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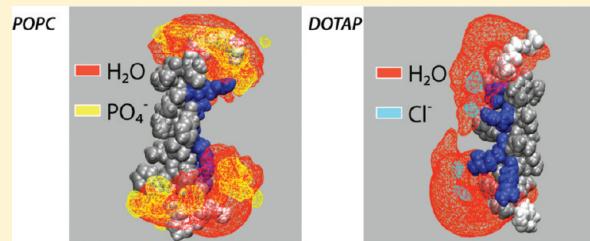
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Structural Dynamics of the S4 Voltage-Sensor Helix in Lipid Bilayers Lacking Phosphate Groups

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ABSTRACT: Voltage-dependent K⁺ (Kv) channels require lipid phosphates for functioning. The S4 helix, which carries the gating charges in the voltage-sensing domain (VSD), inserts into membranes while being stabilized by a protein–lipid interface in which lipid phosphates play an essential role. To examine the physical basis of the protein–lipid interface in the absence of lipid phosphates, we performed molecular dynamics (MD) simulations of a KvAP S4 variant (S4mut) in bilayers with and without lipid phosphates. We find that, in dioleoyltrimethylammoniumpropane (DOTAP) bilayers lacking lipid phosphates, the gating charges are solvated by anionic counterions and, hence, lack the bilayer support provided by phosphate-containing palmitoyloleoylglycerophosphocholine (POPC) bilayers. The result is a water-permeable bilayer with significantly smaller deformations around the peptide. Together, these results provide an explanation for the nonfunctionality of VSDs in terms of a destabilizing protein–lipid interface.



1. INTRODUCTION

The lipid bilayer of the plasma membrane defines the boundary of biological cells and compartmentalizes them by providing a permeability barrier. Membrane proteins embedded in bilayers allow controlled communication with the surrounding environment. The composition of the bilayer is not only important for achieving a proper integration of these membrane proteins, it is also vital for their function. The critical dependence of proteins on the chemical nature of the surrounding lipid bilayer suggests that the two have coevolved.¹ A striking example is the absolute requirement for lipid phosphate groups in the function of voltage-dependent K⁺ (Kv) channels.^{2–4} Molecular dynamics (MD) simulations have identified, and described at atomic detail, a carefully orchestrated protein–lipid interface surrounding the channel gating charges.⁵ By removing negatively charged lipid phosphates, the key players in connecting the protein to the bilayer, an entirely new set of protein–lipid interactions is bound to arise. Here, as a first step toward a detailed understanding as to why Kv channels do not function in the absence of lipid phosphates, we examined the interaction of an arginine-rich S4 helix variant with counterions and water molecules in bilayers lacking phosphate groups, by means of MD simulations.

Kv channels help generate action potentials along the cell membrane of excitable neurons by opening and closing their pore domains in response to changes in the transmembrane electrical potential, thereby controlling the flow of ions through the membrane.^{6–9} Each subunit of these homotetrameric channels contains a voltage-sensing domain (VSD), composed of transmembrane (TM) helices S1–S4, and a pore domain (PD), which includes TM helices S5–S6. The highly charged, arginine-rich S4 helix of the VSD enables Kv channels to detect and

respond to changes in TM potential.^{10–12} Voltage-induced movements of the VSDs cause the PDs to open or close through mechanical coupling provided by the so-called S4–S5 linker helix. Except for this linkage, structural and functional data^{13–18} have revealed that the VSDs move at the protein–lipid interface largely independently of the pore domain. However, the energetic penalty for moving charged Arg residues inside the hydrophobic core of a lipid bilayer is expected to be enormous,^{19,20} and this has fueled a debate regarding the nature of the gating mechanism of Kv channels.²¹ Regardless of the true trajectory of the gating charges within the bilayer, the observation that Kv channels show no activity in dioleoyltrimethylammoniumpropane (DOTAP) bilayers lacking lipid phosphates² indicates that the interplay between the lipid bilayer and the S4 helix lies at the heart of the gating mechanism.

The isolated S4 helix has been reported to adopt a TM configuration in lipid bilayers from mixtures of a synthetic S4 peptide and phospholipids.^{22,23} Also, the propensity for membrane insertion of the S4 helix was addressed in a key experiment showing that a model helix based on the S4 segment (GGPG-LGLFRLVRLLRFLRILLII-GPGG) partitions within the endoplasmic reticulum (ER) membrane with a low apparent free energy penalty.²⁴ A mutated form of S4 (S4mut), in which two Arg residues (R15 and R18) were moved one step closer to the C-terminus, was shown to insert even more efficiently. S4mut was chosen for a computational investigation of the solvation shell that accommodates an S4-like transmembrane helix in lipid

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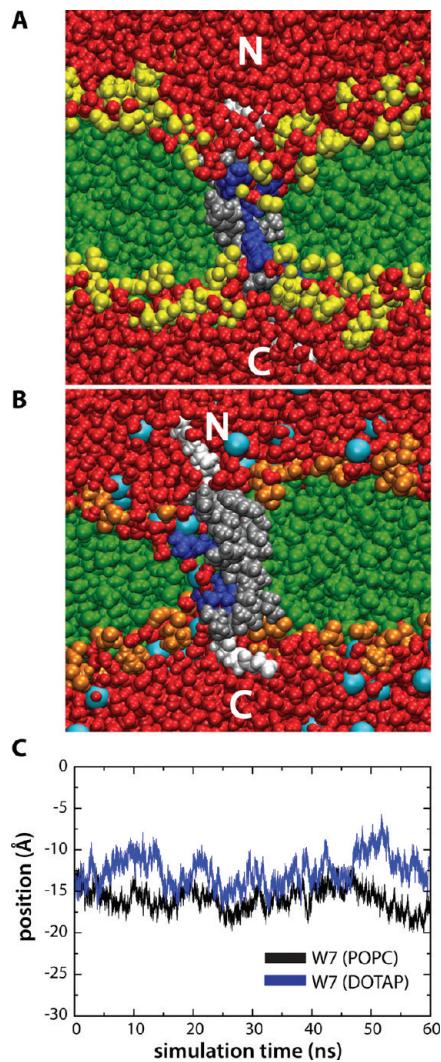


