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Integration Grid Errors for Meta-GGA-Predicted Reaction Energies: Origin of Grid Errors for the M06 Suite of Functionals

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Abstract: We have assessed integration grid errors arising from the use of popular DFT quadrature schemes for a set of 34 organic reaction energies. The focus is primarily on M05-2X and the M06 suite of functionals (M06-L, M06, M06-2X, and M06-HF). M05-2X, M06, and M06-2X outperform popular older DFT functionals for the reaction energies studied and offer accuracies comparable to results from perturbative hybrid DFT functionals. However, these new functionals are more sensitive to the choice of quadrature grid than previous generations of DFT functionals. Errors in predicted reaction energies arising from the use of the popular SG-1 integration grid, which is the default in the Q-Chem package, are significant. In particular, M06-HF reaction energies computed with the SG-1 grid exhibit errors ranging from −6.7 to 3.2 kcal mol⁻¹, relative to results computed with a very fine integration grid. This grid sensitivity is not a general problem for meta-generalized gradient approximation functionals, but is instead due to the specific functional forms used in these functionals. The large grid errors are traced to the kinetic energy density enhancement factor utilized in the exchange component of the M05-2X and the M06 functionals. This term contains empirically adjusted parameters that are of large magnitude for all of these functionals and for M06-HF in particular. The product of these large constants and modest integration errors for the kinetic energy density results in very large errors in individual contributions to the exchange energy. This gives rise to the large errors in reaction energies exhibited by these functionals for certain integration grids.

I. Introduction

Kohn—Sham density functional theory (DFT) has emerged as the preeminent choice for the computational study of organic reactions. The popularity of DFT over traditional ab initio methods in this context stems from a number of factors, including favorable scaling with system size combined with relatively high-accuracy, widespread availability of analytic first and second energy derivatives for efficient geometry optimizations and vibrational frequency computations, and efficient implementations in popular electronic structure theory packages. Because the necessary integrals over exchange—correlation functionals cannot be evaluated in closed form, Kohn—Sham DFT computations typically

$$\int F(\mathbf{r})d\mathbf{r} \approx \sum_{A}^{\text{atoms}} \sum_{g}^{\text{grid}} w_g p_A(\mathbf{r}_g) F(\mathbf{r}_g)$$
 (1)

where w_g is the quadrature weight at the corresponding grid point \mathbf{r}_g , and the atomic partitioning function, $p_A(\mathbf{r}_g)$, is defined such that $\sum_A^{\text{atoms}} p_A(\mathbf{r}_g) = 1$ at each point in space.

Various quadrature methods and atomic partitioning functions have been devised, several of which enjoy widespread use.^{1–5} It has long been known (though not always appreciated!) that the choice of integration grid can significantly affect computed molecular properties.^{6–9} Martin, Bauschlicher, and Ricca⁶ studied the grid sensitivity of B3LYP-

rely on numerical quadrature schemes. In most quantum chemistry programs, these integrals are approximated as a sum of contributions from atom-centered grids:

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computed molecular properties and highlighted problems with popular grids for several third-row transition-metal systems and for simple hydrocarbon radicals. More recently, Papas and Schaefer⁷ compared BLYP and B3LYP energies from the Gaussian, Molpro, NWChem, Q-Chem, and GAMESS packages, using default and finer integration grids, to assess the precision that can be expected from the integration schemes in these packages. Dressler and Thiel⁸ analyzed the effect of integration grids on DFT-computed anharmonic force fields, and Termath and Sauer⁹ examined their effect on DFT-based direct molecular dynamics simulations.

Some new meta-GGAs (meta-generalized gradient approximation), which explicitly depend on the kinetic energy density, offer significant advantages over previous generations of functionals. ^{10,11} A shortcoming of these functionals that is often neglected, however, is the increased sensitivity to the choice of integration grid. By default, popular electronic structure programs utilize quadrature grids that were developed and refined based on previous generations of DFT functionals. Unfortunately, integration grids that proved adequate for these older functionals can lead to significant errors when utilized with some new meta-GGAs. ^{12–17}

In 2004, Johnson et al. 12 demonstrated that potential energy curves for dispersion-bound complexes computed with VS98 18 and other meta-GGA functionals are prone to spurious oscillations unless very large integration grids are used. Gräfenstein and Cremer subsequently proposed 17 the use of locally augmented radial integration grids to combat these issues in a cost-effective way. In 2009, Johnson and co-workers 13 revisited these grid errors and showed that the grid sensitivity originates from singularities near the intermonomer midpoint in kinetic energy density-dependent functional forms present in many meta-GGAs. Similarly, Gräfenstein, Izotov, and Cremer 14 attributed irregularities in certain meta-GGA-predicted energies for stretched covalent bonds to singularities in the self-interaction correction term present in the correlation part of these functionals.

The oscillations in meta-GGA-computed interaction potentials for dispersion-bound complexes¹² were addressed in the design of the M06 suite of functionals through the elimination of problematic terms in the VS98 functional, on which the M06 functionals are in part based.¹⁰ The resulting changes in the M06 functionals offer improved performance in this regard, although with some popular grids problems still arise.^{13,15} For example, Merz and co-workers¹⁵ recently analyzed the performance of a number of methods for the prediction of potential energy curves for model noncovalent interactions. The M06-2X and M06-L DFT functionals outperformed the other methods tested but yielded discontinuous energy curves when used with the popular SG-1 integration grid.¹⁶

Grid issues with these functionals are not limited to dispersion-bound complexes. For example, Csonka and coworkers¹⁹ benchmarked a variety of DFT functionals for the prediction of geometries and conformational energies in a series of saccharides. Although M05-2X yielded the most reliable results of the tested methods when a dense integration

grid was used, results computed using the default integration grid in the Jaguar program package lead to larger errors in energies and problems in geometry optimizations. Similarly, Scuseria and co-workers²⁰ recently reported geometry optimization convergence problems and spurious imaginary frequencies when pairing functionals from the M06 suite with various popular integration grids.

In the present work, we focus on the grid requirements for meta-GGA-predicted energies for a set of 34 organic isomerization reactions recently published by Grimme and co-workers.²¹ This set, shown in Scheme 1, constitutes a diverse collection of reactions that are small enough for the application of accurate benchmark computations yet representative of the diverse changes in bonding that occur in organic reactions. Reference "experimental reaction energies" were derived from standard enthalpies of formation corrected for zero point vibrational and thermal effects by Grimme.²¹ All of the reactions are endothermic as written. Our primary focus is on grid errors for M05-2X and the M06 family of functionals, ¹⁰ which have emerged as promising new functionals for diverse chemical applications.²²

II. Theoretical Methods

Single point energies were computed for the molecular species in Scheme 1 using five meta-GGA DFT functionals paired with various DFT integration grids. The TZV(2df,2pd) basis set 23,24 was used for all computations, which were executed at B3LYP/TZV(p,d) optimized geometries taken from ref 21. The TZV(2df,2pd) basis set comprises the Alrichs TZV triple- ζ quality basis set 23 plus polarization functions from the cc-pVTZ basis set. 24 Grid errors for other popular basis sets are expected to be similar. The meta-GGAs tested are VS98, 18 M05-2X, 25 M06-L, 26 M06-HF, 27 M06, 10 and M06-2X. 10 For comparison, grid errors for B3LYP, 28 PBE, 29,30 and TPSS 31 are also presented.

