See discussions, stats, and author profiles for this publication at: https://www.researchgate.net/publication/239203165

Experimental (FT-IR, FT-RS) and theoretical (DFT) studies of vibrational dynamics and molecular structure of 4- n-pentylphenyl-4'-n-heptyloxythiobenzoate (70S5)

ARTICLE in JOURNAL OF MOLECULAR STRUCTURE · SEPTEMBER 2009

Impact Factor: 1.6 · DOI: 10.1016/j.molstruc.2009.05.051

CITATIONS READS

CHAHONS

14

48

3 AUTHORS, INCLUDING:



Diana Maria Dołęga

Jagiellonian University

6 PUBLICATIONS 23 CITATIONS

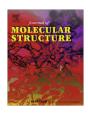
SEE PROFILE

ELSEVIER

Contents lists available at ScienceDirect

Journal of Molecular Structure

journal homepage: www.elsevier.com/locate/molstruc



Experimental (FT-IR, FT-RS) and theoretical (DFT) studies of vibrational dynamics and molecular structure of 4-*n*-pentylphenyl-4′-*n*-heptyloxythiobenzoate (7OS5)

Diana Dołęga a,*, Anna Migdał-Mikuli a, Janusz Chruściel b

ARTICLE INFO

Article history: Received 6 April 2009 Accepted 28 May 2009 Available online 6 June 2009

Keywords: Liquid crystal Vibrational spectroscopy Quantum chemical calculation DFT calculation Potential energy distribution Hartree-Fock

ABSTRACT

Vibrational spectra of 4-n-pentylphenyl-4'-n-heptyloxythiobenzoate (70S5) in the solid phase have been measured using Fourier transform infrared (FT-IR) and Raman scattering (FT-RS). Quantum chemical calculations using density functional theory (DFT) with functionals SVWN, B3PW91, MPW1PW91, PBEPBE, and B3LYP at the 6-31+G* basis set level were performed. Calculations using Hartree–Fock (HF) method with the same basis set were also carried out. Finally the complete assignments of the vibrational modes bands were performed on the basis of B3LYP/6-311+G** calculations. Ultimate geometry of molecule was also fixed. The assignments of the bands was based on the potential energy distribution (PED) data using the GAR2PED program. The general agreement between the observed and calculated frequencies are shown

© 2009 Elsevier B.V. All rights reserved.

1. Introduction

In the present work, we have investigated 4-n-pentylphenyl-4'-n-heptyloxythiobenzoate, with formula: $H_{15}C_7O-C_6H_4-COS-C_6H_4-C_5H_{11}$, using later the acronym 7OS5, by infrared and Raman spectroscopy and by comparison of the experimental spectra with theoretical ones. The mesomorphic behaviour of this compound is characterized by the following enantiotropic phases and phase transitions:

Crystal
$$\xrightarrow{331 \text{ K}}$$
 Nematic $\xrightarrow{353 \text{ K}}$ Isotropic liquid.

The title compound belongs to the group of thermotropic liquid crystals in homologous series of the 4-n-pentylphenyl-4'-n-alkyloxythiobenzoates (mOS5), where m means a number of the carbon atoms in alkoxy-chain [1]. Despite their similar molecular structure the crystals which form this series exhibit very different polymorphism [2]. We believe that theoretical and spectroscopic structural analyzes, combined with crystallographic studies, may help to explain the nature of this polymorphism phenomenon [3,4]. Observation of vibrational bands of the most important functional groups makes it possible to achieve information about change of position both a flexible group and a joining bridge. Calamitic (rod-shaped) molecules forming liquid crystals systems

are usually quite big and contain many atoms. Therefore they have rather rich vibrational spectra. By comparing of experimental spectra with the results of theoretical calculations we can obtain more complete and reliable structural information about their molecular system [1,3]. The theoretical methods based on density functional theory (DFT) are very important tool and widely used for the computation of molecular structure and vibrational frequencies especially for polyatomic systems [4–7].

2. Experimental

The 7OS5 sample was synthesized in the Institute of Chemistry of the University of Podlasie at Siedlce (Poland) by Ossowska-Chruściel [1]. The experimental infrared spectrum in the range of 400–50 cm⁻¹ was recorded using a Fourier transform spectrometer Nicolet Magna-IR 760, equipped with a deuterium triglycine sulphate (DTGS) detector and with a Globar lamp as the light source. The spectrum was obtained with a spectral resolution of 1 cm⁻¹. Sample was in the solid state at room temperature. We have also collected spectrum in the range 4000–400 cm⁻¹. In this case it was used Fourier transform spectrometer Brucker EQUINOX 55 FT-MIR equipped with a highly sensitive mercury cadmium telluride (MCT) detector and a Nernst lamp as the light source. The spectrum was obtained with a spectral resolution of 4 cm⁻¹. Sample was in the solid state. Polarized light scattering spectrum was recorded using Fourier transform Raman scattering (FT-RS)

^a Faculty of Chemistry, Jagiellonian University, Ingardena 3, 30-060 Kraków, Poland

^b Faculty of Sciences, University of Podlasie, 3-go Maja 54, 08-110 Siedlce, Poland

^{*} Corresponding author. Tel.: +48 12 6632265; fax: +48 12 6340515. E-mail address: diana.dolega@uj.edu.pl (D. Dolęga).

spectrometer Thermo-Nicolet NXR 9650, equipped with germanium detector and using Nd-YAG laser, emitting light at 1064 nm. The FT-RS spectrum was collected in the range of 4000–100 cm⁻¹. One hundred fifty scans were performed with a spectral resolution of 4 cm⁻¹ with a laser power of 471 mW for the polycrystalline sample.

3. Calculations

3.1. Molecular structure

In order to give detailed assignment of the vibrational spectra and consequently to solve the molecular structure we applied quantum mechanical calculations in frame of DFT and HF methods. All the calculations were performed using the Gaussian 03 [8] application running on a Baribal supercomputer at the Academic Computer Centre "Cyfronet" in Kraków. All the quantum chemical calculations were carried out with the GaussView molecular visualisation program [9]. A complete optimization of geometry was carried out employing following procedure. At first we had used optimized structure of shorter homologous compound 4OS5 using the molecular mechanics MM+ algorithm [1,10]. In the next step we carried out calculations by Austin model 1 (AM1), semi-empirical method for the obtained structure [10]. The total geometry optimization at the DFT level (SVWN, MPW1PW91, B3PW91, PBEPBE and B3LYP functionals) with basis set 6-31+G* were performed and vibrational frequencies were computed. Ab initio Hartree-Fock (HF) calculations at the same basis set were also done. The molecular geometry was optimized using the Becke three-parameter, Lee, Yang Parr hybrid functional (B3LYP) [7] with the basis set 6-311+G**. Next the optimized structure of 7OS5 was used to compute vibrational frequencies. All the optimized structures showed positive harmonic vibrations only (all eigenvalues are positive, no imaginary frequency indicates a true energy

minimum [7,8]). As far as we know, there are no experimental data for the geometrical structure of 70S5 molecule. Therefore we can not compare the theoretical geometric parameters with experimental values. Rauhut and Pulay [11] show, based on experimental data of 20 organic molecules, that B3LYP functional leads to geometry parameters close to experimental values. We base on these investigations and expect that B3LYP calculated geometry lead us to the most reliable molecular structure. The certain values of calculated bond lengths and angles for 70S5 are included in Table 1. The atom numbering of 70S5 and the molecular structure of investigated molecule 7OS5 is presented in Fig. 1, using program [9]. Main difference is observed in the φ dihedral angle between atoms: C_{19} – C_{18} – S_{17} – C_{15} which defines orientation of both phenyl rings. Phenyl rings are revealed in perpendicular orientation. Moreover, all the calculations predict that geometry in equilibrium state is linear. Additionally, PES (Potential Energy Surface) scan. which has been prepared by AM1 for 8SO5 homologous compound by Drużbicki et al. [8,12] revealed perpendicular geometry of the central group of atoms. We can conclude general agreement with data for homologous compounds presented in literature [1,12].

