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Flattening a Puckered Pentasilacyclopentadienide Ring by Suppression of the Pseudo Jahn-Teller Effect

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We report the theoretical prediction of flattening of the puckered Si_5 ring by suppression of the pseudo-Jahn–Teller effect through coordination of two Mg^{2+} cations to the $Si_5H_5^-$ anion to make an inverse $[Mg^{2+}Si_5H_5^-Mg^{2+}]$ sandwich complex. The pseudo-Jahn–Teller (PJT) effect was suppressed through the OMO–UMO gaps increase in the resultant $[Mg^{2+}Si_5H_5^-Mg^{2+}]$ sandwich complex, as compared to the initial $Si_5H_5^-$ anion. It was the influence of two Mg^{2+} cations that caused the OMO–UMO gaps increase that made this type of PJT effect suppression work. In the three other complexes under the current computational study, namely, $[Li^+Si_5H_5^-Li^+]$, $[Na^+Si_5H_5^-Na^+]$, and $[Be^{2+}Si_5H_5^-Be^{2+}]$, the $Si_5H_5^-$ moiety remains nonplanar. We believe that if the $Mg(Si_5H_5)_2$ solid compound were synthesized, it could have planar $Si_5H_5^-$ building blocks.

Introduction

Cyclic systems such as benzene (C₆H₆), cyclopentadienyl anion (C₅H₅⁻), and the recently synthesized cyclotetrabutadiene dianion derivative $((Me_3Si)_4C_4^{2-})^{1-3}$ are planar due to π -aromaticity, since they all have six π -electrons. These three cyclic systems are prototypical aromatic systems according to the 4n + 2 Huckel rule where n = 1. Out of three silicon analogues, only the tetrasiliconcyclobutadiene dianion derivative $[(^{t}Bu_{2}MeSi)_{4}Si_{4}]^{2-}$, as a part of the $[K(thf)_{2}]_{2}[(^{t}Bu_{2}MeSi)_{4}Si_{4}]$ salt, was synthesized⁴ and shown to have a planar tetraatomic cyclic structure at the core using X-ray diffraction. The Si₆H₆ and Si₅H₅ species have not been obtained experimentally yet. Theoretical calculations of Si₆H₆ and Si₅H₅⁻ showed that both of these molecules are nonplanar. ^{5,6} The deviation from planarity in Si₆H₆ and Si₅H₅⁻ species is caused by the only source of instability of high-symmetry configurations of polyatomic systems, namely, the Jahn-Teller vibronic effects, ⁷⁻⁹

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the pseudo-Jahn-Teller (PJT) effect to be precise. We have demonstrated recently 10 that the PJT effect can be suppressed by flattening the nonplanar "chair" structure of Si_6X_{12} (X = Cl, Br) upon addition of two X^- ions to form the $[Si_6Cl_{14}]^{2-}$ dianion. The resultant $[Si_6Cl_{14}]^{2-}$ dianion was shown to have the restored planar Si₆ hexagon ring "sandwiched" between two apical chlorides. ¹⁰ The comprehensive computational analysis of the Si₆Cl₁₂ electronic structure allowed us to identify the main cause of the suppression of the PJT effect, which originates from filling the initially empty unoccupied molecular orbitals (UMOs) of Si₆Cl₁₂ with electrons of the adduct (i.e., Cl⁻). In the current article we show that there is an alternative mechanism of suppression of the PJT effect through extending the gap between interacting occupied and empty MOs, which participate in the PJT effect. We also demonstrate that the Si₅H₅ structure, which is nonplanar in the isolated state, becomes planar in the [MgSi₅H₅Mg]³⁺ sandwich complex through the OMO-UMO gaps increase caused by positive charge influence of the two extra Mg²⁺ cations.

Computational Results

We initially calculated the planar pentagonal D_{5h} ($^{1}A_{1}'$) structure of $Si_{5}H_{5}^{-}$ (Figure 1, structure I) at two levels of theory: $B3LYP/6-311++G^{**}$ and $CCSD(T)/6-311++G^{**}$.

