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PAPER

Enhanced discrimination of DNA molecules in nanofluidic channels through multiple measurements†

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Nanofluidic sensing elements have been the focus of recent experiments for numerous applications ranging from nucleic acid fragment sizing to single-molecule DNA sequencing. These applications critically rely on high measurement fidelity, and methods to increase resolution are required. Herein, we describe fabrication and testing of a nanochannel device that enhances measurement resolution by performing multiple measurements (>100) on single DNA molecules. The enhanced measurement resolution enabled length discrimination between a mixture of λ -DNA (48.5 kbp) and T7 DNA (39.9 kbp) molecules, which were detected as transient current changes during translocation of the molecules through the nanochannel. As long DNA molecules are difficult to resolve quickly and with high fidelity with conventional electrophoresis, this approach may yield potentially portable, direct electrical sizing of DNA fragments with high sensitivity and resolution.

Introduction

The study of single biomolecules such as nucleic acids and proteins has provided the ability to gain new insights that were previously masked by ensemble measurements.¹ Tools such as nanopore sensors have attracted attention for the study of single biomolecules as they provide direct electrical signals indicative of the structure of a molecule as it passes through the nanopore.^{2–6} These nanopores are typically either biological, such as α -hemolysin inserted into a lipid membrane, or inorganic, such as nanometre-scale holes in thin dielectric membranes. The nanopore separates two reservoirs filled with an electrolyte solution, and upon application of a voltage bias across the membrane, an ionic current is driven through the pore. Translocation of single molecules such as DNA through the pore results in a measurable transient current signal.^{7,8} The translocation signal depends on various factors including the electrokinetic driving force, molecule–pore interactions, and length of the molecule.⁹ A number of novel applications have been investigated using nanopore sensors including DNA sequencing,¹⁰ label-free detection of single nucleotide polymorphisms¹¹ and mapping of DNA-binding proteins.¹²

The ability to accurately measure the nucleic acid length and differentiate nucleic acid fragments with different lengths has important implications for diverse applications such as profiling of short tandem repeat (STR) markers,¹³ biomedical

diagnostics,¹⁴ and biodefense.¹⁵ While fluorescence-based methods for length measurements have achieved high measurement resolution ($\sim 5\%$),^{16,17} the need for sources, optics and detectors limits their portability, particularly for on-field applications. Moreover, the scalability of resolution using optical systems to short DNA fragments is typically restricted by diffraction-limited optics. Electrophoresis-based methods¹⁸ also provide high resolution; however, long DNA molecules require specialized methods for separation and there are limitations on the analysis time and sample size. Nanopore sensors offer the potential for rapid, label-free sizing and mapping of DNA at the single molecule level with a direct electrical readout, which may be advantageous over other methods based on fluorescence or electrophoresis. Potentially, nanopores also offer a platform that is easily integrated with microfluidic devices. However, the resolution of nanopores to size or map DNA molecules is currently limited by variability in DNA translocation time and conformations, as well as the limited signal-to-noise ratio arising from insufficient measurement time during translocation.^{19,20} Integrating the area of the current blockade over time (electron charge deficit) has been shown to account for different folded conformations of the translocating DNA molecule,^{19,21} but till date it has not been possible for nanopore sensors to differentiate DNA molecules with lengths differing by less than 30%. A promising approach to overcome this issue is to repeatedly measure the transient current signal of the same molecule, which can result in improved resolution through statistical averaging.²² Gershow and Golovchenko²³ first showed recapture of 4–10 kbp DNA molecules in solid-state nanopores by reversing the applied voltage bias after detection of a translocation event. The recapture process was shown to depend on the competition between diffusion and electrophoresis outside the nanopore, but the recapture probability

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was less than 70%. Hence, while multiple measurements were demonstrated, the low recapture efficiency precluded large numbers of measurements from being reliably performed on each molecule, and the ability of multiple measurements to improve resolution of the nanopore was not shown.

Herein, we report facile fabrication and testing of a nanofluidic device to interrogate single DNA molecules multiple times as transient changes in ionic current through a nanochannel (Fig. 1). The devices were engineered to provide sufficient electric field around the nanochannel for a high recapture probability (>99%) and enabled large numbers of consecutive measurements (>100) to be performed on single molecules. Consequently, we provide the first experimental validation that statistical averaging of the translocation current amplitude over multiple measurements increases the measurement resolution, and quantify the resolution enhancement as a function of the number of measurements. We use this resolution enhancement from multiple measurements to discriminate between a heterogeneous mixture of DNA molecules with different lengths (48.5 kbp λ -DNA and 39.9 kbp T7-DNA molecules), which was not possible using one measurement per molecule. These nanofluidic devices have the advantages of a simple fabrication strategy (micromolding) and ease of integration with microfluidics upstream of the nanochannel.

Materials and methods

Device fabrication

Nanochannel devices were fabricated by micromolding in polydimethylsiloxane (PDMS). 200 nm thick, 500 nm wide titanium

(Ti) metal lines were patterned using an e-beam lithography lift-off technique to define the nanochannel. A thick layer of photoresist (10 μm , SU-8) was patterned on top of the Ti features to define the connecting microchannels. The length of the nanochannels was 4 μm . PDMS was cast on the mold, cured and peeled off, followed by bonding to a glass slide using oxygen plasma. The device comprised two microfluidic channels with a single nanochannel serving as the only electrical and fluidic connection between the two microchannels. The microchannel and nanochannel dimensions were 0.8 cm \times 1 mm \times 10 μm and 4 μm \times 500 nm \times 200 nm (length \times width \times height), respectively.

