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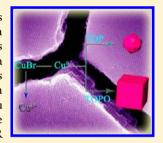


# Shape-Selective Formation of Monodisperse Copper Nanospheres and Nanocubes via Disproportionation Reaction Route and Their **Optical Properties**

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Supporting Information

ABSTRACT: Synthesis of stable and monodisperse Cu nanocrystals of controlled morphology has been a long-standing challenge. In this Article, we report a facile disproportionation reaction approach for the synthesis of such nanocrystals in organic solvents. Either spherical or cubic shapes can be produced, depending on conditions. The typical Cu nanospheres are single crystals with a size of  $23.4 \pm 1.5$  nm, and can self-assemble into three-dimensional (3D) nanocrystal superlattices with a large scale. By manipulating the chemical additives, monodisperse Cu nanocubes with tailorable sizes have also been obtained. The probable formation mechanism of these Cu nanocrystals is discussed. The narrow size distribution results in strong surface plasmon resonance (SPR) peaks even though the resonance is located in the interband transition region. Double SPR peaks are observed in the extinction spectra for the Cu nanocubes with relative large sizes.



Theoretical simulation of the extinction spectra indicates that the SPR band located at longer wavelengths is caused by assembly of Cu nanocubes into more complex structures. The synthesis procedure that we report here is expected to foster systematic investigations on the physical properties and self-assembly of Cu nanocrystals with shape and size singularity for their potential applications in photonic and nanoelectronic devices.

#### **■ INTRODUCTION**

Group IB metallic nanostructures, such as Cu, Ag, and Au nanocrystals, are of particular interest as an important class of photonic components because of the interactions between their free electrons and light. This interaction results in surface plasmon resonances (SPR) at specific wavelengths. The unique optical properties of these nanocrystals, including the SPR effect and its associated light scattering, have enabled them to be widely used. Applications include enhancing light trapping and absorption for improving the performance of solar cells, 1,2 providing substrates for the surface-enhanced Raman spectroscopy of trace chemical species,<sup>3,4</sup> and providing ultrasensitive refractometric biosensors.5

It is well-known that the shape and size play important roles in determining the number, position, and intensity of SPR modes.<sup>6-9</sup> Moreover, a narrow size distribution is vital to induce a strong SPR signal at a specific wavelength. In the past decade, appreciable achievements have been made in the shapecontrolled synthesis of noble metal nanocrystals, such as Au and Ag, with various morphologies.<sup>7,10,11</sup> However, comparable success with Cu has been elusive.<sup>12</sup> Reasons for this include the relatively poorer oxidation resistance of Cu nanostructures and the lack of shape control.

Nevertheless, given that Cu has a relatively low cost compared to noble metals, it may be useful to develop a facile strategy to prepare Cu nanocrystals with controllable shape and size distribution, as well as a better understanding of the role of shape in their optical responses. About one decade ago, Pileni and co-workers had concentrated their efforts on understanding the growth mechanism of Cu nanocrystals of differing shapes <sup>1</sup> as well as tuning their optical properties. 14 Zhong et al. also tried to control the shape and size of Cu nanocrystals in organic solvents with molecular capping agents and reaction temperature. 15 Pastoriza-Santos et al. synthesized single-crystalline Cu nanoplates through reduction of copper salts by hydrazine with polyvinylpyrrolidone as a stabilizer. An intense in-plane SPR band was observed in these Cu nanoplates. 16 Recently, Xia et al. were able to synthesize Cu nanocrystals with controllable shapes.<sup>17</sup> Progress has also been made in fabrication of Cu nanowires, <sup>18–20</sup> nanocubes, <sup>17,20</sup> and nanospheres.<sup>21,22</sup> Nevertheless, preparation of Cu nanocrystals with monodisperse size distribution is still a difficult task. The conflict stems partially from the fact that strong reducing agents or high temperatures are normally required to reduce cupric ions into zerovalent atoms because of the relatively low reduction potential of Cu<sup>2+</sup>. However, under such drastic reaction conditions, control of the shape of nanocrystals may be lost due to aggregative growth and Ostwald ripening. <sup>23,24</sup> Another challenge is the propensity

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of Cu nanocrystals to oxidize. As a consequence, in order to synthesize monodisperse Cu nanocrystals in solution, it is necessary to pay attention to the way by which zerovalent Cu atoms are generated. This may be done by modulating the redox potentials of the metal ions as well as by controlling the growth and ripening rate with stabilizing agents.

