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Electron-impact total ionization cross sections of silicon and germanium hydrides

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Electron-impact total ionization cross sections of some silicon and germanium compounds have been calculated by applying a new theoretical model that has been found to be reliable for a wide range of molecules. The new theory, the binary-encounter-Bethe (BEB) model, combines the binary-encounter theory and the Bethe theory for electron-impact ionization, and uses simple theoretical molecular orbital data—binding energies, average kinetic energies, and occupation numbers—which are readily available from molecular structure codes. Total ionization cross sections of SiH, SiH₂, SiH₃, SiH₄, Si₂H₆, Si(CH₃)₄, GeH, GeH₂, GeH₃, GeH₄, and Ge₂H₆ are presented for incident electron energies T from threshold to 1 keV, and compared to available experimental data. Theory and experiment agree well for SiH _{x} , $x=1-4$, from thresholds to $T < 80$ eV, while theoretical peaks occur at lower T than experimental peaks for SiH _{x} , $x=1-3$. No experimental data are available for germanium hydrides for comparison. The theoretical cross sections are given by a compact analytic form suitable for applications in plasma processing.

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I. INTRODUCTION

It has been shown in previous publications that the new theoretical method, the binary-encounter-Bethe (BEB) model,¹⁻³ produces reliable electron-impact total ionization cross sections for stable molecules and radicals using as inputs simple molecular orbital data for the ground state, all of which can be obtained by using standard molecular wave function codes. There are no adjustable or fitted parameters in the theory.

The BEB model, however, does not provide details such as resonances in the continuum, vibration and/or rotational excitations concomitant with ionization, multiple ionization, and dissociative ionization. The total ionization cross section in this method is the sum of ionization cross sections for ejecting one electron from each of the occupied molecular orbitals in the ground state. The BEB cross sections were found to reproduce reliable experimental data within 5%–15% on a large variety of molecules from the first ionization threshold to several keV in incident electron energies. In principle, this method can be used to calculate the total ionization cross section from any state of a molecule, provided that the state is predominantly a single configuration in character and required molecular orbital quantities are obtainable from molecular structure codes.

In the present paper, we apply the BEB model to 11 molecules of great current interest in the modeling of chemi-

cal vapor deposition (CVD) and plasma enhanced chemical vapor deposition (PECVD) of semiconductors:⁴⁻⁷ SiH, SiH₂, SiH₃, SiH₄, Si₂H₆, Si(CH₃)₄, GeH, GeH₂, GeH₃, GeH₄, and Ge₂H₆. Silane (SiH₄) is widely used for plasma assisted deposition of silicon and amorphous silicon–hydride (a:SiH) films.⁸ Silane has also been detected on Saturn and on several other planets and satellites. As was the case for 28 molecules reported earlier,¹⁻³ the BEB model provides cross sections in good agreement with available experimental data and sometimes enables one to discriminate between conflicting data for the same molecule. To the best of our knowledge, the total ionization cross sections for GeH _{x} , $x=1-4$, and Ge₂H₆ presented here are the first of its kind in the literature.

The main advantages of the BEB model are threefold. First, the model is valid for the entire energy range from threshold to several keV. Second, it provides compact analytic expressions for the total ionization cross section with constants characterizing occupied molecular orbitals. Third, the theory is applicable to a large variety of stable molecules and radicals as long as the ground state wave function can be calculated. The wave function need not be of very high quality.

We outline briefly the underlying theory in Sec. II, present the total ionization cross sections of individual molecules in Sec. III, and the conclusion in Sec. IV.

II. OUTLINE OF THE THEORY

The BEB model is a simpler version of the binary-encounter-dipole (BED) model for electron-impact ioniza-

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tion cross sections for atoms and molecules. The BED model combines the binary-encounter theory⁹ and the Bethe theory.¹⁰ The combination of the binary-encounter theory and the Bethe theory is achieved by requiring that the asymptotic form at high incident energy T of the former, match that of the latter, both in the ionization cross section and in the stopping cross section.¹ The stopping cross section is used to evaluate the stopping power of the target medium. The BED model provides a formula to calculate the singly differential cross section, or the energy distribution of ejected electrons $d\sigma/dW$, with the ejected electron energy W , for each atomic or molecular orbital. To apply the BED model, one needs for each occupied orbital the electron binding energy B , the average kinetic energy $U = \langle \mathbf{p}^2/2m \rangle$, where \mathbf{p} and m are the bound electron momentum and the electron mass, respectively, the electron occupation number N , and the continuum dipole oscillator strength df/dW .