To prevent adhesion of the DNA molecules to the surface of the nanochannel, the nanochannel surface was first passivated with bovine serum albumin (BSA, Sigma Aldrich). 1 mg mL⁻¹ of BSA was dissolved in 1/15 \times (diluted from 1 \times) phosphate buffered saline (PBS, conductivity of $\sigma = 0.1441 \text{ S m}^{-1}$, pH 7.4, composition given in the ESI†) and flowed into the device for 10 min before introduction of DNA. DNA samples (48.5 kbp λ -DNA, New England Biolabs and 39.9 kbp Bacteriophage T7 DNA, Boca Scientific) were dissolved in 1/15 \times PBS to a final concentration of 1 $\mu\text{g mL}^{-1}$ of λ -DNA, or a mixture of 1 $\mu\text{g mL}^{-1}$ of λ -DNA and 1 $\mu\text{g mL}^{-1}$ of T7 DNA. The DNA solution was then injected on one side of the nanochannel, while the other side was injected with only 1/15 \times PBS solution.

Ag/AgCl electrodes (In Vivo Metric, Healdsburg, CA) were inserted into the microchannel reservoirs on opposite sides of the nanochannel for making electrical contact with the solution. A patch clamp amplifier (Axopatch 200B, Molecular Devices) was used to apply a voltage bias, record the corresponding ionic current through the nanochannel, and filter the data (80 dB dec⁻¹) at 1 kHz cutoff frequency. Current measurements were obtained inside a custom-built Faraday cage to shield the devices from any electromagnetic interference, and the measured current data were digitized and recorded using Digidata 1440A and associated software (Molecular Devices). To perform feedback control and multiple measurements on single molecules, current data were digitized at 20 kHz/16 bits using a Data Acquisition Card (DAQ) (PCI-6251M, National Instruments) controlled by LabView (National Instruments). A shielded connector card (SCB-68) was used to send a voltage signal to the Axopatch amplifier to control the driving voltage ($\pm 1 \text{ V}$) applied by the Axopatch amplifier. Translocation signals at a driving voltage of $\pm 1 \text{ V}$ (the maximum voltage that can be applied by the Axopatch 200B amplifier) yielded sufficiently large translocation amplitudes and durations (20–50 pA, 4 ms) for automated real-time detection with LabView.

Real-time feedback control algorithm

A feedback control algorithm for performing multiple measurements on single DNA molecules was implemented in LabView. To detect a DNA translocation event, the last 1 ms interval of digitized ionic current data was averaged and considered as real-time current i_a , while the 3 ms interval of digitized ionic current prior to the last 1 ms was averaged and considered as open nanochannel current i_b . Since the entry of a DNA molecule into the nanochannel caused an increase of ionic current of $\sim 40 \text{ pA}$ for λ -DNA, and $\sim 30 \text{ pA}$ for T7 DNA,

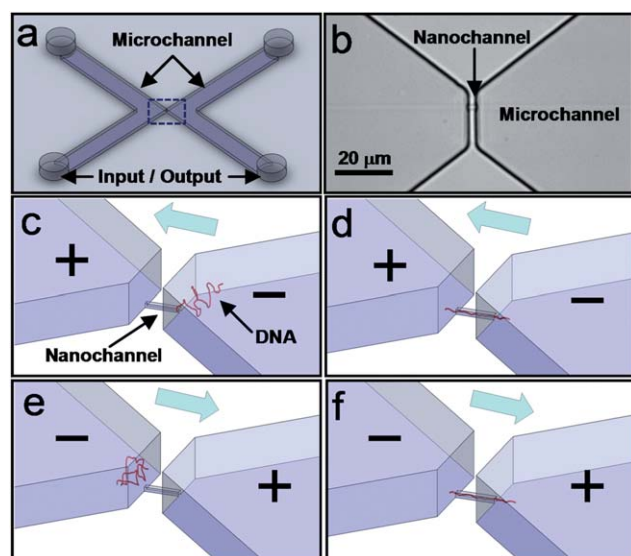


Fig. 1 Nanofluidic device. (a) Schematic of device showing inlet and outlet reservoirs and microchannels. (b) Micrograph of the device (dashed area in (a)) showing a 200 nm \times 500 nm \times 4 μm nanochannel that connects the two microchannels. (c and d) Application of a voltage bias across the nanochannel drives the DNA into the nanochannel. (e and f) Upon completion of the translocation, the voltage bias is reversed to translocate the same molecule in the reverse direction. This process is repeated to obtain multiple translocations of the same DNA molecule through the nanochannel.

the algorithm identified a translocation event when i_a exceeded i_b by a threshold of 10 pA (*i.e.* $\Delta I > 10$ pA). Upon detection of a DNA molecule in the nanochannel, the polarity of the voltage bias was reversed after a delay of 30–40 ms to allow for completion of translocation (duration was 4 ± 2 ms) while avoiding immediate recapture after voltage reversal when the current was rapidly changing. The loop continued, resulting in multiple translocations of the same molecule through the nanochannel.

Data analysis

Translocations were identified and analyzed offline using a second algorithm. To identify a DNA translocation event, a time scale, t_0 , was chosen such that it was small enough that the changing baseline current after voltage reversal could be approximated with a linear regression, but long enough that the translocation signal could not be well approximated. Linear regressions of duration t_0 were fit to the current *vs.* time trace. During the translocation, the linear fit failed to follow the current accurately, thereby increasing the error in the regression (residuals) and allowing for the identification of the translocation. For each segment (current trace between two consecutive voltage reversals), the maximum, initial (first segment right after voltage reversal), and mean residuals (errors) in the linear regression for each segment were computed. A DNA translocation event was defined when the maximum residual in any segment exceeded a threshold value. To accurately determine the amplitude of each translocation, the capacitive transient following voltage reversal was fit and subtracted from the data. Since the translocation waveform had a maximum transient current peak as well as a slight minimum preceding it, the difference between the maximum and minimum current in the vicinity of the translocation was calculated and defined as the translocation amplitude.

