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## Microstructure and high critical current density of *in situ* processed MgB<sub>2</sub> tapes made by WSi<sub>2</sub> and ZrSi<sub>2</sub> doping

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WSi<sub>2</sub>- and ZrSi<sub>2</sub>-doped Fe-sheathed MgB<sub>2</sub> tapes were prepared through the *in situ* powder-in-tube method. Both WSi<sub>2</sub>- and ZrSi<sub>2</sub>-doped tapes were found to have significantly increased critical current density  $J_c$  at 4.2 K in magnetic fields up to 12 T than their undoped counterpart. Scanning electron microscopy investigation revealed that the WSi<sub>2</sub> and ZrSi<sub>2</sub> doping enhanced intergranular connectivity, thus raising  $J_c$  by more than a factor of 2.2 and 3.4, respectively. Moreover, the critical temperature for the doped tapes decreased slightly (less than 0.7 K). It was also found that the improved field dependence of the WSi<sub>2</sub> tapes was due to the pinning by possible segregates or defects caused by the WSi<sub>2</sub> addition. This role of WSi<sub>2</sub> and ZrSi<sub>2</sub> may be very beneficial in the fabrication of MgB<sub>2</sub> tapes for a large range of applications. © 2003 American Institute of Physics. [DOI: 10.1063/1.1600508]

The recent discovery of magnesium diboride (MgB<sub>2</sub>) with its superconducting transition temperature at 39 K has generated much interest in both fundamental research and applications.<sup>1,2</sup> In addition to the chemical simplicity and low cost, a  $T_c$  of 39 K is high enough for use at temperatures (10–20 K) easily accessible to cryocoolers. Therefore, MgB<sub>2</sub> is a possible candidate as a substitute for Nb–Ti or high- $T_c$  oxide superconductors operated at around 20 K. Indeed, although considerable progress has already been made, the best critical current densities  $J_c$  reported so far on MgB<sub>2</sub> wires or tapes fabricated using the powder-in-tube (PIT) method<sup>3–5</sup> are well below those of bulk samples prepared under high pressure,<sup>6,7</sup> due to poor grain connections and the lack of flux pinning centers in this material. Significant improvement of the  $J_c$  is required by either increasing  $B_{c2}$  or introducing pinning centers.

Defects and segregates are known to act as effective flux pinning centers in high- $T_c$  superconductors. Therefore, irradiation defects, chemical doping, the introduction of impurities or substitutions on the Mg or on the B site have been studied in MgB<sub>2</sub> samples.<sup>8–11</sup> Although these were surely effective in increasing inductive critical current density, the improvement in the transport  $J_c$  was not significant. More recently, SiC and Y<sub>3</sub>O<sub>2</sub> doping has been reported to strongly improve the  $J_c$  in MgB<sub>2</sub> superconductors.<sup>12,13</sup> They claimed that a high density of dislocations and nanometer-sized inclusions observed inside MgB<sub>2</sub> grains could act as pinning centers. This leads us to expect that other materials can also be effective in increasing  $J_c$  values. On the other hand, the key applications of MgB<sub>2</sub> superconductors will be for wires and tapes. It is thus highly desirable that synthesis of a highly dense, high critical current capacity phase is possible at ambient pressures to enable large-scale applications.

In this letter, we show that much improvement of  $J_c$  is achieved by the addition of ZrSi<sub>2</sub> and WSi<sub>2</sub> in Fe-sheathed MgB<sub>2</sub> tapes by the PIT process.