Figure 1. Snapshots of S4mut in POPC and DOTAP bilayers. (A) Snapshot from MD simulation in excess water of GPGG...GGPG-flanked (white) S4mut sequence (silver) carrying four charged Arg residues (blue) residing in a POPC bilayer in which the phosphatidyl-choline head groups are depicted in yellow and the acyl chains in green. (B) Simulation snapshot as in (A), except that the bilayer (DOTAP) lacks lipid phosphates in the head-group region (orange). Chloride counterions are depicted as cyan spheres. (C) The positional evolution of W7 centers of mass (COM) along the bilayer normal (*z*-dimension) from equilibrated POPC (black) and DOTAP (blue) MD trajectories.

bilayer membranes as part of a broader experimental research program.

To gain insights into the protein–lipid interactions that allow a TM configuration of the S4 helix, Freites et al.⁵ performed an MD simulation of S4 inserted across a phospholipid bilayer. The simulation revealed that the thickness of the bilayer immediately surrounding the helix was significantly reduced to provide a hydrogen-bonded network of lipid phosphates and water molecules to solvate the snorkeling Arg residues. A ~10 Å hydrophobic gap was observed to separate the head group–arginine complexes on opposite sides of the bilayer, which suggests that the membrane dielectric barrier is only a fraction of the hydrocarbon core thickness. This structural feature is consistent with experimental²⁴ and simulation results^{25–27} showing that

the energetics of arginine-partitioning in a phospholipid bilayer is strongly dependent on the location of the arginine center of charge. The protein–lipid interactions arising in the single S4 helix system are in agreement with those observed in MD simulations of an isolated VSD^{28,29} and an entire Kv channel^{30–32} in bilayers containing lipid phosphates. However, structural studies have revealed that, although the outermost Arg residues are readily accommodated in a protein–lipid interface, Arg residues located deep within the membrane are shielded from a lipid environment.^{14,17,18} Hence, when drawing an analogy between the protein–lipid interactions of an isolated S4 fragment and those of the entire VSD, one needs to consider the caveat that not all Arg residues in the latter are directly exposed to lipids. Nonetheless, an experimental approach combining MD simulations, solid-state NMR, and neutron diffraction³³ showed that the interactions between the exposed charged Arg residues of the VSD and the surrounding bilayer head-group region were strong enough to produce a local deformation of the bilayer, ensuring the stabilization of these charges by means of a water defect. Moreover, it has recently been shown that the charge-containing segments of the VSD interact more strongly with the surrounding bilayer than any other segment of the VSD,³⁴ indicating a crucial role for the solvation shell surrounding the exposed charged Arg residues.

In this paper, we identify and describe, by means of atomistic MD simulations, the solvation shell surrounding an arginine-rich S4 helix variant in bilayers lacking lipid phosphates. We find that chloride counterions serve as ion-pair partners for the charged Arg residues in the absence of lipid phosphates. The counterion–Arg connection translates into an increased mobility of the Arg residues, because they are no longer anchored to the surrounding lipid bilayer. In addition to the chloride counterions, water molecules are also closely associated with the Arg residues, which can now enter deeply into the hydrophobic core of the bilayer, ultimately forming a continuous water wire leading to a collapse of the bilayer integrity. In phosphate-containing bilayers, on the other hand, the Arg–phosphate interaction induces a deformation of the bilayer as it accommodates itself to anchor the peptide. The nature of the solvation shell in the vicinity of the Arg residues is such that water is prevented from entering the hydrophobic gap, thereby stabilizing the protein–lipid complex. As a direct consequence of the failure to form a stable solvation shell around the Arg residues, we observe a significantly smaller deformation of the DOTAP bilayer. Moreover, different rotameric states of the arginines of S4mut were observed to have a strong effect on the secondary structure of the helix.

2. METHODS

2.1. Building the Systems. We performed 100 ns atomistic MD simulations of lipid bilayers containing S4mut (Figure 1). The simulations were carried out in excess water, which denotes a bilayer containing more water molecules than the minimal amount needed for hydration. For a POPC bilayer, a minimal amount of ~9.5 water molecules per lipid is needed for hydration at 93% relative humidity (Mihailescu et al., unpublished and refs 35 and 36). In these simulations, water–lipid ratios of 44.8 and 44.1 were used for POPC and DOTAP, respectively. The sequence of the capped S4mut peptide in our simulations is similar to the wild-type S4:

