

## Vibration Spectra of Cis and Trans C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub> and C<sub>2</sub>D<sub>2</sub>Cl<sub>2</sub>

H. J. Bernstein and D. A. Ramsay

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*J. Chem. Phys.* **5**, 392 (1937); 10.1063/1.1750046

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Vibration Spectra of *Cis* and *Trans*  $C_2H_2Cl_2$  and  $C_2D_2Cl_2$ 

H. J. BERNSTEIN AND D. A. RAMSAY  
National Research Council, Ottawa, Canada  
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The infra-red spectra of gaseous and liquid *cis* and *trans*  $C_2H_2Cl_2$  and  $C_2D_2Cl_2$  have been obtained in the region 3500 to 400  $cm^{-1}$ .

The fine structure of a  $\perp$  band in *trans*  $C_2H_2Cl_2$  was also resolved and the least moment of inertia obtained.

An assignment of all the fundamentals for the four molecules has been made.

Calculations of potential functions and of the equilibrium constants for the reaction *cis*  $\rightleftharpoons$  *trans* support the assignment.

## INTRODUCTION

ALTHOUGH *cis* and *trans* dichloroethylene are relatively simple planar molecules with symmetries  $C_{2v}$  and  $C_{2h}$ , respectively, considerable disagreement exists between the assignments by various authors<sup>1-4</sup> of the twelve fundamental frequencies for each of these molecules. This disagreement is largely caused by the lack of reliable infra-red data, particularly on the contours of the infra-red bands in the vapor. We have therefore re-investigated the infra-red spectra of these molecules and their deuterium analogues. Using these data and the existing Raman data, we have been able to make satisfactory assignments of the twelve frequencies in the *cis* molecule and ten frequencies in the *trans* molecule. The two remaining frequencies in the *trans* molecule are active in the infra-red, but lie beyond the range of a KBr prism. These frequencies have been obtained from combination tones.

Our values for the low frequencies are in agreement with the data of Wood and Stevenson<sup>5</sup> for the equilibrium constants of the *cis*  $\rightleftharpoons$  *trans* reaction at different temperatures. These authors deduced from their results that the sum of the two low frequencies must lie between 420 and 450  $cm^{-1}$ .

We have found it necessary, however, to revise the potential function calculation of Pitzer and Freeman<sup>4</sup> for the out-of-plane vibrations of *cis*- and *trans*-dichloroethylene, since these authors used an erroneous value for the *trans* wagging frequency and assumed that the 620  $cm^{-1}$  band reported by Wu<sup>2</sup> for the *cis* molecule was a combination of the Raman frequency at 350  $cm^{-1}$  with the torsion frequency at 270  $cm^{-1}$ . We have shown, however, that the 620  $cm^{-1}$  band was caused by an impurity.

<sup>1</sup> G. Herzberg, *Infrared and Raman Spectra of Polyatomic Molecules* (D. Van Nostrand Company, Inc., New York, 1945), p. 330.

<sup>2</sup> Ta-You Wu, J. Chem. Phys. **5**, 392 (1937); Phys. Rev. **46**, 465 (1934).

<sup>3</sup> B. Trumpy, Zeits. f. Physik. **98**, 672 (1935).

<sup>4</sup> K. S. Pitzer and N. K. Freeman, J. Chem. Phys. **14**, 586 (1946).

<sup>5</sup> R. E. Wood and D. P. Stevenson, J. Am. Chem. Soc. **63**, 1650 (1941).

## MATERIALS

A mixture of *cis* and *trans*  $C_2H_2Cl_2$  was prepared from Eastman Kodak sym  $C_2H_4Cl_2$  by reaction with Zn, and the isomers separated by fractional distillation.

*trans* b.p. =  $48.1 \pm 0.1^\circ C$ .

*cis* b.p. =  $60.1 \pm 0.1^\circ C$ .

The two samples were judged from their spectra to be about 99 percent pure.

A mixture of *cis* and *trans*  $C_2D_2Cl_2$  was prepared in the following manner:

1.  $CaC_2 + D_2O = CaO + C_2D_2$ ,
2.  $C_2D_2 + 2Cl_2 = C_2D_2Cl_4$ ,
3.  $C_2D_2Cl_4 + Zn = C_2D_2Cl_2 + ZnCl_2$ ,

and the isomers were separated by fractional distillation.

The spectra indicated that these samples were about 99 percent deuterated. Owing to the small quantity of material available, it was not possible to obtain as complete a separation as for the "light" isomers. The *cis*  $C_2D_2Cl_2$  contained about 2 percent of *trans*  $C_2D_2Cl_2$  as impurity and the *trans*  $C_2D_2Cl_2$  contained about 10 percent of *cis*  $C_2D_2Cl_2$  as impurity.

## EXPERIMENTAL

The infra-red spectra of these four molecules were investigated using a Perkin Elmer Model 12B infra-red spectrometer with a G-M amplifier and a Brown Pen Recorder. NaCl and KBr prisms were used for the regions 3500–700  $cm^{-1}$  and 700–400  $cm^{-1}$ , respectively. Vapor spectra were investigated using a 10-cm absorption cell, various pressures being used to resolve all the bands observed; liquid spectra were obtained using 0.1-mm and 0.025-mm absorption cells, the more intense bands being resolved using dilute solutions in  $CS_2$  and  $CCl_4$ .

The infra-red spectra are reproduced in Figs. 1–4 and the infra-red and Raman frequencies given in Tables I–IV.

ASSIGNMENT OF FREQUENCIES<sup>6</sup>(a) *Cis*  $C_2H_2Cl_2$ 

The five strong polarized lines in the Raman spectrum may be assigned to the totally symmetric  $A_1$  vibrations. These frequencies are also active in the infra-red though the  $1183\text{-cm}^{-1}$  band is extremely weak.<sup>7</sup> The two depolarized lines at  $406\text{ cm}^{-1}$  and  $876\text{ cm}^{-1}$  which do not have corresponding infra-red absorptions may be assigned to the two  $A_2$  vibrations since these vibrations are active only in the Raman spectrum. The five remaining vibrations give strong absorptions in the infra-red.

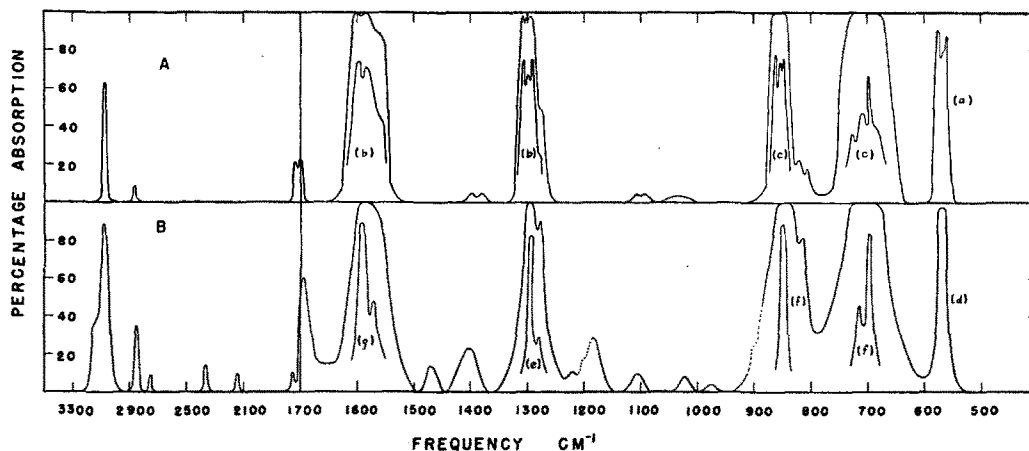
The contours of the infra-red bands may be predicted qualitatively from the curves of Badger and Zumwalt.<sup>8</sup> For vibrations with change of electric moment parallel to the greatest axis of inertia (type  $B_2$ ) the bands should exhibit a  $P, Q, R$  structure with a strong central  $Q$  branch. The  $697\text{-cm}^{-1}$  band may be assigned therefore to the lone vibration of this type,  $\nu_{12}$ . For changes of electric moment parallel to the middle axis of inertia (type  $A_1$ ) a doublet structure is expected. This is in agreement with the above assignment of frequencies to this type, though for the  $714\text{-cm}^{-1}$  band the doublet structure is overlapped by the rotational contour of the  $697\text{-cm}^{-1}$  band. Finally for changes of electric moment parallel to the least axis of inertia (type  $B_1$ ),

the bands should exhibit  $P, Q$ , and  $R$  branches of similar intensity. Such contours are observed for the  $571, 857$ , and  $1303\text{ cm}^{-1}$  bands. The fourth frequency of this type, the  $C-H$  stretching frequency, corresponds to the strong  $3072\text{-cm}^{-1}$  absorption. Further confirmation of the assignment is afforded by the fact that all the remaining weak bands have been interpreted as simple combination tones.

