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Letters to the Editor

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The Thermal Reaction Between Hydrogen and Oxygen at Higher Pressures

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January 27, 1942

In recent experiments we found conditions under which the thermal hydrogen-oxygen reaction ($p=60~{\rm cm}$; T up to $560^{\circ}{\rm C}$) is sufficiently reproducible so that the temperature dependence of the rate could be measured. From the constant activation energy observed over a certain temperature range, we inferred that here the reaction proceeds in non-branched chains. In a recent letter to the editor, Dainton¹ concludes that the chains are branched. From an enlargement of our pressure-time curve representing a typical reaction, he concludes that at the beginning the first derivative of the rate $(\delta^2(\Delta p)/\delta t^2)$ is increasing and that the plot of $\log \Delta p$ against time is linear; applying Semenoff's theory, he argues that in the first stages of the reaction branching of the chains is occurring and, in fact, more frequently than their termination.

This conclusion seems questionable for the following reasons. The beginning of our curve analyzed by Dainton is not reproducible enough to apply Semenoff's criterion; in many cases $\delta^2(\Delta p)/\delta t^2$ seems to begin with its largest value; in still more cases (old vessel) the reaction begins with nearly the maximum rate. The reproducibility of the maximum rate (or rather log rate), however, is evident from our curves (reference 1, Figs. 2–5).

Moreover, immediately after high evacuation near the annealing temperature of Pyrex and subsequent admission of almost one atmosphere of gases, one should expect rapid changes of the surface by adsorption superimposed on the effects considered by Semenoff. This difficulty may be less pronounced in the low pressure explosion (Kowalsky, pressure a few mm, reaction lasting a small fraction of a second); this reaction, which is quantitatively described by Semenoff's theory, takes place under conditions so widely different from those at high pressure that Hinshelwood and Williamson even postulated a mechanism of a different nature. Whether the high pressure reaction should actually be interpreted by Semenoff's theory can possibly be determined by a photochemical investigation; if the induction period is correlated with the time required for the develop-

ment of a chain, a corresponding *decay period* should be observed after the interruption of the illumination. Taylor and Salley, in their photochemical investigation, mention the induction period but not such a decay period.

There is another difficulty in Dainton's interpretation. He assumes that in the beginning of the reaction the branching exceeds the breaking and the explosion is prevented by the change of concentrations occurring during the development of the first chains. He argues that the maximum rate occurs when $\phi = 0(\delta - \beta = 0)$ and that in this region the chains are of constant finite length. While we are unable to find a proof for this general theorem, it seems to us that the condition for the maximum rate would largely depend on the detail of the reaction. In the case thoroughly treated by Semenoff,² "the velocity of the reaction begins to decrease long before $\delta - \beta$ becomes equal to 0."

We applied the conventional procedure (Hinshelwood and Thompson, Taylor and Salley, and Prettre), measuring the nearly constant rate after the induction period. This seems justified in the limiting case of an initial rate which nearly equals and rapidly reaches the maximum rate (many observations in Pyrex, Prettre's observations in the KCl covered vessel); here one should be able to derive an activation energy from the temperature coefficient of the maximum rate. The same procedure leads to well-defined values of the activation energy for various kinds of surface vielding short or long induction periods. Therefore, we consider our interpretation of our curves most plausible, although not inevitable. This procedure does not involve an assumption regarding the nature of the induction period. In order to explore the effects of a pretreatment of the walls and of impurities on this period, we started experiments a year ago which are now interrupted.

¹ F. S. Dainton, J. Chem. Phys. 9, 826 (1941); see F. S. Dainton and R. G. W. Norrish, Proc. Roy. Soc. A177, 393 (1941). For further references see Dainton's paper and our paper, J. Chem. Phys. 9, 432 (1941).

(1941).

² C. N. Hinshelwood and A. T. Williamson, The Reaction between Hydrogen and Oxygen (Clarendon Press, Oxford, 1934), p. 52.

³ N. Semenoff, Chemical Kinetics and Chain Reactions, p. 62.

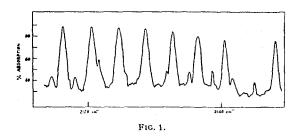
The Infra-Red Absorption of C¹³O¹⁶ at 4.66 Microns

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January 23, 1942

A FTER taking the necessary experimental precautions, it has been found possible to find rotational lines due to the isotopic carbon monoxide molecule C¹⁸O¹⁶. These lines are found interspersed among the lines of the fundamental band of C¹²O¹⁶ at 4.66 microns.

A portion of the absorption curve is shown in Fig. 1. The more intense lines are due to the $C^{12}O^{16}$ molecule. The weaker lines are associated with the $C^{13}O^{16}$ molecule, and comprise part of the R branch of that molecule's fundamental band. At lower frequencies than are shown here the isotopic lines become relatively more intense. As is to be expected, the line spacing for the heavier isotope is less than that for the lighter one.



The grating used was an echelette replica of 2887 lines per inch, ruled for the region 9 to 12 mu and employed in the second order. Readings of percent absorption were made at intervals of 5 seconds of arc, equivalent to about 0.13 cm⁻¹. The slit width was equivalent to 0.22 cm⁻¹.

Measurements are soon to be made on the overtone band and a complete report will be made at a later date.