All calculations were performed using the Gaussian 03 program. The Geometry and molecular orbital visualization was done using the MOLEKEL 4.3 package. We found that the D_{5h} ($^{1}A_{1}'$) structure I is not a minimum, in agreement with the previously reported computational results for $\mathrm{Si}_{5}\mathrm{H}_{5}^{-}$ by Korkin et al. It was found to have two doubly degenerate imaginary frequencies ($\omega_{1,2}(\mathrm{e}_{2}'')$) and $\omega_{3,4}(\mathrm{e}_{1}'')$), leading to out-of-plane distortions, according to the B3LYP/6-311++G** and the CCSD(T)/6-311++G** levels of theory. Geometry optimization (at B3LYP/6-311++G**) following these imaginary frequency modes leads

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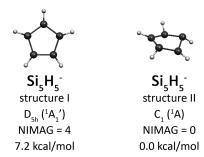


Figure 1. Alternative structures of $\mathrm{Si}_5\mathrm{H}_5^-$ (structure I and structure II), their symmetry, spectroscopic states, number of imaginary frequency modes (NIMAG), and relative energies corrected for zero-point energy, all calculated at the B3LYP/6-311++G** level of theory.

to one isomer: C_1 (¹A) (Figure 1, structure II). We also repeated geometry optimization for the two isomers, which were previously reported by Korkin et al.⁶ The authors found two degenerate nonplanar isomers of $Si_5H_5^-$ of C_s and C_2 symmetry, respectively, according to their calculations performed at the MP2(full)/6-31G* level of theory. We found that these two isomers reported by Korkin et al. are not minima at the B3LYP/6-311++G** level of theory, which is used in the current work. Geometry optimization following imaginary frequency modes of C_s and C_2 isomers leads again to structure II. Hence, there is only one isomer of $Si_5H_5^-$ at B3LYP/6-311++G**, which is of C_1 symmetry in the ¹A electronic state (Figure 1, structure II). The planar D_{5h} (¹A₁') structure I is only 7.2 kcal/mol higher in energy that the puckered C_1 (¹A) structure II at the B3LYP/6-311++G** level of theory with zero-point energy correction.

The planar D_{5h} (${}^{1}A_{1}{}'$) structure I is not a minimum due to the pseudo-Jahn–Teller (PJT) effect, which forces puckering of the Si₅ ring and leads to out-of-plane distortions along the imaginary frequency modes. The PJT effect in the D_{5h} (${}^{1}A_{1}{}'$, $1a_{1}{}'^{2}1e_{1}{}'^{4}1e_{2}{}'^{4}2a_{1}{}'^{2}2e_{1}{}'^{4}2e_{2}{}'^{4}1a_{2}{}''^{2}1e_{1}{}''^{4}3a_{1}{}'^{0}3e_{1}{}'^{0}3e_{2}{}'^{0}$) structure of Si₅H₅ is a consequence of vibronic coupling of pairs of occupied molecular orbitals (OMO) and unoccupied molecular orbitals (UMO). The symmetry rule for the PJT effect is that the symmetry of the reaction coordinate (symmetry of one of the two doubly degenerate imaginary frequency modes) should be the same as the direct product of the OMO and UMO states 13 as follows (eq 1):

$$\psi_{\rm OMO} \times \psi_{\rm UMO} = \Gamma_{\rm RctnCoord}$$
 (1)

Thus, the deformation of the D_{5h} ($^{1}A_{1}'$) structure I of Si₅H₅⁻ along the $\omega_{1,2}(e_{2}'')$ imaginary frequency mode can be caused by interaction between the HOMO ($1e_{1}''$) and LUMO+1 ($3e_{1}'$) (see eq 2 and Figure 2a) and/or by interaction between the HOMO ($1e_{1}''$) and LUMO+2 ($3e_{2}'$) (see eq 3 and Figure 2b):

$$e_1'' \times e_1' = a_1'' + a_2'' + e_2''$$
 (2)

$$e_1'' \times e_2' = e_1'' + e_2''$$
 (3)

The deformation of the D_{5h} ($^{1}A_{1}'$) structure I of $Si_{5}H_{5}^{-}$ along the $\omega_{3,4}(e_{1}'')$ imaginary frequency mode can be caused by interaction between the HOMO ($1e_{1}''$) and LUMO ($3a_{1}'$) (see eq 4 and Figure 2c) and/or by interaction between the HOMO ($1e_{1}''$) and LUMO+2 ($3e_{2}'$) (see eq 3 and Figure 2b):

$$e_1'' \times a_1' = e_1''$$
 (4)

