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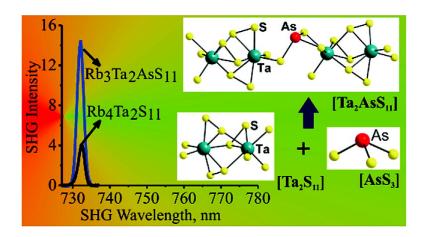
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Strong Second Harmonic Generation from the Tantalum Thioarsenates $A_3Ta_2AsS_{11}$ (A = K and Rb)

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Noncentrosymmetric (NCS) semiconductors can exhibit second harmonic generation (SHG) nonlinear optical (NLO) response in the mid-IR $(2-20 \,\mu\text{m})$ region, a spectral range of importance for molecular spectroscopy, atmospheric sensing, communications and various optoelectronic devices. The use of oxides, e.g., KTiOPO₄ (KTP), LiB₃O₅ (LBO), β-BaB₂O₄ (BBO), LiNbO₃ (LN), etc., which are the benchmark materials in the UV-vis to near-infrared (IR) region of the spectrum, is limited in the mid-IR region because of inadequate optical transparency and relatively low SHG efficiency.2 Chalcogenide semiconductors are more promising in the IR,3 but only a few, for example, AgGaSe2, AgGaS₂, GaSe, possess the necessary SHG susceptibility, optical transparency, chemical stability, etc. for technological applications.⁴ Recent studies in search for new NCS oxides with high SHG susceptibility are aimed at introducing asymmetric building units, such as second-order Jahn-Teller (SOJT) distorted d⁰ metal centers (e.g., Ti⁴⁺, Ta⁵⁺, Mo⁶⁺, etc.), anionic groups with stereochemically active lone pairs [e.g., $(IO_3)^{1-}$, $(TeO_x)^{n-}$, etc.], and noncentrosymmetric π -orbital systems [e.g., (BO₃)³⁻ and (B₃O₆)³⁻].⁷ Corresponding units in the chalcogenides, (e.g., polar [AsS₃]³⁻, [SbS₃]³⁻, [TeS₃]²⁻, etc.)⁸ have been little explored in the context of NLO properties.

Our recent investigation of the A/As/S (A = Li and Na) system revealed the soluble polar direct-band gap material Li_{1-x}Na_xAsS₂ (contains pyramidal [AsS₃]³⁻ units) with high SHG efficiency, which can approach ~30 times relative to the benchmark IR material AgGaSe2.9 Improvements of the SHG efficiency by introducing multiple asymmetric building units into the structure have been reported, e.g., both SOJT distorted d^0 metal centers and anionic groups with stereochemically active lone pairs in BaTeMo₂O₉, TlTeVO₅, etc. ¹⁰ SOJT distorted d^0 metal cations in oxides are abundant, whereas in the chalcogenide structures are less common owing to weaker M-S interactions relative to those of M-O. CsTaS₃ is an example of d⁰ metal chalcogenide, where the Ta⁵⁺ displacement results from SOJT effect.¹¹ Our investigation in the A/Ta/As/S (A = monovalent cations) system was motivated by the idea of introducing two different asymmetric units; $[Ta_mS_n]^{p-}$ and $[AsS_3]^{3-}$ in a single NCS structure. Here we report the synthesis and properties of K₃Ta₂AsS₁₁ (**Ia**), Rb₃Ta₂AsS₁₁ (**Ib**), and $Cs_3Ta_2AsS_{11}$ (II). We find that the NLO SHG of crystals of Ia-bis up to \sim 15 times stronger than that of commercially used AgGaSe₂.