3.2. Molecular vibrations

We performed potential energy distribution (PED) calculations to show a detailed description of vibrational modes. For this reason we used GAR2PED program [13]. It is worth to mention that Gaussian 03 package calculate only Raman activities. To compare experimental and theoretical spectra Raman intensities are required. Therefore Raman intensities were determined from the RAINT [14] program using values of Raman scattering activities obtained from the Gaussian 03 outputs. Obtained vibrational frequencies are reported in Table 2. We employ calibrating procedure. This process is necessary for the HF method to correct for overestimation of vibrational frequencies due to vibrational anharmonicity and

Table 1Selected angles values for equilibrium geometries obtained at the different levels of theory.

Method	$\theta_{C(15)-S(17)-C(18)}$	φ _C (10)-C(9)-O(8)-C(7)	φ _C (16)-C(15)-S(12)-C(13)	φ _C (19)-C(18)-S(17)-C(15)	φ _C (25)–C(24)–C(21)–C(20)
B3LYP 6-311+G(d,p)	100.76	0.35	-179.55	87.24	87.96
B3LYP 6-31+G(d)	100.97	0.08	-178.07	86.38	89.62
PBEPBE 6-31+G(d)	100.90	-0.08	-176.82	69.21	91.89
B3PW91 6-31+G(d)	100.82	0.50	-179.86	87.14	88.13
MPW1PW91 6-31+G(d)	100.63	0.45	-177.54	84.82	90.41
SVWN 6-31+G(d)	100.81	0.02	177.90	54.64	88.67
HF 6-31+G(d)	101.88	0.00	-179.89	91.22	89.31

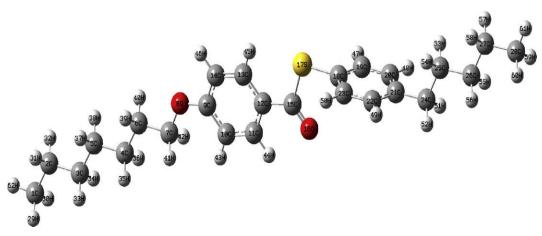


Fig. 1. A view of 70S5 molecule with the atomic labelling.

Table 2Vibrational frequencies, infrared and Raman intensities and their assignments.