The orbital energies of the orbitals involved in the pseudo-Jahn–Teller effect are the following: $\varepsilon(HOMO, 1e_1'') = -0.078$ au; $\varepsilon(LUMO, 3a_1') = 0.164$ au; $\varepsilon(LUMO+1, 3e_1') = 0.165$ au; and $\varepsilon(LUMO+2, 3e_2') = 0.172$ au. Thus, the orbital energy difference

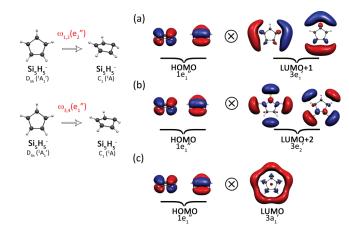


Figure 2. Interaction of the pairs of occupied and unoccupied molecular orbitals of the D_{5h} ($^{1}A_{1}'$) structure I of $Si_{5}H_{5}^{-}$ responsible for the PJT effect: (a) and (b) interactions cause distortion toward the C_{1} (^{1}A) structure upon following the doubly degenerate $\omega_{1,2}(e_{2}'')$ imaginary frequency mode; (b) and (c) interactions cause distortion toward the C_{1} (^{1}A) structure upon following the doubly degenerate $\omega_{3,4}(e_{1}'')$ imaginary frequency mode.

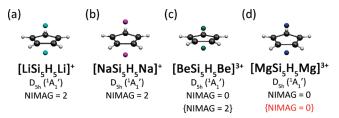


Figure 3. Optimized structures, point group symmetry, spectroscopic states, and number of imaginary frequency modes (NIMAG) of (a) $\text{Li}^+\text{Si}_5\text{H}_5^-\text{Li}^+$, (b) $\text{Na}^+\text{Si}_5\text{H}_5^-\text{Na}^+$, (c) $\text{Be}^{2+}\text{Si}_5\text{H}_5^-\text{Be}^{2+}$, and (d) $\text{Mg}^{2+}\text{Si}_5\text{H}_5^-\text{Mg}^{2+}$ at the B3LYP/6-311++G** level of theory and at the CCSD(T)/6-311++G** level of theory (in curly brackets).

in the $\mathrm{Si}_5\mathrm{H}_5^-$ anion between the HOMO and LUMO is 6.58 eV, while the gap between the HOMO and LUMO+1 is 6.60 eV, which are both somewhat smaller than the energy difference between the HOMO and LUMO+2 (6.80 eV). All the above values are given at the RHF/6-311++G** level of theory since our most reliable data were obtained at the CCSD(T)/6-311++G** level, which is based on RHF/6-311++G**.

In order to test if the PJT effect can be suppressed by external charges, we then calculated four inverse sandwich complexes, namely, $\text{Li}^+\text{Si}_5\text{H}_5^-\text{Li}^+$ (Figure 3a), $\text{Na}^+\text{Si}_5\text{H}_5^-\text{Na}^+$ (Figure 3b), $\text{Be}^{2+}\text{Si}_5\text{H}_5^-\text{Be}^{2+}$ (Figure 3c), and $\text{Mg}^{2+}\text{Si}_5\text{H}_5^-\text{Mg}^{2+}$ (Figure 3d) at the B3LYP/6-311++G** level of theory.

In the following text are the results of testing whether the electrostatic field from the pair of Li⁺, Na⁺, Be²⁺, or Mg²⁺ cations is enough to suppress the PJT effect in the Si₅H₅⁻ anion by shifting apart occupied molecular orbitals (OMO) and unoccupied molecular orbitals (UMO). The addition of two Li⁺ or Na⁺ cations to make the M⁺Si₅H₅⁻M⁺ (M = Li, Na) complex or a pair of Be²⁺ or Mg²⁺ dications to make M²⁺Si₅H₅⁻M²⁺ (M = Be, Mg) does not increase the number of valence electrons in the resultant inverse sandwich complex, as compared to the initial Si₅H₅⁻ anion, although this coordination of the two positively charged ions to Si₅H₅⁻ affects the gap between occupied and unoccupied molecular orbitals responsible for the PJT effect, which manifests itself in the increased stability of the Mⁿ⁺Si₅H₅⁻Mⁿ⁺ complexes. According to our calculations, all four complexes shown in Figure 3 do not have doubly degenerate

