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## Stepwise Hydration of Protonated Carbonic Acid: A Theoretical Study

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The gas-phase geometries, binding energies (BEs), and sequential binding energies (SBEs) of protonated carbonic acid (PCA)–water (W) clusters ( $\text{PCAW}_n$ , where  $n = 1–6$ ) have been calculated using density functional theory (DFT) with Becke's three-parameter hybrid exchange functional and the Lee–Yang–Parr correlation functional (B3LYP) and M05-2X methods. The presence of wirelike structures of protonated water in  $\text{PCAW}_n^x$  clusters is evident from the results. The results indicate that a proton is transferred from PCA to its immediate water molecule in the linear and monohydroxy clusters of PCA. The involvement of the Eigen cation and Grothuss type of mechanism in the proton transport is observed from the sequential hydration energies and from the calculated vibrational spectra. Although geometrical parameters clearly reveal the presence of the Eigen core, calculated lower-energy vibrational modes provide clues about the involvement of the sequence Eigen → Zundel → Eigen in the proton transfer.

### 1. Introduction

Because hydrolysis of  $\text{CO}_2$  is an important process in natural science, several experimental and theoretical studies have been carried out on carbonic acid ( $\text{H}_2\text{CO}_3$ ).<sup>1–11</sup> These studies have revealed that both carbonic acid (CA) and the weakly bound isomer of a  $\text{H}_2\text{O} \cdots \text{CO}_2$  complex are short-lived intermediates in the hydrolysis of  $\text{CO}_2$ . The formation of  $\text{H}_2\text{CO}_3$  from  $\text{CO}_2$  and  $\text{H}_2\text{O}$  has been investigated using a variety of quantum chemical methods.<sup>12</sup> In the same study, the role of microsolvation of  $\text{CO}_2$  in the formation of CA was also addressed.<sup>12</sup> High-resolution spectroscopic techniques have provided indirect evidence for the existence of CA in the gas phase.<sup>13</sup> On the other hand, theoretical studies have shown that isolated  $\text{H}_2\text{CO}_3$  should be stable, although it decomposes into  $\text{H}_2\text{O}$  and  $\text{CO}_2$  in the presence of water.<sup>14</sup> Recently, Andrei et al. detected the protonated carbonic acid  $\text{C}(\text{OH})_3^+$  in the gas phase for the first time using high-resolution infrared (IR) spectroscopy.<sup>15</sup> They also used density functional theory (DFT) with Becke's three-parameter hybrid exchange functional and the Lee–Yang–Parr correlation functional (B3LYP) to predict the structure and stability of different isomers of PCA.<sup>15</sup>

Previous quantum chemical calculations on PCA have predicted the existence of different isomers on the  $[\text{CH}_3\text{O}_3]^+$  potential-energy surface.<sup>16</sup> The linear hydrogen-bonded (H-bonded)  $\text{H}_3\text{O}^+ \cdots \text{CO}_2$  complex (A) is the most stable isomer.<sup>17</sup> The protonation of CA at the carbonyl group leads to two planar isomers: the *anti*-trihydroxycarbenium ion ( $C_{3h}$ ) and the *syn*-trihydroxycarbenium ion ( $C_s$  symmetry). In this work, these two complexes are referred to as B and C, respectively. The fourth isomer of  $[\text{CH}_3\text{O}_3]^+$  arises as a result of the protonation of  $\text{H}_2\text{CO}_3$  at a hydroxyl group. It is a nonplanar molecule with  $C_1$  symmetry. This hitherto undetected complex is designated as

D. Previous DFT calculations on these molecules have shown the stabilities of these isomers to be in the order A > B > C > D.<sup>16</sup>

Egsgaard et al. reported that PCA is generated by the consecutive elimination of a vinyl radical and an alkene from the radical cation of dialkylcarbonates.<sup>16a</sup> They found that protonated carbonic acid forms symmetrical trihydroxy species and that fragmentation of unsymmetrical isomer leads to protonated carbon dioxide and water. Olah and White illustrated that PCA is remarkably stable in superacidic media and might play an important role in the biological carboxylation process, which involves a proton-transfer mechanism.<sup>18a</sup> Nuclear magnetic resonance (NMR), IR, Raman, and X-ray crystallography studies of  $\text{C}(\text{OH})_3^+$  in superacid solutions or corresponding salts confirm the existence of isomer B.<sup>18</sup> Recently, the structure of PCA in the gas phase was studied using mid-IR spectroscopy.<sup>19</sup> Combined evidence gathered by these studies points out that isomer B is stable in both the solution and gas phases.<sup>18,19</sup>

With a view toward understanding proton transfer in various model systems, numerous experimental and theoretical studies have been carried out.<sup>20</sup> The role of the proton in  $\text{H}^+(\text{H}_2\text{O})_n$  clusters ( $n = 1–27$ ) has been investigated using high-resolution IR spectroscopy techniques and various levels of electronic structure calculations.<sup>21</sup> The presence of Eigen<sup>22</sup> ( $\text{H}_3\text{O}^+$ ) and Zundel<sup>23</sup> ( $\text{H}_2\text{O} \cdots \text{H}^+ \cdots \text{H}_2\text{O}$ ) cations in protonated water clusters was observed in earlier investigations. It has been found that small protonated water clusters ( $n \sim 10$ ) form chain structures and that nanometer-scale cages form for  $n > 21$ .<sup>21a</sup> The dynamics of proton transport in water has received widespread attention.<sup>24</sup> It is possible to note that the position of the excess proton in these clusters varies sensitively with the number of water molecules and the geometry of the clusters.<sup>25–27</sup> Voth and co-workers used multistate empirical valence bond theory and ab initio molecular dynamics simulations to study hydrated protons in methanol–water solutions.<sup>28–31</sup> These investigations provided the modern version of Grothuss-like mechanism that is typically described as proton transfer from  $\text{H}_3\text{O}^+$  to  $\text{H}_2\text{O}$ .<sup>25,32</sup>

The transfer of  $\text{H}^+$  along a H-bonded water chain, or proton wire, is thought to be an important mechanism in various

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chemical and biological processes.<sup>33</sup> The structure of a proton wire inside a carbon nanotube was investigated to understand the mechanism of proton conduction in biological systems and the effect of structural confinement on the proton-transport process.<sup>34</sup> Water wires are responsible for the selective proton conductivity in pure lipid bilayers,<sup>32</sup> the ATP synthase complex,<sup>35</sup> bacteriorhodopsin,<sup>36</sup> and pendadecapeptide gramicidin A.<sup>37</sup> In human carbonic anhydrase II, a chain of three water molecules plays a crucial role in the proton transfer.<sup>38</sup> The application of proton wires in molecular switches has also been reported.<sup>39</sup> The proton-transfer mechanism between aqueous Brønsted acids and bases has been studied in real time using ultrafast IR spectroscopy.<sup>40</sup> This study demonstrated the Grotthuss-type sequential proton-hopping mechanism through water bridges in acid–base reactions.

Although numerous studies have been carried out on proton wires, we felt that addressing the question of how a protonated molecule interacts with water molecules and conducts proton would enhance our understanding of proton transfer in other systems. Hence, we undertook a systematic study on the sequential hydration of isomers A–C of PCA. In this context, selection of the appropriate quantum chemical computational method to probe the structure and stability of these hydrated clusters is necessary.

An analysis of various quantum chemical methods for probing molecular hydration has been made.<sup>27,41</sup> In a recent work, hydration of different mono- and divalent metal ions was studied using DFT(B3LYP); Møller–Plesset second-order perturbation (MP2) and coupled-cluster theory with single, double, and triple excitations [CCSD(T)]; and the G3 quantum chemical method.<sup>42a</sup> Double- and triple- $\zeta$  basis sets containing both polarization and diffuse functions were employed in this study. Total and sequential binding energies for hydration of metal ions were calculated for metal–water clusters containing one to six water molecules. In the same study, a systematic benchmarking of the performance of the above-mentioned methods was reported.<sup>42a</sup> Although no single method was found to consistently show excellent performance for the hydration of all metal ions, the B3LYP/6-311++G\*\* method seemed to be the most economical computational method. It was also noted that this method provides correct trends in sequential hydration of energies of different metal ions containing more than four water molecules. Earlier studies on stepwise hydration by the research groups of the authors of the present work, carried out at the HF or DFT level of theory,<sup>42b–d</sup> included substrates such as crown ether, cytosine dimer, and uracil. Recently, several new density functionals have been developed for addressing a variety of issues in chemical and biological systems.<sup>43</sup> For example, Truhlar and co-workers found that the new hybrid meta exchange-correlation functional M05-2X is quite suitable for probing H-bonding and noncovalent interactions as compared to the other functionals.<sup>44</sup> Hence, the present study employs Hartree–Fock (HF), DFT(B3LYP), and DFT(M05-2X) methods for probing the structures and stabilities of hydrated clusters of isomers A–C of PCA.

## 2. Computational Details

The molecular structures and hydration patterns of isomers A–C are presented in Scheme 1. Various initial and corresponding optimized geometries considered in this study are provided in the Supporting Information (Figure S1). The water clusters of isomers A, B, and C are designated, respectively, as AW<sub>n</sub>, BW<sub>n</sub>, and CW<sub>n</sub> (where n = 1–6).

As seen in Scheme 1, three different patterns of hydration of isomer A are considered, corresponding to (i) linear (l), (ii)

branched (b), and (iii) cyclic (c) hydration of the H<sub>3</sub>O<sup>+</sup> fragment in PCA. These clusters are referred to as AW<sub>n</sub><sup>l</sup>, AW<sub>n</sub><sup>b</sup>, and AW<sub>n</sub><sup>c</sup> (n = 1–6), respectively. It can be seen from Figure S1 (Supporting Information) that clusters with CO<sub>2</sub>–W interactions were first thought to be important. However, these initial geometries were found to undergo reorganization during geometry optimization. Further, it may be noted that, after optimization, water molecules tend to interact with the H<sub>3</sub>O<sup>+</sup> moiety rather than with CO<sub>2</sub>. Hence, in this study, the AW<sub>n</sub><sup>l</sup>, AW<sub>n</sub><sup>b</sup>, and AW<sub>n</sub><sup>c</sup> (n = 1–6) types of clusters were considered only for completeness.

Three different patterns of hydration in isomers B and C considered in this study are (i) interaction of water molecules with any one of the hydroxyl groups, referred to as monohydroxy hydration (BW<sub>n</sub><sup>m</sup>, and CW<sub>n</sub><sup>m</sup>); (ii) interaction of water molecules with any two hydroxyl groups, designated as dihydroxy hydration (BW<sub>n</sub><sup>d</sup> and CW<sub>n</sub><sup>d</sup>); and (iii) interaction of water molecules with all of the hydroxyl groups, referred to as trihydroxy hydration (BW<sub>n</sub><sup>t</sup> and CW<sub>n</sub><sup>t</sup>).

