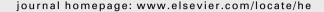
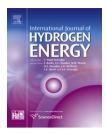


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The reactor design for photoelectrochemical hydrogen production

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ABSTRACT

The heat transfer and flow characteristics of a photoelectrochemical (PEC) hydrogen generation reactor are investigated numerically. Four different reactor designs are considered in this study. The solar irradiation is separated into short and long wavelength parts depending on the energy band gap of the photoelectrode used. While short wavelength part is used to generate electron and hole pairs, the long wavelength part is used to heat the system. Because the energy required for splitting water decreases as temperature is increased, heating the reactor by using the long wave energy increases the system efficiency. Thus, how the long wavelength energy is absorbed by the reactor is very important.

The results show that more long wavelength energy kept inside the reactor can increase the solar-to-hydrogen efficiency, η_{SH} . For Fe₂O₃ photoelectrode, careful reactor design can increase η_{SH} by 11.0%. For design D under 4000 W/m² irradiation and a quantum efficiency of 30%, η_{SH} is found to be 14.1% and the hydrogen volume production rate is 166 L/m² h for Fe₂O₃. Effects of several parameters on the PEC hydrogen reactor are also discussed. Copyright © 2011, Hydrogen Energy Publications, LLC. Published by Elsevier Ltd. All rights

1. Introduction

Hydrogen is regarded as a potential energy carrier for the future. There exist several methods to produce hydrogen from various sources. One promising method is photoelectrochemical (PEC) decomposition of water into hydrogen and oxygen. Both the reactant of water and the energy source from the sun are readily available, renewable, and environmentally clean. Since Fujishima and Honda [1] proved that water can be effectively split into its constituents on a photoelectrode covered with n-type ${\rm TiO}_2$ semiconductor, considerable attention has been devoted to the photodecomposition of water as an alternative solar energy conversion process leading to a non-polluting fuel [2–5].

A PEC cell consists of a semiconductor electrode and a metal counter electrode immersed in an aqueous electrolyte. As light is incident on the semiconductor, it absorbs part of the light and generates electron and hole pairs. These electron and hole pairs react with water and split it into hydrogen and oxygen.

The process can be summarized in the following equations [6,7]:

$$Solar: 2h\nu \rightarrow 2e^- + 2h^+ \tag{1}$$

Anode:
$$2h^+ + H_2O_{(l)} \rightarrow \frac{1}{2}O_{2(g)} + 2H_{(aq)}^+$$
 (2)

Cathode:
$$2e^- + 2H_{(aq)}^+ \rightarrow H_{2(g)}$$
 (3)

Overall:
$$2h\nu + H_2O_{(l)} \rightarrow \frac{1}{2}O_{2(g)} + H_{2(g)}$$
 (4)

Licht [8–10] recently proposed a model to increase the hydrogen production efficiency. While the high frequency solar irradiation is used for photo-electronic conversion, the low frequency sub-band-gap part is used for heating the electrolyte solution to lower the water-splitting potential. This provides a process of efficient water splitting by using the solar energy.

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$ \begin{array}{llllllllllllllllllllllllllllllllllll$	I_{λ} spectral radiative intensity $I_{b\lambda}$ black body spectral radiative intensity k thermal conductivity k_A thermal conductivity of insulator QE quantum efficiency T temperature W distance between anode and glass wall average absorption coefficient spectral absorption coefficient α_{λ} spectral absorption coefficient β thermal expansion coefficient η_{SH} solar-to-hydrogen conversion efficiency λ wavelength ν frequency ρ hydrogen volume production rate
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Grzegorz et al. [11] investigated the effects of electrolyte concentration on the performance of a PEC cell. They made a two-compartment PEC cell consisting of a CdS photoanode immersed in aqueous sulfide solution, Nafion membrane, platinum cathode, and sulphuric acid solution as the dark-compartment electrolyte. Under their optimized conditions, light to hydrogen conversion efficiency up to 12% was completed under sun light illumination.

A semiconductor is characterized by the energy band gap, E_g . Only that part of incident solar radiation having energy higher than E_g can be absorbed by the semiconductor electrode to generate electron—hole pairs. Accounting for the resistance and other losses, the optimum value of E_g is approximately 2.0 eV [12]. Materials having E_g lower than this value are less effective for splitting water.