Results and discussion

Fig. 1a and b depict the nanofluidic devices comprising a single nanochannel ($200 \text{ nm} \times 500 \text{ nm} \times 4 \text{ }\mu\text{m}$) connecting two microchannels ($0.8 \text{ cm} \times 1 \text{ mm} \times 10 \text{ }\mu\text{m}$) on either side. At 1 V voltage bias with low buffer salt concentrations ($1/15 \times \text{PBS}$), the devices were capable of detecting single molecules of λ -DNA (48.5 kbp) and T-7 DNA (39.9 kbp) from the transient increase in ionic current through the channel (ΔI). The transient current increase is due to the introduction of mobile counter-ions screening the inherent charge on the DNA molecules into the nanochannel.^{7,8} For Δn counterions of mobility μ introduced into a channel of length L , with fraction b of counterions that are mobile, electron charge e , and voltage bias V , the change in ionic current is approximately given by²⁴

$$\Delta I = \mu b \Delta n e V / L^2 \quad (1)$$

Assuming that the entire DNA molecule is enclosed in the nanochannel, we expect a current increase of 15.6 pA for λ -DNA and 12.6 pA for T7 DNA ($\mu = 5.4 \times 10^{-8} \text{ m}^2 \text{ V}^{-1} \text{ s}^{-1}$ for Na^+ and $b = 0.3$ for $\sim 10 \text{ mM}$ NaCl salt concentration^{25,26}), which is consistent with the observed translocation current increase of

20–30 pA. At relevant salt concentrations and surface treatments, the electroosmotic mobility of the BSA-coated surfaces can be estimated: $\mu_{\text{EOF}} = 2 \times 10^{-8} \text{ m}^2 \text{ V}^{-1} \text{ s}^{-1}$.^{27,28} The experimentally measured effective DNA translocation duration of ~ 4 ms given by $L/(\mu_{\text{EP}} - \mu_{\text{EOF}})E$ corresponds to an electrophoretic mobility of roughly $\mu_{\text{EP}} = 2.4 \times 10^{-8} \text{ m}^2 \text{ V}^{-1} \text{ s}^{-1}$, which is consistent with the free-solution mobility of DNA in the presence of Na^+ as measured by Ross and Scruggs.²⁹ Thus, while the DNA molecules moved from the negative to the positive electrode, the electroosmotic flow significantly decreased their velocity.

Multiple measurements on single DNA molecules

The real-time feedback control algorithm was designed to reverse the applied voltage bias after each DNA translocation, thereby performing multiple measurements on the same DNA molecule (Fig. 1c–f). A typical current trace pattern for a solution containing a mixture of $1 \text{ }\mu\text{g mL}^{-1}$ λ -DNA and $1 \text{ }\mu\text{g mL}^{-1}$ T7 DNA (Fig. 2) reveals successive events comprising a DNA translocation followed by a voltage reversal. After voltage reversal, the current changes sign and starts decreasing rapidly to its baseline value. Upon detection of a DNA translocation event, the voltage bias is again reversed after a delay of 30–40 ms (Fig. 2a). The translocation events occur alternately at forward and reverse voltage bias, as expected when the DNA molecule is recaptured into the nanochannel and translocated alternately in forward and reverse directions (Fig. 1c–f). Viewed over a longer period of time (Fig. 2c), the current trace reveals that the translocation events are grouped into distinct sets separated by pauses. The pauses between neighboring sets in Fig. 2c indicate that a translocation event was not detected

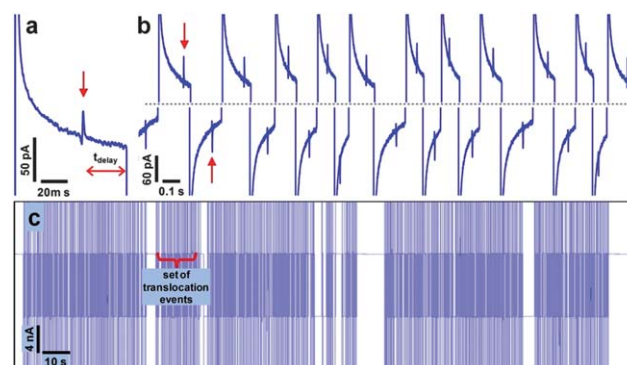


Fig. 2 Current traces with feedback control for a mixture of $1 \text{ }\mu\text{g mL}^{-1}$ λ -DNA and $1 \text{ }\mu\text{g mL}^{-1}$ T7 DNA. (a) Translocation signal of a single DNA molecule (red arrows). Delay time before voltage reversal is also indicated. (b) Current traces showing multiple measurements presumably on the same DNA molecule. When each translocation event (red arrow) is detected, the applied voltage bias is reversed, resulting in successive recapture and translocation of the same molecule. (c) Current traces over a longer period show sets of translocation events separated by pauses that occur whenever the DNA molecule escapes or the translocation is not identified by the real-time algorithm. Vertical lines in the plot correspond to transient current changes that occur when the applied voltage bias is reversed, with each vertical line indicating a voltage reversal following detection of a translocation event.

either following escape of the molecule being recaptured or failure of the algorithm to detect the translocation. Upon detection of another translocation event, the pause terminates and another set of successive measurements begins. The regular translocation events in each set interrupted by much longer pauses suggest that each set corresponds to a single DNA molecule being actively shuttled back and forth through the nanochannel.