The geometries of all clusters were optimized using the HF, DFT(B3LYP), and DFT(M05-2X) methods with the 6-31+G\* and 6-311++G\*\* basis sets. All calculations were performed using the Gaussian 98W<sup>45</sup> and Gaussian 03 (revision E.01) suites of programs.<sup>46</sup> Because the reliability of the B3LYP/6-31+G\* level of calculation in predicting the vibrational frequencies of protonated clusters was confirmed in previous studies,<sup>47,48</sup> the same method was used here. Vibrational frequencies were scaled by a factor of 0.973.<sup>47</sup> The geometries of all clusters were minima on their respective potential energy surfaces at the B3LYP/6-31+G\* level of theory. The binding energies (BEs) of all clusters were calculated using the supermolecule approach and corrected for basis set superposition error (BSSE) using the counterpoise (CP) procedure suggested by Boys and Bernardi<sup>49</sup>

$$\text{BE} = -(E_{\text{cluster}} - \sum_{i=1}^n E_i) \quad (1)$$

where E<sub>cluster</sub> is the total energy of the cluster, E<sub>i</sub> is energy of the monomer, and n is the total number of monomers in the cluster. Specifically, BSSE was estimated for each monomer by computing the energy corresponding to the geometry in the cluster with the n-mer basis set. The BSSE-corrected sequential binding energies (SBEs) of various clusters were calculated as per eq 1 using the procedure explained in an earlier report<sup>42</sup>

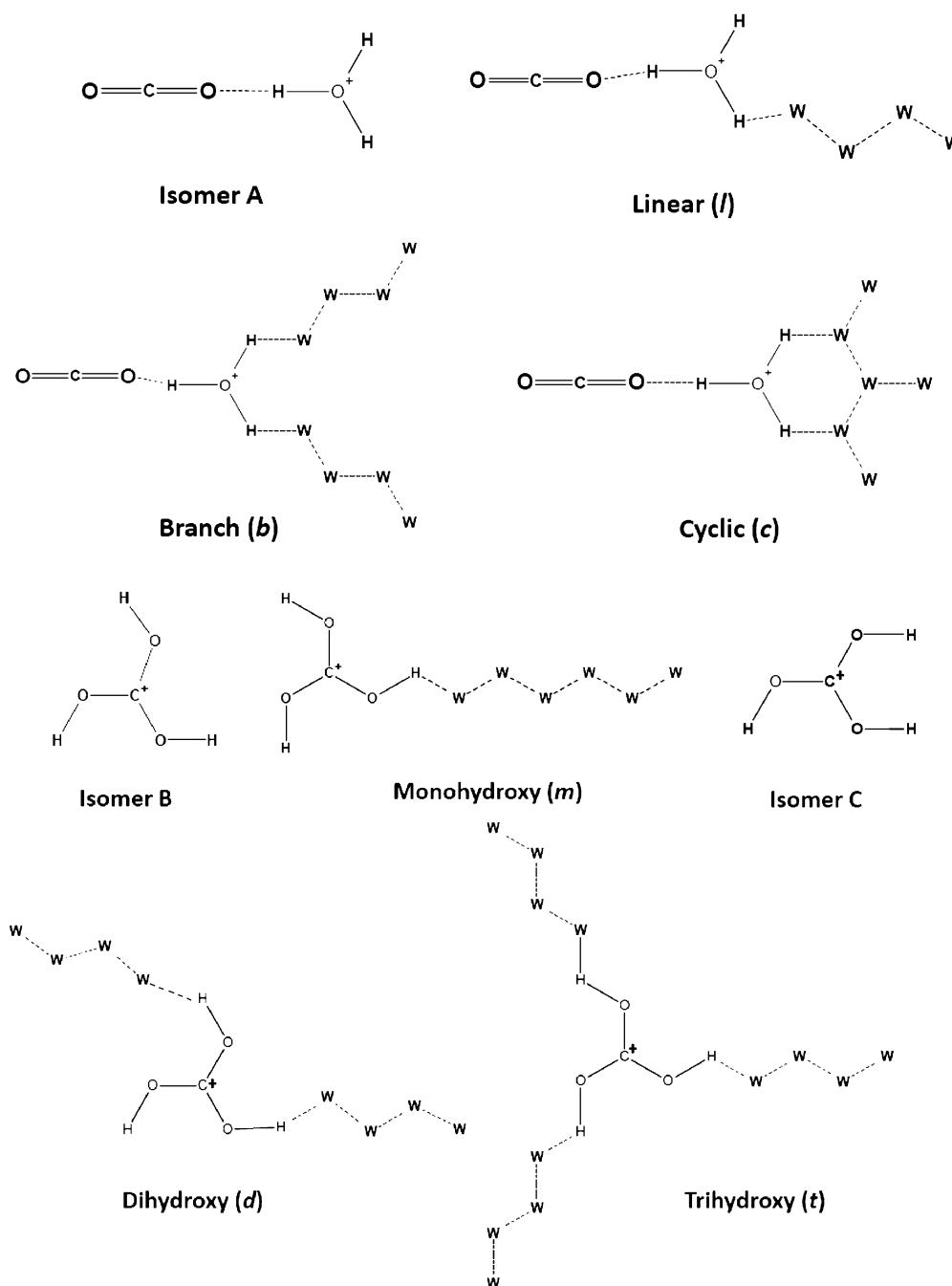
$$E_{\text{SBE}} = -\{E_{(\text{PCAW}_n)} - [E_{(\text{PCAW}_{n-1})} + E_{(\text{H}_2\text{O})}]\} \quad (2)$$

Here, E<sub>SBE</sub> is the sequential binding energy and E<sub>(PCAW<sub>n</sub>)</sub>, E<sub>(PCAW<sub>n-1</sub>)</sub>, and E<sub>(H<sub>2</sub>O)</sub> are the total energies of clusters with n water molecules, n – 1 water molecules, and one water molecule, respectively, obtained at the B3LYP/6-311++G\*\* level.

## 3. Results and Discussion

**3.1. Hydration of Isomer A.** **3.1.1. Linear-Type Clusters (AW<sub>n</sub><sup>l</sup>).** Scheme 1 shows that the most stable structure of isomer A exists as CO<sub>2</sub>•••H<sub>3</sub>O<sup>+</sup>. In this isomer, water molecules strongly interact with the H<sub>3</sub>O<sup>+</sup> moiety of PCA and form structures that are similar to those of protonated water clusters.<sup>21</sup> Figure 1 illustrates the optimized geometries of AW<sub>n</sub><sup>l</sup> (n = 1–6), along with the atom numbering and bond distances obtained

SCHEME 1: Hydration Patterns of Protonated Carbonic Acid (PCA): Isomers A, B, and C

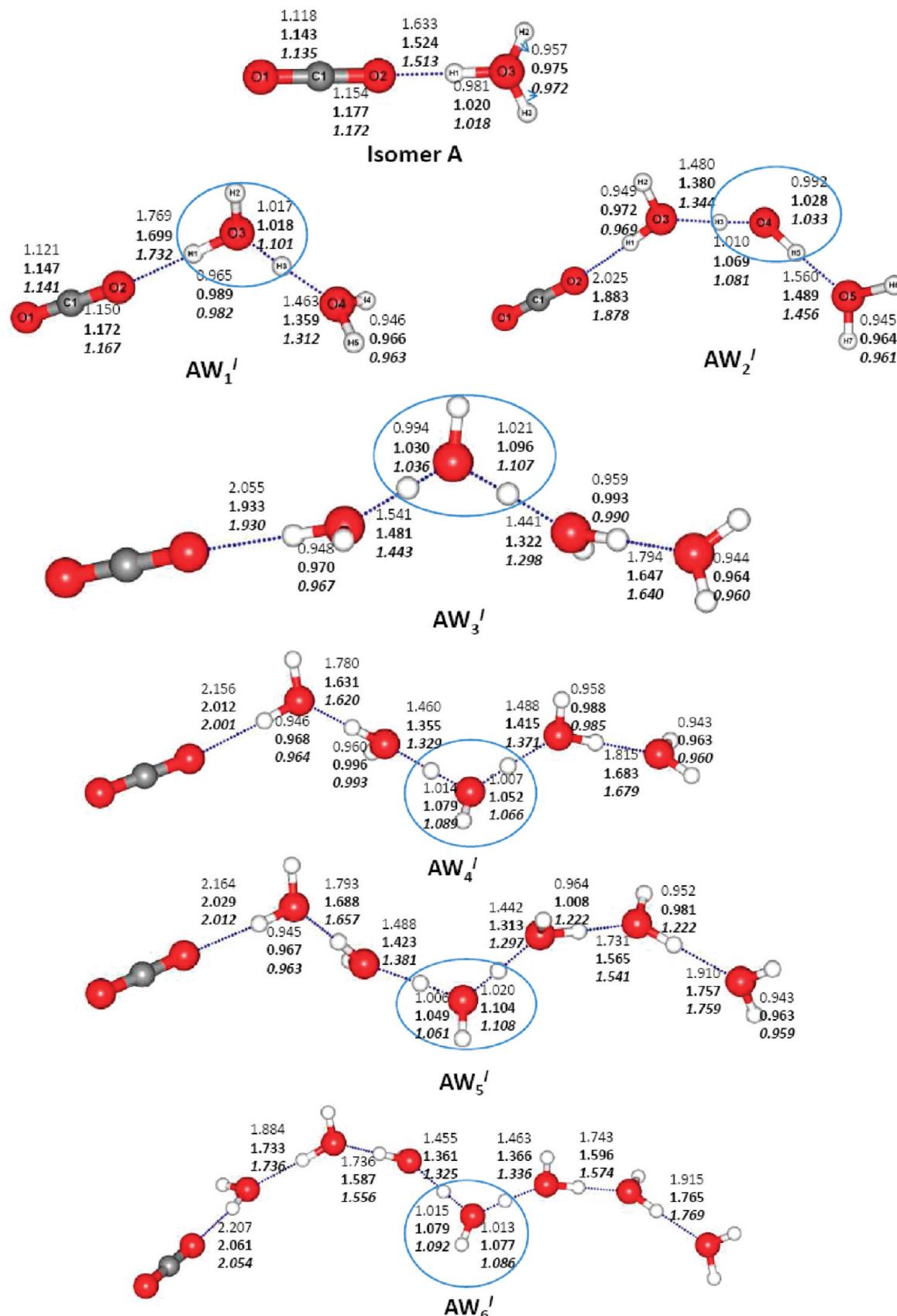


from HF (normal), B3LYP (bold), and M05-2X (bold and italic) calculation using the 6-311++G\*\* basis set. Although calculations were carried out using three different levels of theory, the geometrical parameters obtained from DFT (M05-2X)/6-311++G\*\* are considered for analysis. In isomer A, the O1–C1 and C1–O2 bond lengths are 1.135 and 1.172 Å, respectively. The O2···O3 H-bond distance is 2.531 Å. The distance between O2 and H1 is 1.513 Å. The O–H bond lengths of the H<sub>3</sub>O<sup>+</sup> fragment, which are not involved in the H-bonding, are equal to 0.972 Å.

In cluster AW<sub>1</sub><sup>1</sup>, the water molecule forms a H-bond with the H<sub>3</sub>O<sup>+</sup> fragment. This cluster consists of two H-bonds: O2···O3 (2.714 Å) and O3···O4 (2.413 Å). It can be observed from these H-bonding distances that, compared to the O2···O3 distance in isomer A, the one in AW<sub>1</sub><sup>1</sup> is 0.2 Å larger. This indicates that the H<sub>3</sub>O<sup>+</sup> is closer (H-bonded) to the water

molecule than the CO<sub>2</sub> fragment. Further, increases in the O2···H1 and O3···H3 distances and a decrease in the O3–H1 distance can also be observed when compared to the values for isomer A. The shifting of the proton from isomer A to water is evident from the changes in the geometrical parameters.

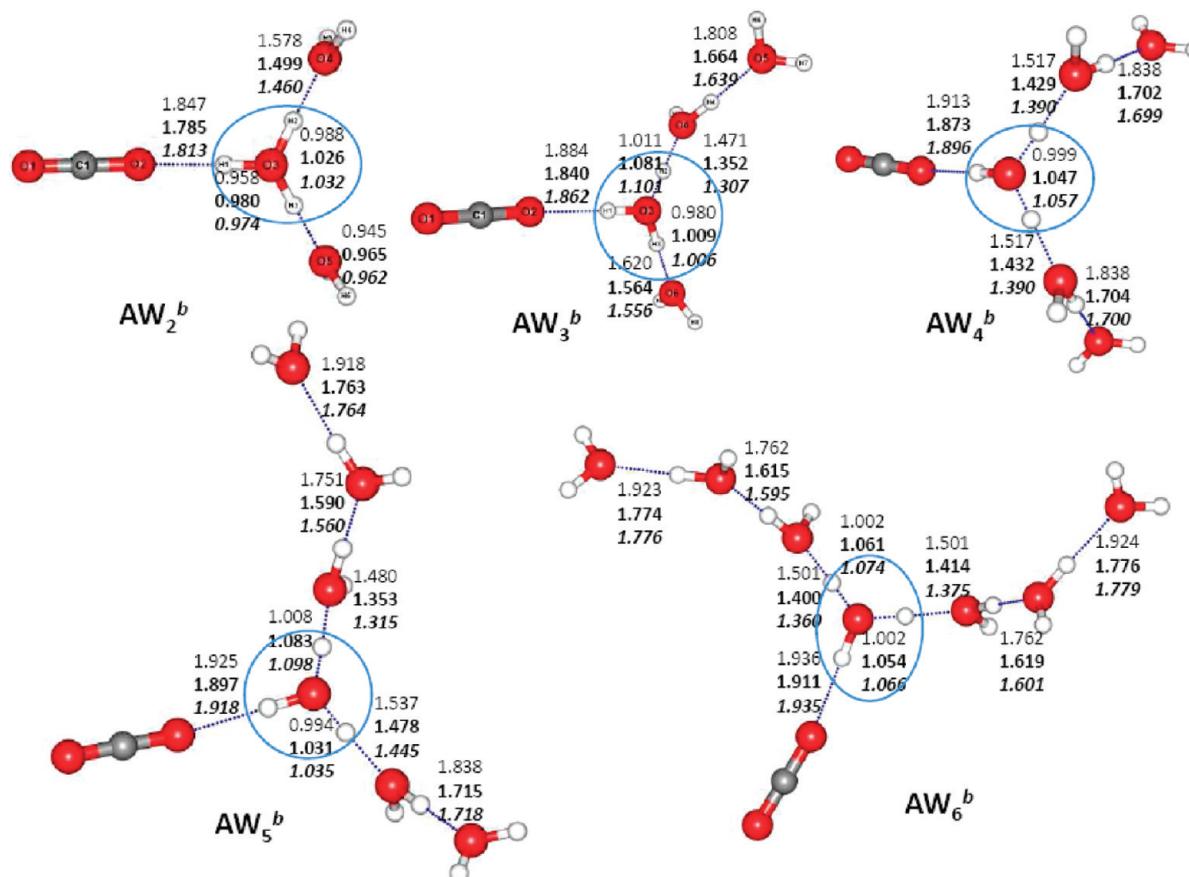
Similar variations in the geometrical parameters of AW<sub>2</sub><sup>1</sup> can be found from Figure 1. This cluster has three H-bonds, viz., O2···O3, O3···O4, and O4···O5. The corresponding distances are 2.847, 2.424, and 2.488 Å. These values clearly reveal that the proton shifts from the H<sub>3</sub>O<sup>+</sup> moiety to the neighboring water molecules. Because of proton transport from isomer A to water, the geometrical parameters of the other parts of the cluster undergo considerable change. As water molecules are sequentially added to AW<sub>1</sub><sup>1</sup> in a linear fashion, proton transport takes place from one water molecule to its immediate neighbor in a manner akin to the Grotthuss mechanism. A similar pattern in



