

## Experimental $pK_a$ Values of Buried Residues: Analysis with Continuum Methods and Role of Water Penetration

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**ABSTRACT** Lys-66 and Glu-66, buried in the hydrophobic interior of staphylococcal nuclease by mutagenesis, titrate with  $pK_a$  values of 5.7 and 8.8, respectively (Dwyer et al., 2000, *Biophys. J.* 79:1610–1620; García-Moreno E. et al., 1997, *Biophys. Chem.* 64:211–224). Continuum calculations with static structures reproduced the  $pK_a$  values when the protein interior was treated with a dielectric constant ( $\epsilon_{in}$ ) of 10. This high apparent polarizability can be rationalized in the case of Glu-66 in terms of internal water molecules, visible in crystallographic structures, hydrogen bonded to Glu-66. The water molecules are absent in structures with Lys-66; the high polarizability cannot be reconciled with the hydrophobic environment surrounding Lys-66. Equilibrium thermodynamic experiments showed that the Lys-66 mutant remained folded and native-like after ionization of the buried lysine. The high polarizability must therefore reflect water penetration, minor local structural rearrangement, or both. When in  $pK_a$  calculations with continuum methods, the internal water molecules were treated explicitly, and allowed to relax in the field of the buried charged group, the  $pK_a$  values of buried residues were reproduced with  $\epsilon_{in}$  in the range 4–5. The calculations show that internal waters can modulate  $pK_a$  values of buried residues effectively, and they support the hypothesis that the buried Lys-66 is in contact with internal waters even though these are not seen crystallographically. When only the one or two innermost water molecules were treated explicitly,  $\epsilon_{in}$  of 5–7 reproduced the  $pK_a$  values. These values of  $\epsilon_{in} > 4$  imply that some conformational reorganization occurs concomitant with the ionization of the buried groups.

### INTRODUCTION

Ionizable residues buried in the interior of proteins play key functional roles in fundamental biochemical processes such as catalysis,  $H^+$  and  $e^-$  transport, photo-activation, and redox reactions. Those buried at the interfaces between molecules modulate recognition specificity and binding affinity. Elucidation of the structural basis of biological function in these processes requires quantitative understanding of electrostatic contributions. This usually entails measurement of  $pK_a$  values or redox potentials, elucidation of their molecular determinants, understanding of the molecular mechanisms whereby buried charges are stabilized, and of the structural and dynamic response of proteins to the ionization of a buried residue. For reasons of size and complexity, these issues are difficult to study experimentally in proteins where buried ionizable residues are functionally important, such as bacteriorhodopsin, cytochrome *c* oxidase, or photosynthetic reaction centers. In systems such as these, structure-based electrostatic calculations are necessary to bridge the gap between structure, energy, and function.

Meaningful calculation of electrostatic energies and  $pK_a$  values of buried ionizable residues remains challeng-

ing. The problems stem from difficulties in capturing quantitatively the dielectric relaxation in the protein interior. Fully microscopic approaches are not yet accurate enough to allow quantitative calculations of  $pK_a$  values (Del Buono et al., 1994; Schutz and Warshel, 2001). Therefore, continuum methods in which some of the contributions to the dielectric relaxation of proteins are captured implicitly through the use of dielectric constants remain useful and desirable, as long as they are calibrated against experimental data.

Most previous experimental and computational studies of  $pK_a$  values and their determinants have focused on surface ionizable residues. These groups are not useful to calibrate continuum models because the self-energies of surface groups in proteins are comparable to those in water (Sham et al., 1998). Buried groups with substantial  $pK_a$  shifts offer considerably greater insight about the nature of dielectric relaxation in the protein interior. They are also useful as benchmarks for testing, calibrating, and improving computational methods (Schutz and Warshel, 2001). Unfortunately, progress has been hindered by the lack of experimental  $pK_a$  values of buried groups. To address this problem, systematic, experimental studies of buried ionizable groups are underway in this laboratory. The approach entails burial of ionizable residues by mutagenesis, measurement of  $pK_a$  values, determination of crystallographic structures to describe microenvironments of buried groups, assessment of the consequences of ionization of buried groups on stability and structure, and analysis with structure-based energy calculations (García-Moreno E. et al., 1997; Dwyer et al., 2000).

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Previously, we reported that, when Lys-66 and Glu-66 are buried in a hydrophobic pocket in the interior of staphylococcal nuclease (SNase) by substitution of Val-66, they titrate with  $pK_a$  values of 5.8 and 8.8, respectively (García-Moreno E. et al., 1997; Dwyer et al., 2000). The  $\Delta pK_a$  of 4.6 and 4.3 relative to the average  $pK_a$  values of Lys and Glu in water are among the largest  $\Delta pK_a$  ever measured experimentally. Analysis with a Born formalism showed that the  $\Delta pK_a$  are energetically equivalent to the transfer of an ion from water to a medium with dielectric constant in the range 9–12. This high apparent polarizability in the interior of SNase was initially puzzling. In independent structures of two mutants with Lys-66, obtained under conditions of pH where the buried Lys is neutral, the ionizable moiety of Lys-66 is encased in an extremely hydrophobic environment, incompatible with such high apparent polarizability (Stites et al., 1991; García-Moreno E. et al., 1997). We conjectured initially that the high polarizability reported by Lys-66 reflected a substantial structural relaxation concomitant with the ionization of the buried Lys-66. However, in two subsequent structures with buried Glu-66, also obtained under conditions of pH where the buried Glu is neutral, internal water molecules were found interacting with the carboxyl group of Glu-66 and connecting it to bulk water (Dwyer et al., 2000). These structures suggested that the high apparent polarizability in the interior of SNase could reflect the presence of internal water molecules near the buried ionizable residues.

Many factors can contribute to the energetics of ionization of a buried group, and thus to the apparent polarizability reported by its  $pK_a$ : interactions with backbone and side chain dipoles, with ionizable residues, and with buried water molecules; changes in the state of ionization of other residues; and relaxation of the protein. One of the aims of this study is to dissect the dielectric response in the interior of SNase with a combined experimental and theoretical approach. Equilibrium thermodynamic and  $^1\text{H-NMR}$  experiments were performed to ensure the accuracy of the  $pK_a$  of Lys-66, and to assess the consequences of the V66K mutation on the structure and stability of a hyperstable  $\Delta$ +PHS nuclease variant. The experiments were also necessary to determine the consequences of the ionization of the buried Lys-66 on the structure of the  $\Delta$ +PHS/V66K protein. Specifically, it was important to establish that this protein remained folded and native-like after ionization of the buried group. This information was needed to understand the meaning of the high polarizability reported by the buried ionizable group, and to ensure that it did not simply reflect a large structural transition such as acid denaturation. Electrostatic interactions between buried and surface ionizable groups were also assessed experimentally to determine their effect on the  $pK_a$  values of the buried groups.

The study has three computational aims. First, to quantify the influence of the buried water molecules on the  $pK_a$  value of Glu-66 by comparing different structure-based  $pK_a$

calculations, some of which omitted the internal water molecules, others which treated them explicitly. Second, to test the hypothesis that the buried ionizable moiety of Lys-66 is also in contact with internal water molecules that are either disordered or only transiently buried, and thus crystallographically invisible. Third, to use the experimental  $pK_a$  of Glu-66 and Lys-66 to determine the empirical values of the dielectric constants in the protein interior useful for structure-based calculations of  $pK_a$  values of buried residues. These studies underscore the important contributions from water penetration to  $pK_a$  values of buried groups, and they contribute insight about mechanisms of dielectric relaxation in the interior of SNase.

## MATERIALS AND METHODS

### Staphylococcal nuclease

The hyperstable  $\Delta$ +PHS variant of nuclease includes mutations P117G, H124L, S128A, G50F, V51N, and a deletion from residues 44–49. Clones of  $\Delta$ +PHS nuclease and its V66K mutant,  $\Delta$ +PHS/V66K, were obtained from Prof. David Shortle (The Johns Hopkins University). Protein was expressed and purified following the method of Shortle and Meeker (1986). The protein was determined to be >98% pure by SDS-PAGE. The concentration was determined using an extinction coefficient of  $15,600 \text{ M}^{-1}\text{cm}^{-1}$  at 280 nm.

### Equilibrium thermodynamic measurements

The protocols that were used for measurement of stability by GdnHCl-denaturation monitored by the intrinsic fluorescence of Trp-140 have been described previously (García-Moreno E. et al., 1997; Dwyer et al., 2000; Whitten and García-Moreno E., 2000). The unfolding of  $\Delta$ +PHS nuclease and of some of its mutants by chemical denaturants is considerably slower than for wild-type SNase. In the automated denaturation experiments, samples were allowed to equilibrate for 40–80 min between the addition of denaturant in the transition region. Five minutes were sufficient to reach equilibrium in the baseline regions. The following buffers were used at a concentration of 25 mM to cover the pH range specified: Acetate, pH 4 to 5.5; MES, 5.5–6.5; HEPES, 7–8; TAPS, 8–9; CHES, 9–10; CAPS, 10–11. The experiments were performed with an ATF-105 automated titration fluorometer from Aviv Inc. (Lakeland, NJ). All data were collected at 25°C, 100 mM NaCl.

Protocols for acid-base titrations monitored by fluorescence have also been described previously (Whitten and García-Moreno E., 2000). The only notable difference with experiments reported previously is that the acid-induced unfolding of  $\Delta$ +PHS nuclease is also slower than that of wild type, and required equilibration times as long as 5 min between addition of titrant. All measurements were performed at 25°C, 100 mM NaCl, in a buffer consisting of 5 mM MES and 5 mM HEPES.

The measurement of proton ( $\text{H}^+$ ) titration data with potentiometric methods has been described previously (Dwyer et al., 2000; Whitten and García-Moreno E., 2000). The data in this paper are measured with concentrations of SNase of 3–4 mg/ml. Reversibility of the titration curves was tested routinely to ensure that the system was at equilibrium. All titration curves were measured in triplicate. All measurements were performed at 25°C, 100 mM KCl or in 6 M GdnHCl and 100 mM KCl.