Previous research normally supposed that surfactants only served as capping agents or adsorbates<sup>25-27</sup> which played important roles in shape-selective synthesis of anisotropic nanocrystals as well as in controlling particle size distribution. However, the possible multiple roles of the surfactant molecules and anions in solvents should not be ignored. For instance, modification of the redox potential of the metal (owing to the coordination effect between metallic ions and surfactant molecules or anions) may have a considerable influence on the synthesized nanocrystals.<sup>27,28</sup> For example, Sun and Peng demonstrated how such complexation could be used to generate single-crystalline Ag nanocubes in a mixed solvent of octyl ether and oleylamine in the presence of Clions (from dimethyldistearylammonium chloride) at 260  $^{\circ}\text{C.}^{29}$ It should be noted that the shape and size of Ag nanocrystals would be difficult to control in the absence of Cl<sup>-</sup> ions at such a high temperature. Recently, Tuan and co-workers reported the synthesis of Cu nanocubes with an average edge length of 75.7 nm using an injection approach that involved the fast reduction of CuCl. 30 A similar result was also demonstrated in our recent research on the shape-controlled synthesis of CuNi alloyed nanoplates<sup>31</sup> and Cu nanowires.<sup>20</sup> The strong coordination effect between Cu<sup>2+</sup> and Cl<sup>-</sup> ions prevented Cu<sup>2+</sup> ions from being reduced by oleylamine. Therefore, Ni<sup>2+</sup> ions were introduced into the reaction to act as catalysts to promote the reduction of Cu<sup>2+</sup>, which enabled a controlled reduction rate of Cu<sup>2+</sup> ions through controlling the reduction rate of Ni<sup>2+</sup> ions by selective addition of surfactant molecules. This phenomenon provides a possible route to tune the nucleation rates of metal nanocrystals by modifying the coordination agents and surface adsorbates in solvents so as to keep the formation of Cu nanocrystals under kinetic control.

In the present work, we exploit these insights to demonstrate a facile disproportionation reaction route in organic solvents for the synthesis of Cu nanocrystals with a tight size distribution. In our syntheses, no Ni2+ species are required to be added to promote the reduction of Cu<sup>2+</sup>. In addition, CuBr is employed as Cu precursor and the reaction temperature (210-260 °C) is much lower than that (330 °C) reported in a previous study. To the best of our knowledge, this is the first time that monodisperse Cu nanocrystals with both spherical and cubic morphologies have been achieved via such a simple one-pot disproportionation reaction. The growth mechanism of asprepared Cu nanocrystals has also been elucidated. We will show that bromide anions along with other capping agents are indispensable for the shape-selective formation of Cu nanocrystals with such a narrow size distribution. Finally, theoretical simulations of the extinction spectra of the as-prepared Cu nanocrystals were undertaken to better understand their optical properties.

## **■ EXPERIMENTAL SECTION**

**Chemicals.** Oleylamine (80–90%) and trioctylphosphine oxide (TOPO, 90%) were purchased from Acros Organics. Trioctylphosphine (TOP, 90%) was purchased from Sigma-Aldrich. Toluene, *n*-hexane, acetone, ethanol, and CuBr were purchased from Sinopharm Chemical Reagent Co., Ltd. All of

these chemicals were used as received without any further purification. It should be noted that it is very important to keep TOP from being oxidized in order to prepare monodispersed Cu nanospheres.

Synthesis of Monodisperse Cu Nanospheres. All of the syntheses were carried out with a protective atmosphere of argon. Typically, oleylamine (7 mL) and CuBr (0.6 mmol) were mixed in a three-necked flask by strong magnetic stirring at 80 °C for 15 min, then TOP (1 mmol) was injected into the obtained green solution, which immediately became colorless. After an additional stirring at 80 °C for 5 min, the mixed solution was heated up to 260 °C quickly and kept at this temperature for 3 h before cooling down to room temperature naturally. Excess hexane was added to the bright reddish solution, and the Cu nanospheres were precipitated out by centrifugation. The Cu nanospheres were purified by three rounds of centrifugation/redispersion in a mixed solution of hexane and acetone prior to being finally dispersed into toluene by ultrasonication for 30 min.

Self-Assembly of Monodisperse Cu Nanospheres into 3D Superlattices. A silicon wafer (5 mm  $\times$  5 mm) which had been rinsed with hydrofluoric acid solution (5%) and acetone continuously was placed on the bottom of a penicillin bottle (5 mL). Toluene solution (2 mL) containing Cu nanospheres was added into the penicillin bottle. After that, the penicillin bottle was kept in a thermostat until all of the toluene solution had evaporated.