In 2006, Kudo [13,14] reported a quantum efficiency of 20% system. Experimentally, Pt-loaded AgInZn₇S₉ ($E_g = 2.3 \text{ eV}$) immersed in 0.5 M K₂SO₃ showed the highest activity for the H₂ production. The quantum efficiency at 420 nm was 20% using a 300 W Xe lamp, and the rate of H₂ production was 944 µmol/h for 33 cm² irradiated area. Its hydrogen volume production rate was 6.41 L/m² h.

Wang et al. [15] used a one-step electrodeposition method to produce $AgInS_2$ electrode. They obtained a maximum photocurrent density of 9.28 mA/cm² with an applied bias of +1.0 V vs. an Ag/AgCl electrode in contact with electrolyte containing 0.25 M K_2SO_3 and 0.35 M Na_2S . The maximum quantum efficiency was 63% at 560 nm.

Tseng et al. [16] performed a thermodynamic analysis of the PEC hydrogen production for air mass 1.5 (AM 1.5) solar irradiation. The maximum hydrogen production rate and solar-to-hydrogen conversion efficiency under various operation conditions were presented and discussed.

So far, most studies on PEC production of hydrogen have focused on the development of photoelectrode materials. Very little attention has been paid to the reactor design and thermal and transport phenomena analysis. Both Licht's work [8–10] and our previous work [16] were based on thermodynamic consideration only. No discussion on the transport phenomena and reactor design was given. It is the purpose of the present work to theoretically investigate the heat transfer and flow characteristics of PEC reactors and study the effects of several

parameters on the temperature of the system, the solar-to-hydrogen efficiency, $\eta_{\rm SH}$, and the hydrogen production rate.

2. Analysis

The schematic diagram of the solar hydrogen reactor considered in this work is depicted in Fig. 1. The solar incidence is directly applied to the reactor. The size of the anode is fixed at 100 mm high and 2 mm thick. The front wall of the reactor is a low absorptive glass. A glass wall is also used as the back wall. The thickness of the back glass wall is designated as G. To hold more thermal energy inside the reactor, an insulation layer of thickness A is used in some cases. The anode is made of fluorine-doped tin oxide (FTO) glass coated with semiconductor films such as Fe₂O₃ or TiO₂. Four designs are studied in this work. Design A: G = 5 mm and A = 0 mm; Design B: G = 5 mm and A = 0 mm, the outside face of the glass is coated opaque; Design C: G = 0 mm and A = 20 mm (thermal conductivity $k_A = 0.05$ W/m K); Design D: G = 5 mm and A=20 mm. The spectral solar radiation (AM 1.5) [17] and the measured spectral absorption coefficients of the FTO and glass are shown in Fig. 2. These values are used in the

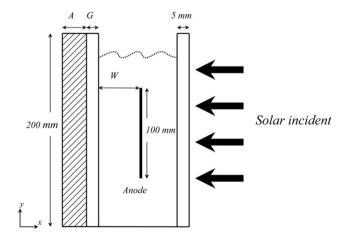


Fig. 1 – The schematic diagram of solar hydrogen reactor. A: thickness of insulator, G: thickness of glass, W: distance between the anode and the glass wall.

numerical simulation. Although electrolytes such as potassium hydroxide solution (KOH) are often used in a practical reactor, water is considered in this study for its availability of optical and thermal properties [18–20].

Steady state, laminar, incompressible and Newtonian flow is considered, and the Boussinesq approximation is employed in this work. The governing equations for this system can be expressed as:

Continuity equation:

$$\frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} = 0 \tag{5}$$

x-momentum equation:

$$u\frac{\partial(\rho u)}{\partial x} + v\frac{\partial(\rho u)}{\partial y} = -\frac{\partial p}{\partial x} + \mu \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}\right) \tag{6}$$

y-momentum equation:

$$u\frac{\partial(\rho v)}{\partial x} + v\frac{\partial(\rho v)}{\partial y} = -\frac{\partial p}{\partial y} + \mu \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2}\right) - \rho g \beta (T - T_0) \tag{7}$$