Identification of DNA translocation events

Identifying translocations required a method for accurately detecting peaks in ionic current during the decaying current transient that followed voltage reversal (Fig. 2a). For a mixture of $1\ \mu\text{g mL}^{-1}$ λ -DNA and $1\ \mu\text{g mL}^{-1}$ T7 DNA, Fig. 3a depicts the maximum, initial, and mean linear regression residuals in each segment of the current trace between two voltage reversals, with each linear fit spanning a time interval $t_0 = 1.5\ \text{ms}$. The residual from the linear regression for the first segment right after voltage reversal, called initial residual, was typically higher than the mean residual. However, the maximum residuals from the linear regressions were clearly distinguished from the initial and mean residuals (Fig. 3b), and corresponded to the DNA translocations. Hence, DNA translocation events were identified based on a threshold residual that cleanly separated residuals corresponding to DNA translocations from the initial and mean residuals. It was manually verified that initial residuals above the threshold were caused by a translocation event very soon after voltage reversal. Each segment with a maximum residual below the threshold value was manually checked to ensure that it did not contain a translocation (which would correspond to a missed translocation). With this criterion, the data analysis missed less than 1 in 10 000 translocations, and did not incorrectly identify any baseline current transients as translocation events. Identification of DNA translocation events enabled extraction of translocation signal amplitudes and analysis of the current traces.

Recapture time distribution: simulation and experiment

The recapture of a DNA molecule by a nanopore or nanochannel depends on the competition between diffusive and electrokinetic transport. This process can be characterized by a recapture radius such that a molecule within this distance from the pore can be recaptured into the pore with a high probability. This recapture radius scales proportionally with the current through the pore, and inversely with the solution conductivity and molecular diffusivity.²³ In the present device, the large cross-sectional area of the nanochannel, large applied voltage (1 V), and the low diffusivity of the λ -DNA and T7 DNA molecules, all favor a larger recapture radius compared to the case of a nanopore in a membrane.²³ Assuming spherical symmetry, the calculated recapture radius for λ -DNA in the device is $88\ \mu\text{m}$ (see ESI†), which exceeds the microchannel height of $10\ \mu\text{m}$. Thus, beyond a distance of $\sim 10\ \mu\text{m}$, the microchannel further confines the electric field.

To investigate the mechanism of DNA recapture and the role of the microchannels, we compared the molecule recapture time distribution with numerical simulations. The recapture time is defined here as the time between voltage reversal and translocation of the molecule. Prior to recapture, transport of DNA molecules around the nanochannel is influenced by diffusion, electrophoresis, and electro-osmosis. Since the cross-sectional area of the nanochannel is more than an order of magnitude smaller than that of the adjacent microchannel, the DNA transport in this region can be approximated as spherically symmetric for length scales smaller than the microchannel height. Simulation of DNA transport is obtained by solving the one-dimensional time-dependent drift-diffusion equation in spherical coordinates:²³

$$\frac{\partial c(r,t)}{\partial t} = \frac{1}{r^2} \frac{\partial}{\partial r} \left[r^2 \left(\mp \mu E c(r,t) + D \frac{\partial c(r,t)}{\partial r} \right) \right] \quad (2)$$

where DNA diffusivity $D = 4 \times 10^{-13}\ (\text{m}^2\ \text{s}^{-1})$, r is the distance away from the nanochannel, c is the concentration (in this case the probability density) of the DNA molecule, and

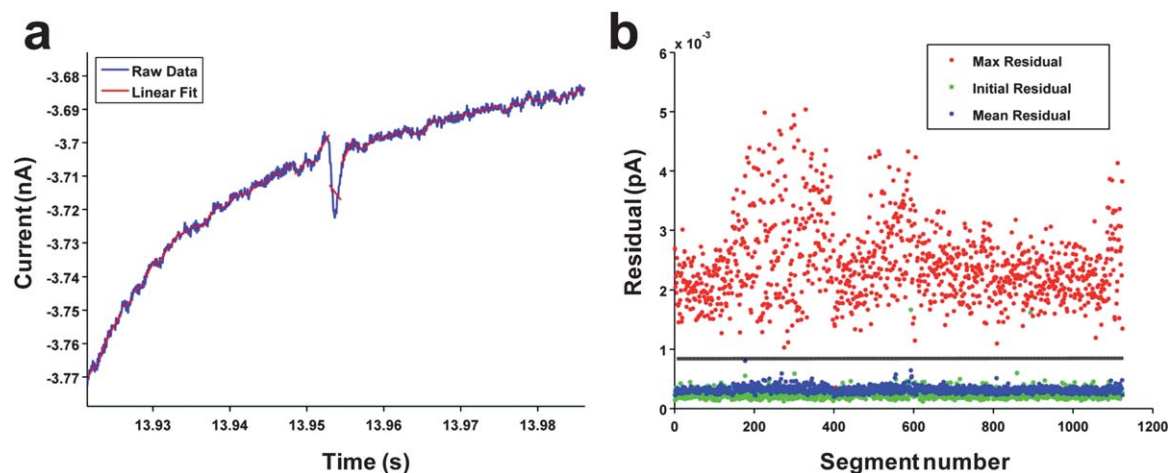


Fig. 3 (a) Visual representation of translocation identification algorithm showing raw data measured by Axopatch 200B (blue) and consecutive linear regressions fit to the data (red). (b) Maximum (red), initial (green) and mean (blue) residuals of consecutive segments of current traces between two voltage reversals. The solid line depicts the threshold residual for determining whether a translocation event occurred in the segment.