(b) *Trans*  $C_2H_2Cl_2$ 

The five strong polarized Raman lines again correspond to the totally symmetric  $A_g$  vibrations. The Raman line of medium intensity at  $758\text{ cm}^{-1}$  is probably the  $B_g$  wagging frequency, though this line should be depolarized.<sup>9</sup> This assignment, however, is supported by the appearance of a corresponding line at  $658\text{ cm}^{-1}$  in the spectrum of *trans*  $C_2D_2Cl_2$  as expected from the Teller-Redlich<sup>10</sup> product rule.

The strong infra-red bands at  $817, 895, 1200$ , and  $3080\text{ cm}^{-1}$  are presumably infra-red active fundamentals. The  $895\text{-cm}^{-1}$  band which shows a clearly resolved rotational fine structure in the vapor spectrum, may be assigned to type  $A_u$  and the  $817$ - and  $1200\text{-cm}^{-1}$  bands which exhibit doublet structures in the vapor spectrum may be assigned to type  $B_u$ . The  $3080\text{ cm}^{-1}$  band is obviously the type  $B_u$   $C-H$  stretching frequency. This leaves two low frequencies

FIG. 1. Infra-red spectra of *cis*-dichloroethylene.

A. Vapor	Cell length 10 cm	Temperature 20°C	
	Pressure (a) 120-mm Hg	(b) 30-mm Hg	(c) 10-mm Hg
B. Liquid	(d) Pure liquid	(e) 10 percent solution in $CS_2$	
	0.1-mm cell	0.1-mm cell	
	(f) 10 percent solution in $CS_2$	(g) 10 percent solution in $CCl_4$	
	0.025-mm cell	0.1-mm cell	
	----- <i>trans</i> impurity		

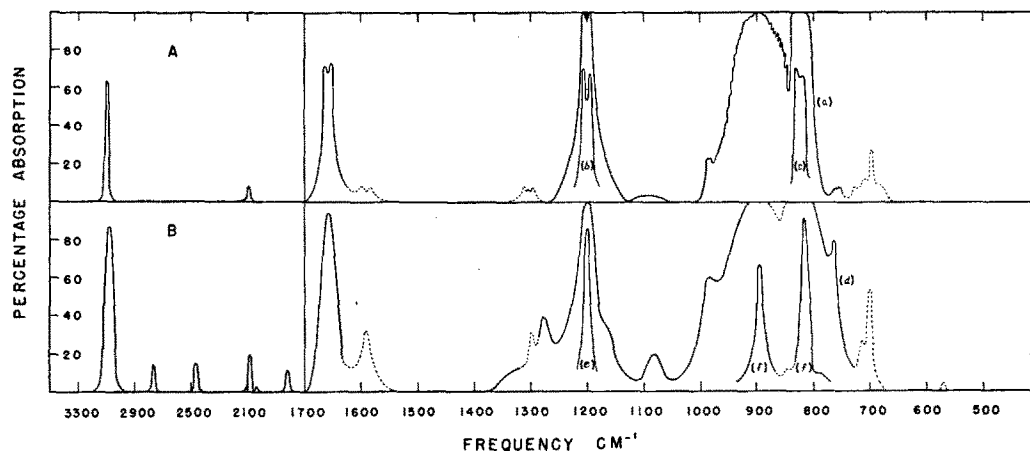
<sup>6</sup> The normal modes of vibration of *cis*- and *trans*-dichloroethylene, their symmetry classes and selection rules, have been given by G. Herzberg (see reference 1) and are not reproduced here. The same system of numbering has been used.

<sup>7</sup> It is interesting to note that the first overtone of this band at  $2367\text{ cm}^{-1}$  has approximately the same intensity as the fundamental. A similar phenomenon is found in the *cis*  $C_2D_2Cl_2$  molecule.

<sup>8</sup> R. M. Badger and L. R. Zumwalt, J. Chem. Phys. **6**, 711 (1938). Using the electron diffraction values for the molecular dimensions (see section on product rule) the parameters  $\rho$  and  $S$ , used by Badger and Zumwalt to express the asymmetry of the molecule, take the values  $\rho = 3.89$ ,  $S = -0.91$ .

<sup>9</sup> The polarization of this line has been discussed by Herzberg (see reference 1).

<sup>10</sup> E. Teller, quoted in J. Chem. Soc. p. 966 (1936). O. Redlich, Zeits. f. Physik chemie **B28**, 371 (1935).

FIG. 2. Infra-red spectra of *trans*-dichloroethylene.

A. Vapor	Cell length 10 cm	Temperature 20°C	
	Pressure (a) 120-mm Hg	(b) 30-mm Hg	(c) 5-mm Hg
B. Liquid	(d) Pure liquid	(e) 10 percent solution in CS <sub>2</sub>	
	0.1-mm cell	(f) 10 percent solution in CS <sub>2</sub>	
	0.025-mm cell	0.1-mm cell	

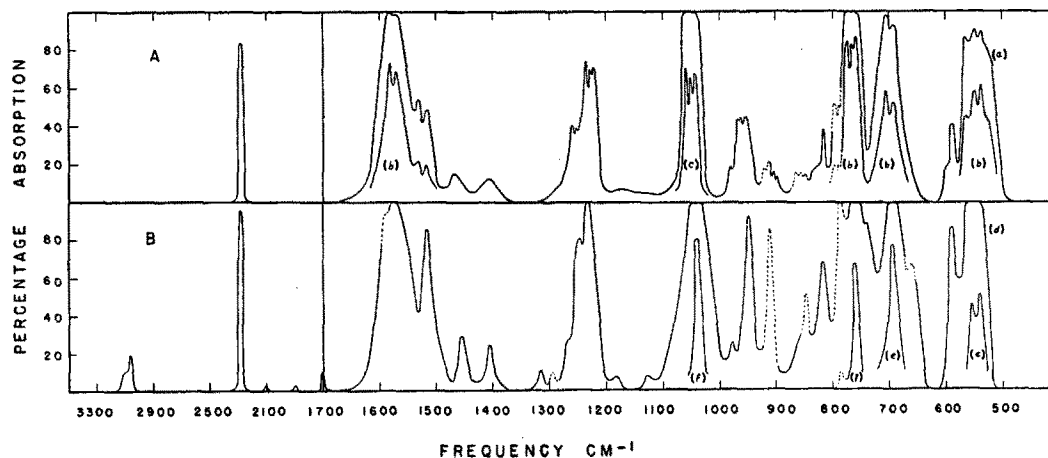
----- *cis* impurity

unassigned, one of type  $A_u$  and the other of type  $B_u$ . All the remaining weak bands may be interpreted as simple combination tones with the exception of the bands at 763, 983, 1080, 1276, and 1816  $\text{cm}^{-1}$ .

The two low frequencies must lie below 400  $\text{cm}^{-1}$  since they were not observed using a KBr prism. Moreover, the results of Wood and Stevenson<sup>5</sup> indicated that the sum of these frequencies should lie between 420 and 450  $\text{cm}^{-1}$ . Pitzer and Freeman<sup>4</sup> assigned a value of 270  $\text{cm}^{-1}$  to the  $A_u$  torsion fre-

quency by assuming that the 620- $\text{cm}^{-1}$  band reported by Wu<sup>2</sup> was a combination tone, 350+270  $\text{cm}^{-1}$ . The 620- $\text{cm}^{-1}$  band, however, must have been caused by an impurity, since we have not observed it using even longer absorption paths than Wu.