**Synthesis of Monodisperse Cu Nanocubes.** All of the syntheses were carried out with a protective atmosphere of argon. Typically, CuBr (0.6 mmol) and TOPO (1.5 mmol) were dissolved into oleylamine (7 mL) in a three-necked flask by strong magnetic stirring at 80 °C for 15 min. Then the resulting solution was heated up to 260 °C quickly and refluxed at this temperature for 1 h before cooling down to room temperature naturally. The remaining steps are the same as those for preparing Cu nanospheres.

In order to obtain Cu nanocubes with a larger edge length (e.g., 76.4 nm), the content of TOPO was reduced to 0.5 mmol, while the aging time was increased to 2.5 h. In contrast, Cu nanocubes with an average edge length of about 31.5 nm were synthesized by decreasing the reaction temperature to 220 °C and aging at this temperature for 1 h. For the synthesis of Cu nanocubes with an average edge length of 24.1 nm, CuBr (0.6 mmol) and TOPO (5 mmol) were dissolved into 2 mL of oleylamine in a three-necked flask by strong magnetic stirring at 80 °C for 15 min. Then, the resulting solution was heated up to 210 °C quickly and refluxed at this temperature for 1 h before cooling down to room temperature naturally. Similarly, monodisperse Cu nanocubes with an average edge length of about 26.8 nm were achieved by prolonging the reaction time to 3 h.

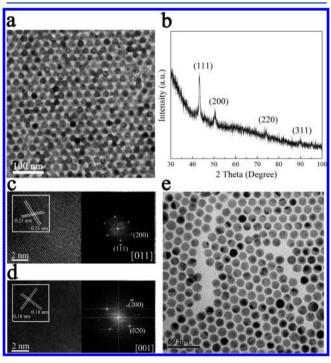
Characterization. The samples for transmission electron microscopy (TEM) analyses were prepared by dropping the particle suspensions in toluene onto a Mo grid coated with carbon film before drying at room temperature under ambient conditions. TEM images and selected area electron diffraction (SAED) patterns were collected on a JEOL JEM-2100 transmission electron microscope operating at 200 kV, while high-resolution TEM (HRTEM) images were collected on a FEI TECNAI F30 transmission electron microscope operating at 300 kV. The scanning electron microscopy (SEM) images were collected on a LEO 1530 scanning electron microscope operating at 15 kV. X-ray diffraction (XRD) measurements

were taken on a Panalytical X'pert PRO diffractometer using Cu K $\alpha$  radiation, operating at 40 kV and 30 mA. Visible-near-infrared (NIR) extinction spectra were obtained at room temperature with a Shimadzu UV-2550 ultraviolet—visible spectrophotometer. Auger spectra were acquired with a PHI-610 scanning Auger microprobe. Fourier transform infrared (FTIR) spectra were obtained with a Thermo Fisher Scientific IS10 FTIR spectrometer.

**Simulations.** The optical extinction properties of single Cu nanocubes and their aggregates were simulated using the discrete dipole approximation (DDA) as implemented in the DDSCAT code of Draine and Flatau. Version 7.1 of the code was used. In this scheme, the shape is represented by an array of polarizable dipoles and the interaction with an incident periodic wave solved numerically. Metallic targets, in which the complex refractive index can have a large magnitude, require very long computation times and closely spaced dipoles in order to obtain converged results. The dielectric data for Cu and toluene were obtained from the literature. <sup>33–35</sup>

#### ■ RESULTS AND DISCUSSION

**Monodisperse Cu Nanospheres.** Monodisperse Cu nanospheres have been synthesized using TOP as a chemical additive. Figure 1a shows the TEM image of the two-

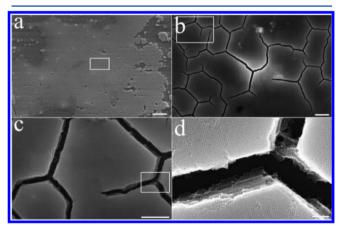


**Figure 1.** (a) TEM image and (b) XRD pattern of the monodisperse Cu nanospheres with an average diameter of 23.4 nm. HRTEM images of a single Cu nanosphere viewed along (c) [011] and (d) [001] zone axes. The right images in (c) and (d) are corresponding FFT patterns acquired from the square marking areas. (e) TEM image of monodisperse Cu nanospheres with an average diameter of 12.6 nm.

dimensional stacking pattern of two to three layers of the assynthesized Cu nanospheres assembling on the amorphous carbon film. The size distribution histogram from the TEM image (Supporting Information Figures S1 and S2) reveals that the spheres have a relatively tight distribution of diameters with an average diameter of 23.4  $\pm$  1.5 nm. Figure 1b displays the XRD pattern of the products prepared in the typical procedure.