No.	brational frequencies, infrared and Raman ir b. Experimental Calculated frequency frequency		Intensities	Assignment (PED(≥ 10%))			
	IR	RS	Unscaled	Scaled	(IR, RS)*		
1.	-	-	10	10	(0.02,29.28)	$\tau C_{15}S_{16}(23) + \tau C_{7}O_{8}(14) + \tau C_{12}C_{15}(12)$	
2.	_	-	14	14	(0.02,7.42)	$\tau S_{17}C_{18}(23) + \tau C_7O_8(12)$	
3. 4.	_	_	17 22	17 22	(0.01,8.72)	$\tau C_{12}C_{15}(32) + \tau C_{15}S_{17}(18)$	
5.	_	_	28	28	(0.01,100) (0.03,5.50)	$\tau S_{17}C_{18}(53) - \tau C_{21}C_{24}(29)$ $\rho C_4C_6C_7(15) + \tau C_{21}C_{24}(12) - \delta C_6C_7O_8(10)$	
6.	_	_	33	33	(0.16,1.13)	$\tau C_{12}C_{15}(36) - \tau S_{17}C_{18}(19)$	
7.	-	-	36	36	(0.01,16.51)	$\tau C_{21}C_{24}(34) - \tau S_{17}C_{18}(31)$	
8.	-	-	42	42	(0.00, 2.17)	$\tau C_6 C_7(23) + \tau C_4 C_5(17) + \tau C_4 C_6(15)$	
9.	-	-	49	49	(0.01,1.96)	$\tau C_{21}C_{24}(37) + \tau C_{24}C_{25}(26) + \tau C_{26}C_{27}(26)$	
10. 11.	55w 69vw	-	55 58	55 58	(0.05,12.08) (0.02,8.25)	$\begin{array}{l} \delta C_{15} S_{17} C_{18}(19) + \delta C_{21} C_{25} C_{24}(19) - \rho C_{22} C_{24} C_{21}(11) \\ \tau O_8 C_9(23) - \tau C_7 O_8(17) - \tau C_{12} C_{15}(11) + \tau C_4 C_6(11) \end{array}$	
12.	75vw	_	88	88	(0.15,3.66)	$\delta C_{11}C_{15}C_{12}(15) + \delta C_{3}C_{4}C_{5}(15) + \delta C_{12}S_{17}O_{16}C_{15}(11)$	
13	95vw	-	95	95	(0.04,0.03)	$\tau C_2 C_3(27) - \tau C_6 C_7(23) + \tau C_9 C_{10} C_{11} C_{12} C_{13} C_{14}(15)$	
14.	102w	-	104	104	(0.02,0.30)	$\tau C_{25}C_{26}(38) - \tau C_3C_5(16)$	
15.	-	-	106	106	(0.02,0.28)	$\tau C_{25}C_{26}(33) + \tau C_3C_5(17) - \tau O_8C_9(15)$	
16 17.	111w 145w	111s -	111 134	111 134	(0.07,4.08)	$\delta C_9 O_8 C_7(15) + \rho C_2 C_5 C_3(12)$	
18.	143W 148 m	_	135	135	(0.22,0.14) (0.21,0.47)	$\tau C_{26}C_{27}(24) + \tau O_8C_9(14)$ $\tau C_{26}C_{27}(24) - \tau O_8C_9(14)$	
19.	-	_	136	136	(0.06,0.81)	$\tau C_{25}C_{26}C_{27}(16) + \tau O_8C_9(14) + \rho C_{28}H3(13)$	
20.	157 m	-	142	142	(0.10,0.14)	$\tau C_4 C_6(40) - \tau C_3 C_5(22) - \tau O_8 C_9(15) + \tau C_{12} C_{15}(12)$	
21.	163w	169 m	159	159	(0.20,0.32)	$\tau C_4 C_5(33) - \tau C_6 C_7(19) - \tau C_2 C_3(18)$	
22.	180w	182w	172	172	(0.15,2.55)	$\delta C_9 O_8 C_7(13) + \delta C_1 C_3 C_2(10)$	
23. 24.	213w -	213vw -	210 237	210 237	(1.23,0.31) (0.00,0.47)	$\gamma^{2}C_{18}C_{19}C_{20}C_{21}C_{22}C_{23}(16) + \delta C_{9}O_{8}C_{7}(13) + \delta C_{15}S_{17}C_{18}(10)$ $\delta C_{4}C_{7}C_{6}(16) + \delta C_{1}C_{3}C_{7}(16)$	
2 4 . 25.	232vw	_	238	238	(0.00,0.47)	$\tau C_{27}C_{28}(89)$	
26.	_	-	248	248	(0.00,0.01)	τC ₁ C ₂ (92)	
27.	256vw	-	251	251	(0.04,0.09)	$\delta S_{17}C_{18}C_{23}(56) + \delta C_{21}C_{22}C_{24}(14)$	
28.	269 m	269w	267	267	(9.53,3.91)	$\delta C_{12}C_{13}C_{15}(14) + \delta C_1C_3C_2(10)$	
29. 30.	287vw 303 m	- 303w	293 303	293 303	(0.07,1.49) (5.50,2.30)	$\gamma^2 C_9 C_{10} C_{11} C_{12} C_{13} C_{14} (36) + \tau C_{12} C_{15} (29)$ $\delta C_{25} C_{26} C_{27} (12)$	
31.	367 m	- -	358	358	(0.52,0.86)	$\delta C_{25}C_{26}C_{27}(12)$ $\delta C_{26}C_{27}C_{28}(42) + \delta C_{12}C_{15}S_{17}(14)$	
32.			362	362	(0.35,0.67)	$\delta C_1 C_2 C_3(22) + \delta C_{24} C_{25} C_{26}(11)$	
33.	377w	-	375	375	(0.02, 0.29)	$\delta C_{22}C_{21}C_{24}(42) - \delta C_{15}S_{17}C_{18}(18)$	
34	403w	412	376	376	(0.07,0.98)	$\delta C_{22}C_{21}C_{24}(11) - \delta C_3C_4C_5(10)$	
35. 36.	412	412w	417 421	417 421	(0.00,0.02) (0.08,2.67)	$\gamma^3 C_{18} C_{19} C_{20} C_{21} C_{22} C_{23} (83)$ $\delta C_{25} C_{26} C_{27} (14) + \beta^2 C_{18} C_{19} C_{20} C_{21} C_{22} C_{23} (11)$	
37.	427w	427w	422	422	(0.00,0.13)	$\gamma^{1}C_{9}C_{10}C_{11}C_{12}C_{13}C_{14}(80)$	
38.	443w	441w	435	435	(1.05,5.16)	$\delta^2 C_{12}C_{15} = O_{16}S_{17}(27) - \rho S_{17}O_{16}C_{15}(14) + \gamma^2 C_{18}C_{19}C_{20}C_{21}C_{22}C_{23}(10)$	
39.	473w	-	458	458	(1.01,0.07)	$\gamma^2 C_{18} C_{19} C_{20} C_{21} C_{22} C_{23} (15)$	
40.	485 m	-	474	474	(0.08,0.88)	$\delta C_3 C_5 C_4(18) + \delta C_4 C_6 C_7(14) + \delta C_1 C_2 C_3(13) + \delta C_5 C_6 C_4(12)$	