Phase pure syntheses of Ia, Ib, and II were achieved with the alkali-metal polythioarsenate flux method using $A_2S/Ta/As/S$ mixtures. The purity of materials was confirmed by comparing powder X-ray diffraction patterns with those simulated from the single crystal X-ray diffraction analysis. The flux basicity (i.e., determined by the A_2S/S ratio) plays a crucial role in the synthesis of pure phase. The relatively basic flux ratios of 1/1/1/4 or 3/1/1/12 did not produce the title phase; instead gave compounds without arsenic namely black

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 $ATaS_3^{11}$ and orange $A_4Ta_2S_{11}.^{14}$ A decrease of the flux basicity, which was achieved by increasing the relative S fraction, led to appreciable proportions of the title quaternary compound and finally 1/1/1/12 was found to be the optimum flux ratio to produce pure $A_3Ta_2AsS_{11}.$ The air and moisture stable orange-red compounds were further characterized with electronic absorption spectroscopy, Raman spectroscopy and differential thermal analysis (see below). A medium size crystal of ${\bf Ib}$ $(1.0\times0.45\times0.24~{\rm mm}^3)$ isolated from the flux reaction is shown in Figure 1a.

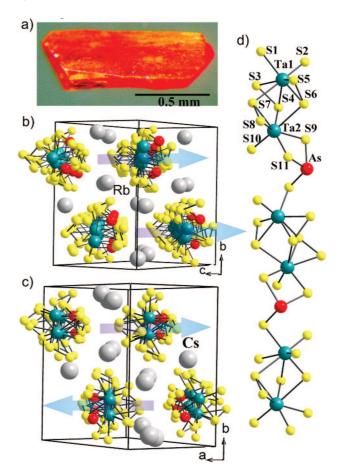


Figure 1. (a) Photo of a medium-size crystal of Rb₃Ta₂AsS₁₁ grown in a polysulfide flux. (b) Noncentrosymmetric packing of the ${}^{1}/_{\infty}$ [Ta₂AsS₁₁^{3−}] chains in **Ia−b**. The in-phase arrangement of chains is shown by the pale blue arrows. (c) Centrosymmetric packing of the ${}^{1}/_{\infty}$ [Ta₂AsS₁₁^{3−}] chain in **II**. (d) A single ${}^{1}/_{\infty}$ [Ta₂AsS₁₁^{3−}] chain, showing the connectivity between AsS₃ trigonal pyramids and asymmetric [Ta₂S₁₁]^{6−} units.

Compounds **Ia** and **Ib** crystallize¹⁵ in the monoclinic space group Cc whereas **II** crystallizes in $P2_1/n$. All compounds contain the same

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parallel ${}^1/_\infty[Ta_2AsS_{11}{}^{3-}]$ polymeric anionic chains, (Figure 1b, 1c). The size of the alkali-metals has a profound effect on the packing of the chains. The K^+ or Rb^+ favor noncentrosymmetric packing (Figure 1b) of the ${}^1/_\infty[Ta_2AsS_{11}{}^{3-}]$ chains, whereas the larger Cs^+ favors the centrosymmetric packing (Figure 1c). The cross-section of a chain is irregular and the noncentrosymmetric packing arises from the in-phase alignment of the AsS_3 pyramids. The structure of the ${}^1/_\infty[Ta_2AsS_{11}{}^{3-}]$ chain is similar to its Se analogue. 16 The chain is a polysulfide species, which is build up of bimetallic $[Ta_2S_{11}]^{6-}$ units linked with AsS_3 pyramids (Figure 1d).

Conceptually, the chain derives from the oxidative insertion of As into the S–S units of the $[Ta_4S_{22}]^{6-}$ core 17 as shown in eq 1. The common dimeric core $[Ta_2S_{11}]^{6-},^{14}$ derives from the trigonal face sharing of two distorted TaS_7 pentagonal bipyramids. It can be best described as $[Ta_2S_5(S_2)_3]^{6-},$ where the S–S distances [2.063(2)-2.098(2)~Å] are within the range of typical S_2^{2-} anion. 18

$$[Ta_4S_4(S_2)_9]^{6-} + 2As \rightarrow [Ta_4S_4(S_2)_6(AsS_3)_2]^{6-}$$
 (1)

The presence of disulfide $(S_2)^{2^-}$ units in TaS₇ was confirmed by the intense peaks in the range 485–525 cm⁻¹ in the Raman spectrum (Figure 2a).¹⁸ All As–S bond distances are normal¹⁹ and A–S (A = K, Rb, and Cs) bonds are mostly ionic in nature.