The value of the average kinetic energy U for each orbital in the initial state can be obtained from any atomic or molecular wave function code that calculates the total energy in an independent particle model. To calculate df/dW , however, both the initial- and continuum-state wave functions are needed, thus making the application of the BED model non-trivial. Alternatively, df/dW can be deduced from experimental photoionization cross sections. Unfortunately, such data are available for only a few stable molecules over a large energy range. Moreover, partial photoionization cross sections are needed to deduce df/dW for each orbital, and the deduction of orbital photoionization cross sections from the total is problematic¹¹ at best.

To circumvent this difficulty, df/dW is approximated in the BEB model by choosing

$$df/dW = N/(B+W)^2. \quad (1)$$

This form of df/dW closely resembles the shape of the df/dW of the hydrogen atom in the ground state.

With this simplified df/dW , the integrated ionization cross section per molecular orbital, σ_{BEB} , is obtained by integrating $d\sigma/dW$ from $W=0$ to $(T-B)/2$. The result is¹

$$\sigma_{\text{BEB}} = \frac{S}{t+u+1} \left[\frac{\ln t}{2} \left(1 - \frac{1}{t^2} \right) + 1 - \frac{1}{t} - \frac{\ln t}{t+1} \right], \quad (2)$$

where $t = T/B$, $u = U/B$, $S = 4\pi a_0^2 N R^2/B^2$, $a_0 = 0.5292 \text{ \AA}$, and $R = 13.61 \text{ eV}$.

In Eq. (2), the term associated with the first logarithmic function on the right-hand side (rhs) represents distant collisions (large impact parameters) dominated by the dipole interaction, and the rest of the terms on the rhs represent close collisions (small impact parameters) as described by the Mott cross section.¹² The second logarithmic function originates from the interference of the direct and exchange scattering also described by the Mott cross section.

We present the values of B , U , and N for the 11 molecules in this article in Tables I and II. The molecular orbital data were generated for the ground states of the molecules by using the molecular structure code GAMESS.¹³ We have omitted deep inner shell orbitals with large binding energies because they contribute little to the total ionization cross sec-

tion. One can use either theoretical or experimental values of B if they are available from photoionization measurements with unambiguous assignments. However, U is a theoretical quantity that cannot be measured directly, although the sum of all U 's is equal to the magnitude of the total energy of the target molecule, according to the virial theorem. Since experimental values of B are often smaller than theoretical B values, which are taken to be equal to the negative of orbital energies using the Koopman theorem, the BEB cross sections obtained using experimental B values are usually higher (by 10%–15% at the cross section peak) than those obtained using theoretical B values. Using the experimental value for the lowest electron binding energy (=first ionization potential, IP) ensures that the cross section starts at the right threshold and also leads to improved agreement of the shape and magnitude of the BEB cross section near the threshold with known experimental cross sections.

For closed-shell molecules, we used the restricted Hartree–Fock (RHF) method with the Gaussian 6-311G basis set available in the GAMESS code. For open-shell molecules, it was found^{1,2} that the unrestricted Hartree–Fock (UHF) method produced more realistic orbital energies for valence orbitals than the restricted open-shell Hartree–Fock (ROHF) method. We used the UHF method for open-shell molecules also with the same Gaussian 6-311G(d) basis set. Although the UHF method has the disadvantage of producing a pair of orbital and kinetic energies (associated with α and β spins), the valence orbital energies from the ROHF method were often unrealistically small compared to experimental IP's, making the corresponding cross section too large. On the other hand, using separate B and U values for molecular orbitals associated with α and β spins from the UHF method, produces BEB cross sections almost identical to those produced by using the average B and U values from matching α and β molecular orbitals. The B and U values presented for open-shell molecules in Tables I and II are these average values.

In practice, the BEB model is insensitive to minor variations in the values of B and U used, except for the lowest B . Thus the use of extended basis sets other than 6-311G(d), leading to slightly altered B and U values, would not change the BEB cross section significantly. However, the BEB cross section is sensitive to the lowest B and U , as the contribution to ionization cross section from the outermost occupied molecular orbital is substantial, if not dominant. For this reason, we used the experimental first IP's for the lowest B , which are available for many molecules,¹⁴ and theoretical values for the remaining orbitals.

III. APPLICATION TO MOLECULES

In this section, we compare the BEB cross sections to the experimental values of 11 molecules. As far as we are aware, no experimental cross sections are available for GeH, GeH₂, GeH₃, GeH₄, and Ge₂H₆. Most experiments measured the “gross” ionization cross section, i.e., the total ion current rather than the number of ions. On the other hand, BEB cross sections are “counting” ionization cross sections,

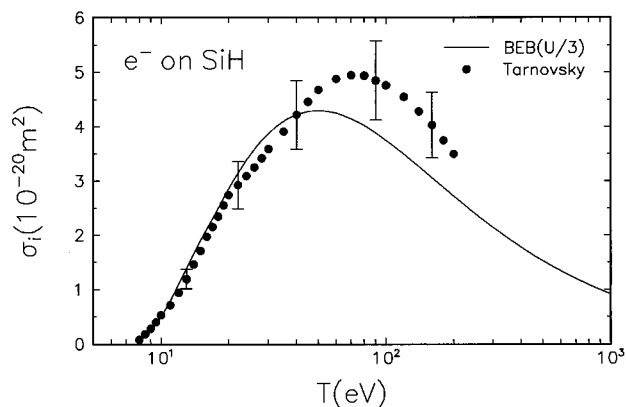


FIG. 1. Comparison of the BEB cross section to experiment for SiH. Solid curve, the present work; circles, experimental data by Tarnovsky *et al.* (Ref. 16).