41. 42.	499vw 506w	506vw 517w	508 517	508 517	(1.36,0.39) (2.20,5.61)	$ \gamma^2 C_9 C_{10} C_{11} C_{12} C_{13} C_{14} (32) + \omega C_7 O_8 C_9 (31) - \gamma C_{15} C_{11} C_{13} C_{12} (15) + \gamma^2 C_9 C_{10} C_{11} C_{12} C_{13} C_{14} (10) \\ \delta C_7 O_8 C_9 (16) + \delta C_4 C_6 C_7 (12) + \delta C_{10} C_9 O_8 (10) $	
43.	527 m	527vw	546	546	(2.03,0.15)	$\gamma^3 C_{18} C_{19} C_{20} C_{21} C_{22} C_{23} (29) + \gamma^2 C_{18} C_{19} C_{20} C_{21} C_{22} C_{23} (16) - \gamma C_{15} S_{17} C_{18} (14)$	
44.	568w	568w	574	574	(0.96,0.74)	$\beta^{1}C_{9}C_{10}C_{11}C_{12}C_{13}C_{14}(14) + \delta O_{16}C_{15}S_{17}(13) + \delta C_{15}S_{17}C_{18}(10) + \gamma^{2}C_{18}C_{19}C_{20}C_{21}C_{22}C_{23}(10)$	
45.	625 m	631 m	642	629		$\beta^2 C_9 C_{10} C_{11} C_{12} C_{13} C_{14} (58)$	
46.	-	-	646 651	633	(3.12,0.33)	$\gamma S_{17}C_{15}O_{16}(50) + \gamma^{1}C_{9}C_{10}C_{11}C_{12}C_{13}C_{14}(20) + \omega C H_{Ph1}(12)$	
47. 48.	- 660 m	- 660 m	662	638 649	(0.01,3.29) (16.28,22.60)	$\beta^{2}C_{18}C_{19}C_{20}C_{21}C_{22}C_{23}(85)$ $\beta^{2}C_{9}C_{10}C_{11}C_{12}C_{13}C_{14}(21) + \beta^{3}C_{9}C_{10}C_{11}C_{12}C_{13}C_{14}(20) + \nu C_{15}S_{17}(11)$	
49.	-	-	674	661	(0.59,0.70)	$vS_{17}C_{18}(26) + \beta^{1}C_{18}C_{19}C_{20}C_{21}C_{22}C_{23}(25)$	
50.	-	-	733	718	(0.75,0.02)	$\rho C_3 H2(29) + \rho C_5 H2(26) + \rho C_2 H2(15) + \rho C_4 H2(10)$	
51.	-	721w	734	719	(0.48,0.01)	$\rho C_{26}H2(33) + \rho C_{25}H2(24) + \rho C_{27}H2(19)$	
52. 53.	-	_	738 746	723 731	(0.34,1.07) (0.03,0.02)	$\gamma^{1}C_{9}C_{10}C_{11}C_{12}C_{13}C_{14}(59) + \gamma O_{8}C_{9}(16) + \gamma C_{12}C_{15}(14)$ $\rho C_{4}H2(29) + \rho C_{7}H2(21) + \rho C_{6}H2(16)$	
53. 54.	– 775vw	_	746 754	739	(1.52,0.64)	$\beta^2 C_{18} C_{19} C_{20} C_{21} C_{22} C_{23} (69) - \gamma C_{21} C_{24} (11)$	
55.	-	_	760	745	(0.04,0.05)	$\rho C_{18} C_{19} C_{20} C_{21} C_{22} C_{23} (30) = \rho C_{21} C_{24} (11)$ $\rho C_{27} H2(31) + \rho C_{26} H2(25) + \rho C_{25} H2(16)$	
56.	-	-	784	768	(0.34,0.07)	$\rho C_6 H2(27) + \rho C_5 H2(17) + \rho C_2 H2(17)$	
57.	-	-	822	806	(1.11,5.16)	$\beta^2 C_{18} C_{19} C_{20} C_{21} C_{22} C_{23} (15) + \nu C_{21} C_{24} (14)$	
58. 59.	- 833s	- 831w	822 833	806 816	(0.09,0.11) (5.80,32.36)	$\gamma H_{43}C_9C_{11}C_{10}(43) + \gamma H_{46}C_9C_{14}C_{13}(21) + \gamma H_{45}C_{14}C_{12}C_{13}(15) + \gamma H_{44}C_{10}C_{12}C_{11}(15)$ $vO_8C_9(16) + \beta^1C_9C_{10}C_{11}C_{12}C_{13}C_{14}(13) + vC_9C_{10}(13)$	
59. 60.	-	831W -	833 844	827	(0.61,0.10)	$\gamma H_{48}C_{19}C_{21}C_{20}(18) + \gamma H_{47}C_{18}C_{20}C_{19}(15) - \gamma H_{49}C_{21}C_{23}C_{22}(17) + \gamma H_{50}C_{22}C_{18}C_{23}(15)$	
61.	-	_	846	829	(4.79,0.04)	$\gamma H_{46}C_{9}C_{13}C_{14}(21) + \gamma H_{45}C_{14}C_{12}C_{13}(13) - \gamma H_{43}C_{9}C_{11}C_{10}(10)$	
62.	-	-	848	831	(8.32,1.18)	$\gamma H_{48}C_{19}C_{21}C_{20}(18) + \gamma H_{47}C_{18}C_{20}C_{19}(15) + \gamma H_{49}C_{21}C_{23}C_{22}(13) + \gamma H_{50}C_{22}C_{18}C_{23}(10)$	
63.	-	-	851	834	(0.55,0.16)	$\rho C_{26}H2(18) + \rho C_{25}H2(17) + \rho C_{28}H2(12) - \gamma H_{49}C_{21}C_{23}C_{22}(11) - \gamma H_{50}C_{22}C_{18}C_{23}(10)$	
64. 65.	-	-	856 899	839 881	(3.93,0.04) (1.38,1.02)	$\gamma H_{43}C_9C_{11}C_{10}(12) + \gamma H_{45}C_{12}C_{13}C_{14}(10)$ $\nu C_2C_3(28) + C_1H3(25)$	
66.	_	_	904	886	(2.54,2.67)	$VC_2C_3(2\delta) + C_1\Pi_3(2\delta) + VC_2C_2(1\delta)$ $VC_2C_2(2\delta) + C_2RH_3(2\delta) + VC_2C_2(1\delta)$	
67.	904vs	899 m	908	890		$\delta^2 C_{12} C_{15} O_{16} S_{17}(29) + v C_{15} S_{17}(14) - v C_{12} C_{15}(11)$	
68.	-	-	939	920	(0.01,0.11)	$\rho C_7 H2(16) + \rho C_5 H2(15) + \rho C_1 H3(12) + \rho C_2 H2(12) + \rho C_4 H2(11)$	
69.	-	-	970	951	(0.04,0.95)	$\gamma H_{50}C_{22}C_{18}C_{23}(21) + \gamma H_{47}C_{18}C_{20}C_{19}(20) + \gamma^2 C_{18}C_{519}C_{20}C_{21}C_{22}C_{23}(18) - \gamma H_{48}C_{19}C_{21}C_{20}(16) - \gamma H_{49}C_{21}C_{23}C_{22}(17)$	
70. 71.	-	_	970 973	951 954	(0.01,0.02) (0.07,0.11)	$\gamma H_{45}C_{14}C_{12}C_{13}(34) - \gamma H_{46}C_{9}C_{14}C_{13}(22) - \gamma H_{44}C_{10}C_{12}C_{11}(16) + \gamma H_{43}C_{9}C_{11}C_{10}(16)$ $\rho C_{26}H2(18) + \rho C_{24}H2(17) + \rho C_{27}H2(15) + \rho C_{28}H2(14) + \rho C_{25}H2(13)$	
, 1.	_	_	313	334	(0.07,0.11)	pe26112(10) · pe24112 (17) · pe27112(13) + pe28112(14) + pe25112(13)	