$\mu = 4 \times 10^{-9} \text{ (m}^2 \text{ V}^{-1} \text{ s}^{-1}\text{)}$ is the measured DNA mobility. The electric field in the microchannels was obtained by solving Gauss' equation for the device geometry and applied voltage.

The initial condition for the probability distribution, c , was set as a delta function located close to the nanochannel entrance. The probability distribution was allowed to propagate under an electric field for 32.5 ms, which corresponds to the period before the applied voltage is reversed, during which the DNA molecule continues to move away from the nanochannel. The direction of the electric field was then reversed to drive the probability distribution towards the nanochannel. The recapture time distribution was calculated by computing the difference in the probability distribution outside the nanochannel between time steps.

Fig. 4 compares the simulated recapture time distribution for a pre-reversal time of 32.5 ms to the experimentally measured recapture time distribution for pre-reversal times in the range of 30–35 ms. The experimental results agree with the simulation result, where most of the molecules were recaptured within 100 ms, and both experimental and simulation peaks of the recapture time were centered around 40–50 ms. The small discrepancies at the right end of the distribution may be attributed to effects of surface charge,³⁰ interaction between DNA molecules and channel walls, the effect of electric field strength on DNA conformation, and the finite DNA size, which were not included in the simulation.

The simulated mean DNA travel distance during the pre-reversal time (1.38 μm) and the characteristic diffusion length scale during recapture (0.3 μm) were both significantly smaller than the microchannel height (10 μm) and the recapture radius (88 μm , ESI†). These results suggest that the primary mechanism for the high recapture probability in the device was the larger recapture radius due to a combination of the larger nanochannel size, higher applied voltage, and lower DNA diffusivity. The results also indicate that the electric field confinement due to the microchannels did not directly aid in DNA recapture.

Enhanced discrimination of DNA molecules

To enable statistical averaging over multiple measurements on single molecules, it is important to identify sets of translocation events that correspond to back-and-forth translocations of the same molecule. A series of measurements on the same molecule was defined as a set of translocation events with: (a) exactly one

voltage reversal between successive translocations and (b) <500 ms delay between neighboring translocation events (compared to a mean arrival time of ~ 10 s under a constant voltage bias and a mean recapture time of 30 ms after voltage reversal). In an experiment with a mixture of 1 $\mu\text{g mL}^{-1}$ λ -DNA and 1 $\mu\text{g mL}^{-1}$ T7 DNA comprising $N_{\text{trans}} = 4200$ translocations, 32 series were identified that satisfied the above criteria and comprised at least 64 translocation events. Each series contained an average of 130 measurements. Arrangement of extracted translocation signal amplitudes in chronological order revealed distinct shifts in the translocation signals (Fig. 5a), which may be expected if one DNA molecule is switched with another molecule of different length between different series (eqn (1)). We also investigated reasons for termination of a series of consecutive measurements by violation of the above criteria. The analysis revealed that $\sim 25\%$ of terminations (8 series) occurred due to detection of a second translocation event prior to voltage reversal (double translocation), $\sim 28\%$ (9 series) were caused by failure of the real-time Labview algorithm to detect a DNA translocation (no trigger), $\sim 25\%$ (8 series) were caused when the voltage was reversed by the algorithm prior to DNA translocation (false trigger), while $\sim 22\%$ ($n_{\text{escape}} = 7$ series) occurred due to the failure of the DNA molecule to be recaptured (escape). Assuming that the probability of escape is independent of the previous measurement, the recapture probability can be calculated as:

$$p_{\text{recap}} \approx 1 - \frac{n_{\text{escape}}}{N_{\text{trans}}} \quad (3)$$

Eqn (3) yields the recapture probability for 7 escape events out of 4200 translocations as 99.83%.

The distinct shifts in the measured translocation signals become more apparent when successive measurements within each series are averaged over 32 measurements (Fig. 5b). More importantly, the translocation signal histogram appeared to have a single peak when no averaging was performed (Fig. 5a), indicating the failure to discriminate between λ -DNA and T7 DNA molecules in the absence of any chronological information. However, the histogram resolved into three distinct peaks after averaging over 32 measurements (Fig. 5b). In accordance with eqn (1), the mean translocation signal amplitude depends on the number of excess charges introduced into the nanochannel. If the entire DNA molecule is enclosed within the nanochannel during translocation, the number of charges introduced into the channel and hence the translocation signal amplitude is expected to be proportional to the DNA length. Thus, the three peaks (33.4, 41.2, and 64.5 pA, respectively) in the histogram may be assigned to T7 DNA, λ -DNA, and a dimer of λ -DNA that can occur due to cohesive ends of the molecule.¹⁹ Interestingly, the ratio of mean translocation signal amplitudes of T7 and λ -DNA matches the ratio of their lengths within 5%, suggesting that the molecules were enclosed within the nanochannels during translocation. Although the length of a translocating DNA molecule is expected to depend on the electric field,³¹ the equilibrium length L_{eq} of a DNA molecule inside a nanochannel with cross-section dimensions larger than twice the persistence length³² is given by the deGennes theory as:³³

$$L_{\text{eq}} = \frac{(\lambda w)^{1/3}}{d^{2/3}} L \quad (4)$$

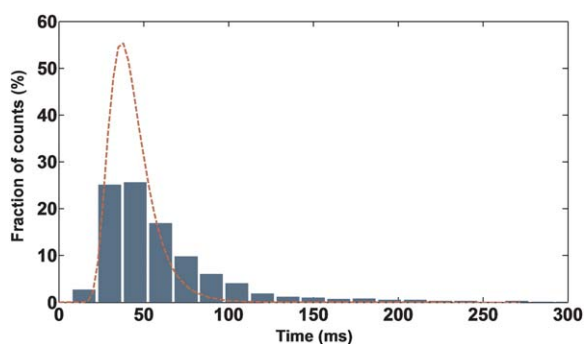


Fig. 4 Experimental histogram and simulation for DNA recapture time distribution (the time between voltage reversal and molecule arrival).