The in-plane bending frequency of type  $B_u$  may be computed from the rotational product rule<sup>11</sup> for *cis-trans* isomers, a value of approximately 280  $\text{cm}^{-1}$  being obtained. The torsion frequency of type  $A_u$  however, cannot be obtained by this means, as the rule is not applicable to out-of-plane vibrations.

FIG. 3. Infra-red spectra of *cis*-dichloroethylene- $d_2$ .

A. Vapor	Cell length 10 cm	Temperature 20°C	
	Pressure (a) 120-mm Hg	(b) 30-mm Hg	(c) 10-mm Hg
B. Liquid	(d) Pure liquid	(e) 10 percent solution in CS <sub>2</sub>	
	0.1-mm cell	(f) 10 percent solution in CS <sub>2</sub>	
	0.025-mm cell	0.1-mm cell	

----- *trans-d2* and *cis* impurities

<sup>11</sup> H. J. Bernstein, J. Chem. Phys. 17, 256 (1949).

TABLE I. Infra-red and Raman frequencies for *cis* C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub>.

Infra-red (cm <sup>-1</sup> )		Raman (cm <sup>-1</sup> ) liquid*	Assignment
Vapor	Liquid		
		173 (s) pol. 406 (s) depol.	$\nu_6(A_1)$ $\nu_7(A_2)$
P 564 Q 571 R 578	571 (s)	563 (m) depol.	$\nu_{11}(B_1)$
P 688 Q 697 R ?	697 (vs)		$\nu_{12}(B_2)$
? 727	714 (s)	711 (s) pol.	$\nu_4(A_1)$
808 823	815 (w)	807 (vw)	406+406 = 812 (A <sub>1</sub> )
P 851 Q 857 R 864	848 (vs)		$\nu_{10}(B_1)$
	975 (vw) 1022 (vw) 1105 (w) 1183 (w) 1218 (w) 1254 (w) 1279 (m)	876 (w) depol. 1179 (s) pol.	$\nu_8(A_2)$ 406+571 = 977 (B <sub>2</sub> ) 173+848 = 1021 (B <sub>1</sub> ) 406+697 = 1103 (B <sub>1</sub> ) $\nu_2(A_1)$ 2×697-173 = 1221 (A <sub>1</sub> ) 406+848 = 1254 (B <sub>2</sub> ) 571+714 = 1285 (B <sub>1</sub> )
P 1296 Q 1303 R 1310	1294 (s)		$\nu_9(B_1)$
	1404 (w) 1470 (w) 1571 (w)		697+714 = 1411 (B <sub>2</sub> ) 173+1294 = 1467 (B <sub>1</sub> ) 697+876 = 1573 (B <sub>1</sub> ) 714+848 = 1562 (B <sub>1</sub> )
1583 1598	1590 (s)	1587 (s) pol.	$\nu_2(A_1)$
	1695 (w)	1689 (w) depol. ?	848+848 = 1696 (A <sub>1</sub> ) 406+1294 = 1700 (B <sub>2</sub> ) 876+876 = 1752 (A <sub>1</sub> ) 571+1183 = 1754 (B <sub>1</sub> ) 173+1590 = 1763 (A <sub>1</sub> ) 848+1294 = 2142 (A <sub>1</sub> ) 1183+1183 = 2366 (A <sub>1</sub> ) 1183+1590 = 2773 (A <sub>1</sub> ) 1294+1590 = 2884 (B <sub>1</sub> )
	1753 (vw)		$\nu_3(B_1)$
	2137 (w) 2367 (w) 2753 (w) 2855 (w) 3072 (s)	3077 (vs) pol. 3160 (w) depol. ?	$\nu_1(A_1)$ 1590+1590 = 3180 (A <sub>1</sub> )
2860 3072	3072 (s)		

\* See reference 1.

Since the sum of the two low frequencies must lie between 420 and 450 cm<sup>-1</sup> the torsion frequency may be expected to lie in the region 150±50 cm<sup>-1</sup>.

Better estimates of these frequencies may be obtained from combination tones. The only simple combinations of the low frequencies which are active in the infra-red are combinations of the Raman frequency at 349 cm<sup>-1</sup> with the two low infra-red frequencies. These combination bands should appear in the region 450 to 650 cm<sup>-1</sup>. The absorption of 1 mm of liquid in this region was investigated using a KBr prism and two weak bands were observed at 541 and 614 cm<sup>-1</sup> (see Fig. 5) in addition to the *cis* impurity band at 571 cm<sup>-1</sup>. If these bands are interpreted in the above manner values of 192 cm<sup>-1</sup> and 265 cm<sup>-1</sup> are obtained for the two low infra-red frequencies. Using these values, simple interpretations are obtained for the five remaining weak bands hitherto unexplained.

(c) *Cis* C<sub>2</sub>D<sub>2</sub>Cl<sub>2</sub>

The five strongest lines in the Raman spectrum, *viz.* 171, 689, 850, 1570, 2325 cm<sup>-1</sup>, correspond to the

five totally symmetric lines for the non-deuterated molecule. Similarly, the 368-cm<sup>-1</sup> Raman line corresponds to the A<sub>2</sub> vibration at 406 cm<sup>-1</sup> in the light molecule. The other frequency of this species has not been observed but may be estimated from the Teller-Redlich product rule to be in the region of 710 cm<sup>-1</sup> (see below).

The lone frequency of type B<sub>2</sub> must correspond to one of the strong infra-red bands at 540 and 558 cm<sup>-1</sup>. The contours of these bands overlap in the vapor spectrum and are not readily separated, but the isotopic product rule indicates that the band at 558 cm<sup>-1</sup> is the B<sub>2</sub> frequency. The band at 540 cm<sup>-1</sup> must be assigned then to type B<sub>1</sub> together with the strong infra-red bands at 761, 1040, and 2280 cm<sup>-1</sup>. The 761 and 1040 cm<sup>-1</sup> bands show well defined PQR structures in the vapor spectrum.

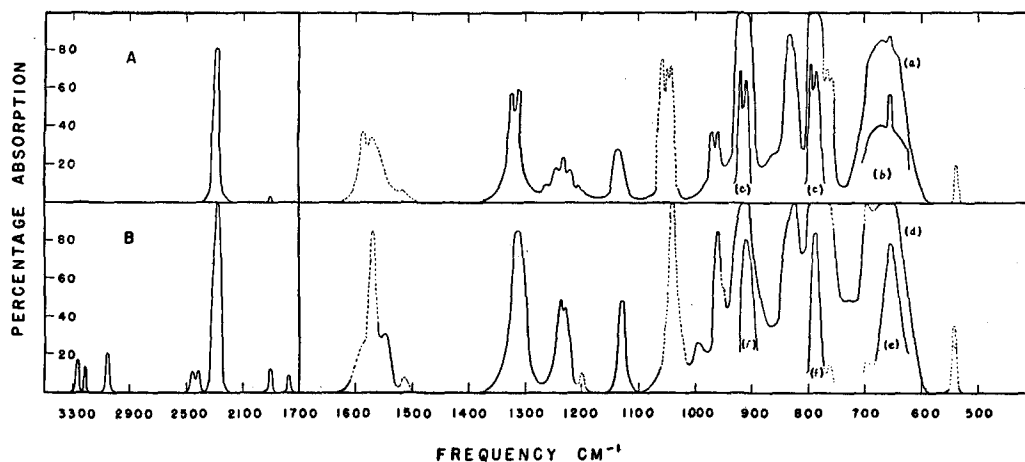
Most of the remaining bands may be interpreted as simple combination tones corresponding to those found for the non-deuterated molecule. A few weak bands, *viz.* 818, 976, 1179, 1314, however, cannot be explained as simple combination tones and probably represent combinations perturbed by Fermi resonance. This is to be expected since many combinations of the same symmetry have approximately the same frequency in the region 800–1400 cm<sup>-1</sup>. Tentative assignments have been given in these cases.

