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Knudsen effusion mass spectrometric determination of the complex vapor composition of samarium, europium, and ytterbium bromides

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RATIONALE: The vaporization of Sm, Eu, and Yb tri- and dibromides is accompanied by decomposition and disproportionation reactions. These result in complex vapor compositions whose analysis is an intricate problem for experimentalists. Approaches have been developed to interpret mass spectra and accurately determine the vapor composition of thermally unstable compounds.

METHODS: A sector type magnet instrument was used. A combined ion source allowed the study of both the molecular and ionic vapor compositions in the electron ionization (EI) and the thermionic emission (TE) modes. The methodological approaches were based on a joint analysis of the ionization efficiency functions, the temperature and time dependences of the ion currents, and special mathematical data evaluation.

RESULTS: The vaporization of SmBr₃, YbBr₃, SmBr₂, EuBr₂, and YbBr₂ was studied in the temperature range of 850–1300 K. An initial stage of incongruent vaporization was observed in the case of the tribromides, SmBr₂, and YbBr₂. This eventually changed to a congruent vaporization stage. Various neutral (Ln, Br, Br₂, LnBr, LnBr₂, LnBr₃, Ln₂Br₄, Ln₂Br₅, and Ln₂Br₆) and charged (Br⁻, LnBr₃, LnBr₄) species were detected at different vaporization stages.

CONCLUSIONS: The quantitative vapor composition of Sm, Eu, and Yb tri- and dibromides was determined. It was found that only EuBr₂ was stable in the studied temperature range. The developed approaches can be useful in the case of other thermally unstable compounds. Copyright © 2013 John Wiley & Sons, Ltd.

It is known that most lanthanide (Ln) tribromides vaporize congruently. Their vapor contains predominantly monomeric molecules LnBr₃, a small fraction (1–10%) of dimer molecules Ln₂Br₆, and a negligibly small amount of larger oligomers (trimers, tetramers, etc.), as well as negative ions (LnBr₃) $_n$ Br⁻ (n=0–4). In contrast, Sm, Eu, and Yb tribromides decompose partly or completely at high temperatures by the reaction:

$$2LnBr_3(s) \rightarrow 2LnBr_2(s) + Br_2(g) \tag{1}$$

In addition, the vapor composition can be more complex and time dependent. This is also the case for the dibromides of Sm, Eu, and Yb which tend to disproportionate at high temperatures by the reaction:

$$3LnBr_2(s) \rightarrow Ln(s)+2LnBr_3(s)$$
 (2)

Unfortunately, reactions (1) and (2) were not taken into account in earlier investigations $^{[6,7]}$ on the vaporization of Sm and Yb bromides by torsion effusion. The authors

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assumed that the vapor consisted mainly of $LnBr_3$ molecules in the case of tribromides and $LnBr_2$ molecules in the case of dibromides. Such assumptions are not justified and may lead to large uncertainties in the derived thermodynamic parameters, e.g. sublimation enthalpies. Therefore, Knudsen effusion mass spectrometry (KEMS) should be used to monitor $in\ situ$ the changes in the vapor composition of these compounds.

In the present work we investigate the vaporization of thermally unstable Sm, Eu, and Yb di- and tribromides by KEMS and establish the occurrence of neutral and charged vapor species under disproportionation/decomposition conditions.

EXPERIMENTAL

Knudsen effusion experiments with mass spectrometric analysis of vapor species were carried out at Ivanovo State University of Chemistry and Technology (Ivanovo, Russia). A MI1201 magnetic sector-type (angle of 90°, curvature radius of 200 mm) mass spectrometer was used. It was supplied by JSC "SELMI" (Sumy, Ukraine) and modified by us for high-temperature studies. A special ion source allowed measurements in the electron ionization (EI) and the thermal ion emission (TE) regimes and the study of both neutral and charged vapor components. In the EI regime, the molecular composition of the equilibrium vapor