The diffraction peaks located at 43.3°, 50.4°, 74.1°, and 90.1° can be assigned to the {111}, {200}, {220}, and {311} planes, respectively, of face-centered-cubic (fcc) Cu (JCPDS #03-1018). No diffraction peaks belonging to Cu<sub>2</sub>O are observed, revealing the pure-phase nature and excellent oxidation resistance of the products, which may be attributed to the protection of surface adsorbed Br ions and surfactants. HRTEM images from a single Cu nanosphere observed along the [011] and [001] zone axes along with the fast Fourier transform (FFT) patterns are shown in Figure 1c and d, respectively. Well-aligned lattices without any twin structures indicate that the Cu nanosphere is a single crystal covered by a mixture of {011}, {111}, and {100} facets on the surface (also see Figure S3). It should be note that it is not easy to tailor the size of monodisperse Cu nanospheres using the recipe described in the Experimental Section. In order to obtain Cu nanospheres with a smaller diameter, CuBr was replaced by copper acetylacetonate as a precursor, while TOPO was used to replace TOP (please find the synthetic details in the Supporting Information). As a result, monodisperse Cu nanospheres with an average diameter of 12.6 nm were obtained (Figures 1e and

Owing to their very narrow size distribution and excellent dispersity in toluene solution, three-dimensional (3D) superlattices of multilayer Cu nanospheres deposited on silicon wafer are acquirable during a slow evaporation of toluene in a thermostat. The self-assembly procedure is schematically illustrated in Figure S5. As can be seen from Figure S6, a lustrous liquid film on the surface of Cu nanosphere colloid forms and it becomes brighter and brighter as the evaporation of toluene proceeds. This is due to the assembling of Cu nanospheres on the air/liquid interface, which will finally land on the silicon substrate. Figure 2 displays the SEM images of



**Figure 2.** SEM images of the 3D superlattices composed of 23.4 nm Cu nanospheres with different magnifications. Images (b), (c), and (d) are taken from the square frame shown in image (a). The bar scales are (a) 20  $\mu$ m, (b) 1  $\mu$ m, (c) 1  $\mu$ m, and (d) 200 nm.

3D superlattice formed by the monodisperse Cu nanospheres with an average diameter of 23.4 nm with different magnifications, demonstrating that the as-synthesized monodispersed Cu nanospheres can self-assemble into a millimeter scaled 3D superlattices. The thickness of the 3D superlattice of Cu nanospheres is estimated to be more than 20 layers through the observation on the broken edges or "cracks" which result from the surface tension stress created during the evaporation of solvent.

**Monodisperse Cu Nanocubes.** Monodisperse Cu nanocubes can be obtained if TOP is replaced by TOPO and the reaction time is shortened to 1 h in the typical procedure for the synthesis of monodisperse Cu nanospheres. A representative TEM image of the as-synthesized Cu nanocubes is presented in Figure 3a, where only cubes are observed,

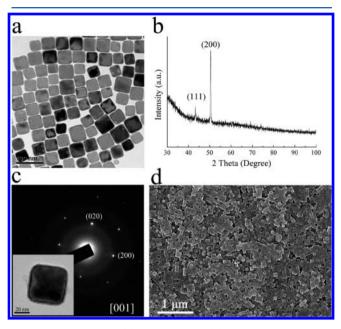


Figure 3. TEM image (a), XRD pattern (b), SAED pattern from a single Cu nanocube (inset) (c), and SEM image (d) of the Cu nanocubes with an average edge length of 46.8 nm.

indicating a good uniformity of shape. The Cu nanocubes have a narrow size distribution with an average edge length of 46.8 nm (see size distribution histogram in Figure S7). Figure 3b shows the XRD pattern of the Cu nanocubes, in which the strongest diffraction peak is (200). This is different from the conventional power XRD pattern of fcc Cu (JCPDS #03-1018). It should be noted that the XRD sample was prepared by dropwise adding the Cu nanocube solution onto a glass slide followed by slow evaporation of toluene in a thermostat. As a result, almost all of the Cu nanocubes have a preferred orientation with {100} facets parallel to the substrate. Tilting experiments were carried out to determine the structure of Cu nanocubes. The diffraction patterns acquired from all of the examined particles can be assigned to the [001] zone axis without exception (Figure 3c), which reveals that the asprepared Cu nanocubes are single crystals covered by six {100} facets. Figure 3d displays the SEM image of the Cu nanocubes with 2-4 layer stacking, where partially ordered structures in which Cu nanocubes are positioned in face-to-face configuration are observed.