**Figure 1.** Optimized geometries of  $\text{AW}_n^1$  linear clusters (where  $n = 1$ –6) at the HF/6-311++G\*\* (normal), B3LYP/6-311++G\*\* (bold), and M05-2X/6-311++G\*\* (bold and italic) levels of theory. The blue circles indicate the presence of the Eigen core cation. Distances are in angstroms.

proton transport has been observed in the other clusters. Figure 1 shows that the H-bond distances in the middle region of the water wire are shorter than those in the peripheral regions.

**3.1.2. Branched-Type Clusters ( $\text{AW}_n^b$ ).** The optimized geometries of  $\text{AW}_n^b$  clusters (where  $n = 2$ –6) are depicted in Figure 2. The geometrical parameters indicate that there is no



**Figure 2.** Optimized geometries of  $\text{AW}_n^b$  branched clusters (where  $n = 2\text{--}6$ ) at the HF/6-311++G\*\* (normal), B3LYP/6-311++G\*\* (bold), and M05-2X/6-311++G\*\* (bold and italic) levels of theory. The blue circles indicate the presence of the Eigen core cation. Distances are in angstroms.

proton movement in these clusters. Further, in  $\text{AW}_2^b$ , the proton is bound to the CO<sub>2</sub> fragment and two water molecules in the first solvation shell. It was found from the calculations that further addition of water molecules does not favor proton transport in any preferential direction.

**3.1.3. Cyclic Clusters ( $\text{AW}_n^c$ ).** Earlier studies on protonated water clusters have provided information about the existence of cyclic clusters.<sup>48a</sup> Similar structures are also observed in isomer A. The optimized geometries of cyclic clusters  $\text{AW}_n^c$  ( $n = 3\text{--}6$ ) are presented in Figure 3. The tetrameric core structure stabilizes the  $\text{AW}_n^c$  clusters ( $n = 3\text{--}6$ ), except in the case of  $\text{AW}_5^{c1}$  (cyclic hexamer). The geometrical parameters confirm the involvement of the Eigen cation in proton transfer in all of the A-type clusters.

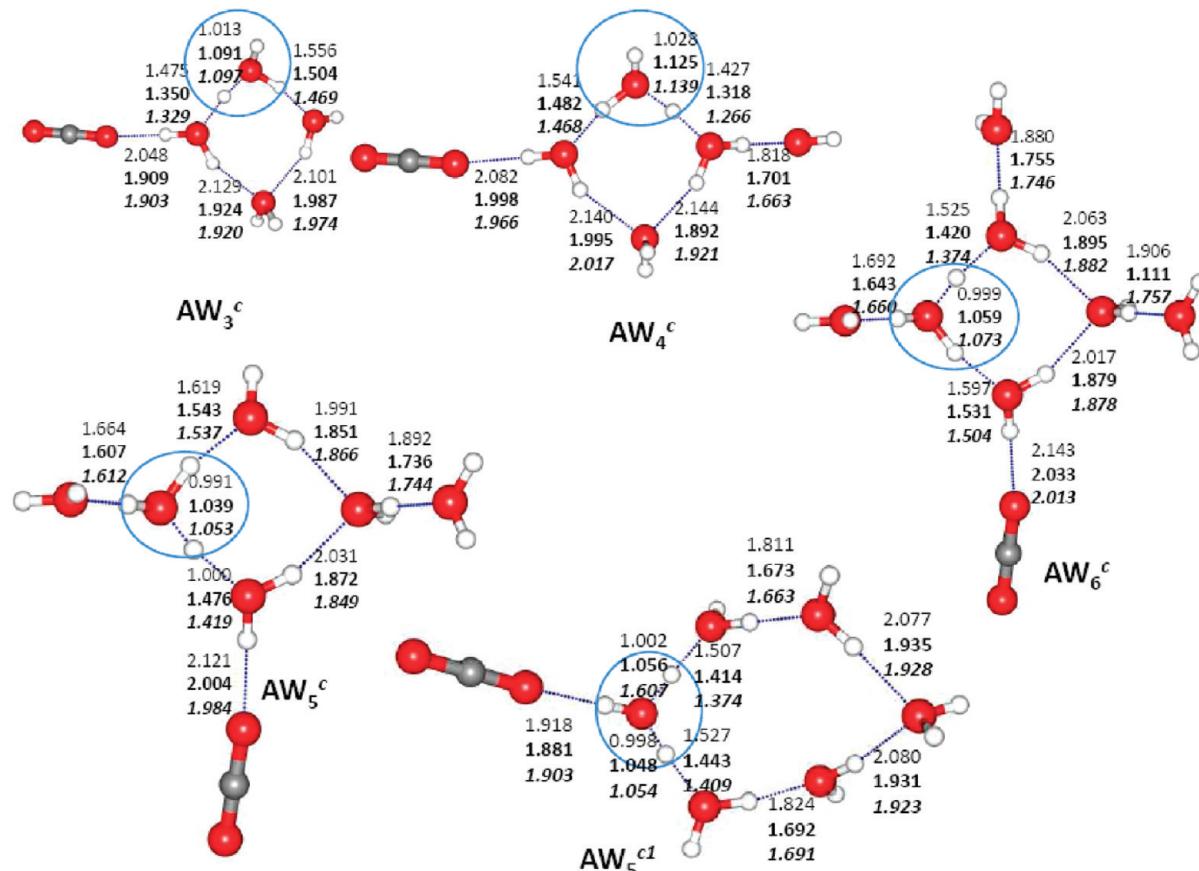
**3.2. Hydration of Isomer B.** **3.2.1. Monohydroxy Hydrated Clusters ( $m$ ).** Figure 4 displays the optimized geometries of isomer B and  $\text{BW}_n^m$  clusters ( $n = 1\text{--}6$ ). The results indicate that the three O—H (0.974 Å) and C—O (1.275 Å) bond lengths and the three C—O—H angles are equal in isomer B. The calculated O—H distance from IR spectroscopy is 0.976 Å.<sup>15</sup> The C—O bond distance in the salt of PCA is 1.231 Å.<sup>19</sup> The results reveal that monohydroxy hydrated clusters of isomer B form linear chainlike structures that are akin to those of  $\text{AW}_n^1$ . Thus, proton-transfer patterns in these clusters are similar to those of  $\text{AW}_n^1$ . In  $\text{BW}_1^m$ , the O1···O4 H-bond distance is 2.430 Å. There is a slight decrease in the length of the C1—O1 bond due to H-bonding. As a result, there is a significant increase in the O1—H1 bond length from 0.974 to 1.093 Å. There are no appreciable changes in the free O—H bond lengths upon formation of a cluster.

In  $\text{BW}_2^m$ , the cluster has two H-bonds: O1···O4 (2.457 Å) and O4···O5 (2.511 Å). It is interesting to note that the O1—

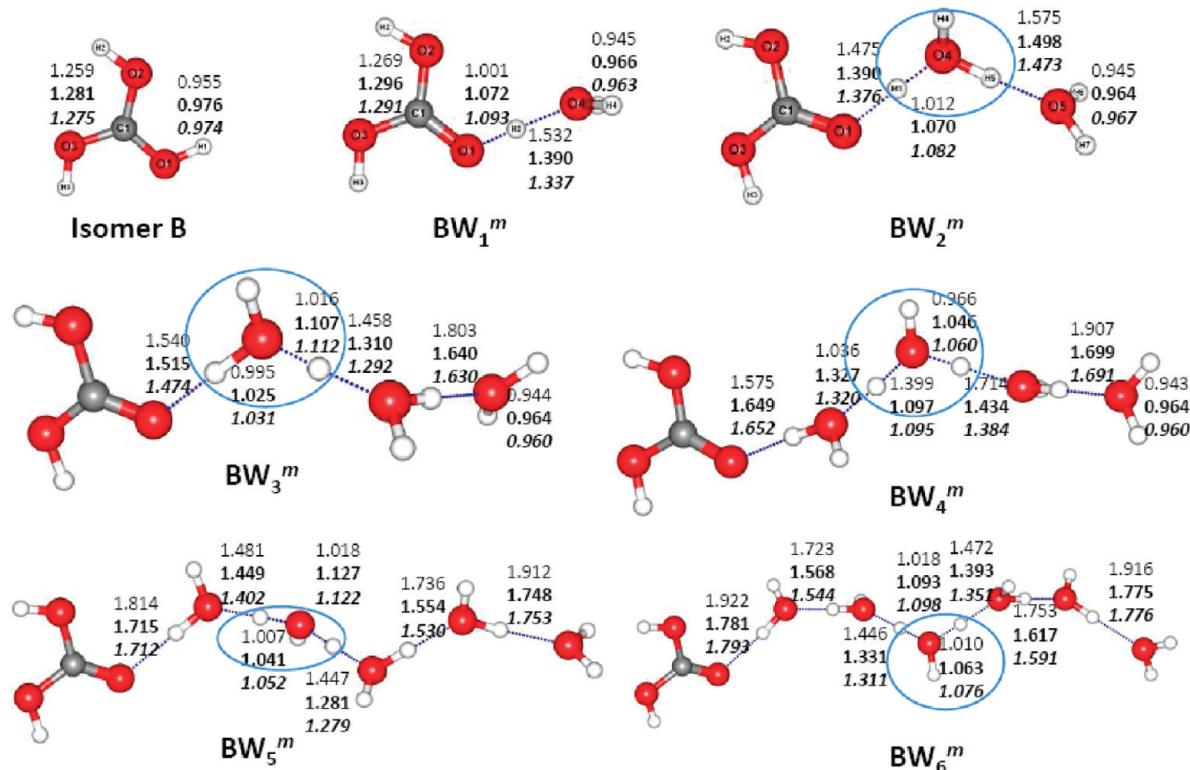
H1 distance increases by 0.283 Å due to the addition of one water molecule to  $\text{BW}_1^m$  and simultaneously the H1—O4 distance decreases by 0.255 Å. This evidence shows that the proton moves slightly away from the CA moiety to the water chain. As the chain length grows from two to six water molecules, the proton shifts from one water molecule to its neighbor in the chain. It can be observed from the geometrical parameters that the distances are shorter for the inner H-bonds than for those present in the periphery. In all of the clusters, most of the H-bond angles are nearly equal to 175°. The presence of the Eigen core in the proton transfer can be seen from the geometrical parameters of all of the clusters.