TABLE II. Infra-red and Raman frequencies for *trans* C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub>.

Infra-red (cm <sup>-1</sup> )		Raman (cm <sup>-1</sup> ) liquid*	Assignment
Vapor	Liquid		
		349 (s) pol.	$\nu_8(A_2)$ 349+192 = 541 (A <sub>u</sub> ) 349+265 = 614 (B <sub>u</sub> )
754 763	763 (w)	758 (m) pol.	$\nu_8(B_2)$ 844+265-349 = 760 (B <sub>u</sub> )
823 832	817 (vs)		$\nu_{11}(B_u)$
		844 (s) pol.	$\nu_4(A_g)$
898	895 (vs) 985 (w) 1080 (w) 1166 (w)		$\nu_6(A_u)$ * { 2×349+265 = 963 (B <sub>u</sub> ) 265+895-192 = 968 (B <sub>u</sub> ) 1270-192 = 1078 (A <sub>u</sub> ) 349+817 = 1166 (B <sub>u</sub> )
1195 1205	1200 (s)		$\nu_{10}(B_u)$
		1270 (s) pol.	$\nu_2(A_g)$
	1276 (w)		192+265+817 = 1274 (A <sub>u</sub> )
		1576 (s) pol. 1626 (vw) pol.	$\nu_2(A_g)$ 817+817 = 1634 (A <sub>g</sub> )
1654 1663	1658 (m)		758+895 = 1653 (B <sub>u</sub> ) 817+844 = 1661 (B <sub>u</sub> ) 844+844 = 1688 (A <sub>g</sub> ) 2×817+192 = 1826 (A <sub>u</sub> ) 844+1200 = 2044 (B <sub>u</sub> )
	1816 (w) 2047 (w)	1692 (w) pol.	817+1270 = 2087 (B <sub>u</sub> ) 1200+1270 = 2470 (B <sub>u</sub> ) 895+1576 = 2471 (A <sub>u</sub> ) 1200+1576 = 2776 (B <sub>u</sub> )
2092	2082 (w) 2460 (w) 2760 (w)	3071 (s) pol.	$\nu_1(A_g)$
3090	3080 (s)		$\nu_9(B_u)$
		3142 (w) depol. ?	1576+1576 = 3152 (A <sub>g</sub> )

\* See reference 1.

\* Fermi resonance gives 985 cm<sup>-1</sup> and 945 cm<sup>-1</sup> the latter being obscured by the intense 895 cm<sup>-1</sup> band.

FIG. 4. Infra-red spectra of *trans*-dichloroethylene- $d_2$ .

A. Vapor Cell length 10 cm Temperature 20°C  
 Pressure (a) 120-mm Hg (b) 30-mm Hg (c) 10-mm Hg  
 B. Liquid (d) Pure liquid (e) 10 percent solution in  $CS_2$   
 0.1-mm cell (f) 10 percent solution in  $CS_2$   
 0.025-mm cell  
 ----- *cis*- $d_2$  impurity

(d) *Trans*  $C_2D_2Cl_2$ 

By comparison with the assignment for *trans*  $C_2H_2Cl_2$ , the Raman frequencies at 346, 765, 992, 1570, and 2325  $cm^{-1}$  may be assigned unambiguously to type  $A_g$  and the 657- $cm^{-1}$  Raman line to type  $B_g$ . The infra-red band at 658  $cm^{-1}$  which exhibits a  $P, Q, R$  structure in the vapor spectrum corresponds to the 895  $cm^{-1}$  band for the light molecule though the resolving power proved insufficient to resolve the rotational fine structure for the heavy molecule. The strong infra-red bands at 784, 912, and 2285  $cm^{-1}$  may be assigned to type  $B_u$ , the doublet structures for the 784- and 912- $cm^{-1}$  bands being resolved in the vapor spectrum. If it is assumed that the two low frequencies in the light molecule, *viz.* 192 and 265  $cm^{-1}$  are not appreciably

affected by substitution of deuterium, a complete assignment is obtained and all the remaining weak bands may be satisfactorily interpreted as simple combination tones.

PRODUCT RULE RATIOS FOR *CIS* AND *TRANS*  $C_2H_2Cl_2$  AND  $C_2D_2Cl_2$ 

Using the dimensions for the *cis* and *trans* molecules obtained by the electron diffraction method<sup>12</sup> *viz.*

	<i>cis</i>	<i>trans</i>
$r_{C-C}$	1.38Å (assumed)	1.38Å (assumed)
$r_{C-Cl}$	$1.67 \pm 0.03Å$	$1.69 \pm 0.02Å$
$Cl-C=C$	$123.5 \pm 1^\circ$	$122.5 \pm 1^\circ$

and further assuming  $r_{C-H} = 1.07Å$ ,  $H-C=C = 120^\circ$  as in ethylene<sup>13</sup> the principal moments of inertia for the light and heavy molecules are evaluated as follows:

	<i>cis</i> $C_2H_2Cl_2$	<i>cis</i> $C_2D_2Cl_2$	<i>trans</i> $C_2H_2Cl_2$	<i>trans</i> $C_2D_2Cl_2$
$I_A$	70.3	82.4	15.8	23.2
$I_B$	329.8	334.8	548.2	548.6
$I_C$	400.1	417.2	564.0	571.8
units of $10^{-40} g cm^2$				

The theoretical product rule ratios for the various symmetry types of the *cis* and *trans* molecules may now be calculated and compared with the observed ratios (Tables V and VI). The agreement, in general, is good though the observed ratio for the *cis*  $B_1$

<sup>12</sup> L. O. Brockway, J. Y. Beach, and L. Pauling, J. Am. Chem. Soc. 57, 2693 (1935).

<sup>13</sup> W. S. Gallaway and E. F. Barker, J. Chem. Phys. 10, 88 (1942).

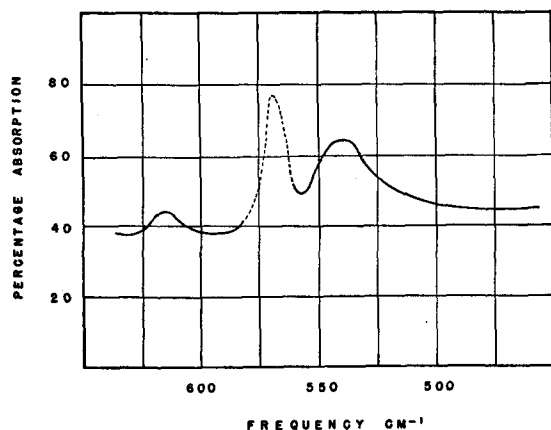


FIG. 5. Infra-red absorption of *trans* dichloroethylene from 450 to 650  $cm^{-1}$  using 1-mm path length of liquid. Dotted band at 571  $cm^{-1}$  caused by *cis*-dichloroethylene impurity.

vibration is slightly high. Liquid frequencies have been used throughout.

### ANALYSIS OF THE FINE STRUCTURE OF THE 898-CM<sup>-1</sup> BAND OF TRANS C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub>

The frequencies of the individual rotation lines were obtained from the ammonia data given by Oetjen, Kao and Randall<sup>14</sup> (see Fig. 6) and are listed in Table VII. The frequencies have been corrected to vacuum and are probably accurate to  $\pm 0.1$  cm<sup>-1</sup>. The *trans* C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub> molecule approximates very closely to a symmetric top whose figure axis passes approximately through the two chlorine atoms. Since the 898-cm<sup>-1</sup> band corresponds to an out-of-plane wagging motion of the CHCl groups, the change in electric moment is perpendicular to this

TABLE III. Infra-red and Raman frequencies for *cis* C<sub>2</sub>D<sub>2</sub>Cl<sub>2</sub>.