The atom-centered grids utilized in popular DFT codes are constructed as a direct product of sets of N^r radial and N^{Ω} angular grid points:

$$\sum_{g}^{\text{grid}} w_g F(\mathbf{r}_g) = \sum_{i=1}^{N^r} w_i^r \sum_{i=1}^{N^{\Omega}} w_j^{\Omega} F(\mathbf{r}_i, \theta_j, \phi_j)$$
(2)

Defining an integration grid requires a choice of atomic partitioning function and the number, weights, and distribution of radial and angular grid points.

Four popular integration grids were tested (see Table 1). Grids labeled Q-Chem, NWChem, and Gaussian are equivalent to the default grids in those packages. The default Q-Chem grid is the popular SG-1 grid of Gill, Johnson, and Pople. All of the tested grids rely on Lebedev's angular quadrature. The radial components of these grids are from either an Euler—Maclaurin quadrature with the coordinate transformation of Murray, Handy, and Laming (Euler) or a modification of the Murray—Handy—Laming scheme published by Mura and Knowles (MK). For the atomic partitioning functions, the tested grids use the scheme of Becke, the modification introduced by Stratmann, Scuseria, and Frisch (SSF), or an unpublished modification of the SSF scheme (Erf1) implemented in NWChem, 33,34 in which

Table 1. Partitioning Function, Radial Quadrature Method, and Number of Radial and Angular Points for Tested Grids^a

	partitioning	radial quadrature	radial points	angular points
Q-Chem NWChem	Becke Frf1	Euler-Maclaurin Mura-Knowles	50 49	194 434
Gaussian03	SSF	Euler-Maclaurin	49 75	302
Fine	Erf1	Mura-Knowles	70	590
Xfine	Erf1	Mura-Knowles	100	1 202

^a All grids are pruned and utilize Lebedev's angular quadrature.¹

the partition function weights are written in terms of products of error functions.

Most program packages utilize automatically pruned integration grids, in which the number of angular points is dependent on the radial coordinate. This leads to significant reductions in computational cost with minimal loss in accuracy compared to the use of the unpruned grid. The effect of pruning was tested for selected reaction energies. Errors arising from pruning were less than 0.1 kcal mol⁻¹ across all grids considered and are considered inconsequential. Presented results are based on pruned grids.

Energies computed using the NWChem "Xfine" grid (see Table 1) were used as a benchmark, i.e., $\Delta E_{\rm grid}^{\rm error} = \Delta E_{\rm grid} - \Delta E_{\rm Xfine}$. To confirm that reaction energies computed using this grid are converged with respect to integration grid density, the M06-2X reaction energies, which are the most sensitive to the choice of grid, were also computed using a very large grid with 300 radial points and 1202 angular

points. The mean absolute deviation in reaction energies between this very large integration grid and the Xfine grid is only 0.0003 kcal mol⁻¹; the maximum deviation is 0.003 kcal mol⁻¹. With these meta-GGA functionals, use of default convergence criteria sometimes lead to convergence problems and resulted in erratic and seemingly irreproducible grid errors. This was particularly true for the less dense grids and the M06-HF functional. Care was taken to ensure that all energies are converged to the correct Kohn-Sham solution to a precision of at least 10^{-10} au. Converging DFT energies to this precision is often hampered by the screening thresholds for integrals and the electron density employed by popular DFT programs. In this work, screening thresholds of at least 10^{-14} au were used. Much larger errors than reported here could be encountered with these functionals if care is not taken to tightly converge results.

All computations were carried out using a locally modified version of NWChem 5.1. 33,34

III. Results and Discussion

The present work is primarily concerned with grid errors in computed reaction energies, not the error in the reaction energies themselves or the grid errors in absolute energies. However, since the accuracy of most of the tested functionals has not been previously assessed for the reactions in Scheme 1, a brief analysis of the reaction energy errors compared to experiment is presented first. This allows the subsequent analysis of grid errors to be put into the perspective of the inherent error in the DFT computations.

Table 2. Experimental Reaction Energies (ΔE , from ref 21) and Errors in Predicted Energy for the Reactions in Scheme 1, Relative to Experiment^a