imaginary frequency mode of e₁" symmetry, which was present in isolated Si₅H₅⁻. That's due to the gap increase between pairs of OMO-UMO molecular orbitals: $1e_1''$ and $3a_1'$; $1e_1''$ and $3e_2'$. The gap between $1e_1^{\prime\prime}$ and $3a_1^{\prime}$ OMO-UMO pair is increased from 6.58 eV to 8.18 eV (in $\text{Li}^+\text{Si}_5\text{H}_5^-\text{Li}^+$); from 6.58 eV to 7.27 eV (in $Na^{+}Si_{5}H_{5}^{-}Na^{+}$); from 6.58 eV to 9.15 eV (in $Be^{2+}Si_{5}H_{5}^{-}Be^{2+}$); and from 6.58 eV to 7.16 eV (in $Mg^{2+}Si_5H_5^-Mg^{2+}$). The gap between 1e₁" and 3e₂' OMO-UMO pair is increased from 6.80 eV to 9.95 eV (in $Li^+Si_5H_5^-Li^+$); from 6.80 eV to 9.42 eV (in $Na^+Si_5H_5^-Na^+$); from 6.80 eV to 11.99 eV (in $Be^{2+}Si_5H_5^-Be^{2+}$); and from 6.80 eV to 10.94 eV (in Mg²⁺Si₅H₅⁻Mg²⁺). Apparently, these gaps' increase was enough to completely suppress the PJT effect along the $\omega_{3,4}(e_1^{"})$ mode. However, two complexes, Li⁺Si₅H₅⁻Li⁺ and Na⁺Si₅H₅⁻Na⁺, retained doubly degenerate imaginary frequency mode of e₂" symmetry upon addition of cations to isolated Si₅H₅. Thus, the presence of two extra Li⁺ or Na⁺ cations does not suppress the PJT effect completely, although it reduces the number of degenerate imaginary frequency modes by one. Both $Be^{2+}Si_5H_5^-Be^{2+}$ and $Mg^{2+}Si_5H_5^-Mg^{2+}$ complexes were found to be minima at the B3LYP/6-311++G** level of theory. However, when we reoptimized the geometry of the Be²⁺Si₅H₅⁻Be²⁺ complex and calculated its frequency modes at the $CCSD(T)/6-311++G^{**}$ level of theory, it was found to have one doubly degenerate imaginary frequency of e2" symmetry. Thus, Be²⁺ cations cannot suppress the PJT effect either. Meanwhile, the Mg²⁺Si₅H₅⁻Mg²⁺ complex was found to be a minimum at both levels of theory, with the lowest $\omega_{1,2}(e_2'')$ frequency mode being 90.0 cm⁻¹ (B3LYP/6-311++G**) and 78.8 cm⁻ $(CCSD(T)/6-311++G^{**})$, which indicates that only the electrostatic field of Mg²⁺ causes the suppression of the PJT effect. The orbital energies of the orbitals of Mg²⁺Si₅H₅⁻Mg²⁺, which correspond to those involved in pseudo-Jahn–Teller effect of Si_5H_5 , are the following: $\varepsilon(\text{HOMO}, 1\text{e}_1^{\prime\prime}) = -0.737 \,\text{au}; \varepsilon(\text{LUMO}, 1\text{e}_1^{\prime\prime}) = -0.737$ $3a_1'$) = -0.474 au; $\varepsilon(LUMO+4, 3e_1')$ = -0.350 au; and $\varepsilon(LUMO+5, 3e_2') = -0.335$ au. Hence, the gap between the HOMO ($1e_1''$) and LUMO ($3a_1'$) is increased from 6.58 eV in $Si_5H_5^-$ to 7.16 eV in $Mg^{2+}Si_5H_5^-Mg^{2+}$. The gaps between other UMO and OMO pairs that participate in the PJT effect are increased by more than 0.60 eV in Mg²⁺Si₅H₅-Mg²⁺: from 6.60 eV to 10.52 eV (in the case of $1e_1^{\prime\prime}$ and $3e_1^{\prime}$) and from 6.80 eVto 10.94 eV (in the case of $1e_1''$ and $3e_2'$). Apparently this extension of UMO-OMO gaps is large enough to attain planarity of Si₅H₅⁻ in the inverse sandwich structure of Mg²⁺Si₅H₅⁻Mg²⁺. As one of the reviewers recommended, we performed additional calculations at B3LYP/6-311++G** to study the planarity or lack of planarity of the underlying five-membered ring in Mg²⁺Si₅H₅-Mg²⁺ as the distance between the magnesium cations and the ring is varied. We performed frequency calculations of Mg²⁺Si₅H₅⁻Mg²⁺ at longer (with two Mg²⁺ cations moved farther from the center of mass by 0.1 Å) and shorter (with two Mg²⁺ cations moved closer to the center of mass by 0.1 Å) cation-ring distances, as well as with one Mg²⁺ cation moved farther from the center of mass by 0.1 Å and the other Mg²⁺ cation moved closer to the center of mass by 0.1 A. In the two latter cases the Mg²⁺Si₅H₅⁻Mg²⁺ retained no imaginary frequencies, while the Mg²⁺Si₅H₅⁻Mg²⁺ with longer cation ring distances has two imaginary frequencies. This elucidates the fact that the cation—ring interactions are mainly Coulombic and that there is a delicate balance between external field of cations and the geometry.