S4	GGPG-LGLFRLVRLLRFLRLLII-GPGG
S4mut	GGPG-LGWFRLLFRLIRLLII-GPGG

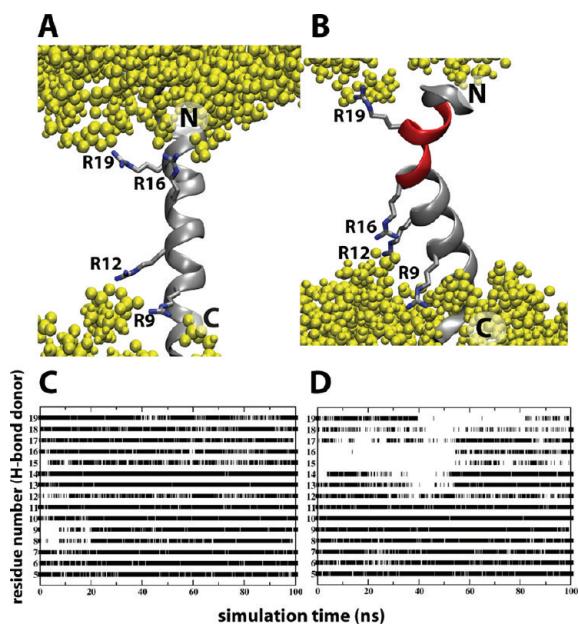


Figure 2. Effect of Arg rotamers on S4mut conformation in POPC bilayers. (A) A representative simulation snapshot of the S4mut helix with a 2-up-2-down distribution of its Arg residues (R9, R12, R16, and R19). Phospholipid head groups are colored yellow. (B) A representative simulation snapshot of the S4mut helix with a 1-up-3-down distribution of its Arg residues (R9, R12, R16, and R19). The snapshot shows the severely kinked region (red) of S4mut. (C) The presence of backbone H-bonds in S4mut from an H-bond donor at position i to an H-bond acceptor at position $i + 4$ for the 1-up-3-down Arg distribution and (D) the 2-up-2-down Arg distribution. The flanking -GPGG- regions were left out of this analysis. The criteria for H-bonds in this analysis were defined by cutoff values 3.5 Å for the donor to acceptor distance and 40° for the donor-H-acceptor supplement angle.

The mutated arginine residues are underlined. In addition, S4mut carries an L7W mutation to facilitate optical detection in ongoing experiments. Hessa et al. found that peptides carrying Trp and Leu at near-terminal equivalent positions inserted equally well into a biological membrane.³⁷ Furthermore, Nishizawa et al.³⁸ could not differentiate the structural dynamics displayed by S4 helices carrying either F8W or F8L or F8V mutations when simulating the behavior of S4-related peptides at the membrane–water interface. The W7L mutation is, therefore, not likely to alter significantly the dynamics displayed by this residue. Indeed, the evolution of the W7 centers of mass in both the POPC and DOTAP simulations indicate no tendency to drift along the bilayer normal (Figure 1C).

To investigate whether two different side-chain configurations, placing one of the S4mut charges at different positions along the bilayer normal, affected the conformational stability of the peptide in the bilayer, we constructed two rotameric states of S4mut, one in which R16 is “up”, toward the C- or “down” toward the N-terminus (Figure 2, A and B). We refer to these configurations as 2-up-2-down and 1-up-3-down, respectively. The 1-up-3-down configuration of S4mut was built as a perfect α -helix using the psfgen plugin of the VMD 1.8.7 software package,³⁹ and the 2-up-2-down configuration was constructed by exploring further rotameric configurations of Arg 16 using the software COOT.⁴⁰ Both configurations of the S4mut helix were then inserted into two separate pre-equilibrated POPC lipid

bilayers by aligning the principal axes of the helix to the bilayer normal, and then making the peptide centers of mass coincident with the bilayer center of mass. Each system was composed of 268 lipid molecules, 12 107 water molecules, 4 counterions, adding up to a total of 72 689 atoms and a water:lipid ratio of 45.2.

The DOTAP chloride salt has been observed to form stable bilayers at different levels of hydration, as we have previously reported in a concerted neutron diffraction and MD simulation study on the structure of pure DOTAP bilayers,⁴¹ the details of which will be provided in a paper that is in preparation. In the present study, S4mut was inserted into a DOTAP bilayer in the 2-up-2-down configuration. The numbers of lipids and water molecules of the DOTAP bilayer system were chosen to have a water:lipid ratio similar to that of the POPC system (44.1). The final DOTAP system contained 284 lipid molecules, 12,515 water molecules, 4 chloride counterions for the protein, and 284 chloride counterions for the lipid bilayer, adding up to a total of 74 347 atoms.

2.2. MD Simulations. Each system was relaxed using a 2500-step conjugate-gradient energy minimization followed by eight 100 ps simulation runs at constant temperature (300 K) and volume during which the peptide was progressively released from its initial configuration using harmonic restraints. The subsequent production simulations were carried out for 100 ns at constant temperature (300 K) and constant pressure (1 atm). All molecular dynamics simulations were run with the NAMD 2.7 software package.^{42,43} The CHARMM22 and CHARMM36 force fields^{44,45} were used for the peptide and lipids, respectively, and the TIP3P model⁴³ was used for the water molecules. A time step of 1 fs was used to integrate the equations of motion, and a reversible multiple-time-step algorithm⁴⁶ of 4 fs was used for the electrostatic forces and 2 fs for short-range, nonbonded forces. The smooth particle mesh Ewald method^{47,48} was used to calculate electrostatic interactions. The short-range interactions were cut off at 12 Å. All bond lengths involving hydrogen atoms were held fixed using the SHAKE⁴⁹ and SETTLE⁵⁰ algorithms. A Langevin dynamics scheme was used for thermostating. Nosé–Hoover–Langevin pistons were used for pressure control.^{51,52} Molecular graphics and simulation analyses were generated with the VMD 1.8.7 software package.³⁹

The evolution of the peptide tilt angle with respect to the bilayer normal, as well as the position of the peptide center of mass along the bilayer normal relative to the bilayer, indicated that all three systems (1-up-3-down and 2-up-2-down distribution of S4mut in POPC, and the 2-up-2-down distribution in DOTAP) were equilibrated after 40 ns (Figure 3). All subsequent analyses were, therefore, based on the last 60 ns of each trajectory.