Table 2 (continued)

R S Oscillot Scaled (R. RS)	No.	•		Calculated frequency		Intensities	Assignment (PED(≥10%))
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$						(IR. RS)*	
73 986 956 960 (19.09.03) 'Hat-Gric-Gric-Grid-19 + Hat-Grid-Grid-19 + Hat-Grid-Grid-Grid-19 - Hat-Grid-Grid-Grid-19 - Hat-Grid-Grid-Grid-19 - Hat-Grid-Grid-Grid-19 - Hat-Grid-Grid-Grid-Grid-19 - Hat-Grid-Grid-Grid-Grid-Grid-Grid-Grid-Grid	72						γH ₅ ρ(22C10C22(23) - γH ₄ γC10C22C12(23) + γH ₄ ρC1ρC21C22(22) - γH ₄ ρC21C22(22)
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
75							
77 1016ww 1025 1005 (191.4x44) $\beta^{1}C_{c}(m_{1}^{2}n_{1}^{2}) + C_{c}(n_{1}^{2}) + C_{c}(c_{1}^{2}n_{1}^{2}) $							
78 1028 1007 (011.003) pcj.42[19]* pcj.42[14]* y cj.25[23]* y cj.25[23]* pcj.42[14]* y cj.25[23]* y cj.25[23]* y cj.25[23]* y cj.25[23]* y cj.25[24]* y cj.							
1014 m 1034 1013 211.043 p C ₁ cccccccccccccccccccccccccccccccccccc							
81. 1035w - 1049 1028							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
83 1060w 1062 1041 (0.064.52) $V_{c,c,c,c,d}(3.3) = V_{c,c,c,c}(2.4) = V_{c,c,c}(2.12) = V_{c,c,c,c}(2.12) = V_{c,c,c,c,c}(2.12) = V_{c,c,c,c,c}(2.12) = V_{c,c,c,c,c}(2.12) = V_{c,c,c,c,c}(2.12) = V_{c,c,c,c,c,c}(2.12) = V_{c,c,c,c,c,c,c,c,c,c,c}(2.12) = V_{c,c,c,c,c,c,c,c,c,c,c,c,c,c,c,c,c,c,c,$							
18.							
86. 1969w 1094m 1094 1092 (0.25.0.10) $\rho_{Ca}H^2(2.25) = \rho^2(r_{11}r_{11}r_{12}r_{12}r_{13}r_{13}) - \nabla_{ca}\rho_{Ca}(1/2)$ 87 1116m 1112w 1126 1103 (2.71.5.13) $\delta_{Ca}\rho_{Ca}$							
87 107 1085 (0.5546.38) $\times 5_1 c_{11}(24) - \nabla c_{11}c_{12}(24) + \nabla c_{12}c_{12}(21) + \nabla c_{12}c_{12}(21)$ 88 1 1180 11120 1120 (0.78.01.0) $(c_1, c_2, 1) + c_2, c_2, 1) + c_2, c_2, 1 + $							
88. 116m 112w 112f 113f 113g 121g 113f 0.75c.jal.42(2)] + $\nabla \nabla_{c} \nabla_{c} \int_{c}^{c} (17)$ 89 1141 1118 (1431.32) $\nabla_{c} \nabla_{c} \int_{c}^{c} \int_{c}$							
90 1 141 1118 (1431,32) $[\rho \hat{c}_{L} + h_{LR} \hat{c}_{L}(3)] + [\hat{c}_{L} \hat{c}_{L} + h_{LR} \hat{c}_{L}(2)]$ 91 1184 1184 1100 (100,004,608) $\hat{c}_{L} \hat{c}_{L} \hat{c}$	88.	1116 m	1112w	1126			
91. 194. 1167 1167 1167 1167 1167 1167 1167 116							
9.2 1698 1167 m 1184 1160 (100.004.6)08) $ c_{11}$							
93 187 1163 (039.037)							
95 1223 1198 (0.13.0.22) pc,H2(16) + pc,H2(12) + pc,H2(12) + pc,H2(11) + pc,H2(13) (97.0.2) (1.10) (1							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
97 1 224 1199 (2.95.84) 127.							
98. 2138 2131 2129 1224 2570.53 370.00 2570.53 370.00 2570.53 370.00 2570.50 2570.00 2570							
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	98.	1213s					
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$\begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$			1306 m				11 12 7 12 12 7 1 1 1 1 1 1 1 1 1 1 1 1
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$			-				
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$\begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$\begin{array}{cccccccccccccccccccccccccccccccccccc$							
$\begin{array}{cccccccccccccccccccccccccccccccccccc$							
$\begin{array}{cccccccccccccccccccccccccccccccccccc$						(0.06, 2.57)	$\delta C_{26}H2(44) + \delta C_{27}H2(41)$
$\begin{array}{cccccccccccccccccccccccccccccccccccc$							
$\begin{array}{llllllllllllllllllllllllllllllllllll$							
$\begin{array}{cccccccccccccccccccccccccccccccccccc$							
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	134.	-		1500	1470	(1.59,1.39)	δC ₁ H3(91)
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$\begin{array}{llllllllllllllllllllllllllllllllllll$							
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	139.	1494 m	1494vw	1519	1489	(9.88,0.17)	$\delta C_7 H2(40) + \delta C_5 H2(20) + \delta C_4 H2(14)$
$ \begin{array}{llllllllllllllllllllllllllllllllllll$							$\beta C_{18}C_{19}H_{47}C_{20}(17) - \beta C_{18}C_{22}H_{50}C_{23}(17) + \beta C_{19}C_{20}H_{48}C_{21}(12) - \beta C_{21}C_{22}H_{49}C_{23}(12) - \nu C_{18}C_{19}(10) - \nu C_{18}C_{23}(10)$
143. 1565vw - 1604 1572 (0.83,0.69) $vC_9C_{10}(18) - vC_9C_{14}(17) + vC_{12}C_{13}(17) - vC_{17}C_{18}(16)$							
144 1635 1602 (3.19,22.38) $vC_{19}C_{20}(21) + vC_{22}C_{23}(20) + \beta^2C_{18}C_{19}C_{20}C_{21}C_{22}C_{23}(10)$			-				$\nu C_{19}C_{20}(21) + \nu C_{22}C_{23}(20) + \beta^2 C_{18}C_{19}C_{20}C_{21}C_{22}C_{23}(10)$
(contin							(continued on next page)

Table 2 (continued)

No.	No. Experimental frequency		Calculated frequency		Intensities	Assignment (PED(≥ 10%))
	IR	RS	Unscaled	Scaled	(IR, RS)*	
145.	1601vs	1597vs	1639	1606	(66.98, 100.00)	$vC_{13}C_{14}(23) + vC_{10}C_{11}(19) + \beta^2C_9C_{10}C_{11}C_{12}C_{13}C_{14}(11)$
146.	1663vs	1662vs	1742	1707	(36.60,33.34)	$vC_{15}O_{16}(83)$
147.	-	-	2993	2873	(0.57,34.29)	$vC_3H_{33}(35) + vC_3H_{34}(34) - vC_4H_{36}(10)$
148.	-	-	2994	2874	(1.73,64.59)	$vC_{26}H_{56}(45) + vC_{26}H_{55}(45)$
149.	-	-	2995	2875	(0.61,25.58)	$\nu C_7 H_{41}(16) + \nu C_7 H_{42}(16) - \nu C_5 H_{38}(14) - \nu C_4 H_{36}(14) - \nu C_4 H_{35}(13) - \nu C_5 H_{37}(13)$
150.	-	2858 m	2998	2878	(1.44,86.87)	$vC_7H_{41}(28) + vC_7H_{42}(27) + vC_5H_{37}(15) + vC_5H_{38}(15)$
151.	-	-	3003	2883	(2.49,76.25)	$vC_4H_{36}(18) + vC_4H_{35}(18) + vC_2H_{32}(17) + vC_2H_{31}(17)$
152.	-	-	3005	2885	(1.07,1.43)	$vC_{27}H_{58}(27) + vC_{27}H_{57}(26) - vC_{25}H_{54}(22) - vC_{25}H_{53}(20)$
153.	2855s	-	3010	2890	(17.25,18.16)	$vC_2H_{32}(28) + vC_2H_{31}(26) + vC_5H_{38}(12)$
154.	-	-	3011	2891	(0.13,89.31)	$\nu C_3 H_{33}(24) - \nu C_3 H_{34}(23) - \nu C_5 H_{37}(17) + \nu C_5 H_{38}(14)$
155.	-	-	3012	2892	(5.13,100.00)	$vC_{25}H_{54}(20) + vC_{27}H_{58}(19) + vC_{25}H_{53}(18) + vC_{27}H_{57}(18)$
156.	-	-	3014	2893	(0.25,68.71)	$vC_{26}H_{55}(33) - vC_{26}H_{56}(33) - vC_{27}H_{59}(11) + vC_{27}H_{58}(11)$
157.	2869s	-	3018	2897	(0.74,75.41)	$vC_{28}H_{60}(36) + vC_{28}H_{59}(35) + vC_{28}H_{61}(25)$
158.	-	2884s	3018	2898	(9.03,6.51)	$vC_4H_{36}(23) - vC_4H_{35}(23) - vC_2H_{31}(10)$
159.	-	-	3019	2898	(9.03,84.20)	$vC_1H_{30}(35) + vC_1H_{29}(32) + vC_1H_{62}(24)$
160.	-	2911s	3023	2902	(69.95,41.22)	$vC_{24}H_{51}(40) + vC_{24}H_{52}(37)$
161.	-	-	3029	2908	(8.85,30.57)	$vC_6H_{40}(50) + vC_6H_{39}(44)$
162.	-	-	3032	2911	(0.46,0.21)	$vC_{25}H_{53}(22) - vC_{29}H_{57}(21) + vC_{27}H_{58}(21) - vC_{25}H_{54}(24)$
163.	-	-	3033	2911	(1.55,44.05)	$\nu C_7 H_{41}(20) - \nu C_4 C_6(20) - \nu C_2 H_{31}(11) + \nu C_2 H_{32}(10) - \nu C_6 H_{39}(10)$
164.	-	-	3035	2914	(1.56,1.01)	$vC_7H_{42}(20) - vC_7C_{42}(66) + vC_2H_{32}(12) - vC_2H_{31}(12)$
165.	-	-	3045	2923	(2.67,29.29)	$vC_{24}H_{52}(16) - vC_{24}H_{51}(14) + vC_{27}H_{57}(12) + vC_{27}H_{58}(12) + vC_{26}H_{56}(12) + vC_{26}H_{55}(12)$
166.	2926s	2926s	3050	2928	(10.63, 3.03)	$vC_5H_{37}(18) - vC_5H_{38}(18) - vC_3H_{34}(14) + vC_3H_{33}(14)$
167.	-	-	3065	2943	(9.00,2.42)	$vC_{24}H_{52}(30) - vC_{24}H_{51}(29) - vC_{25}H_{53}(16) + vC_{25}H_{54}(15)$
168.	-	-	3074	2951	(12.21,6.58)	$vC_6H_{39}(39) - vC_6H_{40}(37)$
169.	-	-	3079	2955	(14.61,16.70)	$vC_{28}H_{60}(44) - vC_{28}H_{59}(42)$
170.	-	-	3079	2955	(16.71,13.56)	$vC_1H_{29}(44) - vC_1H_{30}(43)$
171.	-	2960s	3082	2959	(8.86,58.21)	$vC_{28}H_{61}(73) - vC_{28}H_{59}(14) - vC_{28}H_{60}(12)$
172.	2957s	-	3084	2961	(8.93,55.56)	$vC_1H_{62}(74) - vC_1H_{30}(14) - vC_1H_{29}(11)$
173.	3038w	3038s	3157	3031	(2.44,41.05)	$vC_{22}H_{49}(67) - vC_{20}H_{48}(27)$
174.	_	_	3158	3032	(3.70,40.42)	$vC_{20}H_{48}(66) + vC_{22}H_{49}(26)$
175.	-	_	3185	3058	(1.70,15.78)	$vC_{13}H_{45}(64) - vC_{14}H_{46}(34)$
176.	3058w	_	3189	3062	(1.91,4.14)	$vC_{23}H_{50}(52) - vC_{19}H_{47}(42)$
177.	3060vw	3060w	3191	3063	(0.01,70.50)	$vC_{19}H_{47}(51) + vC_{23}H_{50}(40)$
178.	_	_	3194	3066	(0.04,16.03)	$vC_{11}H_{44}(77) - vC_{10}H_{43}(21)$
179.	_	_	3200	3072	(0.39,54.05)	$VC_{14}H_{46}(65) + VC_{13}H_{45}(34)$
180.	_	_	3210	3082	(2.63,59.38)	$vC_{10}H_{43}(78) + vC_{11}H_{44}(21)$