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was analyzed over the condensed phase. Mass spectra of the molecular beam from the Knudsen effusion cell were recorded using 50 eV ionizing electrons and a cathode emission current of 1 mA. A movable molecular beam shutter, positioned between the effusion cell and the ionization chamber, made it possible to distinguish the species effusing from the cell from those of the background. In the TE regime, the charged species were identified. In this case, ions were generated inside the effusion cell at a high temperature and drawn out by a weak (10⁴-10⁵ V/m) electric field. The voltage applied to the cell was negative with respect to the ground in order to detect the emission of negative ions. The detection system for the ion current included a secondary electron multiplier combined with a picoammeter (Keithley, Cleveland, OH, USA). The sensitivity for the direct current was 10^{-17} A. The sample under investigation was loaded into a graphite or molybdenum cell. The ratio of the cross-sectional area of the cell to the area of the effusion orifice (0.16 mm²) was about 300. The cell was heated by a resistance oven. The cell temperature was measured by a standard tungsten-rhenium thermocouple calibrated by the melting points of pure NaBr and Ag. The accuracy of the temperature measurement is estimated to be within ± 5 K. The instrument was calibrated with metallic silver as internal standard. An automatic program module^[8] recorded the ion current, the cell temperature, and the energy of the ionizing electrons. A more detailed description of the unit and the experimental procedure was given earlier. [9]

The LnBr₃ (Ln=Sm, Yb) samples were synthesized from the respective oxides Ln₂O₃ (99.9%; Fluka, Buchs, Switzerland) following the NH₄Br route. [10,11] It includes the dissolution of the oxide in concentrated (47%) HBr acid, the addition of NH₄Br, the formation of the ternary ammonium rare earth bromide, and its decomposition to LnBr₃ by heating in vacuum. For further purification, the dry LnBr₃ powders were sublimed in vacuum.

The EuBr $_2$ sample was synthesized from Eu $_2$ O $_3$ (99.9%, Fluka). By the procedure described above, a mixture of EuBr $_2$ and EuBr $_3$ was obtained, which was then decomposed to pure EuBr $_2$ by further annealing at 773 K in vacuum.

The LnBr₂ (Ln=Sm, Yb) samples were obtained by a synproportionation reaction from LnBr₃ and the Ln metal (99.99%, Metall Rare Earth Ltd, Sheung Shui, Hong Kong, China). The metal was used in small excess in a LnBr₃/Ln ratio of 2:1.05 to ensure a complete reduction to the divalent state. The reactions were carried out in welded tantalum containers enclosed in silica ampoules. The quality of all the samples was proven by powder X-ray diffraction. All patterns contained reflections of the respective phase only. SmBr₃^[13] and YbBr₃^[14] crystallized in the PuBr₃ and Bil₃ structure types, respectively. At room temperature, SmBr₂^[15] and EuBr₂^[16] adopt the SrBr₂ structure type and YbBr₂^[17] shows the SrI₂ structure type.

RESULTS AND DISCUSSION

Mass spectra and ionization efficiency curves

Ytterbium and samarium tribromides

In the Yb and Sm tribromide vaporization experiments two main stages relating to the different temperature ranges were observed (denoted by I and II in Figs. 1 and 2, and

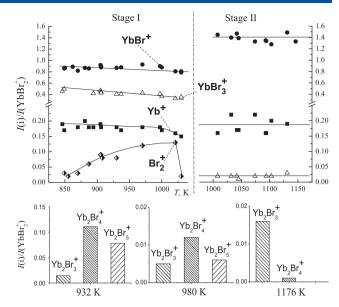


Figure 1. Temperature dependences of the EI mass spectra upon the vaporizations of $YbBr_3$.