3. RESULTS AND DISCUSSION

3.1. Integrity of the Lipid Bilayer. In a palmitoyloleoylglycerophosphocholine (POPC) bilayer, the 2-up-2-down S4mut helix adopted a TM configuration and the simulation revealed an H-bonding pattern where lipid phosphates, water molecules and, to a lesser extent, lipid carbonyls alternated in solvating the Arg residues (Figure 4, A and C), in a similar fashion to that observed in an earlier simulation of S4.⁵ Although the maximum number of H-bonding partners to an Arg residue is five, the average orientation of the Arg side chain was such that four H-bonding partners were preferred to avoid steric clashes. This is reflected in the simulations; the average total number of solvation

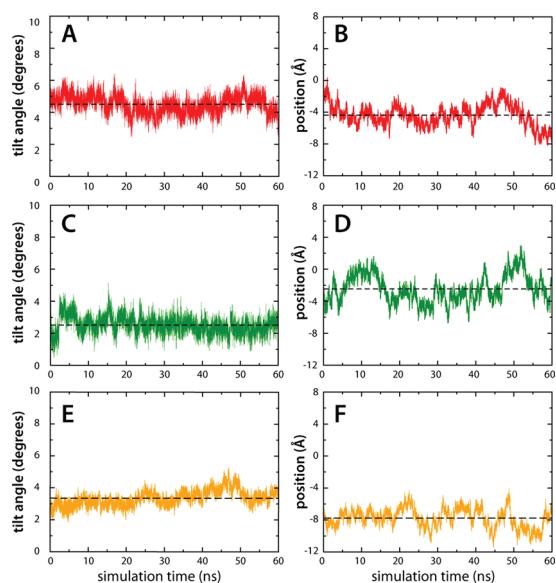


Figure 3. Trajectories of transmembrane tilt angles and positions of different S4mut configurations in POPC and DOTAP bilayers. (A, B) Tilt angles and positions, respectively, for the 2-up-2-down Arg configuration of S4mut in POPC. (C, D) The same as (A) and (B) except the lipid is DOTAP. (D, E) Tilt angles and positions, respectively, for the 1-up-3-down Arg configuration of S4mut in POPC. Tilt angles refer to the angle (in degrees) between the peptide and the bilayer normal. The position refers to the S4mut center of mass (COM) along the bilayer normal (z -dimension).

interactions per Arg residue was around four, with an occasional jump to the maximum of five H-bond partners for Arg9 (Figure 4C). This shows that the Arg residues were essentially fully H bonded throughout the simulation. As a means of observing distortions in the lipid bilayer, we measured the bilayer thickness, defined as the distance between carbonyl carbons in each leaflet. The calculated bilayer thickness for the entire POPC bilayer is 31.5 ± 3.8 Å (Figure 5), whereas close to the peptide, we observed significant distortions of the bilayer. The thickness of the second and first coordination shells, centered at 7 and 4 Å from the center of the peptide, decreased to 25.0 ± 3.0 and 13.7 ± 3.4 Å, respectively. Hence, in order for the structural H-bonding framework to become established around S4mut, the thickness of the lipid bilayer close vicinity of the TM helix is reduced by 6.5 and 17.8 Å for the second and first coordination shells, respectively, thereby bringing the bilayer polar region into close contact with the snorkeling arginines.

The configuration of S4mut in POPC contrasts sharply to that obtained when solvating the peptide in the cationic lipid DOTAP, which lacks phosphate groups. Whereas well-defined positions of the Arg side chains were obtained for S4mut in POPC, the centers of mass of the Arg side chains in DOTAP were widely distributed (Figure 6). These results are consistent with free energy profiles of an isolated arginine side chain displaying relatively high energetics at the DOTAP bilayer interface and more favorable energetics at the same position in bilayers containing lipid phosphates.^{25–27,53}

The lack of lipid phosphate groups in the DOTAP bilayer resulted in charge pairing between the Arg residues and chloride counterions (Figure 4D). We observed no significant difference in the number of Arg–carbonyl interactions between the POPC