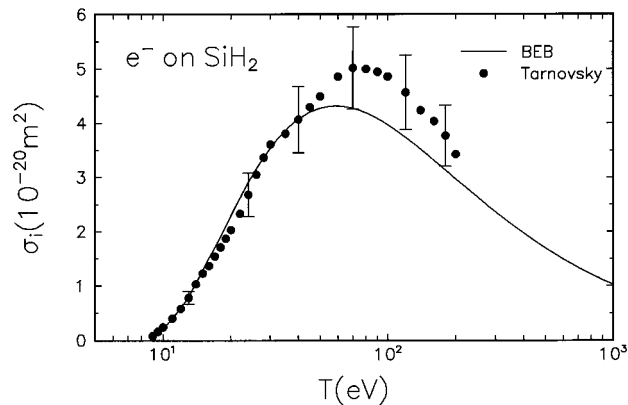


FIG. 2. Comparison of the BEB cross section to experiment for SiH₂. Solid curve, the present work; circles, experimental data by Tarnovsky *et al.* (Ref. 16).

which account for the number of singly charged ions produced. When a significant fraction of the ions produced are multiply charged, the gross ionization cross section will accordingly be larger than the corresponding counting ionization cross section at high T . Thus BEB cross sections provide lower limits to experimental gross ionization cross sections. In mass spectrometric experiments, the goal is to collect both parent molecular ions as well as all fragments at different T . However, since the BEB cross sections are simple sums of cross sections for ionizing one electron from each of the occupied molecular orbitals, the theory cannot provide relative cross sections for the production of different molecular fragments and ions. The separation of BEB cross sections into partial ionization cross sections for different molecular fragments and ions is not straightforward, as the crossing of different potential energy surfaces corresponding to diverse dissociative ionization and fragmentation channels must be taken into account. For instance, the BEB model cannot predict the paucity of SiH₄⁺ in the ionization of SiH₄ due to the preionization into SiH₂⁺ observed by Berkowitz *et al.*¹⁵ For simplicity, we compare our theoretical cross sections to the simple sum of all experimental partial cross sections that produce an ion. Nevertheless, the comparison presented here will clearly demonstrate wide applicability of our theory and suggest that the theoretical cross sections for the other molecules not yet measured should be reliable.

A. Silyldiyne (SiH), Silylene (SiH₂), and Silyl (SiH₃)

For SiH and SiH₃, we used the UHF method with the 6-311G(*d*) basis set and we found that SiH₃ has the doublet C_{3v} ground state. For SiH₂, the singlet C_{2v} structure is the ground state using the same basis set. The only experimental data for these molecules are by Tarnovsky *et al.*,¹⁶ who measured absolute electron-impact ionization cross sections for SiD_{*x*} ($x = 1-3$) from threshold to 200 eV using a fast neutral beam technique. Since the BEB cross sections are insensitive to isotope effects, we have compared their isotopic results with the BEB cross sections with hydrogen atoms. The ions collected by Tarnovsky *et al.* were: SiD⁺ and Si⁺ for SiD,

SiD₂⁺ and SiD⁺ for SiD₂, and SiD₃⁺ and SiD₂⁺ for SiD₃, and hence the sum of these partial cross sections should be a lower limit to the total ionization cross section. Experimental data for these molecules are compared to BEB cross sections in Figs. 1–3. Experimentally, the parent ionization cross section for all three targets SiD_{*x*} ($x = 1-3$), i.e., the production of SiD_{*x*}⁺, has essentially the same maximum value at 70 eV, and the parent ionization is the dominant process.