Table 1). The experimental data at stage I (Figs. 1 and 2) were obtained by a step-by-step heating; at each short isothermal delay the mass spectra were recorded. At this stage the vaporization was incongruent due to the decomposition reaction (1) which is evident due to the presence of Br⁺ and Br₂⁺ ions in the mass spectra. Because of effusion bromine depletion of the system occurred. Reaction (1) was naturally accelerated with temperature unless an irreversible transition to stage II took place. At this stage the Br₂⁺ and Br⁺ ions vanished completely from the mass spectra and the vaporization demonstrated congruency in the heating/cooling cycles up to the end of experiment. Note that these and all subsequent data were obtained from a graphite cell for which stage I lasted longer at the same temperature than in a

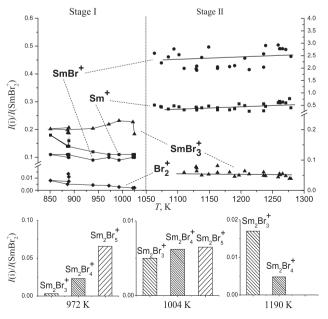


Figure 2. Temperature dependences of the EI mass spectra upon the vaporizations of SmBr₃.



Table 1. EI mass spectra for 50 eV and vapor compositions					
			Relative vapor compositions (%)		
Sample	Vaporization stage, temperature [K]	Ion species in EI mass spectra (relative ion currents)	Neutral species	Ions	
YbBr ₃	Stage I: 850–1050	Yb ⁺ (18), YbBr ⁺ (91), YbBr ₂ ⁺ (100), YbBr ₃ ⁺ (42), Br ⁺ (54), Br ₂ ⁺ (10), Yb ₂ Br ₃ ⁺ (1.4), Yb ₂ Br ₄ ⁺ (11.1), Yb ₂ Br ₅ ⁺ (7.8)	YbBr ₃ (39.1), YbBr ₂ (11.7), Br(35.2), Br ₂ (5.9), Yb ₂ Br ₄ (0.6), Yb ₂ Br ₅ (4.3), Yb ₂ Br ₆ (3.1)	YbBr ₄ (99.9), YbBr ₃ (0.1)	
	Stage II: 1000–1250	Yb ⁺ (18), YbBr ⁺ (139), YbBr ₂ ⁺ (100), YbBr ₃ ⁺ (2), Yb ₂ Br ₃ ⁺ (1.6), Yb ₂ Br ₄ ⁺ (0.1)	YbBr ₂ (95.2), YbBr ₃ (3.8), YbBr(0.5), Yb ₂ Br ₄ (0.48), Yb ₂ Br ₅ (0.02);	YbBr ₄ (83.3) YbBr ₃ (16.7)	
SmBr ₃	Stage I: 850–1050	Sm ⁺ (13), SmBr ⁺ (10), SmBr ₂ ⁺ (100), SmBr ₃ ⁺ (20), Br ⁺ (19), Br ₂ ⁺ (0.5), Sm ₂ Br ₃ ⁺ (0.3), Sm ₂ Br ₄ ⁺ (2.3), Sm ₂ Br ₅ ⁺ (6.6)	SmBr ₃ (57.8), SmBr ₂ (2.3), Br(36.4), Br ₂ (1.2), Sm ₂ Br ₄ (0.1), Sm ₂ Br ₅ (0.6), Sm ₂ Br ₆ (1.7);	SmBr ₄	
	Stage II: 1050–1300	$Sm^+(46)$, $SmBr^+(246)$, $SmBr_2^+(100)$, $SmBr_3^+(6)$, $Sm_2Br_3^+(1.7)$, $Sm_2Br_4^+(0.5)$	SmBr ₂ (91.6), SmBr ₃ (8.2), Sm ₂ Br ₄ (0.1), Sm ₂ Br ₅ (0.1);	SmBr ₄ (83.3), SmBr ₃ (16.7)	
YbBr ₂	960–1300	Yb ⁺ (19), YbBr ⁺ (126), YbBr ₂ ⁺ (100), YbBr ₃ ⁺ (2), Yb ₂ Br ₃ ⁺ (1.6), Yb ₂ Br ₄ ⁺ (0.1)	YbBr ₂ (94.32), Yb(0.94), YbBr(0.47), YbBr ₃ (3.77), Yb ₂ Br ₄ (0.47), Yb ₂ Br ₅ (0.02);	YbBr ₄ (83.3), YbBr ₃ (16.7)	
SmBr ₂	1060–1300	Sm ⁺ (48), SmBr ⁺ (305), SmBr ₂ ⁺ (100), SmBr ₃ ⁺ (3), Sm ₂ Br ₃ ⁺ (1.3), Sm ₂ Br ₄ ⁺ (0.3)	SmBr ₂ (94.04), SmBr ₃ (4.70), SmBr(0.94), Sm(0.14), Sm ₂ Br ₄ (0.14), Sm ₂ Br ₅ (0.04);	SmBr ₄ (80), SmBr ₃ (20)	
EuBr ₂	1049–1261	Eu ⁺ (46), EuBr ⁺ (238), EuBr ₂ ⁺ (100), Eu ₂ Br ₃ ⁺ (2)	EuBr ₂ (99.2), Eu(0.05), EuBr(0.3), Eu ₂ Br ₄ (0.45)	EuBr ₃	