## Measurement of pK<sub>a</sub> values by <sup>1</sup>H-NMR spectroscopy

The procedures of Lecomte and co-workers were used to determine pK<sub>a</sub> values by NMR spectroscopy (Lecomte and Cocco, 1990; Cocco et al., 1992; Bhattacharya and Lecomte, 1997; Kao et al., 2000). Protein samples at concentrations of ~1.2 mM were prepared by solvent exchange in D<sub>2</sub>O containing 100 mM KCl. Amide hydrogen atoms were exchanged by increasing the temperature to the midpoint temperature for 15 min followed by centrifugation. A total volume of at least 1.2 ml was prepared at pH 8.0, and this sample was split into two fractions, one for titration with acid and one for titration with base.

One-dimensional (1D) <sup>1</sup>H-NMR spectra were acquired at 25°C on a Varian Unity Plus 500 MHz spectrometer using a 5-mm triple-resonance probe. Spectra were recorded with low-power water presaturation. A 90° high-power pulse (8 μs) was applied followed by acquisition of 16-K data points with a sweep width of 6024 Hz. Each spectrum contains 256 transients. All spectra were referenced to the chemical shift of the residual HDO line, which in turn was referenced to the external reference sodium 2,2-dimethyl-2-silapentane-5-sulfonate (Wishart et al., 1995). Spectra were recorded at 25°C. The temperature was calibrated based on measurements of the temperature-dependent change in chemical shift of the C and OH group of ethylene glycol (Martin et al., 1980). Temperatures were calibrated to ±1°C and maintained to ±0.1 for both spectral acquisition and pH measurements.

Values of pH during the titration were measured with a combination pH electrode (Ingold 6030-02, Mettler-Toledo, Columbus, OH) and a Radiometer PHM 95-pH meter (Radiometer Analytical, Lyon, France). Measurements were made before and after data acquisition, with the latter value being assigned to the spectra. Estimated errors in pH measurements are 0.10. Uncorrected pH meter readings are reported based on the assumption that the electrode isotope effect is canceled by the isotope effect on the histidinium (Glasoe and Long, 1960; Li et al., 1961). Titrations were carried out with NaOD and DCl (Isotec, Miamisburg, OH). HDO saturation was complete for the ppm range of interest. Ce<sup>1</sup>H signals from each of the four histidine residues were assigned unambiguously in 1D spectra of wild-type SNase (Alexandrescu et al., 1988).

The Ce<sup>1</sup>H resonances obtained from 1D <sup>1</sup>H-NMR were used to monitor the ionization of individual histidines because they are shifted far down field, they are well resolved, and their chemical shift is quite sensitive to the local environment. Multiple resonances are observed for histidines in SNase (Alexandrescu et al., 1989). The dominant resonance is thought to be due to the *cis* form of the peptide bond between Lys-116 and Pro-117 (Evans et al., 1989). This is also the crystallographically observed configuration of this bond (Loll and Lattman, 1989; Hynes and Fox, 1991). Minor resonances are attributed to the *trans* conformation. An additional minor resonance has been identified at high concentrations. This form is concentration dependent, thus it is attributed to a dimeric form or higher order aggregate of SNase (Alexandrescu et al., 1989). It has been shown previously that the titration behavior of the minor resonances of a given histidine are similar (Fox et al., 1986; Alexandrescu et al., 1989). The pK<sub>a</sub> values of histidines in SNase measured by <sup>1</sup>H-NMR spectroscopy are insensitive to protein concentration according to repeated experiments at concentrations ranging from 0.9 to 1.5 mM.

Analysis of 1D data sets was performed with Felix version 97.2 on a Silicon Graphics R10000 workstation. Time-domain data points were fast Fourier-transformed following zero filling and application of a soft sine bell window. Signal peaks were obtained from the Felix peak picking function. The pK<sub>a</sub> values of the histidine Ce<sup>1</sup>H signals were determined by nonlinear least squares fit with (Markley, 1975),

$$\delta(\text{pH}) = \delta_0 + (\delta_+ - \delta_0) \frac{10^{n(\text{pK}_a - \text{pH})}}{1 + 10^{n(\text{pK}_a - \text{pH})}} \quad (1)$$

In this expression, δ<sub>+</sub> and δ<sub>0</sub> refer to the value of the chemical shift (δ) in the acid and basic limits of the transition, and *n* is a phenomenological Hill

coefficient. Error in the calculated pK<sub>a</sub> values is estimated to be less than 0.07 based on three independent experiments.

## pK<sub>a</sub> calculations with semi-macroscopic continuum methods

The single-site ionization method described by Antosiewicz et al. (1994, 1996) was used for all calculations of electrostatic energies and pK<sub>a</sub> values, with some modifications noted below. The University of Houston Brownian Dynamics package (Davis et al., 1991) was used to calculate electrostatic potentials by solution of the linearized form of the Poisson-Boltzmann equation by the method of finite differences. The cluster method of Gilson (1993) was used to calculate ionization energies and mean charges. Polar hydrogen atoms were added to the protein in the neutral state with the HBUILD facility in CHARMm (Accelrys Inc., www.accelrys.com). The position of the hydrogens was energy minimized with 500 steps of steepest descent with CHARMm version 25.3, performed while all heavy atoms were kept static. Hydrogen atoms were placed on OD2 of all Asp, and on the OE2 of all Glu. The tautomeric forms of His were selected from the best fit to the experimental pK<sub>a</sub> values. This placed hydrogen atoms on Ne2 of His-8, on Nδ1 of His-46, on Ne2 of His-121, and on Nδ1 of His-124. Partial charges were taken from the CHARMm polar-hydrogen-only topology file version 21 (CHARMm, Accelrys Inc.; Neria et al., 1996) and atomic radii from the OPLS parameter set (Jorgensen and Tirado-Rives, 1988). The following set of model compound pK<sub>a</sub> values was used in the calculations: C-term, 3.8; Asp Cγ, 4.0; Glu Cδ, 4.4; His Nδ1 or Ne2, 6.3; N-term, 7.5; Tyr OH, 9.6; Lys Nζ, 10.4; Arg Cζ, 12.0. In all the calculations, the temperature was 298 K, the Stern layer was set at 2.0 Å, the external dielectric constant was 78.5, and the protein dielectric constant, ε<sub>in</sub>, was a variable. A Richard's probe-accessible surface was used with a probe radius of 1.4 Å.

The water molecules that were included explicitly in the calculations were treated as TIP3 waters (Jorgensen et al., 1983). Water hydrogen atoms were added with HBUILD. To explore the effects of water reorientation in response to ionization of the buried groups, the position of water hydrogen atoms was relaxed by minimization with the buried group in the charged state. All minimization procedures consisted of 500 steps of steepest descent with CHARMm v25.3. The minimization was carried out in two steps. In the first minimization step, all ionizable groups were kept in the neutral state except for the buried residue of interest, while the hydrogen atoms of protein polar atoms and of buried water molecules were minimized. In the second minimization step, the hydrogen atoms in protein polar atoms were allowed to relax while the buried ionizable residue was in the neutral state, and the water molecules were fixed in the positions that resulted from the first minimization. To investigate further the effects of water relaxation, additional conformations were obtained from minimization as described above, in which the water oxygen atoms were also allowed to reposition. The role of individual water molecules was explored in two ways. First, by progressively removing individual water molecules from an energy-minimized conformation and assessing the consequences in the calculated pK<sub>a</sub> values. Second, by minimizing conformations after removal of individual water molecules, and then evaluating the consequences on the calculated pK<sub>a</sub> values.

The calculations were performed with the structures of PHS/V66E and Δ+PHS/V66K determined previously in this laboratory. The only water molecules that were treated explicitly in the calculations with SNase are the four water molecules that connect the buried Glu-66 with bulk water (Dwyer et al., 2000). To estimate the effects of buried water on the pK<sub>a</sub> of Lys-66, position 66 in the PHS/V66E structure was mutated into a lysine with the InsightII software package (Accelrys Inc.). The resulting structure was energy minimized with a procedure that only allowed relaxation of Cγ to Nζ atoms of Lys-66 and of oxygen and hydrogen atoms of the four water molecules.



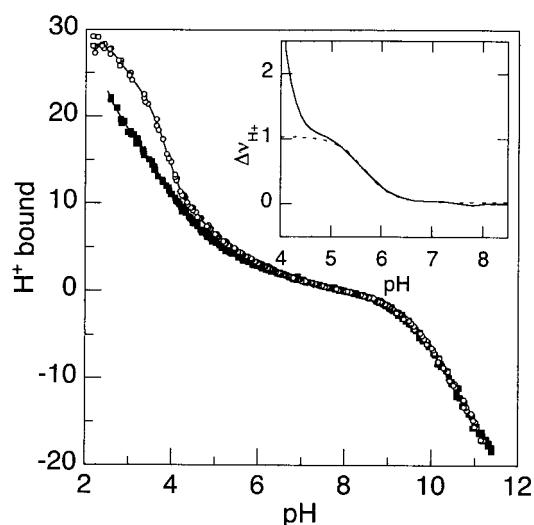


FIGURE 1 Potentiometric  $H^+$  titration of  $\Delta$ +PHS nuclease (filled squares) and its  $\Delta$ +PHS/V66K mutant (open circles) in 100 mM KCl, 25°C. The solid lines through the experimental data represent fits with eighth-order polynomials. The insert includes the difference between these curves (solid line) and the fitted isotherm of a group with a  $pK_a$  of 5.7 (dashed line).

## RESULTS AND DISCUSSION

### $pK_a$ value of the buried Lys-66

Previously, we reported a  $pK_a$  of 6.35 for Lys-66 in wild type SNase and in PHS nuclease, a hyperstable variant that is 3.4 kcal/mol more stable than the wild type (García-Moreno E. et al., 1997). The ionization of the buried Lys-66, even in the hyperstable PHS background, partially unfolds the protein. For this reason the  $pK_a$  of Lys-66 was re-examined in the  $\Delta$ +PHS background, which, at 25°C, is 6.7 kcal/mol more stable than wild type. A preliminary  $pK_a$  of Lys-66 in  $\Delta$ +PHS/V66K was published previously (García-Moreno E. et al., 1997). However, the  $pK_a$  values of these buried residues are measured by indirect thermodynamic methods rather than by a site-specific method based on NMR spectroscopy, thus it was essential to demonstrate that the  $pK_a$  value of Lys-66 in  $\Delta$ +PHS/V66K was accurate. This was accomplished by measuring  $pK_a$  values by two completely independent equilibrium thermodynamic methods. First it was obtained from the difference in potentiometric  $H^+$  titration of  $\Delta$ +PHS nuclease and its V66K mutant ( $\Delta$ +PHS/V66K). Then it was obtained by linkage analysis of the pH dependence of stability of these two proteins. These data were also necessary to assess the effects of the V66K mutation and of the ionization of the buried Lys-66 on the structure and stability of  $\Delta$ +PHS/V66K.