v – Stretching mode, ρ – rocking mode, ω – wagging mode, δ – scissoring mode, β – in-plane-deform, γ – out-of-plane-deform, τ – torsion mode, β^1 – trigonal deformation, β^2 – anti-symmetric deformation the first type, β^3 – anti-symmetric deformation the second type, γ^1 – puckering mode, γ^2 – anti-symmetric torsion the first type, γ^3 – anti-symmetric torsion the second type, low index – labeled atom, vw – very weak, vw – very weak, vw – very weak, vw – very weak, vw – very strong.

neglecting electron correlation the case of DFT methods calibrating factors are used to compensate incomplete incorporation of electron correlation, as well as to account of vibrational anharmonicity [11,15–17]. We did not calibrated the range between 400 and $50~\rm cm^{-1}$. In the wave number range $2000-400~\rm cm^{-1}$ obtained

frequencies were multiplied by factor 0.98. The frequencies from the last area $3500-2000~\rm cm^{-1}$ were multiplied by ratio 0.96. This is consistent with the recommendation Pople et al. [16,17] The obtained experimental spectra (FT-IR, FT-RS) are compared with fitted calculated spectra by five different functionals and $6-31+G^*$

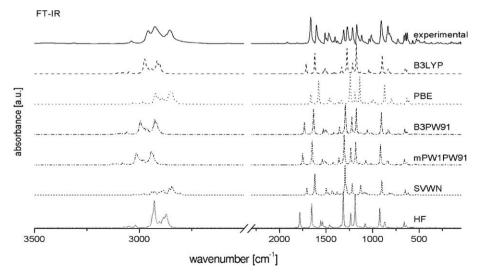


Fig. 2. The comparison of both experimental and scaled theoretical FT-IR spectra for 70S5 in the region between wavenumbers 2260-50 and 3500-2500 cm⁻¹ (different functionals/6-31+G*).

^{*} Percentage relative intensity with respect to the most intensive bands. For FT-IR unscaled calculated spectrum the most intensive band is 1184 cm⁻¹, in FT-RS unscaled spectrum there are bands at 22, 1639 and 3012 cm⁻¹.

basis set are presented in Figs. 2 and 3. Calculated RMS (Root Mean Square) show that hybrid B3LYP functional gives quite good reproduction of experimental frequencies. Intensities relations achieved in the calculations with B3LYP correspond to experimental values pretty well. The experimental and theoretical, calculated with B3LYP functional and higher 6-311+G** basis set, vibrational spectra were shown in Figs. 4 and 5, for FT-IR and FT-RS, respectively.

Figs. 4a, b and 5a, b present calibrated ranges. Far-infrared area of the frequency spectrum is no scaled. Additionally, calculated Dirac delta were shown as spectral lines. We have presented Lorentzian broadening of calculated peaks, which facilitates comparison of calculated spectrum with the experimental one. In the Table 2. We can see the scaled and unscaled values of frequencies calculated on the B3LYP/6-311+G** level of theory. We have

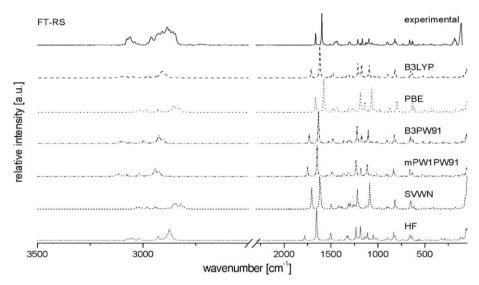


Fig. 3. The comparison of both experimental and scaled theoretical FT-RS spectra for 70S5 in the region between wavenumbers 2260–50 and 3500–2500 cm⁻¹ (different functionals/6-31+G*).

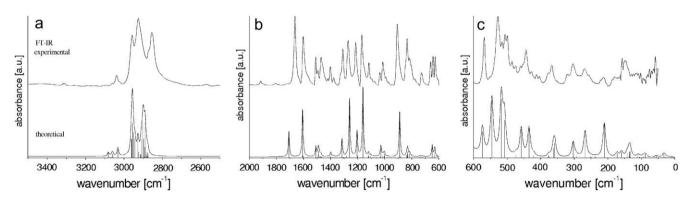


Fig. 4. The comparison of both experimental and theoretical FT-IR spectra for 70S5 (B3LYP/6-311+G**). (a) The unscaled range 50–600 cm⁻¹, (b) the scaled range 600–2000 cm⁻¹ (Scaling factor 0,98), and (c) the scaled range 2000–3500 cm⁻¹ (Scaling factor 0,96).

