CH₃O⁺ and C₂H₅O⁺: High Barriers to Isomerisation and Low Barriers to Symmetry-allowed 1,1-Elimination

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Summary Experimental and theoretical evidence is presented to show that the alkoxy cations CH_3O^+ and $C_2H_5O^+$ have surprisingly high barriers to 1,2-hydride shifts to the oxygen atom, but have very low barriers to 1,1-elimination of H_2 from the α -carbon atom.

In earlier work, we have shown that $CH_2=OH$ undergoes symmetry-forbidden 1,2-elimination of H_2 in a reaction requiring ca. 330 kJ mol⁻¹. The high internal energy required to dissociate $CH_2=OH$ was shown by deuterium

labelling^{1,2} to be insufficient to cause any prior isomerisation of the ion. Yet the reported heat of formation of the methoxy cation (CH₃O⁺, 850 kJ mol⁻¹),³ is only 140 kJ mol⁻¹ higher than that of CH₂=OH (710 kJ mol⁻¹).⁴ These data imply that the barrier to the CH₂=OH \rightarrow CH₃O⁺ isomerisation is >330 kJ mol⁻¹, and that the barrier for the reverse reaction CH₃O⁺ \rightarrow CH₂=OH (occurring through the same channel) is >190 kJ mol⁻¹. The latter barrier appears, at first sight, to be a remarkably high one to a symmetry-allowed 1,2-hydride shift which can lead to a

thermodynamically more stable product. In contrast, in the gas phase a primary carbonium ion appears to rearrange to a secondary carbonium ion via a 1,2-hydride shift essentially without activation energy.5,6

However, unambiguous experimental evidence for a

barrier to the reaction CH₃O⁺ → CH₂=OH is available from the work of Hiraoka and Kebarle.7 Using a pulsed, highpressure mass spectrometer source, these workers were able to show that in a thermally equilibrated system, H2 will add to the formyl cation in a reaction which is exothermic by 16 kJ mol⁻¹. This exothermicity is calculated from the temperature variation of the equilibrium constant for the reaction, i.e. the reaction is readily reversible. The adduct is formulated as CH₃O⁺,7 as it clearly must be since 1,2addition of H₂ to HC+=O would require an activation energy of ca. $220 \ \mathrm{kJ} \ \mathrm{mol^{-1}}$, if occurring by microscopic reversal of 1,2-elimination of H2 from CH2=OH. Moreover,

the reaction $\text{CH}_2\!\!\stackrel{+}{=}\!\!\text{OH} \to \!\!\text{HC}\!\!\stackrel{+}{=}\!\!\text{O} + \text{H}_2$ could not occur in a system in thermal equilibrium at the temperatures employed (-100 to -165 °C). Thus, H₂ adds 1,1 to the formyl ion, and dissociation of the adduct requires less energy than

isomerisation to CH₂=OH.

The product CH_3O^+ from the above $H_2 + HC^+=O$ reaction has ΔH_f 810 kJ mol⁻¹, and this should be an accurate value.7 If it is indeed a more accurate value than the appearance potential value (850 kJ mol-1) quoted earlier, then the barrier (≥230 kJ mol⁻¹) derived for the $CH_3O^+ \rightarrow CH_2=OH$ reaction is even higher than the one quoted at the outset of this paper. We have sought confirmation of the above deductions by

generating CH₃O+ and CH₂=OH in the mass spectrometer by ionisation and fragmentation of selected compounds. Since it appears that RCH₂O+ ions lose H₂ through 1,1elimination with a very small activation energy, precursors of these ions should give very low abundance RCH₂O+ ions and high abundance RC=O ions. In contrast, CH2=OH ions generated from RCH₂OH compounds should give large m/e 31:29 ratios since the m/e 31 \rightarrow 29 reaction has a very large activation energy. The experimental results are given in the Table.

TABLE M/e 31:29 ratios in the mass spectra (70 eV) of selected compounds

Compound	Anticipated ion at threshold	m/e 31:29
CH₃CH₂OH	$CH_2 = \mathring{O}H$	11:1
CF ₃ CH ₂ OH	$CH_2 = \mathring{O}H$	11:1
CH ₂ CH ₂ CH ₂ OH	$CH_2 = \mathring{O}H$	15:1
CH ₃ ONO ₂	CH ₃ O+	0.08:1
CH ₃ OCH ₃	CH ₃ O+ CH ₃ O+	0.08:1

The results strikingly confirm expectations, both CH3-ONO₂ and CH₃OCH₃ producing extremely abundant HC=O ions. Furthermore, while CH2=OH loses H2 in slow reactions (metastable transitions) with a large kinetic energy $release^2$ (140 kJ mol^{-1}) since the reaction is symmetryforbidden,1 the CH₃O+ \rightarrow HC=O reaction does not give a discernible metastable peak. Presumably this is because the symmetry-allowed reaction (see below) has a small activation energy and a relatively fast rise of the rate constant with energy, such that almost all the dissociations are relatively fast.

