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# Nuclear Relaxation of H<sub>2</sub> and H<sub>2</sub>@C<sub>60</sub> in Organic Solvents

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We report an investigation of the nuclear spin-lattice relaxation of H<sub>2</sub> and H<sub>2</sub>@C<sub>60</sub><sup>1</sup> as a function of solvent and temperature. These studies explore and compare the nature of the interactions of a guest H<sub>2</sub> molecule confined transiently within the walls of a solvent cavity and a guest H<sub>2</sub> molecule encapsulated within the walls of the C<sub>60</sub>

The relaxation time  $(T_1)$  of  $H_2$  has been extensively studied in the gas phase and in liquid hydrogen at low temperatures.<sup>2,3</sup> However, there are few reports of the magnetic resonance of H<sub>2</sub> in solution<sup>4,5</sup> and none in ordinary organic solvents. To generate data for comparison with  $H_2@C_{60}$  we measured the  $T_1$  values of  $H_2$  in a range of organic solvents.

The chemical shifts of  $H_2$  and  $H_2@C_{60}$  appear at  $\sim 4.5$  ppm and  $\sim -1.3$  ppm, respectively. The value of  $T_1$  of  $H_2$  at 300 K varies significantly with solvent (Table 1): from 1.44 s (benzene) to 0.84 s (CCl<sub>4</sub>). A somewhat larger variation of  $T_1$  with the same set of solvents was found for H<sub>2</sub>@C<sub>60</sub> (Table 1): from 0.118 s (benzene) to 0.046 s (CCl<sub>4</sub>).<sup>6</sup> The values of  $T_1$  are 10–20 times smaller for  $H_2@C_{60}$  than for  $H_2$  even though the ratios of  $T_1$  for  $H_2$  and  $H_2@C_{60}$ are similar in all the solvents.

The temperature dependences of  $T_1$  for  $H_2$  and  $H_2@C_{60}$  were investigated in detail for toluene- $d_8$  (Figure 1) and for benzene- $d_6$ , 1,1,2,2-tetrachloroethane- $d_2$ , 1,2-dichlorobenzene- $d_4$ , and chloroform $d_1$ . Striking features of the data are the occurrence of a maximum for  $T_1$  at  $\sim$ 240 K for both the H<sub>2</sub> and H<sub>2</sub>@C<sub>60</sub> in toluene- $d_8$  and a ratio of  $T_1$  values which is nearly independent of temperature. A maximum of the value of  $T_1$  with temperature is also found in 1,1,2,2-tetrachloroethane- $d_2$  and chloroform- $d_1$ . For benzene- $d_6$  and 1,2-dichlorobenzene- $d_4$  in the available range of temperatures, only a decrease of  $T_1$  with increasing temperature was observed.

This kind of dependence of  $T_1$  on temperature is uncommon, although a maximum of  $T_1$  has previously been observed for small molecules such as H<sub>2</sub>O,<sup>7</sup> HCl, and HBr in solution, <sup>8,9</sup> and it is consistent with two relaxation mechanisms with different temperature dependences dominating in turn below and above 240 K for both  $H_2$  and  $H_2@C_{60}$ . Since the value of  $T_1$  for both  $H_2$  and  $H_2@C_{60}$ does not significantly change in going from benzene- $h_6$  to benzened<sub>6</sub> (Table 1), the dominating interactions determining H<sub>2</sub> and H<sub>2</sub>@C<sub>60</sub> nuclear relaxation must be *intra*molecular. Furthermore, the intramolecular dipole-dipole interaction and spin-rotation interaction are known<sup>2</sup> to be responsible for the relaxation of gaseous H<sub>2</sub> and their magnitude has been measured for H<sub>2</sub> in molecular beams.  $^{10}$  Therefore it is likely that the relaxation of  $H_2$ 

Table 1. T<sub>1</sub> of H<sub>2</sub> and H<sub>2</sub>@C<sub>60</sub> at 300 K in Various Solvents, at a Field of 11.7 T (500 MHz); Ratio  $T_1(H_2)/T_1(H_2@C_{60})$  and Viscosity of the Solvents, at 25 °C