The amount of TOPO and the reaction temperature were found to play important roles in determining the final size of Cu nanocubes. For example, it was straightforward to increase the edge length of the Cu nanocubes to 76.4 nm (Figure S8) simply by reducing the content of TOPO to 0.5 mmol in the typical synthetic procedure, although a small number of Cu nanowires were also present in the product (Figure S9). On the contrary, monodisperse Cu nanocubes with an average edge length of 24.1 nm were obtained (Figures 4a and S10) when the amount of TOPO was increased to 5 mmol and the reaction temperature was reduced to 210 °C. In the SAED

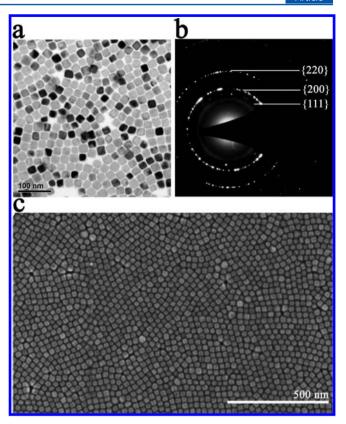
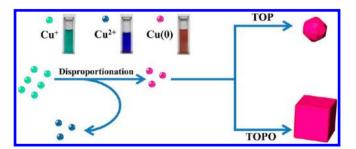


Figure 4. TEM image (a), SAED pattern (b), and SEM image (c) of the Cu nanocubes with an average edge length of 24.1 nm.

pattern acquired from the area shown in Figure 4b, the diffraction ring owing to {111} facets almost disappears, while the diffraction ring owing to {200} facets becomes stronger than all of the others, revealing the preferred orientation of the sample. Because these Cu nanocubes have a monodisperse size distribution and are well dispersed in toluene solution, they can be easily assembled onto silicon wafer (Figure 4c) by the method illustrated in Figure S5. Cu nanocubes with other different sizes are shown in Figures S11 and S12. The above results indicate that not only the shape but also the size of Cu nanocrystals is tunable by adjusting the reaction condition in this facile one-pot procedure.

**Discussion of the Formation Mechanism.** The formation mechanism of the as-prepared Cu nanocrystals was further investigated, and the overall process is illustrated in Scheme 1. CuBr, which is insoluble in water, was dissolved into organic solution by forming coordination complexes with oleylamine. The disproportionation reaction can be determined

Scheme 1. Schematic Illustration of the Shape-Selective Formation of Cu Nanospheres and Nanocubes via the Disproportionation Reaction Route



by the color change of the mixed solution. Green solution was formed by heating the mixture of oleylamine and CuBr at 80 °C, which can be attributed to the presence of Cu<sup>+</sup>-oleylamine complexes. Figure 5 compares the FTIR spectra of oleylamine

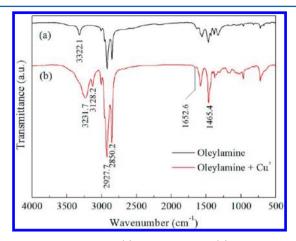
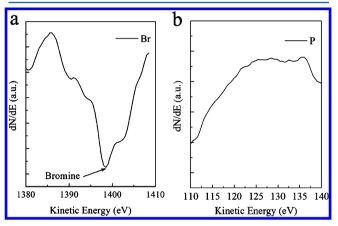


Figure 5. FTIR spectra of (a) oleylamine and (b)  $\mathrm{Cu^+-}$  oleylamine complexes.