Infra-red (cm <sup>-1</sup> )		Raman (cm <sup>-1</sup> )		Assignment
Vapor	Liquid	liquid*		
		171 (s)	$\nu_8(A_1)$	
		368 (m)	$\nu_7(A_2)$	
		515 (w)	$689 - 171 = 518 (A_1)$	
	540 (s)		$\nu_{11}(B_1)$	
	558 (s)		$\nu_{12}(B_2)$	
	590 (m)		$761 - 171 = 590 (B_1)$	
693 } 706 }	695 (s)	689 (s)	$\nu_8(A_1)$	
	741 (w)		$368 + 368 = 736 (A_1)^*$	
P 759 } Q 766 } R 773 }	761 (vs)		$\nu_{10}(B_1)$	
			$\dagger$	
818	818 (w)	850 (m)	$\nu_8(A_1)$	
P 899 } Q 905 } R 912 }	912 (m) #		$368 + 558 = 926 (B_1)^*$	
			Fermi Resonance	
P 955 } Q 962 } R 967 }	949 (m)		$171 + 761 = 932 (B_1)^*$	
	976 (w)		$\dagger$	
P 1045 } Q 1051 } R 1058 }	1040 (vs)		$\nu_9(B_1)$	
	1129 (w)		$368 + 761 = 1129 (B_2)^*$	
	1179 (w)		$\dagger\dagger$	
P 1222 } Q 1228 } R 1235 }	1228 (m)		$540 + 695 = 1235 (B_1)^*$	
	1247 (w)		$558 + 695 = 1253 (B_2)^*$	
	1263 (w)		$558 + 710 = 1268 (B_1)^*$	
	1314 (w)		$\dagger\dagger$	
	1405 (w)		$368 + 1040 = 1408 (B_2)^*$	
	1455 (w)		$695 + 761 = 1456 (B_1)^*$	
1518 } 1532 }	1515 (m)	1507 (w)	$761 + 761 = 1522 (A_1)^*$	
1569 } 1582 }	1572 (s)	1570 (s)	$\nu_2(A_1)$	
	1702 (w)		$850 + 850 = 1700 (A_1)^*$	
	1885 (vw)		$850 + 1040 = 1890 (B_1)$	
	2107 (vw)		$540 + 1572 = 2112 (B_1)$	
2280	2280 (s)		$\nu_8(B_1)$	
		2325 (s)	$\nu_1(A_1)$	
	3060 (w)		$2325 + 761 = 3086 (B_1)$	
	3110 (w)		$2280 + 850 = 3130 (B_1)$	
		3150 ? (vw)	$1572 + 1572 = 3144 (A_1)^*$	

\* See reference 3.

\* The corresponding combination tone was observed for *cis* C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub>.

# *Trans* C<sub>2</sub>D<sub>2</sub>Cl<sub>2</sub> impurity contributes partly to the intensity of this band.

† Fermi resonance between  $2 \times 171 + 558 = 900 (B_2)$ ;  $368 + 540 = 908 (B_2)$ .

†† Fermi resonance between two of the following:  $171 + 368 + 710 = 1249 (A_1)$ ;  $2 \times 710 - 171 = 1249 (A_1)$ ;  $2 \times 540 + 171 = 1251 (A_1)$ .

<sup>14</sup> R. A. Oetjen, C. L. Kao, and H. M. Randall, Rev. Sci. Inst. 13, 515 (1942).

TABLE IV. Infra-red and Raman frequencies for *trans* C<sub>2</sub>D<sub>2</sub>Cl<sub>2</sub>.

Infra-red (cm <sup>-1</sup> )		Raman (cm <sup>-1</sup> )	Assignment
Vapor	Liquid	liquid*	
		346 (s)	$\nu_8(A_g)$
		657 (m)	$\nu_8(B_g)$
P ? } Q 660 } R ? }	658 (s)		$\nu_8(A_u)$
	729 (vw)		$265 + 658 - 192 = 731 (B_u)^*$
		765 (m)	$\nu_4(A_g)$
786 } 796 }	784 (vs)		$\nu_{11}(B_u)$
	826 (m) # 836 (w)		† $657 + 192 = 849 (B_u)$
911 } 921 }	912 (vs)		$\nu_{10}(B_u)$
960 } 969 }	960 (m)		$265 + 2 \times 346 = 957 (B_u)^*$
		992 (s)	$\nu_3(A_g)$
	995 (w)		992 cm <sup>-1</sup> Raman line
	1130 (w)		$346 + 784 = 1130 (B_u)^*$
	1230 (w)		$765 + 657 - 192 = 1230 (B_u)$
	1238 (w)		$192 + 265 + 784 = 1241 (A_u)$
1314 } 1324 }	1314 (m)		$657 + 658 = 1315 (B_u)^*$
		1312 (w) ? 1545 (m)	$658 + 658 = 1316 (A_g)$
	1549 (w)		$765 + 765 = 1530 (A_g)^*$
		1570 (s)	$765 + 784 = 1549 (B_u)^*$
	1780 (w)		$\nu_2(A_g)$
	1900 (w)		$784 + 992 = 1776 (B_u)^*$
2290	2285 (s)		$912 + 992 = 1904 (B_u)^*$
		2325 (s)	$\nu_9(B_u)$
	2417 (w)		$\nu_1(A_g)$
	2462 (w)		$2 \times 765 + 912 = 2442 (B_u)$
	3054 (w)		$912 + 1570 = 2482 (B_u)$
		3154 (w)	$765 + 2285 = 3050 (B_u)$
	3226 (w)		$1570 + 1570 = 3140 (A_g)^*$
	3265 (w)		$912 + 2325 = 3237 (B_u)$
			$992 + 2285 = 3277 (B_u)$

\* See reference 3.

\* The corresponding combination tone was observed for *trans* C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub>.

# Part of the intensity of this band is contributed by *trans* C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub> impurity.

† Probably Fermi resonance between;  $346 + 657 - 192 = 811 (B_u)$ ;  $3 \times 265 = 795 (B_u)$ ; to give 826 cm<sup>-1</sup> and a component which is masked by the intense 784 cm<sup>-1</sup> band.

axis, hence we should expect a perpendicular type band. Moreover since  $I_A \ll I_B$  or  $I_C$  the band will consist mainly of a series of *Q* branches of the various sub-bands present, the *PR* separations in the individual sub-bands being considerably smaller than the separations between the various *Q* branches.

The fact that a strong "*Q* branch" is found in the region of the band origin gives the band the appearance of a  $\parallel$  band. This is altogether out of the question however due to the magnitude of the spacing. The appearance of this *Q* branch is probably due to the slight deviation of the molecule from a perfect symmetric top, c.f., the type *C* band at 949 cm<sup>-1</sup> for ethylene.<sup>15</sup> The spacing of the *Q* branches is given by the equation<sup>16</sup>

$$\nu = \nu_0 + (A' - B') \pm 2(A' - B')K + [(A' - B') - (A'' - B'')]K^2, \quad (1)$$

where  $A'$ ,  $B'$ , are the rotational constants for the upper state,  $A''$ ,  $B''$ , are the rotational constants for the lower state, the positive sign applies to the  $RQ$

<sup>15</sup> See reference 1, p. 483.

<sup>16</sup> See reference 1, p. 424.

TABLE V. Frequency assignments and product rule ratios for *cis* C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub> and *cis* C<sub>2</sub>D<sub>2</sub>Cl<sub>2</sub>.

Type	<i>Cis</i> C <sub>2</sub> H <sub>2</sub> Cl <sub>2</sub>	<i>Cis</i> C <sub>2</sub> D <sub>2</sub> Cl <sub>2</sub>	Theor. ratio	Obs. ratio
<i>A</i> <sub>1</sub>	$\nu_1$ 3077 R	2325 R	1.98	1.94
	$\nu_2$ 1587 R	1570 R		
	$\nu_3$ 1179 R	850 R		
	$\nu_4$ 711 R	689 R		
	$\nu_5$ 173 R	171 R		
<i>A</i> <sub>2</sub>	$\nu_6$ 876 R	710*	1.40	(1.36)
	$\nu_7$ 406 R	368 R		
<i>B</i> <sub>1</sub>	$\nu_8$ 3072 I.R.	2280 I.R.	1.94	1.97
	$\nu_9$ 1294 I.R.	1040 I.R.		
	$\nu_{10}$ 848 I.R.	761 I.R.		
	$\nu_{11}$ 571 I.R.	540 I.R.		
<i>B</i> <sub>2</sub>	$\nu_{12}$ 697 I.R.	558 I.R.	1.29	1.25

\* Calculated from the theoretical ratio after correction for anharmonicity.

branch ( $K=0, 1, 2, \dots$ ) the negative sign applies to the  $^PQ$  branch ( $K=1, 2, \dots$ ).