molybdenum cell. As can be seen in Figs. 1 and 2 and Table 1, the mass spectra differed considerably for the different stages. At both stages, the mass spectra differed from those of the lanthanide tribromides. [18–20] The phenomena that stand out in particular are a very high fraction of LnBr⁺ ions and a wide variety of ions containing two lanthanide atoms.

In order to determine the molecular precursors of the ions at each vaporization stage, we recorded the ionization efficiency curves (IECs), which show the dependence of the mass spectra on the energy of the ionizing electrons. IECs for the Ln⁺, LnBr⁺, LnBr⁺, LnBr⁺, and Br⁺ ions were measured in the course of one experiment and normalized for a fixed electron energy (4–5 eV above the threshold; see Figs. 3 and 4). The energy scale in Figs. 3 and 4 was calibrated using the ionization energy of molecular bromine; $I_0(Br_2) = 10.517 \pm 0.003$ eV.^[21] It should be noted that the ion appearance energies were not determined precisely in this work, because such a determination from the superimposed spectra of the individual molecules is complicated. For the interpretation of the mass spectra, we therefore focus on the analysis of the shapes of the IECs.

It can be seen from Fig. 3 that the shapes of the IECs for Yb⁺ and YbBr₂⁺ ions differ substantially for the temperature ranges of 877–988 K and 1067–1174 K that correspond to the different stages of vaporization (Fig. 1). It is obvious to draw the conclusion that at stage I (877–988 K) these ions are mainly formed from the YbBr₃ molecules. In addition, the low-energy tails extending to about 15 eV (Yb⁺) and 9

eV (YbBr₂⁺) demonstrate the presence of YbBr₂ molecules. These dominate in the vapor of stage II (1067–1174 K). At this stage the intensities of the YbBr₃⁺ ions were too low for IEC measurements. Note that the YbBr₃⁺ ions are exclusively formed from YbBr₃ molecules, in agreement with previous reports on other LnBr₃ molecules. However, the shape of the IEC of the YbBr⁺ ions at stage II changes only insignificantly compared with those of the Yb⁺ and YbBr₂⁺ ions. This observation can be explained if the YbBr₂ molecules are the main precursor for the YbBr⁺ ions at any stage of vaporization. In addition, the low-energy tail, which appears on the YbBr⁺ IEC at 1067–1174 K, indicates some contribution from YbBr molecules.