and Cu<sup>+</sup>-oleylamine complexes. The stretching vibrations of N-H and C-C can be detected at about 3322.1 and 1652.6 cm<sup>-1</sup>, respectively (Figure 5, spectrum a), which are the characteristic peaks of oleylamine.<sup>36</sup> In the FTIR spectrum of Cu<sup>+</sup>-oleylamine (Figure 5, spectrum b), the N-H stretching region shifts to 3230.5 cm<sup>-1</sup> and an obvious broadening in the bandwidth is observed, which may be due to the coordination of Cu<sup>+</sup> with amine group of oleylamine.<sup>37</sup> On the other hand, there was blue solution observed in the mother liquid after reaction, indicating a remnant of bivalent copper ions. For comparison, when CuBr2 was used as precursor instead of CuBr, the particles produced had micrometer-scale size and chaotic shape after aging at 300 °C for 6 h. This is because a high temperature is necessary to reduce Cu<sup>2+</sup> into Cu<sup>0</sup> directly, but under such drastic reaction conditions the Cu crystals grow rapidly and chaotically, as mentioned previously. The disproportionation reaction enables the generation of zerovalent Cu atoms at a relative low temperature so as to keep the growth of Cu nanocrystals under control of the capping agents.

It was found that both TOP and TOPO could promote the disproportionation reaction; in other words, the reaction temperature was brought down by the addition of TOP or TOPO. Visible evidence of the reaction was obtained when the color of the mixed solution began to turn to red, which indicated the formation of Cu clusters at 190 and 240 °C in the presence of TOPO and TOP, respectively. In contrast, this critical temperature increased to about 260 °C in the absence of either TOP or TOPO. The difference in the critical reaction temperature between TOP and TOPO also revealed the stronger coordination effect in the Cu+-TOP complexes compared to Cu<sup>+</sup>-TOPO complexes. The coordination interaction between metallic ions and ligands can be well explained by the HSAB (Hard-Soft-Acid-Base) theory. It is well-known that Cu<sup>+</sup> is a soft acid, while TOP (in the type of R<sub>3</sub>P) is a soft base. However, in TOPO, oxygen (which is a well-known hard base) will act as the coordination atom instead of phosphorus. Therefore, it is reasonable to deduce that Cu<sup>+</sup>-TOP is more stable than Cu<sup>+</sup>-TOPO complexes. The shapeselective effect of TOP and TOPO (illustrated in Scheme 1) is

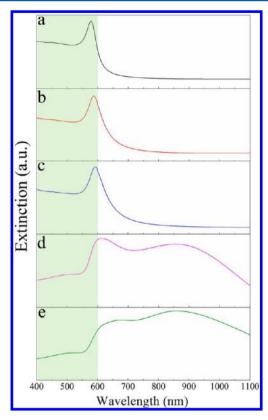
based on their different coordination effects to Cu<sup>+</sup>. <sup>15</sup> First, the much stronger coordination of TOP to Cu+ prevents the metallic ions from being fast reduced to Cu(0), which seems to be necessary to accumulate enough reactant species to reach the critical nucleation concentration to achieve instantaneous nucleation at a higher temperature (260 °C).<sup>38</sup> The high reaction temperature avails the formation of isotropous crystals and weakens the selective adsorption of Br-, which make the seeds grow equally in all directions into spherical-shaped crystals. Second, the synergistic effect of halogen anions and TOPO is necessary for the formation of a cuboidal geometry. Once TOP was replaced by TOPO, which possessed weaker coordination effect to Cu<sup>+</sup>, the nucleation process would happen at a relatively lower temperature (170 °C). The lower reaction temperature enables a slow growth rate. The {100} facets of the Cu crystal seeds would be selectively capped by Br ions. From the Auger spectra acquired from the surface of Cu nanocubes, an obvious peak at 1398.5 eV belonging to the Br LMM transition (Figure 6a) is observed, indicating the



**Figure 6.** Auger spectra from the surface of Cu nanocubes: (a) Br LMM region and (b) P LMM region.

binding of Br<sup>-</sup> ions on Cu surfaces. However, in the Auger spectrum (Figure 6b) that covers P LMM region, no peak is observed. A similar phenomenon is also found from the Auger spectra acquired from Cu nanospheres. This reveals that bromine ions have a stronger binding affinity than that of TOP or TOPO. The face-selective adsorption of halide anions has been well discussed in former reports. <sup>20,27</sup> As a result, the other facets grew at a faster rate than the {100} facets. This kinetic difference led to the formation of Cu nanocubes instead of nanospheres.