In a molecule with a heavy atom, the valence molecular orbitals are often dominated by an atomic valence orbital. In such a case, the U value of the valence orbital tends to be high due to the radial nodes, and a high U value leads to an unusually small orbital cross section. When this happens, we found it necessary to reduce the U value by dividing it by the principal quantum number of the dominant atomic orbital. Since the valence molecular orbitals for the molecules in this article are the silicon 3*p* orbital or the germanium 4*p* orbital, we divided the appropriate U values by the respective principal quantum numbers when the Mulliken population analysis indicated that the valence molecular orbitals were dominated by atomic orbitals (>90%). For instance, we reduced

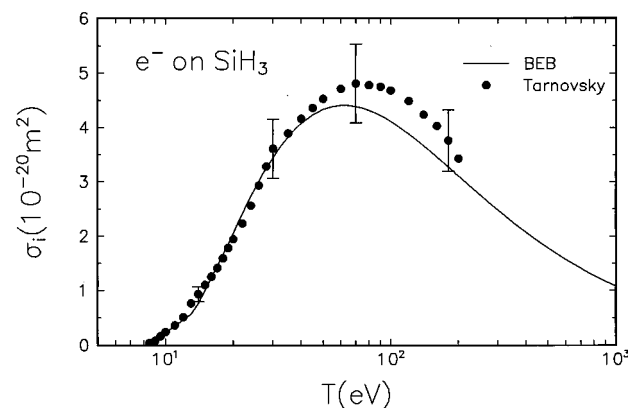


FIG. 3. Comparison of the BEB cross section to experiment for SiH₃. Solid curve, the present work; circles, experimental data by Tarnovsky *et al.* (Ref. 16).

the U value by a factor of 3 in SiH and that for GeH by a factor of 4. The BEB cross sections calculated using the reduced U values are denoted by BEB($U/3$) or BEB($U/4$), depending on the principal quantum number of the dominant atomic orbital. This increases the calculated cross section by 10%–15%.

For germanium hydrides, some inner molecular orbitals are basically $n=3$ atomic orbitals of Ge. The U values of such molecular orbitals were also divided by the principal quantum number, as indicated in Table II. This reduction increases the cross sections from inner orbitals, which in turn increases the total cross section for T beyond the cross section peak by a few percent.

- For SiH, the BEB($U/3$) cross section agrees well with the experimental results for $T < 50$ eV (Fig. 1). For $T > 50$ eV, however, the experimental values are 25%–30% higher than the theory, reaching a maximum between 70 and 80 eV, while the BEB cross section reaches a somewhat lower maximum between 50 and 55 eV.
- For SiH₂, again the BEB cross section agrees well with the experimental result for $T < 50$ eV (Fig. 2). Similar to the case of SiH, the experimental ionization cross section reaches a maximum between 70 and 80 eV, while the experimental values are 25%–30% larger than the theory for $T > 60$ eV. The BEB cross section has a maximum between 60 and 65 eV.
- For SiH₃, again the experimental ionization cross section has a maximum between 70 and 80 eV, while the BEB cross section reaches a maximum between 55 and 60 eV (Fig. 3).

The experimental data by Tarnovsky *et al.*¹⁶ for SiH and SiH₃ show large fluctuations of the order of 20%–25% above $T=50$ eV in the cross sections for the production of the parent molecular ions (SiH⁺ and SiH₃⁺), which dominate the total ionization cross section. The overall experimental accuracy claimed is $\pm 15\%$, which is indicated in Figs. 1–3. The discrepancy between theory and experiment at high T may have resulted from high- T events that have not been accounted for in the theory, such as multiple ionization.

B. Silane (SiH₄) and Disilane (Si₂H₆)

The experimental data for silane and disilane by Chatham *et al.*¹⁷ from threshold to 300 eV, and Krishnakumar and Srivastava¹⁸ from threshold to 1000 eV, are compared to BEB cross sections in Figs. 4 and 5. For SiH₄, recent experimental data by Basner *et al.*¹⁹ agree well with the data by Chatham *et al.*, and also with the BEB cross section calculated using the adiabatic IP of 11.65 eV, while the data by Krishnakumar and Srivastava are significantly lower than the theory at $T > 30$ eV. We found earlier³ that when the adiabatic and vertical IP's are different by more than 1 eV, as is the case for CH₄ and NH₃, using the vertical IP for the lowest binding energy led to a much better agreement between the BEB cross section and experiment near the threshold. However, in the case of SiH₄, the BEB cross sec-

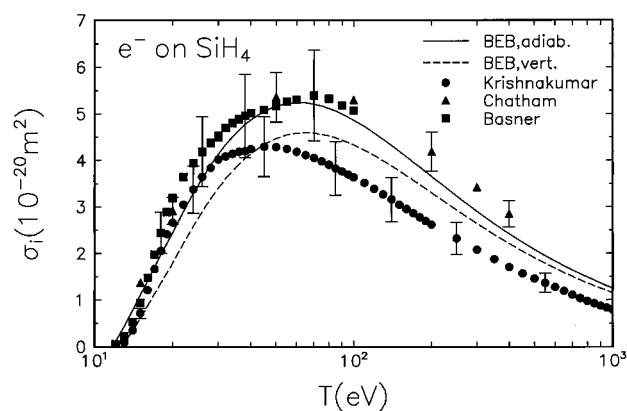


FIG. 4. Comparison of the BEB cross section to experiment for SiH₄. Solid curve, the BEB cross section with adiabatic IP; dashed curve, the BEB cross section with vertical IP; circles, experimental data by Krishnakumar and Srivastava (Ref. 18); triangles, data by Chatham *et al.* (Ref. 17); squares, data by Basner *et al.* (Ref. 19).