**3.2.2. Dihydroxy Hydrated Clusters ( $d$ ).** The dihydroxy hydrated clusters of isomer B are depicted in Figure 5. The presence of the Eigen core is clearly evident from the geometrical parameters displayed therein. In  $\text{BW}_2^d$ , two water molecules are H-bonded to two C—O groups, such as C1—O1 and C1—O2. There are noticeable changes in the H-bond parameters of  $\text{BW}_2^d$  when compared to those of  $\text{BW}_1^m$ . The O1···O4 and O2···O5 H-bond distances in  $\text{BW}_2^d$  are 2.472 and 2.482 Å, respectively, which are higher than that in  $\text{BW}_1^m$ . The geometrical parameters indicate that the proton is equally shared between the two water molecules. A similar trend is seen in other clusters with equal numbers of water molecules on both sides of isomer B. Proton transfer is not observed in these cases.

With a view towards probing proton transport, unequally hydrated clusters of B were considered. These clusters are designated as  $\text{BW}_n^{d(i+j)}$  ( $n = 2\text{--}6$ ,  $i + j = n$ ). In these clusters, one of the O—H groups is H-bonded to  $i$  water molecules and the other O—H group is H-bonded to  $j$  water molecules. The calculated geometrical parameters provide evidence for proton transfer in clusters  $\text{BW}_4^{d1}$ ,  $\text{BW}_5^{d1}$ ,  $\text{BW}_5^{d2}$ ,  $\text{BW}_6^{d1}$ ,  $\text{BW}_6^{d2}$ ,  $\text{BW}_6^{d3}$ ,



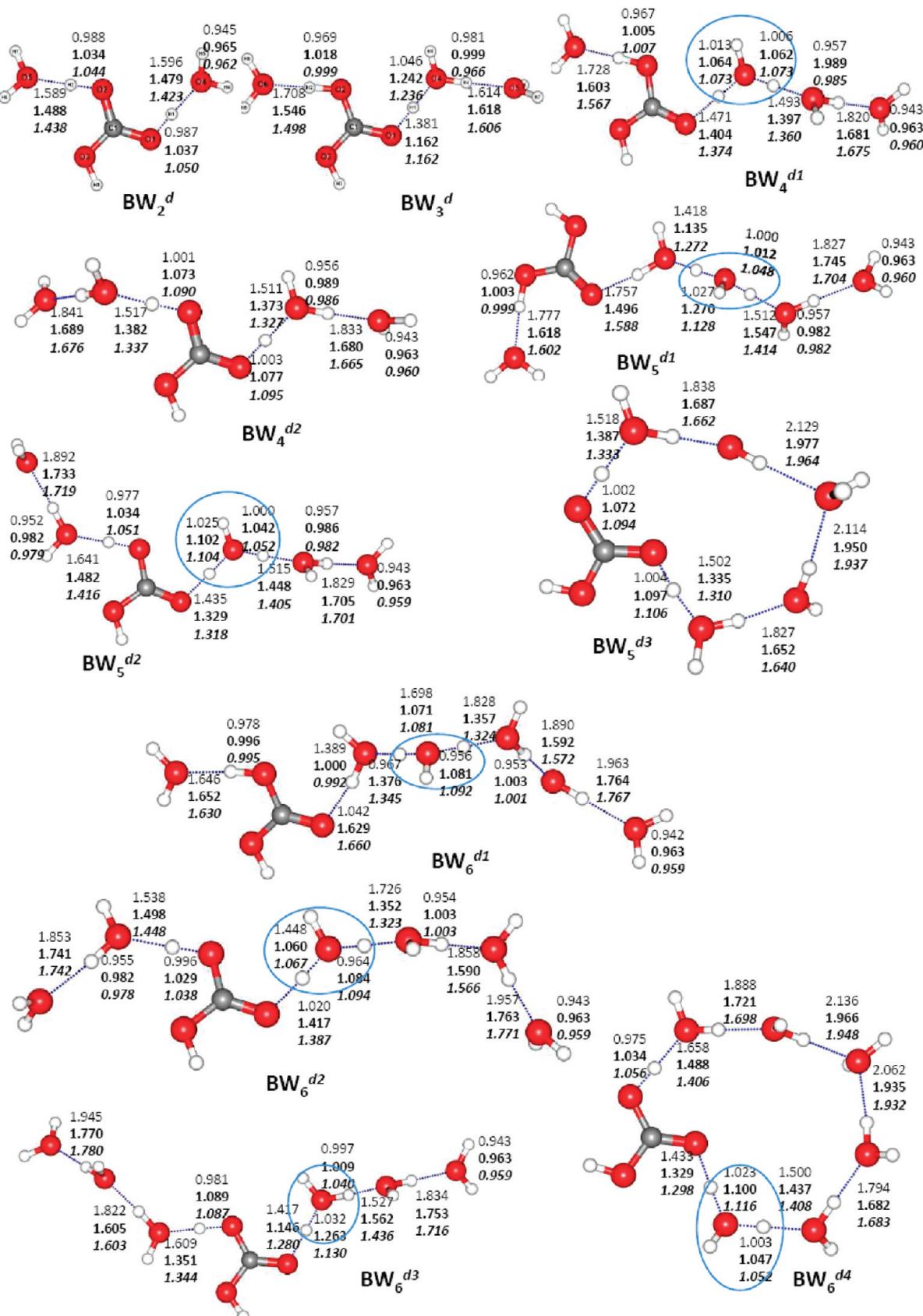
**Figure 3.** Optimized geometries of AW<sub>n</sub><sup>c</sup> cyclic clusters (where n = 3–6) at the HF/6-311++G\*\* (normal), B3LYP/6-311++G\*\* (bold), and M05-2X/6-311++G\*\* (bold and italic) levels of theory. The blue circles indicate the presence of the Eigen core cation. Distances are in angstroms.



**Figure 4.** Optimized geometries of BW<sub>n</sub><sup>m</sup> monohydroxy clusters (where n = 1–6) at the HF/6-311++G\*\* (normal), B3LYP/6-311++G\*\* (bold), and M05-2X/6-311++G\*\* (bold and italic) levels of theory. The blue circles indicate the presence of the Eigen core cation. Distances are in angstroms.

and BW<sub>6</sub><sup>d4</sup>. The unequal environment leads to changes in the charge delocalization arising from the loss of symmetry and

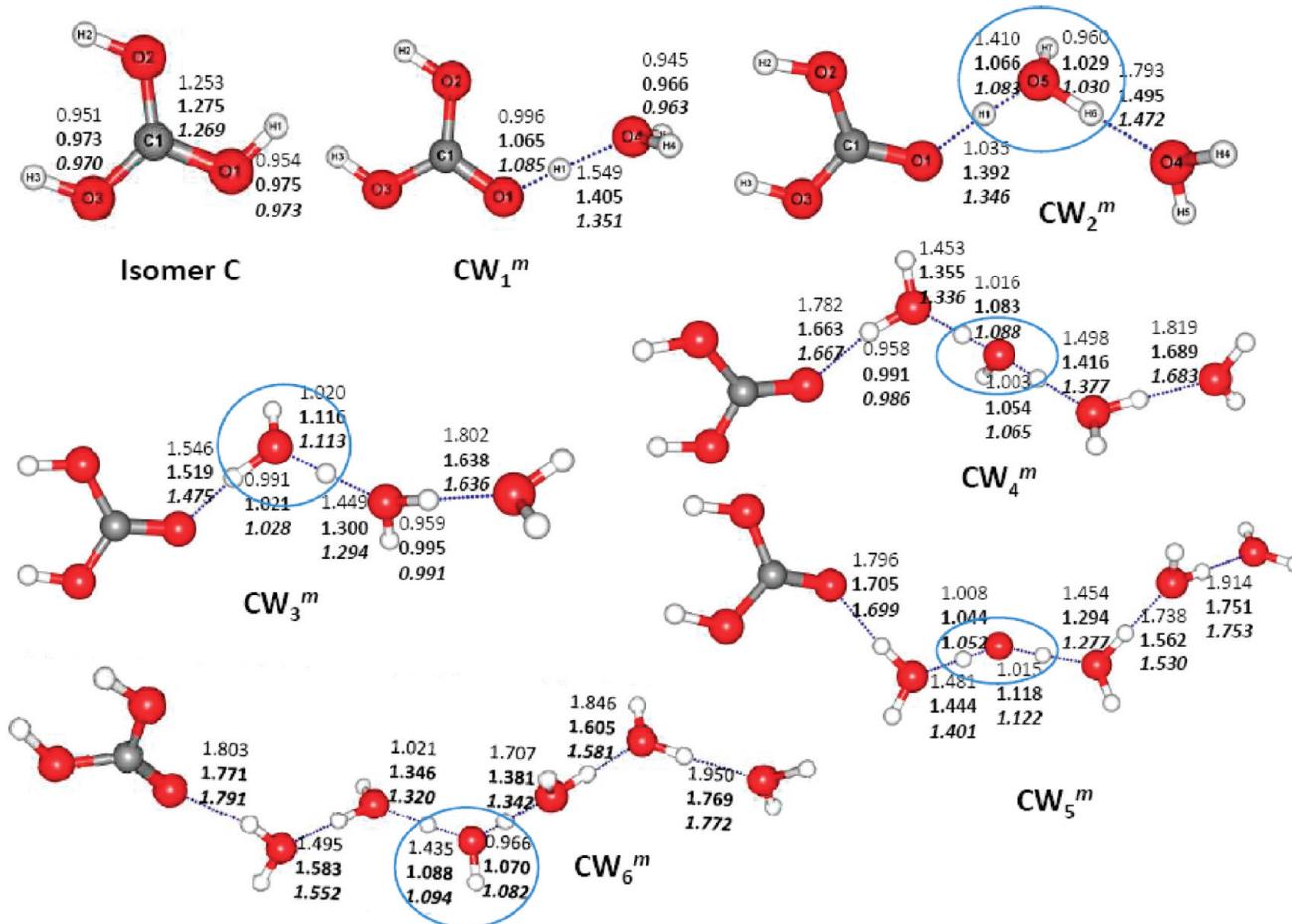
resonance in isomer B. As a result, proton transfer takes place from isomer B to water in the BW<sub>n</sub><sup>di+j</sup>, where i ≠ j type



**Figure 5.** Optimized geometries of  $BW_n^d$  dihydroxy clusters (where  $n = 2$ – $6$ ) at the HF/6-311++G\*\* (normal), B3LYP/6-311++G\*\* (bold), and M05-2X/6-311++G\*\* (bold and italic) levels of theory. The blue circles indicate the presence of the Eigen core cation. Distances are in angstroms.

of clusters. It is observed from Figure 5 that  $n = 5$  and  $6$  clusters form cyclic structures in addition to the linear structures.

**3.3. Hydration of Isomer C. 3.3.1. Monohydroxy Hydrated Clusters ( $m$ ).** The optimized geometries of isomer C and its water clusters ( $CW_n^m$ ,  $n = 1$ – $6$ ) are presented in Figure 6. The



**Figure 6.** Optimized geometries of  $CW_n^m$  monohydroxy clusters (where  $n = 1-6$ ) at the HF/6-311++G\*\* (normal), B3LYP/6-311++G\*\* (bold), and M05-2X/6-311++G\*\* (bold and italic) levels of theory. The blue circles indicate the presence of the Eigen core cation. Distances are in angstroms.

O1–H1 distance of  $CW_1^m$  is longer than that in isomer C, which indicates the movement of proton from isomer C to the water molecule. The movement of the excess proton from isomer C to the neighboring water is evident from the geometrical parameters of  $CW_2^m$ . The subsequent addition of water molecules to  $CW_1^m$  clusters in a linear fashion enables the transfer of the excess proton from one water molecule to other. The proton transfer in water clusters of isomer C is comparable to that in the linear and monohydroxy hydrated clusters of isomers A and B, respectively.