**Optical Properties.** By controlling the geometry of noble metal nanocrystals, such as Ag<sup>11,39</sup> and Au,<sup>40–42</sup> researchers have successfully tailored the optical responses of metallic nanocrystals and set the SPR at a specific wavelength for particular application. Here we study the shape-dependent optical characteristics of Cu nanocrystals. Among group IB metallic nanospheres, Cu displays the largest onset wavelength of interband transitions in the extinction spectrum<sup>43</sup> (about 600 nm, the green shadow area in Figure 7), which results in the weakest SPR intensity because of their location in the interband transition region. However, as can be seen in the extinction spectrum of the monodisperse Cu nanoparticles suspending in toluene solution (Figure 7a), there is still a strong and narrow SPR peak at 576 nm overlapping on the background absorption induced by the interband transitions. Given the monodispersity



**Figure 7.** Vis-NIR extinction spectra of Cu nanospheres with an average diameter of 23.4 nm (a) and nanocubes with an average edge length of 24.1 nm (b), 26.4 nm (c), 30.7 nm (d), and 46.8 nm (e). All of the spectra are acquired from the Cu nanocrystals suspended in toluene.

of the as-synthesized Cu nanospheres, the coherent oscillation of the free conduction electrons could be excited by light within a very narrow wavelength region, which leads to the enhanced SPR absorption. According to Mie scattering theory, there should be a single SPR peak in the extinction spectrum of Cu nanospheres because of their spherically symmetric geometry, which is identical with our spectroscopic observation shown in Figure 7a.

With respect to the extinction spectra of Cu nanocubes (Figure 7b-e), there is a tiny red shift from 576 to 585 nm as the edge length of the nanocubes increases from 24.1 to 26.4 nm (see Figure S11 for TEM image). Interestingly, when the edge length of Cu nanocubes reaches 30.7 nm (see Figure S12 for TEM image), two new SPR peaks appear at about 610 nm and about 855 nm. In the case of even larger cubes, such as the 46.8 nm cubes, these two new peaks are at about 670 and 858 nm. According to former reports, 17 the red shift may be due to the larger dimension of Cu nanocubes compared with Cu nanospheres. Similar phenomena have been observed in gold nanoparticles with well-controlled sizes. 45,46 Furthermore, the double-peaked feature in the spectrum of the 30.7 and 46.8 nm cubes is evidently caused by their aggregation into more complex structures. From the TEM images (Figures 4 and S11), it can be seen that Cu nanocubes with a smaller sizes (24.1 and 26.4 nm) do not contact by each other due to the separation of capping agents. However, the van der Waals attractions between Cu nanocubes with larger sizes will be relatively stronger, and in this case the capping molecules may exert insufficient repulsion. As a result, the formation of aggregates (see Figure S13) such as dimers and trimers may lead to SPR peaks at longer wavelengths.

To better understand the optical properties of the asprepared Cu nanocrystals, we performed theoretical simulations of their extinction spectra, for single cubes, linear chains of up to five cubes, and four cubes in a square arrangement (Figure 8). These shapes were chosen because they had been observed

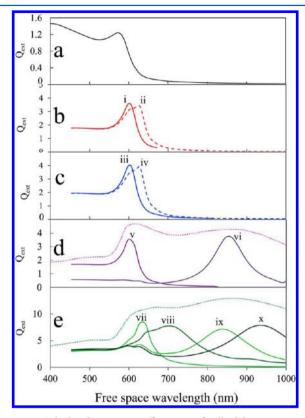


Figure 8. Calculated extinction efficiencies of colloidal suspensions of Cu nanospheres and nanocubes in toluene: (a) isolated 23.4 nm diameter nanospheres; (b) isolated 24.1 nm nanocubes, (i) = rounded edges, (ii) = sharp edges; (c) isolated 26.1 nm nanocubes, (iii) = rounded edges, (iv) = sharp edges; (d) mixture of isolated 30.7 nm cubes and their linear trimers, (v) = single cubes, (vi) = three cubes in a row (experimental data is shown for comparison as a dotted line); (e) calculated extinction spectra of 46.8 nm cubes in various configurations, (vii) = single cubes, (viii) = four cubes in a  $2 \times 2$  arrangement, (ix) = three cubes in a row, (x) = four cubes in row. (Experimental data is shown for comparison as a dotted line.)

in TEM samples of these cubes. The transverse plasmon resonance peaks of the simulation of the single cubes, Figure 8a, are a reasonable match to the higher energy peaks of the experimental data with any residual small discrepancies between the two probably due to a lack of geometric perfection in the experimental nanoparticles. The simulated peaks for single cubes of maximum theoretical edge sharpness are somewhat triangular in shape because the simulation has generated some higher order resonances of the transverse resonance peak. This phenomenon has been observed before in Ag cuboids. Therefore, it was necessary to generate single cubes with slightly rounded edges to better match the experimental data.