Energy equation:

$$u\frac{\partial \left(\rho c_{p} T\right)}{\partial x} + v\frac{\partial \left(\rho c_{p} T\right)}{\partial y} = k\left(\frac{\partial^{2} T}{\partial x^{2}} + \frac{\partial^{2} T}{\partial y^{2}}\right) - \nabla \cdot q^{r} \tag{8}$$

The radiative transfer equation (RTE) [21]:

$$\frac{\mathrm{d}I_{\lambda}}{\mathrm{d}s} = \alpha_{\lambda}I_{b\lambda}(T) - \alpha_{\lambda}I_{\lambda} \tag{9}$$

$$\nabla \cdot q^{r} = \int_{0}^{\infty} \alpha_{\lambda} \left(4\pi I_{b\lambda} - \int_{A_{\tau}} I_{\lambda} d\Omega \right) d\lambda \tag{10}$$

where $I_{b\lambda}$ is the black body intensity. The effect of scattering is not considered in this work.

Except the bottom wall, which is assumed adiabatic, the heat transfer between the reactor and the surroundings is considered in the work. Computational domain is extended 1.5 times the height of the reactor in the lateral and upward directions so that interactions between the reactor and the surroundings can be simulated. The SIMPLEC algorithm [22] is

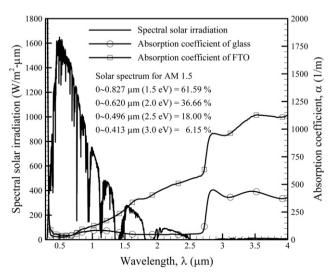


Fig. 2 — The spectral solar irradiation and the absorption coefficient of glass and FTO.

employed in this work. Commercial CFD software package, FLUENT 6.3, is used as the numerical tool in this study. This software solves the Navier—Stokes equations using an unstructured finite-volume method and the RTE using the discrete ordinates method [23,24]. Some numerical techniques are used through the user-defined function (UDF).

Based on the AM 1.5, the fraction of solar energy that is available for water splitting is integrated over the wavelength as a function of E_g , as shown in Fig. 2. The percentage of short wave energy is 61.6% for $E_g=1.5\,$ eV, 36.7% for $E_g=2.0\,$ eV, 18.0% for $E_g=2.5\,$ eV, and 6.15% for $E_g=3.0\,$ eV. Properties of water used in the calculation such as density, heat capacity, thermal conductivity, viscosity and thermal expansion coefficient are functions of temperature. Some modelling techniques [25–29] are used in this work.

The spectral absorption coefficient of water, FTO and glass are averaged for the short and long wave regions respectively using:

$$x_{\text{ave}} = \frac{\int\limits_{\lambda_{1}}^{\lambda_{2}} \alpha_{\lambda} E_{\lambda} d\lambda}{\int\limits_{\lambda}^{\infty} E_{\lambda} d\lambda}$$
 (11)

3. Results and discussion

The solar incidence is divided into short and long wavelength parts according to E_g of the photoelectrode. The short wavelength part is directed to the photoanode for the generation of electro-hole pairs, and the long wavelength part is used for

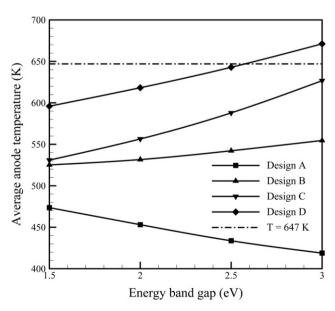


Fig. 3 – Average anode temperature for the four different reactor designs with W=10 mm, solar incident flux = 4000 W/m² and QE = 30%. (1)Design A: G=5 mm and A=0 mm; (2) Design B: G=5 mm and A=0 mm, the outside face of glass is coated opaque; (3) Design C: G=0 mm and A=20 mm ($k_A=0.05$ W/m K); (4) Design D: G=5 mm and A=20 mm.

heating the electrolyte. Four energy band gaps, 1.5 eV, 2.0 eV, 2.5 eV, and 3.0 eV are considered as examples in this work. Representative numerical results of hydrogen volume production rate and $\eta_{\rm SH}$ will be presented and discussed. Because the design of PEC reactors has rarely been studied previously, the attention of this work will be given to the effects of the reactor design. Effects of the solar irradiation, glass thickness G, the distance between the anode and glass wall W and the quantum efficiency QE will also be investigated. The insulator used as an example is the rock wool, with a thermal conductivity of 0.05 W/m K. A is fixed at 20 mm. To prevent the water from boiling, the operating pressure is varied according to the maximum temperature inside the reactor.