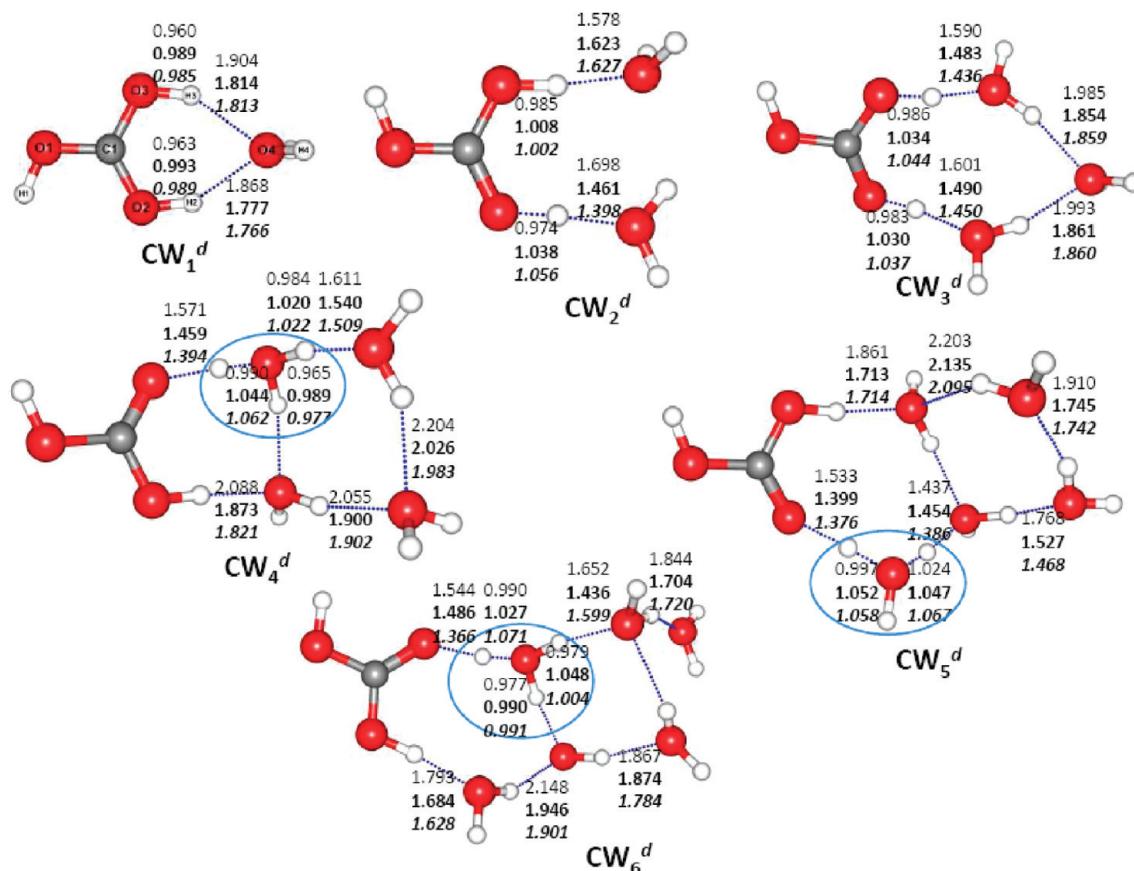
**3.3.2. Dihydroxy Hydrated Clusters (d).** Figure 7 illustrates the geometries of dihydroxy hydrated clusters of isomer C. The existence of cyclic clusters of isomer C (except  $CW_2^d$ ) can be seen from Figure 7. In cluster  $CW_1^d$ , the lone pairs on the oxygen atom interact with isomer C to form a bifurcated H-bond. In  $CW_2^d$ , two water molecules separately bind with the two hydroxyl groups of isomer C. Cyclic motifs such as those found in  $(H_2O)_4$  and  $(H_2O)_5$  are observed in clusters  $CW_3^d$ ,  $CW_4^d$ ,  $CW_5^d$ , and  $CW_6^d$ . There is no significant movement of the excess proton beyond the first hydration shell (i.e., hydration of hydroxyl groups).

**3.3.3. Trihydroxy Hydrated Clusters (t) of Isomers B and C.** Optimized geometries of trihydroxy hydrated clusters of isomers B and C are shown in Figure 8. Only the symmetrically hydrated  $BW_n^t$  and  $CW_n^t$  clusters ( $n = 3$  and 6) are included in this category. The geometrical arrangements of  $BW_3^t$  and  $CW_3^t$  are similar to that of the symmetrically hydrated Eigen cation.  $CW_6^t$  forms a cyclic structure owing to the orientation of the O–H groups in isomer C. It was found from the results that

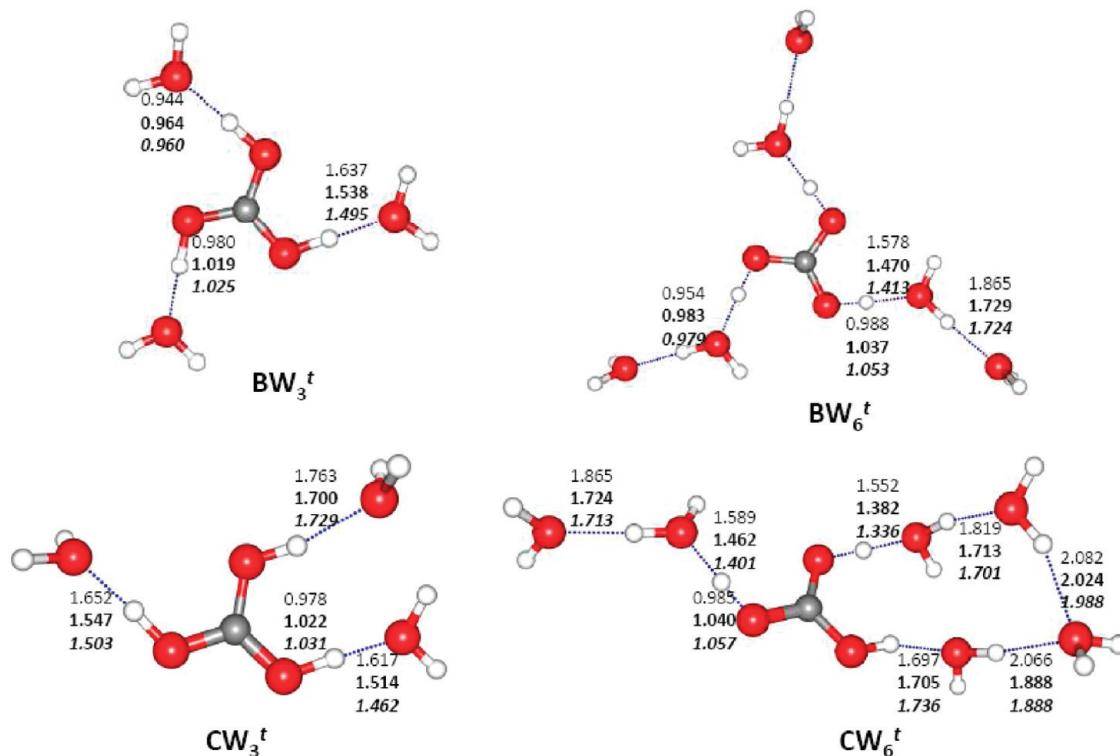
higher-order water clusters of isomers B and C favor the formation of cyclic structures. The calculated geometrical parameters show that there is no significant proton movement from the PCA moiety to the water molecules in these clusters.

**3.4. Comparison of Geometrical Parameters.** In this study, the geometries of various clusters were obtained by three different methods employing three different basis sets. As expected, the geometrical parameters are sensitive to the level of the calculations and the quality of the basis sets. The O2…O3 distance in isomer A at the HF level using the 6-31G\*, 6-31+G\*, and 6-311++G\*\* basis sets was calculated as 2.657, 2.664, and 2.614 Å, respectively. The same distance at the B3LYP level with the above-mentioned basis sets was 2.569, 2.592, and 2.544 Å, respectively. For the same basis sets, the M05-2X method predicts the distance to be 2.579, 2.602, and 2.531 Å, respectively. The calculated O…O distance using the three different basis sets varies as 6-31+G\* > 6-31G\* > 6-311++G\*\* at all calculation levels. The results obtained at different levels of calculation using the 6-311++G\*\* basis set were considered for further analysis. It can be observed from Figure 4 that the O–H distance in isomer B is slightly shorter at the M05-2X level than at the B3LYP level, which is longer than the corresponding value at the HF level. The geometrical parameters of  $CW_n$  clusters are similar to those in clusters with isomer B.

Because the sequential addition of water molecules shifts the proton from one molecule to its neighbor, the effect of calculation level was analyzed for O–H…O distances in various  $PCAW_n^x$  clusters. Comparison of the geometrical



**Figure 7.** Optimized geometries of  $\text{CW}_n^d$  dihydroxy clusters (where  $n = 1-6$ ) at the HF/6-311++G\*\* (normal), B3LYP/6-311++G\*\* (bold), and M05-2X/6-311++G\*\* (bold and italic) levels of theory. The blue circles indicate the presence of the Eigen core cation. Distances are in angstroms.



**Figure 8.** Optimized geometries of  $\text{BW}_n^t$  and  $\text{CW}_n^t$  trihydroxy clusters (where  $n = 3$  and 6) at the HF/6-311++G\*\* (normal), B3LYP/6-311++G\*\* (bold), and M05-2X/6-311++G\*\* (bold and italic) levels of theory. Distances are in angstroms.

parameters indicates that the M05-2X/6-311++G\*\* method yields shorter O-H $\cdots$ O distances than the B3LYP and HF

methods. The general trend in the O-H $\cdots$ O distance obtained using the 6-311++G\*\* basis set is M05-2X < B3LYP < HF.

**TABLE 1:** Sequential Binding Energies (kJ/mol) Calculated at the B3LYP/6-311++G\*\* Level for PCAW<sub>n</sub><sup>x</sup> Clusters<sup>a</sup>, along with Calculated and Experimental Values for Corresponding Clusters with Monovalent Lithium Cation (Li<sup>+</sup>)

n	isomer A			isomer B		isomer C		SBE of Li <sup>+</sup> (H <sub>2</sub> O) <sub>n</sub>	
	linear (l)	branched (b)	cyclic (c)	mono (m)	di (d)	mono (m)	di (d)	calc <sup>b</sup>	expt <sup>c</sup>
1	148	—	—	127	—	122	128	149	142
2	106	104	—	103	100	103	106	127	108
3	71	68	105	71	75	71	90	96	86
4	63	60	77	61	63	61	83	65	69
5	51	49	57	51	51	50	57	59	58
6	48	47	51	48	46	47	64	55	51

<sup>a</sup> PCAW<sub>n</sub><sup>x</sup>, where n = 1–6; x = l, b, c, m, and d. <sup>b</sup> Calculated SBE value taken from ref 42 obtained at the B3LYP/6-311++G\*\* level.

<sup>c</sup> Experimental value taken from ref 50.

However, in clusters BW<sub>2</sub><sup>d</sup>, BW<sub>4</sub><sup>d2</sup>, and BW<sub>6</sub><sup>d3</sup>, there is no similar trend. Evidence from the present calculations reveals that, among the three methods, B3LYP and M05-2X yield consistent geometries.

**3.5. Energetics of PCAW<sub>n</sub><sup>x</sup> Clusters (n = 1–6 and x = l, b, c, m, d, and t).** SBEs of ion solvation are important quantities that can be measured experimentally. The experimental and theoretical values of SBEs of monovalent cation such as Li<sup>+</sup>, Na<sup>+</sup>, and K<sup>+</sup> with water have been reported.<sup>42a</sup> In a previous study, the B3LYP/6-311++G\*\* level of calculation was used to compute the trends in the SBEs of monovalent metal cations.<sup>42a</sup> The same method was selected here to calculate the SBEs of PCAW<sub>n</sub><sup>x</sup> clusters (n = 1–6). The BSSE-corrected SBEs for various clusters are presented in Table 1. Close scrutiny of the results reveals that the addition of the first water molecule to PCA requires more energy than the addition of subsequent water molecules. The incremental SBE gradually decreases with the stepwise addition of water molecules, except in the case of dihydroxy clusters of isomer C (CW<sub>n</sub><sup>d</sup>). This might be due to the changes in the H-bonding pattern in linear and cyclic clusters. It can be seen from Figure 7 that CW<sub>5</sub><sup>d</sup> contains an additional water molecule in the first solvation shell and has the Eigen cation localized at a different position compared to the same in CW<sub>4</sub><sup>d</sup> and CW<sub>6</sub><sup>d</sup>. The general trend observed from this study is similar to that of previous investigation on the solvation of mono- and divalent metal cations.<sup>42</sup> The calculated ranges of SBEs for isomers A, B, and C are 47–148, 46–127, and 47–128 kJ/mol, respectively. The corresponding energies for the monovalent cations obtained from various experimental<sup>50</sup> and theoretical<sup>42</sup> studies are in the ranges 51–142 and 55–149 kJ/mol, respectively. It is interesting to note that the calculated ranges of SBEs of isomers A–C are similar to the Li<sup>+</sup>.

