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# Infra-Red and Raman Spectra of Polyatomic Molecules. III. Ethane

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The Raman and infra-red spectra of ethane have been re-examined. The 993, 2900, and 2955 cm<sup>-1</sup> Raman lines are polarized; there were indications of very faint lines appearing at 620 and 820 cm<sup>-1</sup>. The strong infra-red peak at 1480 cm<sup>-1</sup> is found to be a single minimum; the weak band at 1740 cm<sup>-1</sup> is confirmed.

The available Raman and infra-red data are evaluated, and possible assignments discussed. It is shown that the observed frequencies do not require a 740 cm<sup>-1</sup> fundamental, but are consistent with an "uncertain frequency" of about 1100 cm<sup>-1</sup> and a high restricting potential.

R ECENT work<sup>1</sup> on the problem of internal rotation in ethane has focussed considerable interest on the vibrational frequencies of this molecule. Although the magnitudes of most of the fundamental frequencies are known sufficiently well for thermodynamic calculations, two widely different values have been proposed for the socalled "uncertain frequency." Heat capacity data,<sup>2</sup> including that at low temperatures, seem unequivocally to indicate a value of about 1100 cm<sup>-1</sup> for this frequency; on the other hand, Bartholomé and Karweil<sup>3</sup> have studied the overtone region carefully and find spectral evidence for a value of 740 cm<sup>-1</sup>. In connection with this question, we have considered it worth while to reexamine the Raman and infra-red spectra of ethane, and have also considered possible assignments for the observed frequencies.

#### EXPERIMENTAL

The ethane used was especially prepared from pure hydrogen and ethylene. Both materials were obtained from the Ohio Chemical Co.; the hydrogen was rated 99.9 percent, and the ethylene 99.5 percent. Catalytic combination took place on reduced copper oxide at temperatures not exceeding 120°C. The Raman spectrum showed no traces of the strong lines of the probable impurities.

The Raman spectrograph has been previously described; during exposures, the Raman tube containing the liquid ethane was kept at a temperature of about  $-60^{\circ}$ C by a stream of cold air. Polarization measurements<sup>5</sup> on the four principal lines showed the 993, 2900, and 2955 cm<sup>-1</sup> lines to be polarized, and the 1460 cm<sup>-1</sup> line, depolarized. Microphotometer records of comparable plates made with polarized and unpolarized light are given in Fig. 1. It will be noted that the 2900 cm<sup>-1</sup> line appears to be slightly more polarized than the 2955 cm<sup>-1</sup> line.

The Raman spectrum of ethane has been examined several times with instruments of higher dispersion; it does not seem worth while to report our frequencies for the lines previously found, which checked within experimental error.

On several plates, indications were found of two faint lines, difficult to measure, displaced from the Hg 4358A line by about 620 and 820 cm<sup>-1</sup>. Professor George Glockler has informed us that these lines had appeared on his plates also,<sup>7</sup> and were measured at 617 and 813 cm<sup>-1</sup>. These lines are extremely faint, and moreover the corresponding shifts excited by Hg 4047A were not observed; nevertheless, the appearance of these lines in two independent investigations seems to provide some evidence for their reality.

<sup>\*</sup> National Research Fellow in Chemistry.

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¹ G. B. Kistiakowsky and E. B. Wilson, Jr., J. Am. Chem. Soc. 60, 494 (1938); K. S. Pitzer and J. D. Kemp, ibid. 60, 1515 (1938); and references given therein.

² W. Hunsmann, Zeits. f. physik. Chemie B39, 23 (1938); G. B. Kistiakowsky, J. R. Lacher, and F. Stitt, J. Chem. Phys. 6, 407 (1938); E. B. Wilson, Jr., ibid. 6, 408 (1938).

³ E. Bartholomé and J. Karweil, Zeits. f. physik. Chemie R30 1 (1938).

B39, 1 (1938).