3.1. Effects of reactor design

Fig. 3 shows the average anode temperature versus E_g of the photoelectrode for the four reactor designs with W=10 mm and solar incident flux = $4000 \, \text{W/m}^2$ (concentration ratio of 4). The average anode temperature of design A decreases as E_g increases. On the contrary, the average anode temperature of designs B, C and D increases as E_g increases. For design A, the left wall is semi-transparent to the long wave energy, so portions of the solar irradiation go through the water, the glass wall, and then escape to the outside. As E_g is increased, less

energy is used for the generation of electro-hole pairs, and more energy is lost to the outside through the long wavelength part. Therefore, the average anode temperature decreases. On the other hand, for designs B, C, and D, the left wall is made opaque either by the insulation or coating. Therefore, very few radiant energy is lost to the outside. As E_g is increased, more energy is used for heating, and hence higher average anode temperature.

Since the long wavelength energy is used for heating the reactor, it is an essential subject to keep as much long wavelength energy inside the reactor as possible. The best among the four is design D. It has an opaque insulator to cut both conduction and radiation losses, while the left glass wall acts as a long wavelength energy storage due to its good absorptivity in that wavelength region. For solar incident flux of 4000 W/m² and quantum efficiency QE of 30%, the average anode temperature is 623 K and 686 K respectively for Fe₂O₃ (E_g = 2.1 eV) and TiO₂ (E_g = 3.2 eV) in design D. But with design A, the average anode temperature is only 449 K and 411 K respectively for Fe₂O₃ and TiO₂.

Fig. 4 shows the reactor temperature distributions for designs A, B, C and D with anodes of different E_g for solar incident flux = 4000 W/m² and QE = 30%. In the PEC reactor, the short wavelength energy is directed to the photoanode for splitting water. QE = 30% means the remaining 70% of

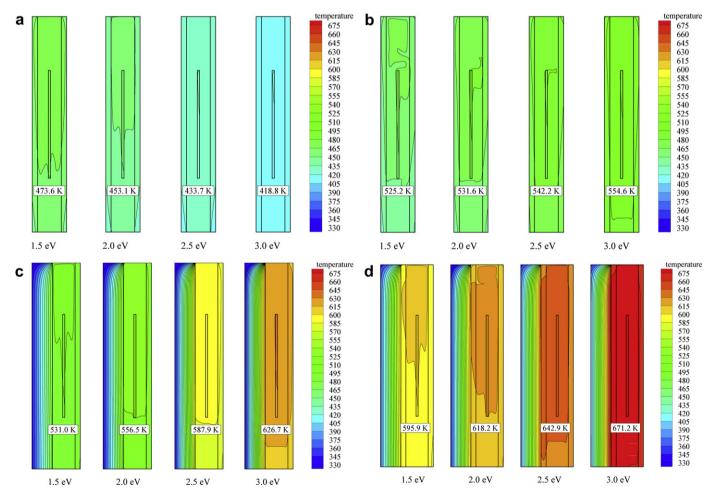


Fig. 4 – Temperature distributions for the four different reactor designs with anodes of different E_g , W = 10 mm, solar incident flux = 4000 W/m² and QE = 30%.