The band center may be determined from the relation

$$^RQ_K + ^PQ_{K+1} = 2\nu_0 + [(A' - B') - (A'' - B'')][K^2 + (K+1)^2]$$

by plotting  $^RQ_K + ^PQ_{K+1}$  against  $K^2 + (K+1)^2$  and extrapolating to zero (see Fig. 7). From this graph it is found that  $\nu_0 = 898.1 \text{ cm}^{-1}$  and  $(A' - B') - (A'' - B'') = -0.0084 \text{ cm}^{-1}$ .

Individual values for  $(A' - B')$  and  $(A'' - B'')$  may be obtained from the combination relations

$$\begin{aligned} ^RQ_K - ^PQ_K &= 4(A' - B')K, \\ ^RQ_{K-1} - ^PQ_{K+1} &= 4(A'' - B'')K. \end{aligned}$$

TABLE VI. Frequency assignments and product rule ratios for *trans* C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub> and *trans* C<sub>2</sub>D<sub>2</sub>Cl<sub>2</sub>.

Type	<i>Trans</i> C <sub>2</sub> H <sub>2</sub> Cl <sub>2</sub>	<i>Trans</i> C <sub>2</sub> D <sub>2</sub> Cl <sub>2</sub>	Theor. ratio	Obs. ratio
<i>A</i> <sub>g</sub>	$\nu_1$ 3071 R	2325 R	1.98	1.89
	$\nu_2$ 1576 R	1570 R		
	$\nu_3$ 1270 R	992 R		
	$\nu_4$ 844 R	765 R		
	$\nu_5$ 349 R	346 R		
<i>A</i> <sub>u</sub>	$\nu_6$ 895 I.R.	658 I.R.	1.40	1.36
	$\nu_7$ 192 I.R.*	192 I.R.*		
<i>B</i> <sub>g</sub>	$\nu_8$ 758 R	657 R	1.17	1.15
<i>B</i> <sub>u</sub>	$\nu_9$ 3080 I.R.	2285 I.R.	1.96	1.85
	$\nu_{10}$ 1200 I.R.	912 I.R.		
	$\nu_{11}$ 817 I.R.	784 I.R.		
	$\nu_{12}$ 265 I.R.*	265 I.R.*		

\* From combination tones.

Since the band positions were determined with an accuracy of  $\pm 0.1 \text{ cm}^{-1}$ , the probable errors in the values of  $(A' - B')$  and  $(A'' - B'')$  will depend on the  $K$  value but will be of the order of  $\pm (0.2/4 \times 15) \sim 0.004 \text{ cm}^{-1}$ . Hence from Table VII we see that

$$\begin{aligned} A' - B' &= 1.674 \pm 0.004 \text{ cm}^{-1}, \\ A'' - B'' &= 1.682 \pm 0.004 \text{ cm}^{-1}, \end{aligned}$$

these values being consistent with the difference found above. Substituting into Eq. (1) the frequencies of the rotation lines were calculated and found to agree with the observed values in most cases to  $\pm 0.1 \text{ cm}^{-1}$ , though occasional differences of  $\pm 0.2 \text{ cm}^{-1}$  were found.

If we take  $I_{B'} = I_{B''} = \text{mean of } I_B \text{ and } I_C$  (obtained from electron diffraction data),

$$B' = B'' = 0.050 \pm 0.001 \text{ cm}^{-1}.$$

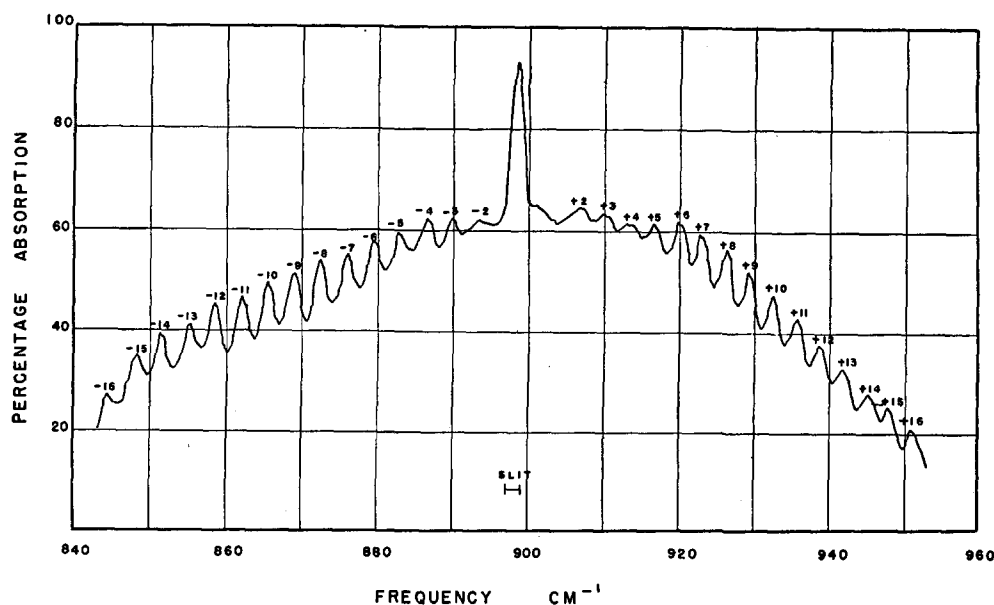


FIG. 6. Rotational fine structure of the  $898.1\text{-cm}^{-1}$  band of *trans*-dichloroethylene. Absorption path 10-cm, pressure 30-mm Hg, Temperature  $20^\circ\text{C}$ .



Hence,

$$A' = 1.724 \pm 0.005 \text{ cm}^{-1}; \quad A'' = 1.733 \pm 0.005 \text{ cm}^{-1},$$

and

$$I_{A'} = 16.23 \pm 0.05 \times 10^{-40} \text{ g cm}^2;$$

$$I_{A''} = 16.15 \pm 0.05 \times 10^{-40} \text{ g cm}^2.$$

These values are considerably more accurate than the value obtained from the electron diffraction results *viz.*  $15.8 \times 10^{-40} \text{ g cm}^2$ .

A slight refinement of the molecular dimensions for the *trans* molecule may be made in the following manner. For small changes in the C—Cl and C=C

bond lengths and the Cl—C=C bond angle, the contribution from the hydrogen atoms to the least moment of inertia is approximately constant. Hence by taking the C—H bond length to be 1.07 Å and the

H—C=C bond angle to be 120°, the hydrogen contribution may be evaluated and subtracted from the moment of inertia. The remaining contributions from the carbon and chlorine atoms depend on three variables, the C—Cl and C=C bond lengths and the

Cl—C=C bond angle. By using the electron diffraction result for the distance between the two chlorine atoms *viz.*  $4.27 \pm 0.02 \text{ Å}$ , the number of independent variables is reduced to two. Moments of inertia are then calculated for reasonable ranges of values of the C=C and C—Cl bond lengths and the values compared with the spectroscopic figure. The ranges used were  $r_{\text{C=C}} = 1.34$  to  $1.40 \text{ Å}$  and  $r_{\text{C-Cl}} = 1.65$  to  $1.75 \text{ Å}$  and it was found that the moments of inertia were consistent with the spectroscopic data only when

$$r_{\text{C=C}} + r_{\text{C-Cl}} = 3.08 \pm 0.01 \text{ Å}.$$

TABLE VII. Rotational fine structure for  $898 \text{ cm}^{-1}$  *trans* C<sub>2</sub>H<sub>2</sub>Cl<sub>2</sub> band.