A similar analysis was performed for the IECs of samarium-containing ions (Fig. 4). In contrast to YbBr⁺, the SmBr⁺ ions are mainly formed from SmBr₃ molecules at stage I and from SmBr₂ at stage II. The larger amount of tribromide molecules at stage I is evidence for the higher thermodynamic stability of SmBr₃ than of YbBr₃, and *vice versa* the higher stability of YbBr₂ than of SmBr₂. In addition, no low-energy tail is observed for the SmBr⁺ curve, i.e. there are no SmBr molecules in the vapor, thus providing support for this conclusion.

The variation of the fractions of two Ln atoms containing ions at the different vaporization stages (see Figs. 1 and 2) correlates with the respective LnBr $_3$ and LnBr $_2$ contributions to the ion currents. This observation allows the Ln $_2$ Br $_3^+$, Ln $_2$ Br $_4^+$, and Ln $_2$ Br $_5^+$ ion currents to be attributed to the Ln $_2$ Br $_4$, Ln $_2$ Br $_5$, and Ln $_2$ Br $_6$ molecules.



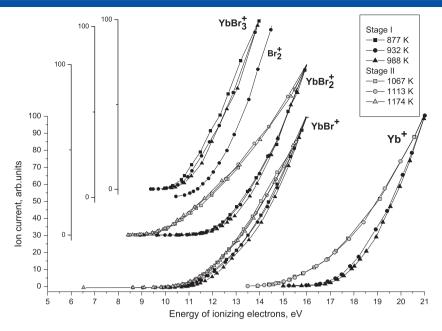


Figure 3. Ionization efficiency curves at stages I and II of vaporization of YbBr₃.

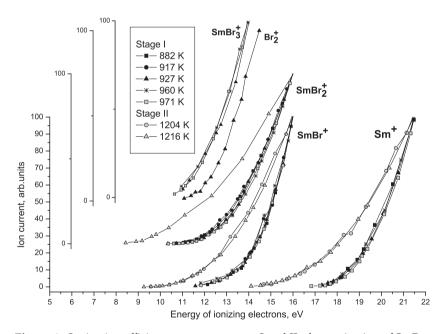


Figure 4. Ionization efficiency curves at stages I and II of vaporization of SmBr₃.

Ytterbium, samarium, and europium dibromides

The ion currents of the species recorded on heating in EI mass spectra are summarized in Table 1. The ion current ratios are shown in Figs. 5, 6, and 7 for YbBr₂, SmBr₂, and EuBr₂, respectively. The observation of YbBr₃⁺ and SmBr₃⁺ ions in the spectra indicates the presence of the LnBr₃ molecules which are formed, in our opinion, by the disproportionation of the dibromides according to reaction (2). It is noteworthy that the mass spectra of the Yb and Sm dibromides are quite similar to those of the corresponding tribromides at stage II.

It can be seen from Fig. 5 that the slope of the $I(Yb^+)/I(YbBr_2^+)$ curve changes significantly around 1030 K. Such a behavior reflects the competition between different contributions to the

 Yb^+ ion current and follows the change in shape of the Yb^+ IECs (see Fig. 8). These contributions are predominately from Yb atoms and $YbBr_2$ molecules. There are also two contributions to the $YbBr^+$ ion current. The main one comes from $YbBr_2$ and the minor one from YbBr molecules. The fraction of YbBr molecules in the vapor is estimated to be about 0.5%. The situation is similar in the case of $YbBr_2^+$ ions which are mainly formed from $YbBr_2$ molecules and to a much smaller fraction of about 4% from $YbBr_3$ molecules.

The IECs of the ions from SmBr₂ (see Fig. 9) are quite similar to those from YbBr₂ with the major difference that the Sm⁺ ion current shows no temperature dependence. A possible reason might be the lower vapor pressure of metallic Sm than of Yb in the investigated temperature range.



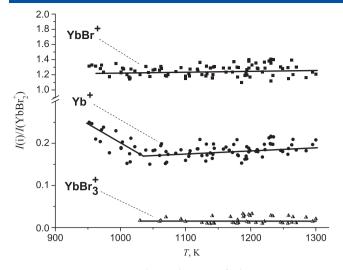


Figure 5. Temperature dependences of the mass spectra upon the vaporization of YbBr₂.