Tiolative to Experiment												
	ΔΕ	B3LYP	PBE	TPSS	VS98	M05-2X	M06-L	M06	M06-2X	BMK	mPW2-PLYP	B2-PLYP
1	1.6	-3.6	-4.7	-4.5	-4.7	-2.1	-2.1	-0.6	-1.1	-2.7	-0.9	-0.9
2	21.9	1.9	-2.7	-2.4	0.2	-0.8	-3.3	-2.6	-1.7	0.8	2.5	2.5
3	7.2	1.8	-1.9	-1.5	1.4	-2.0	-4.4	-4.3	-3.4	-1.1	0.5	0.6
4	1.0	0.3	0.1	0.1	-0.7	0.2	-0.3	0.0	0.1	0.1	0.2	0.2
5	0.9	-0.6	-0.5	-0.7	1.2	0.3	-0.4	0.0	0.1	-0.2	-0.1	-0.2
6	2.6	0.6	1.0	0.7	1.2	0.0	1.4	0.9	0.3	0.3	0.2	0.3
7	11.1	4.2	-0.7	0.5	4.9	1.3	-1.5	-1.2	-0.4	-4.7	1.9	2.1
8	22.9	-3.0	-2.8	-4.0	-5.2	-2.2	-6.1	-4.9	-4.3	0.6	-1.6	-1.8
9	6.9	1.3	2.1	1.9	2.7	-0.1	2.7	1.3	0.0	0.4	0.6	0.7
10	3.6	-2.8	-2.1	-2.8	6.6	0.2	-1.5	0.3	-0.2	-1.2	-1.1	-1.2
11	1.9	-9.8	-7.2	-8.2	29.3	-0.7	-2.0	0.4	-0.8	-3.9	-4.6	-5.0
12	46.9	10.2	4.0	4.5	3.8	4.0	7.3	3.3	3.3	-0.4	6.1	6.0
13	36.0	3.2	3.1	3.9	5.3	3.9	3.9	2.1	2.0	3.4	3.7	3.7
14	24.2	-0.7	0.8	-0.9	1.5	-1.1	0.9	-1.2	-2.2	-4.1	0.3	0.6
15	7.3	0.0	0.3	-1.2	0.8	0.4	-0.1	-0.2	-0.1	0.2	0.4	0.5
16	10.8	1.6	-2.3	-1.0	1.2	-2.3	-4.6	-4.4	-3.6	-1.1	0.4	0.5
17	27.0	-1.4	-1.5	-4.1	2.7	1.0	-0.6	-0.1	-0.1	0.2	0.2	0.1
18	11.2	0.3	-0.6	0.8	-1.4	1.7	-1.2	0.0	0.7	-1.4	0.8	0.6
19	4.6	-0.5	-0.8	-0.7	-0.6	0.4	-0.6	-0.2	0.0	-0.2	-0.2	-0.3
20	20.2	-1.8	-3.1	-2.9	-2.6	-1.3	-1.9	-0.7	-0.9	0.4	-1.6	-1.9
21	0.9	0.3	0.2	0.1	0.5	0.2	0.3	0.4	0.2	0.3	0.2	0.2
22	3.2	0.4	1.2	-1.6	1.8	1.6	-0.8	0.2	0.7	-0.9	0.1	0.0
23	5.3	-0.7	-0.5	-1.9	0.3	-0.1	-0.5	-0.5	-0.6	-0.5	-0.2	-0.3
24	12.5	-1.9	-1.1	-4.3	-1.4	-1.1	-2.4	-2.0	-1.9	-2.0	-1.3	-1.3
25	26.5	1.5	-1.8	-3.6	1.9	-2.9	-2.6	-1.9	-3.7	0.4	1.1	1.1
26	18.2	-2.2	-1.4	-5.3	-1.9	-1.6	-2.1	-1.4	-2.1	-2.2	-1.7	-1.8
27	64.2	-3.5	-5.5	-11.2	-1.1	4.5	-4.3	2.3	2.0	8.0	0.0	-0.9
28	31.2	2.2	-0.2	-2.2	6.6	-1.0	1.6	0.5	-1.4	-1.5	1.5	1.7
29	11.9	-3.0	-0.4	-1.2	-9.6	1.3	-1.9	-0.1	0.0	1.7	-0.5	-0.9
30	9.5	0.1	-0.5	-2.1	0.0	0.3	-0.2	-0.2	-0.3	0.1	0.1	0.0
31	14.0	-3.0	0.6	-0.8	-5.7	1.3	-0.6	0.6	0.2	2.0	-0.3	-0.7
32	7.1	-3.7	-1.5	-1.4	-10.7	2.4	-3.9	-1.4	1.2	-0.6	-1.6	-2.3
33	5.6	4.6	5.0	1.7	-4.2	2.3	2.7	2.7	2.3	2.7	3.5	3.6
34	7.3	-0.4	1.4	0.2	3.0	-0.5	2.4	1.2	0.1	0.5	-0.4	-0.4
MAD		2.3	1.9	2.5	3.7	1.4	2.2	1.3	1.2	1.3	1.2	1.3
MSD		-0.2	-0.7	-1.6	0.8	0.2	-0.8	-0.3	-0.5	-0.4	0.2	0.2
min		-9.8	-7.2	-11.2	-10.7	-2.9	-6.1	-4.9	-4.3	-4.7	-4.6	-5.0
max		10.2	5.0	4.5	29.3	4.5	7.3	3.3	3.3	3.4	6.1	6.0

^a Mean absolute deviations (MAD), mean signed deviations (MSD), and minimum and maximum deviations (min and max, respectively) are also provided for each functional. The Xfine grid and TZV(2df,2pd) basis set were used. BMK, mPW2-PLYP, and B2-PLYP results are from ref 21. All values are in kcal mol⁻¹.

A. Performance of Meta-GGA Functionals for Reaction Energies. Errors in DFT-computed reaction energies, relative to experimental results, ²¹ are shown in Table 2. These values were computed using the Xfine integration grid. Mean signed deviations (MSD), mean absolute deviations (MAD), and the range of deviations from experiment are plotted in Figure 1. Of the functionals tested here, the M05-2X, M06, and M06-2X functionals offer the best performance, although even these functionals exhibit errors approaching ± 5 kcal mol⁻¹ for selected reactions. The mean deviations for the M05-2X, M06, and M06-2X functionals are comparable to those from mPW2-PLYP, 37 B2-PLYP, 38 and BMK 39 for this test set.²¹ The B3LYP, PBE, TPSS, M06-L, and M06-HF functionals perform poorly and are all comparable. VS98 performs the worst for these reactions, delivering the largest mean error and a very large error for reaction 11.

B. Errors with Popular Integration Grids. Integration grid errors have been assessed for the 34 organic isomerizations in Scheme 1. The grid errors for four popular quadrature grids are summarized in Table 3. Details for each reaction are provided in the Supporting Information. Normal-

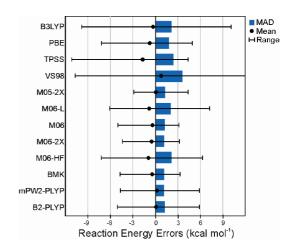


Figure 1. MAD, MSD, and range of deviations from experiment for computed energies of the reactions in Scheme 1 (kcal mol⁻¹).

ized error distributions for the B3LYP, PBE, and TPSS functionals are shown in Figure 2. Analogous plots for VS98, M05-2X, and the four M06 functionals are shown in

Table 3. Analysis of Grid Errors for the Energies of the Reactions in Scheme 1 for Four Popular DFT Integration Grids^a

unus				
	Q-Chem	NWChem	Gaussian	fine
B3LYP MAD MSD min max	0.08 0.04 -0.17 0.69	0.00 0.00 -0.02 0.04	0.01 0.00 -0.02 0.08	0.00 0.00 -0.01 0.01
PBE MAD MSD min max	0.10 0.05 -0.21 0.85	0.00 0.00 -0.02 0.05	0.01 0.00 -0.03 0.10	0.00 0.00 -0.01 0.01
TPSS MAD MSD min max	0.09 0.05 -0.19 0.85	0.01 0.00 -0.03 0.05	0.01 0.00 -0.03 0.09	0.00 0.00 -0.01 0.01
VS98 MAD MSD min max	0.25 0.13 -0.38 2.80	0.03 0.00 -0.13 0.21	0.04 0.00 -0.22 0.37	0.01 0.00 -0.07 0.05
M05-2X MAD MSD min max	0.16 0.02 -0.39 0.52	0.05 -0.02 -0.45 0.10	0.01 0.01 -0.02 0.10	0.01 0.00 -0.03 0.02
M06-L MAD MSD min max	0.20 0.09 -0.84 0.76	0.05 0.03 -0.07 0.65	0.02 0.00 -0.10 0.11	0.01 0.00 -0.02 0.03
M06 MAD MSD min max	0.25 0.06 -1.04 0.75	0.04 0.02 -0.10 0.31	0.02 0.01 -0.08 0.12	0.00 0.00 -0.01 0.02
M06-2X MAD MSD min max	0.29 0.03 -2.02 0.81	0.08 -0.02 -0.39 0.23	0.02 0.00 -0.09 0.11	0.01 0.00 -0.02 0.02
M06-HF MAD MSD min max	1.20 0.05 -6.70 3.21	0.20 -0.07 -1.51 0.42	0.03 0.01 -0.12 0.13	0.02 0.00 -0.05 0.06