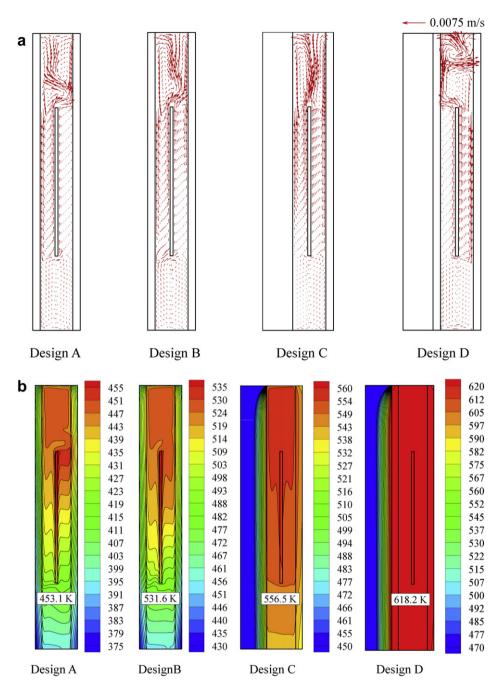


Fig. 5 – The flow and thermal characteristics of the four different reactor designs for $E_g = 2.0$ eV, solar incident flux = 4000 W/m² and QE = 30%, (a) velocity vector; (b) temperature distribution.

absorbed short wavelength energy is converted to heat. As E_g increases, this converted thermal energy decreases. Because in design A, much of the long wavelength energy is lost to the outside, the converted thermal energy in the short wavelength part plays more important role than in other designs. For design A, the average anode temperature is 474 K, 453 K, 434 K and 419 K respectively for $E_q = 1.5$ eV, 2.0 eV, 2.5 eV and 3.0 eV.

In design B, the outside face of the left glass wall is coated opaque to prevent the loss of long wavelength radiant energy through the glass. The average anode temperature increases slightly as E_q is increased. For design B, the average anode

temperature is found to be 525 K, 532 K, 542 K and 555 K respectively for $E_a = 1.5$ eV, 2.0 eV, 2.5 eV and 3.0 eV.

In design C, compared with design B, the left glass wall is replaced by a 20 mm thick insulator ($k_A=0.05~\rm W/m~K$) to further cut the heat loss to the outside by radiation and conduction. As a result, the energy remained inside the reactor is higher than that in design B. For design C, the average anode temperature is 531 K, 557 K, 588 K and 627 K respectively for $E_q=1.5~\rm eV$, 2.0 eV, 2.5 eV and 3.0 eV.

In design D, compared with design C, a 5 mm thick glass is placed inside the insulator to absorb the long wavelength

energy. For $E_g=1.5$ eV, about 85% of the long wavelength energy is absorbed by the glass wall. The absorbed energy raises the temperature of the glass, and re-emits to heat the water and electrode. As E_g is increased, more long wavelength energy is available for heating the reactor. The average anode temperature is 596 K, 618 K, 643 K and 671 K for $E_g=1.5$ eV, 2.0 eV, 2.5 eV and 3.0 eV, respectively.

Fig. 5(a) and (b) presents the flow and thermal characteristics of the four different reactor designs for $E_g=2.0$ eV, solar incident flux = 4000 W/m² and QE = 30%. The photoanode has the highest temperature due to the absorption of solar irradiation.

Comparing with the other three designs, design A has the lowest temperature because most of the long wavelength energy is lost to the outside. The fluids in the two channels separated by the anode have different velocities. The right-hand-side channel (channel R) has smaller velocity due to more pronounced thermally stratified phenomenon. The fluid therefore flows to the right when it passes the top of the anode plate. The flow in design B is similar. However, the velocity in channel R is higher than that in design A due to lesser extent of thermal stratification. The distortion of the flow pattern is not as severe.

For design C, a thermal insulation layer is used, and the temperature inside the reactor is increased and distributed more uniformly. The flow velocities near both sides of the anode plate are almost the same, and two nearly symmetric circulations form in the top region.

For design D, the reactor temperature is further increased. The temperature on the right surface of the anode is slightly higher than that on the left surface due to direct irradiation of solar flux. Therefore, the velocity in the right channel is higher than that in the left channel, and the flow turns to the left as it passes the top of the anode plate.

Fig. 6 shows $\eta_{\rm SH}$ for the four designs. $\eta_{\rm SH}$ is defined as

$$\eta_{SH} = \frac{QE \times E_s \times \Delta H_{H_2}(298 \text{ K}, 1 \text{ bar}) / \Delta G_{H_2O}(T, p)}{E_{in}}$$
(12)

where E_s is the energy higher than E_g , and $E_{\rm in}$ is the total incident solar energy. $\Delta H_{\rm H_2}$ is the heating value of H_2 at the standard state and $\Delta G_{\rm H_2O}$ is the Gibbs free energy of water dissociation reaction at reactor temperature and pressure.