One would expect that Be²⁺ cation being smaller and hence located closer to the Si₅H₅⁻ ring should also form the stable M²⁺Si₅H₅⁻M²⁺ complex similar to Mg²⁺Si₅H₅⁻Mg²⁺. Indeed, the gap between $1e_1^{"}$ and $3e_1^{'}$ pair of OMO-UMO molecular orbitals (involved in PJT effect along the doubly degenerate imaginary frequency mode of e2" symmetry) is increased from 6.60 eV (in $Si_5H_5^-$) to 11.56 eV (in $Be^{2+}Si_5H_5^-Be^{2+}$), hence, by 4.96 eV. This gap increase in $Be^{2+}Si_5H_5^-Be^{2+}$ is even higher than there in $Mg^{2+}Si_5H_5^-Mg^{2+}$ (which is 3.72 eV). We believe, that the suppression of the PJT distortion in $Be^{2+}Si_5H_5^-Be^{2+}$ doesn't occur due to the nature of chemical bonding. Natural Bonding Orbital atomic charges calculated at B3LYP/6- $311++G^{**}$ are +1.30 |e| on each of the beryllium atoms and +0.40 |e| on Si_5H_5 ring in $Be^{2+}Si_5H_5^-Be^{2+}$, while in the $Mg^{2+}Si_5H_5^-Mg^{2+}$ complex they are +1.65 |e| on each of the magnesium atoms and -0.30 |e| on Si₅H₅ ring. This elucidates that the nature of the chemical bonding in the Be²⁺Si₅H₅⁻Be²⁺ complex is more covalent than there in the Mg²⁺Si₅H₅⁻Mg²⁺ complex. The Li⁺Si₅H₅⁻Li⁺ and Na⁺Si₅H₅⁻Na⁺ complexes are more ionic as expected: the natural charges in Li⁺Si₅H₅⁻Li⁺ are +0.87 |e| (on each of the lithium atoms) and -0.74|e| (on Si₅H₅ ring); and there in the $Na^+Si_5H_5^-Na^+$ complex they are +0.93|e| (on each of the sodium atoms) and -0.86|e| (on Si₅H₅ ring). The electrostatic field of a singly charged cation is weaker than that of doubly charged one. The gap between $1e_1''$ and $3e_1'$ OMO-UMO pair is increased by 2.85 eV (in Li⁺Si₅H₅⁻Li⁺), by 2.00 eV (in $Na^+Si_5H_5^-Na^+$), as compared to 3.72 eV (in $Mg^{2+}Si_5H_5^-Mg^{2+}$). This explains why two Li⁺ cations or two Na⁺ cations cannot suppress the PJT effect in $Si_5H_5^-$.

The PJT effect is a qualitative model of chemical bonding, and it is not capable of predicting to what extent the OMO-UMO gaps should be increased in order to suppress its distortive nature. However, it can provide direction in the search for the species where the high symmetry can be attained. In our case we studied four types of counter cations, and it was enough to find the appropriate combination of cations to suppress the PJT effect.