^a Grid errors are relative to Xfine results. All values are in kcal

Figure 3. For B3LYP, PBE and TPSS, grid errors arising from the use of the NWChem, Gaussian, and fine grids are negligible, never exceeding ± 0.1 kcal mol⁻¹. The errors are slightly larger for the Q-Chem grid, and there is one outlier (reaction 11) at 0.7, 0.9, and 0.9 kcal mol⁻¹ for B3LYP, PBE, and TPSS, respectively. Overall, the grid requirements for these three functionals are modest, and the errors resulting from any of these grids are far less than the errors in the computed reaction energies.

For VS98, M05-2X, and the four M06 functionals (Figure 3), the grid errors are significantly larger than for B3LYP, PBE, or TPSS. Even so, the errors arising from the use of the Gaussian and fine integration grids are tightly grouped around 0 and never exceed 0.15 kcal mol⁻¹ for M05-2X or

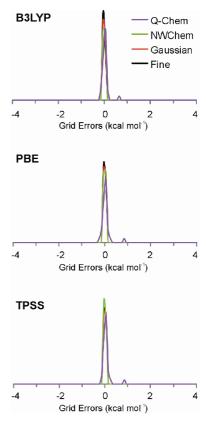


Figure 2. Normalized distributions of grid errors (kcal mol⁻¹) for the energies of the reactions in Scheme 1, computed using four popular quadrature grids paired with the B3LYP, PBE, and TPSS functionals. Grids labeled Q-Chem, NWChem, and Gaussian are equivalent to the default grids in those packages. All computations were carried out using the NWChem program.33,34

the M06 suite of functionals. For the NWChem grid, the errors are only slightly larger than the Gaussian and fine grid results, although the magnitude exceeds 0.5 kcal mol⁻¹ in several cases. The most troubling results arise from use of the Q-Chem (SG-1) grid, ¹⁶ which leads to significant errors in computed reaction energies. The problems are most severe for the M06-HF functional, for which the grid errors range from -6.7 to 3.2 kcal mol⁻¹. For M06-HF computations employing this popular grid, these grid errors will be the dominant source of error, and the predicted reaction energies will generally be qualitatively different than those computed with finer integration grids.

Across all of these functionals, the largest grid errors occur for reaction 10, followed by 11. However, not all functionals behave uniformly. For example, with the Q-Chem grid, the M06-HF predicted energy for reaction 11 is within 1.7 kcal mol⁻¹ of the Xfine reference, while the energy for reaction 10 deviates by 5.6 kcal mol⁻¹. For the M05-2X functional, the energy for reaction 11 falls 0.4 kcal mol⁻¹ from the reference value, even though the value for reaction 10 is in error by only -0.1 kcal mol⁻¹. This nonsystematic behavior is indicative of an underlying cancelation of more sizable errors, which is discussed in detail below. Reaction 10 has been highlighted previously 16 as a case prone to grid errors. Gill and co-workers¹⁶ attributed this to the large difference in shape of the two isomers of pentane. Essentially, for more

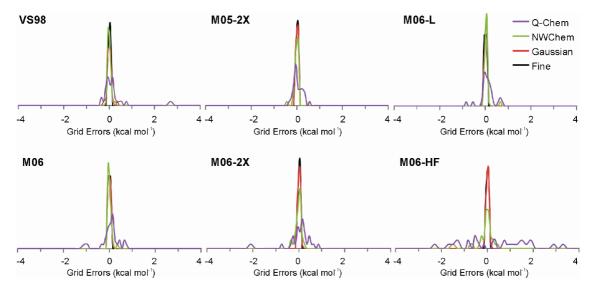


Figure 3. Normalized distributions of grid errors (kcal mol⁻¹) for the energies of the reactions in Scheme 1 using four popular quadrature grids.

"compact" molecules (e.g., neopentane), the same number of grid points covers a smaller amount of space, so the grid error in the absolute energy should be reduced compared to that of the error for less compact species (e.g., *n*-pentane).

C. Effects of Atomic Partitioning Function and Radial Quadrature Methods. To quantify the effect of different atomic partitioning functions and radial quadrature schemes, the errors associated with grids comprising 50 radial points and 194 angular points are examined more closely. Errors for six combinations of atomic partitioning function and radial quadrature scheme are summarized in Table 4. The left-most column (Becke-Euler) is the Q-Chem (SG-1) grid. 16 Normalized error distributions for the six meta-GGAs paired with these six grids are shown in Figure 4. The Euler and MK radial quadrature schemes^{2,3} perform similarly, regardless of the choice of partitioning function. The MK scheme provides a minimal reduction in mean grid errors compared to that of the Euler method. On the other hand, the choice of partitioning function has a significant impact on grid errors for all functionals considered. The SSF or Erf1 partitioning functions⁴ result in significantly smaller mean errors compared to that of the Becke results.⁵ These results are in accord with previous findings of Martin et al.,6 and the assertion by Scuseria and co-workers⁴ that the SSF scheme⁴ is numerically more stable than that of Becke.⁵ However, for the M06-HF functional, none of the grids tested delivers reaction energies with errors consistently below 0.5 kcal mol^{-1} .