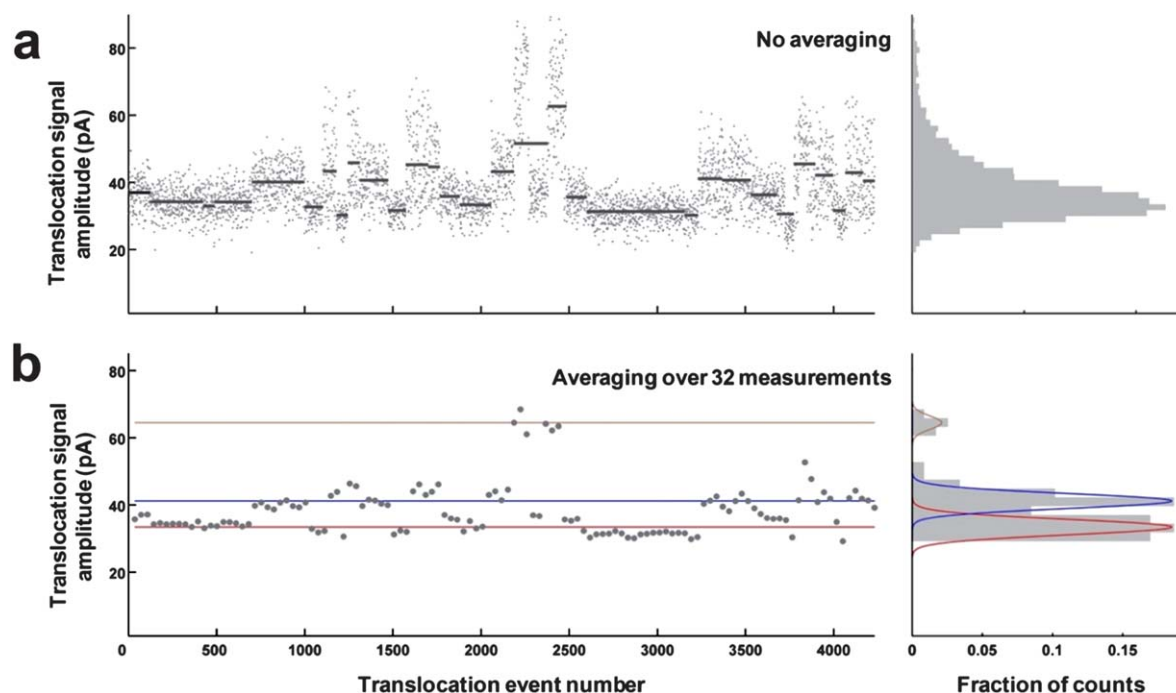


Fig. 5 (a) Chronological sequence of translocation signal amplitudes exhibits distinct shifts in the mean amplitude between series. Horizontal bars denote the mean value of the translocation signal for each series. The histogram of the translocation signal amplitude without any averaging (right). (b) Chronological sequence of translocation current amplitudes after averaging over 32 consecutive measurements within each series accentuates the shifts in the mean amplitude between series. The corresponding histogram of translocation signal amplitudes after averaging over 32 consecutive measurements reveals distinct groups of translocation signals (right). Fit to two Gaussian distributions is depicted; the outlying set is excluded from the fit.

where $d = 316$ nm is the equivalent width of the nanochannel, L is the contour length of the DNA, λ is the persistence length of the DNA, and w is the molecule width. The equilibrium lengths of T7 and λ -DNA molecules inside the nanochannels were calculated to be 1.7 and 2.1 μm , respectively, which is consistent with the above observations. Larger DNA molecules above ~ 100 kbp have an equilibrium length that exceeds the nanochannel length. We observed that the translocation amplitude of the λ -DNA dimer (97 kbp) was only 56% greater than that of λ -DNA, suggesting that it is not completely enclosed in the nanochannel during translocation. Sizing of longer DNA molecules could be easily achieved by increasing the length of the nanochannel so that the whole DNA molecule fits in. On the other hand, although smaller DNA molecules result in proportionally lower translocation current amplitudes, noise in the feedback loop prevents multiple measurements of DNA molecules smaller than about 25 kbp. Sizing of smaller DNA molecules could be achieved by decreasing the channel length, or modulating the cross-sectional area of the nanochannel.³⁴

The experimental results displayed in Fig. 5 also agreed with control experiments with a solution of only λ -DNA molecules, whereby the translocation signal histogram narrowed upon averaging, and did not resolve into multiple peaks (data not shown). These results indicate that averaging over multiple measurements enhanced the measurement resolution and enabled discrimination between T7 DNA and λ -DNA molecules, which was not possible with single measurements. In the ESI†, we demonstrated the ability to detect the switching of a DNA molecule with another molecule of different length.

Furthermore, we illustrate that the specific time sequence of translocation events is fundamental to the improved resolution between the two molecules, and that all differences in the translocation amplitudes between series are lost upon scrambling of the chronological sequence of the translocations. In conjunction with the enhanced discrimination facilitated by averaging over multiple measurements, the results in the ESI† support the conclusion that the majority of the series represent multiple measurements on the same molecule.