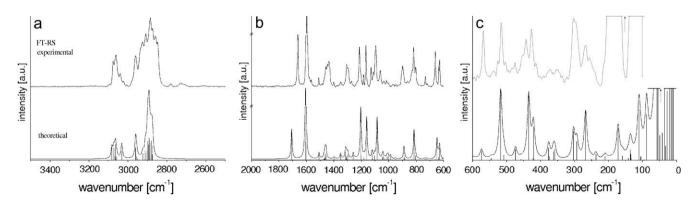


Fig. 5. The comparison of both experimental and theoretical FT-RS spectra for 70S5 (B3LYP/6-311+ G^{**}). (a) The unscaled range 50–600 cm⁻¹, (b) the scaled range 600–2000 cm⁻¹ (Scaling factor 0,98), and (c) the scaled range 2000–3500 cm⁻¹ (Scaling factor 0,96).

normalized theoretical FT-IR spectrum to the most intensive frequency. The calculated FT-RS spectrum has been normalized in three ranges 3500–1800, 1800–600 and 600–50 cm⁻¹.

4. Results and discussion

According to the DFT and group theory results 70S5 has 180 fundamental vibrations. The calculated frequencies of normal vibrations of 70S5 from the B3LYP functional with 6-311+G** basis set have been compared with infrared and Raman results for compound in the solid state. All calculations were performed assuming C1 point group symmetry for 70S5 molecule. It should be emphasized ones more that the simulated spectra refer to the isolated molecule at 293 K, whereas the assignments are made for the vibrational spectra of the solid state at room temperature. Thus, some differences in the calculated versus experimental frequencies or intensities are most probably caused by intermolecular interactions (crystal packing) and/or by anharmonicity of the vibrations, limited range of interactions between atoms and basis set. However, we can see that experimental and theoretical spectra are very similar. Hence, we conclude that intermolecular interactions in this solid state are very weak.

The computed frequencies, compared with the experimental data are given in Table 2. With our calculations we predict the torsion motions in the range 55-232 cm⁻¹ especially in the alkylo- and alkoxychains. At 252 $\,\mathrm{cm^{-1}}$ deformative mode of –CSC-molecular group can be observed. As we can see in Figs. 2 and 3, fundamentals achieved from experimental far-infrared and Raman measures are relatively low intense. Comparison of both 7OS5 and 8OS5 compounds lead us to small differences in the far-infrared range, below 252 cm⁻¹. However, it can be observed there torsion motions of the chains [12]. Calculations at the $300-400 \text{ cm}^{-1}$ region allow us to assign scissoring motions of these chains. The corresponding bands in 8OS5 in the range of skeletal groups are shifted of several cm⁻¹ [12]. Moreover, the calculations predict out-of-plane motions at 417 and 427 cm⁻¹ of the both phenyl rings, which correspond to weak peaks at 412 and 427 cm⁻¹ of the experimental infrared and Raman spectra. The fundamental at 443 cm⁻¹ (IR experimental spectrum) and 441 cm⁻¹ (RS experimental spectrum) we assigned mainly to the bridge vibrations. The next bands, at 500 cm⁻¹ (IR) and 506 cm⁻¹ (RS), are found to be out-of-plane motions of phenyl ring associated with alkoxy-chain. Similar kind of vibrations of the second ring can be observed both in IR and RS experimental spectra. Bands in the range of 420-530 cm⁻¹ of 70S5 compared with these of 80S5 compound have small differences between intensities relations [12]. Vibrations of the joining bridge and phenyl ring appear in the range of 568-660 cm⁻¹ at the experimental spectra. The intensities and frequencies of bands of the 8OS5 compound are overlapped is with a maximal deviation of 3 cm^{-1} [12].

Our calculations predict that the medium intensity fundamentals at 660 cm⁻¹ are found to deform-in-plane phenyl rings. Weak band at 721 cm⁻¹ on Raman spectrum is especially caused by the rocking CH₂ motions in alkoxy-chain. The same band is observed for 8OS5 compound [12]. Very weak intense band in infrared spectrum at 775 cm⁻¹ corresponds with of out-of-plane vibration of phenyl ring from our calculations. Signals in area 833 cm⁻¹ (IR spectrum) and 831 cm⁻¹ (RS spectrum) represent stretching motions of alkoxy-chain and associated phenyl ring. We have noticed that respective bands for 8OS5 are shifted and have different integral intensity [12]. One of the most intensive peaks in the infrared spectrum is related with rocking motion of joining bridge, namely thioester group (-COS-). We can also note stretching vibration of sulphur atom associated with carbon atom from phenyl ring. Frequencies from $974-990 \, \text{cm}^{-1}$ to $1035-1077 \, \text{cm}^{-1}$ ranges in the infrared spectrum represent stretching motions from heptyloxy terminal group. There are no significant quantitative differences at 80S5 spectrum. However, band recorded at about 1016 and 1060 cm⁻¹ has been assigned to the stretching vibrations of terminal group $(-C_5H_{11})$ of investigated molecular system [12]. Also this part of spectrum gives characteristic band because of motions in phenyl rings. Signals at 1096 and 1094 cm⁻¹ in infrared and Raman spectra are assigned to the rocking motions of two atoms of hydrogen associated with carbon atom no. 24 (Fig. 1). However, Drużbicki et al. [12] show that this bands belong to the stretching motion of -C-S-. There is also small contribution of relevant phenyl ring. Predicate upon potential energy distribution calculations, small intensive peak on Raman spectrum, at frequency 1128 cm⁻¹ represent skeletal modes of heptyloxy-chain and deform-in-plane motions of hydrogen atoms in associated phenyl ring. The most intensive band in infrared spectrum at 1169 cm⁻¹ corresponding to that at 1167 cm⁻¹ on Raman spectrum and is caused by deformation motions of carbon atoms with numbers: from 9 to 14 in the first phenyl ring. Weak intensive mode at frequency 1184 cm⁻¹ in Raman experimental spectrum mainly has been assigned to the hydrogen in-plane motions of the second phenyl ring. Towards higher frequency both Raman and infrared experimental spectra are present peaks at frequency 1213 cm⁻¹ strong and weak intensity, respectively, alike at [12]. Those bands are caused by stretching modes between the carbon atom in thioester group and associated carbon from phenyl ring, especially. There are also some contributions of vibrations in the phenyl ring. The band recorded at 1232 cm⁻¹ on Raman spectrum corresponds to wagging motions of hydrogen atoms in terminal (-C₅H₁₁) alkylo-group. However, this band is absent in 8OS5 spectra [12]. While the following band on this spectrum give much contributions by rocking motions of hydrogen atoms in the second alkoxy-terminal group. Intensive band at 1271 cm⁻¹ in the FT-IR spectrum and very weak band in the FT-RS spectrum comes from the stretching mode. The vibrations are connected with alkoxyphenyl group O₈-C₉ (45% PED). Rocking motions of CH2 alkylo-group has been ascribed to the medium intensive band in Raman spectrum. While no significant differences are observed for the related bands in the homologous compound 80S5 [12]. Stretching vibrations of phenyl ring correspond to strong band in the infrared spectrum at frequency 1308 and 1306 cm⁻¹ in the FT-RS spectrum. Wagging CH₂ motions of heptyloxy-fragment appear at 1324 cm⁻¹ in the infrared spectrum. The same kind of vibrations appears in the Raman spectrum at 1350 cm^{-1} . It is concerned with alkylo- $(-C_5H_{11})$ group. Motions of two methyl (-CH₃) groups in 7OS5 give rather big contributions to the bands in the region of 1377 cm⁻¹ (both spectra). Further around 1402 cm⁻¹ there is present peak, which comes from stretching vibrations of oxygen atom in hepyloxy- (C_7-O_8) tail. There are also some contributions of wagging motions (CH₂) and stretching motions between carbon atoms in this molecular fragment. Medium intensive signals in the FT-RS spectrum at region 1437-1468 cm⁻¹ is assigned to deform scissoring motions of hydrogen atoms in alkylo-group. Calculations show that observed bands both infrared and Raman spectra at 1494 cm⁻¹ comes from rocking modes in heptyloxy tail. Deform-in-plane vibrations of both phenyl rings cause emerging signals in the range 1509-1609 cm⁻¹ in the vibrational spectra. Stretching bending of carbonyl group C₁₅=O₁₆ can be observed as a very strong band in both spectra at 1662 cm⁻¹. With our calculations we predict respectively rich contribution (83% PED). The v(C=0) peak for analogous 80S5 which informs about similar interaction was recorded at 1665 cm⁻¹ (in both spectra) [12]. The last region is connected with stretching vibrations between carbon atoms and hydrogen atoms both terminal fragments of 70S5 in the investigated compound. Theoretical calculations also show stretching modes of CH groups in both phenyl rings. According to Drużbicki et al. [12] there were not observed significant differences in this spectral range.