The potentiometric  $H^+$  titration curves of  $\Delta$ +PHS and  $\Delta$ +PHS/V66K are shown in Fig. 1. The two titration curves are superimposed arbitrarily at pH 9; the difference between them ( $\Delta\nu_{H^+}$ ) is shown in the insert. Between pH 6.5 and

11.5 the titration curves are identical. This demonstrates that the V66K mutation had no structural consequences that affected the  $pK_a$  values of surface groups that titrate in this pH range, where the buried Lys-66 is neutral. The two titration curves diverged at pH < 6.5 because of the ionization of the buried Lys-66.

The  $pK_a$  of Lys-66 extracted by fitting a single-site isotherm to the  $\Delta\nu_{H^+}$  curve in Fig. 1 is 5.6 (5.52, 5.66). This  $pK_a$  is considerably lower than the normal  $pK_a$  of 10.4 of lysine in water (Matthew et al., 1985). It is also lower than the value of 6.38 measured previously in the wild type and in the PHS background (García-Moreno E. et al., 1997). The fit of the isotherm to the difference titration curve is excellent at pH > 4.5. Below this pH value the isotherm does not fit the experimental data because the difference  $H^+$  binding curve also reflects proton uptake concomitant with the acid unfolding of  $\Delta$ +PHS/V66K.

The  $H^+$  titration behavior of  $\Delta$ +PHS/V66K was also measured potentiometrically under unfolding conditions (6 M GdnHCl) to assess further the ionization properties of Lys-66 and its effect on the global stability of  $\Delta$ +PHS nuclease. The  $H^+$  titration curves in native and denatured states are compared in Fig. 2 A. The distance between these two curves was determined explicitly with batch experiments, which measured the number of  $H^+$  bound by SNase when GdnHCl was added to a solution of SNase in 100 mM KCl to a final concentration of 6 M. The difference (unfolded minus native) between these  $H^+$  binding curves is shown in Fig. 2 B. This  $\Delta\nu_{H^+}$  curve has three notable features. First, the convergence toward  $\Delta\nu_{H^+} = 0$  above pH 9 shows that, in the unfolded state, Lys-66 ionized with a  $pK_a$  close to 10, in agreement with the known  $pK_a$  of 10.4 of lysine in water. Second, the value of  $\Delta\nu_{H^+} \approx 1$  between pH 9.5 and 6.3 shows that, in this pH range, there is one extra  $H^+$  bound to the unfolded protein that is not bound to the folded protein. This is consistent with the difference between the  $pK_a > 10$  for Lys-66 in the unfolded state, and the  $pK_a$  of 5.6 in the native state. Third, the drop in  $\Delta\nu_{H^+}$  toward 0 that occurred at pH < 6.3 is consistent with the  $pK_a$  of 5.6 for Lys-66 in the native state obtained from the data in Fig. 1. The  $pK_a$  of Lys-66 in the native state could not be resolved from the  $\Delta\nu_{H^+}$  curve in Fig. 2 B because, below pH 5, preferential  $H^+$  uptake by the GdnHCl-unfolded state relative to the native state masked the  $H^+$  binding reaction of Lys-66.

The second method used to determine the  $pK_a$  of Lys-66 involved linkage analysis of the pH dependence of stability of  $\Delta$ +PHS nuclease and the  $\Delta$ +PHS/V66K mutant. The free energy differences between native and denatured SNase ( $\Delta G_{H_2O}^0$ ) measured by GdnHCl denaturation monitored by fluorescence are shown in Fig. 3.  $\Delta$ +PHS is more than twice as stable as wild-type SNase at neutral pH. The stability of  $\Delta$ +PHS was not pH sensitive in the range 4.5–9.5, but it decreased at pH < 4.5. In contrast, the stability of  $\Delta$ +PHS/V66K decreased markedly between pH

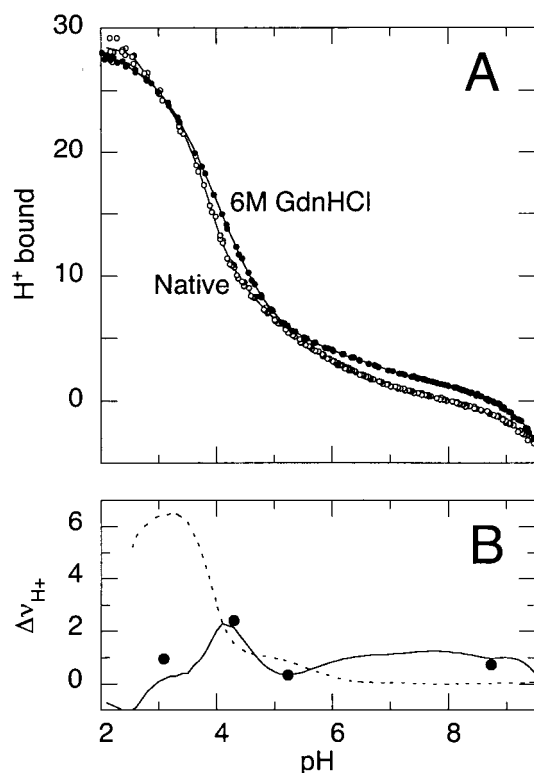


FIGURE 2 (A) Potentiometric H<sup>+</sup> titrations of Δ+PHS/V66K under native conditions (*open circles*) and in 6 M GdnHCl (*filled circles*) at 25°C. (B) Preferential H<sup>+</sup> binding of GdnHCl-denatured protein obtained as the difference in the continuous titration curves in (A) (*solid line*), and with batch experiments (*filled circles*). The Δν<sub>H<sup>+</sup></sub> obtained by batch experiments represent the average of three independent experiments. The error is smaller than the diameter of the symbols. The difference in the H<sup>+</sup> titration curves in Fig. 1 is included for comparison (*dashed line*).

9.5 and 4. According to the titration curves in Fig. 1, the H<sup>+</sup> binding properties of these two proteins were identical between pH 11.5 and 6.5. Therefore, the loss of stability of Δ+PHS/V66K in this range of pH must reflect the depressed pK<sub>a</sub> of Lys-66 when it is buried in the hydrophobic core of the folded protein. The difference between the two stability curves (ΔΔG<sub>H<sub>2</sub>O</sub><sup>o</sup>) is also shown in Fig. 3. In the linear region, the slope of the difference stability curve was 1.25 kcal/mol, close to the theoretical value of 1.36 kcal/mol, equivalent to ΔpK<sub>a</sub> = 1 at 25°C.

The pK<sub>a</sub> values of Lys-66 in the native and unfolded states were obtained by fitting the ΔΔG<sub>H<sub>2</sub>O</sub><sup>o</sup> versus pH curve in Fig. 3 with (Stites et al., 1991)

$$\Delta\Delta G_{\text{H}_2\text{O}}^{\text{o}}(\text{pH}) = \Delta\Delta G_{\text{H}_2\text{O}}^{\text{o}}(\text{mut}) - RT \ln \frac{1 + e^{2.303(\text{pK}_a^{\text{D}} - \text{pH})}}{1 + e^{2.303(\text{pK}_a^{\text{N}} - \text{pH})}} \quad (2)$$

The rightmost term describes the pH-dependent component of ΔΔG<sub>H<sub>2</sub>O</sub><sup>o</sup>. The pH-independent term, ΔΔG<sub>H<sub>2</sub>O</sub><sup>o</sup>(mut), reflects the energetic consequences of the mutation other

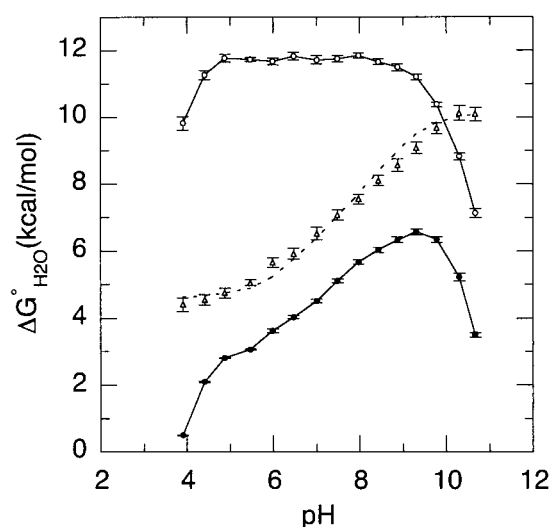


FIGURE 3 Stability of Δ+PHS (*open circles*) and Δ+PHS/V66K (*filled circles*) nuclease determined by fluorescence-monitored GdnHCl denaturation at 25°C. The solid lines through the data are only meant to guide the eye. The difference between these curves (*open triangles*) and the fit of Eq. 1 (*dashed line*) are shown. For display purposes the ΔΔG<sub>H<sub>2</sub>O</sub><sup>o</sup> curve was shifted along the ordinate; information about the pK<sub>a</sub> of Lys-66 lies in the shape, not in the absolute magnitude of this curve.

than the electrostatic effects related to shifts in pK<sub>a</sub> values. Fitting with Eq. 2 yields a denatured state pK<sub>a</sub> value (pK<sub>a</sub><sup>D</sup>) of 9.76 (9.26, 10.32), slightly lower than the pK<sub>a</sub> value of 10.4 in water (Matthew et al., 1985). The pK<sub>a</sub> value in the native state (pK<sub>a</sub><sup>N</sup>) is 5.76 (5.52, 5.96), in excellent agreement with the pK<sub>a</sub> value that was determined potentiometrically.

Analysis of the Δν<sub>H<sup>+</sup></sub> versus pH curve in Fig. 1 with a single-site H<sup>+</sup> binding isotherm, and of the ΔΔG<sub>H<sub>2</sub>O</sub><sup>o</sup> versus pH in Fig. 3 with Eq. 1, both assume that the V66K mutation does not affect the ionization properties of other groups. This assumption is validated by the data in Fig. 1, and by the slope of ΔΔG<sub>H<sub>2</sub>O</sub><sup>o</sup> versus pH in Fig. 3, which suggests that the dependence of stability on pH is determined primarily by the shift in the pK<sub>a</sub> of a single group. The excellent agreement between pK<sub>a</sub> values determined by independent thermodynamic methods also supports the validity of this assumption.