Mg (325 mesh, 99.8% in purity), B (325 mesh, amorphous, 99.99%) and 5-at. % WSi<sub>2</sub> or ZrSi<sub>2</sub> powders (2–5  $\mu$ m) were used for the fabrication of tapes by the *in situ* PIT method. The sheath materials chosen for this experiment were commercially available pure Fe. Then the mixture was filled into a Fe tube of 6-mm outside diameter and 1.25-mm wall thickness. After packing, the tubes were subsequently cold-rolled into a rectangular rod using a groove-rolling machine. The rods were subsequently rolled to tapes of 4×0.5 mm. Finally, the tapes were sintered at 600 °C for 1 h in argon. Undoped tapes were also similarly prepared for comparative study. The phase composition and microstructure were investigated by x-ray diffraction (XRD) and scanning electron microscope (SEM). Dc magnetization measurements were performed with a superconducting quantum interference device magnetometer. Magnetization curves were also measured with a vibrating sample magnetometer. The transport  $J_c$  at 4.2 K and its magnetic field dependence were evaluated by a standard four-probe technique with a criterion of 1  $\mu$ V/cm. The  $I_c$  measurement was performed for several tape samples to check reproducibility.

Figure 1 shows the XRD patterns for three tapes. These are undoped, WSi<sub>2</sub>- and ZrSi<sub>2</sub>-doped heat-treated tapes, respectively. Note that the peaks of Fe were contributed from the Fe sheath. It can be seen that all the peaks are attributed to pure MgB<sub>2</sub> in the case of the undoped MgB<sub>2</sub> sample, suggesting the good structure quality. The same was found to be the case for the WSi<sub>2</sub>- and ZrSi<sub>2</sub>-doped samples. However, the WSi<sub>2</sub>-doped MgB<sub>2</sub> tape has additional peaks besides MgB<sub>2</sub> phase, which have been indexed to the WSi<sub>2</sub> compound. Due to the high x-ray scattering factor of the heavy element W, the relative intensities of MgB<sub>2</sub> peaks are smaller than those of WSi<sub>2</sub> peaks. By contrast, the addition of ZrSi<sub>2</sub> leads to the formation of Zr<sub>3</sub>Si<sub>2</sub> and Mg<sub>2</sub>Si as the major impurity phases; there are no peaks corresponding to pure ZrSi<sub>2</sub>, suggesting that there were reaction between MgB<sub>2</sub> and ZrSi<sub>2</sub>.

Figure 2 shows transport  $I_c$  and  $J_c$  as a function of mag-

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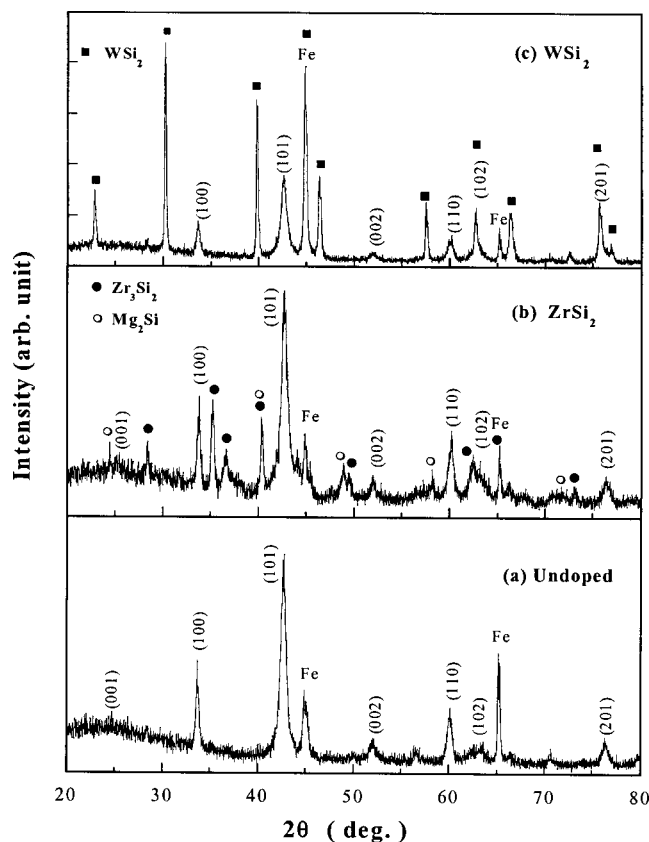