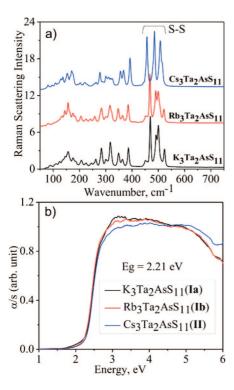


Figure 2. (a) Raman spectra of **Ia**, **Ib**, and **II**, showing the characteristic shift from disulfide anion $(S_2)^{2^-}$. (b) Solid-state UV—vis optical absorption spectrum of **Ia**, **Ib**, and **II**, shows sharp absorption edge at 2.21 eV.

Electronic absorption spectroscopy of the solid samples showed a sharp absorption edge at 2.21 eV for all three compounds consistent with their orange color (Figure 2b). The alkali-metal ions have negligible effect on the energy gap suggesting purely ionic interactions with the $^{1}/_{\infty}[\mathrm{Ta}_{2}\mathrm{AsS}_{11}^{3-}]$ chains. The optical excitation is believed to originate from S-based *p*-orbital dominating valance band to the Tabased orbitals in the conduction band.²⁰

Differential thermal analysis (DTA) confirmed that all three compounds melt incongruently (see the Supporting Information). Both melting and crystallization points are within 400–500 °C for all three compounds and those values decease on going from ${\bf Ia}$ (K) to ${\bf Ib}$ (Rb)

to \mathbf{II} (Cs). Though single melting and crystallization peaks were observed in the DTA plots, powder X-ray diffraction done after DTA, showed extra diffraction peaks along with the title compounds, suggesting thermal decomposition. For \mathbf{Ib} , the black decomposition product observed after DTA was identified as RbTaS₃,²¹ whereas that for \mathbf{II} was CsTaS₃.¹¹ Since the syntheses of the compounds were achieved at 550 °C, which is above their melting points, it is apparent that the polysulfide flux allows for the formation of the incongruently melting phases in pure crystalline form either during the soaking time or on cooling. Thus, large optical quality crystals could best be obtained by the flux technique as in the case of $A_2Hg_3M_2S_8$ (A = K, Rb; M = Ge, Sn).²²

The excellent nonlinear optical SHG response is the most significant property of the polar compounds ${\bf Ia}$ and ${\bf Ib}$. It was measured with a variable wavelength (1000–2000 nm) laser source using a modified Kurtz and Perry method (see the Supporting Information).²³ A comparison of the SHG response relative to the benchmark material ${\rm AgGaSe_2}$ was also done using the same particle size ranges mainly in the wavelength >700 nm, Figure 3. Efficiency comparison below 700

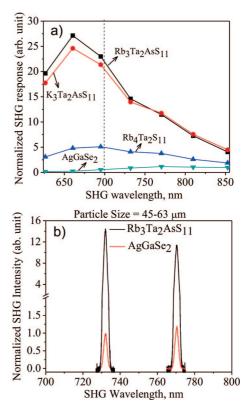


Figure 3. (a) Comparison of the SHG response of Ia, Ib, and Rb₄Ta₂S₁₁, relative to AgGaSe₂ at different SHG wavelength, using powder samples (average particle size of 54 (± 9) μm). The solid lines between the points are a guide to the eye. The vertical dotted line represents the band-edge (1.80 eV) of AgGaSe₂. (b) Overlay of SHG signals from Rb₃Ta₂AsS₁₁ and AgGaSe₂ at two different wavelengths. In both cases the average particle size was 54 (± 9) μm.