D. Origin of Grid Errors. In order to unravel the origins of the large grid errors arising from the use of the Q-Chem (SG-1) grid, ¹⁶ the contributions to these errors are examined in more detail. Normalized error distributions for the exchange (E_x) and correlation (E_c) components of the reaction energies are shown in Figure 5. ⁴⁰ Mean grid errors in E_x and E_c are given in the Supporting Information (Table S2). For all but the VS98 functional, the errors in E_x swamp those arising from E_c , for which the MADs are less than 0.1 kcal mol⁻¹. For the VS98 functional, the situation is reversed; in this case, E_c exhibits a larger MAD than E_x . Regardless, the

Table 4. Analysis of Grid Errors for the Energies of the Reactions in Scheme 1 for Combinations of Three Partitioning Functions and Two Radial Quadrature Schemes^a

001101111	50						
	Eul	er-Macla	urin	Mura-Knowles			
	Becke	SSF	Erf1	Becke	SSF	Erf1	
VS98							
MAD	0.26	0.18	0.19	0.27	0.15	0.16	
MSD	0.14	0.10	0.12	0.13	0.05	0.05	
min	-0.38	-0.24	-0.27	-0.35	-0.28	-0.33	
max	2.80	1.19	1.37	3.12	0.93	0.94	
M05-2>	(
MAD	0.16	0.06	0.06	0.16	0.06	0.07	
MSD	0.03	-0.01	-0.01	0.04	-0.01	-0.02	
min	-0.39	-0.13	-0.14	-0.32	-0.18	-0.21	
max	0.52	0.19	0.25	0.37	0.11	0.15	
M06-L							
MAD	0.21	0.09	0.08	0.20	0.07	0.07	
MSD	0.09	0.05	0.04	0.08	0.02	0.02	
min	-0.84	-0.13	-0.21	-0.67	-0.16	-0.17	
max	0.76	0.54	0.34	0.71	0.30	0.27	
M06							
MAD	0.25	0.08	0.08	0.23	0.09	0.09	
MSD	0.06	0.02	0.01	0.07	0.01	0.01	
min	-1.04	-0.17	-0.19	-0.82	-0.18	-0.16	
max	0.75	0.35	0.32	0.69	0.48	0.43	
M06-2	<						
MAD	0.29	0.07	0.08	0.26	0.07	0.07	
MSD	0.05	0.00	0.00	0.05	-0.02	-0.02	
min	-2.02	-0.36	-0.49	-1.85	-0.20	-0.28	
max	0.81	0.30	0.33	0.66	0.14	0.17	
M06-HI	F						
MAD	1.24	0.24	0.30	1.13	0.20	0.25	
MSD	0.05	-0.04	-0.03	0.05	-0.05	-0.04	
min	-6.70	-1.35	-1.71	-6.57	-1.24	-1.61	
max	3.21	0.69	0.89	2.92	0.51	0.65	

 $^{^{\}rm a}$ Results for B3LYP, PBE, and TPSS are available in the Supporting Information. All grids have 50 radial points and 194 angular points. Errors are relative to Xfine results. All values are in kcal $\rm mol^{-1}$.

unsettling grid errors exhibited by M05-2X and M06 functionals arise from the exchange energies.

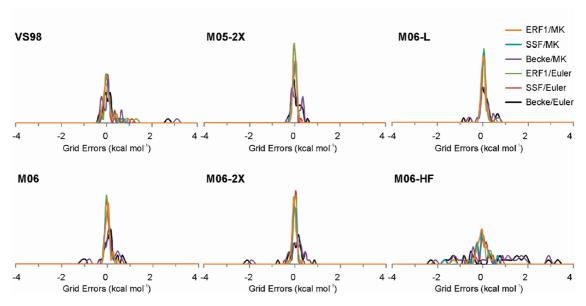


Figure 4. Normalized distributions of grid errors (kcal mol⁻¹) for the energies of the reactions in Scheme 1 using six combinations of partitioning functions (Becke, SSF, or Erf1) and radial quadrature schemes [Euler-Maclaurin (Euler)² or Mura-Knowles (MK)³]. All grids are pruned and have 50 radial points and 194 angular points. The Becke/Euler combination corresponds to the default Q-Chem (SG-1) grid.16

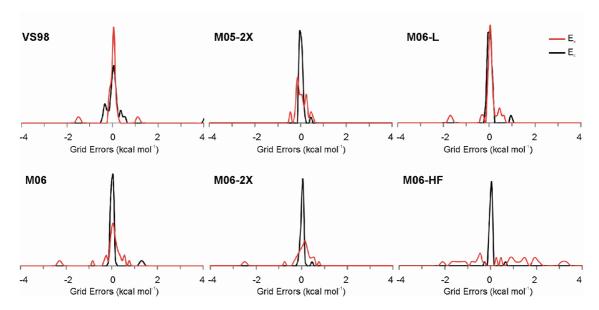


Figure 5. Normalized distributions of grid errors (kcal mol⁻¹) for the contribution of E_x and E_c to the total energies for the reactions in Scheme 1, computed with the Q-Chem grid.

The exchange functional utilized for the M06 suite¹⁰ is a linear combination of the functional forms of the M05-2X²⁵ and VS98¹⁸ exchange functionals:

$$E_{\rm x}^{\rm M06} = E_{\rm x}^{\rm M05-2X} + E_{\rm x}^{\rm VS98} \tag{3}$$

 $E_x^{\text{M05-2X}}$ is the PBE exchange functional, $F_{X\sigma}^{\text{PBE}}$, multiplied by a kinetic energy density-dependent term, $f(w_{\sigma})$. This function is referred to by Truhlar and co-workers 10 as the "kinetic energy density enhancement factor" and is written as a power series in w_{σ} . The w_{σ} , in turn, is a function of the spin kinetic energy density. 41 The form for the M06 exchange functional is given in eqs 4–11, where ρ_{σ} , $\nabla \rho_{\sigma}$, and τ_{σ} are the spin density, gradient, and kinetic energy density, respectively, and d_i and a_i are empirically

determined parameters. Setting a_0 to 1.0 and the other a_i and d_i constants to zero gives the standard PBE GGA exchange functional, while setting the a_i constants to zero (and the d_i to the appropriate values) yields VS98 exchange. The M05-2X exchange functional is obtained by setting the d_i constants to zero, and the a_i to the appropriate values.

$$E_{\mathbf{x}}^{\mathbf{M06}} = \sum_{\sigma}^{\alpha,\beta} \int d\mathbf{r} \left[F_{X\sigma}^{\mathbf{PBE}}(\rho_{\sigma}, \nabla \rho_{\sigma}) f(w_{\sigma}) + \frac{d_{1}x_{\sigma}^{2} + d_{2}z_{\sigma}}{\gamma^{2}(x_{\sigma}, z_{\sigma})} + \frac{d_{1}x_{\sigma}^{2} + d_{2}z_{\sigma}}{\gamma^{2}(x_{\sigma}, z_{\sigma})} + \frac{d_{3}x_{\sigma}^{4} + d_{4}x_{\sigma}^{2}z_{\sigma} + d_{5}z_{\sigma}^{2}}{\gamma^{3}(x_{\sigma}, z_{\sigma})} \right]$$

$$(4)$$

$$\varepsilon_{X\sigma}^{\text{LSDA}} = \frac{3}{2} \left(\frac{3}{4\pi} \right)^{1/3} \rho_{\sigma}^{4/3} \tag{6}$$

$$x_{\sigma} = \frac{|\nabla \rho_{\sigma}|}{\rho_{\sigma}^{4/3}} \tag{7}$$

$$z_{\sigma} = \frac{2\tau_{\sigma}}{\rho^{5/3}} - \frac{3}{5} (6\pi^2)^{2/3} \tag{8}$$

$$\gamma(x_{\sigma}, z_{\sigma}) = 1 + \alpha(x_{\sigma}^2 + z_{\sigma}) \tag{9}$$

$$w_{\sigma} = (t_{\sigma} - 1)/(t_{\sigma} + 1) \tag{10}$$

$$t_{\sigma} = \frac{3(6\pi^2)^{2/3} \rho_{\sigma}^{5/3}}{10\tau_{\sigma}}$$
 (11)