### Ionization of the buried Lys-66 does not unfold nuclease

To understand the mechanisms of dielectric relaxation upon ionization of the buried groups, it is necessary to characterize the effects of the titration of the buried group on the conformation of the protein. Previously, it was demonstrated that PHS/V66D is mostly folded after ionization of the buried Glu-66 (Dwyer et al., 2000). However, the ionization of Lys-66 in the wild type and in PHS nuclease disrupts the native conformation and leads to measurable

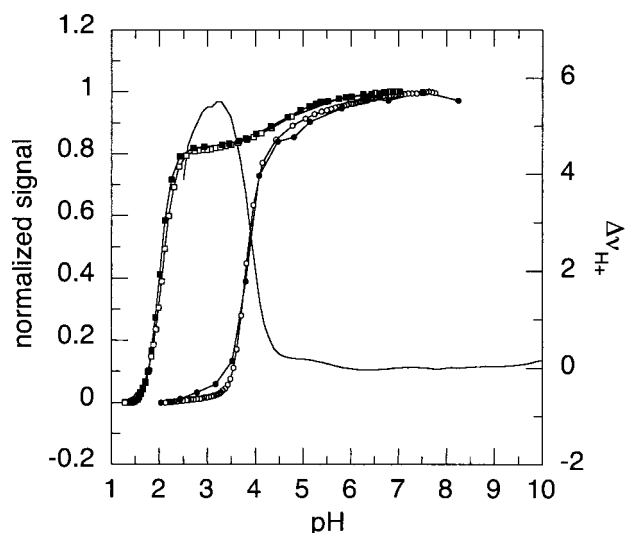


FIGURE 4 Acid/base titrations of  $\Delta$ +PHS (squares) and  $\Delta$ +PHS/V66K (circles) nucleases monitored by fluorescence (filled symbols) and far-ultraviolet CD (open symbols). The lines through the data are only meant to guide the eye. Superimposed, with reference to the right axis, is the difference  $H^+$  binding curve from Fig. 1 (thin line) from which the titration of a group with a  $pK_a$  of 5.7 has been subtracted.

global or partial unfolding (Stites et al., 1991; García-Moreno E. et al., 1997). Thus, it is important to assess the conformational state of  $\Delta$ +PHS/V66K SNase under conditions of pH where Lys-66 is charged.

Five thermodynamic and spectroscopic observations demonstrate conclusively that, in the hyperstable  $\Delta$ +PHS background, the ionization of the buried Lys-66 does not unfold the protein or perturb its global conformation in a detectable manner.

1. The uptake of  $H^+$  that signals the onset of acid denaturation occurs at  $pH < 4.2$ , more than 1 pH unit below the  $pK_a$  of 5.7 for Lys-66 (insert in Fig. 1). This is also evident in the  $\Delta H^+$  curve in the lower panel in Fig. 2. The increase in  $\Delta \nu_{H^+}$  at pH 5.1 and below reflects greater  $H^+$  binding to  $\Delta$ +PHS/V66K in 6 M GdnHCl than in 100 mM KCl. This implies that the  $pK_a$  values of the acid groups in  $\Delta$ +PHS/V66K are still depressed and native-like after the ionization of the buried Lys-66. The maximum in the  $\Delta \nu_{H^+}$  curve near pH 4.2 identifies the pH at which the acid unfolding transition begins.
2. The stability of  $\Delta$ +PHS/V66K is approximately 3 kcal/mol at the pH corresponding to the  $pK_a$  of Lys-66 (Fig. 3). Therefore, at this pH, the protein is predominantly in the native state. Even at pH 4.4, more than a pH unit below the  $pK_a$  of Lys-66,  $\Delta G_{H_2O}^o$  is still greater than 2 kcal/mol. The steep change in  $\Delta G_{H_2O}^o$  below this pH signals the onset of acid-denaturation.
3. The acid/base titrations of  $\Delta$ +PHS/V66K monitored by far-ultraviolet CD and by the intrinsic fluorescence of Trp-141 are shown in Fig. 4. The fluorescence-monitored titration of  $\Delta$ +PHS is also included in this figure for comparison. Acid denaturation of  $\Delta$ +PHS nuclease began at pH 2.5. The cooperative acid-unfolding transition of the  $\Delta$ +PHS/V66K mutant occurred at  $pH \leq 4.2$ , more than 1.5 pH units below the  $pK_a$  of Lys-66. Figure 4 also illustrates the excellent agreement between the acid-unfolding transition monitored by CD, by fluorescence, and by preferential  $H^+$  binding to the acid denatured state. The latter curve is obtained by subtracting from the titration curve of  $\Delta$ +PHS/V66K the  $H^+$  titration curves of  $\Delta$ +PHS and the titration curve of a single group with a  $pK_a$  of 5.7. All the different types of titration data show that the major acid-induced conformational transition began near pH 4.2, and ended near pH 3.4.
4. The pH dependence of 1D  $^1H$ -NMR spectra of  $\Delta$ +PHS/V66K is shown in Fig. 5. These spectra confirmed that  $\Delta$ +PHS/V66K exists primarily in the native conformation at  $pH \geq 4.2$ . The acid-induced conformational transition between pH 4.2 and 3.6 can be observed most

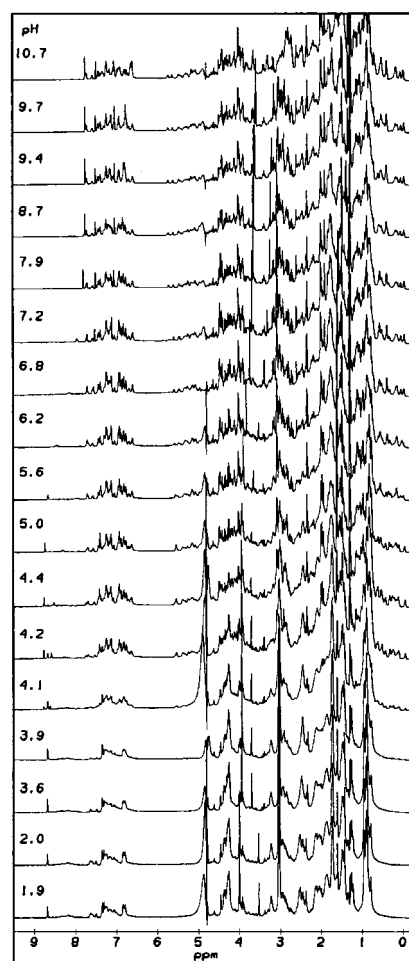


FIGURE 5 Acid/base titration of 1D  $^1H$ -NMR spectra of  $\Delta$ +PHS/V66K at 25°C in 100 mM KCl.

tored titration of  $\Delta$ +PHS is also included in this figure for comparison. Acid denaturation of  $\Delta$ +PHS nuclease began at pH 2.5. The cooperative acid-unfolding transition of the  $\Delta$ +PHS/V66K mutant occurred at  $pH \leq 4.2$ , more than 1.5 pH units below the  $pK_a$  of Lys-66. Figure 4 also illustrates the excellent agreement between the acid-unfolding transition monitored by CD, by fluorescence, and by preferential  $H^+$  binding to the acid denatured state. The latter curve is obtained by subtracting from the titration curve of  $\Delta$ +PHS/V66K the  $H^+$  titration curves of  $\Delta$ +PHS and the titration curve of a single group with a  $pK_a$  of 5.7. All the different types of titration data show that the major acid-induced conformational transition began near pH 4.2, and ended near pH 3.4.

**TABLE 1** pK<sub>a</sub> Values of surface histidines and buried Lys-66 and Glu-66

	His-8*	His-46	His-121	His-124	Lys-66 or Glu-66 <sup>†</sup>	Lys-66 or Glu-66 <sup>‡</sup>
Wild type <sup>§</sup>	6.61 ± 0.07	5.94 ± 0.02	5.29 ± 0.04	5.81 ± 0.07		
PHS	6.53 ± 0.02	5.90 ± 0.02	5.31 ± 0.02			
Δ+PHS	6.50 ± 0.02		5.25 ± 0.04			
PHS/V66K	6.51 ± 0.05	6.38 ± 0.03	6.49 ± 0.17			6.35 ± 0.10 <sup>¶</sup>
Δ+PHS/V66K	6.56 ± 0.01		5.38 ± 0.01		5.6 ± 0.01	5.76 ± 0.30
PHS/V66E					8.8 ± 0.10 <sup>  </sup>	8.5 ± 0.30 <sup>  </sup>

\*Values determined by <sup>1</sup>H-NMR.<sup>†</sup>Values determined potentiometrically.<sup>‡</sup>Values determined from the pH dependence of ΔG<sub>H<sub>2</sub>O</sub>.<sup>§</sup>Errors are standard deviations from three separate experiments. All other errors in NMR-measured pK<sub>a</sub> values are errors of the fit.<sup>¶</sup>García-Moreno et al., 1997.<sup>||</sup>Dwyer et al., 2000.

dramatically in the pH-induced changes in the resonances in the aliphatic region (0–1 ppm) and in the aromatic region (6–9 ppm). The latter resonances, corresponding to the histidine residues, coalesced into a single resonance near 8.7 ppm at pH 3.6, suggesting that, in the acid-unfolded state, the histidine residues sampled, on average, identical microenvironments. The titration curves obtained by plotting the area under the upfield resonances of Val-74 against pH also paralleled the acid/base titration monitored spectroscopically or potentiometrically (data not shown).

- Finally, the pK<sub>a</sub> values of the surface histidines His-8 and His-121 are very similar in Δ+PHS, in Δ+PHS/V66K, and in the wild-type protein (see Table 1). This suggests that Δ+PHS/V66K is native-like down to pH 4. In contrast, the pK<sub>a</sub> values of histidines in PHS/V66K are the same as the values for histidine in water, as expected, because the titration of the buried Lys-66 induces a substantial conformational transition (García-Moreno E. et al., 1997).

The experimental observations demonstrate unequivocally that the Δ+PHS/V66K mutant exists primarily in the fully folded, native conformation at pH values more than a full pH unit below the pK<sub>a</sub> of Lys-66. This does not imply that local rearrangement or minor conformational changes do not occur upon ionization of the buried residues, but if structural changes do occur, they are invisible to the spectroscopic probes that were used. Note that, even if local rearrangements take place, the interpretation of ΔpK<sub>a</sub> values in terms of dielectric properties of the protein is meaningful. The ΔpK<sub>a</sub> can still be used to determine the highest value of ε<sub>in</sub> needed to reproduce the pK<sub>a</sub> values with a given computational method.