The BSSE-corrected BEs of different PCAW<sub>n</sub><sup>x</sup> clusters of isomers A–C are listed in Tables 2–4. The BSSE-corrected BEs of PCAW<sub>n</sub><sup>x</sup> clusters calculated using the HF method with the 6-31+G\* and 6-311++G\*\* are provided in the Supporting Information (Table S1). It can be seen that the BE increases with increasing degree of solvation. However, the incremental BE decreases increasing number of water molecules in the cluster. As an example, the incremental BE of AW<sub>2</sub><sup>l</sup> with respect to AW<sub>1</sub><sup>l</sup> is 34.6 kcal/mol at the M05-2X/6-311++G\*\* level of calculations, whereas the same for AW<sub>3</sub><sup>l</sup> with respect to AW<sub>2</sub><sup>l</sup> is 19.1 kcal/mol. A similar trend has also been observed in the hydration of Eigen and Zundel cations.<sup>51</sup>

It can be seen from the optimized geometries that isomer A prefers to form linear, branched, and cyclic clusters. Among these clusters, BE values of linear protonated wire-type clusters of isomer A are higher than those for the branched and cyclic types. Both B and C isomers form mono-, di-, and trihydroxy clusters. Comparison of BEs of these clusters shows that dihydroxy clusters are more stable than the other types. Even

**TABLE 2: Binding Energies (kcal/mol) of Isomer A Clusters (AW<sub>n</sub><sup>x</sup>, where n = 1–6; x = l, b, and c) Calculated by the B3LYP and M05-2X Methods Using the 6-31+G\* and 6-311++G\*\* Basis Sets**

	B3LYP		M05-2X	
	6-31+G*	6-311++G**	6-31+G*	6-311++G**
AW <sub>1</sub> <sup>l</sup>	34.6	<b>35.4</b>	38.8	<b>39.5<sup>a</sup></b>
AW <sub>2</sub> <sup>l</sup>	68.8	70.5	71.8	74.2
AW <sub>3</sub> <sup>l</sup>	88.0	88.3	91.3	93.4
AW <sub>4</sub> <sup>l</sup>	100.0	101.1	105.0	107.5
AW <sub>5</sub> <sup>l</sup>	112.9	114.8	118.3	121.4
AW <sub>6</sub> <sup>l</sup>	124.4	125.5	130.5	133.8
AW <sub>2</sub> <sup>b</sup>	54.6	55.5	56.4	58.6
AW <sub>3</sub> <sup>b</sup>	73.6	74.5	79.2	80.7
AW <sub>4</sub> <sup>b</sup>	86.5	87.4	91.1	93.4
AW <sub>5</sub> <sup>b</sup>	100.4	100.8	106.7	108.3
AW <sub>6</sub> <sup>b</sup>	110.6	111.5	117.1	119.7
AW <sub>3</sub> <sup>c</sup>	85.9	85.7	91.1	90.6
AW <sub>4</sub> <sup>c</sup>	98.7	98.9	108.6	110.0
AW <sub>5</sub> <sup>c</sup>	112.4	112.2	117.0	122.2
AW <sub>5</sub> <sup>cl</sup>	99.8	99.9	106.1	107.0
AW <sub>6</sub> <sup>c</sup>	118.6	124.6	129.8	131.3

<sup>a</sup> BEs corresponding to the most stable clusters among isomers A–C given in bold.

though cyclic clusters are observed by isomer C, a linear hydration pattern is more favorable than the other types for this isomer. A comparison of the BEs of corresponding water clusters for all isomers reveals that AW<sub>1</sub><sup>l</sup>, CW<sub>2</sub><sup>m</sup>, BW<sub>3</sub><sup>d</sup>, BW<sub>4</sub><sup>d1</sup>, BW<sub>5</sub><sup>d2</sup>, and BW<sub>6</sub><sup>d3</sup> are more stable than the other clusters with same numbers of water molecules.

It can be noted that the basis sets have a marginal effect on the calculated BEs of various hydrated clusters. However, the inclusion of electron correlation has a considerable effect on the calculated BEs. BEs of various clusters obtained from M05-2X calculations are higher than the corresponding B3LYP values. This might be due to the improved treatment of intermolecular interactions at the M05-2X level of theory. The importance of the correlation energy in the BEs was assessed by computing the differences in BEs obtained from HF calculations compared to B3LYP and M05-2X calculations. The difference in the BEs was taken as the electron correlation contribution. The results are presented in Table S2 of the Supporting Information. The percentage of electron correlation energy contributed to the BE at the M05-2X level of calculation for various clusters ranges from 15% to 50%. These results confirm that electrostatic interactions contribute significantly to the stability of various clusters in addition to the dispersive interaction. Table S2 (Supporting Information) shows that the percentage electron correlation contribution to the BE is higher for the M05-2X method than for the B3LYP method. As explained by Truhlar and co-workers, the parameters used in

**TABLE 3: Binding Energies (kcal/mol) of Isomer B Clusters ( $BW_n^x$ , where  $n = 1\text{--}6$ ,  $x = \text{m}$ ,  $\text{d}$ , and  $\text{t}$ ) Calculated by the B3LYP and M05-2X Methods Using the 6-31+G\* and 6-311++G\*\* Basis Sets**

	B3LYP		M05-2X	
	6-31+G*	6-311++G**	6-31+G*	6-311++G**
$BW_1^{\text{m}}$	29.3	30.4	33.0	34.3
$BW_2^{\text{m}}$	69.6	71.5	72.9	75.3
$BW_3^{\text{m}}$	88.8	89.3	93.6	93.6
$BW_4^{\text{m}}$	100.9	101.7	105.2	106.7
$BW_5^{\text{m}}$	119.3	116.4	121.9	122.0
$BW_6^{\text{m}}$	124.0	124.9	130.5	132.9
$BW_2^{\text{d}}$	49.2	50.4	52.3	55.2
$BW_3^{\text{d}}$	86.7	<b>110.9</b>	103.7	<b>115.0<sup>a</sup></b>
$BW_4^{\text{d1}}$	102.9	<b>104.9</b>	107.9	<b>111.5</b>
$BW_4^{\text{d2}}$	83.6	84.7	91.6	93.8
$BW_5^{\text{d1}}$	121.8	121.3	123.4	126.4
$BW_5^{\text{d2}}$	117.2	<b>120.9</b>	124.5	<b>129.4</b>
$BW_5^{\text{d3}}$	97.3	97.6	106.7	107.9
$BW_6^{\text{d1}}$	126.1	127.7	133.3	136.3
$BW_6^{\text{d2}}$	130.4	131.8	135.7	139.8
$BW_6^{\text{d3}}$	130.9	<b>132.3</b>	142.1	<b>146.9</b>
$BW_6^{\text{d4}}$	128.8	131.5	126.4	143.1
$BW_3^{\text{t}}$	66.4	67.4	70.1	72.8
$BW_6^{\text{t}}$	107.8	108.3	117.2	118.8

<sup>a</sup> BEs corresponding to the most stable clusters among isomers A–C given in bold.

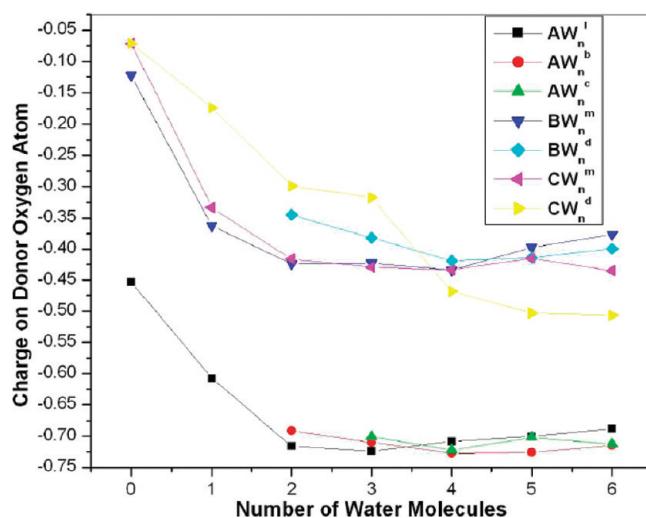
**TABLE 4: Binding Energies (kcal/mol) of Isomer C Clusters ( $CW_n^x$ , where  $n = 1\text{--}6$ ,  $x = \text{m}$ ,  $\text{d}$ , and  $\text{t}$ ) Calculated by the B3LYP and M05-2X Methods Using the 6-31+G\* and 6-311++G\*\* Basis Sets**

	B3LYP		M05-2X	
	6-31+G*	6-311++G**	6-31+G*	6-311++G**
$CW_1^{\text{m}}$	28.4	29.1	31.6	32.8
$CW_2^{\text{m}}$	71.6	<b>73.2</b>	75.3	<b>77.1<sup>a</sup></b>
$CW_3^{\text{m}}$	90.8	91.3	95.3	94.9
$CW_4^{\text{m}}$	101.4	102.6	106.0	108.2
$CW_5^{\text{m}}$	116.6	116.8	121.9	123.4
$CW_6^{\text{m}}$	125.2	126.2	131.1	133.5
$CW_1^{\text{d}}$	31.1	30.7	32.8	32.5
$CW_2^{\text{d}}$	51.5	51.4	55.5	56.2
$CW_3^{\text{d}}$	68.7	69.0	73.5	75.6
$CW_4^{\text{d}}$	95.4	95.4	100.5	101.6
$CW_5^{\text{d}}$	109.4	109.7	120.2	119.0
$CW_6^{\text{d}}$	121.5	121.8	129.4	130.9
$CW_3^{\text{t}}$	69.0	68.8	72.9	74.2
$CW_6^{\text{t}}$	113.1	112.5	123.9	123.9

<sup>a</sup> BEs corresponding to the most stable clusters among isomers A–C given in bold.

the M05-2X method are generated by the simultaneous optimization of exchange and correlation functionals with the kinetic energy density included in both of them, so that an improved treatment of medium-range correlation effects is obtained by this exchange-correlation functional.<sup>43,44</sup>

Some of the problems encountered in deriving thermochemical parameters for hydrated clusters using static quantum chemical calculations include the floppiness of the molecules involved and the existence of many isomers with small energy differences. However, static quantum chemical calculations can provide useful trends. The thermochemical parameters were calculated for the structures obtained from geometry optimization at the B3LYP/6-31+G\* level using the methodology



**Figure 9.** Correlation of the number of water molecules in PCAW<sub>n</sub><sup>x</sup> (isomers A, B, and C) versus total Mullikan charge on the proton-donating oxygen atoms in the clusters at the M05-2X/6-311++G\*\* level.

implemented in the Gaussian package. The calculated thermodynamic parameters at 298 K for various hydrated clusters are presented in Table S3 of the Supporting Information. The calculated free energy values reveal that hydration of PCA is favorable in general and, in particular, dihydroxy hydration of isomer C is more favorable than hydration of isomers B and A.

**3.6. Charge Analysis.** Charges on the proton-donating oxygen atom of isomers A–C and their hydrated clusters were obtained from Mulliken population analysis (MPA). The variation of the charge with the number of water molecules is plotted in Figure 9. The charge on the proton-donor oxygen decreases as a result of proton transfer. For linear clusters ( $AW_n^l$ ), the charge on the oxygen atom decreases dramatically up to  $n = 3$ , and thereafter, the variation in the charge is constant. These results show that proton movement is localized in the middle of the linear chain, which is in accordance with the findings of an earlier report.<sup>33b</sup> A similar trend can be seen from the plot for monohydroxy clusters of isomers B and C.

**3.7. Vibrational Frequencies.** Recent developments in spectroscopic measurements and theoretical calculations have successfully unveiled structures of protonated water clusters.<sup>21,52–58</sup> Scaled vibrational frequencies of all of the hydrated clusters studied here, calculated at the B3LYP/6-31+G\* level of theory using the harmonic approximation, are listed in Tables 5–7. Both experimental and theoretical vibrational frequencies of protonated water clusters from previous studies<sup>21c,55</sup> are included as reference for analysis of the values for PCAW<sub>n</sub><sup>x</sup> clusters.