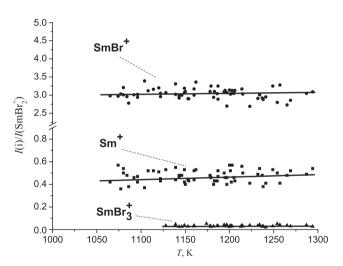


Figure 6. Temperature dependences of the mass spectra upon the vaporization of $SmBr_2$.

The mass spectrum of europium dibromide is distinctly different from those of the ytterbium and samarium dibromides (see Table 1), in that the undetectable $EuBr_3^+$ and $Eu_2Br_4^+$ ions prove the absence of $EuBr_3$ molecules in the vapor within instrumental sensitivity. Accordingly, the decomposition of $EuBr_2$ by reaction (2) is negligible. Nevertheless, the IECs of the Eu^+ and $EuBr^+$ ions demonstrate low-energy tails similar to the other dibromides (see Fig. 10). They correspond to small fractions of gaseous Eu and EuBr, of about 0.05% and 0.3%, respectively.

Vapor composition

The observations described above and our interpretations suggest that the vapor composition of all the investigated bromides is complex. The next stage of data processing therefore involves the separation of the contributions from different molecular precursors to the ion currents. Let us introduce the concept of fragmentation coefficients:

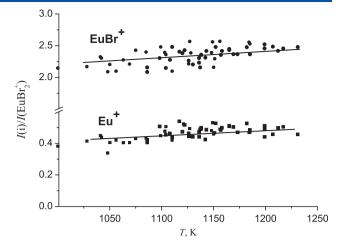


Figure 7. Temperature dependences of the mass spectra upon the vaporization of EuBr₂.

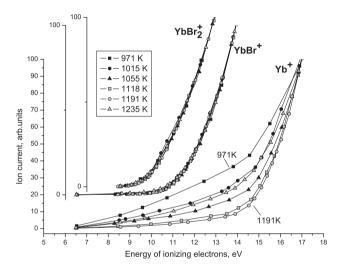


Figure 8. Ionization efficiency curves upon the vaporization of YbBr₂.

$$f_{ij} = I_{ij}/I_{jj},\tag{3}$$

which determine the ratio between the $LnBr_i^+$ fragment and $LnBr_j^+$ molecular ion currents formed from a $LnBr_j$ molecule (i < j). The ion currents I_{03} , I_{13} , and I_{23} refer to those of the fragment ions Ln^+ , $LnBr^+$, and $LnBr_2^+$ originating from a $LnBr_3$ molecule. Accordingly, I_{33} refers to $LnBr_3^+$ from $LnBr_3$. This results in Eqn. (4) for the fragment ions:

$$I_{i3} = f_{i3}I_{33}, (i = 0, 1, 2).$$
 (4)

Likewise, for LnBr₂ molecules, we obtain the expressions:

$$I_{i2} = f_{i2}I_{22}, (i = 0, 1).$$
 (5)

Thus, the coefficients f_{03} , f_{13} , f_{23} , f_{02} , and f_{12} attribute the ion currents to molecular precursors. With this in mind, we consider the balance equations of ion currents measured upon vaporization of LnBr₃ at stage I:

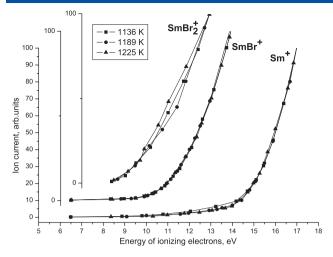


Figure 9. Ionization efficiency curves upon the vaporization of ${\sf SmBr}_2$.