### Surface ionizable groups do not sense the positive charge of Lys-66

Electrostatic interactions between surface ionizable groups and the buried Lys-66 could significantly influence the pK<sub>a</sub>

value of the buried group. Were this the case, some surface basic and acidic residues should exhibit more depressed pK<sub>a</sub> values in the Δ+PHS/V66K mutant than in the Δ+PHS protein. The titration curves show no evidence of this. Once corrected to account for the potentiometric H<sup>+</sup> titration curve of Lys-66, the H<sup>+</sup> titration properties of Δ+PHS and Δ+PHS/V66K were nearly identical down to pH 4.2, where the acid unfolding transition of Δ+PHS/V66K begins (See Figs. 1 and 4). According to structure-based pK<sub>a</sub> calculations discussed ahead, the pK<sub>a</sub> of His-121 would have been depressed significantly in Δ+PHS/V66K if the electrostatic interactions between Lys-66 and His-121 were strong. Instead, the pK<sub>a</sub> values of histidines listed in Table 1 show that His-121 titrated with almost identical pK<sub>a</sub> values in Δ+PHS and in Δ+PHS/V66K. The slight shift in the pK<sub>a</sub> of His-121 in these two proteins reflects our inability to collect the entire acid limit of the titration curve of this histidine because of the acid denaturation of the protein. The chemical shifts of His-121 in Δ+PHS/V66K and in Δ+PHS are identical down to the lowest pH in which the protein exists mainly in the native state. Overall, the data suggest that the electrostatic interactions between the buried Lys-66 and surface ionizable residues are negligible.

### Meaning of dielectric constants in continuum models

It is of interest to identify the values of the protein dielectric constant, ε<sub>in</sub>, needed to reproduce pK<sub>a</sub> values quantitatively with a continuum model for structure-based pK<sub>a</sub> calculations. To appreciate the significance of the values of ε<sub>in</sub> thus resolved, it is necessary to review the meaning of ε<sub>in</sub>.

The semi-macroscopic continuum model is based on the macroscopic model of Warwicker and Watson (1982), which treats the protein interior as a medium with low dielectric constant. However, this model also considers the microscopic features of the protein permanent dipoles (Warshel and Russell, 1984) by representing them explicitly in terms of partial charges (Klapper et al., 1986; Warshel et



al., 1989; Bashford and Karplus, 1990). When this model is used with a static structure to calculate the  $pK_a$  values of surface groups, ad hoc use of  $\epsilon_{in} \approx 20$  is required to reproduce  $pK_a$  values (Antosiewicz et al., 1994). However, as demonstrated in the present study, high values of  $\epsilon_{in}$  grossly underestimate the self-energies of buried groups; considerably lower values of  $\epsilon_{in}$  are required to capture the  $pK_a$  values of buried groups consistently (Antosiewicz et al., 1994; Demchuk and Wade, 1996; Sham et al., 1997; Schutz and Warshel, 2001). Self-consistent calculation of  $pK_a$  values of surface and buried groups using a single value of  $\epsilon_{in}$  is impossible with semi-macroscopic methods unless protein flexibility is taken into account explicitly. When conformational relaxation is treated explicitly, low  $\epsilon_{in}$  can yield reasonable results (Langsetmo et al., 1991; Antosiewicz et al., 1994, 1996; You and Bashford, 1995; Alexov and Gunner, 1997; Zhou and Vijayakumar, 1997; Rabenstein et al., 1998; van Vlijmen et al., 1998; Havranek and Harbury, 1999; Scharnagl et al., 1999; Ullmann and Knapp, 1999), especially when the neutral and the ionized states of the group of interest are treated separately to capture contributions by the relaxation of local dipoles upon ionization (Langen et al., 1992; Yang et al., 1993; Alexov and Gunner, 1997; Sham et al., 1997; Zhou and Vijayakumar, 1997; Rabenstein and Knapp, 2001). In microscopic methods, such as the protein dipole–Langevin dipole (PDL) method developed by Warshel and co-workers, the relaxation of protein dipoles has always been included explicitly (Warshel and Russell, 1984). Microscopic models are not yet fully convergent, thus a semi-macroscopic PDL/S model (Warshel et al., 1989) was developed that uses a dielectric constant to account implicitly for energy contributions that do not converge completely in the fully microscopic simulations (Schutz and Warshel, 2001).

The appropriate choice of  $\epsilon_{in}$  in semi-macroscopic continuum models is not obvious, as discussed previously by Warshel (Warshel and Russell, 1984; King et al., 1991). A single value of  $\epsilon_{in}$  is not necessarily appropriate because the protein interior is highly heterogeneous and anisotropic; a single value implies that dielectric relaxation is the same in all proteins, and uniform throughout any one protein. This is unlikely (Demchuk and Wade, 1996; Sham et al., 1997; Simonson et al., 1999; Gunner and Alexov, 2000).  $\epsilon_{in}$  need not be the same in the hydrophobic core, where electronic polarizability makes the dominant contribution, near the surface, where the reaction field from bulk water and flexibility of charged side chains are more dominant, and in intermediate regions, where reorganization of permanent dipoles are important (Simonson and Perahia, 1995; Simonson and Brooks, 1996). Furthermore, different dielectric constants are necessary to describe formally the two different components of the dielectric response upon ionization of a buried group: one to account for the static equilibrium charge distribution, and a higher value to account for relaxation (Krishtalik et al., 1997; Simonson et al., 1999). Sim-

ilarly, different values of  $\epsilon_{in}$  are needed to capture self-energies and coulombic energies (Warshel and Papazyan, 1998), although these could be captured by a single value of  $\epsilon_{in}$  if contributions by protein relaxation were accounted for consistently and explicitly (Sham et al., 1997).

It is also useful to recognize that the value of  $\epsilon_{in}$  used in a calculation depends on the level of physical detail represented explicitly in a model (Schutz and Warshel, 2001). In fully microscopic simulations,  $\epsilon_{in} = 1$  because all contributions to dielectric relaxation in the protein–water system are treated explicitly. In calculations with continuum models  $\epsilon_{in}$  is always  $>1$ . However, in these models  $\epsilon_{in}$  is not strictly a dielectric constant. It is a scaling parameter meant to represent contributions that are not included explicitly in the model (Simonson et al., 1999; Schutz and Warshel, 2001). Implicit treatment of induced dipoles (electronic polarization) requires  $\epsilon_{in} \approx 2$ . When induced dipoles and protein relaxation (nuclear relaxation or reorientation of dipoles) are implicit,  $\epsilon_{in} \approx 4$ – $10$  should be used (Warshel et al., 1997; Rabenstein et al., 1998; Simonson et al., 1999). Even higher values of  $\epsilon_{in} \geq 20$  are necessary to capture correctly  $pK_a$  values of surface residues. These high values of  $\epsilon_{in}$  account implicitly for relaxation of permanent dipoles upon charging (Rabenstein et al., 1998; Sham et al., 1998). The essential point is that protein dielectric constants used in continuum models are not fundamental parameters. They are empirical parameters that need to be calibrated against microscopic simulations (Sham et al., 1998; Warshel and Papazyan, 1998; Simonson et al., 1999), or ideally against experimental data in carefully controlled situations, as in the present study.

### Comparison of measured and calculated $H^+$ titration behavior of surface residues

To determine the range of values of  $\epsilon_{in}$  needed to reproduce  $pK_a$  values, we focus first on the surface residues. The potentiometric  $H^+$  titration curves of  $\Delta$ +PHS nuclease and of the  $\Delta$ +PHS/V66K mutant are compared with the curves calculated with the continuum method in Fig. 6. Calculated curves obtained with  $\epsilon_{in}$  of 4, 10, and 20 are shown. As expected, the agreement between measured and calculated behavior was very poor with  $\epsilon_{in} = 4$ , and satisfactory with  $\epsilon_{in} \geq 10$ . The slight differences between the curves calculated with  $\epsilon_{in} = 20$  and with  $\epsilon_{in} = 10$  in the pH range 5.7–9.0 (Fig. 6B) reflect differences in the state of ionization of the buried Lys-66, which titrates with a high  $pK_a$  when  $\epsilon_{in} = 20$  and with a much lower one when  $\epsilon_{in} = 10$ . Comparison of calculated and measured titration behavior for  $\Delta$ +PHS/V66K is only meaningful at pH  $> 4.2$ . The divergence between the measured and calculated  $H^+$  titration curves at pH  $\leq 4.2$  reflects preferential  $H^+$  binding by the acid denatured form of  $\Delta$ +PHS/V66K in the experimental curve. On the other hand, the divergence between measured and calculated  $H^+$  titration curves of  $\Delta$ +PHS



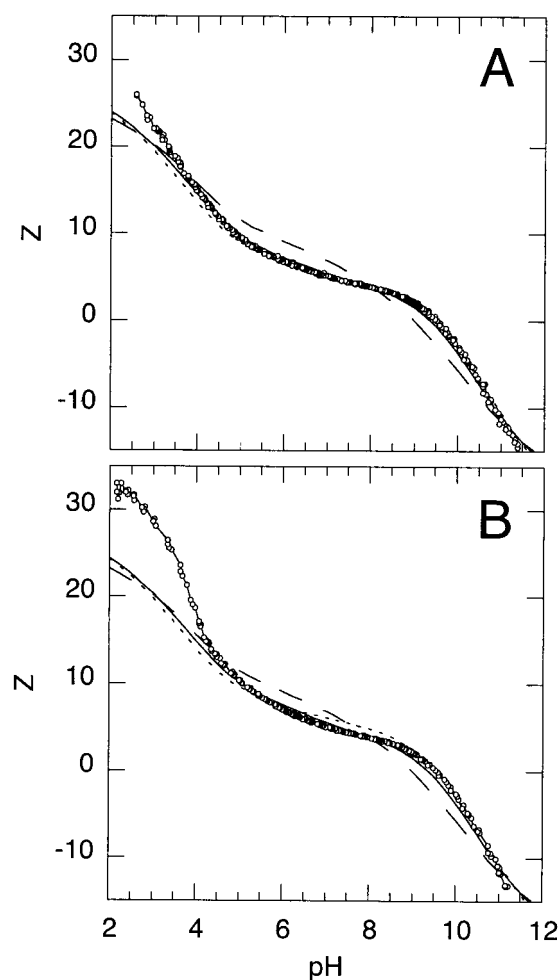


FIGURE 6 (A) pH Dependence of the net charge of  $\Delta$ +PHS nuclease measured potentiometrically (open circles) and calculated using  $\epsilon_{in} = 4$  (dashed line), 10 (solid line), and 20 (dotted line). All data were measured or calculated at 25°C in 100 mM KCl. The experimental data were shifted arbitrarily along the ordinate to superimpose them artificially with the data calculated with  $\epsilon_{in} = 10$ . (B) Same as in (A) but for  $\Delta$ +PHS/V66K nuclease.

nuclease at low pH values is meaningful. It reflects inaccuracies in the calculated pK<sub>a</sub> values of acidic residues not related to inherent problems with the algorithm, but to the use of a static structure under conditions of pH where considerable structural relaxation takes place (C. Fitch, S. Whitten and B. García-Moreno, in preparation).