tion with the vertical IP of 12.7 eV deduced from the photoionization spectrum,²⁰ is significantly smaller than all experimental data near the threshold. The BEB cross section with the vertical IP is marked “BEB, vert.” in Fig. 4. We should wait until we have more examples before rendering any conclusion about the merit of using adiabatic versus vertical IP's.

For Si₂H₆, the experimental results of Chatham *et al.*¹⁷ and Krishnakumar and Srivastava¹⁸ differ markedly in the peak region (Fig. 5). The Chatham results resemble in general the shape of the BEB cross section but are higher in magnitude. However, if we reduce the Chatham data by a factor of 0.78, then the renormalized Chatham data agree well with BEB cross section in magnitude over the entire range of T . The data by Krishnakumar and Srivastava form an unusually sharp peak at $T \sim 30$ eV, unlike the shape of ionization cross sections for a large number of molecules, although their data at $T > 100$ eV are in excellent agreement with the BEB cross section.

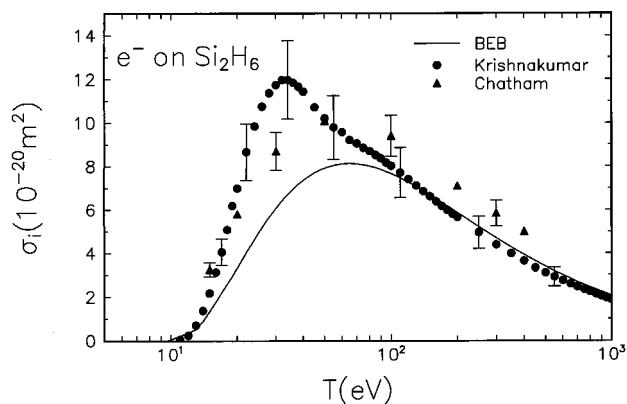


FIG. 5. Comparison of the BEB cross section to experiment for Si₂H₆. Solid curve, the present work; circles, experimental data by Krishnakumar and Srivastava (Ref. 18); triangles, data by Chatham *et al.* (Ref. 17).

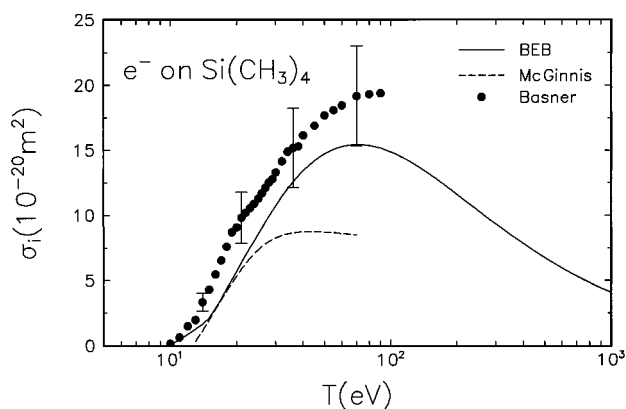


FIG. 6. Comparison of the BEB cross section to experiment for TMS, $\text{Si}(\text{CH}_3)_4$. Solid curve, the present work; dashed curve, experimental data by McGinnis *et al.* (Ref. 21); circles, data by Basner *et al.* (Ref. 22).

C. Tetramethylsilane (TMS) $\text{Si}(\text{CH}_3)_4$

Two sets of experimental data are available for this molecule: by McGinnis *et al.*²¹ from the threshold to $T = 70$ eV, and by Basner *et al.*²² from the threshold to $T = 90$ eV. The BEB cross section agrees in shape with the data by Basner *et al.*, but the theory is smaller in magnitude, while the data by McGinnis *et al.* near the peak at $T \sim 70$ eV are almost one-half of the theory (Fig. 6). The premature flattening of the data by McGinnis *et al.* at $T > 30$ eV suggests that their experiment may not have collected as many fragment ions as Basner *et al.* did. For instance, the combined cross section for ions of mass/charge ≤ 42 collected by Basner *et al.*, but not by McGinnis *et al.*, at $T = 70$ eV is $2.87 \times 10^{-20} \text{ m}^2$, which will increase the peak cross section by McGinnis *et al.* to within $\sim 20\%$ of the BEB cross section peak.