For solar incident flux of 4000 W/m² and QE of 30%, η_{SH} is 14.1% and 2.6% respectively for Fe₂O₃ and TiO₂ with design D, and 12.7% and 2.2% with design A. Comparing designs A and D, the enhancement of η_{SH} is respectively 11% and 18% for Fe₂O₃ and TiO₂.

3.2. Effects of solar incident flux

To study the effects on the hydrogen volume product rate and solar-to-hydrogen efficiency, the solar incident flux is assumed to be 1000 W/m² for the unconcentrated case, and 2000 W/m², 3000 W/m² and 4000 W/m² for the cases using a solar concentrator. The average anode temperatures under different solar incident flux versus E_g with W=10 mm, G=10 mm and QE=30% for design D are shown in Fig. 7. As expected, as solar incident flux increases, the average anode temperature increases. For solar incident flux = 4000 W/m², the average anode temperature is close to 647 K.

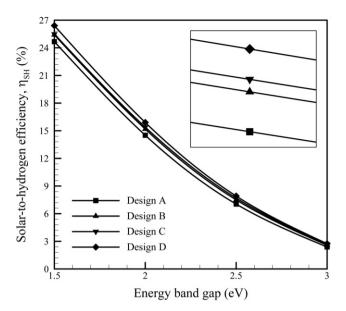


Fig. 6 – Solar-to-hydrogen efficiency for the four different reactor designs with W = 10 mm, solar incident flux = 4000 W/m^2 and QE = 30%.

Fig. 8 depicts the hydrogen volume production rate for different solar incident flux. The hydrogen volume production rate is calculated from the following equation

$$\phi = \frac{\text{QE} \times \text{E}_{\text{s}}}{\Delta G_{\text{H}_2\text{O}}(\text{T}, p)} \times 3600 \text{ s} \times 22.4 \text{ L}$$
 (13)

At the standard state, the gas volume of 1 mol is 22.4 L.

Because the energy available for splitting water decreases as E_g is increased, the hydrogen production rate decreases. For solar incident flux = 1000 W/m² and QE = 30% with design D, the hydrogen volume production rate is 35.0 L/m² h and 5.7 L/m² h respectively for Fe₂O₃ and TiO₂. If solar incident flux is raised to

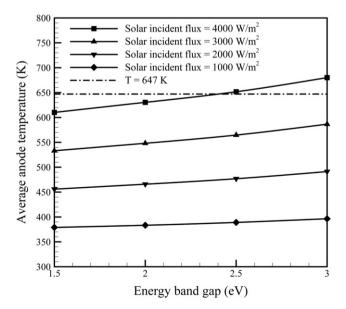


Fig. 7 – Average anode temperature for different solar incident flux in design D with W=10 mm, G=10 mm, A=20 mm and QE=30%.

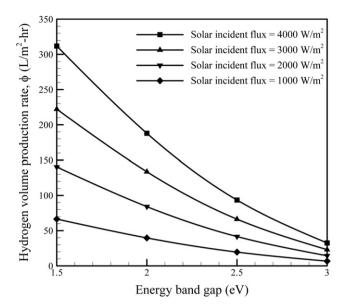


Fig. 8 – Hydrogen volume production rate for different solar incident flux with W = 10 mm, G = 10 mm, A = 20 mm and QE = 30%.

4000 W/m² while keeping other parameters unchanged, the hydrogen volume production rate is increased to $166 \, \text{L/m}^2 \, \text{h}$ and $31.5 \, \text{L/m}^2 \, \text{h}$ for Fe₂O₃ and TiO₂ respectively. Therefore, increasing solar incident flux is an effective way to raise the hydrogen volume production rate.