### Calculated pK<sub>a</sub> values of the buried Lys-66 and Glu-66

Figure 7 shows pK<sub>a</sub> values of Lys-66 and Glu-66 calculated as a function of  $\epsilon_{in}$  with the continuum method using a static structure. The calculated pK<sub>a</sub> values of Lys-66 and Glu-66 were only slightly dependent on  $\epsilon_{in}$  when  $\epsilon_{in} > 20$ , but steeply dependent at lower values of  $\epsilon_{in}$ . Calculations with

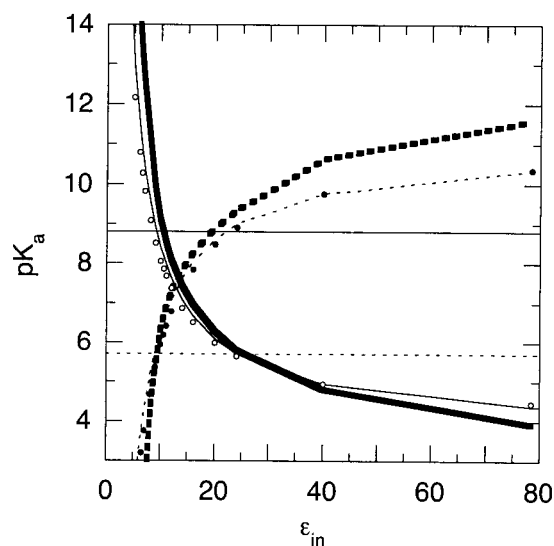


FIGURE 7 pK<sub>a</sub> values of Lys-66 in  $\Delta$ +PHS/V66K (thick dashed) and Glu-66 in PHS/V66E (thick solid), calculated as a function of  $\epsilon_{in}$  with the semi-macroscopic methods based on finite-difference solution of the linearized Poisson–Boltzmann equation. The buried water molecules were not included explicitly in these calculations. Also shown are the Born energy (open and closed circles) calculated with the finite difference calculation, and the Born energy (thin solid and dashed lines) calculated with Eq. 3 in Dwyer et al. (2000). The thin horizontal solid lines identify the experimental pK<sub>a</sub> values of Lys-66 (dashed) and Glu-66 (solid). Calculations and experimental data are at 25°C and 100 mM NaCl.

$\epsilon_{in} = 4$ , a value commonly used to represent the dielectric properties of the protein interior, failed dramatically to capture the experimental pK<sub>a</sub> values. The empirical values of  $\epsilon_{in}$  that best reproduced the experimental pK<sub>a</sub> values of Glu-66 and Lys-66 were 10.5 and 9.5, respectively. Almost identical apparent dielectric constants were obtained by analysis of  $\Delta$ pK<sub>a</sub> with a simple Born formalism (García-Moreno E. et al., 1997; Dwyer et al., 2000), also shown in Fig. 7.

The dominant determinant of the pK<sub>a</sub> values was the desolvation of the charged group when it is buried in the protein. This was established by dissection of the calculated pK<sub>a</sub> values into energetic contributions from coulombic, Born, and background terms (Bashford and Karplus, 1990; Sham et al., 1997). The dependence of the Born energy on  $\epsilon_{in}$  closely parallels that of the pK<sub>a</sub> itself (Fig. 7). In contrast, the contributions to pK<sub>a</sub> values from interactions with polar or surface charged residues were negligible when  $\epsilon_{in} \geq 10$ , consistent with the experimental observation that coulombic interactions between the buried Lys-66 or Glu-66 and surface ionizable residues are weak. The interactions between surface and buried ionizable groups become increasingly significant with  $\epsilon_{in} \leq 10$ . Those with polar atoms remain small even at low values of  $\epsilon_{in} \approx 4$  because the buried charged group makes minimal contact with side chain or backbone polar atoms.

In general,  $\epsilon_{\text{in}} = 20$  is the empirical value of choice that maximizes agreement between calculated and measured  $\text{pK}_a$  values of surface groups (Antosiewicz et al., 1994). The data in Fig. 6 show that this holds true for SNase. However,  $\epsilon_{\text{in}} = 20$  is too high a value to capture the self-energies of Glu-66 and Lys-66 (Fig. 7).  $\epsilon_{\text{in}} \approx 10$  is the empirical dielectric constant that captures  $\text{pK}_a$  values of buried groups at this location in SNase. Coincidentally, in the case of SNase, this value also captures the behavior of surface groups quantitatively (see Fig. 6). The empirical value  $\epsilon_{\text{in}} = 10$  is unlikely to be of general use for estimation of  $\text{pK}_a$  values of buried groups. This high apparent polarizability probably reflects, at least partly, contributions by water penetration to the dielectric response of the protein that were not treated explicitly in the calculations, and it remains to be established that water penetration is a general mechanism of dielectric relaxation in proteins. Experimental studies are underway to determine empirically the values of  $\epsilon_{\text{in}}$  required by this continuum method to capture the experimental  $\text{pK}_a$  values of groups buried at other locations in SNase and in other proteins.

### Influence of internal, site-bound water on $\text{pK}_a$ values of the buried Glu-66

To evaluate contributions by buried water molecules to the  $\text{pK}_a$  value of Glu-66, a continuum calculation was performed in which some of the internal water molecules were treated explicitly, as done previously by others (Langen et al., 1992; Yang et al., 1993; Gibas and Subramaniam, 1996; Alexov and Gunner, 1999; Schutz and Warshel, 2001). Only four waters were treated explicitly, the two that are hydrogen bonded directly to the carboxyl oxygen atoms of Glu-66, and the two that connect these with bulk water (Dwyer et al., 2000). Results from these calculations are presented in Fig. 8. The buried water molecules had no effect on  $\text{pK}_a$  values when the system was energy minimized with Glu-66 in the neutral state. In contrast, when the minimization was performed with Glu-66 in the charged state, the buried water molecules had a considerable effect on the  $\text{pK}_a$  value. In these calculations, the experimental  $\text{pK}_a$  values were reproduced with  $\epsilon_{\text{in}} = 6.4$ . However, with  $\epsilon_{\text{in}}$  in this range, the calculated coulombic interactions between the buried Glu-66 and surface-charged residues are exaggerated and too large to be consistent with the experimental observations. When the coulombic interactions between buried and surface ionizable residues were artificially turned off, the experimental  $\text{pK}_a$  of Glu-66 was reproduced with  $\epsilon_{\text{in}} = 4.3$ . When, in addition to the repositioning of hydrogen atoms, the water oxygen atoms were also allowed to reposition during the minimization, the  $\text{pK}_a$  values were reproduced with  $\epsilon_{\text{in}} = 5.0$ .

The pattern of hydration of Glu-66 present in crystallographic structures at  $-178^\circ\text{C}$  need not represent the state of hydration in solution at  $25^\circ\text{C}$ . To assess how the value of  $\epsilon_{\text{in}}$

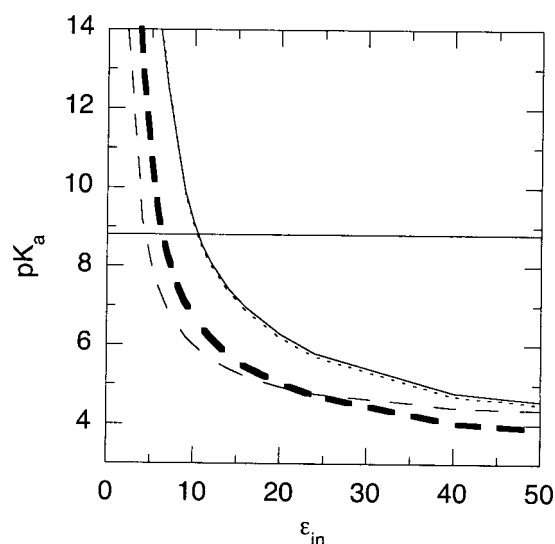


FIGURE 8  $\text{pK}_a$  Values of Glu-66 in PHS/V66E calculated as a function of  $\epsilon_{\text{in}}$  with buried water molecules omitted (*solid*) and with buried water molecules included after energy minimization of hydrogens with Glu-66 in the neutral state (*dotted*) or in the charged state and with coulombic interactions included (*thick dash*), or excluded (*thin dash*). The horizontal line identifies the experimental  $\text{pK}_a$  of Glu-66. Calculations and experimental data are at  $25^\circ\text{C}$  and 100 mM NaCl.

that reproduced  $\text{pK}_a$  values depended on the number of water molecules that were treated explicitly, calculations were also performed with only a subset of the explicit internal water molecules included. The positions and orientation of the internal water molecules after minimization were largely insensitive to the number of crystallographic water molecules included in the minimization procedure. When only the most deeply buried water molecule was treated explicitly, and when coulombic interactions were omitted,  $\epsilon_{\text{in}} = 7.4$  reproduced the experimental  $\text{pK}_a$  value. When the two most deeply buried water molecules that are directly hydrogen bonded to the carboxyl group of Glu-66 were included, the  $\text{pK}_a$  values were reproduced with  $\epsilon_{\text{in}} = 5.4$ . These calculations have two important implications. They show that the two innermost water molecules directly hydrogen bonded to the carboxyl are the ones responsible for the dramatic effect of internal water on the  $\text{pK}_a$  value. The value of  $\epsilon_{\text{in}} = 7.4$ , or more exactly, the fact that  $\epsilon_{\text{in}} > 4$ , further suggests that there might be structural reorganization concomitant with ionization of Glu-66, that is not captured implicitly in the calculations with a static structure. This was also suggested by previous calculations with semi-microscopic methods (Schutz and Warshel, 2001).