K	$\nu$ obs.	$\nu$ calc.	Diff.	K	$RQ_K - P Q_K$	$4K(A' - B')$	$(A' - B')$
-16	844.3	844.1	0.2	2	906.7-893.1	13.6	(1.700)
-15	847.8	847.7	0.1	3	909.7-889.6	20.1	1.675
-14	851.4	851.3	0.1	4	913.0-886.3	26.7	1.669
-13	854.9	854.8	0.1	5	916.4-882.8	33.6	1.680
-12	858.3	858.4	-0.1	6	919.7-879.3	40.4	1.683
-11	861.9	861.9	0.0	7	922.9-875.9	47.0	1.679
-10	865.4	865.5	-0.1	8	926.0-872.4	53.6	1.675
-9	868.9	869.0	-0.1	9	929.2-868.9	60.3	1.675
-8	872.4	872.5	-0.1	10	932.4-865.4	67.0	1.675
-7	875.9	875.9	0.0	11	935.6-861.9	73.7	1.675
-6	879.3	879.4	-0.1	12	938.7-858.3	80.4	1.675
-5	882.8	882.8	0.0	13	941.7-854.9	86.8	1.669
-4	886.3	886.3	0.0	14	945.0-851.4	93.6	1.671
-3	889.6	889.7	-0.1	15	947.9-847.8	100.1	1.668
-2	893.1	893.1	0.0	16	951.0-844.3	106.7	1.667
-1		896.4					
0		899.8					
1		903.1					Mean = 1.674
2	906.7	906.5	0.2	K	$RQ_{K-1} - P Q_{K+1}$	$4K(A'' - B'')$	$(A'' - B'')$
3	909.7	909.8	-0.1	3	906.7-886.3	20.4	(1.700)
4	913.0	913.0	0.0	4	909.7-882.8	26.9	1.681
5	916.4	916.3	0.1	5	913.0-879.4	33.6	1.680
6	919.7	919.6	0.1	6	916.4-875.9	40.5	1.687
7	922.9	922.8	0.1	7	919.7-872.5	47.2	1.686
8	926.0	926.0	0.0	8	922.9-869.0	53.9	1.684
9	929.2	929.2	0.0	9	926.0-865.5	60.5	1.681
10	932.4	932.4	0.0	10	929.2-861.9	67.3	1.683
11	935.6	935.6	0.0	11	932.4-858.4	74.0	1.682
12	938.7	938.8	-0.1	12	935.6-854.8	80.8	1.683
13	941.7	941.9	-0.2	13	938.7-851.3	87.4	1.681
14	945.0	945.0	0.0	14	941.7-847.7	94.0	1.679
15	947.9	948.1	-0.2	15	945.0-844.1	100.9	1.682
16	951.0	951.2	-0.2				Mean = 1.682

#### POTENTIAL FUNCTION FOR OUT-OF-PLANE FREQUENCIES

##### Type A<sub>2</sub> (Cis)

A general quadratic potential function was used (Fig. 8) *viz.*

$$V = k_{\theta}\theta^2 + k_{\phi}\phi^2 + k'\theta\phi,$$

where  $\theta$  = angle of wag between the CHCl group and the extended C=C axis,  $\phi$  = angle of twist about the extended C=C referred to the equilibrium plane,  $k_{\theta}$ ,  $k_{\phi}$ ,  $k'$  are the appropriate force constants.

The equations for the frequencies were derived by

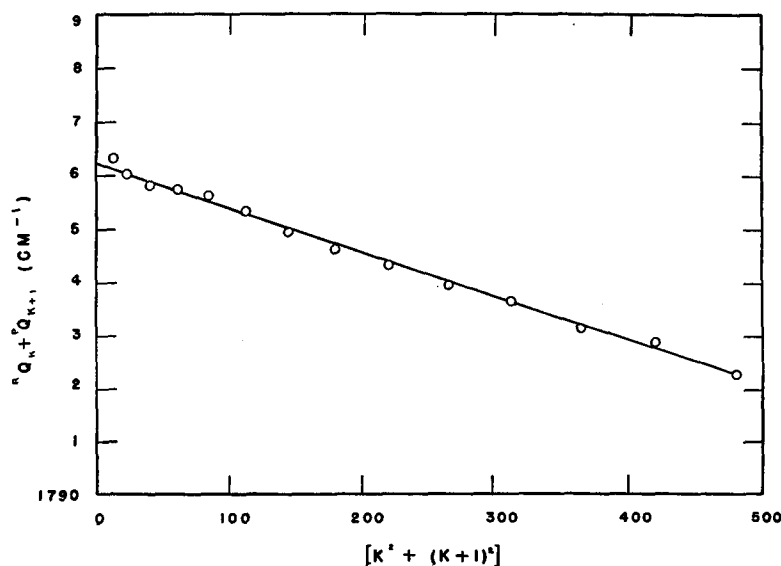


FIG. 7. Determination of the band center for the  $898\text{-cm}^{-1}$  band of *trans* dichloroethylene.

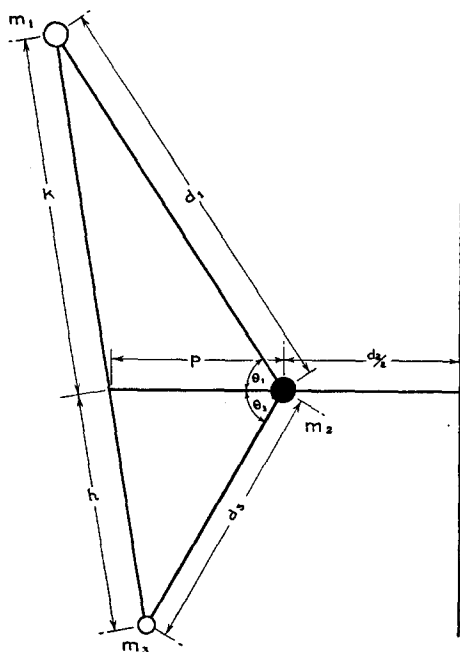


FIG. 8. ● = C, ○ = Cl, ○ = H. Structural parameters for one half of the dichloroethylene molecule.

the method of Eliashevich<sup>17</sup> and may be expressed as follows:

$$\lambda_1^2 + \lambda_2^2 = 2(ak_\theta + ck_\phi - bk'),$$

$$\lambda_1^2 \lambda_2^2 = (4k_\theta k_\phi - k'^2)(ac - b^2),$$

where

$$\lambda = 2\pi c\nu,$$

$$a = \left( \frac{h'^2}{m_1} + \frac{f^2}{m_2} + \frac{k'^2}{m_3} \right),$$

$$b = \left( -\frac{h'}{m_1 d_1 \sin \theta_1} + \frac{fg}{m_2} + \frac{k'}{m_3 d_3 \sin \theta_3} \right),$$

$$c = \left( \frac{1}{m_1 d_1^2 \sin^2 \theta_1} + \frac{g^2}{m_2} + \frac{1}{m_3 d_3^2 \sin^2 \theta_3} \right);$$

$$h' = \frac{h}{p(h+k)}, \quad k' = \frac{k}{p(h+k)}, \quad f = \frac{1}{p} + \frac{2}{d_2},$$

$$g = -\frac{1}{d_1 \sin \theta_1} + \frac{1}{d_3 \sin \theta_3}.$$

The equations for the other out-of-plane vibrations for the *cis* and *trans* molecules may be obtained from the above as follows:

$$\begin{aligned} \text{type } B_2 \text{ (cis): } & \text{let } d_2 = \infty, \quad k_\phi = 0, \quad k' = 0, \\ \text{type } A_u \text{ (trans): } & \text{let } d_2 = \infty, \\ \text{type } B_g \text{ (trans): } & \text{let } k_\phi = 0, \quad k' = 0. \end{aligned}$$

The condition that  $d_2 = \infty$  eliminates the contribution to the potential energy from the relative motion of the carbon atoms.