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The Vibrational Structure of Electronic Transitions for Some Complex Ions

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February 9, 1942

POR several compounds of some of the transition elements, absorption spectra consisting of groups of approximately equidistant bands have been reported. The positions of the bands and the intervals between them depend only upon the metal ion and upon the identity of the complex of which the metal ion is a part. These spectra may be interpreted as originating in electronic transitions in which the vibrational structure characteristic of the complex ion present is resolved. The experimental basis for this interpretation is summarized in the following paragraphs.

Dreisch and Trommer¹ have found in dilute aqueous solutions of CoCl₂, CoBr₂ and CoI₂ an absorption band at 1.23μ consisting of six maxima with an average separation of 248 cm⁻¹. This band may be attributed to the Co(OH₂)₆⁺⁺ ion. In aqueous solutions of NiCl2, NiSO4 and Ni(NO8)2 they have found at 1.17μ a band characteristic of the Ni(OH₂)₆⁺⁺ ion. This band is resolved into four maxima with a separation of 264 cm⁻¹. In ammoniacal solutions of Ni(NH₈)₆Cl₂ the band is shifted to 0.940 µ and the separation becomes 390 cm⁻¹, the number of maxima remaining the same. For solutions of a number of ferrous salts, Dreisch and Kallscheuer² have reported a series of twelve or thirteen bands extending from about 0.7μ to about 1.27μ . The intervals separating the five bands of highest frequency vary considerably from salt to salt. The separations of the next seven bands are approximately the same for all of the salts measured. This constant separation, 324 cm⁻¹, is then characteristic of the Fe(OH₂)₆⁺⁺ ion.

The absorption spectra, measured at liquid air temperatures, of solid hydrates of some of the above salts have been

reported.³ These spectra consist of many bands, from about 20 cm⁻¹ to about 150 cm⁻¹ in width, extending throughout the visible region. Analysis indicates that these spectra may, in part, be interpreted in the same manner as those from the solutions. With CoCl₂·6H₂O, the average wave number difference for twelve pairs of bands is 226 cm⁻¹. With CoBr₂·6H₂O, the difference for ten pairs of bands is 220 cm⁻¹. For NiCl₂·6H₂O, the difference for nine pairs of bands is 237 cm⁻¹. For Co(OH₂)₆++, in passing from solution to the crystal, there is a decrease of 25 cm⁻¹ in the separation of the vibrational bands. For Ni(OH₂)₆++, there is a corresponding decrease of 27 cm⁻¹.

The interpretation presented here is consistent with the information obtained from the Raman effect of amino complex ions. Cu(NH₃)₄++, Zn(NH₃)₆++, and Cd(NH₃)₆++ give Raman displacements of the same order of magnitude as the separations of the vibrational levels of Ni(NH₃)₆++ and the aquo complexes (see Table I). If one makes the

TABLE I.

Complex ion	Vibration frequency cm ⁻¹	Force constant dynes/cm
Cu(NH ₃) ₄ ⁺⁺ (solution)	415 (Raman)	1.16×105
Cd(NH ₃) ₆ ++ (solution)	340 (Raman)	0.85
Zn(NH ₃)6 ⁺⁺ (solution)	423 (Raman)	1.22
Ni(NH ₃) ₆ ⁺⁺ (solution)	390 (Absorption)	1.01
Ni(OH ₂)6 ⁺⁺ (solution)	264 (Absorption)	0.52
Ni(OH ₂) ₆ ++ (crystal)	237 (Absorption)	0.42
Co(OH ₂) ₆ ⁺⁺ (solution)	248 (Absorption)	0.46
Co(OH ₂)6 ⁺⁺ (crystal)	223 (Absorption)	0.37
Fe(OH ₂) ₆ ++ (solution)	324 (Absorption)	0.77

approximation that the vibration of the complex ion may be represented by a simple harmonic oscillator, the force constant of the bond between the central metal ion and the coordinated group may be estimated from the vibration frequency. The values so obtained are listed in the table.

Some evidence for anharmonicity is obtained from a series of five bands for solid CoCl₂·6H₂O.³ The wave numbers of these bands may be represented with an accuracy of one wave number by:

$$\nu = 24504 + 222n - 1.6n^2 \quad (n = 0, \dots, 4).$$

From the constants of this equation, the heat of dissociation of $Co(OH_2)_6^{++}$ in crystals of $CoCl_2 \cdot 6H_2O$ is estimated to be 22.000 cal. The heat of the reaction:

$$CoX_2 \cdot 6H_2O(s) \rightarrow CoX_2(s) + 6H_2O(l)$$

determined from thermochemical data (I.C.T.) is 21,300 cal. for X = Cl, 19,800 cal. for X = Br, and 16,700 cal. for $X = \text{NO}_3$. The decrease in the thermochemical value for the bromide as compared with the chloride parallels the decrease in the average separations of the vibrational levels, 226 cm⁻¹ for the chloride as compared with 220 cm⁻¹ for the bromide. The nitrate has not yet been studied spectroscopically.

¹Th. Dreisch and W. Trommer, Zeits. f. physik. Chemie B37, 37 (1937).

² Th. Dreisch and O. Kallscheuer, Zeits. f. physik. Chemie B45, 19 (1939).

³ Gieleszen, Ann. d. Physik [5] 22, 537 (1935).

4 I. Damaschun, Zeits. f. physik. Chemie B16, 81 (1932).