	T	T <sub>1</sub> (s) <sup>a</sup>		
solvent	H <sub>2</sub>	H <sub>2</sub> @C <sub>60</sub>	T <sub>1</sub> ratio	$\eta$ (cP)
CCl <sub>4</sub>	0.84	0.046	18	$0.884^{b}$
chloroform- $d_1$	1.21	0.082	15	$0.539^{c}$
1,2-dichlorobenzene-d <sub>4</sub>	1.27	0.101	13	$1.322^{c}$
toluene-d <sub>8</sub>	1.28	0.104	12	$0.548^{d}$
$1,1,2,2$ -tetrachloroethane- $d_2$	1.37	0.108	13	$1.629^{c}$
benzene- $d_6$	1.42	0.118	12	$0.636^{e}$
benzene	1.44	0.118	12	$0.599^{e}$

 $^a \pm$  5%.  $^b$  Reference 22.  $^c \eta$  of protonated solvents, ref 22.  $^d$  Reference

in solution also depends on the competition between intramolecular dipole—dipole interaction and spin-rotation interaction.

The contribution to  $1/T_1$  (in extreme narrowing conditions) from intramolecular dipolar and spin-rotation interaction may be estimated by eq  $1^2$  and eq  $2,^{11,12}$  respectively:

$$\frac{1}{T_{\text{Idip}}} = \frac{3 \gamma_{\text{H}}^4 \hbar^2}{2r^6} \tau_{\text{dip}} \tag{1}$$

$$\frac{1}{T_{\rm 1sr}} = \frac{4Ik_{\rm B}TC^2}{3\hbar^2} \tau_{\rm sr}$$
 (2)

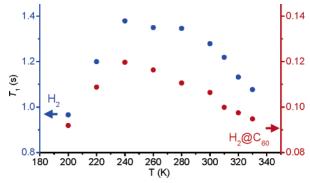
where  $\gamma_{\rm H}$  is the magnetogyric ratio for the proton, r is the equilibrium internuclear distance of H<sub>2</sub> (0.74 Å), C is the spinrotation coupling constant  $(7.16 \times 10^5 \text{ rad s}^{-1})$ , <sup>10</sup> *I* is the moment of inertia of H<sub>2</sub> (4.6  $\times$  10<sup>-48</sup> kg m<sup>2</sup>), and  $k_{\rm B}$  is the Boltzmann constant. The correlation times  $au_{\text{dip}}$  and  $au_{\text{sr}}$  are measures of the timedependent fluctuations in the orientation and angular velocity of H<sub>2</sub>, respectively. Both correlation times are expected to be functions of viscosity and temperature which depend on the details of the motion of H<sub>2</sub> molecules and the surrounding medium.<sup>11</sup>

Qualitatively, the dipole-dipole interaction (eq 1) might account for the observed increase of  $T_1$  with temperature because faster molecular reorientations correspond to shorter  $au_{
m dip}$ . On the other hand, the spin-rotation interaction might account for the observed decrease of  $T_1$  with temperature both through the explicit temperature dependence (eq 2) and through  $\tau_{\rm sr}$  temperature dependence. 13 Assuming that the relaxation at the lowest temperature is dominated by the dipolar interaction, from the value of T<sub>1</sub> at 200 K and eq 1 it is possible to estimate the correlation time  $au_{
m dip}$  of the process that modulates this interaction:  $\tau_{\text{dip}}(H_2) = 0.20 \text{ ps}$  and  $\tau_{\text{dip}}(H_2@C_{60})$ = 2.1 ps. In a similar way, assuming that at the highest temperature the relaxation is dominated by the spin-rotation interaction, from

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**Figure 1.**  $T_1$  (s) of  $H_2$  (blue dots) saturated at room temperature in deoxygenated toluene- $d_8$  and  $H_2@C_{60}$  (red dots) dissolved in toluene- $d_8$ , at temperatures ranging from 200 to 330 K, at a field of 11.7 T.