It was found in previous experimental studies that the asymmetric ( $E_{\text{as}}$ ) and symmetric ( $E_{\text{ss}}$ ) stretching frequencies of bare H<sub>3</sub>O<sup>+</sup> are 3530 and 3390 cm<sup>-1</sup>, respectively.<sup>59</sup> The corresponding values calculated in the present investigation using the B3LYP/6-31+G\* method are 3508 and 3414 cm<sup>-1</sup>, which are in good agreement with the above-mentioned experimental values. In protonated water clusters, these  $E_{\text{as}}$  and  $E_{\text{ss}}$  modes occur at 2665 and 2420 cm<sup>-1</sup>, respectively, because of H-bonding.<sup>21c,55</sup> In addition, H–O–H bending modes of the Eigen core (major component) and dangling water molecules (minor component) are found at 1900 and 1760 cm<sup>-1</sup>.<sup>21c,55</sup> In the same study, the H–O–H intramolecular bend of dangling water molecules was observed at 1620 cm<sup>-1</sup>. Spectral features

**TABLE 5: Calculated OH-Stretching Frequencies of the Eigen (E) Core Present in  $\text{AW}_n^x$  Clusters<sup>a</sup> at the B3LYP/6-31+G\* Level along with Experimental Values from Protonated Water Clusters**

$\text{AW}_n^x$	description of OH stretching <sup>b,c</sup>	$\nu_{\text{calc}}$ (cm <sup>-1</sup> )	$\nu_{\text{expt}}^d$ (cm <sup>-1</sup> )
isomer A	$\text{E}_{\text{as}}$	3508	3530 <sup>e</sup>
	$\text{E}_{\text{ss}}$	3414	3390 <sup>f</sup>
	$\text{E}_{\text{as}}$	3568	3566 <sup>g</sup>
	$\text{E}_{\text{ss}}$	3497	3545 <sup>g</sup>
$\text{AW}_1^1$	H-bonded O—H stretch to $\text{CO}_2$	2863	2869 <sup>g</sup>
	H-bonded O—H stretch to $\text{CO}_2$	3320	
	E-core free —OH stretch	3615	3580
	H—O—H bend of E core and dangling $\text{H}_2\text{O}$	1958	1900, 1760
$\text{AW}_2^1$	E-core stretching	2066	~2100
	$\text{E}_{\text{as}}$	2616	2665
$\text{AW}_3^1$	$\text{E}_{\text{as}}$ stretch to AD-type $\text{H}_2\text{O}^e$	2660	2665
	O—H <sup>+</sup> —O stretching mode	1302	1317 <sup>h</sup> , 1337 <sup>i</sup>
$\text{AW}_4^1$	H—O—H bend of E core and dangling $\text{H}_2\text{O}$	1908	1900
	$\text{E}_{\text{ss}}$ stretch to AD-type $\text{H}_2\text{O}^e$	2228	2420
	E-core free —OH stretch	3652	3580
	O—H <sup>+</sup> —O bending mode	1465	~1510 <sup>j</sup>
$\text{AW}_5^1$	H—O—H bend of E core and dangling $\text{H}_2\text{O}$	1802	1900
	$\text{E}_{\text{ss}}$ stretch to AD-type $\text{H}_2\text{O}$	2261	2420
	H—O—H bend of E core and dangling $\text{H}_2\text{O}$	1821	1900
	$\text{E}_{\text{ss}}$ stretch to AD-type $\text{H}_2\text{O}$	2078	2420
$\text{AW}_2^b$	$\text{E}_{\text{as}}$	2525	2665
	E-core stretch to A-type $\text{H}_2\text{O}$	2715	2830 <sup>anh k</sup>
$\text{AW}_3^b$	H-bonded O—H stretch with $\text{CO}_2$	3446	
	H—O—H bend of E core and dangling $\text{H}_2\text{O}$	1891	1900
$\text{AW}_4^b$	E-core HB stretch to A-type $\text{H}_2\text{O}$	2913	2889 <sup>anh</sup>
	AD-type $\text{H}_2\text{O}$ H-bond stretch	3200	3195
$\text{AW}_5^b$	$\text{E}_{\text{as}}$ stretch to AD-type $\text{H}_2\text{O}$	2202	2665
	$\text{E}_{\text{ss}}$ stretch to AD-type $\text{H}_2\text{O}$	2364	2420
$\text{AW}_6^b$	H—O—H bend of E core and dangling $\text{H}_2\text{O}$	1826	1900
	$\text{E}_{\text{as}}$ stretch to AD-type $\text{H}_2\text{O}$	2600	2665
$\text{AW}_3^c$	E-core stretch to AD-type $\text{H}_2\text{O}$	1996	1885
	$\text{E}_{\text{ss}}$ to AD-type $\text{H}_2\text{O}$	2282	2420
$\text{AW}_4^c$	H—O—H bend of E core and dangling $\text{H}_2\text{O}$	1852	1900
	$\text{E}_{\text{as}}$ stretch to AD-type $\text{H}_2\text{O}$	2605	2665
$\text{AW}_5^c$	O—H <sup>+</sup> —O bending mode	1470	~1510 <sup>j</sup>
	E-core HB stretch to A-type $\text{H}_2\text{O}$	3030	2889
$\text{AW}_6^c$	$\text{E}_{\text{as}}$ stretch to ADD-type $\text{H}_2\text{O}$	2730	2665
	$\text{E}_{\text{ss}}$ to AD-type $\text{H}_2\text{O}$	2388	2420
$\text{AW}_5^{cl}$	E-core stretching	2111	2100
	$\text{E}_{\text{ss}}$	2360	2420
$\text{AW}_6^c$	E-core dangling $\text{H}_2\text{O}$	2225	2100
	$\text{E}_{\text{as}}$ stretch to AD-type $\text{H}_2\text{O}$	2613	2665

<sup>a</sup>  $\text{AW}_n^x$  clusters, where  $n = 1–6$ ,  $x = 1$ , b, and c, use harmonic calculations (B3LYP with 6-31+G\* basis set) and are scaled by 0.973. <sup>b</sup> E, Eigen cation; PCA, protonated carbonic acid; CA, carbonic acid;  $\text{E}_{\text{as}}$ , Eigen asymmetric stretch;  $\text{E}_{\text{ss}}$ , Eigen symmetric stretch. <sup>c</sup> A-type  $\text{H}_2\text{O}$  molecules accept a hydrogen bond from the Eigen core. AD-type  $\text{H}_2\text{O}$  molecules accept and donate a hydrogen-bond. ADD-type  $\text{H}_2\text{O}$  molecules accept one and donate two hydrogen bonds (HBs). <sup>d</sup> Unless otherwise indicated, experimental O—H stretching frequencies of protonated water clusters taken from refs 21c and 55. <sup>e</sup> Taken from ref 59a. <sup>f</sup> Taken from ref 59b. <sup>g</sup> Experimental O—H stretching frequency value taken from ref 15. <sup>h</sup> Taken from ref 53. <sup>i</sup> Taken from ref 54. <sup>j</sup> Taken from ref 57. <sup>k</sup> anh, anharmonic frequencies.

of Zundel cations reported in protonated water cluster appear in the range of 600–1900 cm<sup>-1</sup>. Asmis et al. reported an O—H<sup>+</sup>—O stretching mode at 1317 cm<sup>-1</sup> in protonated water

clusters.<sup>53</sup> Fridgen and co-workers observed the same band at 1337 cm<sup>-1</sup>.<sup>54</sup> Recently, very weak O—H<sup>+</sup>—O bending frequencies (~1510 and ~1370 cm<sup>-1</sup>) were reported.<sup>57</sup> The same study provided evidence for the existence of an (O—H<sup>+</sup>—O)<sub>as</sub> stretching frequency at ~1000 cm<sup>-1</sup>.<sup>57</sup> This spectral information was used as the basis for analysis of the vibrational spectral frequencies of PCAW<sub>n</sub><sup>x</sup> clusters.

Tables 5–7 reveal the presence of vibrational frequencies in the ranges of 2202–2730 and 2078–2566 cm<sup>-1</sup> in PCAW<sub>n</sub><sup>x</sup> clusters. A comparison of these values with those of protonated water clusters reveals that these signals can be assigned to the  $\text{E}_{\text{as}}$  and  $\text{E}_{\text{ss}}$  modes of the Eigen cation. Analysis of the results indicates the presence of the Eigen core in all clusters except in the case of trihydroxy hydrated clusters. The calculated asymmetric ( $\text{E}_{\text{as}}$ ) and symmetric ( $\text{E}_{\text{ss}}$ ) stretching frequencies of free O—H of the Eigen core in isomer A are 3568 and 3497 cm<sup>-1</sup>, respectively. The corresponding experimental values for the same obtained from IR photodissociation (IRPD) studies<sup>15</sup> on PCA are 3566 and 3545 cm<sup>-1</sup>. The calculated stretching frequency of the O—H group that is H-bonded to the  $\text{CO}_2$  fragment of isomer A is 2863 cm<sup>-1</sup>, which is in close agreement with the experimental value of 2869 cm<sup>-1</sup>.<sup>15</sup> In  $\text{AW}_1^1$ , the free O—H stretching frequency of the Eigen core occurs at 3615 cm<sup>-1</sup>. In the same cluster, a value of 1958 cm<sup>-1</sup> has been observed. This might be due to the H—O—H bend of the Eigen core and dangling water molecules, in accordance with a previous report.<sup>21c,55</sup>

In  $\text{AW}_2^1$ , the O—H stretching frequencies of the Eigen core occur at 2066 (2100) and 2616 (2665) cm<sup>-1</sup>. The calculated  $\text{E}_{\text{as}}$  frequency of the  $\text{AW}_3^1$  cluster is 2660 cm<sup>-1</sup>. These values are in close agreement with the respective experimental frequencies of protonated water clusters. The calculated O—H<sup>+</sup>—O stretching mode in the  $\text{AW}_3^1$  cluster is observed at 1302 cm<sup>-1</sup>. The corresponding experimental value for protonated water clusters varies from 1317 to 1337 cm<sup>-1</sup>. It is interesting to see that all of the clusters with isomer A have spectral signatures of the Eigen core. As mentioned, the spectral features in the lower-energy region below 1900 cm<sup>-1</sup> might arise as a result of Eigen-to-Zundel and Zundel-to-Eigen transitions during proton transport. Although geometrical parameters obtained form the present static quantum chemical calculations do not provide any signatures of Zundel ions, the calculated vibrational frequencies in the lower-energy region yield some clues about the involvement of Zundel ions during proton transfer.

It is well-known that the vibrational frequencies of individual molecules undergo a considerable red shift upon formation of H-bonded complexes, and the same type of red shift has been used to characterize the H-bonding interactions in protonated water clusters. Analysis of all of the values listed in the tables reveals larger red shifts (~790–2600 cm<sup>-1</sup>) and higher intensity values (~2000–6000 arbitrary units) in all of the clusters. Similar observations (very large red shifts of ~1100–1630 cm<sup>-1</sup> and intensities of ~3000–5000 arbitrary units) have been made in earlier experimental<sup>21</sup> and theoretical studies on hydrated Eigen or Zundel ions.<sup>52–58</sup>

The experimental and calculated O—H stretching frequencies of isomer B are 3520 and 3511 cm<sup>-1</sup>, respectively.<sup>15</sup> It may be noticed from Table 6 that the O—H stretching frequency of  $\text{BW}_1^{\text{m}}$ , which is H-bonded to a water molecule, is 2022 cm<sup>-1</sup> and the corresponding red shift is 1489 cm<sup>-1</sup>. The stretching of free O—H occurs at 3545 cm<sup>-1</sup>. The O—H<sup>+</sup>—O<sub>as</sub> stretching modes of clusters  $\text{BW}_5^{\text{m}}$  and  $\text{BW}_3^{\text{d}}$  occur at 869 and 860 cm<sup>-1</sup>, respectively. This is due to the

**TABLE 6: Calculated OH-Stretching Frequencies of the Eigen (E) Core Present in  $BW_n^x$  Clusters<sup>a</sup> at the B3LYP/6-31+G\* Level along with Experimental Values from Protonated Water Clusters**