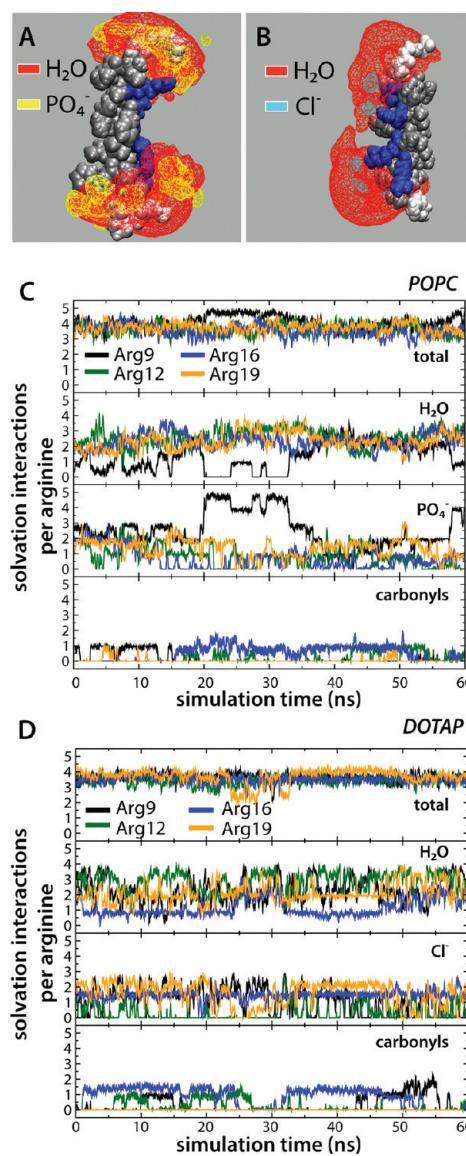


Figure 4. Isosurfaces and solvation-interaction counts for S4mut (2-up-2-down) in POPC and DOTAP bilayers. (A) Isosurfaces of lipid phosphates (yellow) and water molecules (red) based on 60 ns equilibrated simulation of S4mut in POPC. (B) Isosurfaces of water molecules (red) and chloride counterions (cyan) based on 60 ns equilibrated simulation of S4mut in DOTAP. For clarity, the isosurfaces are depicted at 10% occupancy. (C) The total number of solvation interactions (first panel), the number of water H-bonds (second panel), the number of lipid phosphate interactions (third panel), and the number of H-bonds to the lipid carbonyl oxygens (fourth panel), formed by the S4mut helix in POPC. (D) The total number of solvation interactions (upper panel), the number of water H-bonds (second panel), the chloride ion interactions (third panel), and the number of H-bonds to the lipid carbonyl oxygens (fourth panel), formed by S4mut helix in DOTAP. The criteria for the presence of a solvation interaction were defined by cutoff values 3.5 Å for the donor to acceptor distance and 40° for the donor-H-acceptor supplement angle.

and DOTAP systems and, hence, the Arg residues did not seem to use the lipid carbonyl oxygens as additional anchor points to a greater extent when lipid phosphates were not present. Importantly, by interacting electrostatically with the mobile DOTAP chloride counterions, the Arg residues are not anchored to the

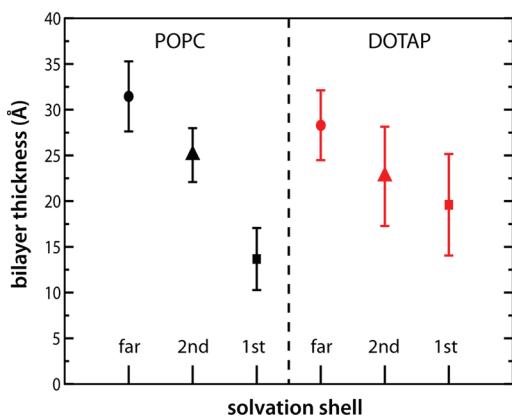


Figure 5. Bilayer thickness measurements for POPC (black symbols) and DOTAP bilayers (red symbols) at various distances from S4mut (2-up-2-down). Measurements represent transbilayer carbonyl-to-carbonyl distances. At large distances from the peptide, the thicknesses are 31.5 ± 3.8 Å for POPC and 28.3 ± 3.8 Å for DOTAP (solid circles). For the second coordination shell, the thicknesses are 25.0 ± 3.0 and 22.7 ± 5.4 Å for POPC and DOTAP, respectively (triangles). For the first coordination shell, the distances decrease to 13.7 ± 3.4 Å for POPC and 19.6 ± 5.6 Å for POPC and DOTAP, respectively (solid squares).

polar region of the DOTAP bilayer (Figure 4B). The resulting increase in conformational flexibility of the Arg side chains affects the water distribution. In POPC, the water molecules are distributed around the snorkeling Arg residues and the hydrophobic gap between the two solvation shells is maintained (Figure 4A). The dynamical feature of this structural framework is illustrated by the number-density plot in Figure 7B, which shows the water profile along the normal of the POPC bilayer. The symmetrically distributed peak regions of the water profile show the preference of water molecules for the polar head groups of the lipid bilayer. The Gaussian character of the water distribution indicates that water penetrates quite deep into the bilayer, although it never enters the hydrophobic gap of the bilayer. In DOTAP, on the other hand, the hydrophobic gap collapses, because water molecules are being dragged through the bilayer by the unanchored Arg residues (Figure 7B). In several instances during the S4mut–DOTAP simulation, we observed the formation of a water wire stretching across the entire bilayer. The occurrence of such a water wire was strictly confined to the immediate vicinity of the peptide at all times during the simulations (Figure 7A). At no instance during the simulation did we observe a chloride ion crossing over the bilayer. Although ions do enter the hydrophobic core in the neighborhood of S4mut, the charge–charge interactions to the arginines prevent any flux of ions over the bilayer. Moreover, because the phosphate–Arg solvation shells were disrupted, bilayer distortion in the vicinity of S4mut was reduced. The average thickness of the DOTAP bilayer was found to be 28.3 ± 3.8 Å, while the second and first coordination shells around the peptide were 22.7 ± 5.4 and 19.6 ± 5.6 Å, respectively (Figure 5). This amounts to a distortion of the DOTAP bilayer in the vicinity of the S4mut peptide of $\sim 5.6\text{--}8.7$ Å, which is significantly less than the distortion of the POPC system.