Fig. 9 presents the effect of solar incident flux on $\eta_{\rm SH}$. As solar incident flux is increased, the average anode temperature is increased. Therefore, the energy required for splitting water is decreased. For a solar incident flux of 1000 W/m² and QE = 30% with design D, $\eta_{\rm SH}$ is 12.3% and 2.3% respectively for Fe₂O₃ and TiO₂. If the solar incident flux is raised to 4000 W/m²,

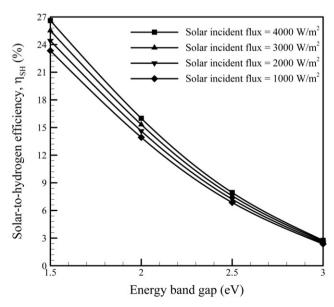


Fig. 9 – Solar-to-hydrogen efficiency for different solar incident flux with W=10 mm, G=10 mm, A=20 mm and QE=30%.

 $\eta_{\rm SH}$ is increased to 14.1% and 2.6% for Fe₂O₃ and TiO₂ respectively. The enhancement is 14.6% and 13.0% respectively for Fe₂O₃ and TiO₂ as solar incident flux is increased from 1000 W/m² to 4000 W/m² with design D.

3.3. Effects of G and W

Figs. 10 and 11 show the effects of the back glass thickness G and the distance between anode and the glass wall W for design D with solar incident flux = 4000 W/m^2 and QE = 30%. As shown in Fig. 10, the average anode temperature increases with increasing G and decreases with increasing G. For the case of $E_g = 1.5$ eV, $\alpha_{ave} = 381.5$ 1/m in the long wavelength region, and the absorptivity of the glass is 32%, 85% and 98% respectively for G = 1 mm, 5 mm and 10 mm. This means that as G is increased, more long wavelength energy is absorbed by the glass. This absorbed energy is re-emitted to heat the water. On the other hand, as G is increased, more water exists in the reactor. The absorbed solar energy is dispersed into larger amount of water, and the average anode temperature decreases. Therefore, larger G and smaller G are good for heating the reactor.

In Fig. 11, the variation of η_{SH} with G and W is depicted. For G=10 mm and W=10 mm, η_{SH} is 14.1% and 2.7% respectively for Fe₂O₃ and TiO₂. But for G=1 mm and W=100 mm, η_{SH} is 13.0% and 2.4% respectively for Fe₂O₃ and TiO₂. The enhancement using the former design parameters over the latter ones is 8.5% and 12.5% for Fe₂O₃ and TiO₂, respectively.

3.4. Effects of quantum efficiency

Fig. 12 shows the effects of QE on the average anode temperature for solar incident flux = 4000 W/m^2 with design D. As explained previously, the average anode temperature increases as E_g is increased. Increasing QE, the average anode temperature decreases due to more short wavelength energy

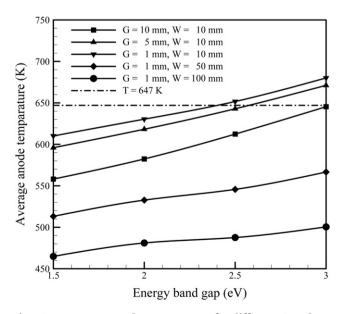


Fig. 10 – Average anode temperature for different G and W with A = 20 mm, solar incident flux = 4000 W/m² and QE = 30%.

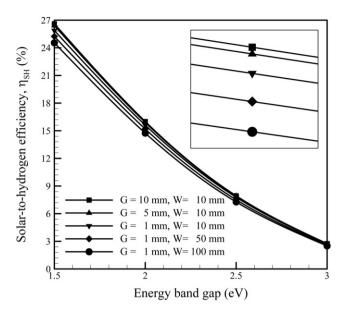


Fig. 11 – Solar-to-hydrogen efficiency for different G and W with A = 20 mm, solar incident flux = 4000 W/m^2 and QE = 30%.

being used to produce electron—hole pairs and less short wavelength energy remained to be converted into heat. The remaining short wavelength energy is larger for lower E_g . Therefore, the difference in anode temperature caused by QE is more obvious for lower E_g .