nm is not reliable because of the band-edge absorption of AgGaSe₂. In the range 700–900 nm the SHG efficiency was \sim 15 times stronger than that of AgGaSe₂ and comparable to the recently reported CsZrPSe₆.²⁴ Particle size dependent SHG measurements indicated that both **1a** and **1b** are type-I nonphase-matchable at 770 nm. Generally, when a powder sample is nonphase-matchable, SHG sensitivity peaks near the coherence length, which is typically 1–20 μ m, and it start to diminish through larger particle sizes. Despite this such materials can still be used through 'random' quasi-phase-matching techniques.²⁵

phase-matchability also varies with wavelength, for example AgGaSe₂ is phase-matchable only in the range of 3800–12400 nm.²⁶ So, it is possible that **Ia**-**b** are phase-matchable at a different wavelength range beyond the one studied here. The highly polar structure that results from the alignment of dipoles in Ia-b (cf. Figure 1b) is likely responsible for the enhanced SHG intensity compared to AgGaSe₂, which is only weakly polar. Unlike the Li_{1-x}Na_xAsS₂ system, where we observed a systemetic increase of the SHG efficiency with x, the efficiency of both end-members (Ia and Ib) of K_{3-x}Rb_xTa₂AsS₁₁ system are similar (cf. Figure 3), suggesting negligible contribution from the alkali-metal polarizability.

To better assess the relative contribution of the pyramidal [AsS₃]³⁻ and $[Ta_2S_{11}]^{6-}$ units toward the total polarizability, we examined the NLO properties of Rb₄Ta₂S₁₁, a related NCS compound with similar chains containing the polysulfide unit of $[Ta_2S_{11}]^{4-}$ but no $[AsS_3]^{3-}$ units. 14 The SHG response of Rb₄Ta₂S₁₁ was found to be strong but only $\sim 4 \times$ that of AgGaSe₂, Figure 3. This result is suggestive of the predominant role of the pyramidal [AsS₃]³⁻ unit on the total polarity and consequently on the SHG response of the Ib.

In conclusion, the polysulfide flux is critical in stabilizing the strongly anisotropic thioarsenates A₃Ta₂AsS₁₁. The combination of two asymmetric units, [Ta₂S₁₁] and [AsS₃], in a single strand coupled with polar packing of strands appears to lead to strong NLO SHG response. This implies that the approach of combing different asymmetric fragments (e.g., chalcogenides) to impart strong polarity in extended structures could be promising in finding exceptional candidate materials for NLO applications. The results of this study also justify future synthetic investigations in the A/M/As/Q systems (A = monovalentcations; M = V, Nb, Ta; Q = S, Se, Te).

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Supporting Information Available: X-ray crystallographic files (CIF) and experimental details for K₃Ta₂AsS₁₁, Rb₃Ta₂AsS₁₁, and Cs₃Ta₂AsS₁₁. This material is available free of charge via the Internet at http:// pubs.acs.org.