Force constants may be evaluated as follows. Using  $\nu_8(B_g) = 758 \text{ cm}^{-1}$  for the *trans* molecule,

$$k_\theta = 1.12 \times 10^{-12} \text{ dyne cm/radian.}$$

Assuming that this force constant may be taken over for type  $A_2$  (*cis*) as has been done by Pitzer and Freeman<sup>4</sup> and using  $\nu_6 = 876 \text{ cm}^{-1}$ ,  $\nu_7 = 406 \text{ cm}^{-1}$  in the above equations,

$$\begin{aligned} k_\phi &= 1.46 \times 10^{-12} \text{ dyne cm/radian,} \\ k' &= 0.62 \times 10^{-12} \text{ dyne cm/radian.} \end{aligned}$$

Using these force constants to calculate the two frequencies for *cis*  $\text{C}_2\text{D}_2\text{Cl}_2$  we obtain

$$\nu_6 = 678 \text{ cm}^{-1}, \quad \nu_7 = 374 \text{ cm}^{-1}.$$

The low frequency agrees well with the Raman frequency observed at  $368 \text{ cm}^{-1}$  while the other frequency is in tolerable agreement with the value deduced from the product rule *viz.*  $710 \text{ cm}^{-1}$ .

In a similar manner the wagging force constant for type  $B_2$  (*cis*) may be used in type  $A_u$  (*trans*). Using  $\nu_{12}(B_2) = 697 \text{ cm}^{-1}$  for the *cis* molecule, we obtain

$$k_\theta = 1.51 \times 10^{-12} \text{ dyne cm/radian.}$$

Only one of the  $A_u$  frequencies for the *trans* molecule, however, has been observed, *viz.* at  $895 \text{ cm}^{-1}$  for  $\text{C}_2\text{H}_2\text{Cl}_2$  and  $658 \text{ cm}^{-1}$  for  $\text{C}_2\text{D}_2\text{Cl}_2$ . Since the unobserved frequency is expected to be low, it will change very little from  $\text{C}_2\text{H}_2\text{Cl}_2$  to  $\text{C}_2\text{D}_2\text{Cl}_2$  and may be assumed to be unchanged to a very good approximation. The three unknowns for type  $A_u$ , *viz.*  $k_\phi$ ,  $k'$ , and  $\nu_7$  may then be evaluated from the three equations, for the sums of the frequencies in  $\text{C}_2\text{H}_2\text{Cl}_2$  and  $\text{C}_2\text{D}_2\text{Cl}_2$  and the product of the frequencies in either  $\text{C}_2\text{H}_2\text{Cl}_2$  or  $\text{C}_2\text{D}_2\text{Cl}_2$ . Solution of these equations gives

$$\begin{aligned} k_\phi &= 1.00 \times 10^{-12} \text{ dyne cm/radian,} \\ k' &= 0.20 \times 10^{-12} \text{ dyne cm/radian,} \\ \nu_7 &= 196 \text{ cm}^{-1}. \end{aligned}$$

This value for the low frequency is in good agreement with the value  $192 \text{ cm}^{-1}$  deduced from combination tones.

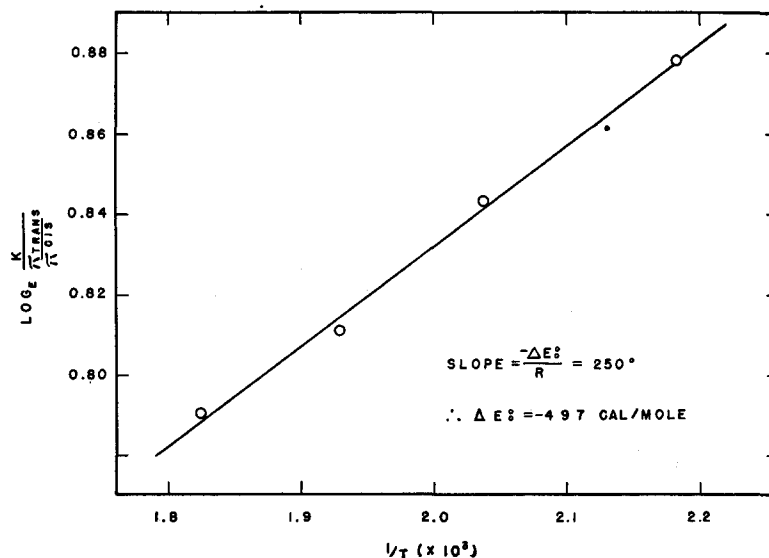
#### THE ENERGY DIFFERENCE BETWEEN *CIS* AND *TRANS* DICHLOROETHYLENE

For an equilibrium of the type *cis*  $\rightleftharpoons$  *trans* the equilibrium constant<sup>18</sup>

$$\begin{aligned} K &= \frac{\text{Number of } cis \text{ molecules}}{\text{Number of } trans \text{ molecules}} \\ &= \frac{f_{cis}}{f_{trans}} \cdot \exp(-\Delta E_0^0/RT), \quad (2) \end{aligned}$$

<sup>17</sup> M. A. Eliashevich, Comptes. Rendus. U.R.S.S. 28, 604 (1940).

<sup>18</sup> R. H. Fowler and E. A. Guggenheim, *Statistical Thermodynamics* (The Macmillan Company, New York, 1939), p. 157.

FIG. 9. Plot of  $\log_e K/(\Pi_{trans}/\Pi_{cis})$  against  $1/T$ .

where  $f_{cis}$ ,  $f_{trans}$  are the partition functions of *cis*- and *trans*-dichloroethylene, respectively,  $R$  is the gas constant, and  $\Delta E_0^0$  is the energy difference between the two forms at 0°K.

Assuming the rigid rotator and harmonic oscillator approximations<sup>19</sup>

$$\frac{f_{cis}}{f_{trans}} = \left[ \frac{(I_A I_B I_C)_{cis}}{(I_A I_B I_C)_{trans}} \right]^{\frac{1}{2}} \frac{\prod_{i=1}^{12} (1 - e^{-h\nu_i/kT})_{trans}}{\prod_{i=1}^{12} (1 - e^{-h\nu_i/kT})_{cis}} \quad (3)$$

Substituting in Eq. (2) and taking logarithms, we get

$$\log_e \frac{K}{\Pi_{trans}/\Pi_{cis}} = \frac{1}{2} \log_e \left[ \frac{(I_A I_B I_C)_{cis}}{(I_A I_B I_C)_{trans}} \right] - \Delta E_0^0 / RT \quad (4)$$

Hence by plotting  $\log_e K/(\Pi_{trans}/\Pi_{cis})$  against  $1/T$  a straight line is obtained whose slope is  $-\Delta E_0^0/R$  and intercept  $\frac{1}{2} \log_e (I_A I_B I_C)_{cis}/(I_A I_B I_C)_{trans}$ . In this way the value of  $\Delta E_0^0$  is not dependent on experimental errors in the determination of the principal moments of inertia of the two molecules.

<sup>19</sup> See reference 1, p. 503, *et. seq.*

Wood and Stevenson<sup>5</sup> have obtained equilibrium data for the isomerization of *cis* and *trans* dichloroethylene at 185°C, 215°C, 245°C, and 275°C. Using their values of  $K$  and our values of  $\Pi_{trans}/\Pi_{cis}$  and plotting  $\log_e K/(\Pi_{trans}/\Pi_{cis})$  against  $1/T$  (see Fig. 9) the best straight line obtained by the method of least squares gives

$$\Delta E_0^0 = -497 \text{ cal./mole.}$$

The principal source of inaccuracy in this determination lies in the assignment to the *trans* molecule of the two low frequencies obtained from combination tones. If we assume that these frequencies *viz.* 192 and 265 cm<sup>-1</sup>, are accurate to  $\pm 5$  cm<sup>-1</sup>, the error introduced into  $\Delta E_0^0$  is  $\pm 30$  cal./mole. Hence for dichloroethylene the *cis* isomer is more stable than the *trans* isomer by  $500 \pm 30$  cal./mole.

The moment of inertia ratio  $[(I_A I_B I_C)_{cis}/(I_A I_B I_C)_{trans}]^{\frac{1}{2}}$  may also be evaluated from Eq. (4) a range of values of  $1.39 \pm 0.04$  being obtained. This is consistent with the value calculated from electron diffraction results *viz.* 1.38.

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