Scaling of the standard deviation of the averaged translocation signal amplitudes

Finally, we examined the scaling of the standard deviation of the mean value of the translocation signal amplitude after averaging over n measurements (SDOM or σ_n). Ideally, for a given series of multiple measurements on the same molecule, all measurements sample the same distribution. In this case, for a large number of measurements, the standard deviation of the mean after averaging over n measurements (σ_n) is expected to be smaller than the standard deviation corresponding to single measurements (σ_1) in accordance with the Central Limit Theorem:

$$\sigma_n = \frac{\sigma_1}{\sqrt{n}} \quad (5)$$

Thus, the resolution enhancement after performing multiple measurements arises from the reduced SDOM, and the \sqrt{n}^{-1} scaling, in principle, determines the number of measurements per molecule needed to resolve its length to a given degree. To

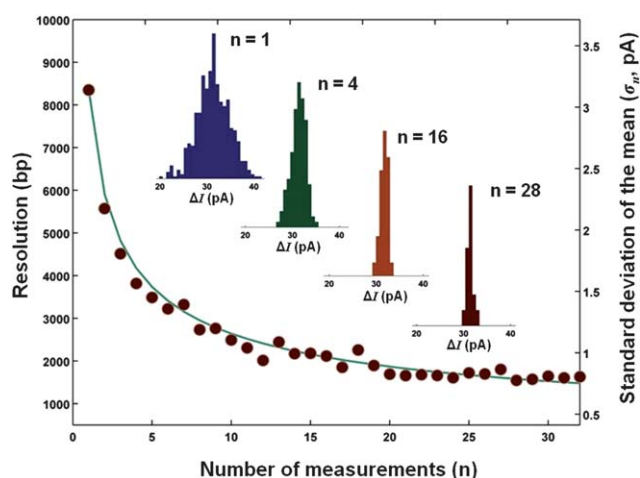


Fig. 6 The effect of averaging over multiple measurements on the resolution and standard deviation of the averaged translocation signals. Symbols denote experimental data while the solid line denotes a $1/\sqrt{n}$ fit. Histograms of the translocation signal amplitude (ΔI) corresponding to $n = 1$ and averaging over $n = 4, 16$, and 28 are shown in the insets.

quantify σ_n , a single series with $N > 500$ measurements was divided into $\text{floor}(N/n)$ smaller sub-series. The standard deviation of the mean for each sub-series of n measurements (up to 32) was computed. Within this series, σ_n scaled inversely with \sqrt{n} , decreasing from 3.4 pA for a single measurement to 0.6 pA for $n = 28$ (Fig. 6). This scaling behavior was also typical for other series. Assuming that the DNA translocation signal scales linearly with the length of the molecule, we can estimate the resolution as follows:

$$\frac{R}{L_{\text{DNA}}} = \frac{2\sigma_n}{\Delta I_n} \quad (6)$$

Here, the resolution R is expressed in terms of the DNA length L_{DNA} (39.9 kbp), and ΔI_n is the mean translocation signal (31.4 pA). Based on eqn (6), the resolution is expected to decrease from 8.5 kbp for $n = 1$ to 1.5 kbp for $n = 32$, as depicted in Fig. 6.

The histogram consisting of all DNA molecules measured (Fig. 5b), however, displayed a weaker scaling with the number of measurements. For $n = 32$, while the λ -DNA and T7 DNA molecules were easily distinguished, the standard deviations of the two Gaussian curves fit were 2.2 and 2.3 pA, corresponding to a measurement resolution of 4.5 kbp. The histogram over all DNA molecules includes translocations that occur over approximately two hours. As demonstrated in Fig. 6, the standard deviation scaling within the timescale of a single series (5–10 seconds) is consistent with a $1/\sqrt{n}$ scaling. Thus, the weaker standard deviation scaling for the cumulative set of all translocations is a result of temporal fluctuations in the mean translocation amplitude between series, which occurs on a timescale longer than that of a single series. Since the fluctuations in mean translocation current displayed no functional dependence on the baseline nanochannel current, it is likely that such fluctuations were caused by secondary effects, such as fluctuations in the nanochannel temperature, surface charge, or buffer conditions between series, similar to sources that lead to $1/f$ noise. Better control of these variables could mitigate these effects in future devices.

Conclusions

In conclusion, we have demonstrated enhancement of the ability to discriminate between DNA molecules of different lengths using multiple measurements in a nanochannel device. The device can recapture DNA molecules with a probability of greater than 99%, and consistently perform >100 measurements on single DNA molecules. The DNA sizing resolution improved from ~ 8.5 kbp to ~ 4.5 kbp by statistical averaging over multiple measurements, enabling discrimination between 48.5 kbp λ -DNA and 39.9 kbp T7 DNA molecules, which was not possible with a single measurement per molecule. The devices can be rapidly fabricated using micromolding and easily integrated with microfluidic channels. This work demonstrates the feasibility of enhancing the ability of nanochannel or nanopore sensors for DNA sizing via multiple measurements on the same molecule. Further development of this approach may enable rapid, label-free, high-resolution electrical sizing of DNA lengths by nanofluidic devices, which has hitherto been largely confined to electrophoresis.