FIG. 1. XRD patterns of undoped,  $\text{WSi}_2$ -, and  $\text{ZrSi}_2$ -doped samples heated at  $600^\circ\text{C}$  for 1 h. The data were obtained after peeling off the Fe-sheath. The XRD peaks of  $\text{MgB}_2$  are indexed, and the peaks of  $\text{WSi}_2$ ,  $\text{ZrSi}_2$ , and  $\text{Mg}_2\text{Si}$  are marked by squares, solid circles and open circles, respectively. The peaks of Fe were contributed from the Fe sheath.

netic field at 4.2 K for our undoped,  $\text{WSi}_2$ - and  $\text{ZrSi}_2$ -doped tapes. Only data above 4 T are shown, because at lower field region,  $I_c$  was so high that strong heating at the current contacts precluded the collection of reliable data. Clearly, both  $\text{WSi}_2$ - and  $\text{ZrSi}_2$ -doped tapes had a much higher  $J_c$  than the undoped tape in magnetic field. The best result was achieved by the  $\text{ZrSi}_2$  doping, at 4.2 K and in a field of 8 T, the  $J_c$  of

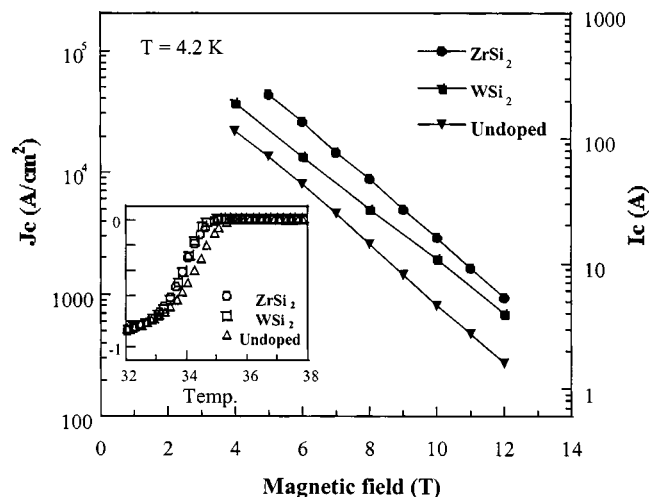


FIG. 2.  $J_c$  and  $I_c$  vs applied field at 4.2 K, for undoped,  $\text{WSi}_2$ -, and  $\text{ZrSi}_2$ -doped samples heated at  $600^\circ\text{C}$  for 1 h, as measured by transport. The measurements were performed in magnetic fields parallel to the tape surface. Inset: ZFC dc magnetization as a function of temperature for all three samples.

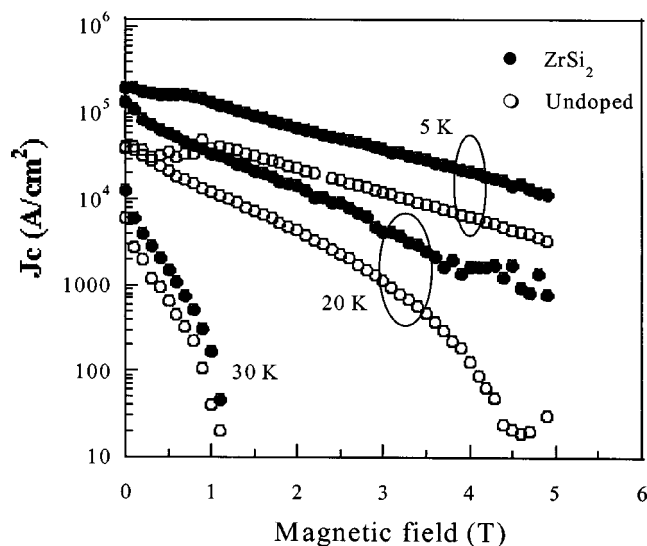


FIG. 3. Magnetic field dependence of  $J_c$  for the undoped and  $\text{ZrSi}_2$ -doped samples heated at  $600^\circ\text{C}$  for 1 h at various temperatures, as measured by magnetization.