As mentioned, experimental samples with cube sizes of 30.7 and 46.8 nm also showed a second broad resonance centered on about 860 nm. The simulations suggest that even two cubes

in a side-by-side cannot generate such a deeply red-shifted resonance, and it appears necessary to invoke the existence of aggregates containing at least three, four and five cubes to explain them. Spectra very similar in nature to the measured ones could be obtained by producing an ensemble average of such aggregated cubes, with the broader peaks of the experimental data being probably due to the presence of a small spread in cube dimensions. A breakdown of the shapes necessary to match the experimental spectrum is shown in Figure 8e for the case of the 46.8 nm cubes. A similar breakdown (not shown) applied in the case of the 30.7 nm cubes. In these simulations an interparticle gap about 1.9 nm is necessary in order to provide the correct peak positions.

The numerical simulations show that the electric field intensity on the individual Cu nanospheres is relatively low compared to that of the nanocubes, Figure 9. The cubes are not

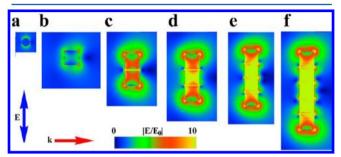


Figure 9. Calculated electric field intensities (shown as the norm, ||E||) at the plasmon resonance frequencies of selected Cu nanocrystals or nanocrystal aggregates in toluene. Crystals and fields are shown at the same scales. (a) 23.4 nm diameter sphere irradiated at 573 nm, (b) single 46.8 nm cube (edges have been rounded) irradiated at 613 nm, (c) longitudinal plasmon resonance of two 46.8 nm cubes in a line irradiated at 647 nm (d) longitudinal plasmon resonance of three 46.8 nm cubes in a line irradiated at 709 nm, (e) longitudinal plasmon resonance of four 46.8 nm cubes in a line irradiated at 934 nm, and (f) longitudinal plasmon resonance of five 46.8 nm cubes in a line irradiated at 1022 nm. (A separation gap of 1.9 nm was used for the multicube geometries.)

only better nanoantennas for geometric reasons, but, importantly, their characteristic resonance frequencies are also well away from the damping effect of copper's interband transitions in the mid to lower-wavelengths of visible light. Given that aggregation of the larger cubes occurred, it was interesting to also explore the electric field intensities of the assembled structures of two, three, four, and five cubes. The increase in overall electric field intensity as the longitudinal resonance of the aggregates is red-shifted is quite apparent in Figure 9. Notice also that the fields at the sharp edges and in the gaps of these structures are very high relative to those of an isolated sphere or cube. Therefore, the aggregated structures may be more useful for some applications, such as SERS, than unaggregated particles.

#### CONCLUSION

In summary, we have demonstrated a facile one-pot method that is based on a disproportionation reaction in oleylamine solvent. This technique provides a shape-selective synthesis of Cu nanocrystals with spherical and cubic morphology as well as with very narrow size distribution. Owing to the relative mild reaction conditions and capping agents, the growth of Cu nanocrystals was kept under kinetic control. Monodisperse Cu

nanospheres were generated when TOP was introduced as a capping agent, while TOPO could be used for the synthesis of monodisperse Cu nanocubes with tunable sizes. Furthermore, the optical properties of the as-prepared Cu nanocrystals were investigated both experimentally and theoretically. A strong and narrow SPR peak at 576 nm was observed for the Cu nanospheres, while the SPR peaks of Cu nanocubes exhibited a red-shift that increased with their sizes. Theoretical simulations on the extinction spectra agree well with the observed spectra. While the smaller cubes evidently stayed dispersed as single entities, the optical properties of suspensions of the cubes with larger sizes (30.7 and 46.8 nm) suggested that they had aggregated into triple, quadruple or even pentuple structures. The reported simple strategy for the synthesis of Cu nanocrystal in this study not only provides a convenient access to the further investigation of their physical properties due to their shape and reasonable size singularity, but also may promote the application of Cu nanocrystals in photonic and nanoelectronic devices due to their comparatively lower cost than Au or Ag.

#### ASSOCIATED CONTENT

## **S** Supporting Information

Synthetic details, TEM and HRTEM images, size distribution histograms, and self-assembly scheme. This material is available free of charge via the Internet at http://pubs.acs.org.

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#### Notes

The authors declare no competing financial interest.

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