the values of  $T_1$  at 330 K and eq 2, the correlation time  $\tau_{\rm sr}$  of the process which modulates this interaction is obtained:  $\tau_{\rm sr}(H_2) = 0.72$  ps and  $\tau_{\rm sr}(H_2@C_{60}) = 8.2$  ps.<sup>14</sup>

It is instructive to compare these values to the characteristic free rotation times,  $^{15}$   $\tau_{FR}=(2\pi/9)(I/k_BT)^{1/2},$  of 0.02 and 3 ps calculated for  $H_2$  and  $C_{60}$ , respectively, at 330 K, and the values for the reorientation time measured for  $C_{60}$  in toluene at room temperature:  $\tau_{dip}(C_{60})=7{-}10$  ps and extrapolated at 200 K:  $\sim\!30$  ps.  $^{16,17}$  An estimate of  $\tau_{dip}$  can be made for  $H_2$  in toluene- $d_8$  at 200 K using the Stokes–Einstein–Debye relationship ( $\tau=4\pi r^3\eta/3k_BT$ ), the viscosity at 200 K (4.5 cP),  $^{18}$  and the van der Waals radius (1.38 Å  $^{19}$ ):  $\tau_{SED}\approx18$  ps.

Comparisons of the various estimates of correlation times for  $H_2$  and  $H_2@C_{60}$  in toluene- $d_8$  lead to the following qualitative conclusions

Modulation of the angular velocity of  $H_2$  in both environments occurs somewhat more slowly than the frequency of reorientation  $(\tau_{sr} > \tau_{dip})$ , since at 330 K  $\tau_{dip}$  may be safely assumed even shorter than those calculated at 200 K and in the case of  $\tau_{dip}(H_2)$  close to  $\tau_{FR}$  of  $H_2$ . This implies that the  $H_2$  rotates through large angles between collisions with the solvent or walls of the  $C_{60}$  cavity.  $^{20}$ 

The value of  $au_{dip}$  for  $H_2$  in solution is much smaller than expected from simple Stokes–Einstein–Debye behavior ( $au_{dip}(H_2) \ll au_{SED}$ ). This is consistent with the reduced influence of viscous forces on reorientation expected for a nearly spherical molecule rotating under "slip" conditions.<sup>21</sup>

 $au_{
m dip}(H_2@C_{60})$  is larger than  $au_{
m dip}(H_2)$ . This suggests sufficient interaction between  $H_2$  and the walls of the  $C_{60}$  cavity to entice the smaller molecule to follow the rotation of the larger, although  $au_{
m dip}(C_{60})$  is still one order of magnitude larger than  $au_{
m dip}(H_2@C_{60})$ . The absence of correlation between solvent viscosity and the value of  $T_1$  in the different solvents (Table 1) is expected for  $H_2$  in solvents under slip conditions and does not conflict with the latter hypothesis on  $H_2@C_{60}$  because  $au_{
m dip}(C_{60})$  is much smaller than its  $au_{
m SED}$ , being close to its  $au_{
m FR}$  at room temperature  $^{16}$  and showing no dependence on viscosity when measured in different solvents.  $^{17}$ 

The value of  $T_1$  of  $H_2$  and  $H_2@C_{60}$  varies with solvent (Table 1), although a correlation between the values of  $T_1$  and viscosity or dielectric constant of the solvent could not be found. These observations and the above considerations on the correlation times support an approximate model in which it is the motions of the  $H_2$  and  $H_2@C_{60}$  in the solvent cages that are significant, along with the collisions or interactions of  $H_2$  with the concave walls of  $C_{60}$ . The similarity of the ratio of the values of  $T_1$  for all of the solvents investigated suggests that the motions of  $H_2$  and  $H_2@C_{60}$  generate fluctuating fields with corresponding correlation times. Further theoretical studies are required to determine the validity of these speculations.

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