The data in Figs. 8 and 9 provide a clear example of how it is impossible to capture the behavior of surface and buried groups with a single value of  $\epsilon_{\text{in}}$  in calculations with semi-macroscopic methods using a static structure. The problem is inherent to calculations with continuum methods, in which contributions by protein relaxation are subsumed in

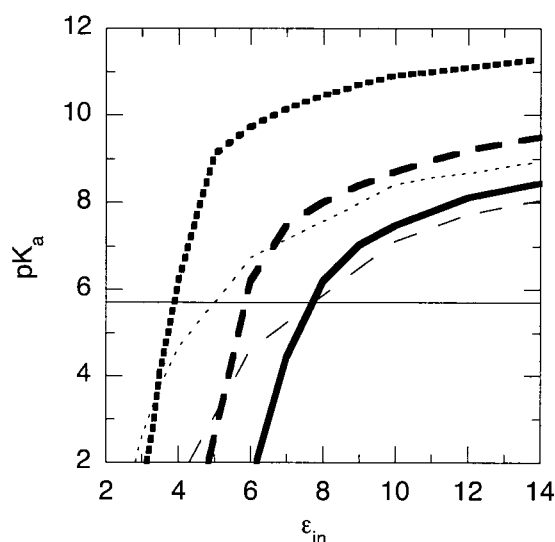


FIGURE 9 pK<sub>a</sub> values of Lys-66 in PHS/V66K calculated as a function of  $\epsilon_{in}$  with buried water molecules omitted (*thick solid*), with four buried water molecules included after energy minimization (allowing for movement of water molecules) with Lys-66 in the charged state, and with coulombic interactions included (*thick dotted*), or excluded (*thin dotted*). Also shown are results of calculations with only the two innermost buried waters included after minimization with Lys-66 in the charged state and coulombic energies included (*thick dashed*), or excluded (*thin dashed*). The horizontal line identifies the experimental pK<sub>a</sub> of Lys-66. Calculations and experimental data are at 25°C and 100 mM NaCl.

$\epsilon_{in}$ , which need be neither uniform throughout the protein, nor the same for calculations of self-energies and coulombic energies. The problem can be avoided if, instead of using a single conformation in the calculations, the microscopic reorganization of the specific environments around the charges is accounted for explicitly. This can be done by the linear response approximation formulation of Warshel and coworkers, which requires averaging over the charge and uncharged configurations (Langen et al., 1992), and by a similar treatment in the more recent methods (Alexov and Gunner, 1997; van Vlijmen et al., 1998; Schutz and Warshel, 2001). These problems are also avoided in the method developed by Mehler and co-workers, which is based on the use of screened coulombic potentials coupled with a hydrophobicity parameter to account for variations in local microenvironments of ionizable groups (Mehler and Guarnieri, 1999).

### Test of the hypothesis that Lys-66 in the buried state is hydrated

The only noteworthy difference between the crystallographic structures with Lys-66 and Glu-66 are the buried water molecules hydrogen bonded to the carboxyl of Glu-66. They are not present in two structures of V66K mutants of wild type and  $\Delta$ +PHS nuclease, obtained at 25°C and

−178°C, respectively. The absence of buried water molecules in the V66K structures is surprising because the  $\Delta pK_a$  for Glu-66 and Lys-66 are virtually identical—these two groups experience equivalent net polarizability in the ionized state. The side chain of the buried Lys-66 is in an extremely hydrophobic environment, incompatible with the high apparent polarizability reflected in the  $\Delta pK_a$ .

With the experimental data at hand, it is not possible to exclude the possibility that ionization of Glu-66 or Lys-66 is accompanied by local unfolding or minor restructuring of their local environment. Close examination of the structures and side chain rotamers suggests that side-chain extrusion toward bulk water would be very difficult in the case of Glu-66 except through local unfolding. The longer side chain of Lys-66 could reach the surface of the protein, but only by assuming a highly unfavorable conformation. In energy-minimization and molecular-dynamics simulations, the buried Lys-66 never abandons its buried position (Stites et al., 1991). In the absence of any direct, experimental evidence for structural conformational changes concomitant with ionization of the buried Lys-66, we entertain the hypothesis that the high polarizability reported by the pK<sub>a</sub> of this group reflects contact with internal water similar to those seen in the structures with Glu-66. These buried waters could be transient, or disordered, and therefore crystallographically invisible. This is partly expected based on the observation that amines are weakly hydrated, especially compared to the strong hydration of carboxylic groups (Collins, 1997).

To determine how the pK<sub>a</sub> of Lys-66 would be affected if it were in contact with internal water molecules, we performed calculations with a model of the structure of PHS/V66K, made by introducing an Lys at position 66 into the PHS/V66E structure with the conformation that Lys-66 has in  $\Delta$ +PHS/V66K. The results are presented in Fig. 9. When the orientation of the buried water molecules was energy minimized with Lys-66 in the charged state, and when the coulombic interactions between buried and surface ionizable residues were included, the experimental pK<sub>a</sub> value was reproduced with  $\epsilon_{in} = 3.9$ . When the coulombic interactions were omitted,  $\epsilon_{in} = 5.0$  reproduced the pK<sub>a</sub> values. Again, the two most deeply buried water molecules were the ones that influenced the pK<sub>a</sub> of Lys-66 the most. The value of  $\epsilon_{in}$  needed to reproduce the experimental pK<sub>a</sub> increased toward 7.5 when only one of the innermost water molecules was included in the calculation.

Detailed views of the buried side chains of Glu-66 and Lys-66 and the positions and orientation of the water molecules minimized with the neutral and charged form of the buried group are shown in Fig. 10. The hydrogen-bonding patterns between the buried water molecules, the buried side chain, and backbone polar atoms are substantially different in the cases of Glu-66 and Lys-66. Despite these differences, the pK<sub>a</sub> values of Glu-66 and Lys-66 were captured with nearly identical values of  $\epsilon_{in}$  in the

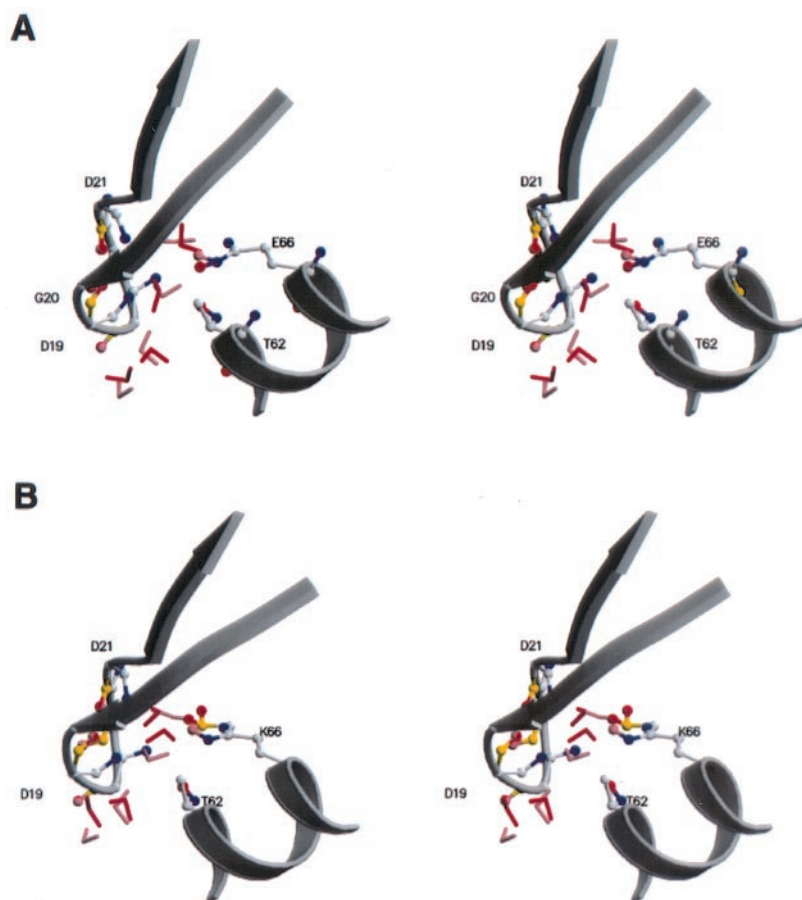


FIGURE 10 Microenvironments around (A) Glu-66 in PHS/V66E. Water molecules and some hydrogen ions are shown after minimization with Glu-66 in the neutral (*pink*) and charged (*red*) states. (B) Lys-66 in the artificial PHS/V66K with the side chain of Glu-66 superimposed for comparison. Shown are the position and orientation of water molecules after minimization with Glu-66 in the neutral state (*pink*), or with Lys-66 in the charged state (*red*). Images drawn with MOLSCRIPT and with Raster 3D (Kraulis, 1991; Merritt and Bacon, 1997).

range 4–5 when the water molecules were allowed to relax. This figure shows that, in the Lys-66 mutant, there is a sufficiently large volume to accommodate the water molecules. However, the pentagonal structures formed by the carboxylic moiety of Glu-66, the backbone polar atoms, and the water molecules, cannot be formed with Lys-66. This might be another reason why the internal waters are more disordered in the presence of Lys-66, and therefore crystallographically invisible.

In calculations with low values of  $\epsilon_{in}$ , the details of the calculations matter greatly: the manner in which explicit water molecules are treated, the specific procedure for energy minimization, and the choice of charge set used in the calculations to represent permanent dipoles can have significant effects. Even for surface residues, the agreement between experimental and calculated  $pK_a$  values in calculations with  $\epsilon_{in} \approx 4$  can be improved simply by redistributing charge on the ionizable side chain (Antosiewicz et al., 1996). More sophisticated treatment of multiple conformers and different placement of hydrogen

atoms in the region around the buried residues is warranted when  $\epsilon_{in} \approx 4$  (Alexov and Gunner, 1997). These issues will be addressed in the future, when additional experimental  $pK_a$  values of buried residues are available to critically assess the performance of the computational models. However, the conclusions from the calculations in Figs. 8 and 9, regarding the important role of buried water molecules in determining  $pK_a$  values of buried groups, are independent of the details of the calculations. They demonstrate that internal water molecules can contribute significantly to the high apparent polarizability reported by Glu-66. They support the hypothesis that the buried Lys-66 interacts with internal waters. They illustrate the profound effects that water penetration can have on the electrostatic properties of ionizable groups buried in the protein interior or at interfaces between macromolecules. One or two buried water molecules in the vicinity of a buried Lys or Glu can modulate its  $pK_a$  within the physiologically relevant range of pH. According to these calculations, internal water molecules in proteins can be



only slightly less effective than bulk water in solvating buried groups and in stabilizing charge.