$\text{ZrSi}_2$ -doped tapes reached a higher value around  $10 \text{ kA/cm}^2$ , a factor of 3.4 higher than for the undoped one, while for the  $\text{WSi}_2$ -doped tape, at 8 and 10 T,  $J_c$  reached 5 and  $2 \text{ kA/cm}^2$ , respectively (about a factor of 2.2 in the  $J_c$  values). It can also be seen that the  $\text{WSi}_2$  sample revealed a significant improvement in  $J_c$  field performance compared to the undoped tape. Such a difference can be explained in term of higher flux pinning effect. This is demonstrated by the field dependence of volume pinning forces for the undoped and  $\text{WSi}_2$  tapes (not shown here). Although the positions of the maximum pinning force ( $B_{Fp\text{-max}}$ ) for both tapes are the same, the pinning force over  $B_{Fp\text{-max}}$  is apparently larger in the  $\text{WSi}_2$  sample, suggesting that effective pinning centers were possibly introduced with increasing field and thus resulting in the improved field dependence of  $J_c$  in a high-field region. It is worth noting that compared to the undoped tape, the critical temperature for the  $\text{WSi}_2$ - and  $\text{ZrSi}_2$ -doped samples slightly decreased by only 0.5 and 0.7 K, respectively, as shown in the inset of Fig. 2. This shift suggests that incorporation of additive elements into the  $\text{MgB}_2$  actually occurred.

Figure 3 presents  $J_c$  for undoped and  $\text{ZrSi}_2$ -doped samples as a function of magnetic field at various temperatures, as calculated from magnetization loops  $M(H)$  using the critical state Bean model and assuming uniform  $J_c$  flowing through the entire sample. Compared with the transport and magnetization  $J_c$  data, two significant observations are recorded. First, the  $\text{ZrSi}_2$  sample revealed a higher  $J_c$  than the undoped sample at all temperatures and in the entire field region, which is well consistent with the transport  $J_c$ - $B$  curves shown in Fig. 2. Second, the transport  $J_c$  was larger than the magnetic one at a fixed  $T$  and  $H$  (e.g., at 4.2 K and 5 T, respectively) and can be understood in terms of criteria fixed for two independent experiments.

In order to investigate the reason for the  $J_c$  improvement, we studied the differences in the microstructures of the tapes with and without doping. Figure 4 shows the typical SEM images of the fractured core layers for undoped,  $\text{WSi}_2$ - and  $\text{ZrSi}_2$ -doped samples. Clearly, well-developed grains can be seen in all three samples. However, the undoped sample



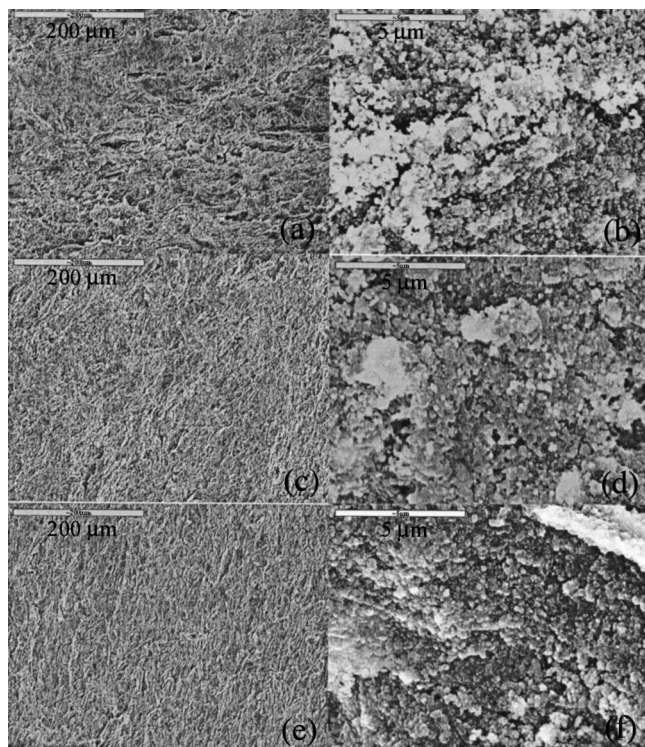


FIG. 4. SEM images of the fractured  $\text{MgB}_2$  core layers of Fe-sheathed tapes, where (a) and (b) are from undoped, (c) and (d) are from  $\text{WSi}_2$ -doped, and (e) and (f) are from  $\text{ZrSi}_2$ -doped samples.