D. GeH, GeH_2 , GeH_3 , GeH_4 , Ge_2H_6

For GeH and GeH_3 , the UHF approximation was used to obtain molecular wave functions with an internally supplied double-zeta valence (DZV) basis set with one d polarization function. The binding energies and kinetic energies reported

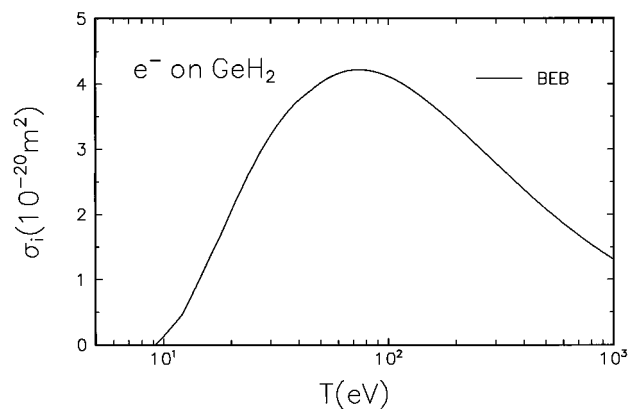


FIG. 8. The BEB cross section for GeH_2 .

in Table II are the average of the quantities of the corresponding α and β molecular orbitals characterized by the same irreducible representation. The ground state geometry of GeH_3 was found to be C_{3v} rather than D_{3h} . The singlet C_{2v} state of GeH_2 was found to be the ground state with the same basis set, i.e., DZV plus one d polarization function in the RHF approximation. For GeH_4 and Ge_2H_6 , we used the DZV basis set with three d polarization functions in the RHF approximation with the T_d and D_{2d} point groups. As far as we are aware, our BEB total ionization cross sections are the only results available for these molecules. In view of the importance of these radicals and molecules in the chemical vapor deposition of Ge, experimental verification of our results for these molecules would be worthwhile. Our BEB total ionization cross sections are shown in Figs. 7–11.

We noted earlier^{2,3} that the UHF approximation sometimes produces an unrealistically low binding energy for an unpaired valence electron. For instance, the theoretical binding energy for the 4π electron of GeH is only 4.008 eV, which is almost certainly too low. Such a low binding energy in turn generates an unrealistically large ionization cross section. Similar low theoretical binding energy was also obtained for SiH, but in this case an experimental IP was available. Since no experimental IP is known for GeH, we

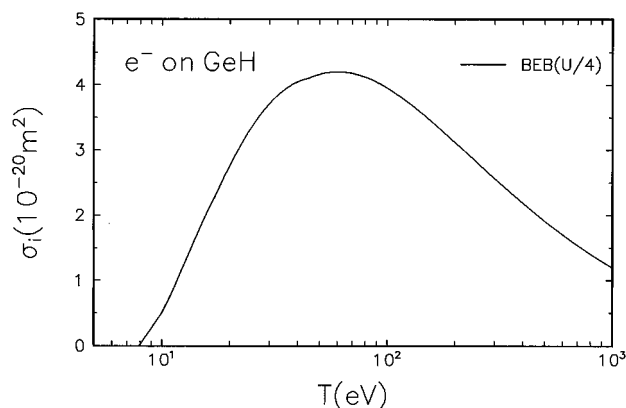


FIG. 7. The BEB cross section for GeH.

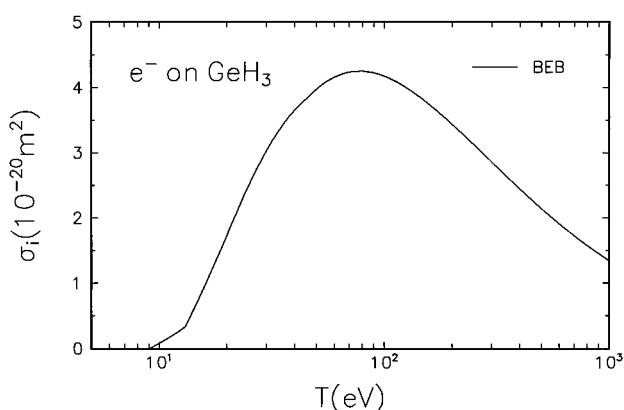


FIG. 9. The BEB cross section for GeH_3 .