Errors arising from the use of the Q-Chem grid for each term in eq 4 are summarized in the Supporting Information (Tables S2 and S3). These were computed using converged densities from Xfine computations with a locally modified version of NWChem $5.1.^{33,34}$ Distributions of grid errors arising from the two main components of the M06-HF exchange functional, $E_{\rm x}^{\rm M05-2X}$ and $E_{\rm x}^{\rm VS98}$, are shown in Figure 6a. The grid errors arise primarily from the M05-2X component.

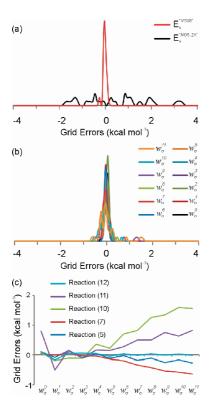


Figure 6. (a) Normalized Q-Chem grid error distributions for the $E_{\rm x}^{\rm M05-2X}$ and $E_{\rm x}^{\rm VS98}$ components of M06-HF. (b) Q-Chem grid error distribution for each power of w_{σ} in eq (12). (c) Q-Chem grid errors for each power of w_{σ} in eq (12) for representative reactions from Scheme 1.

Table 5. Coefficients in the Kinetic Energy Density Enhancement Factor (eq 5) in M05-2X (see ref 25) and the M06 Suite of Functionals (see ref 10)

	M05-2X	M06-L	M06	M06-2X	M06-HF
a_0	1.0	0.3987756	0.5877943	0.46	0.1179732
a_1	-0.56833	0.2548219	-0.1371776	-0.2206052	-1.066708
a_2	-1.30057	0.3923994	0.2682367	-0.09431788	-0.1462405
a_3	5.5007	-2.103655	-2.515898	2.164494	7.481848
a_4	9.06402	-6.302147	-2.978892	-2.556466	3.776679
a_5	-32.21075	10.97615	8.710679	-14.22133	-44.36118
a_6	-23.73298	30.97273	16.88195	15.55044	-18.30962
a_7	70.22996	-23.18489	-4.489724	35.98078	100.3903
a_8	29.88614	-56.7348	-32.99983	-27.22754	38.6436
a_9	-60.25778	21.60364	-14.4905	-39.24093	-98.06018
a ₁₀	-13.22205	34.21814	20.43747	15.22808	-25.57716
a ₁₁	15.23694	-9.049762	12.56504	15.22227	35.90404

Normalized grid error distributions for each power of w_{σ} in $f(w_{\sigma})$ are plotted in Figure 6b with all of the a_i constants set to unity, i.e.:

$$\sum_{\sigma}^{\alpha,\beta} \int d\mathbf{r} \mathbf{F}_{X\sigma}^{\text{PBE}}(\rho_{\sigma}, \nabla \rho_{\sigma}) w_{\sigma}^{i}$$
 (12)

The grid errors for each term are modest and exceed 0.5 kcal mol⁻¹ for only a few reactions. The tendency is for the magnitude of the integration error to gradually increase with the power of w_{σ} , although for some reactions the integration errors are negligible for all powers of w_{σ} (errors for representative reactions are shown in Figure 6c). The empirical constants a_i , on the other hand, vary considerably in magnitude and oscillate between negative and positive values (see Table 5). The grid errors for each term in $E_{\rm x}^{\rm M05-2X}$ will be the product of the a_i constant and the integration error for the corresponding power of w_{σ} . Because the magnitudes of the a_i constants are large, grid errors for the individual contributions to $f(w_{\sigma})$ are enormous. For example, for reaction 10 the errors in the w_{σ}^{7} and w_{σ}^{9} terms are +82and -132 kcal mol⁻¹, respectively. These sizable errors mostly cancel due to the almost monotonic change in grid errors with each power of w_{σ} (see Figure 6c) combined with the oscillations in the a_i constants. This cancelation is incomplete, of course, and gives rise to the troubling spread of grid errors discussed above. The substantial grid errors exhibited by the M06-HF functional, in particular, are a result of the very large a_i constants defining that functional (e.g.: see the a_7 and a_9 constants for M06-HF in Table 5).

The attribution of these errors to the particular functional form used in the M06 suite of functionals is further supported by the qualitatively different grid-dependence exhibited by the meta-GGA functional TPSS. This functional, which does not contain a polynomial expansion in terms of the kinetic energy density, exhibits very modest dependence on the choice of integration grid; TPSS grid errors are on par with those from popular GGAs.

These grid errors are qualitatively different from those underlying the discontinuities in meta-GGA-computed potential energy curves for dispersion-bound complexes, 12 which arise from singularities in the τ -dependent functional forms. 13 The grid errors in reaction energies arise from modest errors in the integration of w_{σ} amplified by the large empirical a_i constants. One consequence is that these errors

vary smoothly across potential energy surfaces. This is demonstrated for neopentane in the Supporting Information (Figures S1 and S2), in which grid errors in the M06-HF energy computed with the Q-Chem grid are plotted as a function of the C-C and C-H bond lengths.

IV. Summary and Conclusions

DFT is invaluable in the computational study of organic reactions and can provide accurate energies for large molecular systems that are beyond the reach of traditional ab initio methods. New DFT functionals offer increased accuracy and broader applicability compared to those of previous generations of functionals and have enabled the application of DFT to myriad new problems in organic chemistry and molecular biology. However, these new functionals are not without drawbacks, one of which is the increased sensitivity of energies and other properties to the choice of integration grid. Previously documented 12-16 grid sensitivities exhibited by these functionals include the prediction of potential energy curves for dispersion-bound complexes with spurious oscillations as well problems with predicted energies, geometries, and vibrational frequencies. 19,20

We have quantified integration grid errors for six meta-GGA functionals (VS98, M05-2X, M06-L, M06, M06-2X, and M06-HF) paired with popular integration grids for 34 organic isomerization energies. The popular SG-1 grid, ¹⁶ which is the default in the Q-Chem package, leads to large errors for all of these functionals and very large errors for M06-HF. This grid should not be used with any of these functionals. By contrast, the grid errors in B3LYP, PBE, and TPSS computed energies are small for all of the grids tested. Use of the SSF⁴ or Erf1 atomic partitioning functions reduces grid errors compared to that of the standard SG-1 grid, 16 which utilizes the partitioning function of Becke. ⁵ However, the M06-HF functional still exhibits grid errors exceeding 0.5 kcal mol⁻¹ for several of the reactions.