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References

- W. E. Moerner, *Proc. Natl. Acad. Sci. U. S. A.*, 2007, **104**, 12596–12602.
- S. Howorka and Z. Siwy, *Chem. Soc. Rev.*, 2009, **38**, 2360–2384.
- M. Karhanek, J. T. Kemp, N. Pourmand, R. W. Davis and C. D. Webb, *Nano Lett.*, 2005, **5**, 403–407.
- O. A. Saleh and L. L. Sohn, *Biophys. J.*, 2002, **82**, 166a.
- O. A. Saleh and L. L. Sohn, *Proc. Natl. Acad. Sci. U. S. A.*, 2003, **100**, 820–824.
- O. A. Saleh and L. L. Sohn, *Rev. Sci. Instrum.*, 2001, **72**, 4449–4451.
- R. Fan, R. Karnik, M. Yue, D. Y. Li, A. Majumdar and P. D. Yang, *Nano Lett.*, 2005, **5**, 1633–1637.
- R. M. M. Smeets, U. F. Keyser, D. Krapf, M. Y. Wu, N. H. Dekker and C. Dekker, *Nano Lett.*, 2006, **6**, 89–95.
- A. J. Storm, C. Storm, J. H. Chen, H. Zandbergen, J. F. Joanny and C. Dekker, *Nano Lett.*, 2005, **5**, 1193–1197.
- D. Branton, D. W. Deamer, A. Marziali, H. Bayley, S. A. Benner, T. Butler, M. Di Ventra, S. Garaj, A. Hibbs, X. H. Huang, S. B. Jovanovich, P. S. Krstic, S. Lindsay, X. S. S. Ling, C. H. Mastrangelo, A. Meller, J. S. Oliver, Y. V. Pershin, J. M. Ramsey, R. Riehn, G. V. Soni, V. Tabard-Cossa, M. Wanunu, M. Wiggin and J. A. Schloss, *Nat. Biotechnol.*, 2008, **26**, 1146–1153.
- S. M. Iqbal, D. Akin and R. Bashir, *Nat. Nanotechnol.*, 2007, **2**, 243–248.
- S. W. Kowalczyk, A. R. Hall and C. Dekker, *Nano Lett.*, 2010, **10**, 324–328.
- C. R. Reedy, K. A. Hagan, D. J. Marchiarullo, A. H. Dewald, A. Barron, J. M. Bienvenue and J. P. Landers, *Anal. Chim. Acta*, 2011, **699**, 126.
- H. C. Fan, Y. J. Blumenfeld, U. Chitkara, L. Hudgins and S. R. Quake, *Clin. Chem.*, 2010, **56**, 1279–1286.
- J. B. H. Tok, *Nano and Microsensors for Chemical and Biological Terrorism Surveillance*, Royal Society of Chemistry, Cambridge, 2008.

- 16 H. P. Chou, C. Spence, A. Scherer and S. Quake, *Proc. Natl. Acad. Sci. U. S. A.*, 1999, **96**, 11–13.
- 17 M. Foquet, J. Korlach, W. Zipfel, W. W. Webb and H. G. Craighead, *Anal. Chem.*, 2002, **74**, 1415–1422.
- 18 J. Herschleb, G. Ananiev and D. C. Schwartz, *Nat. Protoc.*, 2007, **2**, 677–684.
- 19 D. Fologea, E. Brandin, J. Uplinger, D. Branton and J. Li, *Electrophoresis*, 2007, **28**, 3186–3192.
- 20 D. W. Deamer and M. Akeson, *Trends Biotechnol.*, 2000, **18**, 147–151.
- 21 D. Fologea, M. Gershow, B. Ledden, D. S. McNabb, J. A. Golovchenko and J. L. Li, *Nano Lett.*, 2005, **5**, 1905–1909.
- 22 L. I. Berge, J. Feder and T. Jossang, *Rev. Sci. Instrum.*, 1989, **60**, 2756–2763.
- 23 M. Gershow and J. A. Golovchenko, *Nat. Nanotechnol.*, 2007, **2**, 775–779.
- 24 Y. H. Sen and R. Karnik, *Anal. Bioanal. Chem.*, 2009, **394**, 437–446.
- 25 T. Vuletic, S. D. Babic, D. Grgicin, D. Aumiler, J. Radler, F. Livolant and S. Tomic, *Phys. Rev. E: Stat., Nonlinear, Soft Matter Phys.*, 2011, **83**, 041803.
- 26 G. S. Manning, *J. Chem. Phys.*, 1969, **51**, 924.
- 27 W. Schrott, Z. Slouka, P. Cervenka, J. Ston, M. Nebyla, M. Pribyl and D. Snita, *Biomicrofluidics*, 2009, **3**, 044101.
- 28 F. Zhou, W. Wang, W. Y. Wu, J. R. Zhang and J. J. Zhu, *J. Chromatogr., A*, 2008, **1194**, 221–224.
- 29 P. D. Ross and R. L. Scruggs, *Biopolymers*, 1964, **2**, 231–236.
- 30 D. Stein, Z. Deurvorst, F. H. J. van der Heyden, W. J. A. Koopmans, A. Gabel and C. Dekker, *Nano Lett.*, 2010, **10**, 765–772.
- 31 J. T. Mannion, C. H. Reccius, J. D. Cross and H. G. Craighead, *Biophys. J.*, 2006, **90**, 4538–4545.
- 32 W. Reisner, K. J. Morton, R. Riehn, Y. M. Wang, Z. N. Yu, M. Rosen, J. C. Sturm, S. Y. Chou, E. Frey and R. H. Austin, *Phys. Rev. Lett.*, 2005, **94**, 196101.
- 33 J. O. Tegenfeldt, C. Prinz, H. Cao, S. Chou, W. W. Reisner, R. Riehn, Y. M. Wang, E. C. Cox, J. C. Sturm, P. Silberzan and R. H. Austin, *Proc. Natl. Acad. Sci. U. S. A.*, 2004, **101**, 10979–10983.
- 34 D. Huh, K. L. Mills, X. Y. Zhu, M. A. Burns, M. D. Thouless and S. Takayama, *Nat. Mater.*, 2007, **6**, 424–428.