Comparing the simulations of S4mut in POPC and DOTAP, the nature of the solvation shell in the protein–lipid interface emerges as the most prominent discriminatory feature. Bilayers lacking lipid phosphates are not capable of charge pairing to the

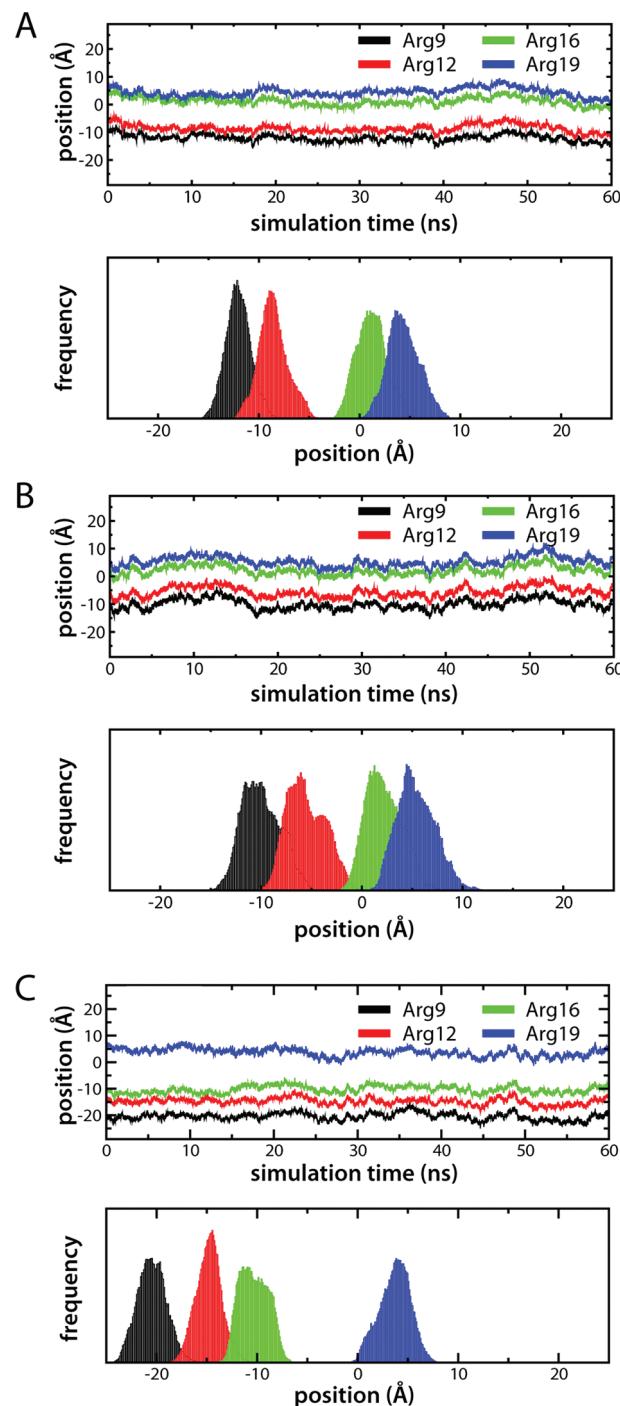


Figure 6. Transbilayer distributions of Arg residues for S4mut in POPC and DOTAP bilayers. The equilibrated positions of the Arg residues are represented by the center-of-mass (COM) trajectories along the bilayer normal (upper panels) and their corresponding distribution histograms (lower panels). In all cases of R9 (black), R12 (red), R16 (green), and R19 (blue), the COM were taken to be the position of the guanidinium carbon atom. (A) The Arg residues (2-up-2-down) in S4mut in POPC. (B) The Arg residues (2-up-2-down) in DOTAP. (C) Arg residues (1-up-3-down configuration) for S4mut in POPC.

Arg residues of the S4mut helix, and as a consequence the bilayer deformation around the peptide is significantly smaller in the DOTAP bilayer. The protein–lipid interactions in these

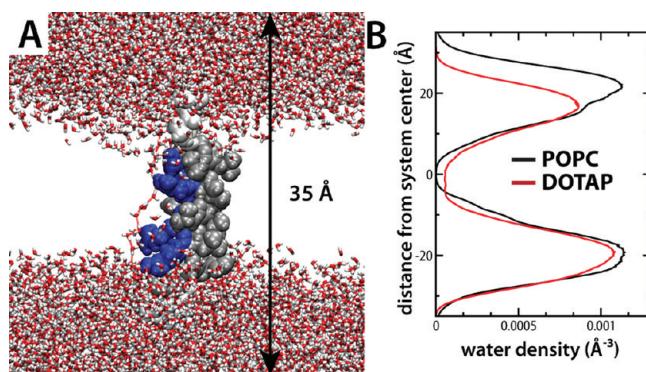


Figure 7. Formation of water wires in DOTAP bilayers containing S4mut (2-up-2-down). (A) A snapshot from the DOTAP simulation showing the formation of a water wire that extends from one side of the bilayer to the other. Connecting water molecules were defined by cutoff values 3.5 Å for the donor to acceptor distance and 40° for the donor-H-acceptor supplement angle. (B) The number densities of water in the first and second solvation shells around the Arg residues in POPC (black line) and DOTAP (red line). Notice the absence waters in the center of the POPC bilayer due to the so-called hydrophobic gap (see text).

simulations provide an atomistic view of the lipid phosphate requirement by Kv channels.