5. Conclusions

The spectra of both 70S5 and 80S5 compounds are very similar (see Ref. [12]). The experimental vibrational spectra of 70S5 molecule are rather rich. It is caused by overtones and combinational bands. Theoretical computations could not reproduce exactly experiment by using of selected model. Results of calculations show that B3LYP functional and relatively high 6-311+G** basis set is quiet good to be compared with experiment. In summary, a complete vibrational assignments of investigated compound is obtained from these studies in addition to supplying normal mode descriptions and PED values for the vibrational frequencies. We can also conclude, that good agreement between the experiment and theory, which was obtained despite the fact that the spectra were recorded in the crystalline state and the computations were based on the isolated molecule model, suggests existence of very weak interactions, probably dispersive, between the molecules of 4-n-pentylphenyl-4'-nheptyloxythiobenzoate in the crystal lattice.

Acknowledgments

Calculations were done at the Academic Computer Centre "Cyfronet", Kraków, Poland which is acknowledged for computing time (Grant number: MNiSW/SGI3700/UI/009/2009).

References

 M.D. Ossowska-Chruściel, Z. Karczmarzyk, J. Chruściel, Mol. Cryst. Liq. Cryst. 382 (2002) 37.

- [2] J. Chruściel, PhD disseration, Jagiellonian University, Cracow, 1978.
- [3] R. Korlacki, K. Merkel, J.K. Vij, R. Wrzalik, A. Kocot, M.D. Ossowska-Chruściel, J. Chruściel, S. Zalewski, Liq. Cryst. 33 (2006) 219.
- 4] A. Kocot, G. Kruk, R. Wrzalik, J.K. Vij, Liq. Cryst. 12 (1992) 1005.
- [5] D. Bauman, E. Mykowska, A. Zięba, E. Chrzumnicka, G. Czechowski, J. Jadżyn, Phase Trans. 80 (2007) 599.
- [6] R. Wrzalik, K. Merkel, A. Kocot, B. Cieplak, J. Chem. Phys. 117 (2002) 10.
- [7] J.B. Foresman, A.E. Frish, Exploring Chemistry with Electronic Structure Methods, Gaussian, Pittsburgh, 1998.
- [8] M.J. Frisch, G.W. Trucks, H.B. Schlegel, G.E. Scuseria, M.A. Robb, J.R. Cheeseman, J.A. Montgomery Jr., T. Vreven, K.N. Kudin, J.C. Burant, J.M. Millam, S.S. Iyengar, J. Tomasi, V. Barone, B. Mennucci, M. Cossi, G. Scalmani, N. Rega, G.A. Petersson, H. Nakatsuji, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, M. Klene, X. Li, J.E. Knox, H.P. Hratchian, J.B. Cross, V. Bakken, C. Adamo, J. Jaramillo, R. Gomperts, R.E. Stratmann, O. Yazyev, A.J. Austin, R. Cammi, C. Pomelli, J.W. Ochterski, P.Y. Ayala, K. Morokuma, G.A. Voth, P. Salvador, J.J. Dannenberg, V.G. Zakrzewski, S. Dapprich, A.D. Daniels, M.C. Strain, O. Farkas, D.K. Malick, A.D. Rabuck, K. Raghavachari, J.B. Foresman, J.V. Ortiz, Q. Cui, A.G. Baboul, S. Clifford, J. Cioslowski, B.B. Stefanov, G. Liu, A. Liashenko, P. Piskorz, I. Komaromi, R.L. Martin, D.J. Fox, T. Keith, M.A. Al-Laham, C.Y. Peng, A. Nanayakkara, M. Challacombe, P.M.W. Gill, B. Johnson, W. Chen, M.W. Wong, C. Gonzalez, J.A. Pople, Gaussian 03, Revision D.01, Gaussian, Inc., Wallingford, CT, 2004.
- [9] W. Gaussview, Æ. Frish, A.B. Nielsen, A.J. Holder, Gaussian Inc., Carnegie Office Park, Building 6, Pittsburg, PA 15106, USA.
- [10] A.R. Leach, Molecular Modelling. Priciples and Application, second ed., Pearson Education EMA, Prentice-Hall, 2001, p. 773.
- [11] G. Rauhut, P. Pulay, J. Phys. Chem. 99 (1995) 3093.
- [12] K. Drużbicki, E. Mikuli, M.D. Ossowska-Chruściel, Vib. Spectrosc. in press.
- [13] J.M.L. Martin, C. Van Alsenoy, Gar2ped, University of Antwerp, 1995.
- [14] R. Wysokiński, D. Michalska, Chem. Phys. Lett. 403 (2005) 211.
- [15] M. Barańska, K. Chruszcz, B. Poduszek, L.M. Proniewicz, Vib. Spectrosc. 31 (2003) 291.
- [16] J.A. Pople, A.P. Scott, M.W. Wong, L. Radom, Isr. J. Chem. 33 (1993) 345.
- [17] J.A. Pople, H.B. Schlegel, R. Krishnan, D.J. Defrees, J.S. Binkley, M.J. Frisch, R.A. Whiteside, R.F. Hout, W.J. Hehre, Int. J. Quant. Chem. S15 (1981) 269.