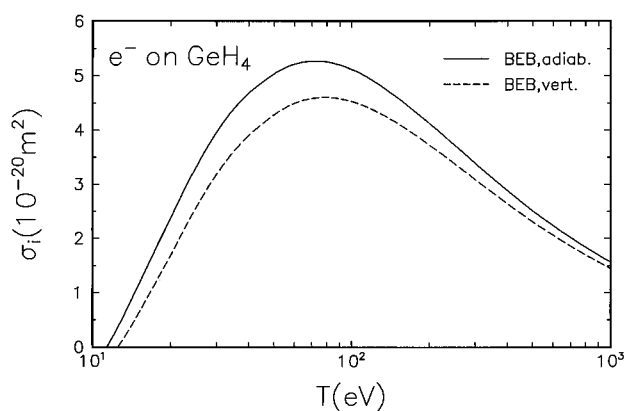


FIG. 10. The BEB cross section for GeH_4 . Solid curve, the BEB cross section with adiabatic IP; dashed curve, the BEB cross section with vertical IP.

obtained an “estimated” binding energy of its 4π electron by multiplying the ratio of the experimental to theoretical binding energies of SiH to the theoretical binding energy of GeH , which put the 4π binding energy at 7.906 eV.

IV. CONCLUSION

For the molecules where experimental ionization cross sections are available, the BEB cross sections are in good agreement with experiments at low incident energies, $T < 50$ eV. At higher T , however, BEB cross sections tend to be lower than the available experimental data. This discrepancy may have come from wider fluctuations in the experimental data at higher T than the general error bounds claimed by the experimenters, or some unidentified high- T events, such as multiple ionization, that have not been accounted for in the theory. For unstable molecules and radicals, when only one set of measurements is available, additional independent measurements would be desirable.

When vertical and adiabatic IP's are different by more than 1 eV, such as CH_4 and NH_3 , often as a result of the Jahn–Teller effect, we found that the use of vertical IP's significantly improved the agreement between BEB cross

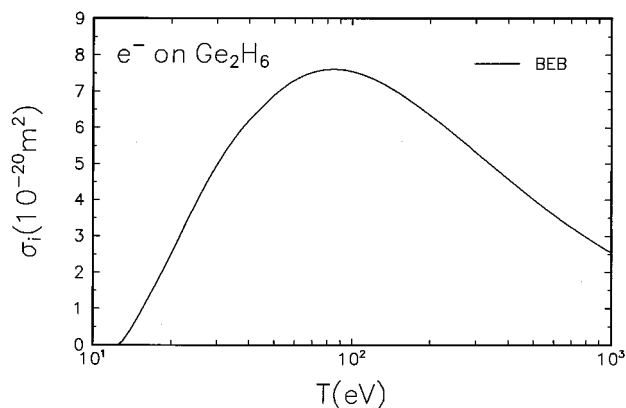


FIG. 11. The BEB cross section for Ge_2H_6 .

TABLE I. Molecular orbitals, electron binding energy B in eV, kinetic energy U in eV, and electron occupation number N for SiH , SiH_2 , SiH_3 , SiH_4 , Si_2H_6 and $\text{Si}(\text{CH}_3)_4$. All B and U values are theoretical, except that those marked by an asterisk are experimental adiabatic IP's.

Molecule	MO	B	U	N
$\text{SiH } ^2\Pi$ average of α and β orbital values	2σ	167.91	360.49	2
	1π	116.24	331.62	4
	3σ	116.23	332.69	2
	4σ	17.457	31.174	2
	5σ	10.482	28.523	2
	6σ	18.75	58.81	2
$\text{SiH}_2 C_{2v} ^1A_1$	2π (Si $3p$)	7.89*	21.287/3	1
	$2a_1$	167.24	360.26	2
	$1b_2$	115.63	330.80	2
	$1b_1$	115.59	332.16	2
	$3a_1$	115.57	332.72	2
	$4a_1$	18.414	29.633	2
$\text{SiH}_3 C_{3v} ^2A_1$	$2b_1$	12.582	22.138	2
	$5a_1$	8.92*	28.095	2
	$2a_1$	166.88	359.80	2
	$3a_1$	115.35	331.76	2
	$1e$	115.35	331.92	4
	$4a_1$	19.253	29.560	2
$\text{SiH}_4 T_d ^1A_1$	$2e$	13.304	22.883	4
	$5a_1$	8.14*	26.787	1
	$2a_1$	166.74	359.73	2
	$1f_2$	115.20	331.82	6
	$3a_1$	19.921	28.986	2
	$2f_2$	11.65*	22.782	6
$\text{Si}_2\text{H}_6 D_{2d} ^1A_1$	$2a_{1g}$	166.71	359.68	2
	$2a_{2u}$	166.71	359.78	2
	$1e_u$	115.19	331.79	4
	$3a_{1g}$	115.19	331.68	2
	$1e_g$	115.19	331.83	4
	$3a_{2u}$	115.20	332.09	2
	$4a_{1g}$	20.604	29.637	2
	$4a_{2u}$	18.536	29.153	2
	$2e_u$	13.619	22.212	4
	$2e_g$	12.912	23.100	4
	$5a_{1g}$	9.7*	29.580	2
	$1f_2$	304.62	435.95	6
$\text{Si}(\text{CH}_3)_4 T_d ^1A_1$ (TMS)	$2a_1$	304.62	435.96	2
	$3a_1$	165.93	360.08	2
	$2f_2$	114.41	331.91	6
	$4a_1$	26.477	32.346	2
	$3f_2$	25.070	33.786	6
	$5a_1$	17.007	35.471	2
	$4f_2$	15.407	24.504	6
	$1e$	14.925	25.177	4
	$1f_1$	14.376	26.810	6
	$5f_2$	9.8*	34.868	6