<sup>&</sup>lt;sup>4</sup> H. Gershinowitz and E. B. Wilson, Jr., J. Chem. Phys. **6**, 247 (1938).
<sup>5</sup> J. T. Edsall and E. B. Wilson, Jr., J. Chem. Phys. **6**, 124

<sup>(1938).</sup> 

<sup>&</sup>lt;sup>6</sup> P. Daure, Ann. de physique **12**, 375 (1929); S. Bhagavantam, Ind. J. Phys. **6**, 596 (1931); C. M. Lewis and W. V. Houston, Phys. Rev. **44**, 903 (1933); G. Glockler and M. M. Renfrew, J. Chem. Phys. **6**, 295 (1938).

We are indebted to Professor Glockler for sending us these data; he had not reported them because on his plates, as on ours, these shifts excited by the Hg 4047A line did not appear.

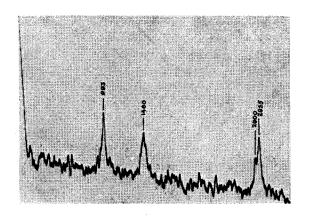


Fig. 1 (a). Microphotometer records of plates made with polarized light, showing the shifts at 993 (P), 1460 (D), 2900 (P), and 2955 (P) cm<sup>-1</sup>. To conform with Table I the lines are labeled with the corresponding gas-phase frequencies instead of the actual liquid-phase frequencies.

The infra-red absorption of gaseous ethane was studied with the spectrometer previously described;8 a fluorite prism was used for the region from 3 to  $9\mu$ , and a rocksalt prism from 9 to  $15\mu$ . Transmission curves of the absorption regions beyond  $5\mu$  are given in Fig. 2.

The infra-red spectrum has been studied by several investigators.9 The band at 827 cm<sup>-1</sup> is well established. The 1740 cm<sup>-1</sup> band was reported by Benedict, Morikawa, Barnes and Taylor as lying at 1770 cm<sup>-1</sup>; we believe our value of the frequency to be more reliable than theirs, since they used a rocksalt prism, which gives lower resolution.

Our results for the  $7\mu$  region are of particular interest. This region was examined under high resolution by Levin and Meyer; their measurements show a peak at 1379 cm<sup>-1</sup>, a region of weak absorption between 1390 and 1430 cm<sup>-1</sup>, and a strong band with center about 1480 cm<sup>-1</sup> which shows regular rotational spacing from 1442 to 1515 cm<sup>-1</sup>, and gives every appearance of a single band of the perpendicular type. Bartholomé and Sachsse, using a fluorite prism, found a single weak minimum at 1365 cm<sup>-1</sup>, and found the strong peak to be double, with minima at

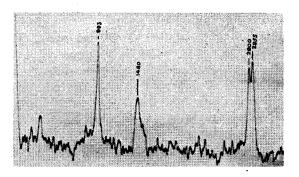


Fig. 1 (b). Microphotometer records of plates made with unpolarized light, showing the shifts at 993 (P), 1460 (D), 2900 (P), and 2955 (P) cm<sup>-1</sup>. To conform with Table I the lines are labeled with the corresponding gasphase frequencies instead of the actual liquid-phase fre-

1465 and 1495 cm<sup>-1</sup>. Our curve (Fig. 2) shows a single strong minimum at 1477 cm<sup>-1</sup>, and two weaker minima at 1380 and 1414 cm<sup>-1</sup>. The resolution of our spectrometer in this region is more than adequate; in studies of cyclopropane, two minima at 1425 and 1445 cm<sup>-1</sup> are clearly resolved. It is difficult to explain the double minimum found by Bartholomé and Sachsse; it is possible that they did not properly correct for the strong atmospheric water bands which overlie this region. Considering our own data and those of Levin and Meyer, we believe that the evidence is strongly in favor of a single strong band at 1480 cm<sup>-1</sup>, with weaker bands at 1380 and 1415 cm<sup>-1</sup>, as shown in our curve.