Figs. 13 and 14 show the effects of QE on the hydrogen volume production rate and η_{SH} for solar incident flux of 4000 W/m². For TiO₂, the hydrogen volume production rate is 9.8 L/m² h, 20.7 L/m² h and 31.5 L/m² h respectively for QE = 10%, 20% and 30%. For Fe₂O₃, the hydrogen volume production rate is 56.8 L/m² h, 112 L/m² h and 166 L/m² h respectively for QE = 10%,

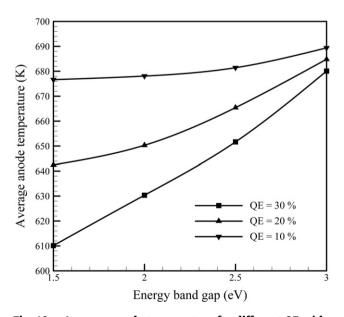


Fig. 12 – Average anode temperature for different QE with W=10 mm, G=10 mm, A=20 mm and solar incident flux = 4000 W/m².

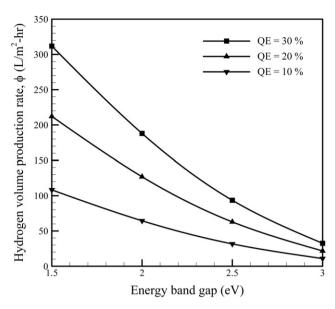


Fig. 13 – Hydrogen volume production rate for different QE with W=10 mm, G=10 mm, A=20 mm and solar incident flux = 4000 W/m².

20% and 30%. Furthermore, for TiO $_2$, $\eta_{\rm SH}$ is 0.9%, 1.7% and 2.6% respectively for QE = 10%, 20% and 30%, and for Fe $_2$ O $_3$, $\eta_{\rm SH}$ is 4.8%, 9.5% and 14.1% respectively for QE = 10%, 20% and 30%. As a comparison, Kudo [13,14] used Pt-loaded AgInZn $_7$ S $_9$ (E $_g=2.3$ eV) as the photoanode, and experimentally measured the H $_2$ production rate under a 300 W Xe lamp. The rate of H $_2$ production was 944 μ mol/h at 33 cm 2 (equivalent to 6.41 L/h m 2) and the quantum efficiency at 420 nm was 20%. Using desing D of the present study with QE = 20% and a 300 W visual light incidence, the hydrogen volume production rate is 7.10 L/h m 2 . The result is very close to Kudo's result.

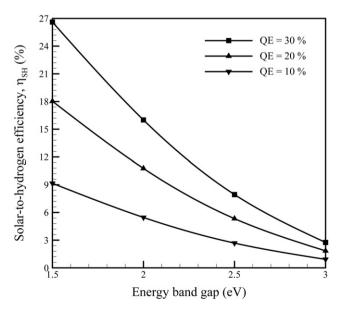


Fig. 14 – Solar-to-hydrogen efficiency for different QE with W=10 mm, G=10 mm, A=20 mm and solar incident flux = 4000 W/m².

4. Conclusions

The heat transfer and flow characteristics of PEC hydrogen reactors have been investigated numerically. Four different reactor designs are studied in this work. Solar irradiation is separated spectrally into short and long wavelength parts depending on the energy band gap of the photoelectrode used. Effects of several parameters on the hydrogen volume production rate and solar-to-hydrogen efficiency are studied.

Of the four designs, design D has the best efficiency because the long wavelength energy is used effectively. Comparing with design A, the enhancement in η_{SH} of design D for solar incident flux = 4000 W/m^2 and QE = 30% is 11.0% and 18.2% respectively for Fe₂O₃ and TiO₂.

As the solar incident flux is increased from 1000 W/m² to 4000 W/m², for Fe₂O₃ with QE = 30% and design D, the hydrogen volume production rate is increased from 35.0 L/m² h to 166/m² h and the enhancement in η_{SH} is 14.6%.

Larger G and smaller W are good for heating the reactor. As G is increased from 1 mm to 10 mm and W is decreased from 100 mm to 10 mm, for Fe₂O₃ with solar incident flux = 4000 W/m², QE = 30% and design D, the enhancement in η_{SH} is 8.5% and 12.5%, respectively.

QE directly affects the hydrogen volume production rate and $\eta_{\rm SH}$. As QE is increased, both hydrogen volume production rate and $\eta_{\rm SH}$ are increased.

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