sections and experimental data between the threshold and the peak.² However, we see (Fig. 4) that this is not the case for SiH_4 , whose vertical IP is 12.7 eV,²⁰ while the adiabatic IP is 11.65 eV. One feature that distinguishes SiH_4 from CH_4 and NH_3 is that few parent ions, SiH_4^+ , are generated in the former,^{15,17} while the parent ions of the latter are not difficult to produce.

Intuitively, vertical IP's appear to be the right ones to use in electron-impact ionization because the ionization takes place faster than the time a molecule takes to rearrange its ionic configuration. Vertical IP's are also used in interpreting photoionization events. A comparison of the ionization cross

TABLE II. Molecular orbitals, electron binding energy B in eV, kinetic energy U in eV, and electron occupation number N for GeH, GeH₂, GeH₃, GeH₄, and Ge₂H₆. All B and U values are theoretical, except that those marked by an asterisk are experimental adiabatic IP's.

Molecule	MO	B	U	N
GeH $^2\Pi$	4σ	193.34	587.95/3	2
	2π	141.12	547.90/3	4
	5σ	140.99	548.13/3	2
	1δ	45.160	409.88/3	4
	3π	44.975	409.68/3	4
	6σ	44.975	408.74/3	2
	7σ	17.434	44.324	2
	8σ	10.388	35.737	2
GeH ₂ C_{2v} 1A_1	4π (Si $4p$)	7.9068	26.370/4	1
	$4a_1$	192.80	587.75/3	2
	$2b_2$	140.70	547.69/3	2
	$2b_1$	140.52	548.02/3	2
	$5a_1$	140.45	548.18/3	2
	$6a_1$	44.744	409.25/3	2
	$1a_2$	44.654	409.66/3	2
	$3b_2$	44.567	409.65/3	2
	$3b_1$	44.414	408.14/3	2
	$7a_1$	44.308	409.42/3	2
	$8a_1$	18.357	41.292	2
	$4b_1$	12.063	27.228	2
	$9a_1$	9.132	35.813	2
	$4a_1$	192.73	587.37/3	2
	$5a_1$	140.56	548.09/3	2
	$2e$	140.55	548.04/3	4
GeH ₃ C_{3v} 2A_1	$6a_1$	44.538	409.09/3	2
	$3e$	44.578	408.04/3	4
	$4e$	44.477	409.30/3	4
	$7a_1$	19.343	40.919	2
	$5e$	12.957	28.568	4
	$8a_1$	9.029	34.153	1
	$3a_1$	192.92	575.24/3	2
	$2f_2$	140.34	551.86/3	6
	$3f_2$	44.346	407.41/3	6
	$1e$	44.238	409.24/3	4
GeH ₄ T_d 1A_1	$4a_1$	20.071	39.214	2
	$4f_2$	11.33*	29.609	6
	$4a_{1g}$	192.94	575.25/3	2
	$4a_{2u}$	192.93	575.34/3	2
	$2e_u$	140.35	551.93/3	4
	$2e_g$	140.35	551.86/3	4
Ge ₂ H ₆ D_{2d} 1A_1	$5a_{1g}$	140.35	551.83/3	2
	$5a_{2u}$	140.34	551.64/3	2
	$6a_{1g}$	44.433	405.82/3	2
	$3e_u$	44.379	406.93/3	4
	$3e_g$	44.352	407.56/3	4
	$4e_u$	44.276	408.78/3	4
	$4e_g$	44.224	410.05/3	4
	$6a_{2u}$	44.202	411.02/3	2
	$7a_{1g}$	20.773	40.906/3	2
	$7a_{2u}$	18.645	39.843	2
	$5e_u$	13.334	28.632	4
	$5e_g$	12.553	30.408	4
	$8a_{1g}$	12.5*	39.824	2

section of GeH₄ to the BEB cross section may provide further clues on the question of whether to use vertical or adiabatic IP's. For this reason, we presented two theoretical cross

sections for GeH₄ in Fig. 10. The BEB cross section obtained by using the adiabatic IP (11.33 eV) is marked "BEB, adiab.," while the curve marked "BEB, vert." was obtained by using the vertical IP (12.5 eV). Experimental vertical IP's are often not available, although adiabatic IP's are available for many molecules and radicals.

The results for germanium hydrides are new and no experimental results are available for comparison.

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