Our results above 2000 cm<sup>-1</sup> are consistent with the grating measurements of Bartholomé and Karweil in this region, except for one very large discrepancy. Their curves show a band at 2303 cm<sup>-1</sup>, while we find this band at 2353 cm<sup>-1</sup>, as shown in Fig. 2. While our resolution was con-

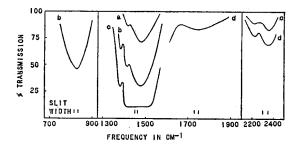


Fig. 2. Infra-red transmission of ethane gas. Path length, 30 cm; pressures, (a) 60 mm, (b) 125 mm, (c) 355 mm, (d) 760 mm.

<sup>&</sup>lt;sup>8</sup> H. Gershinowitz and E. B. Wilson, Jr., J. Chem. Phys.

<sup>6, 197 (1938).

9</sup> A. Levin and C. F. Meyer, J. Opt. Soc. Am. 16, 137 (1928); E. Bartholomé and H. Sachsse, Zeits. f. physik. Chemie B30, 40 (1935); W. S. Benedict, K. Morikawa, R. B. Barnes, and H. S. Taylor, J. Chem. Phys. 5, 1 (1937); E. Bartholomé and J. Karweil, Zeits. f. physik. Chemie **B39**, 1 (1938).

siderably less than that of Bartholomé and Karweil, this difference of 50 cm<sup>-1</sup> lies far outside the experimental error.

#### Discussion

After a study of the various investigations referred to above, we have arrived at the observed frequencies given in Table I as being the most complete and reliable list available at present. The Raman frequencies are principally the gasphase values of Lewis and Houston; we have rejected the weak 1344 cm<sup>-1</sup> line, which is almost certainly due to ethylene present as an impurity, and have added the 620 and 820 cm<sup>-1</sup> lines discussed above, the 1463 and 2963 cm<sup>-1</sup> shifts observed in the liquid by Glockler and Renfrew, and the 975 cm<sup>-1</sup> shift found by Bhagavantam. Of the infra-red frequencies, those lying below

2000 cm<sup>-1</sup> have been discussed above. The higher frequencies are taken from Bartholomé and Karweil; we have preferred to regard their peak at 2230 cm<sup>-1</sup> as a single perpendicular band rather than as two bands 15 cm<sup>-1</sup> apart. With regard to the discrepancy mentioned above, we have included their value of 2303 cm<sup>-1</sup> as well as our own value, 2350 cm<sup>-1</sup>.

Glockler and Renfrew<sup>10</sup> have recently suggested that the 813 and 1463 cm<sup>-1</sup> Raman lines found in liquid ethane may be the infra-red fundamentals at 827 and 1480 cm<sup>-1</sup>, appearing in the Raman effect through a breakdown of the selection rules due to intermolecular forces active in the liquid. These frequencies would actually be permitted in the Raman effect if the selection

<sup>10</sup> G. Glockler and M. M. Renfrew, J. Chem. Phys. 6, 409

TABLE	T	Observed	frequencies	οf	ethane	with a	possible	assignment.*
LADLE	1.	Coscioca	HEGINETICIES	$v_I$	condine.	$\omega u u u a$	20331010	assignment.