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- $(12)\ \ K_{3}Ta_{2}AsS_{11}(\textbf{Ia}).\ A\ mixture\ of\ K_{2}S\ (0.050\ g,\ 0.45\ mmol),\ Ta\ (0.082\ g,\ 0.45\ mmol),\ Ta\ (0.$ mmol), As (0.034 g, 0.45 mmol), and S (0.175 g, 5.45 mmol) was loaded into a fused-silica tube in a nitrogen-filled glovebox. It was flame-sealed under vacuum ($\sim 10^{-4}$ mbar) and then heated to 550 °C in 6 h. After 48 h of soaking it was cooled down to 250 °C in 60 h followed by rapid cooling to room temperature. An orange-red color compound was isolated as a single phase in >60% yield after dissolution of the excess flux with degassed DMF and unreacted sulfur with CS₂. Semiquantitative energy dispersive analysis (EDS) gave an average composition of K_{3.1}Ta_{2.2}As_{1.2}S_{12.1}; $Rb_3Ta_2AsS_{11}(\mathbf{1b})$. A mixture of Rb_2S (0.081 g, 0.40 mmol), Ta (0.052 g, 0.40 mmol), As (0.030 g, 0.40 mmol), and S (0.154 g, 4.80 mmol) was loaded into a fused-silica tube in a nitrogen-filled glovebox. It was flamesealed under vacuum ($\sim 10^{-4}$ mbar) and then application of a heating and cooling cycle similar to that above, produced orange-red crystalline product as a single phase in >65% yield. Semiquantitative energy dispersive analysis (EDS) gave an average composition of Rb_{3.1}Ta_{2.1}As_{1.2}S_{11.8}; Cs₃Ta₂AsS₁₁(II). It was isolated from the reaction of Cs₂S (0.107 g, 0.36 mmol), Ta (0.065 g, 0.36 mmol), As (0.027 g, 0.36 mmol), and S (0.139 g, 4.32 mmol) using a procedure similar to that above. The compound was obtained as a single phase in >60% yield after dissolving the excess flux with degassed DMF and un reacted sulfur with CS₂. EDS analysis gave an average composition
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- (15) Single-crystal X-ray diffraction data were collected using a STOE imagingplate diffraction system (IPDS-2T) with graphite-monochromatized Mo Ka radiation. An analytical absorption correction was applied. Direct methods and full-matrix least squares refinements against F² were performed with and full-matrix least squares refinements against F^2 were performed with the SHELXTL package. Crystal data for $K_3Ta_2AsS_{11}(Ia)$: monoclinic Cc, Z=4, a=14.2571(7) Å, b=12.4154(7) Å, c=9.4505(5) Å, $\beta=100.120(4)^\circ$, V=1646.78(15) Å 3 at 100 K, $\theta_{\rm max}$ (Mo K α) = 29.16° , total reflections = 7683, unique reflections $F_0^2 > 2\sigma$ (F_0^2) = 4151, number of variables = 155, $\mu=17.397$ mm $^{-1}$, $D_c=3.657$ g cm $^{-3}$, $R_{\rm int}=4.38\%$, GOF = 1.034, $R_1=2.15\%$, $R_{\rm w}=2.23$ for $I>2\sigma(I)$. Crystal data for Rb $_3$ Ta $_2$ AsS $_{11}(Ib)$. monoclinic Cc, Z=4, a=14.4397(8) Å, b=12.8013(9) Å, c=9.5687(5) Å, $\beta=99.289(4)^\circ$, V=1745.55(18) Å 3 at 100 K, $\theta_{\rm max}$ (Mo K α) = 29.22° , total reflections = 8079, unique reflections $F_0^2 > 2\sigma$ (F_0^2) = 4136, number of variables = 155, $\mu=24.009$ mm $^{-1}$, $D_c=3.390$ g cm $^{-3}$, $R_{\rm int}=4.80\%$, GOF = 1.047, $R_1=3.10\%$, $R_{\rm w}=3.31$ for $I>2\sigma(I)$. Crystal data for Cs_3 Ta $_2$ AsS $_{11}(II)$. monoclinic P_2/h , Z=4, a=9.5967(7) Å, b=13.5926(7) Å, c=14.7160(10) Å, $\beta=98.876(6)^\circ$, V=1896.63(21) Å 3 at 100 K, $\theta_{\rm max}$ (Mo K α) = 29.14° , total reflections = 14682, 1896.63(21) Å 3 at 100 K, $\theta_{\rm max}$ (Mo K α) = 29.14° , total reflections = 14682,

- 9.5967(7) Å, b = 13.5926(7) Å, c = 14.7160(10) Å, $\beta = 98.876(6)^{\circ}$, V = 1896.63(21) ų at 100 K, θ_{\max} (Mo K α) = 29.14 $^{\circ}$, total reflections = 14682, uniquelections $F_0^2 > 2\sigma$ (F_0^2) = 5058, number of variables = 155, $\mu = 20.122$ mm⁻¹, $D_c = 4.161$ g cm⁻³, $R_{\text{int}} = 4.96\%$, GOF = 1.168, $R_1 = 4.19\%$, $R_{\text{w}} = 5.59$ for $I > 2\sigma(I)$.

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