3.2. Anchoring Affects Helical Conformation. The anchoring by lipid phosphates in the head-group region of the bilayer also has a strong influence on the helical conformation of the peptide. This becomes clear by comparing the two rotameric states of S4mut in POPC bilayers. When R16 was pointing up in the 2-up-2-down configuration, the integrity of the helix was maintained throughout the simulation (Figure 2A). Although the 1-up-3-down rotamer distribution (R16 pointing down) was seen earlier to generate a conformationally stable S4 helix in POPC,⁵ our 100 ns MD simulation of S4mut with the same rotamer distribution in POPC produced a strongly kinked peptide, suggesting a conformationally unstable state (Figure 2B). A similar behavior, where the rotameric state of a charged side chain dictates the secondary structure of a TM segment, was observed in simulations of the membrane-spanning domain (MSD) of the HIV-1 glycoprotein gp41.⁵⁴ Furthermore, Nishizawa et al. reported highly tilted helix orientations in MD simulations where the S4 helix was placed at different initial depths in a dipalmitoylphosphatidylcholine (DPPC) bilayer.³⁸ Interestingly, in this latter work the rotameric distribution of the Arg residues is either in a 1-up-3-down or a 3-up-1-down configuration, thereby confirming our observation that strain is imposed on the helix when the arginines are not in a 2-up-2-down configuration. The movement of the side chain of either R12 or R16, in the 3-up-1-down and 1-up-3-down configurations, respectively, toward the edge of the bilayer reduces the energetic penalty of charge burial in the nonpolar interior of the membrane, and results in a conformationally unstable state. The structural consequences seem to be either a pronounced helical tilt³⁸ or, as seen in our simulations, a strong kink in the helix that disturbs the normal α -helical H-bonding pattern in the upper part of the helix (Figure 2D).

Because switching through different rotameric states of charged residues within the lipid bilayer involves the crossing of high energy barriers,²⁶ the most energetically favorable set of side-chain configurations is maintained throughout a simulation.

This analysis prompted us to use the 2-up-2-down configuration of S4mut in the simulation of the DOTAP bilayer.

4. SUMMARY

Taken together, our simulations show that in bilayers lacking lipid phosphates, the charged Arg residues of S4mut are solvated by chloride counterions, water molecules, and, to a much lesser degree, the carbonyl groups of the surrounding lipids. In such a protein–lipid interface, the Arg residues lack anchor points to the lipid bilayer and, hence, display an increased structural flexibility along the bilayer normal. As a direct consequence of the nature of this protein–lipid interface, we make two observations that might explain the nonfunctionality of Kv channels in DOTAP bilayers. First, because of the lack of lipid anchors, Arg residues, solvated by chloride counterions and water molecules, move freely into the bilayer core and ultimately disrupt the integrity of the hydrophobic gap. Second, the absence of strong lipid connections leads to a significantly less distorted DOTAP bilayer around the charged Arg residues. Moreover, this set of simulations show that phosphate–Arg interactions have a strong influence on the secondary structure of the TM segment.

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■ REFERENCES

- (1) Lee, A. G. *Biochim. Biophys. Acta* **2004**, *1666*, 62–87.
- (2) Schmidt, D.; Jiang, Q.-X.; MacKinnon, R. *Nature* **2006**, *444*, 775–779.
- (3) Ramu, Y.; Xu, Y.; Lu, Z. *Nature* **2006**, *442*, 696–699.
- (4) Xu, Y.; Ramu, Y.; Lu, Z. *Nature* **2008**, *451*, 826–830.
- (5) Freites, J. A.; Tobias, D. J.; von Heijne, G.; White, S. H. *Proc. Natl. Acad. Sci. U.S.A.* **2005**, *102*, 15059–15064.
- (6) Hodgkin, A. L.; Huxley, A. F. J. *Physiol. (London)* **1952**, *116*, 449–472.
- (7) Hodgkin, A. L.; Huxley, A. F. J. *Physiol. (London)* **1952**, *116*, 473–496.
- (8) Hodgkin, A. L.; Huxley, A. F. J. *Physiol. (London)* **1952**, *116*, 497–506.
- (9) Hodgkin, A. L.; Huxley, A. F. J. *Physiol. (London)* **1952**, *117*, 500–544.
- (10) Seoh, S. A.; Sigg, D.; Papazian, D. M.; Bezanilla, F. *Neuron* **1996**, *16*, 1159–1167.
- (11) Aggarwal, S. K.; MacKinnon, R. *Neuron* **1996**, *16*, 1169–1177.
- (12) Swartz, K. J. *Nature* **2008**, *456*, 891–897.