	Int'y.; Type	Ass.	CALC. FREQ.	Symmetry		
FREQ. cm-1				$D_{3h}$	$D_{3d}$	
Infra-Red: 827	⊥s	Fund.	(827)	$E' \\ A_2'' \\ E' \\ E' \\ E' \\ E' \\ A_1' + A_2' + E'$	$E_u$	
1380	w ~	Fund.	(1380)	$A_2^{\prime\prime}$	$egin{array}{c} A_{2u}^{2} \ E_{u} \ E_{u} \ E_{u} \end{array}$	
1414	⊥w	1120 + 300	1420	<i>E'</i>	$E_u$	
1480	⊥s ?w	Fund.	(1480)	E'	$E_u$	
1740	/ ?w	1460+300	1760	E'	$E_u$	
2230	?⊥w	$\begin{cases} 1120^2 \\ 1275 + 827 \end{cases}$	2240	$A_1'+A_2'+E'$	$\overline{E_u}$	
		$\begin{array}{c} 1375 + 827 \\ 827 + 1480 \end{array}$	2202		$E_u$	
2303	?][w	827 + 1460	2307	$A_1'+A_2'+E'$	4 1 4 1 5	
2350	?  w	993+1380	2287 2373	<u> </u>	$A_{1u}+A_{2u}+E_u$	
2660	±w	2970 - 300	2670	$A_2^{\prime\prime} \ E^{\prime}$	$\stackrel{A}{E}_{u}^{2u}$	
2755	l lm	1380 + 1375	2755	4."	$A_{2u}$	
2895	m -	Fund.	2926	$A_2^{\prime\prime} \\ A_2^{\prime\prime}$	$\left( egin{array}{c} A_{2u} \ A_{2u} \end{array}  ight)$	
2955	ls	$\{1480+1460$	2940	$A_1'' + A_2'' + E''$	$A_{1u}+A_{2u}+E_u$ resonance	
3007	m	993 + 827 + 1120	2940	$A_1'' + A_2'' + E''$ $A_1'' + A_2'' + E''$	$A_{1u} + A_{2u} + E_u$ triplet	
2950-3000	l ïw l	Fund.	2980	E'	$E_u$	
3100-3150	⊥w	$827^2 + 1480$	3134	$A_{1}' + A_{2}' + 2E'$	$\begin{array}{c} E_u \\ A_{1u} + A_{2u} + 2E_u \end{array}$	
3186	w	$\int 1460^2 + 300$	3220	$A_1' + A_2' + 2E'$ $A_1'' + A_2'' + E''$	$A_{1u}+A_{2u}+E_u$ resonance	
3223	w	$1120+1480+300^{2}$	3200	$A_1^{\prime\prime}+A_2^{\prime\prime}+E^{\prime\prime}$	$A_{1u}+A_{2u}+E_u\int doublet$	
Raman:	1 1	2000		_		
(620)?	v.w.	$300^{2}$	600	${A_1'} \ E'$	$\stackrel{A_{1arrho}}{E_u\dagger}$	
(813)	v.w.	Fund. $C^{12}-C^{13}$ Fund.	(827)	E'	$E_u\dagger$	
975 993	w	Fund.	(002)	$A_{1}'$	$A_{1g}$	
993	s, P	rund.	(993)	$A_1$	$rac{A_{1q}}{E_{q}}$	
(1463)	m, D	Fund.	{ (1460) 1480	$A_1' \ A_1' \ E'' \ E'$	$E_{m{o}} \ E_{m{u}}^{m{t}}$	
2744	w	( 1375 <sup>2</sup>	2750	$A_{1}'$	$A_{1g}$ resonance	
2778	w	13802	2760	$A_1'$	$A_{1g}$ doublet	
2900	s, P	Fund.	2927	$A_1'$	$A_{1g}$ resonance	
2955	s, P	\ 1460 <sup>2</sup>	2920	$A_{1}'$	$A_{1q}$ doublet	
2940	w	∫ <b>1480+146</b> 0	2940	$A_1'' + A_2'' + E''$		
	, "	14802	2960	· —	$A_{1g} + A_{2g} + E_g$ \ resonance	
(2963)	w	Fund.	2970	E''	$E_g$ $\int$ doublet in $D_{3d}$ case	

<sup>\*</sup> Liquid-phase frequencies in parentheses: s = strong, m = medium, w = weak, v.w. = very weak, P = polarized, D = depolarized. † Breakdown of  $D_{ad}$  selection rules.

rules appropriate to the  $D_{3h}$  configuration apply to ethane.<sup>11</sup>

If the 620 cm<sup>-1</sup> line is real, there is the interesting possibility that it is the first overtone of the internal torsional motion associated with a high potential barrier. The selection rules permit the appearance of this overtone in the Raman effect. Professor E. B. Wilson, Jr., has calculated certain levels of this motion; <sup>12</sup> for a restricting potential of 3000 cal./mole magnitude and the usually assumed cosine shape, the two lowest excited levels lie at about 300 and 565 cm<sup>-1</sup>. While this value for the overtone is too low to fit the 620 cm<sup>-1</sup> Raman line, it is quite possible that the actual potential barrier has a different shape which would give a somewhat higher second excited level when adjusted to fit the thermodynamic data, an idea which is confirmed by a study of a square potential barrier. Another possible assignment for the 620 cm<sup>-1</sup> line would be the difference tone 1463-813 (liquid-phase frequencies); this calculated value lies 30 cm<sup>-1</sup> too high.