$BW_n^x$	description of O–H stretching <sup>b</sup>	$\nu_{\text{calc}}$ (cm <sup>-1</sup> )	$\nu_{\text{expt}}^c$ (cm <sup>-1</sup> )
isomer B	O–H <sub>ss</sub>	3511	3520 <sup>d</sup>
$BW_1^m$	PCA O–H stretch to A-type H <sub>2</sub> O (corresponding red shift, 1489 cm <sup>-1</sup> )	2022	
	free O–H stretching in PCA	3545	
$BW_2^m$	E-core stretching	2148	~2100
	E <sub>as</sub>	2636	2665
	E-core free –OH stretch	3629	3580
$BW_3^m$	O–H <sup>+</sup> –O stretching mode	1300	1317, <sup>e</sup> 1337 <sup>f</sup>
	H–O–H bend	1723	1756, <sup>e</sup> 1768
	E-core H-bond stretch to CA (CA…H <sub>3</sub> O <sup>+</sup> )	2747	2730 <sup>h</sup>
$BW_4^m$	O–H <sup>+</sup> –O stretching mode	1417	1317, <sup>e</sup> 1337 <sup>f</sup>
	E <sub>ss</sub> to AD-type H <sub>2</sub> O	2414	2420
$BW_5^m$	O–H <sup>+</sup> –O <sub>as</sub> stretching mode	869	921, <sup>e</sup> 990, <sup>f</sup> 1000 <sup>g</sup>
	E <sub>ss</sub> to AD-type H <sub>2</sub> O	2566	2420
$BW_6^m$	O–H <sup>+</sup> –O stretching mode	1441	1317, <sup>e</sup> 1337 <sup>f</sup>
	E <sub>ss</sub> to AD-type H <sub>2</sub> O	2160	2420
$BW_2^g$	CA O–H stretch to A-type H <sub>2</sub> O (corresponding red shift, 1055 cm <sup>-1</sup> )	2456	
$BW_3^g$	O–H <sup>+</sup> –O <sub>as</sub> stretching mode	860	921, <sup>e</sup> 990, <sup>f</sup> 1000 <sup>g</sup>
	CA O–H stretch to A-type H <sub>2</sub> O	2839	
$BW_4^{d1}$	H–O–H bend of E core and dangling H <sub>2</sub> O	1994	1900
	E <sub>ss</sub> stretch to AD-type H <sub>2</sub> O	2324	2420
	CA O–H stretch to A-type H <sub>2</sub> O	2974	
$BW_4^{d2}$	PCA O–H stretch to AD-type H <sub>2</sub> O (corresponding red shift, 1589 cm <sup>-1</sup> )	1922	
$BW_5^{d1}$	O–H <sup>+</sup> –O stretching mode	1039	1043, <sup>e</sup> 1000 <sup>g</sup>
	E <sub>ss</sub> stretch to AD-type H <sub>2</sub> O	2548	2420
	CA O–H stretch to A-type H <sub>2</sub> O	3064	
$BW_5^{d2}$	H–O–H bend of E core and dangling H <sub>2</sub> O	1951	1900
	E <sub>ss</sub> stretch to AD-type H <sub>2</sub> O	2318	2420
	CA O–H stretch to AD-type H <sub>2</sub> O	2552	
$BW_5^{d3}$	PCA O–H stretch to AD-type H <sub>2</sub> O	1368 and 2038	
$BW_6^{d1}$	H–O–H bend of E core and dangling H <sub>2</sub> O	1842	1900
	E <sub>ss</sub> stretch to AD-type H <sub>2</sub> O	2100	2420
	CA O–H stretch to A-type H <sub>2</sub> O	3085	
$BW_6^{d2}$	O–H <sup>+</sup> –O stretching mode	1458	1317, <sup>e</sup> 1337 <sup>f</sup>
	E <sub>ss</sub> to AD-type H <sub>2</sub> O	2373	2420
	CA O–H stretch to AD-type H <sub>2</sub> O	2673	
$BW_6^{d3}$	O–H <sup>+</sup> –O stretching mode	1326	1317 <sup>e</sup> , 1337 <sup>f</sup>
	PCA O–H stretch to AD-type H <sub>2</sub> O (corresponding red shift, 1326 cm <sup>-1</sup> )	2185	
$BW_6^{d4}$	PCA O–H stretch to AD-type H <sub>2</sub> O	1490 and 2536	
	H–O–H bend of E core and dangling H <sub>2</sub> O	1960	1900
	E <sub>ss</sub>	2243	2420
$BW_3^t$	PCA O–H stretch to A-type H <sub>2</sub> O (corresponding red shifts, 805 and 801 cm <sup>-1</sup> )	2706 and 2710	
$BW_6^t$	PCA O–H stretch to AD-type H <sub>2</sub> O (corresponding red shifts, 1169 and 1150 cm <sup>-1</sup> )	2342 and 2361	

<sup>a</sup>  $BW_n^x$  clusters, where  $n = 1–6$ ,  $x = m$ , d, and t, use harmonic calculations (B3LYP with 6-31+G\* basis set) and are scaled by 0.973. <sup>b</sup> E, Eigen cation; PCA, protonated carbonic acid; CA, carbonic acid; O–H<sub>ss</sub>, O–H symmetric stretching; E<sub>as</sub>, Eigen asymmetric stretch; E<sub>ss</sub>, Eigen symmetric stretch. <sup>c</sup> Unless otherwise indicated, experimental O–H stretching frequencies of protonated water clusters taken from refs 21c and 55. <sup>d</sup> Experimental O–H stretching frequency value taken from ref 15. <sup>e</sup> Taken from ref 53. <sup>f</sup> Taken from ref 54. <sup>g</sup> Taken from ref 57. <sup>h</sup> Taken from ref 58b.

oscillation of the proton between the two water molecules. The presence of the Eigen core in clusters  $BW_n^m$  ( $n = 2–6$ ) and  $BW_n^d$  ( $n = 4–6$ ) is evident from the spectral information. It is necessary to mention that  $BW_n^d$  clusters with unequal numbers of water molecules have the geometric and spectral signatures of the Eigen core. Vibrational frequencies of lower-energy modes are presented in monohydroxy clusters ( $BW_n^m$ , where  $n = 3–6$ ) and dihydroxy clusters ( $BW_3^d$ ,  $BW_5^{d1}$ ,  $BW_5^{d3}$ ,  $BW_6^{d2}$ ,  $BW_6^{d3}$ , and  $BW_6^{d4}$ ).

The calculated O–H stretching frequencies of  $CW_n^x$  clusters are presented in Table 7. The O–H stretching frequency of isomer C is 3556 cm<sup>-1</sup>. In  $CW_1^m$ , the O–H stretching frequency is 2122 cm<sup>-1</sup>. The corresponding red shift is 1434 cm<sup>-1</sup>. The O–H stretching frequencies and Eigen signatures of  $CW_n^x$  clusters are analogous to those of  $BW_n^x$  clusters. Vibrational frequencies of lower-energy modes are presented in monohydroxy clusters only ( $CW_n^m$ , where  $n = 3, 5$ , and 6). It can be seen that fully hydrated Eigen cores are present in  $AW_5^c$ ,  $AW_5^{cl}$ ,  $AW_6^c$ ,  $CW_4^d$ ,  $CW_5^d$ , and  $CW_6^d$ . These O–H stretching frequen-

**TABLE 7: Calculated OH-Stretching Frequencies of the Eigen (E) Core Present in  $CW_n^x$  Clusters<sup>a</sup> at the B3LYP/6-31+G\* Level along with Experimental Values from Protonated Water Clusters**

$CW_n^x$	description of O–H stretching <sup>b</sup>	$\nu_{\text{calc}}$ ( $\text{cm}^{-1}$ )	$\nu_{\text{expt}}^c$ ( $\text{cm}^{-1}$ )
isomer C	O–H <sub>ss</sub>	3556	
$CW_1^m$	PCA O–H stretch to A-type H <sub>2</sub> O (corresponding red shift, 1434 $\text{cm}^{-1}$ )	2122	
	free O–H stretching in PCA	3618	
$CW_2^m$	E-core stretching	2180	~2100
	E <sub>as</sub>	2615	2665
$CW_3^m$	E-core free –OH stretch	3638	3580
	O–H <sup>+</sup> –O stretching mode	1281	1317, <sup>d</sup> 1337 <sup>e</sup>
	H–O–H bend	1753	1756, <sup>d</sup> 1768
	E-core H-bond stretch to CA (CA•••H <sub>3</sub> O <sup>+</sup> )	2794	2730 <sup>g</sup>
$CW_4^m$	H–O–H bend of E core and dangling H <sub>2</sub> O	1901	1900
	E <sub>ss</sub> to AD-type H <sub>2</sub> O	2271	2420
$CW_5^m$	O–H <sup>+</sup> –O stretching mode	1232	1317, <sup>d</sup> 1337 <sup>e</sup>
	H–O–H bend	1761	1756, <sup>d</sup> 1768
	E <sub>ss</sub> to AD-type H <sub>2</sub> O	2432	2420
$CW_6^m$	O–H <sup>+</sup> –O stretching mode	1508	1317, <sup>d</sup> 1337 <sup>e</sup>
	H–O–H bend of E core and dangling H <sub>2</sub> O	1805	1900
	E-core stretching	2090	~2100
$CW_1^f$	PCA O–H stretch to AA-type H <sub>2</sub> O	3279	
	free O–H stretching in PCA	3563	
$CW_2^d$	PCA O–H stretch to A-type H <sub>2</sub> O (corresponding red shifts, 1135 and 611 $\text{cm}^{-1}$ )	2421 and 2945	
$CW_3^d$	PCA O–H stretch to AD-type H <sub>2</sub> O (corresponding red shift, 935 $\text{cm}^{-1}$ )	2621	
$CW_4^d$	H–O–H bend	1806	1756, <sup>d</sup> 1768
	E <sub>as</sub>	2501	2665
$CW_5^d$	E-core stretch to CA and AD-type H <sub>2</sub> O	2754	
	E-core stretching	2198	~2100
$CW_6^d$	E <sub>ss</sub> to AAD-type H <sub>2</sub> O	2438	2420
	E <sub>ss</sub> to AD-type H <sub>2</sub> O	2308	2420
	E-core stretch to CA and AD-type H <sub>2</sub> O	2715	2730 <sup>g</sup>
	E-core stretch to AAD-type H <sub>2</sub> O	3125	
$CW_3^t$	PCA O–H stretch to A-type H <sub>2</sub> O (corresponding red shifts, 900 and 444 $\text{cm}^{-1}$ )	2656 and 3112	
$CW_6^t$	PCA O–H stretch to AD-type H <sub>2</sub> O (corresponding red shifts, 1670 and 1158 $\text{cm}^{-1}$ )	1886 and 2398	

<sup>a</sup>  $CW_n^x$  clusters, where  $n = 1–6$ ,  $x = m$ ,  $d$ , and  $t$ , use harmonic calculations (B3LYP with 6-31+G\* basis set) and are scaled by 0.973. <sup>b</sup> E, Eigen cation; PCA, protonated carbonic acid; CA, carbonic acid; O–H<sub>ss</sub>, O–H symmetric stretching; E<sub>as</sub>, Eigen asymmetric stretch; E<sub>ss</sub>, Eigen symmetric stretch. <sup>c</sup> Unless otherwise indicated, experimental O–H stretching frequencies of protonated water clusters taken from refs 21c and 55. <sup>d</sup> Taken from ref 53. <sup>e</sup> Taken from ref 54. <sup>f</sup> Taken from ref 57. <sup>g</sup> Taken from ref 58b.

cies of the Eigen core are in good agreement with the available experimental and theoretical values for protonated water clusters.<sup>21,52–58</sup>

#### 4. Summary and Conclusions

In this study, an attempt has been made to probe the sequential hydration of three different isomers of PCA. The geometrical parameters of clusters reveal the occurrence of proton transfer from PCA to water in linear and monohydroxy hydrated clusters. The presence of the Eigen core in proton transfer is evident from the geometries and IR spectral signatures of these clusters. The lower-energy modes illustrate the involvement of Eigen → Zundel → Eigen transitions in proton transfer in PCAW<sub>n</sub><sup>x</sup> clusters. One of the important results emerging from the present study is that the symmetrically hydrated clusters do not favor proton transport. The SBEs of all of the clusters are similar to those for the hydration of the Li<sup>+</sup> ion. Comparison of BEs of all of the clusters containing a single water molecule reveals that the AW<sub>1</sub><sup>l</sup> cluster has the maximum stability. In clusters with two water molecules, monohydroxy hydrated isomer C has the highest stability. For  $n = 3–6$ , the clusters of isomer B, viz., BW<sub>3</sub><sup>d</sup>, BW<sub>4</sub><sup>d1</sup>, BW<sub>5</sub><sup>d2</sup>, and BW<sub>6</sub><sup>d3</sup> have the maximum BEs

when compared to the respective clusters of isomers A and C with the same numbers of water molecules. The calculated BEs indicate that the electrostatic interaction is the predominant interaction in the stabilization of these clusters and the contribution of intermolecular dispersive interaction (correlation energy) at M05-2X/6-311++G\*\* varies from 15% to 50% depending on the intrinsic nature of the clusters.

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**Supporting Information Available:** BEs calculated by the HF method, contribution of electron correlation to the BEs, and thermochemical parameters of PCAW<sub>n</sub><sup>x</sup> clusters are given Tables S1–S3. The starting and final optimized geometries of AW<sub>n</sub> clusters (where  $n = 1$ , 3, and 6) are provided in Figure S1. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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