### Structural relaxation upon charging of the buried ionizable residues

To interpret the meaning of the  $\epsilon_{\text{in}}$  values obtained from the different simulations shown in Figs. 8 and 9, it is useful to remember that the equilibrium dielectric constants of dried proteins and peptides measured experimentally are in the range 2–4 (Harvey and Hoekstra, 1972). In the limit of extreme dehydration, dried proteins are presumably still folded in native-like conformations, rigid, uncharged, with all polar groups hydrogen bonded among themselves (Rupley and Careri, 1991). Under these conditions, and in the high-frequency limit, the static dielectric constant is actually closer to 2. Therefore,  $\epsilon_{\text{in}} \approx 2$  reflects mainly electronic polarizability.  $\epsilon_{\text{in}} \approx 4$  is thought to include contributions by relaxation of permanent dipoles in addition to electronic polarizability.

In the calculations in which four water molecules were treated explicitly,  $\epsilon_{\text{in}} \approx 4$ –5 captured pK<sub>a</sub> values of both Lys-66 and Glu-66 with a single conformation. This suggests that, if the pattern of hydration observed crystallographically is the one populated in solution, then the buried groups can ionize without triggering any significant structural relaxation in SNase beyond the fluctuation of permanent dipoles already implicit in  $\epsilon_{\text{in}} = 4$ –5. This is consistent with the experimental data showing that SNase is folded and indistinguishable from the native protein at pH values below the pK<sub>a</sub> of Lys-66 and above the pK<sub>a</sub> of Glu-66. The implications are that SNase behaves as a rigid body that minimizes the energetic consequences of burial and ionization of a polar group in a highly nonpolar environment by allowing the penetration of water into the core, rather than by undergoing significant structural reorganization. The response is different from that expected from regions of proteins that have evolved to solvate buried charges. In the active sites of enzymes, for example, the dominant mechanisms of dielectric relaxation might involve a more significant reorientation of polar groups (Warshel, 1987). In these cases, the values of  $\epsilon_{\text{in}}$  used to represent protein relaxation are higher, ranging from 4 to 8 (Sham et al., 1997; Simonson et al., 1999; Schutz and Warshel, 2001).  $\epsilon_{\text{in}} \approx 4$  are more consistent with the microscopic fluctuations in the core of several proteins studied by molecular-dynamics simulations (Smith et al., 1993; Simonson and Brooks, 1996; Simonson and Perahia, 1996), and with values of  $\epsilon_{\text{in}}$  found to reproduce reorganizational energies in electron transfer systems (Muegge et al., 1997; Sharp, 1998).

One reason why the significance of  $\epsilon_{\text{in}} = 4$ –5 should not be over-interpreted is that the relaxation of the water molecules in the calculations in Figs. 8 and 9 was handled through an energy-minimization procedure, instead of through more physically realistic Monte Carlo or molecular-

dynamics simulations. The energy-minimized structures bias the orientation of water molecules toward the energetically most favorable state relative to the charged group. The natural result of this procedure is to minimize the values of  $\epsilon_{\text{in}}$  needed to capture the experimental pK<sub>a</sub> values. More explicit simulations are underway to determine whether the same effect could be achieved if configurational averaging over neutral and ionized states based on molecular dynamics were used to capture the reorganization of water molecules and protein dipoles explicitly.

The values of  $\epsilon_{\text{in}} \approx 7.5$  obtained when the pK<sub>a</sub> values are analyzed with only the one or two innermost water molecules treated explicitly suggest that conformational reorganization concomitant with ionization of the buried group can occur. Conformational changes were not detected by a variety of spectroscopic probes. Therefore, if they occur, they must be minor, and local rather than global. They might involve local and transient unfolding to expose the charged Lys-66 or Glu-66 to water, or a slight change in the interactions between the helix where position 66 is found and the rest of the protein, to enhance hydration of the charged group.

### CONCLUSIONS

When Lys or Glu are buried at position 66 in SNase, the equilibrium between the ionized and the neutral states shifts substantially in favor of the neutral state, thus the pK<sub>a</sub> values of acidic residues are elevated, and those of basic residues are depressed relative to pK<sub>a</sub> values in water. The dominant determinant of the pK<sub>a</sub> values of these buried residues is the difference in the Born energy of the charged species in bulk water and in the protein interior. Although the pK<sub>a</sub> shifts of Lys-66 and Glu-66 are among the largest ever measured, the remarkable aspect is that they are actually very modest compared to the energetic cost of transferring a charge from water to a nonpolar environment. In SNase, water penetration can apparently contribute significantly to the stabilization of ionizable and charged groups in the protein interior. Buried water molecules can ameliorate the energetic penalty of removal of a polar side chain from bulk water, and they can reduce the magnitude of the shift in pK<sub>a</sub> value that a buried group would experience. Modulation of pK<sub>a</sub> values of buried groups in the physiological range of pH through interactions with internal water molecules is a recurring motif encountered in many membrane proteins involved in H<sup>+</sup> and e<sup>−</sup> transport, and in active sites of enzymes.

Dielectric constants used in continuum methods to describe polarizability in the protein interior are empirical parameters meant to account implicitly for contributions not considered explicitly in computational models (Schutz and Warshel, 2001). These parameters need to be calibrated against microscopic calculations or against experimental data. In the case of SNase, the pK<sub>a</sub> value of residues buried

at position 66 were reproduced with a continuum model when  $\epsilon_{\text{in}} \approx 10$ . This value of  $\epsilon_{\text{in}}$  is high because it reflects significant contributions by water penetration and minor conformational reorganization that are not handled explicitly in the calculations. When four internal water molecules were treated explicitly, the value of  $\epsilon_{\text{in}}$  that reproduced  $\text{pK}_a$  values of buried residues dropped to 4–5. This should not be construed as evidence that  $\epsilon_{\text{in}} = 4\text{--}5$  is the true macroscopic dielectric constant of SNase in particular, or of proteins in general. It should also not be assumed that this value of  $\epsilon_{\text{in}}$  is relevant for calculations of  $\text{pK}_a$  values of buried residues with these methods in other proteins or even at other sites in SNase. The spatial dependence of  $\epsilon_{\text{in}}$  in the interior of proteins remains unknown.

The consequence of treating the buried, site-bound water molecules explicitly was to depress the value of  $\epsilon_{\text{in}}$  used in the calculation because, in general, the less implicit the model, the lower the dielectric constant required. The value of  $\epsilon_{\text{in}} \approx 4\text{--}5$  required to reproduce  $\text{pK}_a$  values when all internal waters were treated explicitly implies that the ionization of these buried residues does not result in a significant conformational reorganization. However, when fewer water molecules were treated explicitly, this value of  $\epsilon_{\text{in}}$  increases to values near 7.5. This means that, if the buried ionizable groups are less hydrated than suggested by the crystallographic structures, then conformational reorganization must take place upon ionization of the buried residues.

The meaningful interpretation of  $\Delta\text{pK}_a$  in terms of  $\epsilon_{\text{in}}$  in SNase was possible because position 66 is unique in several respects. The ionizable moieties of Glu-66 and Lys-66 sit in one of the most hydrophobic environments in which buried ionizable groups have ever been found (Mehler et al., 2002). There are no backbone or side-chain polar groups nearby that could affect the  $\text{pK}_a$  values significantly. The ionizable moieties of Lys-66 and Glu-66 are also buried too deeply to sense the reaction field of bulk water, and they are too far from surface charges to be affected through coulombic interactions. The response of SNase as a rigid body to the ionization of a group buried in its hydrophobic core is also noteworthy. Proteins that have evolved to solvate buried charges, for example, at active sites, are expected to be more pliable than SNase and might require  $\epsilon_{\text{in}} > 4$  to capture the energetics of ionization of buried groups, depending on the level of relaxation included explicitly in the calculations.

Without the extensive experimental constraints available for these SNase mutants, it would have been impossible to reach many of the conclusions of this study. For example, interactions between the buried and surface ionizable groups were excluded arbitrarily in some of the calculations because the experimental evidence indicated that these interactions are negligible. The evidence that water penetration played a significant role was based on direct crystallographic observation of water penetration; it is not yet obvious that the presence of buried water molecules could have been predicted from physical principles. The experi-

mental demonstration that the structural consequences of the ionization of buried groups are minimal was also necessary for the interpretation of the values of  $\epsilon_{\text{in}}$  that reproduce the experimental data.

The thermodynamic data and the crystallographic structures of mutants with Lys-66 and Glu-66 raise several interesting and challenging problems related to water penetration in proteins. The data suggest that water penetration can play a significant role in the stabilization of charge in the protein interior. They also suggest that crystallographic structures are not necessarily dependable sources of information about patterns of hydration of buried groups. Without the structure of mutants with Glu-66 to identify the hydration motif, any attempts based solely on the structure with Lys-66 to elucidate microscopic origins of polarizability in the protein interior would have failed to include the possible significant contributions of transient, buried water molecules in the hydrophobic core. The physical origins of permittivity in the protein interior would have been misinterpreted. To improve the accuracy and utility of structure-based calculations of electrostatic energy in proteins, methods will have to be developed to identify water-binding sites, to calculate the occupancy of the binding sites, and to calculate conformational relaxation of internal water molecules (Warshel and Papazyan, 1998).

SNase appears to be unable to retain a charged group in a nonpolar environment without undergoing at least local rearrangement or without water penetration, as proposed originally by Warshel for proteins in general (Warshel, 1981; Warshel and Papazyan, 1998). The extent to which water penetration occurs in other proteins or in other locations in SNase is unknown at this time. The magnetic relaxation dispersion studies by Halle and Denisov suggest that water penetration is common, and that internal water molecules exchange rapidly with bulk water (Denisov and Halle, 1996; Denisov et al., 1997). Water penetration has also been observed and characterized in molecular-dynamics simulations (García and Hummer, 2000). Given the significant consequences of internal water molecules on  $\text{pK}_a$  values of buried groups, it will be necessary to improve our understanding of water penetration in proteins and of properties of internal water molecules, to elucidate structural origins of function in many key biochemical systems.

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