Since a study of the infra-red and Raman spectra of heavy ethane, C<sub>2</sub>D<sub>6</sub>, is now in progress in this laboratory, it seems best to postpone any final vibrational analysis until data for the two isotopic molecules may be considered together. However, it may be well to point out that the overtone region, studied by Bartholomé and Karweil, does not afford conclusive evidence on the magnitude of the "uncertain frequency." These authors account for their observed frequencies with the use of a 740 cm<sup>-1</sup> fundamental and the selection rules appropriate to the  $D_{3h}$ configuration. The double minimum, 1465 and 1495 cm<sup>-1</sup>, reported by Bartholomé and Sachsse, would suggest a 740 cm<sup>-1</sup> fundamental; as discussed above, however, the evidence seems to be in favor of a single peak at 1480 cm<sup>-1</sup>. When we accept this single peak, it is possible to account for the observed frequencies without the use of a 740 cm<sup>-1</sup> frequency. To illustrate this, we give in

TABLE II. A possible set of fundamental frequencies for ethane.

(Symmetry symbols are given for point groups $D_{3h}',^{11}$ $D_{3h},$ and $D_{3d}$ , in that order.)					
$A_1, A_1', A_{1g}$ : 993	$E, E', E_u$ : 827				
(1375)	1480				
2927	2980				
$\bar{A}_2, A_2'', A_{2u}$ : 1380	$\bar{E}, E'', E_g$ :(1120)				
2926	1460				
$\bar{A}_1, A_1'', A_{1u}$ : (300)	2970				

Table I possible assignments of the observed frequencies, using the fundamentals listed in Table II. These fundamentals are consistent with the existence of a 3000 cal./mole restricting potential, an "uncertain frequency" of about  $1100 \, \mathrm{cm^{-1}}$ , and an internal torsional oscillation of about  $300 \, \mathrm{cm^{-1}}$ . We have given assignments satisfying the selection rules for both the  $D_{3h}$  and the  $D_{3d}$  equilibrium configurations; the  $D_{3h}$  assignment is slightly more satisfactory in the three cases where different assignments are required. It is to be emphasized that this assignment is not to be regarded as a final one, but is given merely to show that the observed overtones do not require a 740 cm<sup>-1</sup> fundamental.

In concluding, we should like to express our gratitude to Professor E. B. Wilson, Jr., who suggested this project and has constantly contributed to its progress.\*

The Raman work of Goubeau and Karweil shows four shifts differing significantly from previous work—786, 1335, 1491, and 2791 cm<sup>-1</sup>. The values 786 and 2791 cm<sup>-1</sup> deviate from the measurements of Glockler and Renfrew further than one would expect. The evidence that the 1335 cm<sup>-1</sup> line arises from ethane seems insufficient. The 1491 cm<sup>-1</sup> line may well be the second degenerate fundamental of this magnitude (see Table I).

The analysis now given by Karweil and Schäfer, which differs from the former analysis of Bartholomé and Karweil, is based on essentially the same set of fundamentals as ours. However, we feel that the available evidence is insufficient to exclude the  $D_{3d}$  equilibrium configuration, contrary to the conclusion of Karweil and Schäfer.

<sup>&</sup>lt;sup>11</sup> J. B. Howard, J. Chem. Phys. **5**, 442 (1937). <sup>12</sup> E. B. Wilson, Jr., J. Chem. Phys. **6**, 408 (1938).

<sup>\*</sup> Note added in proof: Since this paper was submitted, two papers of the same series as that of Bartholomé and Karweil<sup>3</sup> have been received, one by Goubeau and Karweil (Zeits. f. physik. Chemie **B40**, 376 (1938)) on the Raman effect, and one by Karweil and Schäfer (ibid., **B40**, 382) on the assignment.