28] S. Hayashi, F. Hoshi, T. Ishikura, M. Yumura, S. Ohshima, Carbon 41 (2003) 3047.
- [29] A. Kleinhammes, S.-H. Mao, X.-J. Yang, X.-P. Tang, H. Shimoda, J.P. Lu, O. Zhou, Y. Wu, Phys. Rev. B 68 (2003) 754181.
- [30] M.J. Frisch, G.W. Trucks, H.B. Schlegel, G.E. Scuseria,

- M.A. Robb, J.R. Cheeseman, V.G. Zakrzewski, J.A. Montgomery, R.E. Stratmann, J.C. Burant, S. Dapprich, J.M. Millam, A.D. Daniels, K.N. Kudin, M.C. Strain, O. Farkas, J. Tomasi, V. Barone, M. Cossi, R. Cammi, B. Mennucci, C. Pomelli, C. Adamo, S. Clifford, J. Ochterski, G.A. Petersson, P.Y. Ayala, Q. Cui, K. Morokuma, D.K. Malick, A.D. Rabuck, Raghavachari, J.B. Foresman, J. Cioslowski, J.V. Ortiz, A.G. Baboul, B.B. Stefanov, G. Liu, P. Liashenko, A. Piskorz, I. Komaromi, R. Gomperts, R.L. Martin, D.J. Fox, T. Keith, M.A. Al-Laham, C.Y. Peng, A. Nanayakkara, C. Gonzalez, M. Challacombe, P.M.W. Gill, B. Johnson, W. Chen, M.W. Wong, J.L. Andres, C. Gonzalez, M. Head-Gordon, E.S. Replogle, J.A. Pople, Gaussian, Gaussian Inc., Pittsburgh, P.B., 2003.
- [31] S. Grimme, J. Comp. Chem. 27 (2006) 1787.
- [32] J.T. Frey, D.J. Doren TubeGen 3.3; University of Delaware: Newark, DE, web interface, http://turin.nss.udel.edu/research/tubegenonline.html, 2005.
- [33] F.A. Hamprecht, A. Cohen, D.J. Tozer, N.C. Handy, J. Chem. Phys. 109 (1998) 6264.
- [34] E.A.C. Lucken, Nuclear Quadrupole Coupling Constants, Academic Press, London, 1992.
- [35] P. Pyykkö, Mol. Phys. 99 (2001) 1617.
- [36] G. Vidali, G. Ihm, H.Y. Kim, M.W. Cole, Surf. Sci. Rep. 12 (1991) 133.
- [37] M.R. Johnson, S. Rols, P. Wass, M. Muris, M. Bienfait, P. Zeppenfeld, N. Dupont-Pavlovsky, Chem. Phys. 293 (2003) 217.
- [38] S. Talapatra, V. Krugleviciute, A.D. Migone, Phys. Rev. Lett. 89 (2002) 2461061.
- [39] B.K. Agrawal, S. Agrawal, S. Singh, R. Srivastava, J. Phys. Condens. Matter 18 (2006) 4649.
- [40] L. Lai, J. Lu, W. Song, M. Ni, L. Wang, G. Luo, J. Zhou, W.N. Mei, Z. Gao, D. Yu, J. Phys. Chem. C 112 (2008) 16417.
- [41] E. Zurek, J. Autschbach, J. Am. Chem. Soc. 126 (2004) 13079.
- [42] E. Zurek, C.J. Pickard, B.J. Autschbach, J. Phys. Chem. C 112 (2008) 9267.

S118 www.SID.ir