The grid errors exhibited by the M06 suite of functionals arise from integration errors in the exchange component of the energy. In particular, the significant errors arising from the use of the Q-Chem (SG-1) grid are due to the large empirical constants in the kinetic energy density enhancement factor. Some of these constants are of considerable size and amplify modest errors in the integration of the kinetic energy density. This is not a general weakness of meta-GGAs but a problem arising from the particular functional form used in the M06 suite of functionals.

Zhao and Truhlar recently published¹¹ the M08-HX and M08-SO functionals, which incorporate more flexible functional forms than members of the M06 suite and contain an altered self-interaction correction term that avoids the numerical instabilities discussed by Gräfenstein, Izotov, and Cremer. 14 Although the grid errors associated with these new functionals were not tested here, the kinetic energy density enhancement factor in M08-HX and M08-SO is the same as in the M06 functionals. Moreover, the empirical coefficients in this factor are larger than in any of the M06 suite of functionals, so M08-HX and M08-SO grid errors are expected to be even more severe than observed for M06-HF.

The popularity of DFT and ease with which many computational chemistry program packages can be used has led to a continued increase in the application of DFT to chemical problems by nonspecialists. While this is certainly a welcome development and a testament to the maturity of the field of Kohn-Sham DFT, the present results offer a poignant reminder of the dangers of employing "default" options in any program package. These defaults are not suitable for all applications, and in the case of DFT integration grids, the defaults in some cases are woefully inadequate for some meta-GGA functionals. In particular, use of the SG-1 integration grid with M05-2X or the M06 suite of functionals can result in significant errors in predicted reaction energies.

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Supporting Information Available: Tables of grid errors for individual reactions and absolute energies. Plots of grid errors for neopentane as a function of the C-C and C-H bond lengths. This material is available free of charge via the Internet at http://pubs.acs.org.

References

- (1) Lebedev, V. I.; Laikov, D. N. Dokl. Math. 1999, 59, 477.
- (2) Murray, C. W.; Handy, N. C.; Laming, G. L. Mol. Phys. 1993, *78*, 997.
- (3) Mura, M. E.; Knowles, P. J. J. Chem. Phys. 1996, 104, 9848.
- (4) Stratmann, R. E.; Scuseria, G. E.; Frisch, M. J. Chem. Phys. Lett. 1996, 257, 213.
- (5) Becke, A. J. Chem. Phys. 1988, 88, 1053.
- (6) Martin, J. M. L.; Bauschlicher, C. W., Jr.; Ricca, A. Comput. Phys. Commun. 2001, 133, 189.
- (7) Papas, B. N.; Schaefer, H. F. J. Mol. Struct. THEOCHEM **2006**, 768, 175.
- (8) Dressler, S.; Thiel, W. Chem. Phys. Lett. 1997, 273, 71.
- (9) Termath, V.; Sauer, J. Chem. Phys. Lett. 1996, 255, 187.
- (10) Zhao, Y.; Truhlar, D. G. Theo. Chem. Acc. 2008, 120, 215.
- (11) Zhao, Y.; Truhlar, D. G. J. Chem. Theory Comput. 2008, 4, 1849.
- (12) Johnson, E. R.; Wolkow, R. A.; DiLabio, G. A. Chem. Phys. Lett. 2004, 394, 334.
- (13) Johnson, E. R.; Becke, A.; Sherrill, C. D.; DiLabio, G. A. J. Chem. Phys. 2009, 131, 034111.
- (14) Gräfenstein, J.; Izotov, D.; Cremer, D. J. Chem. Phys. 2007, *127*, 214103.
- (15) Fusti-Molnar, L.; He, X.; Wang, B.; Merz, K. M., Jr. J. Chem. Phys. 2009, 131, 065102.
- (16) Gill, P. M. W.; Johnson, B. G.; Pople, J. A. Chem. Phys. Lett. 1993, 209, 506.