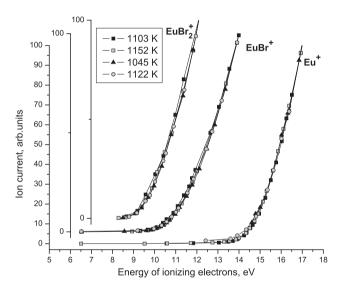


Figure 10. Ionization efficiency curves upon the vaporization of ${\rm EuBr_2}$.

$$I_2 = I_{22} + f_{23}I_3, (6)$$

$$I_1 = f_{13}I_3 + f_{12}I_{22}, (7)$$

$$I_0 = f_{03}I_3 + f_{02}I_{22}, (8)$$

where I_i is the measured ion current.

Fragmentation coefficients for $LnBr_2$ molecules can be expressed as:

$$f_{12} = \frac{I_1 - f_{13} I_3}{I_2 - f_{23} I_3},\tag{9}$$

$$f_{02} = \frac{I_0 - f_{03} I_3}{I_2 - f_{23} I_3}. (10)$$

Equations (9) and (10) are valid for each of the experimental data sets at stage I if one neglects the temperature dependence of f_{02} and f_{12} . They can therefore be written for two various temperatures as:

$$f_{02}(T_1) = f_{02}(T_2), \tag{11}$$

$$f_{12}(T_1) = f_{12}(T_2). (12)$$

Then

$$\frac{I'_{1}f_{13}I'_{3}}{I'_{2}f_{23}I'_{3}} = \frac{I''_{1}f_{13}I''_{3}}{I''_{2}f_{23}I''_{3}},\tag{13}$$

$$\frac{I'_{0}f_{03}I'_{3}}{I'_{2}f_{23}I'_{3}} = \frac{I''_{0}f_{03}I''_{3}}{I''_{2}f_{23}I''_{3}},\tag{14}$$

where I_i and I_i are the ion currents for the first and second random points.

A redundant system of type (13), (14) was devised for the experimental data on ion currents measured at different temperatures, and solved using Mathcad software. The values f_{03} , f_{13} , f_{23} , f_{02} , and f_{12} were found by a least-squares procedure (see Table 2).

The fractions of vapor species (Table 1) were calculated using the relationship:

Table 2. Fragmentation coefficients of LnBr ₂ and LnBr ₃ molecules					
	Vi	alue**			
Coefficient*	Yb	Sm			
$f_{02} = I(Ln^+, LnBr_2)/I(LnBr_2^+, LnBr_2)$ $f_{12} = I(LnBr^+, LnBr_2)/I(LnBr_2^+, LnBr_2)$ $f_{03} = I(Ln^+, LnBr_3)/I(LnBr_3^+, LnBr_3)$ $f_{13} = I(LnBr^+, LnBr_3)/I(LnBr_3^+, LnBr_3)$ $f_{23} = I(LnBr_2^+, LnBr_3)/I(LnBr_3^+, LnBr_3)$	0.2 ± 0.1 1.2 ± 0.2 0.4 ± 0.1 1.2 ± 0.3 2.0 ± 0.3	$\begin{array}{c} 0.6 \pm 0.2 \\ 3.5 \pm 0.6 \\ 0.7 \pm 0.1 \\ 0.5 \pm 0.1 \\ 5.0 \pm 0.5 \end{array}$			

^{*}The term I(A, B) denotes the intensity of species A formed from species B, e.g. $I(Ln^+, LnBr_2)$ is the intensity of the current of Ln^+ ions formed from $LnBr_2$ molecules.