Finally, we wished to establish that an acyl ion RC=O, when formed from a compound RCH2OX, does arise specifically by 1,1-elimination of H₂ from the oxygenbound carbon even where other processes are possible. Ethyl nitrate produces in its 70 eV mass spectrum peaks at m/e 43 (B) and 45 (A) in the ratio 20:1 (Scheme 1). [1,1- 2 H₂]-Ethyl nitrate produces under similar conditions peaks at m/e 43 (B) and 47 (C) in the ratio 12:1, with peaks at m/e 44 and 45 being of negligible abundance relative to that at m/e 43. The CH₃C=O ion is therefore produced by

a specific 1,1-elimination.

$$\begin{array}{c} -\mathrm{e} \\ -\mathrm{H_{3}CH_{2}ONO_{2}} \xrightarrow{-\mathrm{e}} \mathrm{CH_{3}CH_{2}O^{+}} \xrightarrow{-\mathrm{H_{2}}} \mathrm{CH_{3}C=O} \\ -\mathrm{NO_{2}} \\ (\mathbf{A}), \ m/e \ 45 \qquad (\mathbf{B}), \ m/e \ 43 \\ \\ \mathrm{CH_{3}CD_{2}ONO_{2}} \xrightarrow{-\mathrm{e}} \mathrm{CH_{3}CD_{2}O^{+}} \xrightarrow{-\mathrm{D_{2}}} \mathrm{CH_{3}C=O} \\ -\mathrm{NO_{2}} \\ (\mathbf{C}), \ m/e \ 47 \qquad (\mathbf{B}), \ m/e \ 43 \\ \\ \end{array}$$

The extremely small barrier to the reaction CH₃O⁺ → HC=O + H2 is supported by earlier orbital symmetry arguments.8 1,1-Elimination from an sp3 carbon adjacent to a vacant p-orbital is a symmetry-allowed process.

It remains to explain the high barrier (≥ 230 kJ mol⁻¹) to the reaction $CH_3O^+ \rightarrow CH_2=OH$. The hydride shift should occur so that bonding on to the vacant p-orbital on oxygen is possible in the transition state $(1) \rightarrow (2)$ (Scheme

2). However, this results in the formation of CH2OH in which the lone-pairs of electrons on the oxygen atom are orthogonal to the vacant π -orbital on the carbon atom (3). Thus (3), formed in this manner, lacks all the π -stabilisa-

tion of CH₂=OH, but suffers all the destabilisation of the CH₂OH cation which exists through σ-electron withdrawal by the electronegative oxygen atom; it may plausibly represent the highest point on the potential surface of the

 $CH_3O^+ \rightarrow CH_2=OH$ isomerisation. The heat of formation of (3) may be estimated in the following manner. The difference in heats of formation of CH₄ and CH₃⁺ is 1155 kJ mol-1;9 the analogous change performed on methanol $[CH_3OH \rightarrow (3)]$ should require a similar energy difference if π back-donation to the cationic centre did not occur, nor did inductive destabilisation by oxygen. Lack of π -donation is precisely the situation which we wish to examine in (3). Inductive destabilisation will however be present in (3); calculations by Pople and his co-workers 10 give an inductive destabilisation of a carbocation of 42 kJ mol⁻¹ when the oxygen atom is separated by one carbon atom from the cationic centre, and 13 kJ mol-1 when it is separated by two carbon atoms. In view of the rapid decrease of inductive effects due to an increase in the number of intervening σ -bonds, it seems probable that the inductive destabilisation due to a directly bonded oxygen will not be

less than $100 \,\mathrm{kJ} \,\mathrm{mol}^{-1}$. We therefore conclude that ΔH_{f} (3) is $\geq 1055 \text{ kJ mol}^{-1}$. Thus, a plausible transition state for the reversible reaction $CH_2=OH \rightleftharpoons CH_3=O$ appears to be energetically inaccessible at the high internal energies necessary to induce the reaction $CH_2 \stackrel{+}{=} OH \rightarrow HC \stackrel{+}{=} O+ H_2$; this is in accord with the experimental facts.

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