- (17) Gräfenstein, J.; Cremer, D. J. Chem. Phys. 2007, 127, 164113.
- (18) Van Voorhis, T.; Scuseria, G. J. Chem. Phys. 1998, 109, 400.
- (19) Csonka, G. I.; French, A. D.; Johnson, G. P.; Stortz, C. A. J. Chem. Theory Comput. 2009, 5, 679.
- (20) Jiménez-Hoyos, C. A.; Janesko, B. G.; Scuseria, G. E. Phys. Chem. Chem. Phys. 2008, 10, 6621.
- (21) Grimme, S.; Steinmetz, M.; Korth, M. J. Org. Chem. 2007, 72, 2118.
- (22) Zhao, Y.; Truhlar, D. G. Acc. Chem. Res. 2008, 41, 157.
- (23) Schafer, A.; Huber, C.; Ahlrichs, R. J. Chem. Phys. 1994, 100, 5829.
- (24) Dunning, T. H., Jr. J. Chem. Phys. 1989, 90, 1007.
- (25) Zhao, Y.; Schultz, N. E.; Truhlar, D. G. J. Chem. Theory Comput. 2006, 2, 364.
- (26) Zhao, Y.; Truhlar, D. G. J. Chem. Phys. 2006, 125, 194101.
- (27) Zhao, Y.; Truhlar, D. G. J. Phys. Chem. A 2006, 110, 13126.
- (28) Becke, A. D. J. Chem. Phys. 1993, 98, 5648.
- (29) Perdew, J. P.; Burke, K.; Ernzerhof, M. Phys. Rev. Lett. 1996, 77, 3865.
- (30) Perdew, J. P.; Burke, K.; Ernzerhof, M. Phys. Rev. Lett. 1997, 78, 1396.
- (31) Tao, J.; Perdew, J. P.; Staroverov, V.; Scuseria, G. Phys. Rev. Lett. 2003, 91, 146401.
- (32) Shao, Y.; Fusti-Molnar, L.; Jung, Y.; Kussmann, J.; Ochsenfeld, C.; Brown, S. T.; Gilbert, A. T. B.; Slipchenko, L. V.; Levchenko, S. V.; O'Neill, D. P.; DiStasio, R. A., Jr.; Lochan, R. C.; Wang, T.; Beran, G. J. O.; Besley, N. A.; J. M., H.; Lin, C. Y.; Voorhis, T. V.; Chien, S. H.; Sodt, A.; Steele, R. P.; Rassolov, V. A.; Maslen, P. E.; Korambath, P. P.; Adamson, R. D.; Austin, B.; Baker, J.; Byrd, E. F. C.; Dachsel, H.; Doerksen, R. J.; Dreuw, A.; Dunietz, B. D.; Dutoi, A. D.; Furlani, T. R.; Gwaltney, S. R.; Heyden, A.; Hirata, S.; Hsu, C.-P.; Kedziora, G.; Khalliulin, R. Z.; Klunzinger, P.; Lee, A. M.; Lee, M. S.; Liang, W.; Lotan, I.; Nair, N.; Peters, B.; Proynov, E. I.; Pieniazek, P. A.; Rhee, Y. M.; Ritchie, J.; Rosta, E.; Sherrill, C. D.; Simmonett, A. C.; Subotnik, J. E.; Woodcock, H. L., III; Zhang, W.; Bell, A. T.; Chakraborty, A. K.; Chipman, D. M.; Keil, F. J.; Warshel, A.; Hehre, W. J.; Schaefer, H. F.; Kong, J.; Krylov, A. I.; Gill, P. M. W.; Head-Gordon, M. Phys. Chem. Chem. Phys. 2006, 8, 3172.
- (33) Kendall, R. A.; Apra, E.; Bernholdt, D. E.; Bylaska, E. J.; Dupuis, M.; Fann, G. I.; Harrison, R. J.; Ju, J.; Nichols, J. A.; Nieplocha, J.; Straatsma, T. P.; Windus, T. L.; Wong, A. T. Comput. Phys. Commun. 2000, 128, 260.
- (34) Bylaska, E. J.; de Jong, W. A.; Govind, N.; Kowalski, K.; Straatsma, T. P.; Valiev, M; Wang, D.; Aprà, E.; Windus, T. L.; Hammond, J.; Nichols, P; Hirata, S.; Hackler, M. T.; Zhao, Y.; Fan, P.-D.; Harrison, R. J.; Dupuis, M.; Smith, D. M. A.; Nieplocha, J.; Tipparaju, V.; Krishnan, M.; Wu, T; Van Voorhis, T; Auer, A. A.; Nooijen, M.; Brown, E; Cisneros, G.; Fann, G. I.; Fruchtl, H; Garza, J; Hirao, K; Kendall, R.; Nichols, J. A.; Tsemekhman, K; Wolinski, K.; Anchell, J.; Bernholdt, D;

- Borowski, P; Clark, T; Clerc, D.; Dachsel, H; Deegan, H.; Dyall, K.; Elwood, D.; Glendening, E.; Gutowski, M.; Hess, A.; Jaffe, J.; Johnson, B.; Ju, J.; Kobayashi, R.; Kutteh, R.; Lin, Z.; Littlefield, R; Long, X.; Meng, B.; Nakajima, T; Niu, S; Pollack, L.; Rosing, M.; Sandrone, G.; Stave, M.; Taylor, H.; Thomas, G.; van Lenthe, J.; Wong, A.; Zhang, Z. NWChem, A Computational Chemistry Package for Parallel Computers, Version 5.1; Pacific Northwest National Laboratory: Richland, Washington, 2007.
- (35) Frisch, M. J.; Trucks, G. W.; Schlegel, H. B.; Scuseria, G. E.; Robb, M. A.; Cheeseman, J. R.; Montgomery, J. A., Jr.; Vreven, T.; Kudin, K. N.; Burant, J. C.; Millam, J. M.; Iyengar, S. S.; Tomasi, J.; Barone, V.; Mennucci, B.; Cossi, M.; Scalmani, G.; Rega, N.; Petersson, G. A.; Nakatsuji, H.; Hada, M.; Ehara, M.; Toyota, K.; Fukuda, R.; Hasegawa, J.; Ishida, M.; Nakajima, T.; Honda, Y.; Kitao, O.; Nakai, H.; Klene, M.; Li, X.; Knox, J. E.; Hratchian, H. P.; Cross, J. B.; Bakken, V.; Adamo, C.; Jaramillo, J.; Gomperts, R.; Stratmann, R. E.; Yazyev, O.; Austin, A. J.; Cammi, R.; Pomelli, C.; Ochterski, J. W.; Ayala, P. Y.; Morokuma, K.; Voth, G. A.; Salvador, P.; Dannenberg, J. J.; Zakrzewski, V. G.; Dapprich, S.; Daniels, A. D.; Strain, M. C.; Farkas, O.; Malick, D. K.; Rabuck, A. D.; Raghavachari, K.; Foresman, J. B.; Ortiz, J. V.; Cui, Q.; Baboul, A. G.; Clifford, S.; Cioslowski, J.; Stefanov, B. B.; Liu, G.; Liashenko, A.; Piskorz, P.; Komaromi, I.; Martin, R. L. Fox, D. J.; Keith, T.; Al-Laham, M. A.; Peng, C. Y.; Nanayakkara, A.; Challacombe, M.; Gill, P. M. W.; Johnson, B.; Chen, W.; Wong, M. W.; Gonzalez, C.; and Pople, J. A. Gaussian 03, Revision C.02; Gaussian, Inc.: Wallingford, CT, 2004.
- (36) Note that as implemented in NWChem 5.1.1, the density screening threshold used by other functionals is overridden by the VS98, M05-2X, and M06 suite of functionals. Instead, in NWChem 5.1.1, density screenings for the exchange component of these functionals are fixed at 10^{-7} or 10^{-8} au.
- (37) Grimme, S.; Schwabe, T. Phys. Chem. Chem. Phys. 2006, 8, 4398.
- (38) Grimme, S. J. Chem. Phys. 2006, 124, 034108.
- (39) Boese, A. D.; Martin, J. M. L. J. Chem. Phys. **2004**, 121, 3405.
- (40) The decomposition of grid errors for DFT energies into exchange and correlation components is complicated by the secondary effects of the grid errors on the SCF procedure. Even though the other components of the electronic energy (Coulomb and one-electron terms and HF exchange) are evaluated analytically and should exhibit no grid errors, when the Kohn—Sham orbitals are optimized, grid the errors in E_x and E_c lead to differences in all of the components of the energy for different grids. As such, to evaluate the errors arising from E_c and E_x , energies for the Q-Chem grid were evaluated using converged Kohn—Sham orbitals from Xfine grid computations. This has only a minor effect on the grid errors in the total reaction energies.
- (41) Becke, A. J. Chem. Phys. 2000, 112, 4020.

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