^{**}Standard deviation follows the \pm sign.



$$P_{j} \sim \frac{T}{\sigma_{j}} \sum_{i} \frac{I_{i}}{\gamma_{i} a_{i}},\tag{15}$$

where p_j is the partial pressure, T is the cell temperature, σ_j^{mol} is the total ionization cross section of the jth molecule with the working energy of ionizing electrons (calculated on the basis of ionization cross sections σ_n^{at} of atoms^[22] n using the expression $\sigma_j^{\text{mol}} = 0.75 \sum_n \sigma_n^{\text{at}}$, σ_j^{at} is the total ion current of ions i of all types formed from molecule j (calculated on the basis of the resultant fragmentation coefficients, Table 2), σ_j^{at} is the coefficient taking into account the natural abundance of isotopes of the measured ion, and σ_j^{at} is the coefficient of ion-electron conversion (it is assumed that $\sigma_j^{\text{at}} \sim M_i^{-1/2}$, σ_j^{at}) where σ_j^{at} is the mass of the ion).

Negative ions in the vapor of Yb, Sm, and Eu bromides

The negative ions in the TE mass spectra of Yb, Sm, and Eu bromides are listed in Table 1. If one considers the $LnBr_4^-$ and $LnBr_3^-$ ions as a combination of $LnBr_3^-$ and $LnBr_2^-$ molecules and Br^- anions, the relative content of these ions in the vapor may, correspondingly, be used as an additional instrument in the analysis of the molecular constituents of the vapor phase.

In the vaporization of YbBr₃ and SmBr₃ the LnBr₄ ions were discovered at stage I. Upon transition to stage II their ion currents dropped drastically (see Fig. 11). This observation confirmed the occurrence of reaction (1) as was inferred above throughout the analysis of neutral vapor moieties. LnBr₃ ions were seen in small amounts at stage I only in the case of ytterbium tribromide. Their fraction gradually increased towards stage II. The concentrations of LnBr₃ and LnBr₄ ions at stage II are of the same order of magnitude and correspond to those observed for the vaporization of LnBr₂. These changes in the temperature dependence of negative ion mass spectra are shown in Fig. 11. An additional argument in parallel with the molecular analysis in the SmBr₃ case is the detection of SmBr $_3$ ions only at stage II (T > 1200 K). This is a result of the higher decomposition resistance of SmBr₃ than of YbBr₃.

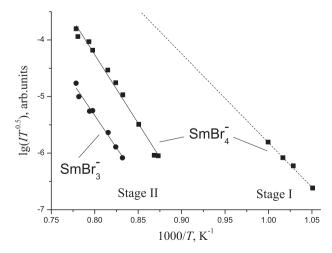


Figure 11. Temperature dependences of the TE mass spectra and their change over time upon the vaporations of SmBr₃.

Although there were molecules of ytterbium and samarium monobromides in the vapor, none of the experiments with LnBr₃ and LnBr₂ revealed the presence of LnBr₂ ions. Also, these ions were not recorded in additional experiments with LnBr₂–Ln systems.

In the case of europium dibromide vaporization only $EuBr_3^-$ ions were recorded. Again this is a confirmation of the above conclusion that the disproportionation of $EuBr_2$ is negligible.

CONCLUSIONS

The vaporization of SmBr₃, YbBr₃, SmBr₂, EuBr₂, and YbBr₂ was studied in the temperature range of 850-1300 K. The mass spectra were interpreted by the joint analysis of the ionization efficiency functions, the temperature and the time dependences of the ion currents. It was found that only EuBr₂ vaporized congruently. In the case of the tribromides, SmBr₂ and YbBr2, an initial stage of incongruent vaporization was observed. Various neutral (Ln, Br, Br2, LnBr, LnBr2, LnBr3, Ln₂Br₄, Ln₂Br₅, and Ln₂Br₆) and charged (Br⁻, LnBr₃⁻, and LnBr₄) species were detected at different vaporization stages. The quantitative vapor composition of Sm, Eu, and Yb triand dibromides was determined. A technique for mass spectra interpretation and accurate determination of the vapor composition of the thermally unstable bromides was developed. This technique is based on a solution of the equations incorporating the fragment ion currents and corresponding fragmentation coefficients. The developed approaches can be extended for characterization of the vaporization of other thermally unstable compounds.

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