Aromaticity in Si₅H₅⁻ and Mg²⁺Si₅H₅⁻Mg²⁺

The Si₅H₅⁻ anion is aromatic in the planar form according to the molecular orbital chemical bonding analysis and the Si-Si bond equalization. We performed a nuclear-independent chemical shift (NICS) analysis 14 for both $\mathrm{Si}_5\mathrm{H}_5^-$ and its planar carbon analogue, $\mathrm{C}_5\mathrm{H}_5^-$. The calculated NICS $_{zz}$ values of the $Si_5H_5^-$ anion were found to be -11.5 ppm (Z = 0.0 Å), where Z represents the distance from the center of the $Si_5H_5^-$ moiety along the z axis), -12.1 ppm (Z = 0.2 Å), -13.5 ppm (Z = 0.4 Å), -15.6 ppm (Z = 0.6 Å), -17.7 ppm(Z = 0.8 Å), and -19.3 ppm (Z = 1.0 Å) at the optimized planar B3LYP/6-311++ G^{**} geometry. The above NICS_{zz} values of the Si₅H₅ anion can be compared to those of the $C_5H_5^-$ anion at the same level of theory: -16.2 ppm (Z=0.0A, where Z represents the distance from the center of the $C_5H_5^-$ moiety along the z axis), -18.7 ppm (Z = 0.2 Å), -24.6 ppm (Z = 0.4 Å), -30.3 ppm (Z = 0.6 Å), -33.5 ppm $(Z = \hat{0.8} \text{ Å})$, and -33.7 ppm (Z = 1.0 Å). It is clear that Si₅H₅⁻ is less aromatic than C₅H₅⁻, according to the calculated NICS_{zz} values. Since, the global minimum structure of the Si₅H₅⁻ anion is nonplanar, the PJT effect overwhelms the stabilization due to aromaticity. When the planarity of the Si₅H₅⁻ anion is restored in the sandwich Mg²⁺Si₅H₅⁻Mg²⁺ structure, aromaticity in the Si₅H₅⁻ anion as a part of the inverse sandwich complex is also restored. The NICS_{zz} values calculated for the ${\rm Mg^{2+}Si_5H_5}^{-}{\rm Mg^{2+}}$ complex are -15.1 ppm (Z = 0.0 Å), -15.9 ppm (Z = 0.2 Å), -18.2ppm (Z = 0.4 Å), -21.7 ppm (Z = 0.6 Å), -26.3 ppm (Z = 0.8 Å)Å), and -32.1 ppm (Z = 1.0 Å).

Conclusions

In this article we have shown that the nonplanarity of Si₅H₅ caused by the pseudo-Jahn-Teller effect can be eliminated through coordination of two Mg²⁺ cations to it to make the

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 ${\rm Mg^{2+}Si_5H_5^-Mg^{2+}}$ inverse sandwich complex. The PJT effect was suppressed in the ${\rm Si_5H_5^-}$ anion through the OMO–UMO gaps increase in the ${\rm Mg^{2+}Si_5H_5^-Mg^{2+}}$ complex because of the charge influence of two ${\rm Mg^{2+}}$ cations. We also presented the data on the [Li $^+$ Si $_5$ H $_5^-$ Li $^+$], [Na $^+$ Si $_5$ H $_5^-$ Na $^+$], and [Be 2 +Si $_5$ H $_5^-$ Be 2 +] complexes. None of them was shown to contain a planar Si $_5$ H $_5^-$ moiety. It was concluded that out of all the studied cations only two ${\rm Mg^{2+}}$ ions are able to increase the OMO–UMO gaps to an appreciable extent to suppress the PJT effect and restore the planarity of Si $_5$ H $_5^-$. We believe that suppressing the PJT effect through coordination of extra cations can be a useful tool in restoring high-symmetry structures of numerous molecular species with otherwise distorted cyclic geometries.

The salts containing $C_5H_5^-$ anions are known to form stacking structures, where $C_5H_5^-$ anions together with counterions form "infinite" piles of sandwich-type structures. We hope that our theoretical prediction of planarization of the $Si_5H_5^-$ anion in the $Mg^{2+}Si_5H_5^-Mg^{2+}$ complex stimulates

the synthesis of the first compounds containing the planar $Si_5H_5^-$ anion.

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Note Added after ASAP Publication: This paper was published on the Web on August 6, 2010, with several errors in acronyms. The corrected version was reposted on August 12, 2010.

Supporting Information Available: Reference 11 in full and Cartesian coordinates of $Si_5H_5^-$ stuctures as well as those of the $Li^+Si_5H_5^-Li^+$, $Na^+Si_5H_5^-Na^+$, $Be^{2+}Si_5H_5^-Be^{2+}$, and $Mg^{2+}Si_5H_5^-Mg^{2+}$ complexes are available free of charge via the Internet at http://pubs.acs.org.

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