- (13) Jiang, Y. X.; Ruta, V.; Chen, J. Y.; Lee, A.; MacKinnon, R. *Nature* **2003**, *423*, 42–48.
- (14) Cuello, L. G.; Cortes, D. M.; Perozo, E. *Science* **2004**, *306*, 491–495.
- (15) Ruta, V.; Chen, J.; MacKinnon, R. *Cell* **2005**, *123*, 463–475.
- (16) Alabi, A. R.; Bahamonde, M. I.; Jung, H. J.; Kin, J. I.; Swartz, K. J. *Nature* **2007**, *450*, 370–376.
- (17) Long, S. B.; Tao, X.; Campbell, E. B.; Mackinnon, R. *Nature* **2007**, *450*, 376–382.
- (18) Chakrapani, S.; Cuello, L. G.; Cortes, D. M.; Perozo, E. *Structure* **2008**, *16*, 398–409.
- (19) Parsegian, A. *Nature* **1969**, *221*, 844–846.
- (20) Grabe, M.; Lecar, H.; Jan, Y. N.; Jan, L. Y. *Proc. Natl. Acad. Sci. U.S.A.* **2004**, *101*, 17640–17645.
- (21) Tombola, F.; Pathak, M. M.; Isacoff, E. Y. *Annu. Rev. Cell Dev. Biol.* **2006**, *22*, 23–52.
- (22) Castro-Román, F.; Fernandez-Vidal, M.; Mihailescu, M.; White, S. H. *Biophys. J.* **2007**, *92*, 294a.
- (23) Doherty, T.; Su, Y.; Hong, M. J. *Mol. Biol.* **2010**, *401*, 642–652.
- (24) Hessa, T.; White, S. H.; von Heijne, G. *Science* **2005**, *307*, 1427.
- (25) MacCallum, J. L.; Bennett, W. F. D.; Tielemans, D. P. *Biophys. J.* **2008**, *94*, 3393–3404.
- (26) Dorairaj, S.; Allen, T. W. *Proc. Natl. Acad. Sci. U.S.A.* **2007**, *104*, 4943–4948.
- (27) Schow, E. V.; Freites, J. A.; Cheng, P.; Bernsel, A.; von Heijne, G.; White, S. H.; Tobias, D. J. *J. Membr. Biol.* **2011**, *239*, 35–48.
- (28) Freites, J. A.; Tobias, D. J.; White, S. H. *Biophys. J.* **2006**, *91*, L90–L92.
- (29) Sands, Z. A.; Sansom, M. S. P. *Structure* **2007**, *15*, 235–244.
- (30) Treptow, W.; Tarek, M. *Biophys. J.* **2006**, *90*, L64–L66.
- (31) Jogini, V.; Roux, B. *Biophys. J.* **2007**, *93*, 3070–3082.
- (32) Bjelkmar, P.; Niemelä, P. S.; Vattulainen, I.; Lindahl, E. *PLoS Comput. Biol.* **2009**, *5*, 1–14.
- (33) Krepkiy, D.; Mihailescu, M.; Freites, J. A.; Schow, E. V.; Worcester, D. L.; Gawrisch, K.; Tobias, D. J.; White, S. H.; Swartz, K. J. *Nature* **2009**, *462*, 473–479.
- (34) Butterwick, J. A.; MacKinnon, R. *J. Mol. Biol.* **2010**, *403*, 591–606.
- (35) Hristova, K.; White, S. H. *Biophys. J.* **1998**, *74*, 2419–2433.
- (36) McIntosh, T. J.; Magid, A. D.; Simon, S. A. *Biophys. J.* **1989**, *55*, 897–904.
- (37) Hessa, T.; Kim, H.; Bihlmaier, K.; Lundin, C.; Boekel, J.; Andersson, H.; Nilsson, I. M.; White, S. H.; von Heijne, G. *Nature* **2005**, *433*, 377–381.
- (38) Nishizawa, M.; Nishizawa, K. *Biophys. J.* **2008**, *95*, 1729–1744.
- (39) Humphrey, W.; Dalke, W.; Schulten, K. *J. Mol. Graph.* **1996**, *14*, 33–38.
- (40) Emsley, P.; Cowtan, K. *Acta Crystallogr.* **2004**, *D60*, 2126–2132.
- (41) Keer, H. S.; Freites, J. A.; Mihailescu, E.; White, S. H.; Tobias, D. J. *Biophys. J.* **2010**, *98* (Suppl. 1), 492a.
- (42) Kalé, L.; Skeel, R.; Bhandarkar, M.; Brunner, R.; Gursoy, A.; Krawetz, N.; Phillips, J.; Shinozaki, A.; Varadarajan, K.; Schulten, K. *J. Comput. Phys.* **1999**, *151*, 283–312.
- (43) Jorgensen, W. L.; Chandrasekhar, J.; Madura, J. D.; Impey, R. W.; Klein, M. L. *J. Chem. Phys.* **1983**, *79*, 926–935.
- (44) MacKerell, A. D., Jr.; Bashford, D.; Bellott, M.; Dunbrack, R. L., Jr.; Evanseck, J. D.; Field, M. J.; Fischer, S.; Gao, J.; Guo, H.; Ha, S.; Joseph-McCarthy, D.; Kuchnir, L.; Kuczera, K.; Lau, F. T. K.; Mattos, C.; Michnick, S.; Ngo, T.; Nguyen, D. T.; Prodhom, B.; Reiher, W. E., III; Roux, B.; Schlenkrich, M.; Smith, J. C.; Stote, R.; Straub, J.; Watanabe, M.; Wiórkiewicz-Kuczera, J.; Yin, D.; Karplus, M. *J. Phys. Chem. B* **1998**, *102*, 3586–3616.
- (45) Klauda, J. B.; Venable, R. M.; Freites, J. A.; O'Connor, J. W.; Tobias, D. J.; Mondragon-Ramirez, C.; Vorobyov, I.; MacKerell, A. D.; Pastor, R. W. *J. Phys. Chem. B* **2010**, *114*, 7830–7843.
- (46) Grubmüller, H.; Heller, H.; Windemuth, A.; Schulten, K. *Mol. Simul.* **1991**, *6*, 121–142.
- (47) Darden, T.; York, D.; Pedersen, L. *J. Chem. Phys.* **1993**, *98*, 10089–10092.
- (48) Essmann, U.; Perera, L.; Berkowitz, M. L.; Darden, T.; Lee, H.; Pedersen, L. G. *J. Chem. Phys.* **1995**, *103*, 8577–8593.
- (49) Ryckaert, J.-P.; Ciccotti, G.; Berendsen, H. J. C. *J. Comput. Phys.* **1977**, *23*, 327–341.
- (50) Miyamoto, S.; Kollman, P. *J. Comput. Chem.* **1992**, *13*, 952–962.
- (51) Martyna, G. J.; Tobias, D. J.; Klein, M. L. *J. Chem. Phys.* **1994**, *101*, 4177–4189.
- (52) Feller, S. E.; Zhang, Y.; Pastor, R. W.; Brooks, B. R. *J. Chem. Phys.* **1995**, *103*, 4613–4621.
- (53) Johansson, A. C.; Lindahl, E. *J. Chem. Phys.* **2009**, *130*, 185101–185108.
- (54) Gangopomu, V. K.; Abrams, C. F. *Biophys. J.* **2010**, *99*, 3438–3444.