has more voids formed by the evaporation of Mg and the grains are not well connected [Figs. 4(a) and 4(b)]. In other words, weak links exist at the grain boundaries of this sample, and thus introduce a strong limitation to the flow of currents. In contrast, both  $\text{WSi}_2$ - and  $\text{ZrSi}_2$ -doped samples have a higher density with fewer voids, resulting in the better connectivity between the  $\text{MgB}_2$  grains. The best microstructure is observed in the  $\text{ZrSi}_2$  sample [Figs. 4(e) and 4(f)]. However, in contrast to previous work on doping for reducing grain size,<sup>14,15</sup> our high magnification images for all three tapes indicated that the  $\text{MgB}_2$  grain size is almost the same ( $\sim 0.2 \mu\text{m}$ ), which further confirmed that the  $J_c$  difference in all three tapes is not due to the grain-size difference, but due to the improved grain coupling. We should note here that the initial grain size of  $\text{WSi}_2$  and  $\text{ZrSi}_2$  was  $2\text{--}5 \mu\text{m}$ ; however, such large grains were not detected in either sample. From the microstructure data we thus conclude that the main beneficial effect of the doping is the improvement of the inter-grain connectivity. Also, the enhanced field dependence of the  $\text{WSi}_2$  sample is due to the pinning by possible segregates or defects caused by the  $\text{WSi}_2$  addition.

Our results are in consistent with recent reports, in which doping  $\text{MgB}_2$  with Ti and Zr showed an improvement of  $J_c$ , likely due mainly to the densification effect.<sup>14–16</sup> As were demonstrated by several authors,<sup>2,17</sup> the high-density samples have high superconducting homogeneity and strong inter-granular current flow as determined by magneto-optical studies. The fact was also corroborated by recent result of Serquis *et al.*<sup>18</sup> They found that for the  $\text{MgB}_2$  wires prepared by PIT, they could either increase or decrease the  $J_c$  greatly, depending on the certain conditions. Based on microstruc-

tural analyses and magnetization data, they claimed that the main  $J_c$  limiting factor is connectivity, while intragrain pinning, or grain size are issues of secondary order. Clearly, our data of the high performance of  $J_c$  in the whole range of magnetic fields up to 12 T in the  $\text{WSi}_2$  and  $\text{ZrSi}_2$  samples further supports this viewpoint. We should note that the  $\text{WSi}_2$  and  $\text{ZrSi}_2$  grain size used was  $2\text{--}5 \mu\text{m}$ . In the SiC and  $\text{Y}_2\text{O}_3$  addition, which is significantly effective in increasing  $J_c$  in the high-field region, nanometer-sized particles were employed.<sup>12,13</sup> Therefore, further improvement in  $J_c$  is expected upon either utilization of fine  $\text{WSi}_2$  and  $\text{ZrSi}_2$  particles or optimization of the level of  $\text{WSi}_2$  and  $\text{ZrSi}_2$  doping.

In summary, we have investigated the effect of  $\text{WSi}_2$  and  $\text{ZrSi}_2$  addition to  $\text{MgB}_2$  tapes.  $\text{WSi}_2$  and  $\text{ZrSi}_2$  doping improved intergranular connection thus raising  $J_c$  by more than a factor of 2 and 3.4, respectively. For the  $\text{ZrSi}_2$  sample, the  $J_c$  value was up to about  $10 \text{ kA cm}^{-2}$  at 4.2 K and 8 T. Furthermore, the improved field dependence of the  $\text{WSi}_2$  tapes was due to the pinning by possible segregates or defects caused by the  $\text{WSi}_2$  addition. This role of  $\text{WSi}_2$  and  $\text{ZrSi}_2$  may be very beneficial in the fabrication of  $\text{MgB}_2$  tapes for a large number of applications.

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